





# Constraining PDFs with $J/\psi \& Y$ inclusive photoproduction

#### Yelyzaveta Yedelkina

#### work done in collaboration with Carlo Flore, Jean-Philippe Lansberg, Hua-Sheng Shao, Alice Colpani Serri, Yu Feng, Melih A. Ozcelik in IJCLab (Orsay)

Kickoff meeting - GDR-QCD / WG1, Quarkonium & QCD meeting, Aussois, June 23, 2021

Inclusive photoproduction of  $J/\psi \& Y$ 

June 23, 2021 1/19

(日) (周) (日) (日)

### Part I

### Introducing $J/\psi$ & Y photoproduction

Y. Yedelkina (IJCLab)

Inclusive photoproduction of  $J/\psi \& Y$ 

**১ ব ট ১ ট ৩**৫৫ June 23, 2021 2/19

Image: A math the second se

C.-H. Chang, NPB172, 425 (1980); R. Baier & R. Rückl Z. Phys. C 19, 251(1983);

In this work we discuss inclusive  $J/\psi(Y)$  photoproduction:

C.-H. Chang, NPB172, 425 (1980); R. Baier & R. Rückl Z. Phys. C 19, 251(1983);

In this work we discuss inclusive  $J/\psi(Y)$  photoproduction:

• as a reminder,  $J/\psi$  (Y) is a  $c\bar{c}$  ( $b\bar{b}$ ) bound state with

J = 1, L = 0, S = 1; vector particle

(日) (周) (日) (日)

C.-H. Chang, NPB172, 425 (1980); R. Baier & R. Rückl Z. Phys. C 19, 251(1983);

In this work we discuss **inclusive**  $J/\psi(Y)$  **photoproduction**:

- as a reminder,  $J/\psi$  (Y) is a  $c\bar{c}$  ( $b\bar{b}$ ) bound state with J = 1, L = 0, S = 1; vector particle
- inclusive photoproduction:

$$\gamma + p \rightarrow J/\psi + X.$$
 (1)

(日) (周) (日) (日)

C.-H. Chang, NPB172, 425 (1980); R. Baier & R. Rückl Z. Phys. C 19, 251(1983);

In this work we discuss **inclusive**  $J/\psi(Y)$  **photoproduction**:

- as a reminder,  $J/\psi$  (Y) is a  $c\bar{c}$  ( $b\bar{b}$ ) bound state with J = 1, L = 0, S = 1; vector particle
- inclusive photoproduction:

$$\gamma + p \rightarrow J/\psi + X.$$
 (1)

(日) (周) (日) (日)

We will discuss the photoproduction at high & low P<sub>T</sub>.

C.-H. Chang, NPB172, 425 (1980); R. Baier & R. Rückl Z. Phys. C 19, 251(1983);

In this work we discuss **inclusive**  $J/\psi(Y)$  **photoproduction**:

- as a reminder,  $J/\psi$  (Y) is a  $c\bar{c}$  ( $b\bar{b}$ ) bound state with J = 1. L = 0. S = 1: vector particle
- inclusive photoproduction:

$$\gamma + p \rightarrow J/\psi + X.$$
 (1)

- We will discuss the photoproduction at high & low P<sub>T</sub>.
- We do not consider in this work:
  - hadroproduction (pp collisions);
  - resolved-photon contributions (low z);
  - UPC (pA and AA collisions);
  - diffractive (large z).

< ロ > < 同 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ >

C.-H. Chang, NPB172, 425 (1980); R. Baier & R. Rückl Z. Phys. C 19, 251(1983);

In this work we discuss inclusive  $J/\psi(Y)$  photoproduction:

- as a reminder,  $J/\psi$  (Y) is a  $c\bar{c}$  ( $b\bar{b}$ ) bound state with J = 1. L = 0. S = 1: vector particle
- inclusive photoproduction:

$$\gamma + p \rightarrow J/\psi + X.$$
 (1)

- We will discuss the photoproduction at high & low P<sub>T</sub>.
- We do not consider in this work:
  - hadroproduction (pp collisions);
  - resolved-photon contributions (low z);
  - UPC (pA and AA collisions);
  - diffractive (large z).

Inclusive photoproduction allows one to probe PDFs.

(日) (周) (日) (日)

C.-H. Chang, NPB172, 425 (1980); R. Baier & R. Rückl Z. Phys. C 19, 251(1983);

**CSM:** we use the Taylor series expansion of the amplitude in the  $Q\overline{Q}$  relative momentum to the first non-vanishing term.

(日) (周) (日) (日)

C.-H. Chang, NPB172, 425 (1980); R. Baier & R. Rückl Z. Phys. C 19, 251(1983);

**CSM:** we use the Taylor series expansion of the amplitude in the  $Q\overline{Q}$  relative momentum to the first non-vanishing term.

One supposes factorisation:

 collinear, in which the hadronic cross section can be written as the convolution of the PDFs with the partonic cross section;



**NB:**The quark and anti-quark attached to the ellipsis are taken as on-shell and their relative velocity v is set to zero.

C.-H. Chang, NPB172, 425 (1980); R. Baier & R. Rückl Z. Phys. C 19, 251(1983);

**CSM:** we use the Taylor series expansion of the amplitude in the  $Q\overline{Q}$  relative momentum to the first non-vanishing term.

#### One supposes factorisation:

- collinear, in which the hadronic cross section can be written as the convolution of the PDFs with the partonic cross section;
- between the hard part (a perturbative amplitude) and the soft part (a non-perturbative matrix element, which describes hadronisation):
  - Perturbative creation of 2 quarks Q and  $\bar{Q}$



**NB:**The quark and anti-quark attached to the ellipsis are taken as on-shell and their relative velocity v is set to zero.

C.-H. Chang, NPB172, 425 (1980); R. Baier & R. Rückl Z. Phys. C 19, 251(1983);

**CSM:** we use the Taylor series expansion of the amplitude in the  $Q\overline{Q}$  relative momentum to the first non-vanishing term.

#### One supposes factorisation:

- collinear, in which the hadronic cross section can be written as the convolution of the PDFs with the partonic cross section;
- between the hard part (a perturbative amplitude) and the soft part (a non-perturbative matrix element, which describes hadronisation):
  - Perturbative creation of 2 quarks Q and  $\bar{Q}$  BUT
    - ► on-shell (×)
    - in a colour singlet state
    - with a vanishing relative momentum
    - in a <sup>3</sup> $S_1$  state (for  $J/\psi$ ,  $\psi'$  and Y)
  - Non-perturbative binding of quarks

NB: The quark and anti-quark attached to the ellipsis are taken as on-shell and their relative velocity v is set to zero.



C.-H. Chang, NPB172, 425 (1980); R. Baier & R. Rückl Z. Phys. C 19, 251(1983);

**CSM:** we use the Taylor series expansion of the amplitude in the  $Q\overline{Q}$  relative momentum to the first non-vanishing term.

#### One supposes factorisation:

- collinear, in which the hadronic cross section can be written as the convolution of the PDFs with the partonic cross section;
- between the hard part (a perturbative amplitude) and the soft part (a non-perturbative matrix element, which describes hadronisation):
  - Perturbative creation of 2 quarks Q and  $\overline{Q}$  BUT
    - ► on-shell (×)
    - in a colour singlet state
    - with a vanishing relative momentum
    - in a <sup>3</sup> $S_1$  state (for  $J/\psi$ ,  $\psi'$  and Y)

#### • Non-perturbative binding of quarks $\rightarrow$ Schrödinger wave function

NB: The quark and anti-quark attached to the ellipsis are taken as on-shell and their relative velocity v is set to zero.

LO  $J/\psi$  photoproduction

araph

Singularities at NLO [and how they are removed]:

**NB1:** At large  $P_T (J/\psi + c)$  & QED contributions can also be important.

A D N A (P) N A B N A B N



Singularities at NLO [and how they are removed]:

- Real emission
  - Infrared divergences: Soft [cancelled by loop Infrared contribution]
  - Infrared divergences: Collinear



Singularities at NLO [and how they are removed]:

- Real emission
  - Infrared divergences: Soft [cancelled by loop Infrared contribution]
  - Infrared divergences: Collinear



Singularities at NLO [and how they are removed]:

- Real emission
  - Infrared divergences: Soft [cancelled by loop Infrared contribution]
  - Infrared divergences: Collinear
    - initial emission [subtracted by AP-CT in the factorised PDFs]
    - final emission [phase-space integration (the KLN theorem)]



Singularities at NLO [and how they are removed]:

- Real emission
  - Infrared divergences: Soft [cancelled by loop Infrared contribution]
  - Infrared divergences: Collinear
    - initial emission [subtracted by AP-CT in the factorised PDFs]
    - final emission [phase-space integration (the KLN theorem)]

Virtual (loop) contribution



Singularities at NLO [and how they are removed]:

- Real emission
  - Infrared divergences: Soft [cancelled by loop Infrared contribution]
  - Infrared divergences: Collinear
    - initial emission [subtracted by AP-CT in the factorised PDFs]
    - final emission [phase-space integration (the KLN theorem)]
- Virtual (loop) contribution
  - Ultraviolet divergences: [removed by renormalisation]
  - Infrared divergences: [cancelled by real Infrared contribution]

A (10) F (10)

### Part II

# Photoproduction at mid and high $P_T$ at HERA

Y. Yedelkina (IJCLab)

Inclusive photoproduction of  $J/\psi \& Y$ 

→ < ∃ > ∃ June 23, 2021 6/19

Image: A match the second s

C. Flore, J.P. Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926



C. Flore, J.P. Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926

#### Notes:

All the computations were done with HELAC-ONIA. The scale and mass uncertainties are shown by the hatched and solid bands. H.S. Shao, CPC198 (2016) 238; See also https://nloaccess.in2p3.fr

[The quark and antiquark attached to the ellipsis are taken as on-shell and their relative velocity v is set to zero.]



C. Flore, J.P. Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926

#### Notes:

All the computations were done with HELAC-ONIA. The scale and mass uncertainties are shown by the hatched and solid bands. H.S. Shao, CPC198 (2016) 238; See also https://nloaccess.in2p3.fr

[The quark and antiquark attached to the ellipsis are taken as on-shell and their relative velocity v is set to zero.]



C. Flore, J.P. Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926

#### Notes:

All the computations were done with HELAC-ONIA. The scale and mass uncertainties are shown by the hatched and solid bands. H.S. Shao, CPC198 (2016) 238; See also https://nloaccess.in2p3.fr

[The quark and antiquark attached to the ellipsis are taken as on-shell and their relative velocity v is set to zero.]

Inclusive photoproduction of  $J/\psi \& Y$ 



C. Flore, J.P. Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926

#### Notes:

All the computations were done with HELAC-ONIA. The scale and mass uncertainties are shown by the hatched and solid bands. H.S. Shao, CPC198 (2016) 238; See also https://nloaccess.in2p3.fr The quark and antiquark attached to the ellipsis are taken as on-shell and their relative velocity vis set to zero.] NLO\* only contains the real-emission contributions with an IR cut-off and is expected to account for the leading  $P_T$ contributions at NLO ( $P_T^{-6}$ ). It has been succesfully checked against full NLO computations for  $P_T > 3$  GeV.

Inclusive photoproduction of  $J/\psi \& Y$ 

C.Flore, JP Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926



A D N A (P) N A B N A B N

C.Flore, JP Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926



• LO QCD : OK at low P<sub>T</sub>

(日)

C.Flore, JP Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926



LO QCD : OK at low P<sub>T</sub>
LO QED small but much harder

(日)

C.Flore, JP Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926



- LO QCD : OK at low P<sub>T</sub>
- LO QED small but much harder
- $J/\psi$ +charm: matter at high  $P_T$

(日) (周) (日) (日)





- LO QCD : OK at low P<sub>T</sub>
- LO QED small but much harder
- $J/\psi$ +charm: matter at high  $P_T$
- NLO<sup>(\*)</sup> close the data, the overall sum nearly agrees with them

(日) (周) (日) (日)





- LO QCD : OK at low P<sub>T</sub>
- LO QED small but much harder
- $J/\psi$ +charm: matter at high  $P_T$
- NLO<sup>(\*)</sup> close the data, the overall sum nearly agrees with them
- Agreement with the last bin when the expected  $B \rightarrow J/\psi$  feed down (in gray) is subtracted





- LO QCD : OK at low P<sub>T</sub>
- LO QED small but much harder
- $J/\psi$ +charm: matter at high  $P_T$
- NLO<sup>(\*)</sup> close the data, the overall sum nearly agrees with them
- Agreement with the last bin when the expected  $B \rightarrow J/\psi$  feed down (in gray) is subtracted

#### The CSM up to $\alpha \alpha_s^3$ reproduces photoproduction at HERA

#### $\rightarrow$ we will restrict to CSM for our EIC predictions

Y. Yedelkina (IJCLab)

Inclusive photoproduction of  $J/\psi \& Y$ 

## Part III

# Photoproduction at mid and high $P_T$ at the Electron-Ion Collider

Y. Yedelkina (IJCLab)

Inclusive photoproduction of  $J/\psi \& Y$ 

June 23, 2021 9/19

C. Flore, J.P. Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926



• At  $\sqrt{s_{ep}} = 45$  GeV, one gets into valence region

(日)

Y. Yedelkina (IJCLab)

June 23, 2021 10/19

C. Flore, J.P. Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926



- At  $\sqrt{s_{ep}} = 45$  GeV, one gets into valence region
- Yield steeply falling with P<sub>T</sub>
- Yield can be measured up to  $P_T \sim 11 \text{ GeV}$  with  $\mathcal{L} = 100 \text{ fb}^{-1}$

[using both ee and  $\mu\mu$  decay channels and  $\varepsilon_{J/\psi}\simeq$  80%]

(日)

C. Flore, J.P. Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926



- At  $\sqrt{s_{ep}} = 45$  GeV, one gets into valence region
- Yield steeply falling with P<sub>T</sub>
- Yield can be measured up to  $P_T \sim 11 \text{ GeV}$  with  $\mathcal{L} = 100 \text{ fb}^{-1}$

[using both *ee* and  $\mu\mu$  decay channels and  $\varepsilon_{J/\psi} \simeq$  80%]

- QED contribution leading at the largest reachable P<sub>T</sub>
- photon-quark fusion contributes more than 30 % for P<sub>T</sub> > 8 GeV
- contamination for probing gluon PDFs

(日) (周) (日) (日)

C. Flore, J.P. Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926



• At  $\sqrt{s_{ep}} = 140$  GeV, larger  $P_T$  range up to approx. 18 GeV

Y. Yedelkina (IJCLab)

(日)

C. Flore, J.P. Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926



- At  $\sqrt{s_{ep}} = 140 \text{ GeV}$ , larger  $P_T$  range up to approx. 18 GeV
- QED contribution also leading at the largest reachable P<sub>T</sub>
- photon-gluon fusion contributions dominant up to approx. 15 GeV

(日)

potential probe of gluon PDFs

C. Flore, J.P. Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926



- At √sep = 140 GeV, larger P<sub>T</sub> range up to approx. 18 GeV
- QED contribution also leading at the largest reachable P<sub>T</sub>
- photon-gluon fusion contributions dominant up to approx. 15 GeV
- potential probe of gluon PDFs
- $J/\psi$  + 2 hard partons [*i.e.*  $J/\psi$  + {*gg*, *qg*, *cc*}] dominant for  $P_T \sim 8 - 15 \text{ GeV}$

(日)

C. Flore, J.P. Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926



- At  $\sqrt{s_{ep}} = 140 \text{ GeV}$ , larger  $P_T$  range up to approx. 18 GeV
- QED contribution also leading at the largest reachable P<sub>T</sub>
- photon-gluon fusion contributions dominant up to approx. 15 GeV
- potential probe of gluon PDFs
- $J/\psi$  + 2 hard partons [*i.e.*  $J/\psi$  + {*gg*, *qg*, *cc*}] dominant for  $P_T \sim 8 - 15 \text{ GeV}$
- It could lead to the observation of  $J/\psi + 2$  jets with moderate  $P_T^{\text{jet}}$

C. Flore, J.P. Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926



- At  $\sqrt{s_{ep}} = 140 \text{ GeV}$ , larger  $P_T$  range up to approx. 18 GeV
- QED contribution also leading at the largest reachable P<sub>T</sub>
- photon-gluon fusion contributions dominant up to approx. 15 GeV
- potential probe of gluon PDFs
- $J/\psi$  + 2 hard partons [*i.e.*  $J/\psi$  + {*gg*, *qg*, *cc*}] dominant for  $P_T \sim 8 - 15 \text{ GeV}$
- It could lead to the observation of  $J/\psi + 2$  jets with moderate  $P_T^{\text{jet}}$

C. Flore, J.P. Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926



- At  $\sqrt{s_{ep}} = 140 \text{ GeV}$ , larger  $P_T$  range up to approx. 18 GeV
- QED contribution also leading at the largest reachable P<sub>T</sub>
- photon-gluon fusion contributions dominant up to approx. 15 GeV
- potential probe of gluon PDFs
- $J/\psi$  + 2 hard partons [*i.e.*  $J/\psi$  + {*gg*, *qg*, *cc*}] dominant for  $P_T \sim 8 - 15 \text{ GeV}$
- It could lead to the observation of  $J/\psi + 2$  jets with moderate  $P_T^{\text{jet}}$

### Part IV

# $J/\psi$ +charm associated production at the EIC

Y. Yedelkina (IJCLab)

Inclusive photoproduction of  $J/\psi \& Y$ 

June 23, 2021 11/19

C. Flore, J.P. Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926



C. Flore, J.P. Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926

- Same LO VFNS computation previously shown in green except for the charm-detection efficiency ε<sub>c</sub>: σ<sup>VFNS</sup> = σ<sup>3FS</sup> × (1 (1 ε)<sup>2</sup>) + (σ<sup>4FS</sup> σ<sup>CT</sup>) × ε
   At √s<sub>ep</sub> = 45 GeV, yield limited to low P<sub>T</sub> even with L = 100 fb<sup>-1</sup>
- But it is clearly observable if  $\epsilon_c = 0.1$  with  $\mathcal{O}(500, 50, 5)$  events for  $\mathcal{L} = (100, 10, 1)$  fb<sup>-1</sup>



C. Flore, J.P. Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926

 Same LO VFNS computation previously shown in green except for the charm-detection efficiency *c*<sub>c</sub>: σ<sup>VFNS</sup> =

 $\sigma^{3FS} \times (1-(1-\epsilon)^2) + (\sigma^{4FS} - \sigma^{CT}) \times \epsilon$ 

- At  $\sqrt{s_{ep}} = 45 \text{ GeV}$ , yield limited to low  $P_T$  even with  $\mathcal{L} = 100 \text{ fb}^{-1}$
- But it is clearly observable if  $\epsilon_c = 0.1$  with  $\mathcal{O}(500, 50, 5)$  events for  $\mathcal{L} = (100, 10, 1)$  fb<sup>-1</sup>
- At  $\sqrt{s_{ep}} = 140 \text{ GeV}$ ,  $P_T$  range up to 10 GeV with up to thousands of events with  $\mathcal{L} = 100 \text{ fb}^{-1}$
- Could be observed via charm jet



C. Flore, J.P. Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926

 Same LO VFNS computation previously shown in green except for the charm-detection efficiency *ε<sub>c</sub>*: *σ<sup>VFNS</sup>* =

 $\sigma^{3FS} \times (1 - (1 - \epsilon)^2) + (\sigma^{4FS} - \sigma^{CT}) \times \epsilon$ 

- At  $\sqrt{s_{ep}} = 45 \text{ GeV}$ , yield limited to low  $P_T$  even with  $\mathcal{L} = 100 \text{ fb}^{-1}$
- But it is clearly observable if  $\epsilon_c = 0.1$  with  $\mathcal{O}(500, 50, 5)$  events for  $\mathcal{L} = (100, 10, 1)$  fb<sup>-1</sup>
- At  $\sqrt{s_{ep}} = 140 \text{ GeV}$ ,  $P_T$  range up to 10 GeV with up to thousands of events with  $\mathcal{L} = 100 \text{ fb}^{-1}$
- Could be observed via charm jet

• 4FS  $\gamma c \rightarrow J/\psi c$  depend on c(x) and could be enhanced by intrinsic charm



C. Flore, J.P. Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926

 Same LO VFNS computation previously shown in green except for the charm-detection efficiency ε<sub>c</sub>: σ<sup>VFNS</sup> =

 $\sigma^{3\textit{FS}} \times (1-(1-\epsilon)^2) + (\sigma^{4\textit{FS}} - \sigma^{\textit{CT}}) \times \epsilon$ 

- At  $\sqrt{s_{ep}} = 45 \text{ GeV}$ , yield limited to low  $P_T$  even with  $\mathcal{L} = 100 \text{ fb}^{-1}$
- But it is clearly observable if  $\epsilon_c = 0.1$  with  $\mathcal{O}(500, 50, 5)$  events for  $\mathcal{L} = (100, 10, 1)$  fb<sup>-1</sup>
- At  $\sqrt{s_{ep}} = 140 \text{ GeV}$ ,  $P_T$  range up to 10 GeV with up to thousands of events with  $\mathcal{L} = 100 \text{ fb}^{-1}$
- Could be observed via charm jet

• 4FS  $\gamma c \rightarrow J/\psi c$  depend on c(x) and could be enhanced by intrinsic charm • Small effect at  $\sqrt{s_{ep}} = 140 \text{ GeV}$  [We used IC c(x) encoded in CT14NNLO]



C. Flore, J.P. Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926

• Same LO VFNS computation previously shown in green except for the charm-detection efficiency  $\epsilon_c: \sigma^{VFNS} =$ 

 $\sigma^{3FS} \times (1 - (1 - \epsilon)^2) + (\sigma^{4FS} - \sigma^{CT}) \times \epsilon$ 

- At  $\sqrt{s_{ep}} = 45 \text{ GeV}$ , yield limited to low  $P_T$  even with  $\mathcal{L} = 100 \text{ fb}^{-1}$
- But it is clearly observable if  $\epsilon_c = 0.1$  with  $\mathcal{O}(500, 50, 5)$  events for  $\mathcal{L} = (100, 10, 1)$  fb<sup>-1</sup>
- At  $\sqrt{s_{ep}} = 140 \text{ GeV}$ ,  $P_T$  range up to 10 GeV with up to thousands of events with  $\mathcal{L} = 100 \text{ fb}^{-1}$
- Could be observed via charm jet

• 4FS  $\gamma c \rightarrow J/\psi c$  depend on c(x) and could be enhanced by intrinsic charm • Small effect at  $\sqrt{s_{ep}} = 140 \text{ GeV}$  [We used IC c(x) encoded in CT14NNLO]



C. Flore, J.P. Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926

 Same LO VFNS computation previously shown in green except for the charm-detection efficiency *ε<sub>c</sub>*: *σ<sup>VFNS</sup>* =

 $\sigma^{3FS} \times (1-(1-\epsilon)^2) + (\sigma^{4FS} - \sigma^{CT}) \times \epsilon$ 

- At  $\sqrt{s_{ep}} = 45 \text{ GeV}$ , yield limited to low  $P_T$  even with  $\mathcal{L} = 100 \text{ fb}^{-1}$
- But it is clearly observable if  $\epsilon_c = 0.1$  with  $\mathcal{O}(500, 50, 5)$  events for  $\mathcal{L} = (100, 10, 1)$  fb<sup>-1</sup>
- At  $\sqrt{s_{ep}} = 140 \text{ GeV}$ ,  $P_T$  range up to 10 GeV with up to thousands of events with  $\mathcal{L} = 100 \text{ fb}^{-1}$
- Could be observed via charm jet

• 4FS  $\gamma c \rightarrow J/\psi c$  depend on c(x) and could be enhanced by intrinsic charm • Small effect at  $\sqrt{s_{ep}} = 140 \text{ GeV}$  [We used IC c(x) encoded in CT14NNLO] • Measurable effect at  $\sqrt{s_{ep}} = 45 \text{ GeV}$ 



C. Flore, J.P. Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926

 Same LO VFNS computation previously shown in green except for the charm-detection efficiency *ε<sub>c</sub>*: *σ<sup>VFNS</sup>* =

 $\sigma^{3FS} \times (1-(1-\epsilon)^2) + (\sigma^{4FS} - \sigma^{CT}) \times \epsilon$ 

- At  $\sqrt{s_{ep}} = 45 \text{ GeV}$ , yield limited to low  $P_T$  even with  $\mathcal{L} = 100 \text{ fb}^{-1}$
- But it is clearly observable if  $\epsilon_c = 0.1$  with  $\mathcal{O}(500, 50, 5)$  events for  $\mathcal{L} = (100, 10, 1)$  fb<sup>-1</sup>
- At  $\sqrt{s_{ep}} = 140 \text{ GeV}$ ,  $P_T$  range up to 10 GeV with up to thousands of events with  $\mathcal{L} = 100 \text{ fb}^{-1}$
- Could be observed via charm jet

• 4FS  $\gamma c \rightarrow J/\psi c$  depend on c(x) and could be enhanced by intrinsic charm • Small effect at  $\sqrt{s_{ep}} = 140 \text{ GeV}$  [We used IC c(x) encoded in CT14NNLO] • Measurable effect at  $\sqrt{s_{ep}} = 45 \text{ GeV}$ 



C. Flore, J.P. Lansberg, H.S. Shao, YY, PLB 811 (2020) 135926

 Same LO VFNS computation previously shown in green except for the charm-detection efficiency ε<sub>c</sub>: σ<sup>VFNS</sup> =

 $\sigma^{3FS} \times (1-(1-\epsilon)^2) + (\sigma^{4FS} - \sigma^{CT}) \times \epsilon$ 

- At  $\sqrt{s_{ep}} = 45 \text{ GeV}$ , yield limited to low  $P_T$  even with  $\mathcal{L} = 100 \text{ fb}^{-1}$
- But it is clearly observable if  $\epsilon_c = 0.1$  with  $\mathcal{O}(500, 50, 5)$  events for  $\mathcal{L} = (100, 10, 1)$  fb<sup>-1</sup>
- At  $\sqrt{s_{ep}} = 140 \text{ GeV}$ ,  $P_T$  range up to 10 GeV with up to thousands of events with  $\mathcal{L} = 100 \text{ fb}^{-1}$
- Could be observed via charm jet

• 4FS  $\gamma c \rightarrow J/\psi c$  depend on c(x) and could be enhanced by intrinsic charm • Small effect at  $\sqrt{s_{ep}} = 140 \text{ GeV}$  [We used IC c(x) encoded in CT14NNLO] • Measurable effect at  $\sqrt{s_{ep}} = 45 \text{ GeV}$ : BHPS valence-like peak visible !

### Part V

# Study of the impact of the NLO corrections to $P_T$ -integrated cross section

#### Negative partonic limit



• The partonic high-energy limit:  $\lim_{\hat{s}\to\infty}(\hat{\sigma}^{NLO})\sim \frac{\alpha_s}{\pi}\hat{\sigma}_0^{LO}(\log\frac{M^2}{\mu_E^2}+\hat{A});$ 

< 回 > < 三 > < 三 >

#### Negative partonic limit



• The partonic high-energy limit:  $\lim_{\hat{s}\to\infty}(\hat{\sigma}^{NLO}) \sim \frac{\alpha_s}{\pi} \hat{\sigma}_0^{LO}(\log \frac{M^2}{\mu_F^2} + \hat{A});$ 

• large 
$$\mu_F \rightarrow \sigma < 0$$
.

< 回 > < 三 > < 三 >

#### $J/\psi$ : $\mu_R \& \mu_F$ choice



• The choice of factorisation scale to avoid possible negative hadronic cross-section:  $\hat{\mu}_F = 0.84M$  (J.P. Lansberg, M.A. Ozcelik: Eur.Phys.J.C 81 (2021) 6, 497;)

A > < > > < >

#### $J/\psi$ : $\mu_R \& \mu_F$ choice



- The choice of factorisation scale to avoid possible negative hadronic cross-section:  $\hat{\mu}_F = 0.84M$  (J.P. Lansberg, M.A. Ozcelik: Eur.Phys.J.C 81 (2021) 6, 497;)
- The natural scale in this case is not 3GeV, rather  $\sim$  5GeV ( $\hat{s}$  for  $\gamma g \rightarrow J/\psi g$ ):  $\mu_B = [2.5; 10] GeV$

A > < > > < >

#### $J/\psi$ : CT14nlo PDF uncertainties of $\sigma(\sqrt{s_{\gamma p}})$



・ 同 ト ・ ヨ ト ・ ヨ

#### $J/\psi$ : CT14nlo PDF uncertainties for $d\sigma/dz$



- We can improve PDFs in the region with low *x*, which corresponds to large *z*.
- PDF uncertainties are increasing with the growth of CM energy.

#### Y photoproduction



• For Y we can improve PDFs in the whole range of CM energy. NB:  $\mu_R$  uncertainties are not visible.

nclusive photoproduction of  $J/\psi \& Y$ 

June 23, 2021 18/19

A D N A (P) N A B N A B N

• For quarkonium production, QCD corrections with  $P_T$ -enhanced topologies are known to be important: we have revisited  $J/\psi$  photoproduction at HERA

イロト イポト イヨト イヨト

- For quarkonium production, QCD corrections with  $P_T$ -enhanced topologies are known to be important: we have revisited  $J/\psi$  photoproduction at HERA
- CSM can describe the latest HERA photoproduction data

イロト イポト イヨト イヨト

- For quarkonium production, QCD corrections with P<sub>T</sub>-enhanced topologies are known to be important: we have revisited J/ψ photoproduction at HERA
- CSM can describe the latest HERA photoproduction data
- I have presented the QCD-correction study to inclusive J/ψ photoproduction at the EIC

イロト イポト イヨト イヨト

- For quarkonium production, QCD corrections with P<sub>T</sub>-enhanced topologies are known to be important: we have revisited J/ψ photoproduction at HERA
- CSM can describe the latest HERA photoproduction data
- I have presented the QCD-correction study to inclusive J/ψ photoproduction at the EIC
- $\sqrt{s_{ep}} = 140$  GeV,
  - gamma-quark QED contribution leading at high P<sub>T</sub>
  - gluon-fusion mostly dominant
  - $J/\psi$ +charm jet accessible
  - $J/\psi$  + 2 jets accessible

- For quarkonium production, QCD corrections with P<sub>T</sub>-enhanced topologies are known to be important: we have revisited J/ψ photoproduction at HERA
- CSM can describe the latest HERA photoproduction data
- I have presented the QCD-correction study to inclusive J/ψ photoproduction at the EIC
- $\sqrt{s_{ep}} = 140$  GeV,
  - gamma-quark QED contribution leading at high P<sub>T</sub>
  - gluon-fusion mostly dominant
  - $J/\psi$ +charm jet accessible
  - $J/\psi$  + 2 jets accessible
- $\sqrt{s_{ep}} = 45$  GeV,
  - gamma-quark QED contribution leading at high P<sub>T</sub>
  - $J/\psi$ +charm sensitive to charm PDFs

- For quarkonium production, QCD corrections with P<sub>T</sub>-enhanced topologies are known to be important: we have revisited J/ψ photoproduction at HERA
- CSM can describe the latest HERA photoproduction data
- I have presented the QCD-correction study to inclusive J/ψ photoproduction at the EIC
- $\sqrt{s_{ep}} = 140$  GeV,
  - gamma-quark QED contribution leading at high P<sub>T</sub>
  - gluon-fusion mostly dominant
  - $J/\psi$ +charm jet accessible
  - $J/\psi$  + 2 jets accessible
- $\sqrt{s_{ep}} = 45$  GeV,
  - gamma-quark QED contribution leading at high P<sub>T</sub>
  - $J/\psi$ +charm sensitive to charm PDFs
- Potential to improve PDFs using  $J/\psi \& Y$  photoproduction.

- For quarkonium production, QCD corrections with P<sub>T</sub>-enhanced topologies are known to be important: we have revisited J/ψ photoproduction at HERA
- CSM can describe the latest HERA photoproduction data
- I have presented the QCD-correction study to inclusive J/ψ photoproduction at the EIC
- $\sqrt{s_{ep}} = 140$  GeV,
  - gamma-quark QED contribution leading at high P<sub>T</sub>
  - gluon-fusion mostly dominant
  - $J/\psi$ +charm jet accessible
  - $J/\psi$  + 2 jets accessible
- $\sqrt{s_{ep}} = 45$  GeV,
  - gamma-quark QED contribution leading at high P<sub>T</sub>
  - $J/\psi$ +charm sensitive to charm PDFs
- Potential to improve PDFs using  $J/\psi \& Y$  photoproduction.
- To avoid negative  $\sigma^{NLO}$  we have used  $\hat{\mu}_F$  scale prescriptions and have shown that it works.

イロト 不得 とうき とうとう ヨー

- For quarkonium production, QCD corrections with P<sub>T</sub>-enhanced topologies are known to be important: we have revisited J/ψ photoproduction at HERA
- CSM can describe the latest HERA photoproduction data
- I have presented the QCD-correction study to inclusive J/ψ photoproduction at the EIC
- $\sqrt{s_{ep}} = 140$  GeV,
  - gamma-quark QED contribution leading at high P<sub>T</sub>
  - gluon-fusion mostly dominant
  - $J/\psi$ +charm jet accessible
  - $J/\psi$  + 2 jets accessible
- $\sqrt{s_{ep}} = 45$  GeV,
  - gamma-quark QED contribution leading at high P<sub>T</sub>
  - $J/\psi$ +charm sensitive to charm PDFs
- Potential to improve PDFs using  $J/\psi \& Y$  photoproduction.
- To avoid negative  $\sigma^{NLO}$  we have used  $\hat{\mu}_F$  scale prescriptions and have shown that it works.

#### Thank you for attention!