

The origin of anisotropy from pp to AA collisions

by

GIULIANO GIACALONE

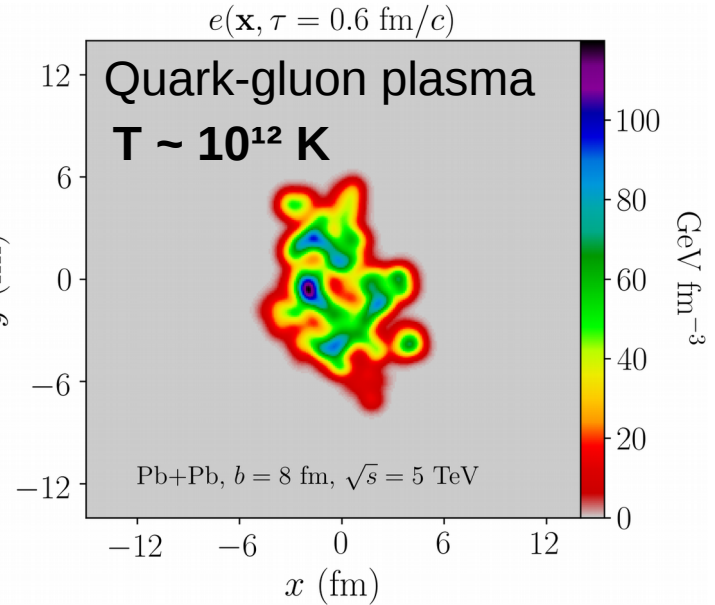
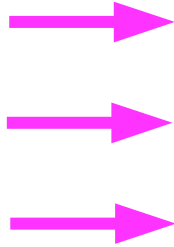
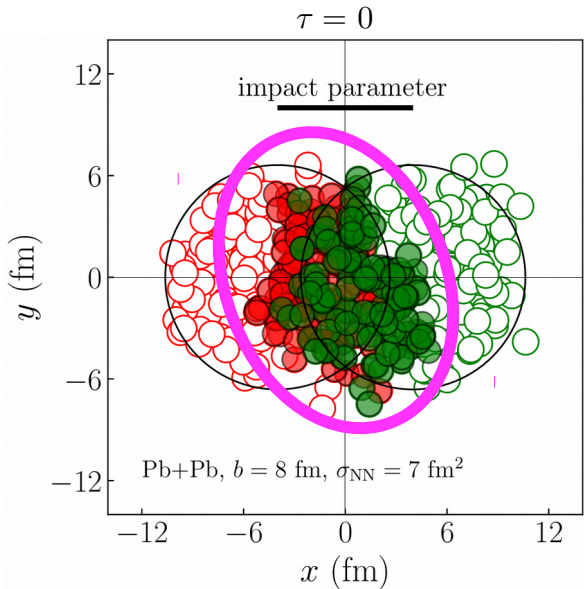
01 / 07 / 2021



UNIVERSITÄT
HEIDELBERG
ZUKUNFT
SEIT 1386



HEAVY-ION COLLISIONS = THE EARLY UNIVERSE IN THE LAB



\implies **Effective description: relativistic fluid.** [Romatschke & Romatschke, [1712.05815](#)]

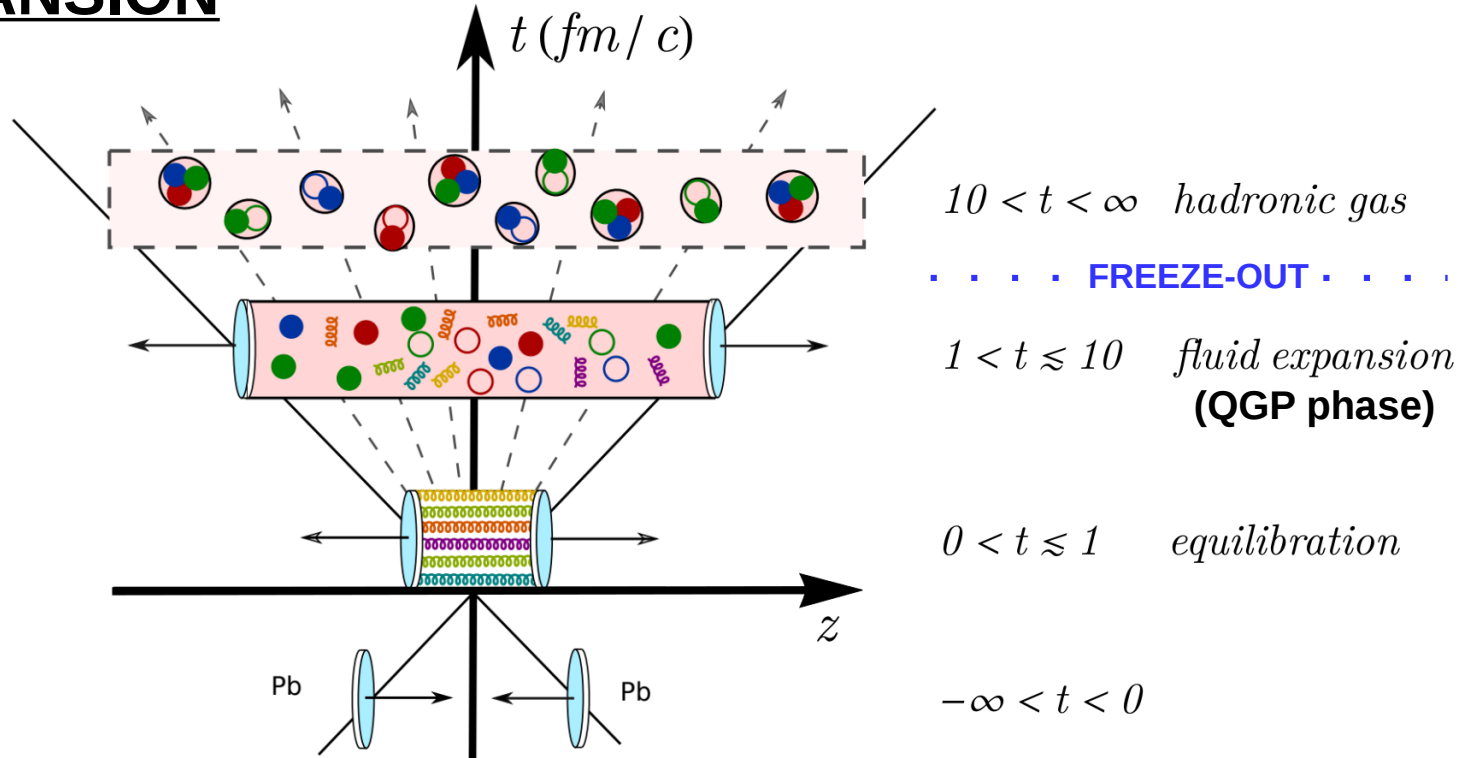
$$T^{\mu\nu} = (\epsilon + P)u^\mu u^\nu - P g^{\mu\nu} + \text{small viscous corrections } (\eta/s, \zeta/s, \dots) + \partial_\mu T^{\mu\nu} = 0$$

Equation of state from lattice QCD ($T > 156$ MeV). Large number of **DOF (~ 40): QGP.**

[HotQCD collaboration, [1407.6387](#)]

THE EXPANSION

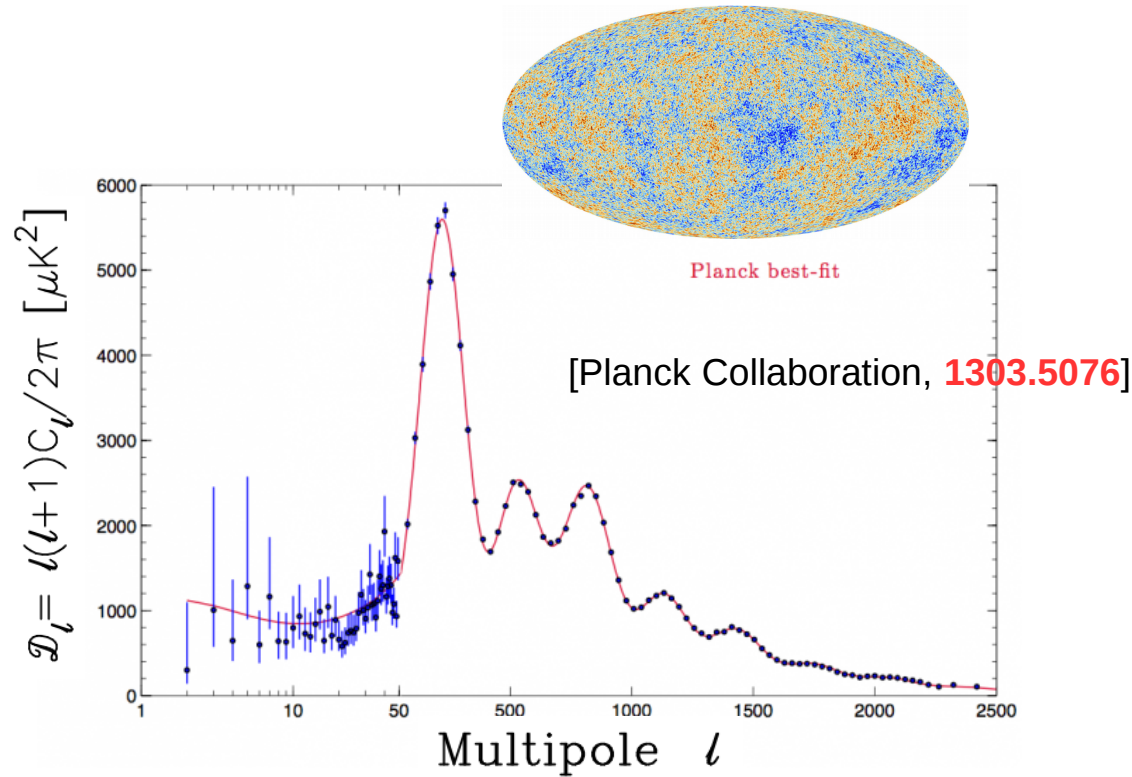
(courtesy A. Mazeliauskas)



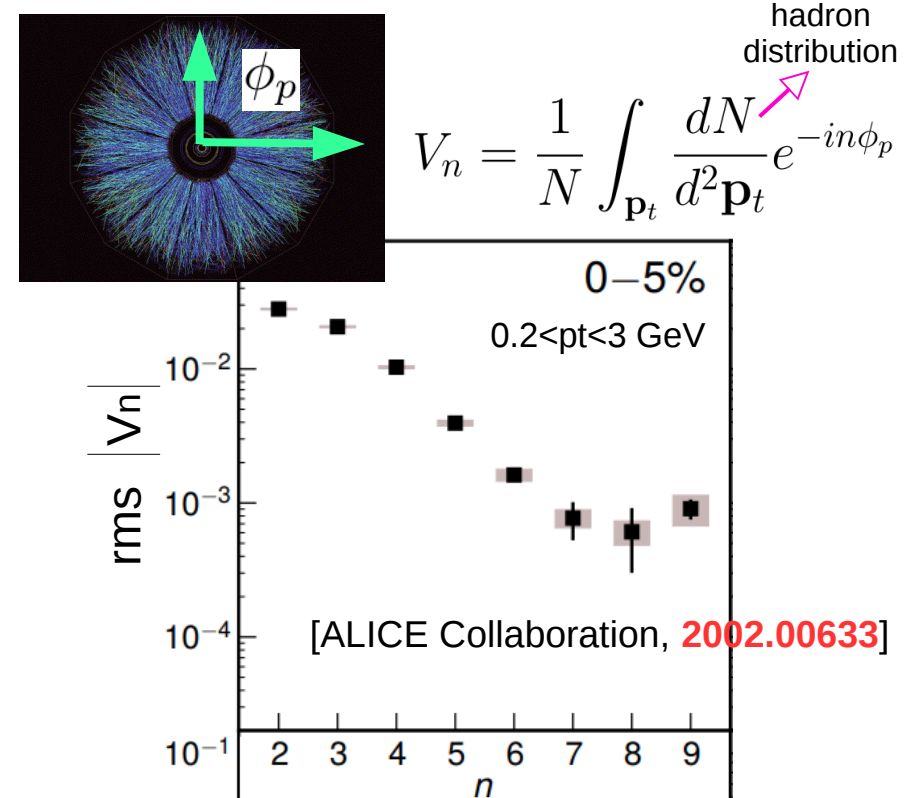
$t = \infty$ → $\frac{dN}{d^3\mathbf{p}}$ or $\frac{dN}{dyd^2\mathbf{p}_t}$ **Hadron spectrum in momentum space.**

The final state presents anisotropy in azimuthal angle.

A great Big Bang.



A billion Little Bangs.



What is the origin of anisotropy?

Structure of nuclei across length scales → primordial anisotropy → observed anisotropy

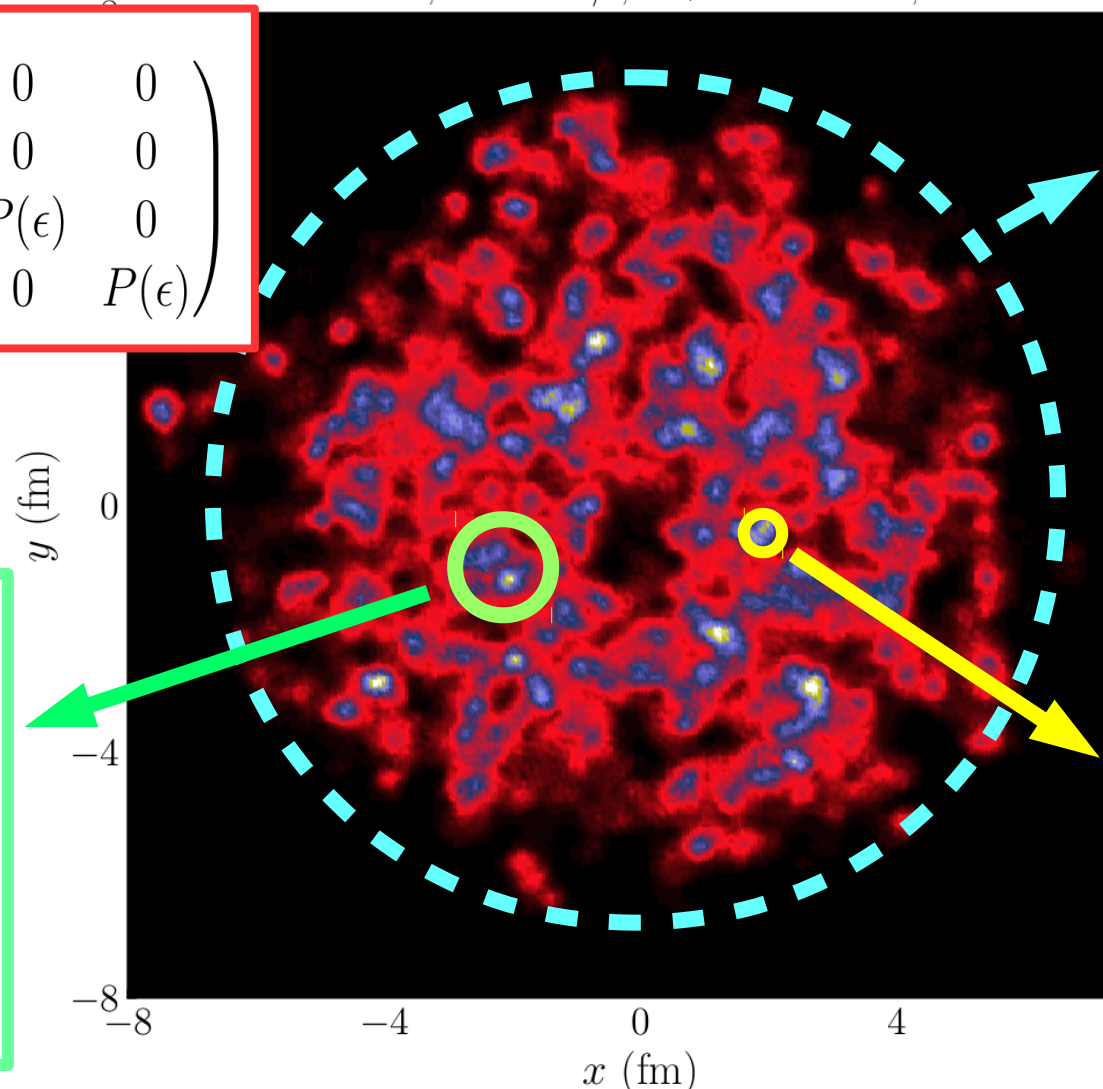
IP-Glasma, $\tau = 0.1$ fm/c, Pb+Pb @ 2.76 TeV, $b = 0$

$$T^{\mu\nu}(\tau_0) \approx \begin{pmatrix} \epsilon & 0 & 0 & 0 \\ 0 & P(\epsilon) & 0 & 0 \\ 0 & 0 & P(\epsilon) & 0 \\ 0 & 0 & 0 & P(\epsilon) \end{pmatrix}$$

Collective features

$\beta \approx 0$ $\beta \approx +0.3$

Scale ~ RA



Individual nucleons

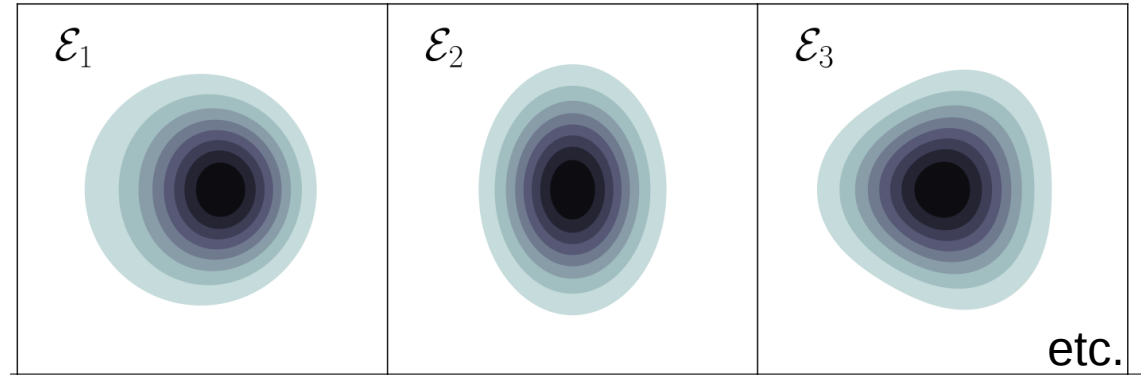
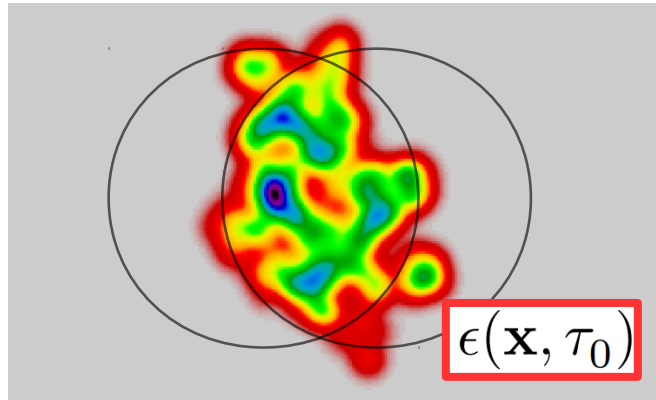
Scale ~ 1fm - RA

Subnucleonic structure ?

Scale < 1fm

Anisotropic flow from spatial anisotropy.

$$F = -\nabla P, \quad [\text{Ollitrault, 1992}]$$



[Alver, Roland, 1003.0194]

In the QGP, all multi-pole moments are nonzero:

$$\mathcal{E}_n = - \frac{\int r dr d\phi r^n e^{in\phi} \epsilon(r, \phi)}{\int r dr d\phi r^n \epsilon(r, \phi)}$$



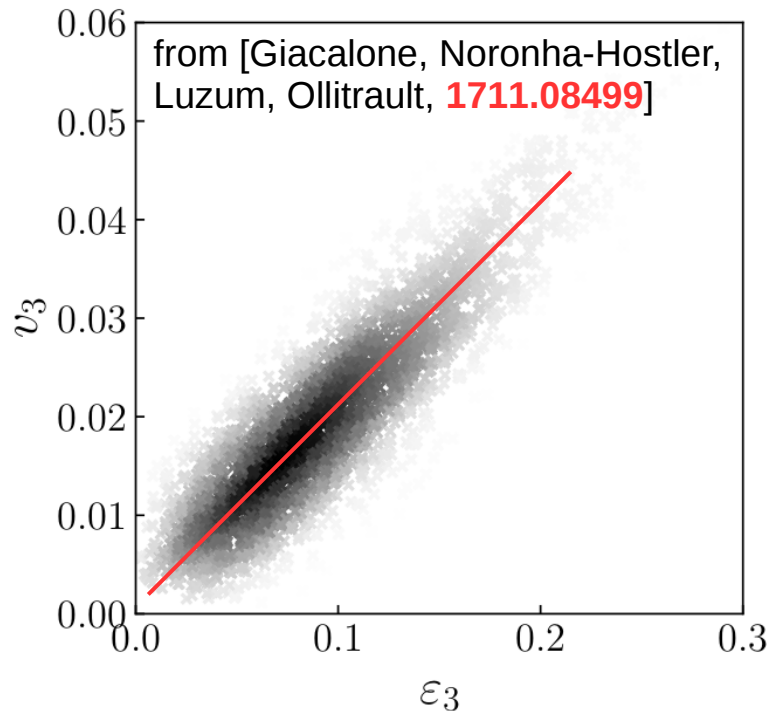
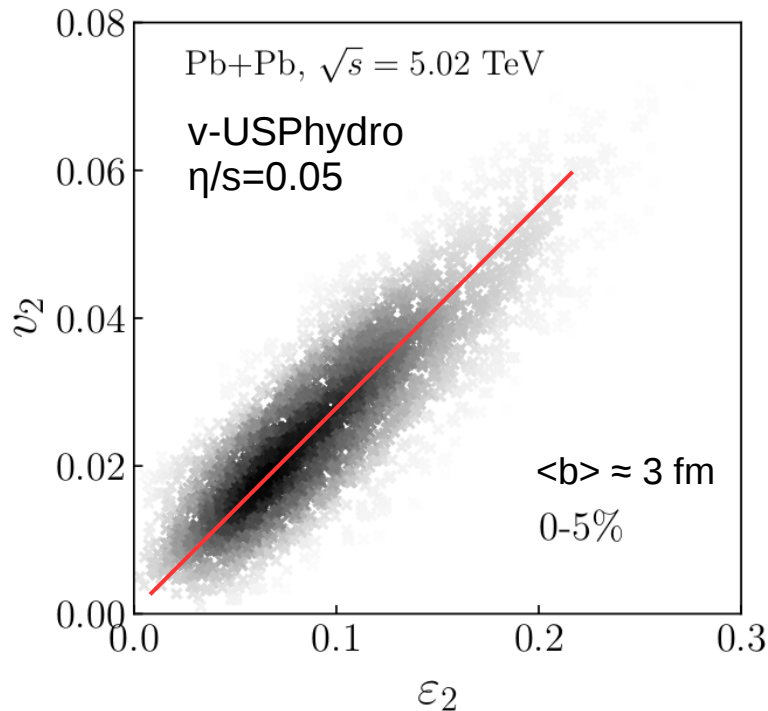
$$V_n \propto \mathcal{E}_n$$

[Teaney, Yan, 1010.1876]

Each \mathcal{E}_n in the initial state leads to V_n in the final state.

Relation is simple: $V_n \propto \mathcal{E}_n$

Verified in full hydrodynamic simulations ($\varepsilon_n = |\mathcal{E}_n|$, $v_n = |V_n|$)



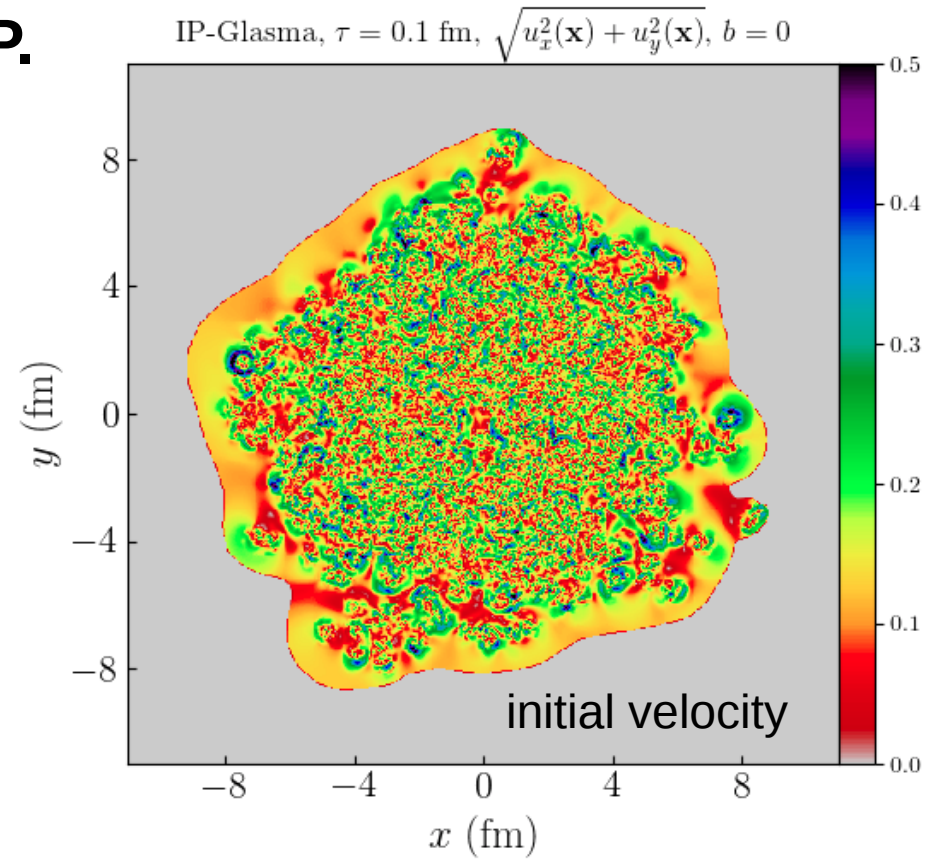
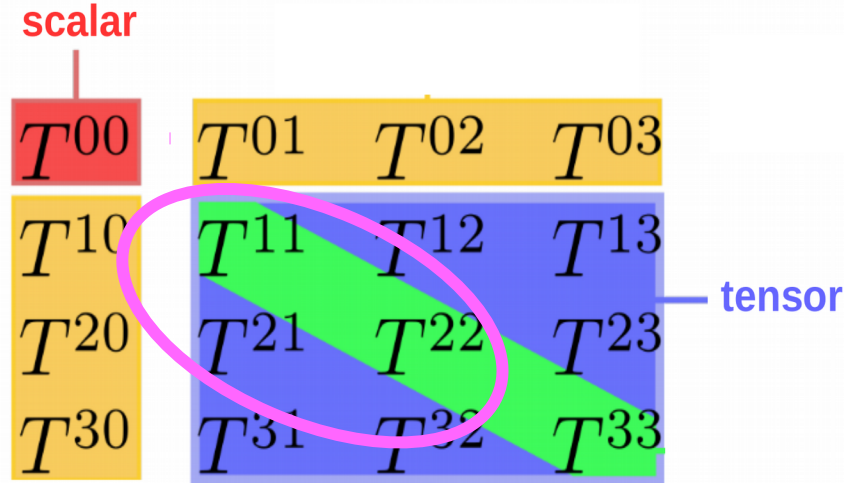
Explains experimental data in both large and small systems.

The importance of initial conditions.

[Giacalone, Noronha-Hostler, Ollitrault, [1702.01730](#)]

New frontier: Going beyond $\mathbf{F} = -\nabla P$.

Including the full stress-energy tensor.



PRIMORDIAL “MOMENTUM” ANISOTROPIES

$$\mathcal{E}_{2p} \propto \langle T^{xx} - T^{yy} + 2iT^{xy} \rangle$$

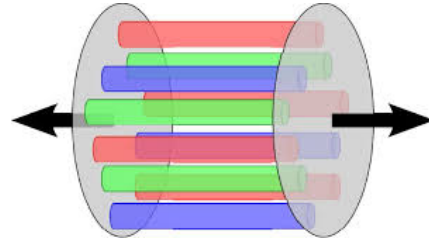
$$\mathcal{E}_{3p} \propto \left\langle r e^{i\phi} (T^{xx} - T^{yy} + 2iT^{xy}) \right\rangle$$

Evidence in data?

... evidence in theoretical calculations? Natural framework: CGC

$$T_{\text{LO}}^{00} = \frac{1}{2} \underbrace{[\mathbf{E}^2 + \mathbf{B}^2]}_{\text{class. fields}}, \quad T_{\text{LO}}^{0i} = [\mathbf{E} \times \mathbf{B}]^i$$

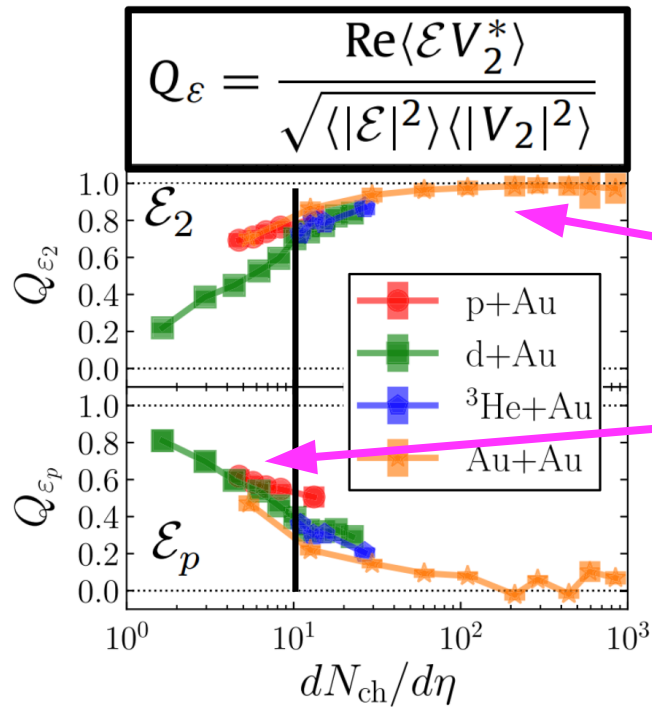
$$T_{\text{LO}}^{ij} = \frac{\delta^{ij}}{2} [\mathbf{E}^2 + \mathbf{B}^2] - [\mathbf{E}^i \mathbf{E}^j + \mathbf{B}^i \mathbf{B}^j],$$



[Gelis, [2102.07604](#)]

[upcoming talk by
A. Czajka]

Hybrid CGC+Hydro framework: IP-Glasma+MUSIC.



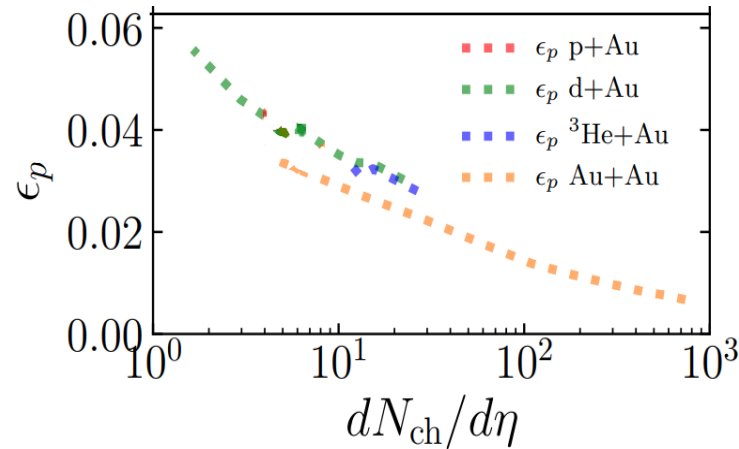
[Schenke, Shen, Tribedy, [1908.06212](#)]

\mathcal{E}_2 dominant contribution to V_2 for $dN/d\eta \geq 20$.

V_2 in a stronger correlation with E_p at low $dN/d\eta$.

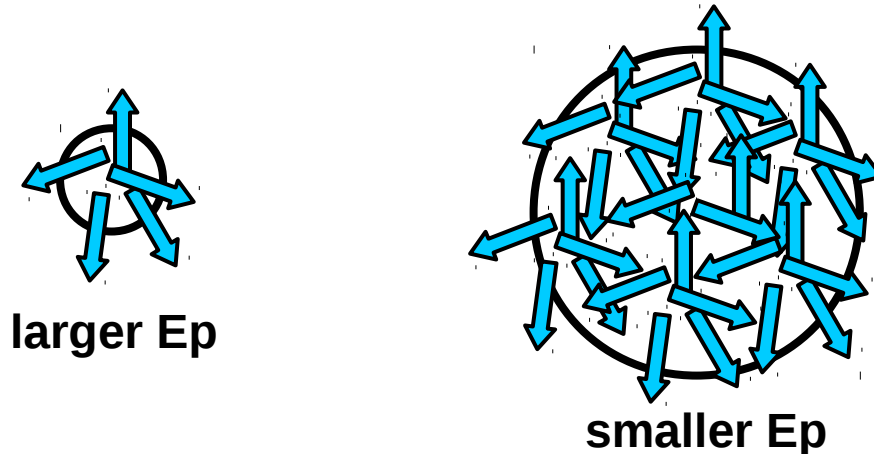
Can we reveal this transition in experiments?

Highly nontrivial... idea: exploit the dependence on system size.



[Schenke, Shen, Tribedy, [1908.06212](#)]

Net velocity field is naturally larger if the system has a smaller size.



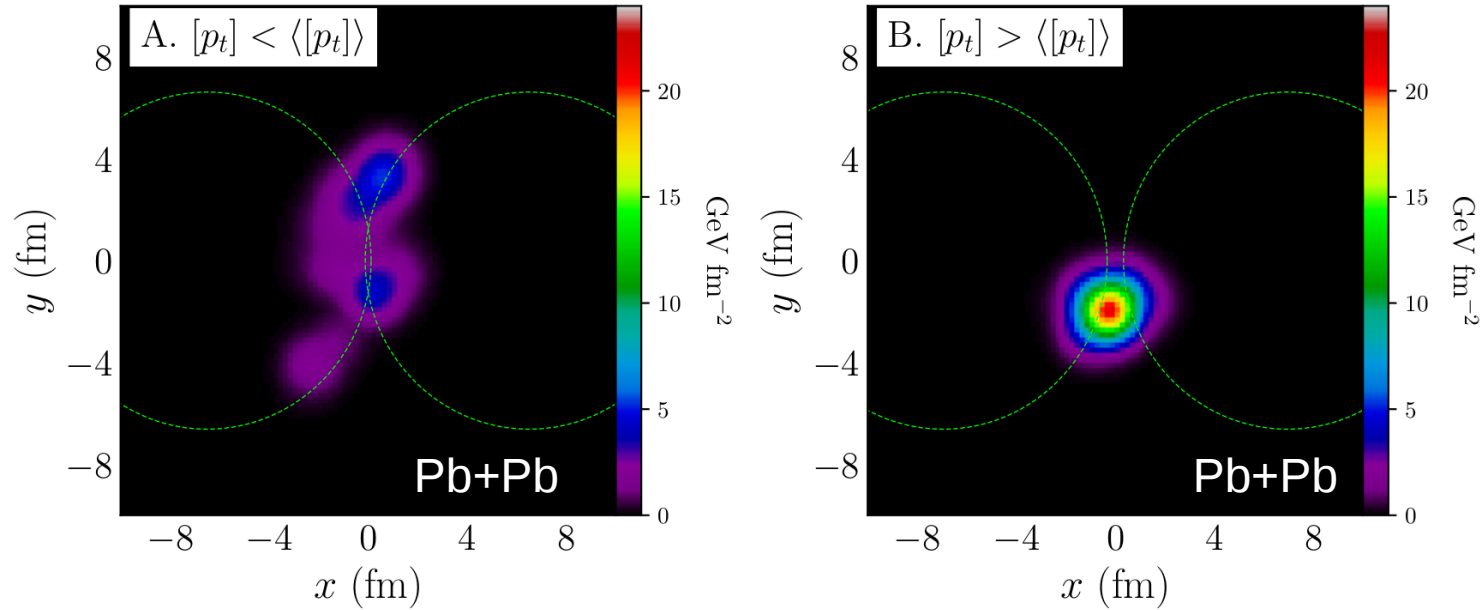
Can we change the system size in experiments?

The mean transverse momentum of the hadrons gives a handle on the system size.

[Gardim, Giacalone, Noronha-Hostler, Ollitrault [2004.01765](#)]

[Schenke, Shen, Teaney, [2004.00690](#)]

[Božek, Mehrabpour, [2002.08832](#)]



event	A	B
N_{ch}	134	134
b (fm)	13.0	13.9
$[p_t] / \langle [p_t] \rangle$	0.907	1.143
R (fm)	2.97	1.34
ε_2	0.675	0.133
ε_3	0.229	0.067

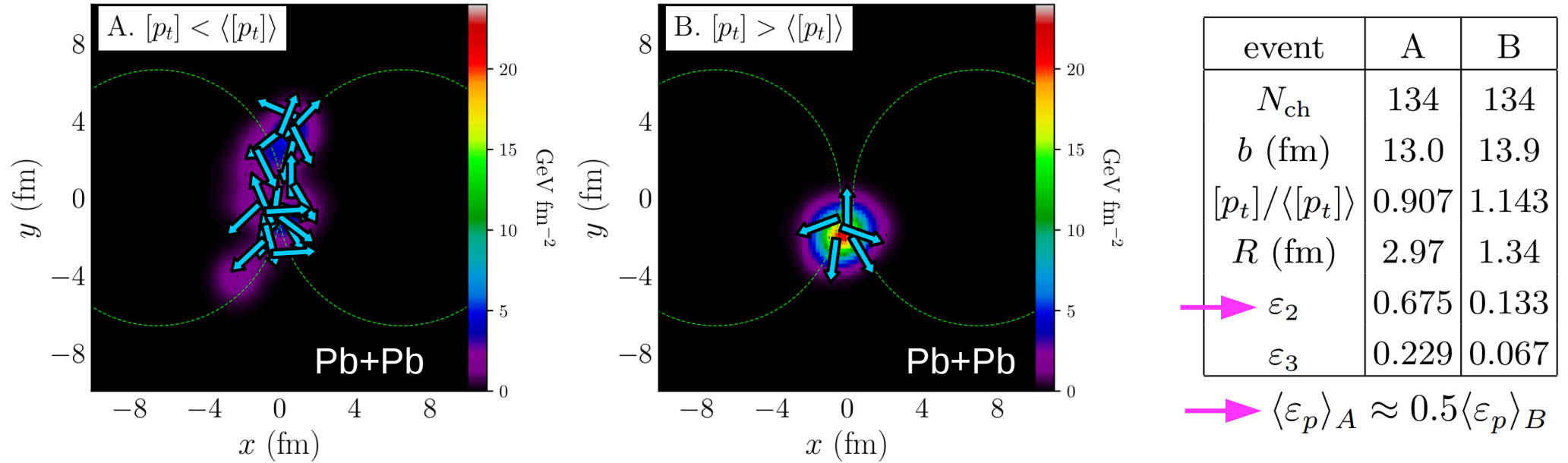
Small systems: at large $\langle p_t \rangle$ hot spots clustered around one transverse point.

The mean transverse momentum of the hadrons gives a handle on the system size.

[Gardim, Giacalone, Noronha-Hostler, Ollitrault [2004.01765](#)]

[Schenke, Shen, Teaney, [2004.00690](#)]

[Bożek, Mehrabpour, [2002.08832](#)]



Small systems: at large $\langle p_t \rangle$ hot spots clustered around one point (at fixed centrality).

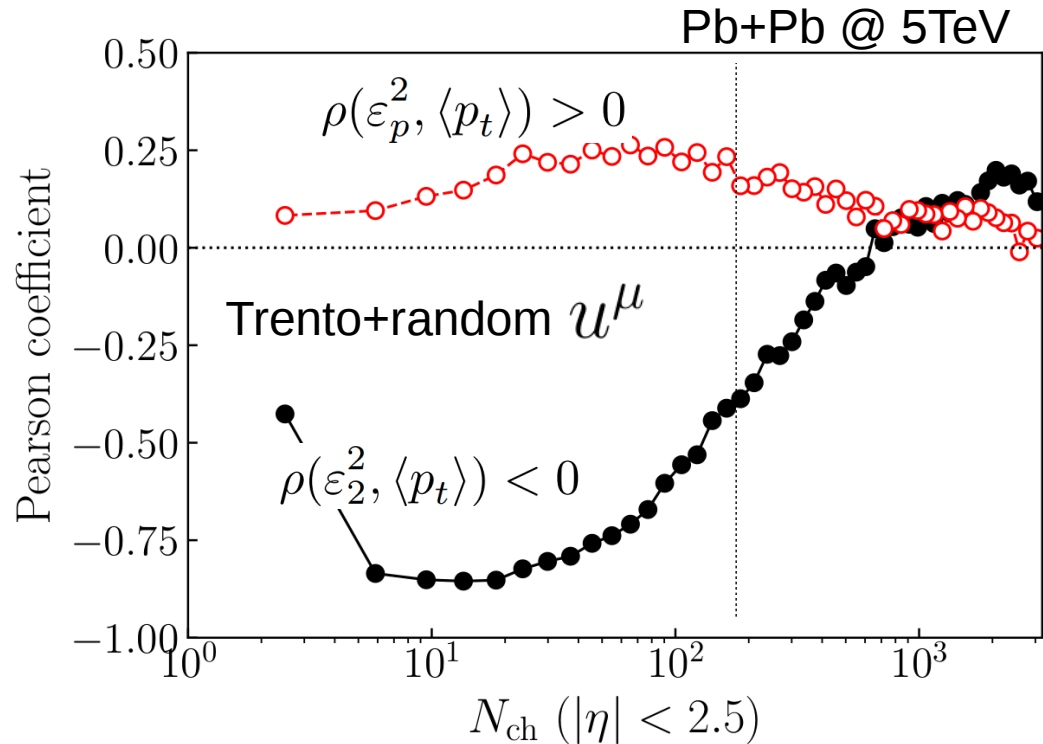
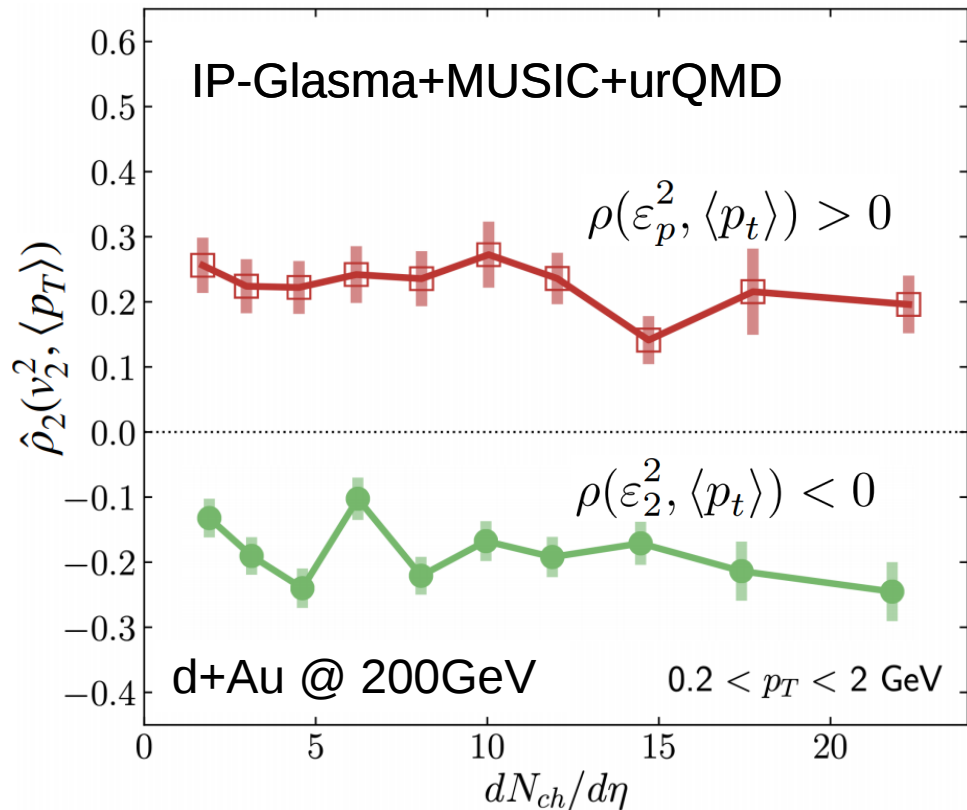
Prediction:

statistical correlation
(Pearson coefficient)

$$\rho(\varepsilon_2^2, \langle p_t \rangle) < 0$$

$$\rho(\varepsilon_p^2, \langle p_t \rangle) > 0$$

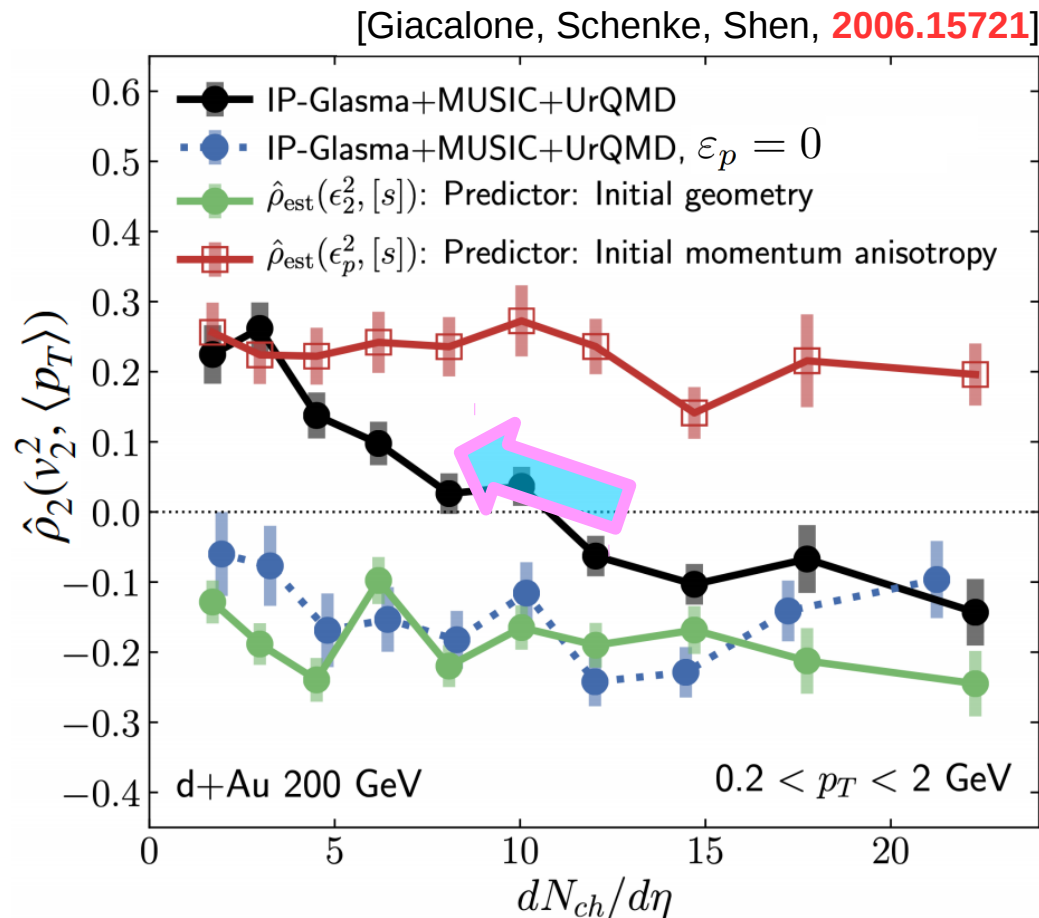
Validation in models. Generic feature of small systems (pp, pA, peripheral AA).



The contributions are qualitatively different.

[Giacalone, Schenke, Shen, [2006.15721](#)]

Turning to elliptic flow. Momentum anisotropy is dominant for $dN/d\eta < 10$.



– Sign change around $dN/d\eta=10$.
A neat prediction for experiments.

– No sign change if we set $E_p=0$.

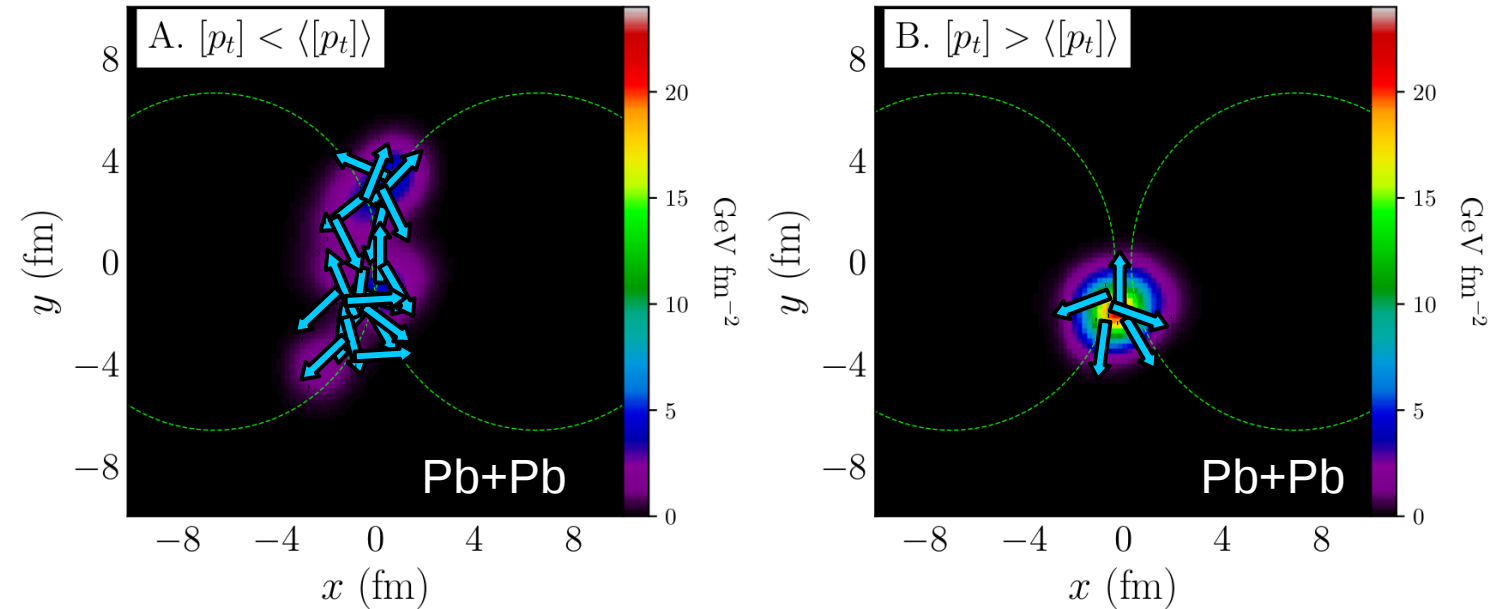
– Measurement is however difficult.
Non-flow (PYTHIA) mimics the signal.

[Zhang, Behera, Bhatta, Jia, 2102.05200]

[Lim, Nagle, 2103.01348]

SOMETHING NEW FOR THIS CONFERENCE

In small systems, ϵ_2 and ϵ_3 have the same geometric origin.



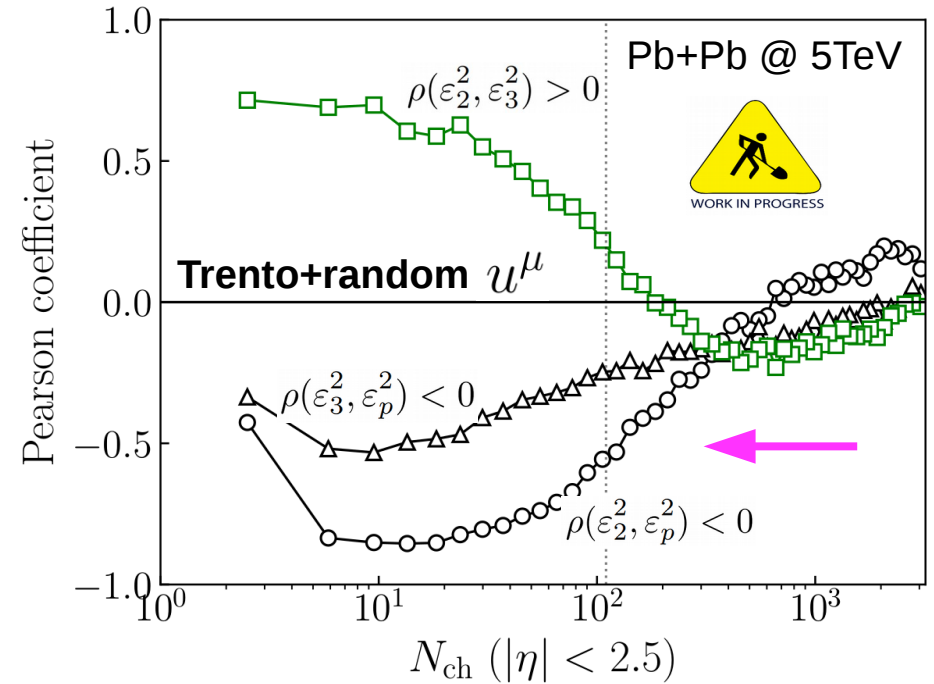
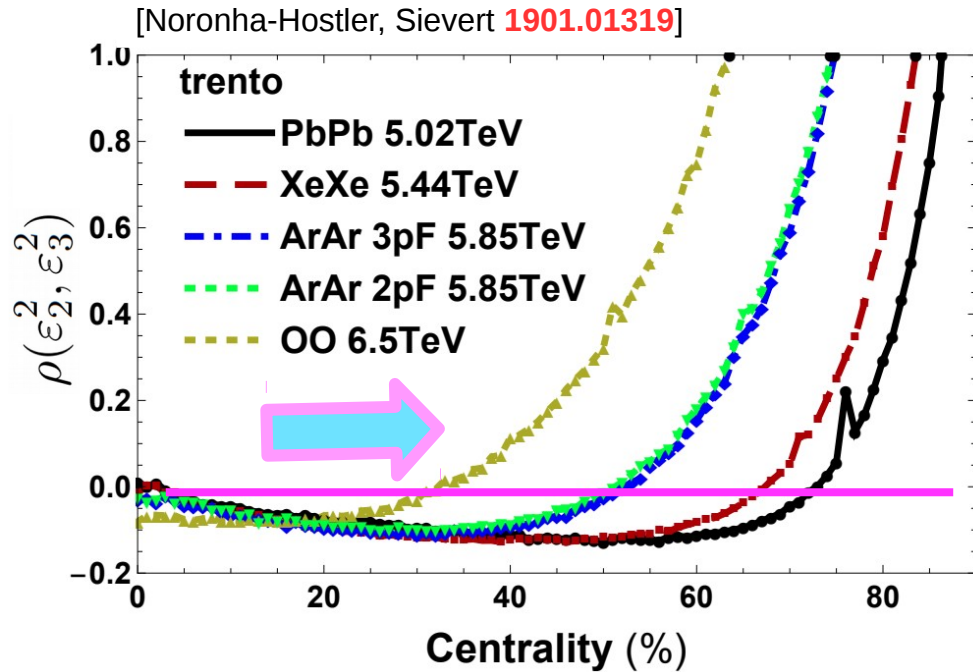
event	A	B
N_{ch}	134	134
b (fm)	13.0	13.9
$[p_t]/\langle [p_t] \rangle$	0.907	1.143
R (fm)	2.97	1.34
ϵ_2	0.675	0.133
ϵ_3	0.229	0.067
$\langle \epsilon_p \rangle_A \approx 0.5 \langle \epsilon_p \rangle_B$		

Remarkable prediction for small systems:

$$\rho(\epsilon_2^2, \epsilon_3^2) > 0$$

$$\rho(\epsilon_2^2, \epsilon_p^2) < 0, \quad \rho(\epsilon_3^2, \epsilon_p^2) < 0$$

Predictions verified in a toy calculation + geometry-only results.



Negative contributions appear due to initial momentum anisotropy.

Replace eccentricities with flow coefficients. Sign of $\rho(v_2^2, v_3^2)$?

With initial momentum anisotropy:

$$\rho(v_2^2, v_3^2) \propto \langle v_2^2 v_3^2 \rangle - \langle v_2^2 \rangle \langle v_3^2 \rangle =$$

$$\kappa_2^2 \kappa_3^2 (\langle \varepsilon_2^2 \varepsilon_3^2 \rangle - \langle \varepsilon_2^2 \rangle \langle \varepsilon_3^2 \rangle) + \kappa_2^2 \kappa_{3p}^2 (\langle \varepsilon_2^2 \varepsilon_{3p}^2 \rangle - \langle \varepsilon_2^2 \rangle \langle \varepsilon_{3p}^2 \rangle) + \kappa_{2p}^2 \kappa_3^2 (\langle \varepsilon_{2p}^2 \varepsilon_3^2 \rangle - \langle \varepsilon_{2p}^2 \rangle \langle \varepsilon_3^2 \rangle) + \kappa_{2p}^2 \kappa_{3p}^2 (\langle \varepsilon_{2p}^2 \varepsilon_{3p}^2 \rangle - \langle \varepsilon_{2p}^2 \rangle \langle \varepsilon_{3p}^2 \rangle)$$

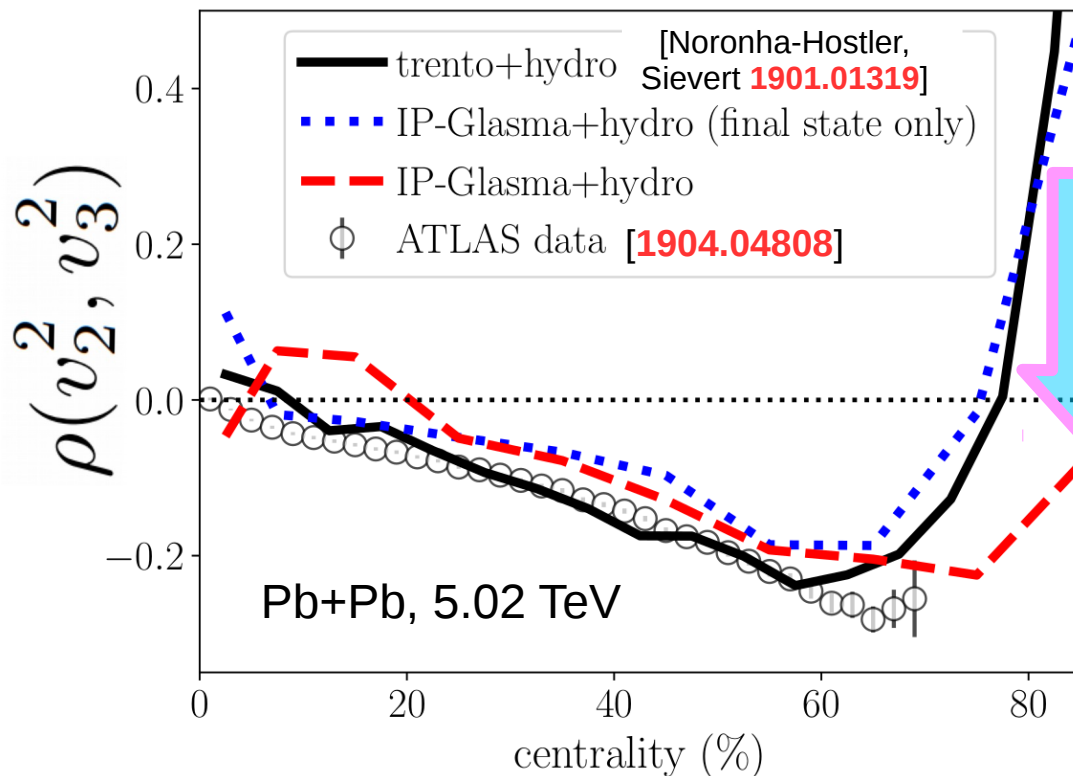
POSITIVE

NEGATIVE

NEGATIVE

?

Assume, e.g., $V_2 = \kappa_2 \mathcal{E}_2 + \kappa_p \mathcal{E}_p$



Initial E_p keeps the correlation negative.

Effect at $dN/d\eta \sim 30$, very promising.

Prediction for experiments.



- Conclusion.
- Cosmological model for pp, pA, AA collisions at high energy.
- Origin of anisotropy: mostly $F = -\nabla P$.
- Beyond response to geometry: primordial momentum anisotropy.
- Effects visible in theory results: Phenomenology is possible!
- **Natural handle: vary system size with $\langle p_t \rangle$.**
- **Natural handle: correlation between v_2 and v_3 .**
- **Hunt is ON**: more observables for unambiguous discovery.

THANK YOU!