Warm dense QCD matter in strong magnetic fields

Helena Kolešová

University of Stavanger

Joint work with Tomáš Brauner, Naoki Yamamoto and Georgios Filios

Outline

 \bullet Motivation: \langle QCD phase diagram

Inhomogeneous phases of QCD matter

- **2** Chiral soliton lattice phase at LO
- **3** Chiral soliton lattice phase at finite temperature
- ⁴ Chiral soliton lattice phase on lattice?

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Motivation: Inhomogeneous phases of matter

[Alford,Bowers,Rajagopal(2001)], review: [Buballa,Carignano(2015)]

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- 1D modulation in spherically symmetric 3D system unstable at any non-zero temperature! [Peierls (1934), Landau] (but quasi-long-range order may be sustained [Hidaka et al. (2015)])

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- **If the spherical symmetry broken by magnetic field, 1D modulations** stable! [Tatsumi,Nishiyama,Karasawa(2015)] [Ferrer,Incera(2020)]

[Son,Stephanov(2008)][Brauner,Yamamoto(2017)]

Ground state of QCD matter for sufficiently strong magnetic field and large enough baryon chemical potential: inhomogeneous condensate of neutral pions

Shown using chiral perturbation theory!

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[Brauner,Yamamoto(2017)]

- In strong magnetic fields charged pions get large effective masses ⇒ only neutral pions remain relevant degrees of freedom!
- Ground state configuration of the neutral pion field? Minimize the energy functional (based on chiral perturbation theory):

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\mathcal{H}_{\textrm{eff}}=\frac{f_\pi^2}{2}(\boldsymbol{\nabla}\phi)^2+m_\pi^2f_\pi^2(1-\cos\phi)-\frac{\mu}{4\pi^2}\boldsymbol{B}\cdot\boldsymbol{\nabla}\phi \qquad \big(\phi=\frac{\pi^0}{f_\pi}\big)
$$

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Ground state for $\mu B \geq 16\pi m_\pi f_\pi^2$:

$$
\cos\frac{\phi(z)}{2}=\operatorname{sn}(\frac{m_{\pi}z}{k},k),
$$

with elliptic modulus k fixed as

$$
\frac{E(k)}{k} = \frac{\mu B}{16\pi m_\pi f_\pi^2}
$$

Chiral soliton lattice phase at finite temperature

[Brauner,H.K.,Yamamoto (soon)]

Chiral soliton lattice phase stabilized by NLO contributions for moderate magnetic fields!

Chiral soliton lattice phase at finite temperature

• LO: domain wall formed at the phase transition \Rightarrow

1-loop free energy for domain wall background $\mathcal{F}^{NLO}(B, \mathcal{T})$ calculated 2 phase transition occurs at such points in parameter space where $\mathcal{F}^{LO+NLO}(\mu, B, \mathcal{T})=0$

$$
\begin{split}\n\text{LO:} \quad & \left[\frac{\mathcal{E}_{\text{wall}}}{S} = 8m_{\pi}f_{\pi}^{2} - \frac{\mu B}{2\pi} \right] \qquad \text{NLO, } \text{T=0} \\
& \frac{\mathcal{F}_{\text{wall,fin}}^{T=0}}{S} = \frac{B^{3/2}}{\sqrt{2}\pi} \left\{ \zeta \left(-\frac{1}{2}, \frac{1}{2} - \frac{3m_{\pi}^{2}}{2B} \right) + \zeta \left(-\frac{1}{2}, \frac{1}{2} \right) \right. \\
& \qquad \left. - \int \frac{\mathrm{d}P}{2\pi} \frac{2}{1+P^{2}} \left[\zeta \left(-\frac{1}{2}, \frac{1}{2} + \frac{m_{\pi}^{2}(1+P^{2})}{2B} \right) + \frac{2}{3} \left(\frac{m_{\pi}^{2}}{2B} \right)^{3/2} (1+P^{2})^{3/2} \right] \right. \\
& \qquad \left. - \int \frac{\mathrm{d}P}{2\pi} \frac{4}{4+P^{2}} \left[\zeta \left(-\frac{1}{2}, \frac{1}{2} + \frac{m_{\pi}^{2}(1+P^{2})}{2B} \right) + \left(\frac{m_{\pi}^{2}}{2B} \right)^{3/2} \left(\frac{2}{3} (4+P^{2})^{3/2} - 3 (4+P^{2})^{1/2} \right) \right] \right\}.\n\end{split}
$$

NLO, finite T:

$$
\frac{\mathscr{F}_{\text{wall}}^{T, \text{(m^a)}}}{S} = -\frac{\zeta(3) T^3}{2 \pi} - \frac{m_\pi^2 T}{\pi^2} \int_0^\infty Q \arctan Q \quad \left[\frac{\mathscr{F}_{\text{wall}}^{T, \text{(r^a)}}}{S} = \frac{B T}{\pi} \sum_{m=0}^\infty \left\{ \log \left[1 - e^{-\beta \sqrt{(2m+1)B - 3m_\pi^2}} \right] + \log \left[1 - e^{-\beta \sqrt{(2m+1)B}} \right] \right. \\ \left. \qquad \qquad \times \log \left(1 - e^{-x \sqrt{1 + Q^2}} \right) \mathrm{d} Q, \right\} = \frac{B T}{\pi} \sum_{m=0}^\infty \left\{ \log \left[1 - e^{-\beta \sqrt{(2m+1)B - 3m_\pi^2}} \right] + \log \left[1 - e^{-\beta \sqrt{(2m+1)B}} \right] \right\}.
$$

NB: Connection to Chiral Density Wave?

[Tatsumi,Nishiyama,Karasawa(2015)] [Ferrer,Incera(2020)]

- QCD at finite B, μ, τ studied within NJL-like models, CDW found to be preferred in certain parameter range
- **•** Important effect of chiral anomaly observed
- Only chiral limit considered
- CSL at chiral limit is equivalent to CDW! [Brauner, Yamamoto(2017)]

CSL phase in lattice simulations?

[Tomáš Brauner, Georgios Filios, H.K.; Phys. Rev. Lett. 123(2019), JHEP 1912 (2019) 029]

- \bullet In certain QCD-like theories (e.g., two-color QCD) the sign problem is absent \Rightarrow lattice simulations possible
- CSL-like phase present for sufficiently large magnetic fields! (Shown using chiral perturbation theory.)
- Conjecture of [Splittorff, Son, Stephanov (2001)] that the inhomogeneous phases exist only in theories with the sign problem disproved!

Conclusions

- "New" inhomogeneous phase of QCD matter appears for strong magnetic field and moderate baryon chemical potential
- Chiral soliton lattice is stable under thermal fluctuations!
- Chiral soliton lattice phase in QCD-like theories may be in principle seen in lattice simulations!

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Thank you for your attention!

Backup Slides

Relevance of CSL for heavy ion collisions?

