

Towards the Consistent Dark Matter exploration

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Southampton University & Rutherford Appleton Laboratory

LIO international conference on Future colliders and the origin of mass

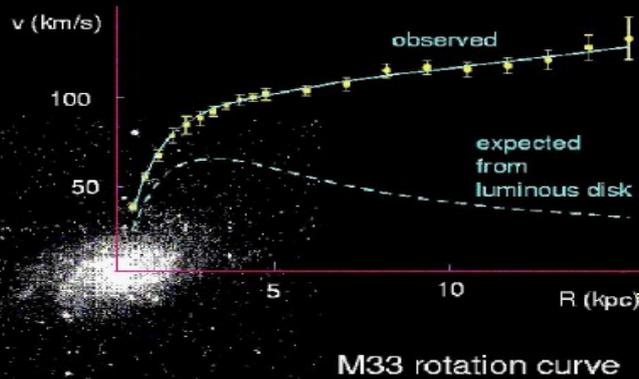
21-25 June 2021

Collaborators & Projects

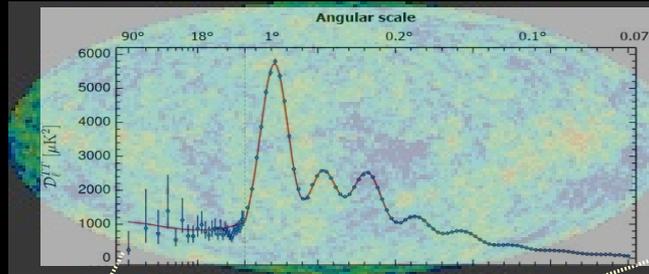
- U.Blumenschein, A. Freegard, D.Gupta, S.Moretti, AB arXiv:**2107.xxxxx**
- **G.Cacciapaglia, D.Locke, AB** arXiv:**2107.xxxxx**
- I.Ginzburg, D.Locke, A. Freegard, T. Hosken, A.Pukhov, AB arXiv:**2107.xxxxx**
- S.Prestel, F.Rojas-Abate, J.Zurita, AB arXiv:**2008.08591**
- G.Cacciapaglia, J.McKay, D. Marin, A.Zerwekh, AB arXiv:**1808.10464**

The existence of Dark Matter is confirmed by several independent observations at cosmological scale

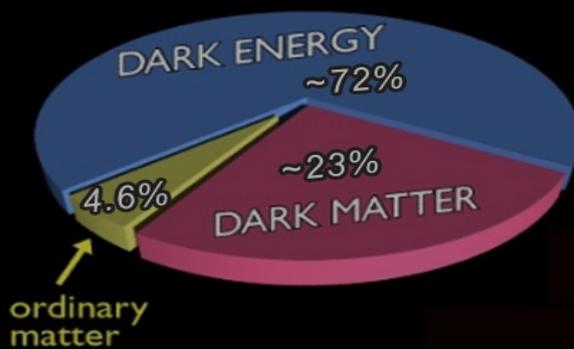
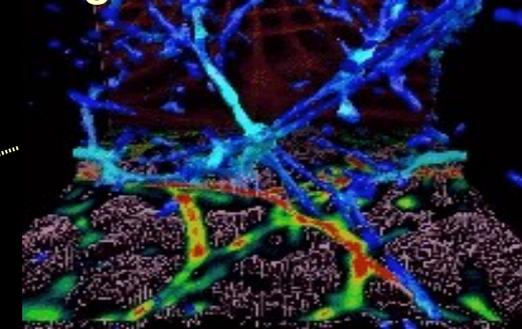
Galactic rotation curves



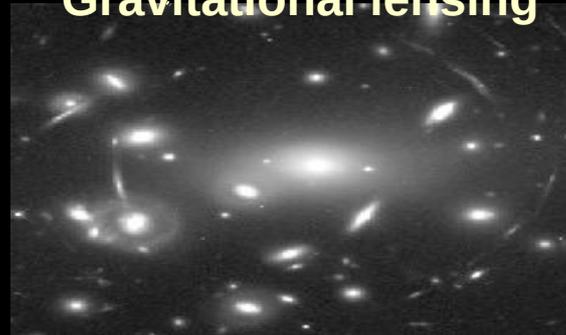
CMB: WMAP and PLANCK



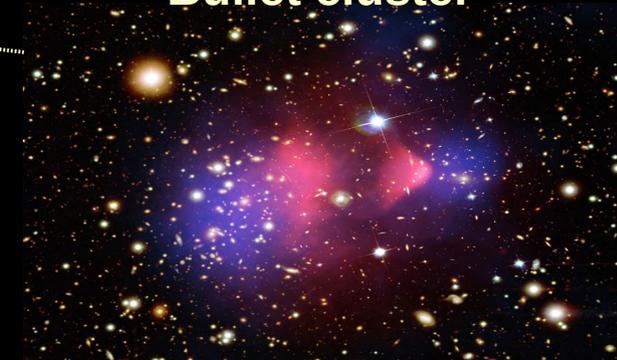
Large Scale Structures



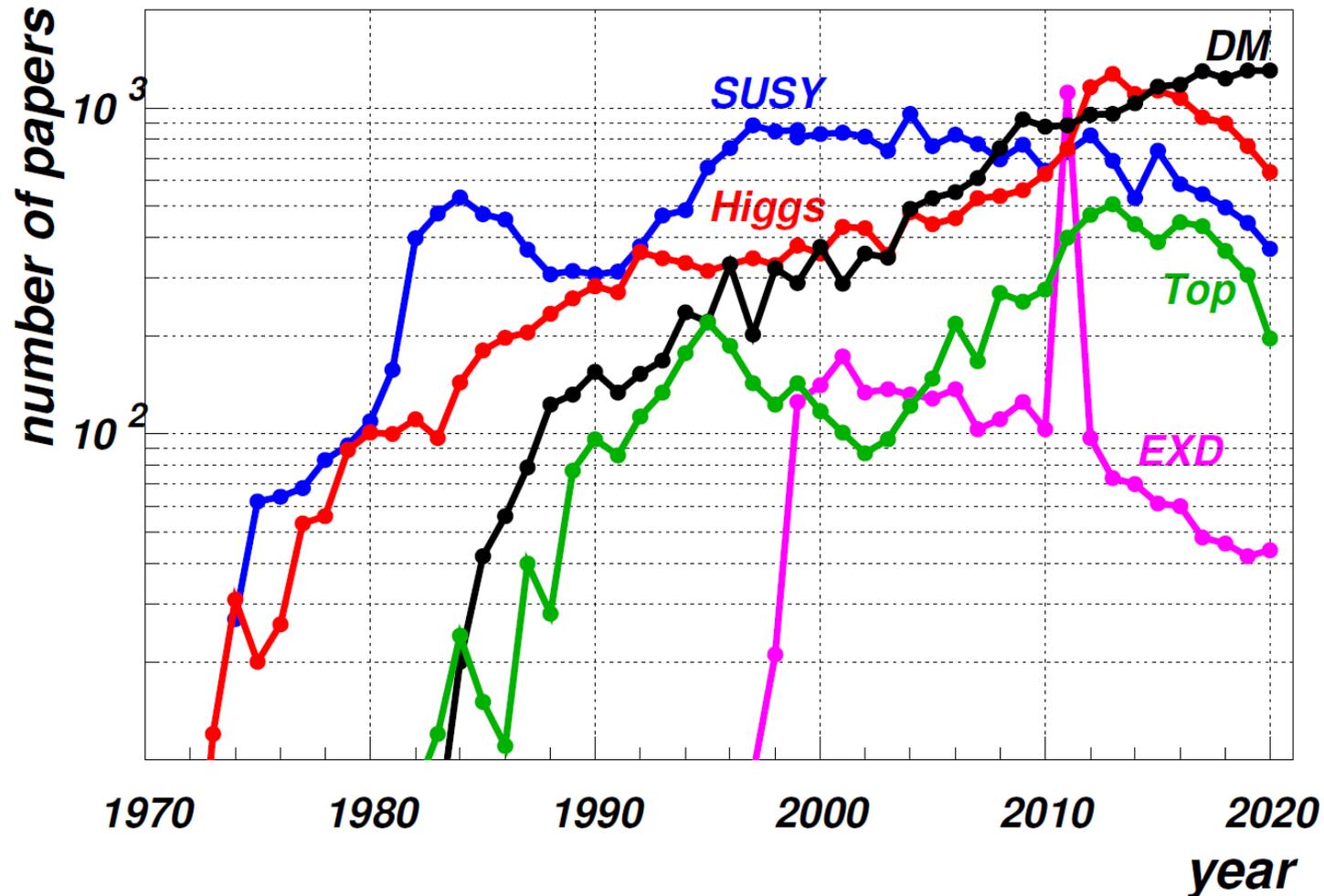
Gravitational lensing



Bullet cluster



The evidence for Dark Matter is a very appealing argument for BSM



DM is very appealing even though we know almost nothing about it!

Spin

Mass

Stable

Yes

No

symmetry behind
stability

Couplings

gravity

weak

higgs

quarks/gluons

leptons

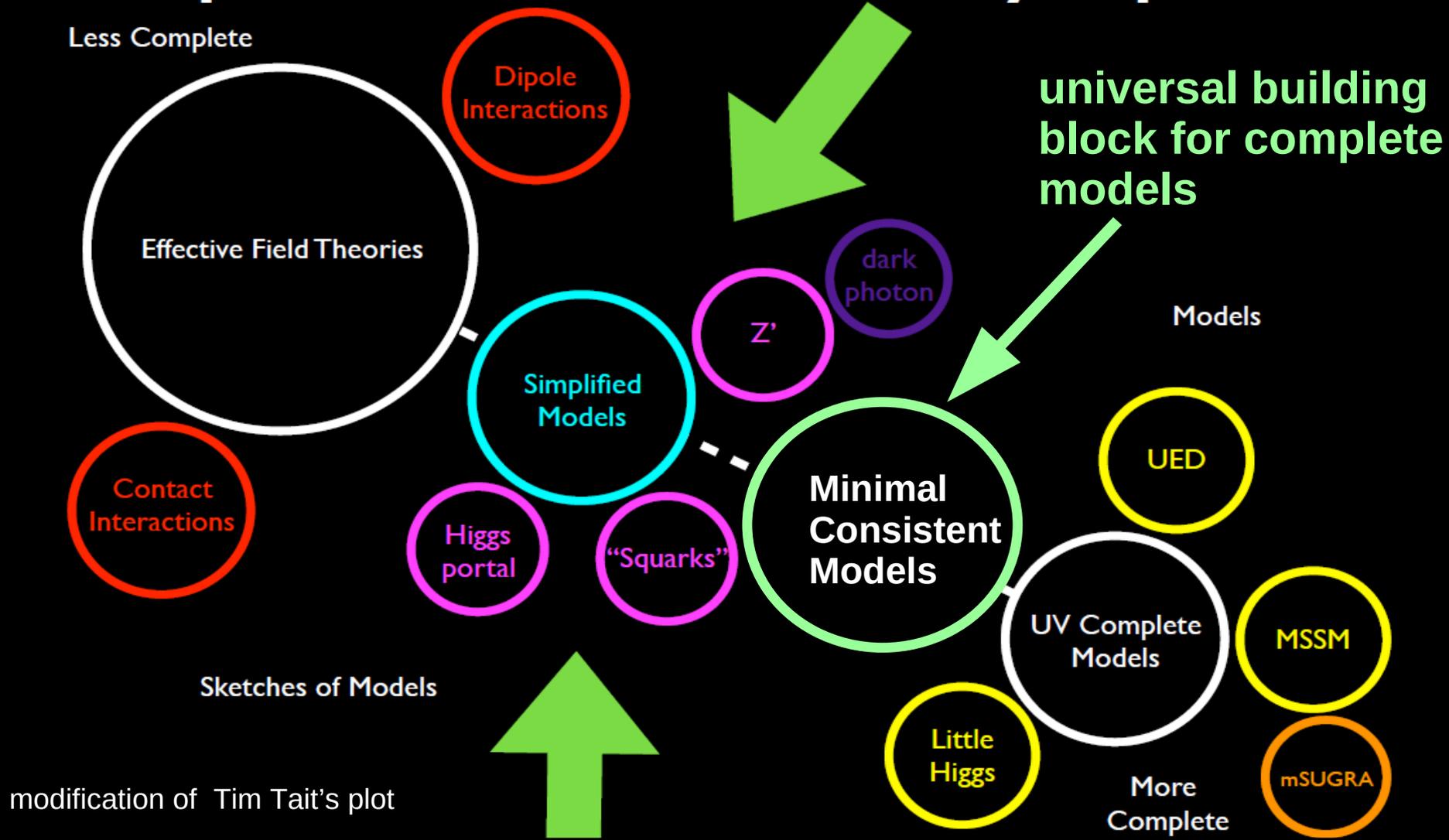
New mediators

Thermal relic

Yes

No

Spectrum of Theory Space



Minimal Consistent DM (MCDM) Models

Properties

- gauge-invariant
- renormalisable
- anomaly-free
- can also be a building block of a bigger theory (e.g. SUSY)

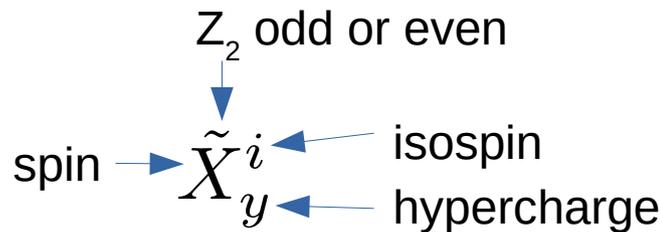
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Classification is important for systematic DM exploration

- DM is a part of EW multiplet
- at most one mediator multiplet
- very important for consistent exploration of DM theory space



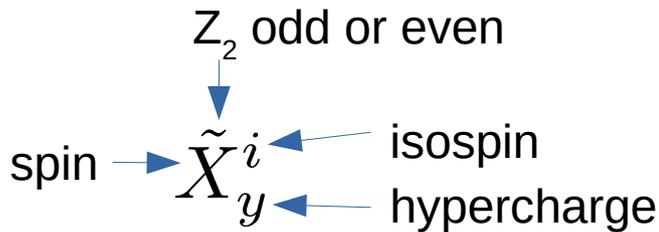
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Spin of Dark Matter \ Spin of Mediator	0	1/2	1
spin 0 even mediator	$\tilde{S}_Y^I S_{Y'}^{I'}$	$\tilde{F}_Y^I S_0^{I'}$	$\tilde{V}_Y^I S_{Y'}^{I'}$
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G.Cacciapaglia, D.Locke, AB to appear

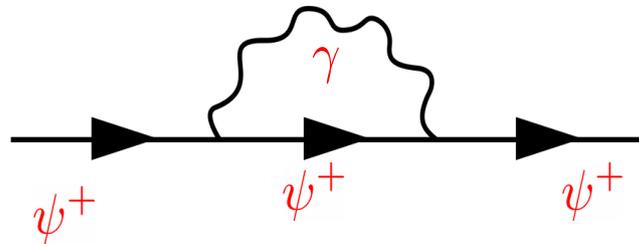
DM multiplet only

$$\mathcal{L} = i\bar{\psi}\gamma^\mu D_\mu\psi - m_D\bar{\psi}\psi$$

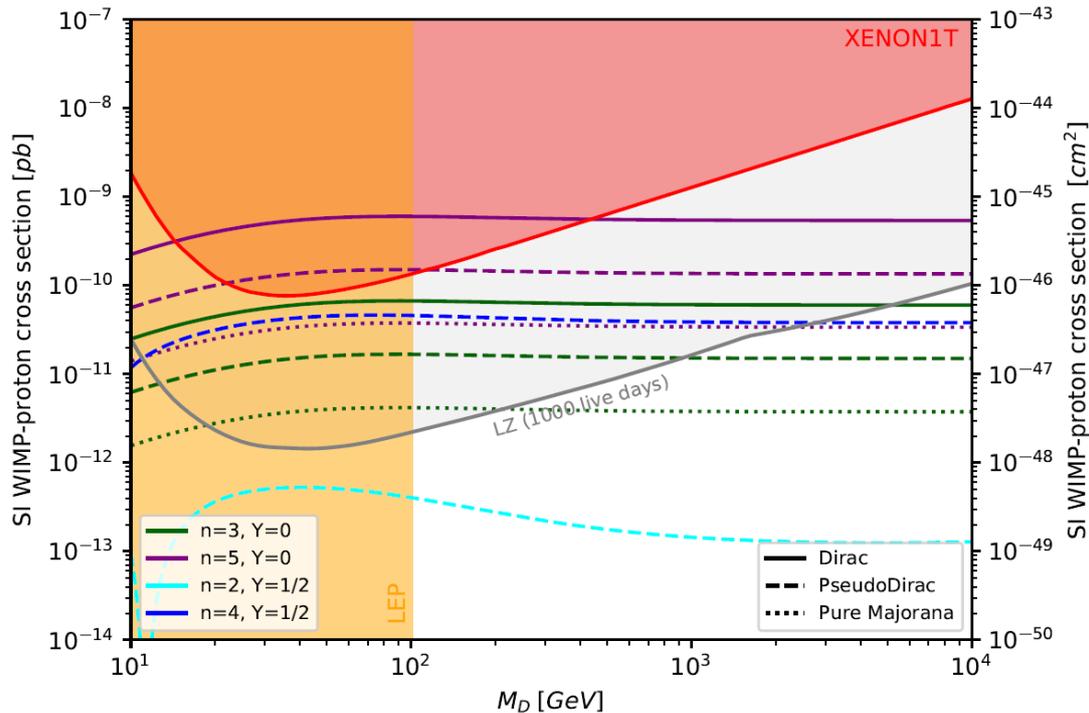
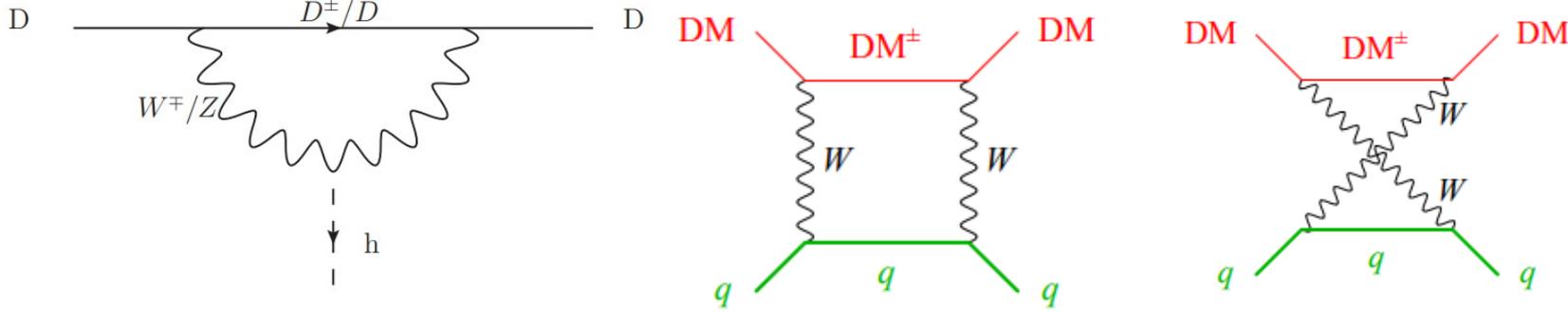
Cirelli,Fornengo,Strumia hep-ph/0512090 (Minimal Dark Matter)

$$\psi = \begin{pmatrix} \psi^{n+} \\ \vdots \\ \psi^+ \\ \psi_0 \\ \psi^- \\ \vdots \\ \psi^{m-} \end{pmatrix}$$

- $\{I, Y\} = \{0, 0\}, \{1/2, 1/2\}, \{1, 0\}$
- Z_2 forbids yukawa couplings
- $\{0, 0\}$ – no gauge-interactions – invisible to direct detection and collider but over(under) abundant if thermal (non-thermal)
- $Y \neq 0$ (Dirac DM) Is excluded by direct detection or requires additional sector – which splits the mass of ψ
- Radiative mass split – very important for the phenomenology



The role of loops in DM DD

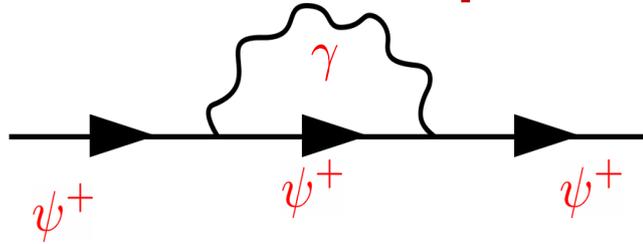


Y=0 minimal candidates may be discovered or ruled out at next generation of DD experiments

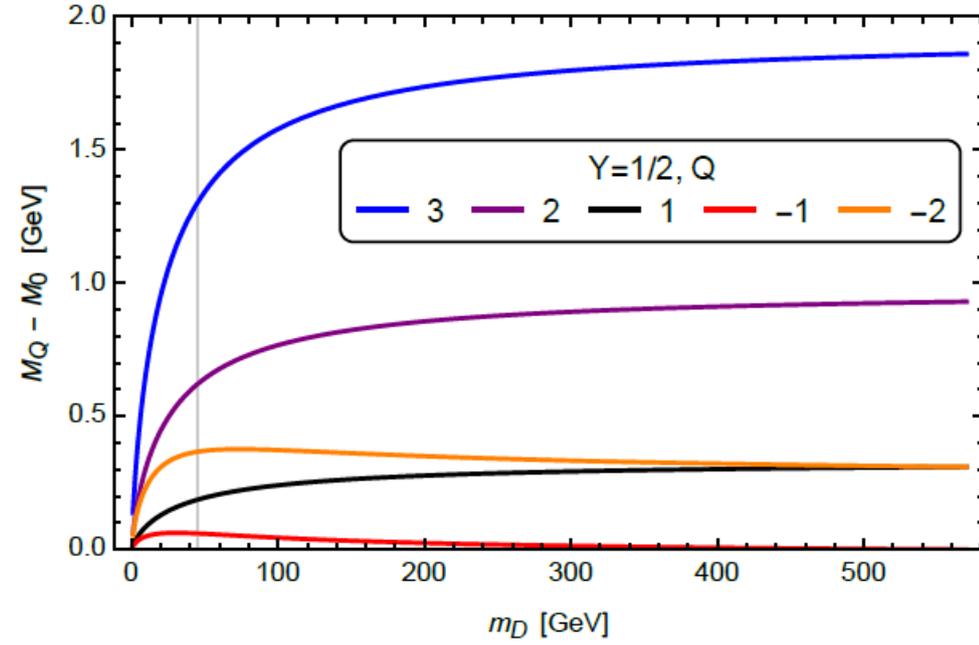
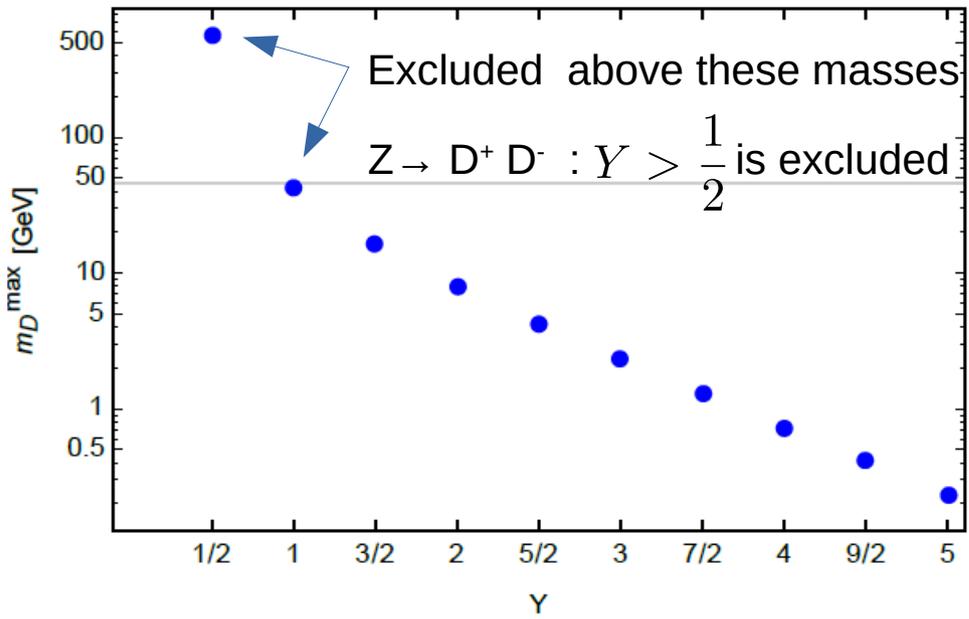
G.Cacciapaglia, D.Locke, AB to appear (preliminary)

Radiative mass Split

simplest models with $Y > 1/2$ are excluded



$$M_Q - M_{Q'}|_{m_D \gg m_W} \approx \frac{\alpha m_W}{2(1 + c_W)} \left[(Q^2 - Q'^2) + \frac{2Y(Q - Q')}{c_W} \right]$$



Left: maximum value of m_D above which the lightest particle has charge $Q = -1$ for various values of Y

Right: spectrum for a generic multiplet with $Y = 1/2$, with $m_D < 570$ GeV.

The vertical line shows $m_D \sim m_Z/2$, below which the model is excluded by the Z decays

Long Lived Particles (LLPs)

- LLPs appear in the minimal DM models with DM being the part of the EW multiplet: **the radiative mass split** of charged and neutral components is of the order of pion mass
- The **hypercharge of the multiplet**
 - a) should be zero, otherwise the model is excluded by DM DD constraints from Z-boson exchange
 - b) or neutral component (DM) of the multiplet should be split by additional (e.g. Yukawa) interactions, which eliminate DM-DM-Z
 - c) multiplet for non-zero hypercharge can not be large – negatively charged component becomes the lightest particle

$$\begin{pmatrix} D^+ \\ D^0 \\ D^- \end{pmatrix} \longrightarrow \Delta M = M_{D^\pm} - M_{D^0} \sim m_\pi$$

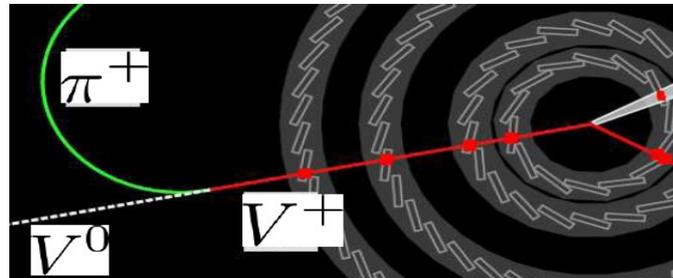
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Cirelli, Fornengo, Strumia 2005 (scalar and fermion DM)

$$\Delta M = \frac{5g_W^2 (M_W - c_W^2 M_Z)}{32\pi}$$

AB, Cacciapaglia, McKay, Marin, Zerwekh 2018 (vector DM)

$D^+ \rightarrow D^0 \pi^+$ is the dominant decay, D^+ is LLP

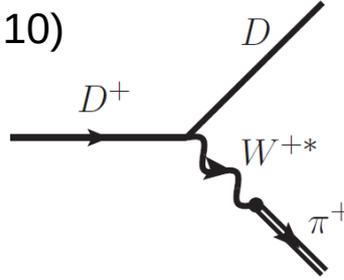


This small mass gap (~ pion mass) provides **disappearing track** signature

D⁺ (charged partner of DM multiplet) decay

- D⁺ life-time should be properly evaluated using **W-pion mixing** (otherwise overestimated by factor of 10)

$$\mathcal{L}_{W\pi} = \frac{gf_\pi}{2\sqrt{2}} W_\mu^+ \partial^\mu \pi^- + \text{h.c.}$$

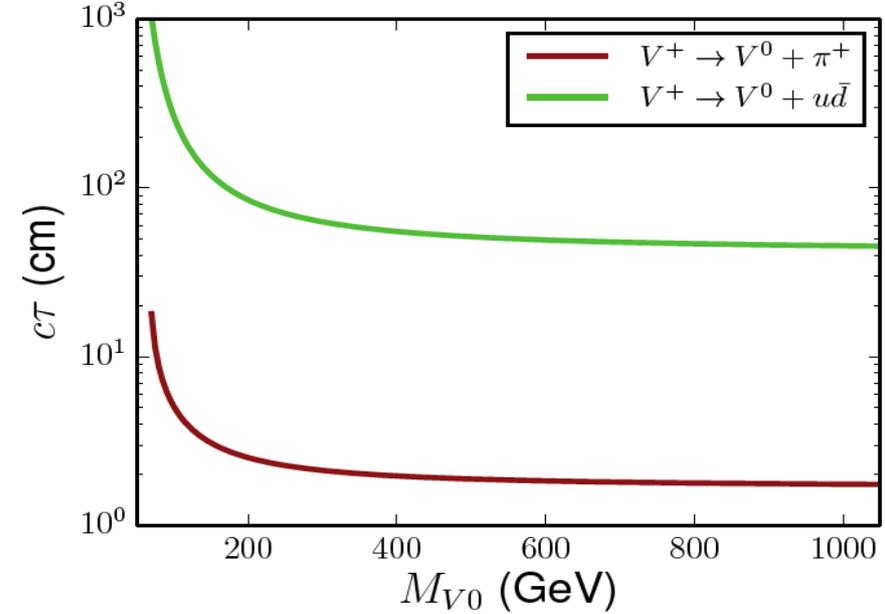


$$\mathcal{L}_{D^+D\pi^-}^{\text{i2HDM}} = -\frac{g^2 f_\pi}{4\sqrt{2}M_W^2} [(p_D - p_D^+) \cdot p_{\pi^-}] D^+ D \pi^- + \text{h.c.}$$

$$\mathcal{L}_{D^+D\pi^-}^{\text{MFDM}} = -\frac{g^2 f_\pi}{4\sqrt{2}M_W^2} \cos(\theta_{DD_3}) p_{\pi^-}^\mu D^+ \gamma^\mu D \pi^- + \text{h.c.}$$

$$\mathcal{L}_{D^+D\pi^-}^{\text{MSSM}} = -\frac{g^2 f_\pi}{4\sqrt{2}M_W^2} p_{\pi^- \mu} D^+ [g_L \gamma^\mu P_L + g_R \gamma^\mu P_R] D \pi^- + \text{h.c.}$$

$$\mathcal{L}_{D^+D\pi^-}^{\text{VDM}} = -\frac{g^2 f_\pi}{2\sqrt{2}M_W^2} [(p_D - p_{D^+})^\mu g^{\nu\rho} - p_D^\nu g^{\mu\rho} + p_{D^+}^\rho g^{\mu\nu}] p_{\pi^- \mu} D_\nu^+ D_\rho \pi^- + \text{h.c.}$$



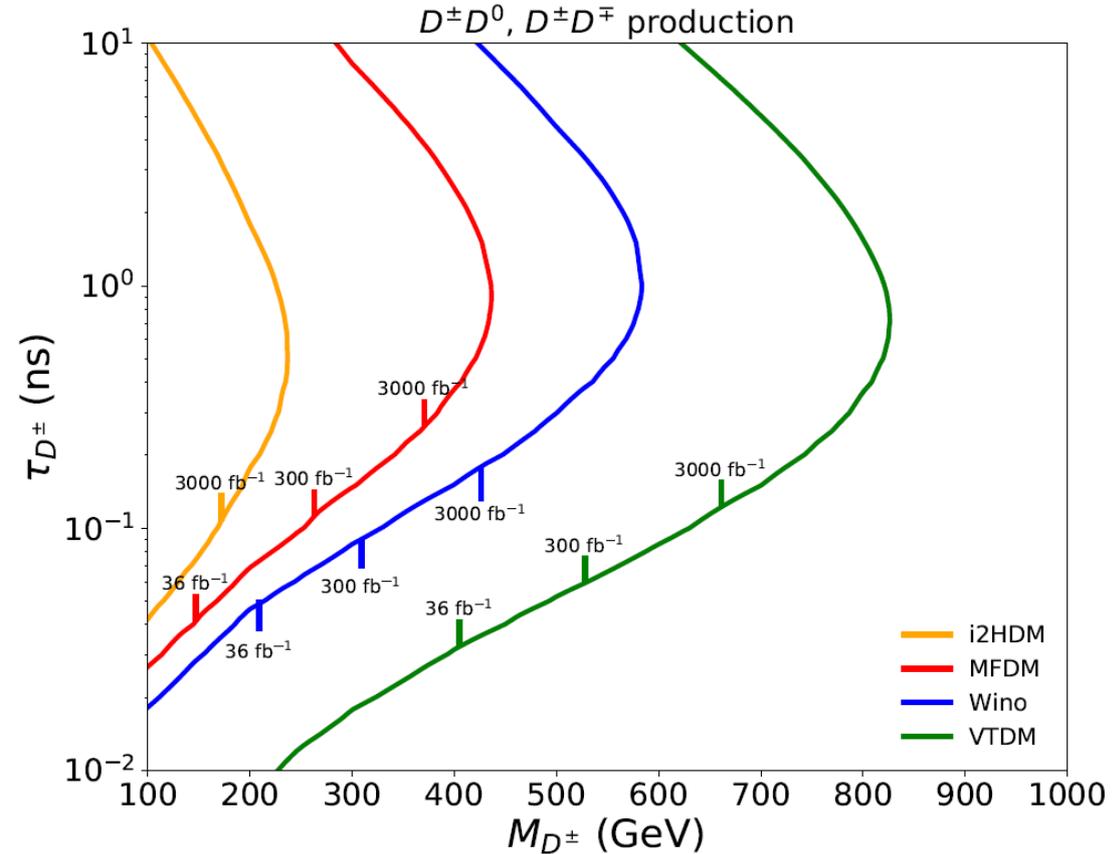
AB, Prestel, Rojas, Zurita [arXiv 2008.08581]

The power of DT for DM probe versus mono-jet limits

- New DT limits for DM models with different spin
- The limits are well beyond those from mono-jet signature analysis for $\tau \sim 1$ ns

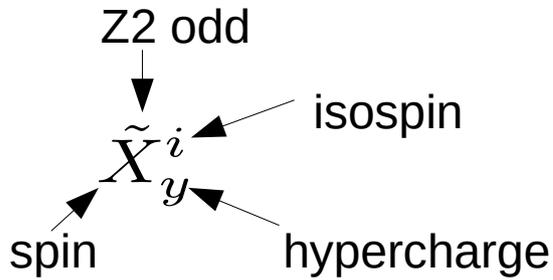
Models	Mass (GeV)	tau (ns)
i2HDM	237	0.5
MFDM	436	0.9
VTDM	822	0.7
WINO	587	1.0

- VTDM \rightarrow MFDM \rightarrow i2HDM hierarchy is defined by CS and PT



AB, Prestel, Rojas, Zurita [arXiv 2008.08581]

$$\tilde{F}_{1/2}^{1/2} \tilde{M}_0^0$$



Minimal fermion DM model

+Mediator

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- **ILC:**
D.Locke,
A.Freegard,
I.Ginzburg, T.Hosken,
A.Pukhov, AB (to appear)
- **Drell Yan, VBF:**
U.Blumenschein,
A.Freegard, S.Moretti, AB
(to appear)

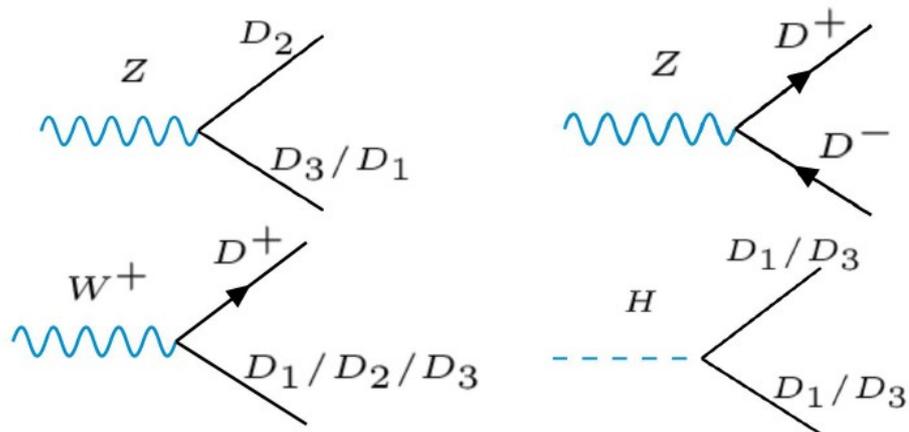
Minimal fermion DM model (MFDM)

gives 2/3 -lepton signatures at the LHC

$$\widetilde{F}_{1/2}^{1/2} \widetilde{M}_0^0$$

$$\mathcal{L}_{FDM} = \mathcal{L}_{SM} + \bar{\psi}(i\not{D} - m_\psi)\psi + \frac{1}{2}\chi_s^0(i\not{D} - m_s)\chi_s^0 - (Y(\bar{\psi}\Phi\chi_s^0) + h.c.)$$

$$\psi = \begin{pmatrix} \chi^+ \\ \frac{1}{\sqrt{2}}(\chi_1^0 + i\chi_2^0) \end{pmatrix}, \text{ additional Majorana singlet fermion } \chi_s^0$$



$$[M_{D1}, \Delta M_{D+}, \Delta M_{D3}]$$

only three parameters (effectively two for the LHC)

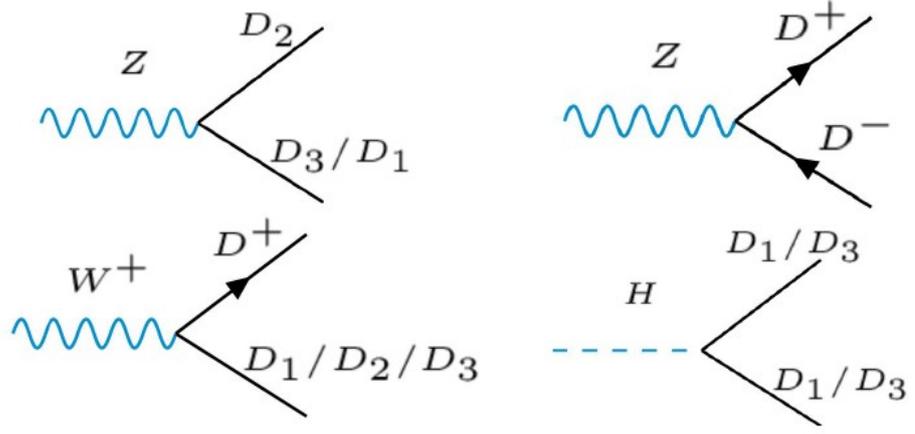
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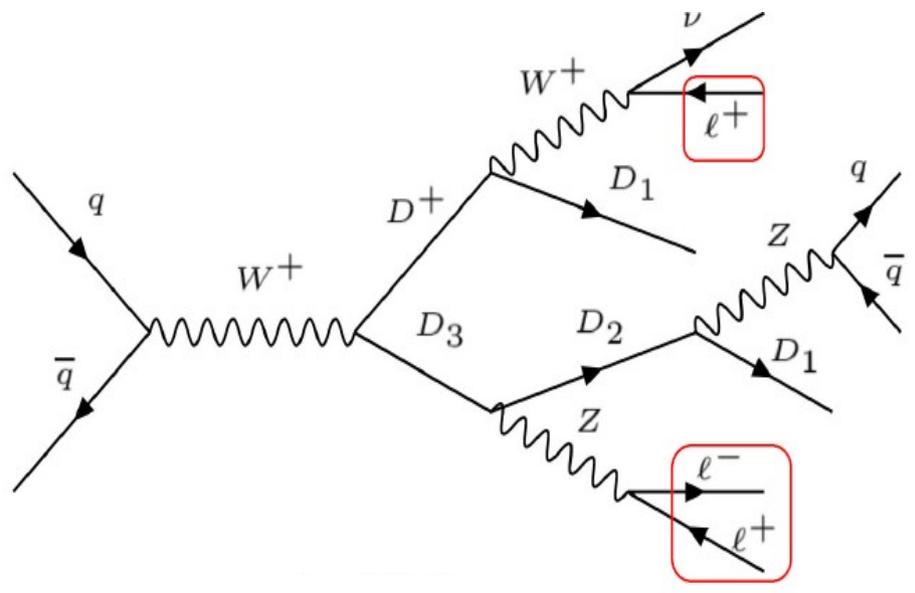
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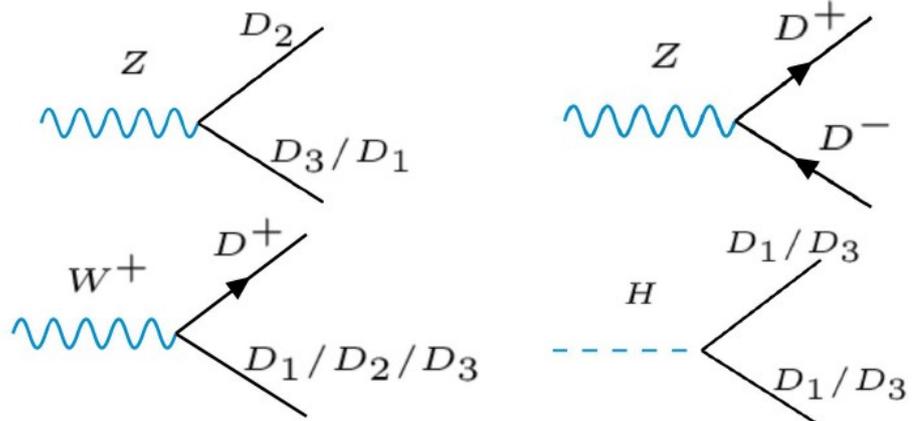
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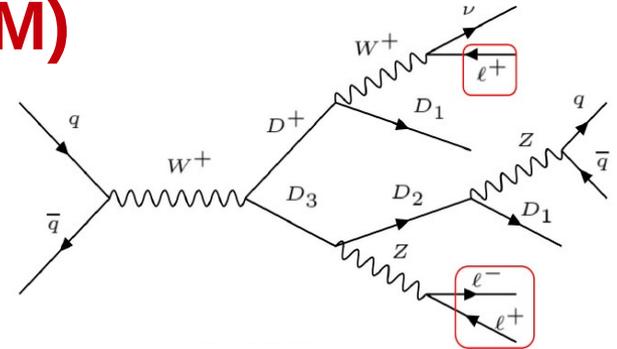
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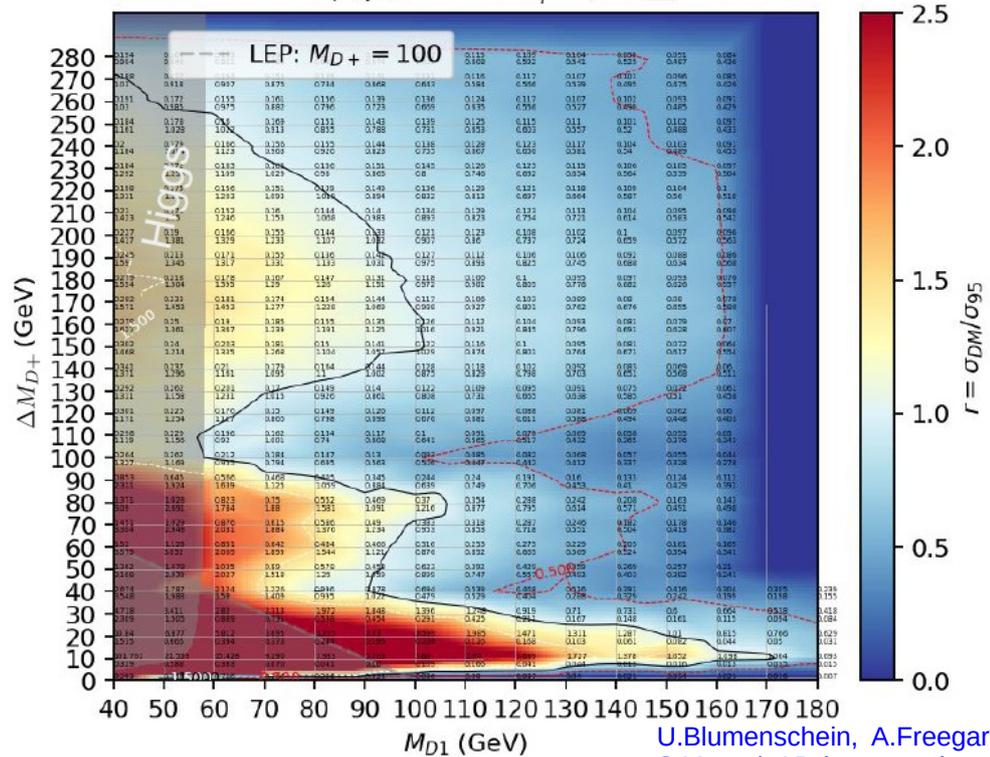


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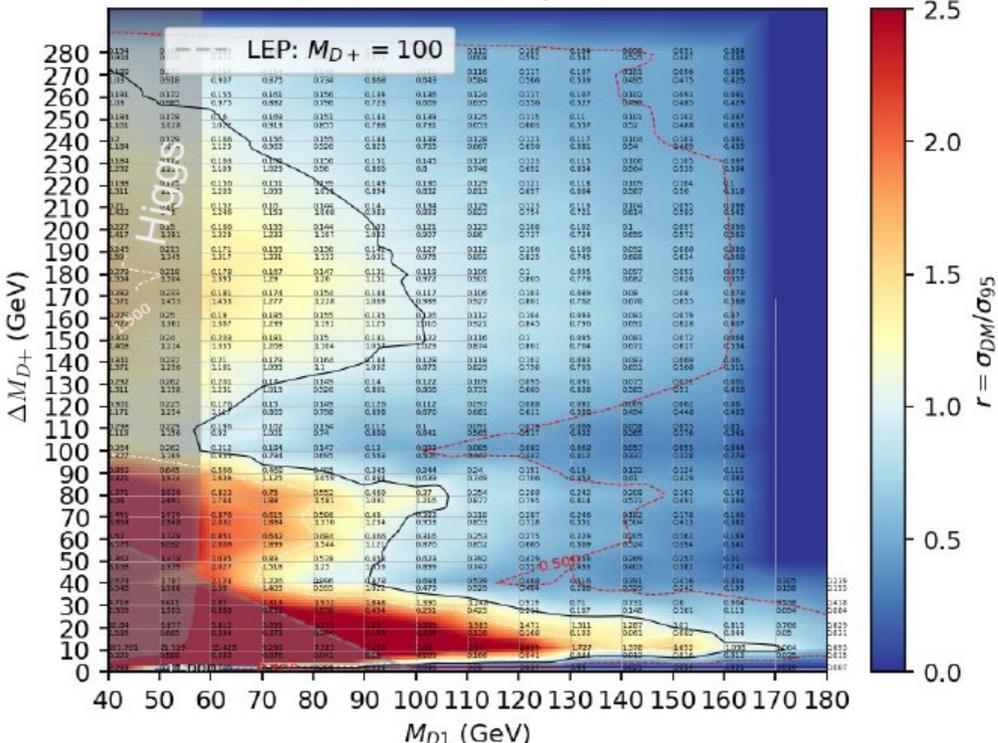
$pp \rightarrow l/2l/3l + DM, DM, \mathcal{L} = 13.3 \text{ fb}^{-1}$
 MFDM 13TeV: $p, p, l^+ l^- + E_T^{miss}, \Delta M_{D3} = 10 \text{ GeV}$



MFDM Results

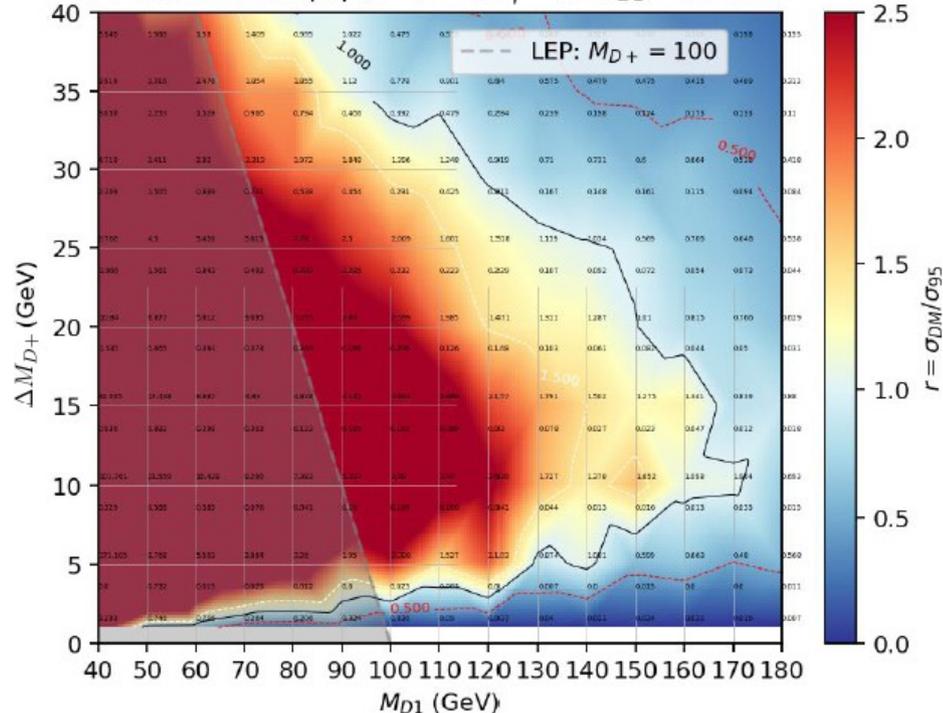
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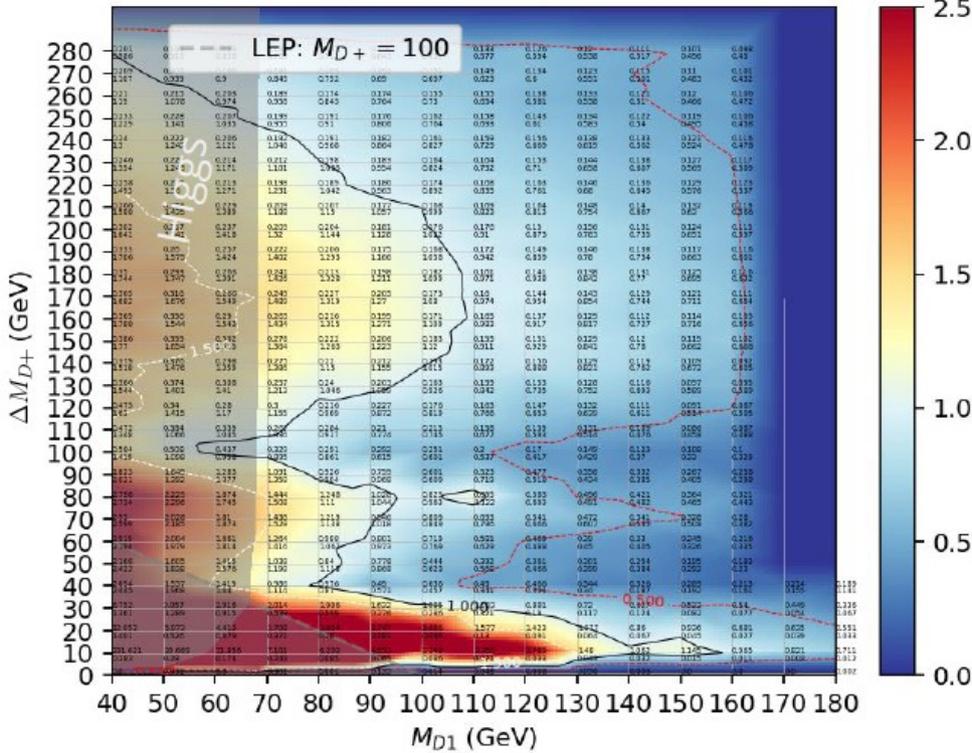


- As ΔM_{D3} increases, coupling between $D_1 - D^\pm$ increases, while heavy D_3 leads to suppressed production cross-section - 'no-lose' theorem

MFDM Results

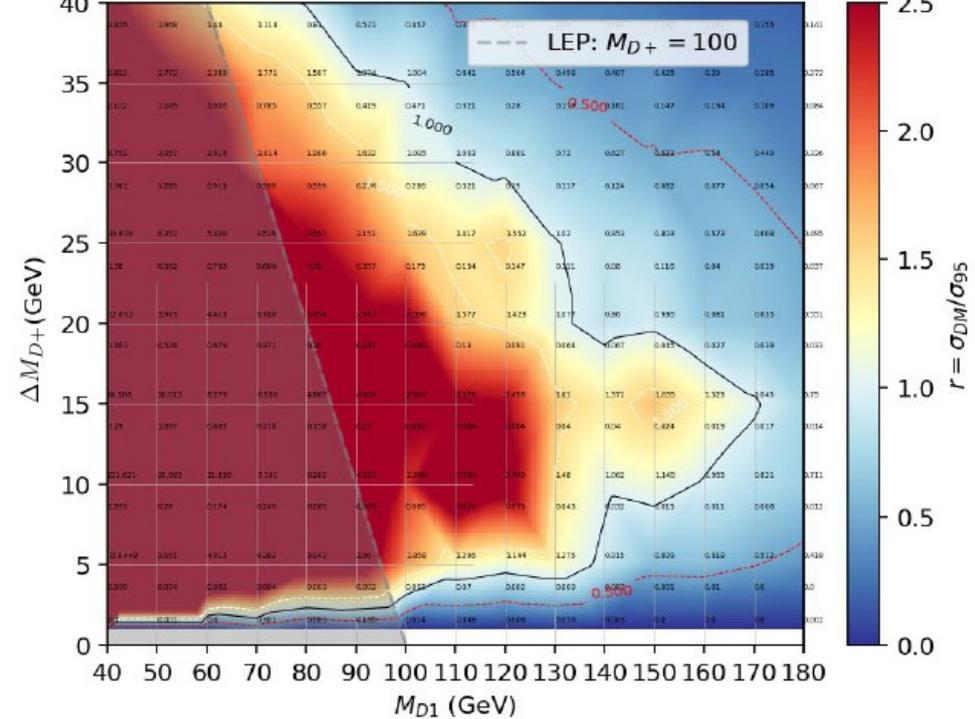
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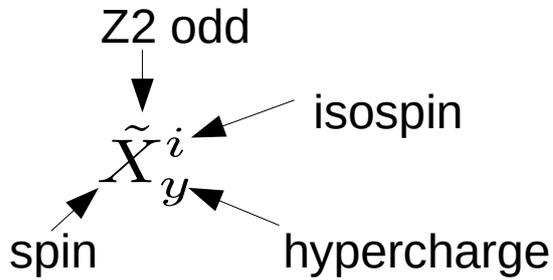
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MFDM 13TeV: $p, p, \ell^+ \ell^- + E_T^{\text{miss}}, \Delta M_{D3} = 100 \text{ GeV}$



- With increasing ΔM_{D3} , Higgs to invisible limit covers larger M_{D1} upto $M_{D1} = M_H/2$

$$\tilde{F}_0^0 S_0^0 (CP - \text{odd})$$



Minimal fermion DM model with pseudo-scalar mediator

new model, has not been explored previously
two-component DM model (pseudoscalar is accidentally stable)

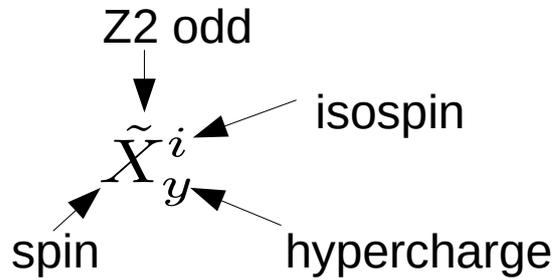
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G.Cacciapaglia, D.Locke, AB arXiv:2104.xxxxx

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$$\mathcal{L} \supset i Y_\psi a \bar{\psi} \gamma^5 \psi - \frac{\lambda_{aH}}{4} |a|^2 \phi_H^\dagger \phi_H$$

Fermion DM Singlet pseudoscalar SM Higgs doublet

- a does not acquire VEV \rightarrow no linear coupling to Higgs
- $m_a < 2m_\psi \rightarrow$ "secluded DM"
- Model implemented in **LanHEP**, and numerical scan performed using **micrOMEGAs**.

4 relevant parameters:

$$m_\psi, Y_\psi, m_a, \lambda_{aH}$$

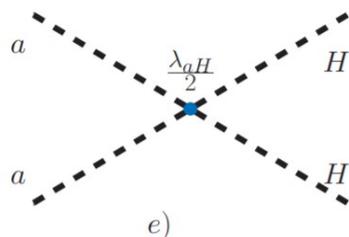
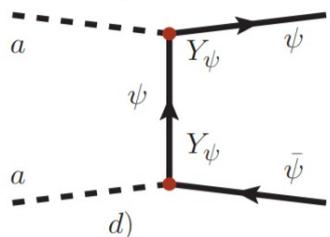
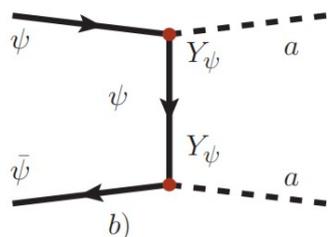
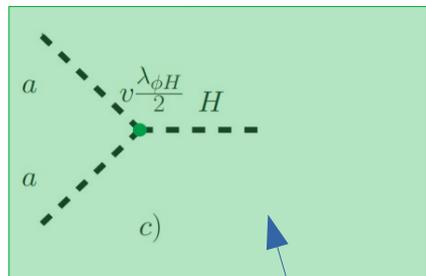
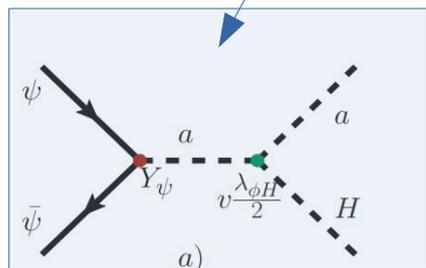
Spin of Dark Matter \ Spin of Mediator	0	1/2	1
spin 0 even mediator	$\tilde{S}_Y^I S_{Y'}^{I'}$	$\tilde{F}_Y^I S_0^{I'}$	$\tilde{V}_Y^I S_{Y'}^{I'}$
spin 0 odd mediator	$\tilde{S}_Y^I \tilde{S}_{Y'}^{I'}$	$\tilde{F}_Y^I \tilde{S}_{Y'}^{I'}$, $\tilde{F}_Y^I \tilde{S}_{Y'}^{I'c}$	$\tilde{V}_Y^I \tilde{S}_{Y'}^{I'}$
spin 1/2 even mediator			
spin 1/2 odd mediator	$\tilde{S}_Y^I \tilde{F}_{Y'}^{I'}$, $\tilde{S}_Y^I \tilde{F}_{Y'}^{I'c}$	$\tilde{F}_Y^I \tilde{F}_{Y\pm 1/2}^{I\pm 1/2}$	$\tilde{V}_Y^I \tilde{F}_{Y'}^{I'}$, $\tilde{V}_Y^I \tilde{F}_{Y'}^{I'c}$
spin 1 even mediator	$\tilde{S}_Y^I V_0^{I'}$	$\tilde{F}_Y^I V_0^{I'}$	$\tilde{V}_Y^I V_{Y'}^{I'}$
spin 1 odd mediator	$\tilde{S}_Y^I \tilde{V}_{Y'}^{I'}$	$\tilde{F}_Y^I \tilde{V}_{Y'}^{I'}$, $\tilde{F}_Y^I \tilde{V}_{Y'}^{I'c}$	$\tilde{V}_Y^I \tilde{V}_{Y'}^{I'}$

G.Cacciapaglia, D.Locke, AB arXiv:2104.xxxxx

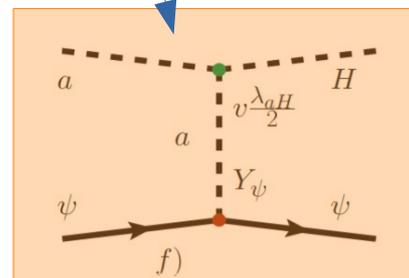
Minimal fermion DM model with pseudo-scalar mediator: rich phenomenology: relic density, DD, colliders

(co)Annihilation channels

$$m_\psi \gtrsim \frac{m_a + m_h}{2} \gtrsim \frac{3m_H}{4} \sim 90\text{GeV}$$



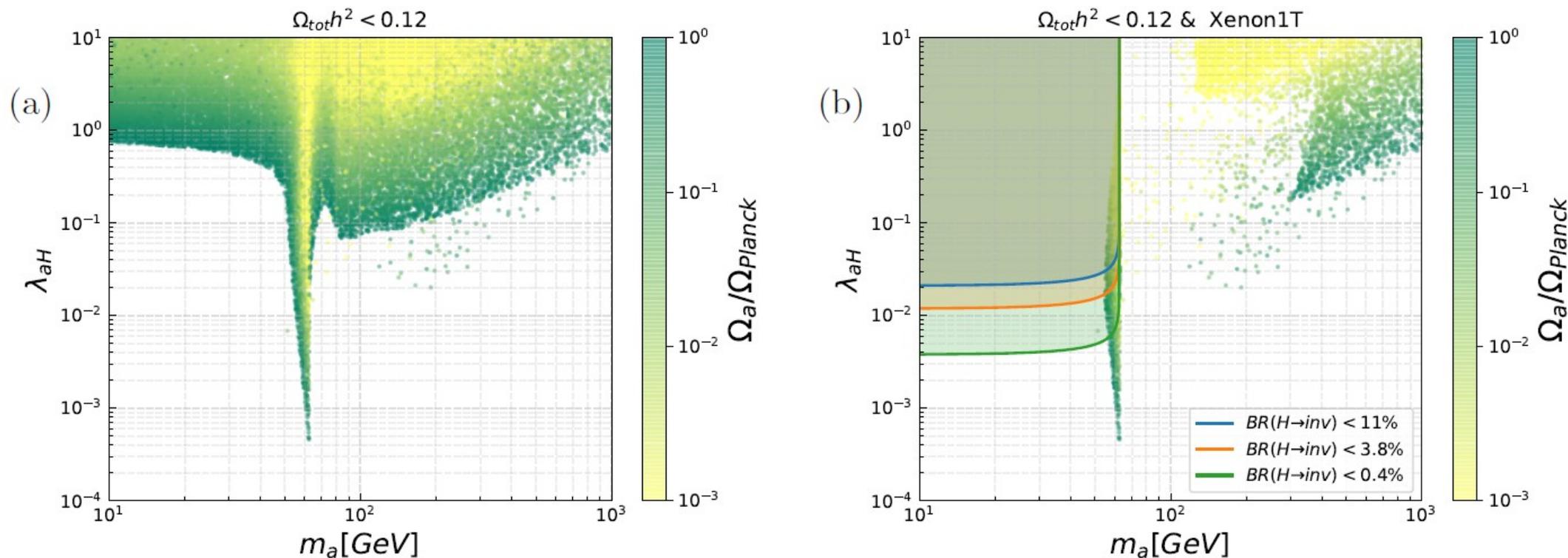
$$m_a \sim m_\psi \gtrsim m_H$$



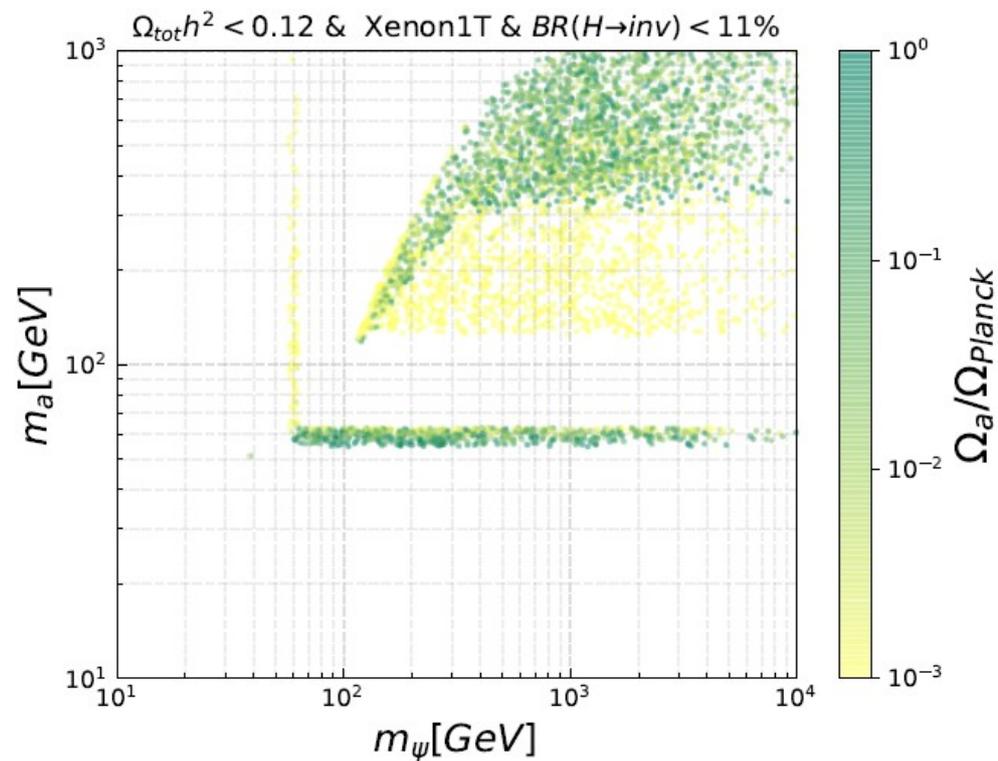
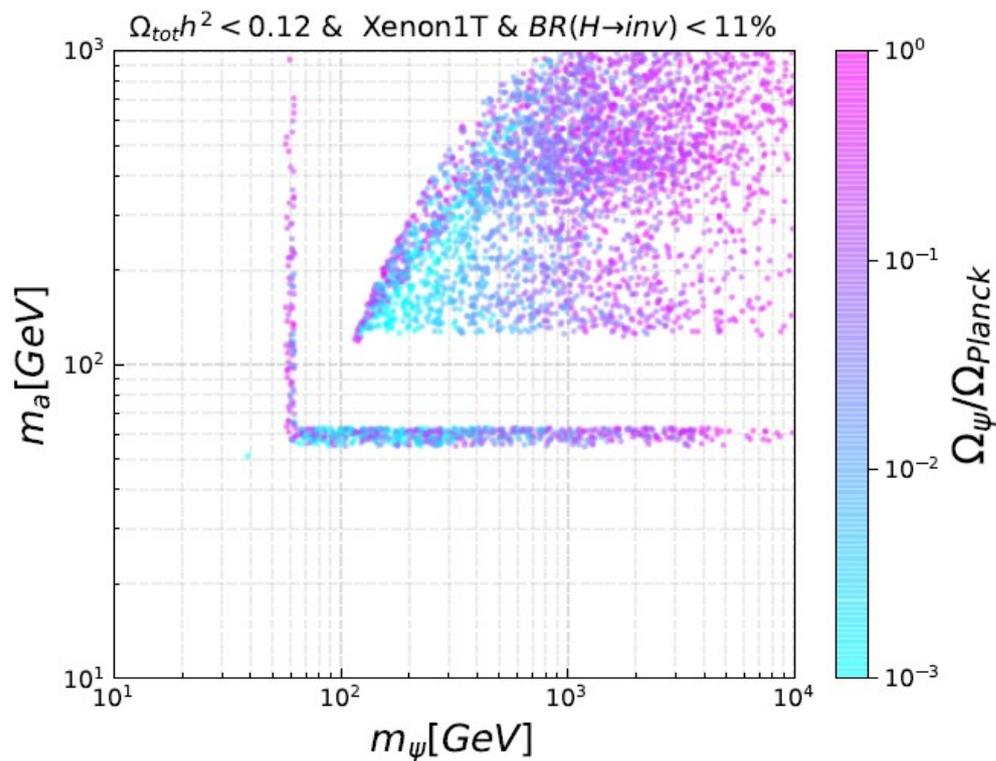
$aa \rightarrow WW$
 $ma > 80\text{GeV}$
 $aa \rightarrow ZZ$
 $ma > 90\text{GeV}$
 $aa \rightarrow tt$
 $ma > 173\text{GeV}$

$$\sigma_{aa \rightarrow ff}^{ann} v \sim \frac{\lambda^2 m_f^2}{(4m_a^2 - m_H^2)^2 + m_H^2 \Gamma_H^2}$$

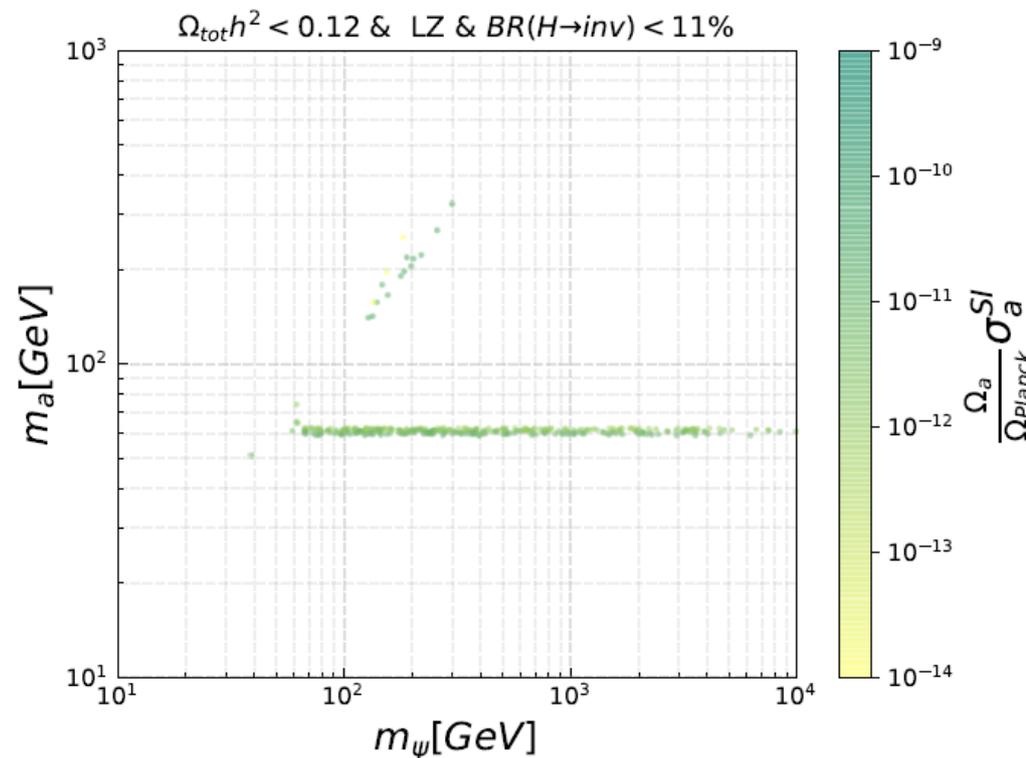
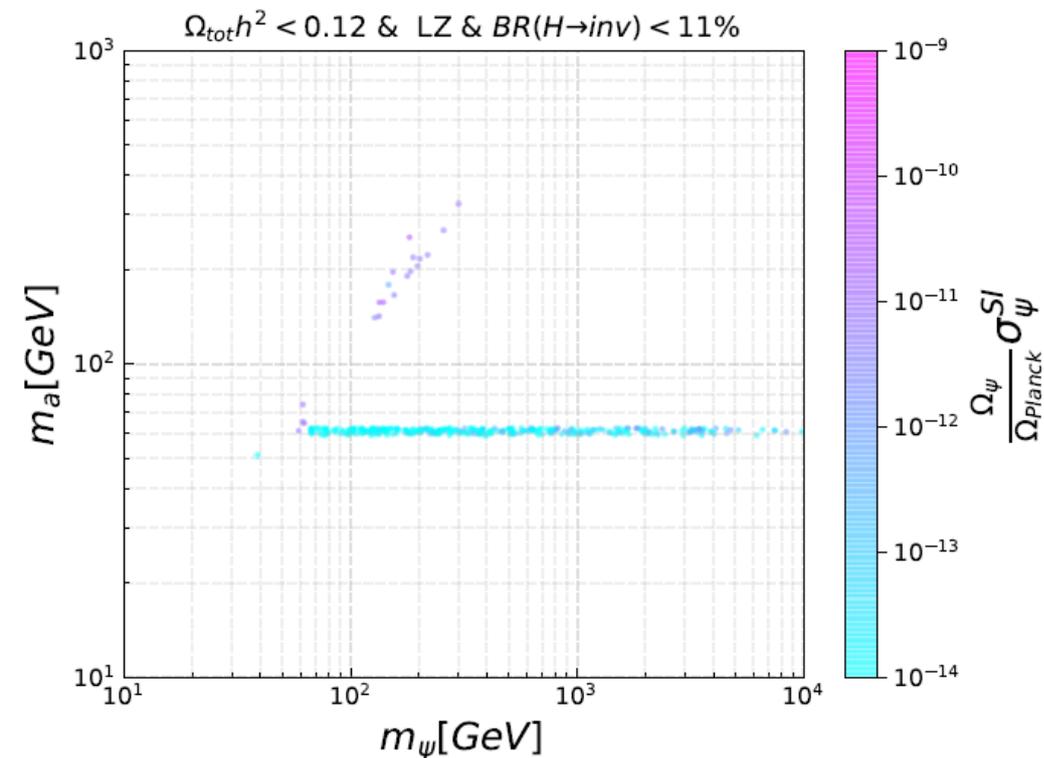
Minimal fermion DM model with pseudo-scalar mediator: rich phenomenology: relic density, DD, colliders



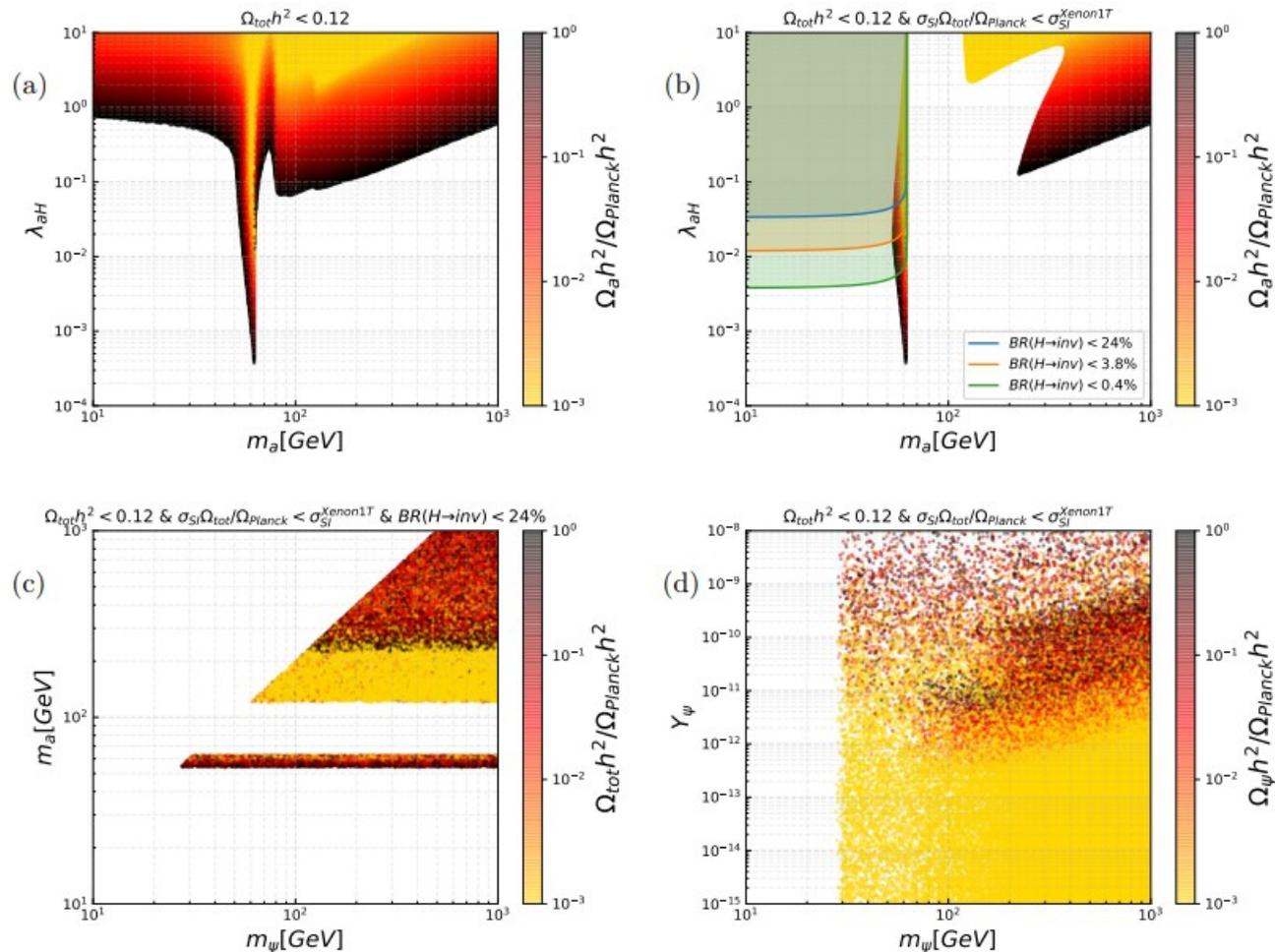
Minimal fermion DM model with pseudo-scalar mediator: Xenon1T vs LZ exclusion



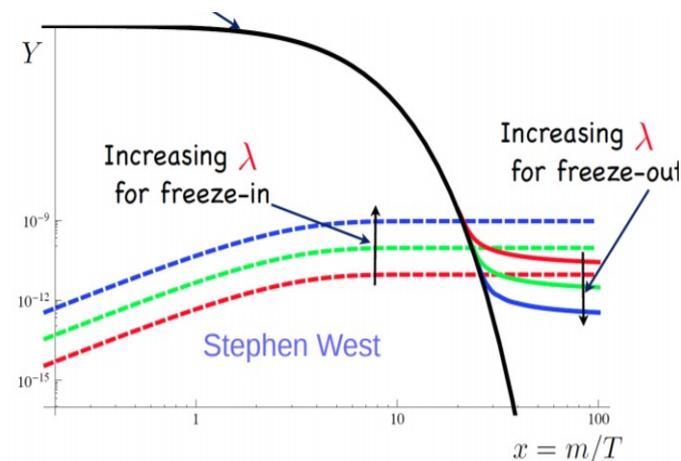
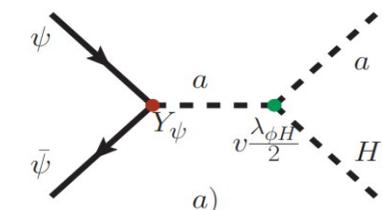
Minimal fermion DM model with pseudo-scalar mediator: Xenon1T vs LZ exclusion



Minimal fermion DM model with pseudo-scalar mediator: non-thermal ψ



$$m_a < 2m_{1/2}$$



Small Y – FDM initially decoupled from SM+a bath, freezes in whilst a freezes out

Conclusions and Outlook

- **To decode the nature of DM** we need a signal first! But at the moment we can
 - understand what kind of DM is already excluded
 - systematically explore theory/parameter space and prepare to DM decoding
- **MCDM models:** consistent but simple – one can explore the entire parameter space
- **Systematic classification** is suggested: new models can be found even for simplest cases, e.g. 2-component DM with pseudoscalar mediator in $\tilde{F}_0^0 S_0^0 (CP - odd)$ model
- **Probing DM space**
 - non-singlets can be probed via DT searches or multi-lepton signatures at colliders.
 - DM DD is sensitive to the loop-induced diagrams but not exclude all models
 - sensitivity is highly dependent on mass-split
 - rich phenomenology
- HEPMDB (hepmdb.soton.ac.uk): the database for BSM models, will include MCDMs

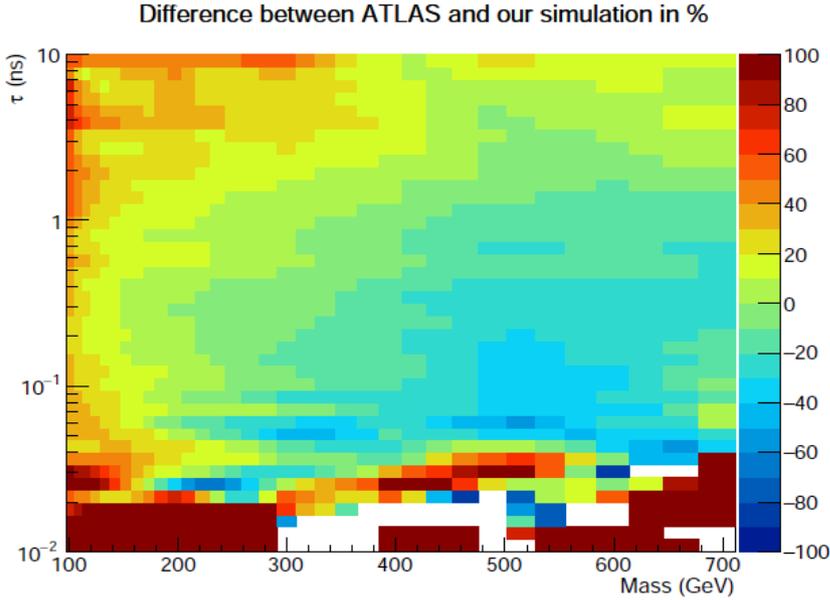
Thank you!

Backup slides

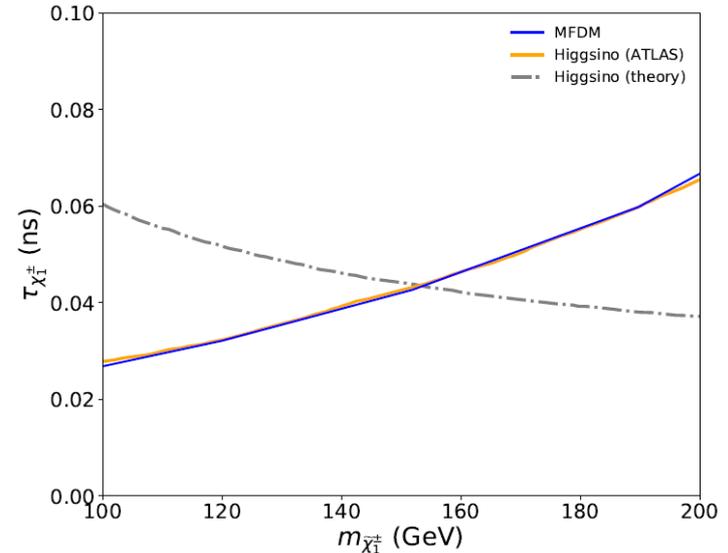
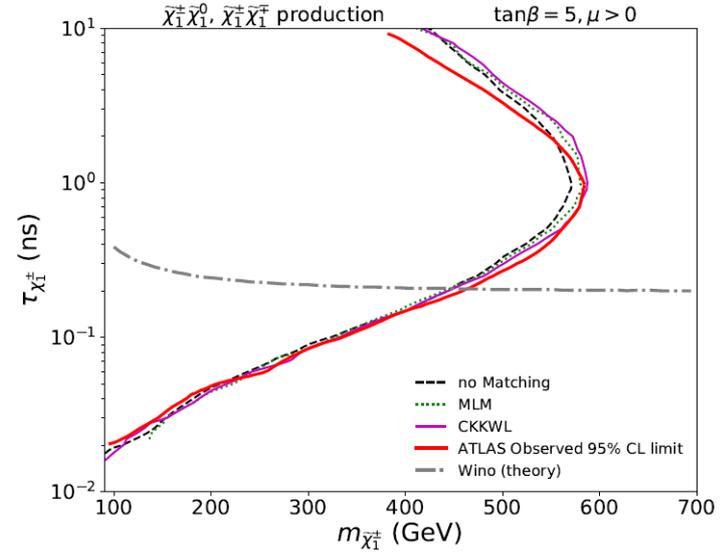
Details on DT studies

Validation of our code/results against ATLAS

- Total acceptance x efficiency check in MDM- τ plane



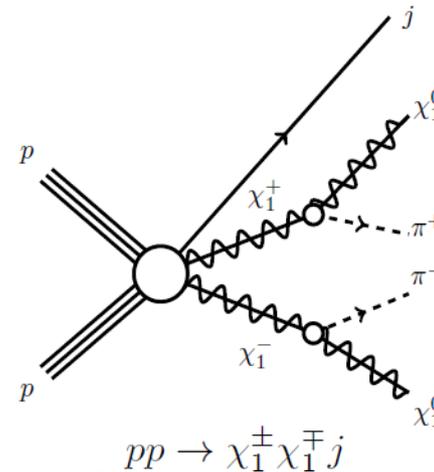
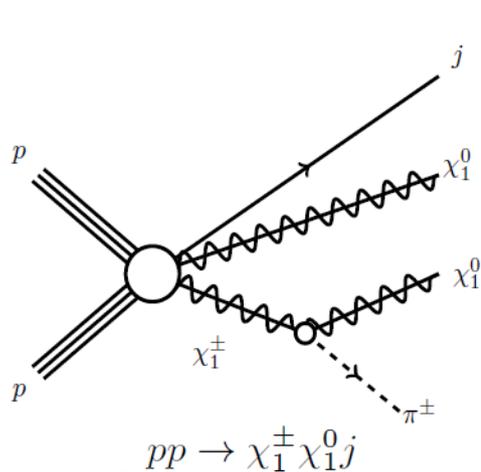
- Compare Limits in MDM- τ plane for pure wino and higgsino models
- The difference in limits on chargino mass is less than 5%



Reinterpretation of the LHC analysis

We apply our validated analysis to minimal consistent models

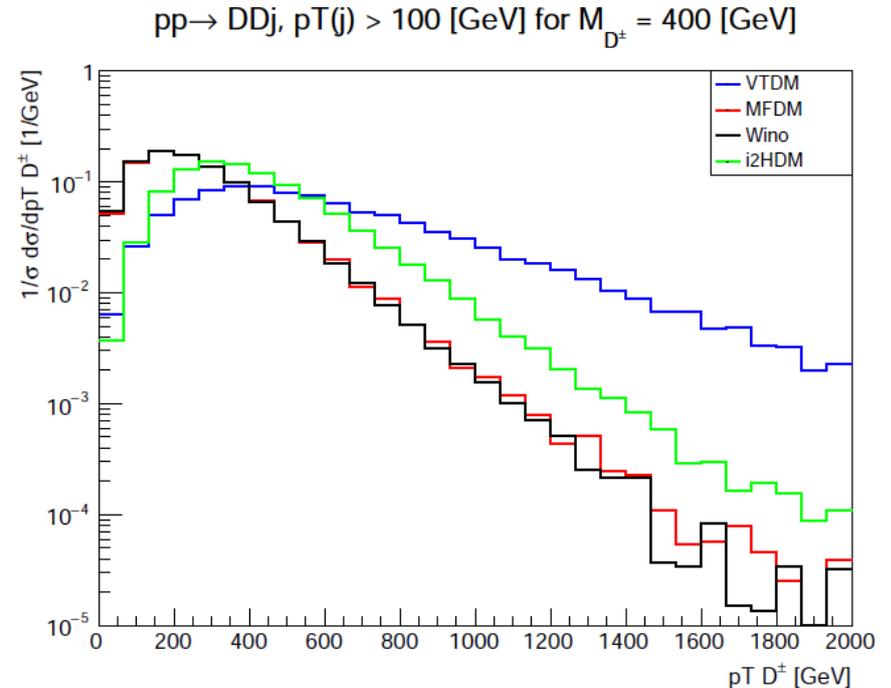
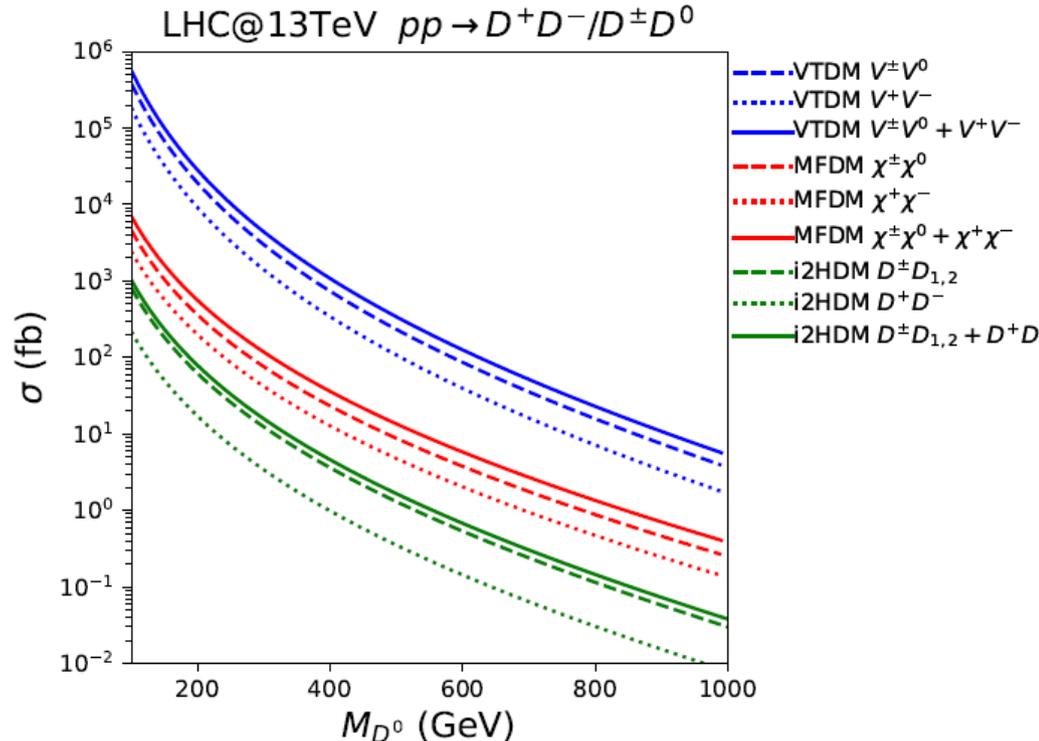
- Scalar: Inert two-Higgs doublet model (i2HDM)
 - Minimal Fermion Dark Matter model (MFDM)
 - Vector: Minimal Vector Triplet Dark Matter model (VTDM)
-
- Two classes of processes: D^+D^- and D^+D^0/D^-D^0 production mediated by s-channel Z/γ and W^+/W^- respectively



Reinterpretation of the LHC analysis

We apply our validated analysis to minimal consistent models

- Scalar: Inert two-Higgs doublet model (i2HDM)
- Minimal Fermion Dark Matter model (MFDM)
- Vector: Minimal Vector Triplet Dark Matter model (VTDM)
- Cross section and Transverse momentum distribution hierarchy: VTDM \rightarrow MFDM \rightarrow i2HDM defines the respective hierarchy of the efficiencies and the LHC sensitivity

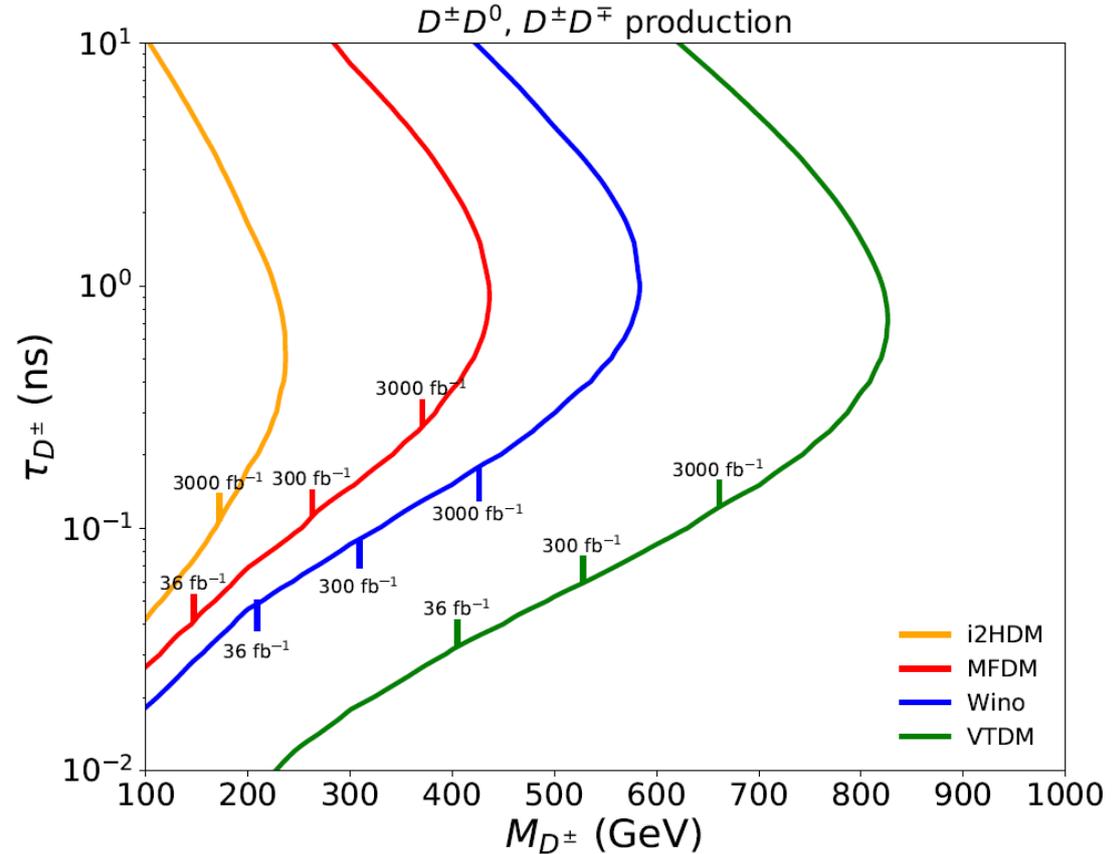


The power of DT for DM probe versus mono-jet limits

- New DT limits for DM models with different spin
- The limits are well beyond those from mono-jet signature analysis for $\tau \sim 1$ ns

Models	Mass (GeV)	tau (ns)
i2HDM	237	0.5
MFDM	436	0.9
VTDM	822	0.7
WINO	587	1.0

- VTDM \rightarrow MFDM \rightarrow i2HDM hierarchy is defined by CS and PT



Public source for the interpretation

- The reinterpretation code is public at <https://github.com/lprecasting/recastingCodes/> [reads root file after LHE → PYTHIA → Delphes simulation]
- [Tables of efficiencies and limits in MDM- \$\tau\$ plane](#) allow to quickly find the reach for your own model in a simple code

tau (ns)	Mass (GeV)						
	100	200	300	400	500	600	700
0.01	1.37e-06	1.90e-07	5.64e-08	1.86e-08	1.17e-08	2.59e-11	2.41e-09
0.02	2.31e-05	9.19e-06	4.13e-06	2.26e-06	1.46e-06	6.29e-07	3.84e-07
0.03	8.67e-05	5.20e-05	3.10e-05	2.06e-05	1.43e-05	8.99e-06	6.72e-06
0.04	1.90e-04	1.43e-04	1.02e-04	7.52e-05	5.61e-05	4.06e-05	3.24e-05
0.05	3.19e-04	2.83e-04	2.27e-04	1.77e-04	1.42e-04	1.10e-04	9.33e-05

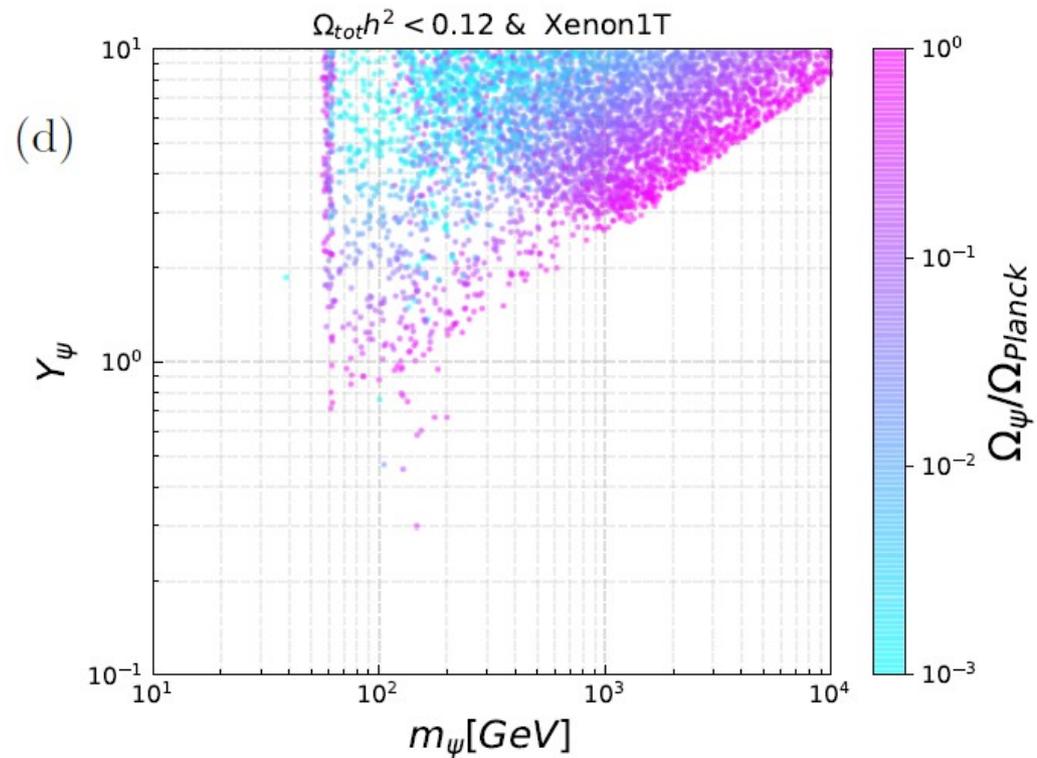
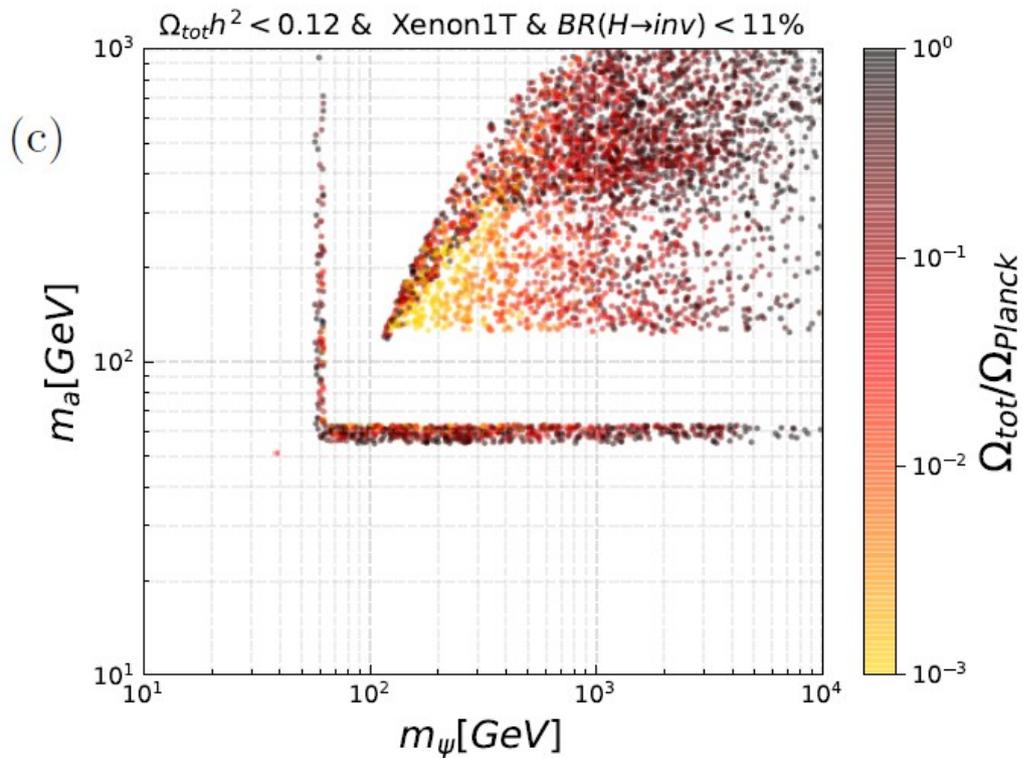
efficiencies

tau (ns)	Mass (GeV)							
	91	200	300	400	500	600	700	800
0.01	968.4	10390	63800	318700	1.44e+06	4.17e+06	2.08e+07	1.993e+09
0.02	187.4	753.3	2580	6434	15530	31210	64850	1.272e+05
0.03	99.06	256.7	649.0	1246	2324	3940	7094	11360
0.04	70.91	142.5	293.7	482.5	768.2	1179	1909	2814
0.05	58.26	97.35	173.7	259.1	377.6	538	797.9	1107
0.06	51.03	74.99	120.8	167.5	227.3	305.9	427.8	568.8

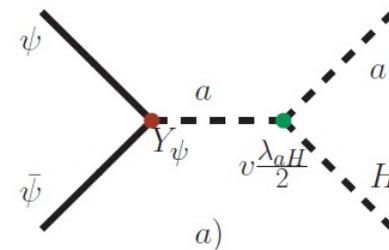
Limits in fb

- available at zenodo <https://zenodo.org/record/4288736> (thanks to Sabine for idea about zenodo)
- efficiencies for [separate channels of \$D^+D^-\$ and \$D^+D^0/D^-D^0\$ production](#) are important for more general interpretation – being produced now (thanks to Felipe Rojas)

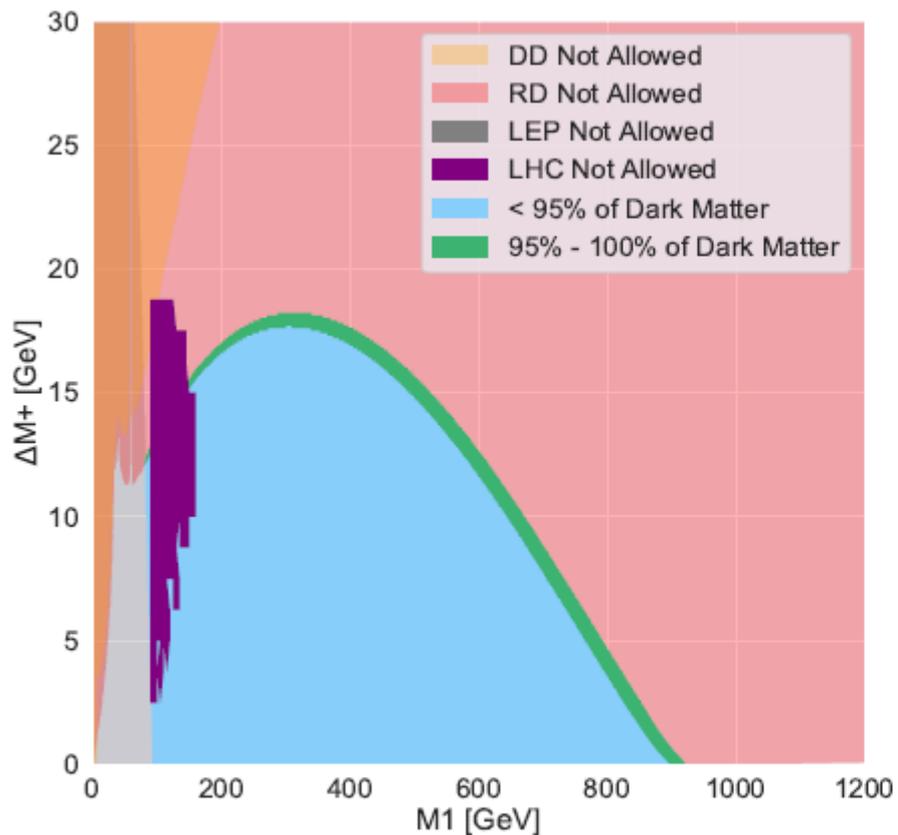
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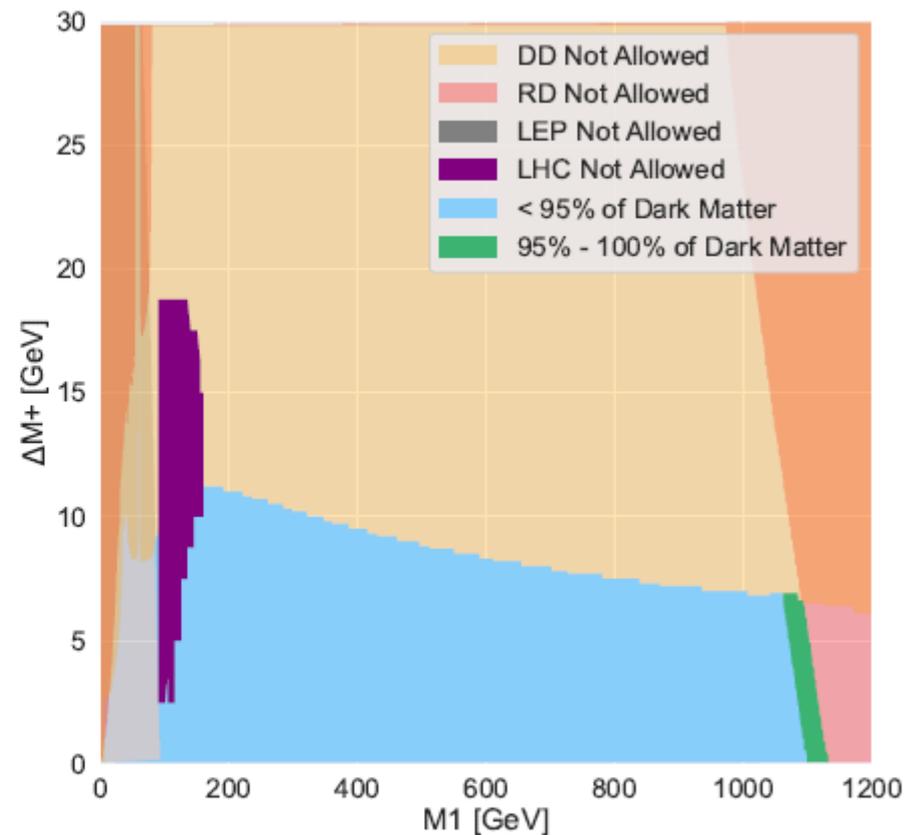
Parameter similar to that of typical portal models, however stable mediator causes “leakage” of points with FDM mass above $\sim 90\text{GeV}$ caused by channel a) opening



MFDM parameter space: the current status



((a)) $\Delta m_3 = 1\text{GeV}$



((b)) $\Delta m_3 = 1000\text{GeV}$