## Warm dense QCD matter in strong magnetic fields

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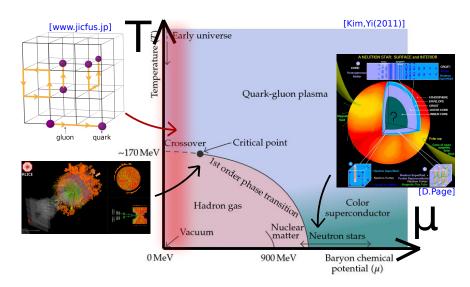
Joint work with Tomáš Brauner, Naoki Yamamoto and Georgios Filios

#### Outline

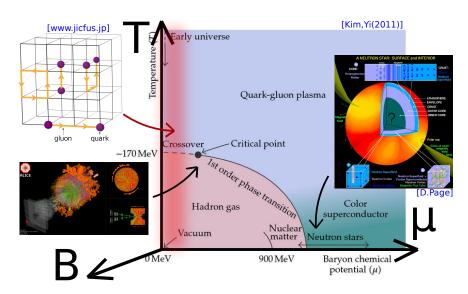
- $\textbf{ Motivation: } \left\langle \begin{array}{c} \mathsf{QCD} \ \mathsf{phase} \ \mathsf{diagram} \\ \mathsf{Inhomogeneous} \ \mathsf{phases} \ \mathsf{of} \ \mathsf{QCD} \ \mathsf{matter} \end{array} \right.$
- Chiral soliton lattice phase at LO
- Ohiral soliton lattice phase at finite temperature
- Chiral soliton lattice phase on lattice?



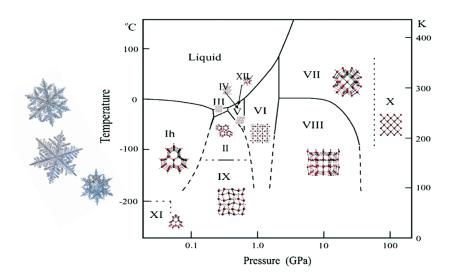
## Motivation: QCD phase diagram



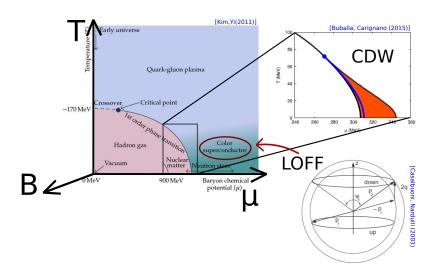
## Motivation: QCD phase diagram



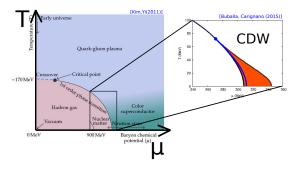
#### Motivation: Inhomogeneous phases of matter



[Alford, Bowers, Rajagopal (2001)], review: [Buballa, Carignano (2015)]

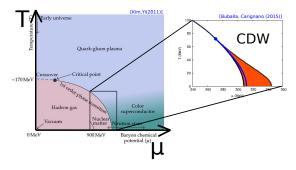


Review: [Buballa, Carignano (2015)]



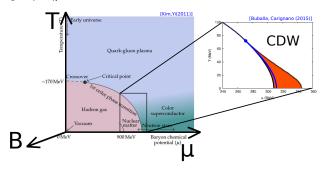
• Chiral Density Wave:  $\langle \overline{\psi}\psi \rangle = \Delta \cos{(\vec{q} \cdot \vec{x})}, \ \langle \overline{\psi}i\gamma_5\tau_3\psi \rangle = \Delta \sin{(\vec{q} \cdot \vec{x})}$ 

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- 1D modulation in spherically symmetric 3D system unstable at any non-zero temperature! [Peierls (1934), Landau]
   (but quasi-long-range order may be sustained [Hidaka et al. (2015)])

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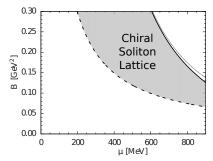


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- If the spherical symmetry broken by magnetic field, 1D modulations stable! [Tatsumi,Nishiyama,Karasawa(2015)] [Ferrer,Incera(2020)]

[Son,Stephanov(2008)][Brauner,Yamamoto(2017)]

Ground state of QCD matter for sufficiently strong magnetic field and large enough baryon chemical potential: inhomogeneous condensate of neutral pions

#### Shown using chiral perturbation theory!

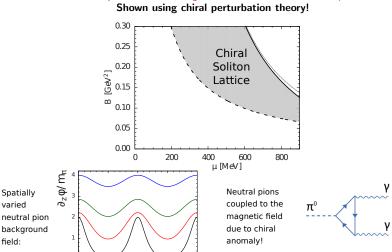


[Son, Stephanov (2008)] [Brauner, Yamamoto (2017)]

varied

field:

Ground state of QCD matter for sufficiently strong magnetic field and large enough baryon chemical potential: inhomogeneous condensate of neutral pions



2.0 Z/L

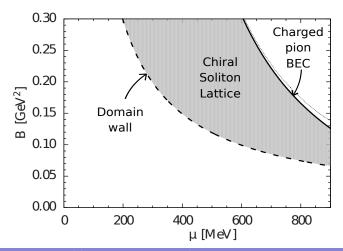
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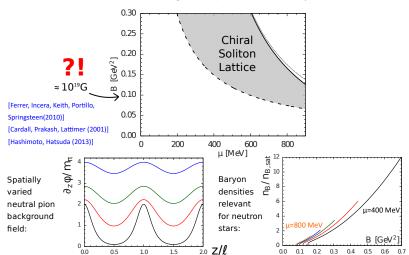
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#### Shown using chiral perturbation theory!



[Brauner, Yamamoto (2017)]

- In strong magnetic fields charged pions get large effective masses ⇒ only neutral pions remain relevant degrees of freedom!
- Ground state configuration of the neutral pion field? Minimize the energy functional (based on chiral perturbation theory):

$$\mathcal{H}_{ ext{eff}} = rac{f_\pi^2}{2}(oldsymbol{
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$$\mathcal{H}_{\text{eff}} = \frac{f_{\pi}^2}{2} (\nabla \phi)^2 + m_{\pi}^2 f_{\pi}^2 (1 - \cos \phi) - \frac{\mu}{4\pi^2} \mathbf{B} \cdot \nabla \phi \quad \leftarrow \text{Chiral anomaly!}$$

 $\Rightarrow \partial_z^2 \phi = m_\pi^2 \sin \phi$  (magnetic field in z direction)

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Ground state for  $\mu B \geq 16\pi m_{\pi} f_{\pi}^2$ :

$$\cos\frac{\phi(z)}{2}=\sin(\frac{m_\pi z}{k},k),$$

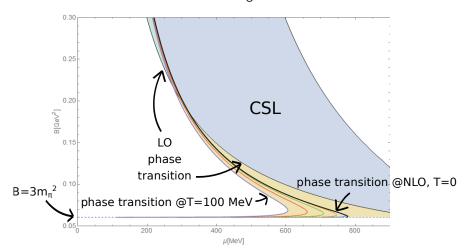
with elliptic modulus k fixed as

$$\frac{E(k)}{k} = \frac{\mu B}{16\pi m_{\pi} f_{\pi}^2}$$

#### Chiral soliton lattice phase at finite temperature

[Brauner, H.K., Yamamoto (soon)]

Chiral soliton lattice phase stabilized by NLO contributions for moderate magnetic fields!



## Chiral soliton lattice phase at finite temperature

- LO: domain wall formed at the phase transition ⇒
  - 1-loop free energy for domain wall background  $\mathcal{F}^{NLO}(B,T)$  calculated
  - phase transition occurs at such points in parameter space where  $\mathcal{F}^{LO+NLO}(\mu, B, T) = 0$

LO: 
$$\frac{\mathscr{E}_{\text{wall}}}{S} = 8m_{\pi}f_{\pi}^2 - \frac{\mu B}{2\pi}$$
. NLO, T=0

NLO, 
$$T=0$$

$$\begin{split} \frac{\mathscr{F}_{\text{wall,fin}}^{T=0}}{S} &= \frac{B^{3/2}}{\sqrt{2\pi}} \bigg\{ \zeta(-\frac{1}{2}, \frac{1}{2} - \frac{3m_z^2}{2B}) + \zeta(-\frac{1}{2}, \frac{1}{2}) \\ &- \int \frac{\mathrm{d}P}{2\pi} \frac{2}{1 + P^2} \bigg[ \zeta(-\frac{1}{2}, \frac{1}{2} + \frac{m_\pi^2(1 + P^2)}{2B}) + \frac{2}{3} \bigg( \frac{m_\pi^2}{2B} \bigg)^{3/2} (1 + P^2)^{3/2} \bigg] \\ &- \int \frac{\mathrm{d}P}{2\pi} \frac{4}{4 + P^2} \bigg[ \zeta(-\frac{1}{2}, \frac{1}{2} + \frac{m_\pi^2(1 + P^2)}{2B}) + \bigg( \frac{m_\pi^2}{2B} \bigg)^{3/2} \bigg( \frac{2}{3} (4 + P^2)^{3/2} - 3(4 + P^2)^{1/2} \bigg) \bigg] \bigg\}. \end{split}$$

$$\frac{\mathscr{F}_{\text{wall,ren}}^{T=0}}{S} = \frac{m_\pi^3}{24\pi^2} \left( -17\log\frac{4\pi\Lambda_{\text{RG}}^2}{m_\pi^2} + 40\log2 - \frac{121}{3} + 17\gamma_{\text{E}} \right) + m_\pi^3 \left( -\frac{64}{3}\bar{\ell}_1 - \frac{8}{3}\bar{\ell}_2 + \frac{8}{3}\bar{\ell}_3 \right).$$

#### NLO. finite T:

$$\begin{split} \frac{\mathscr{F}_{\text{wall}}^{T,(\pi^0)}}{S} &= -\frac{\zeta(3)T^3}{2\pi} - \frac{m_\pi^2 T}{\pi^2} \int_0^\infty Q \arctan Q \\ &\qquad \times \log \left(1 - e^{-x\sqrt{1+Q^2}}\right) \mathrm{d}Q, \end{split}$$

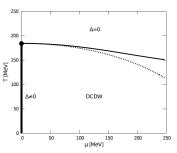
$$\frac{\mathscr{F}_{\text{wall}}^{T,(\tau^0)}}{S} = -\frac{\zeta(3)T^3}{2\pi} - \frac{m_\pi^2 T}{\pi^2} \int_0^\infty Q \arctan Q \\ \times \log \left(1 - e^{-x\sqrt{1+Q^2}}\right) \mathrm{d}Q, \\ \left[ \frac{\mathscr{F}_{\text{wall}}^{T,(\tau^\pm)}}{S} = \frac{BT}{\pi} \sum_{m=0}^\infty \left\{ \log \left[1 - e^{-\beta\sqrt{(2m+1)B} - 3m_\pi^2}\right] + \log \left[1 - e^{-\beta\sqrt{(2m+1)B}}\right] - \frac{1}{\pi} \int_0^\infty \mathrm{d}P \left(\frac{2}{1+P^2} + \frac{4}{4+P^2}\right) \log \left[1 - e^{-\beta\epsilon(m,P^2)}\right] \right\}.$$

## NB: Connection to Chiral Density Wave?

[Tatsumi, Nishiyama, Karasawa (2015)] [Ferrer, Incera (2020)]

- QCD at finite  $B, \mu, T$  studied within NJL-like models, CDW found to be preferred in certain parameter range
- Important effect of chiral anomaly observed
- Only chiral limit considered
- CSL at chiral limit is equivalent to CDW! [Brauner, Yamamoto(2017)]



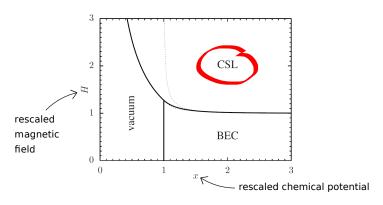


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#### CSL phase in lattice simulations?

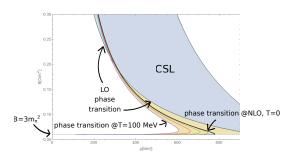
[Tomáš Brauner, Georgios Filios, H.K.; Phys. Rev. Lett. 123(2019), JHEP 1912 (2019) 029]

- $\bullet$  In certain QCD-like theories (e.g., two-color QCD) the sign problem is absent  $\Rightarrow$  lattice simulations possible
- CSL-like phase present for sufficiently large magnetic fields! (Shown using chiral perturbation theory.)
- Conjecture of [Splittorff, Son, Stephanov (2001)] that the inhomogeneous phases exist only in theories with the sign problem disproved!



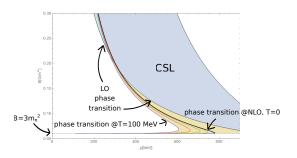
#### Conclusions

- "New" inhomogeneous phase of QCD matter appears for strong magnetic field and moderate baryon chemical potential
- Chiral soliton lattice is stable under thermal fluctuations!
- Chiral soliton lattice phase in QCD-like theories may be in principle seen in lattice simulations!



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## Thank you for your attention!

# Backup Slides

#### Relevance of CSL for heavy ion collisions?

