

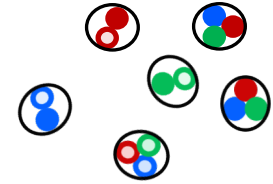
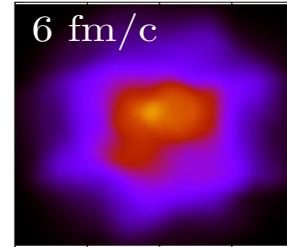
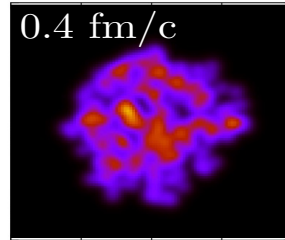
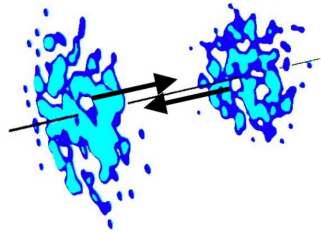
Slow modes in rapidly-expanding quark–gluon plasma

Jasmine Brewer



Based on [1910.00021] and ongoing work
with Weiyao Ke, Li Yan, and Yi Yin

Equilibration of strongly-interacting matter in heavy-ion collisions



far from equilibrium initial state

hydrodynamics

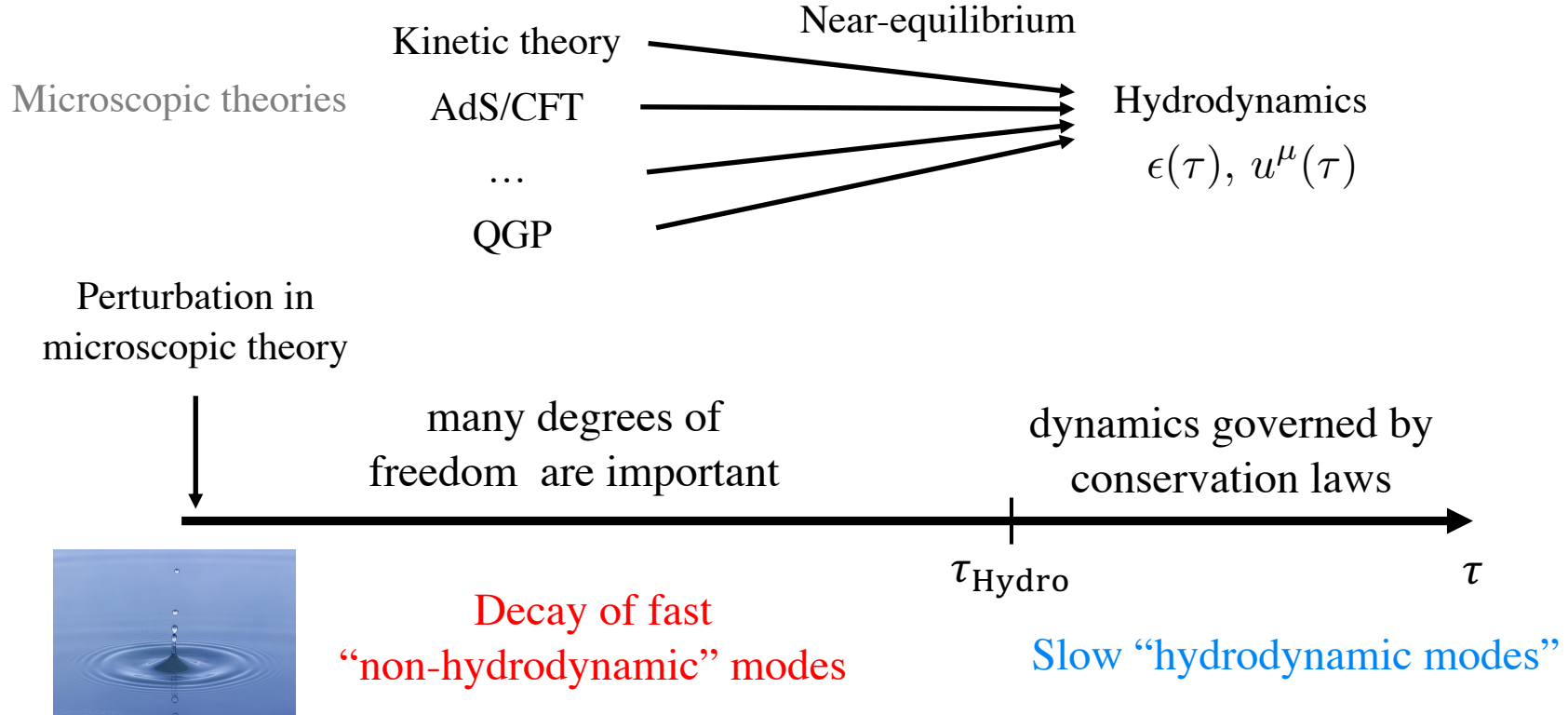
hadron gas

$\sim 1 \text{ fm}/c$

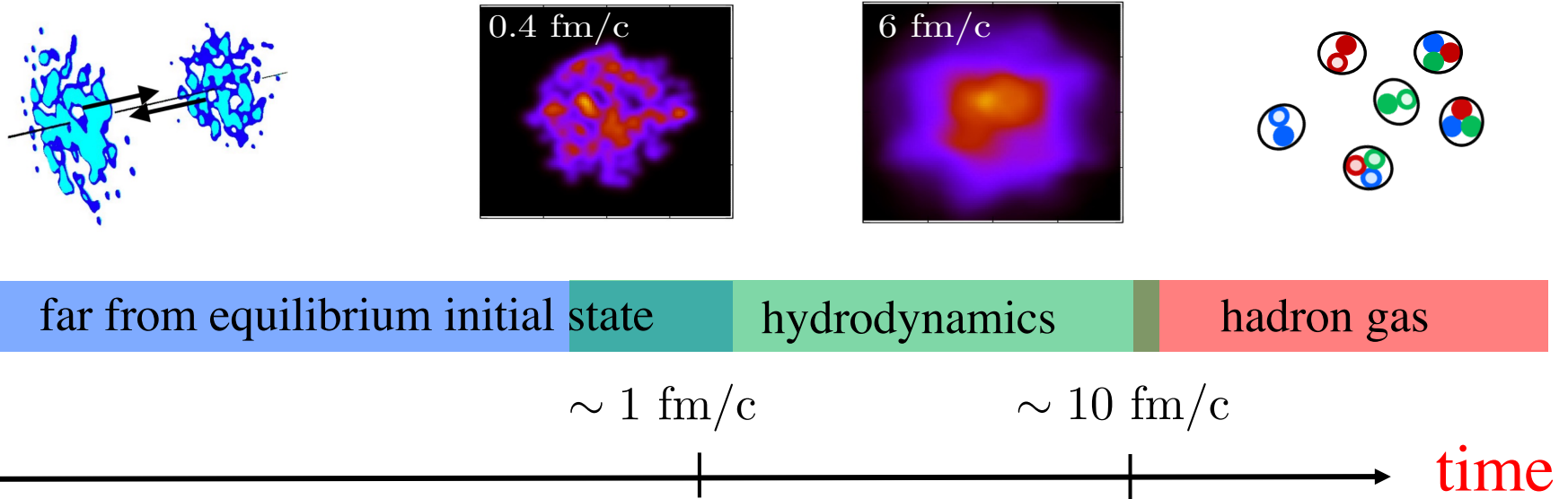
$\sim 10 \text{ fm}/c$

time

Hydrodynamization near equilibrium



Equilibration of strongly-interacting matter in heavy-ion collisions



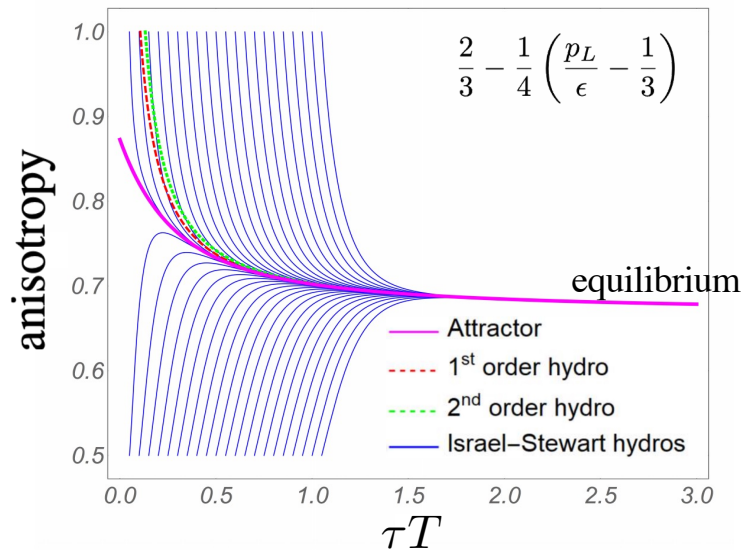
Evolution of initial state to hydrodynamics in kinetic theory

A. Mazeliauskas, next talk
X. Du 15:20

This talk: description of far-from-equilibrium systems and onset of hydrodynamics

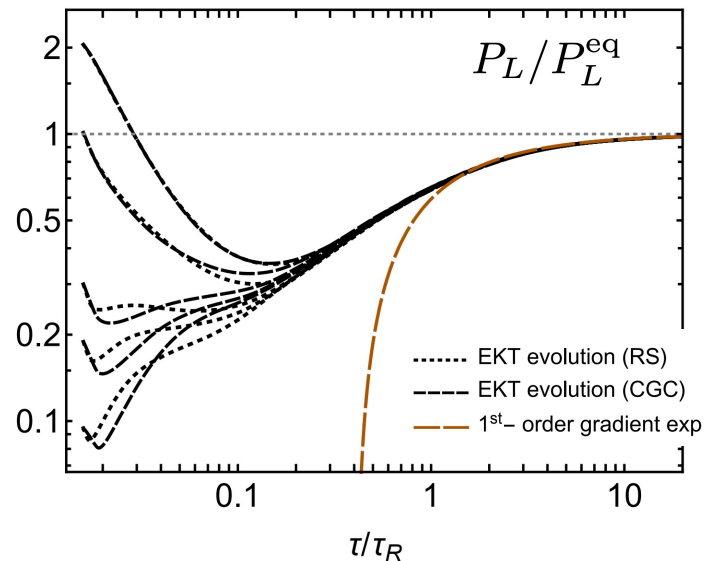
Onset of hydrodynamics far-from-equilibrium?

Far-from-equilibrium universality in hydrodynamics



Heller, Spalinski [1503.07514]

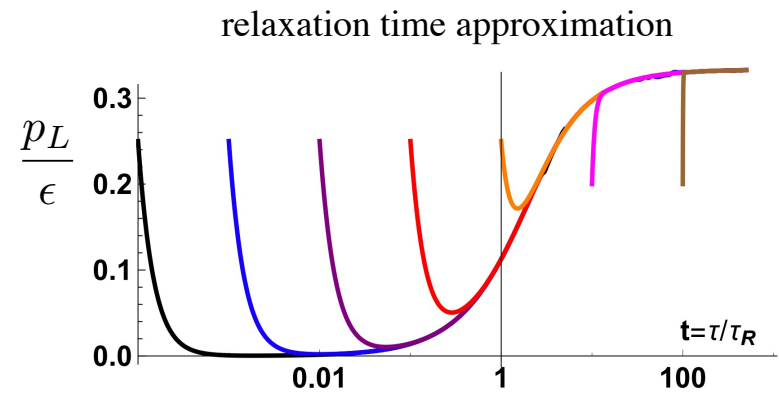
QCD effective kinetic theory



Almaalol, Kurkela, Strickland [2004.05195]

Extensive work: Heller, Spalinski, Svensson, Romatschke, Strickland, Noronha, Denicol, Kurkela, Wiedemann, Wu, Martinez,

Free-streaming and hydrodynamic attractors



(qualitatively similar results for Israel-Stewart)

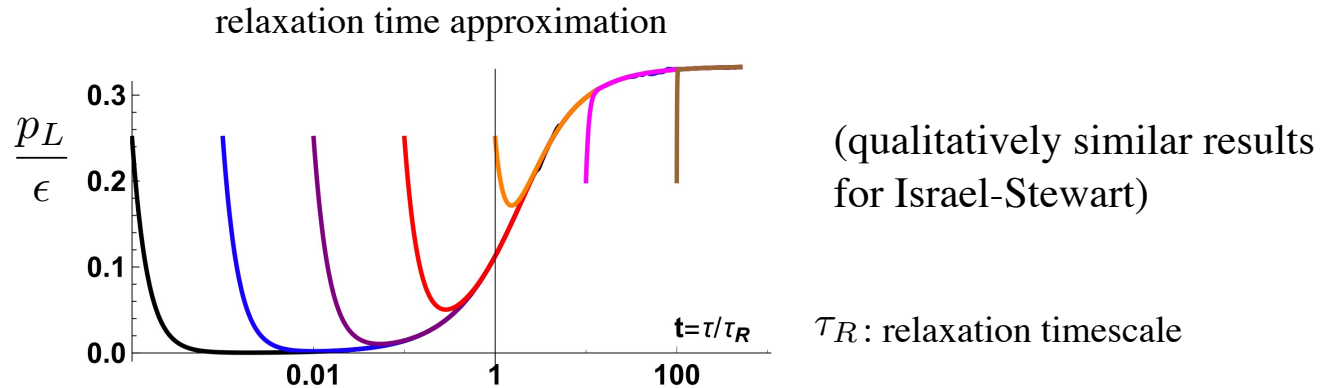
τ_R : relaxation timescale

Kurkela, van der Schee, Wiedemann, Wu [1907.08101]

Rapid expansion
without collisions

Collisions

Free-streaming and hydrodynamic attractors



Kurkela, van der Schee, Wiedemann, Wu [1907.08101]

Rapid expansion
without collisions

Collisions

Collectivity can be generated by small number of final-state interactions R. Tornkvist, 15:00

Borghini, Feld, Kerstig [1804.05729], Kurkela, Wiedemann, Wu [1905.05139], Kurkela, Mazeliauskas, Tornkvist [2104.08179]

Motivates understanding the transition from free-streaming to hydrodynamics

Hydrodynamization in kinetic theory

$$\partial_\tau f - \underbrace{\frac{p_z}{\tau} \partial_{p_z} f}_{\text{longitudinal expansion}} = - \underbrace{C[f]}_{\text{collisions}}$$

Hydrodynamization in kinetic theory

$$\int_p p \left(\underbrace{\partial_\tau f}_{\partial_\tau F} - \frac{p_z}{\tau} \underbrace{\partial_{p_z} f}_{\frac{1}{\tau}(\dots)F} = - \underbrace{C[f]}_{(\dots)F} \right)$$

(sometimes)

Here: RTA

Nonlinear kernels:

JB, Scheihing-Hitschfeld, Yin (in progress)

F : angular distribution with dimension of the stress tensor

Truncate using moment expansion

Bjorken expansion

$$F(\cos\theta; \tau) = \epsilon(\tau) + \sum_{n=1} \frac{4n+1}{2} \mathcal{L}_n(\tau) P_{2n}(\cos\theta) \iff \psi = (\epsilon, \mathcal{L}_1, \mathcal{L}_2, \dots)$$

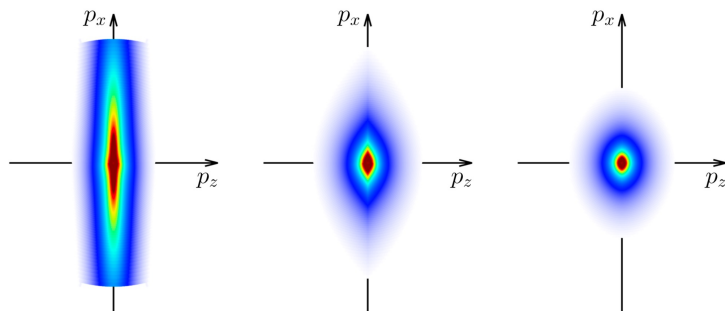
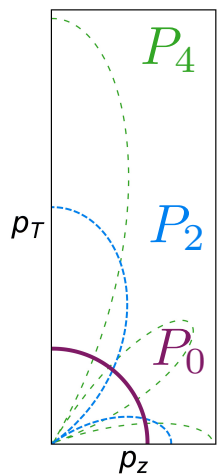
$$\text{Evolution of } F \iff \tau \partial_\tau \psi = -\mathcal{H}(\tau) \psi$$

Truncate using moment expansion

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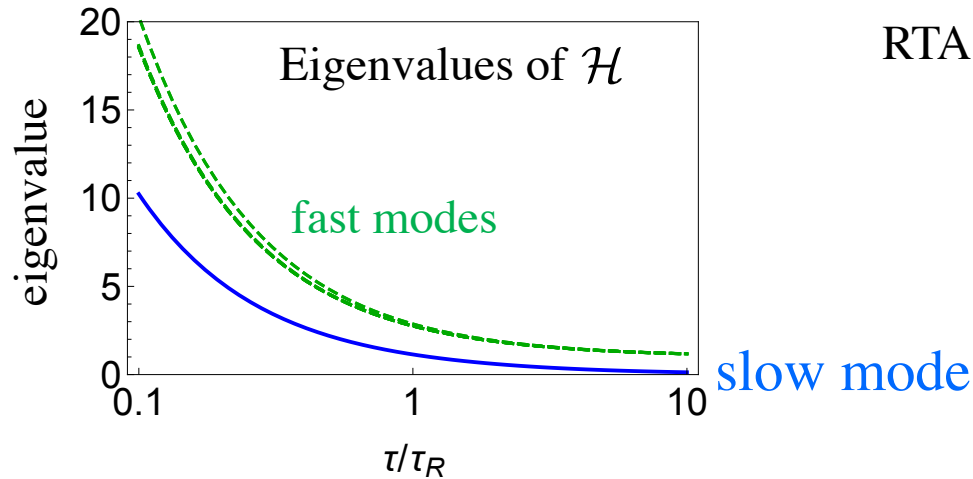


$$(\epsilon, \mathcal{L}_1, \mathcal{L}_2, \dots) \longrightarrow (\epsilon, 0, 0, \dots)$$

hydrodynamization

Fig adapted from KoMPoST [1805.00961]

Ground state: far-from-equilibrium slow mode

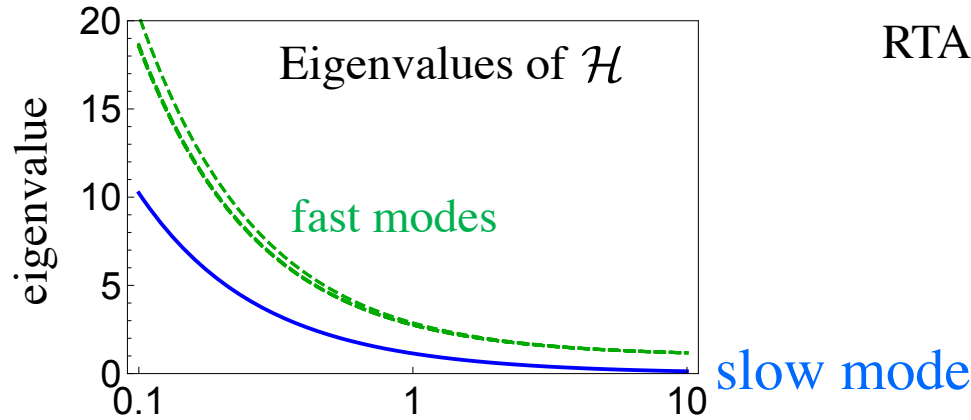


$$\text{RTA: } \mathcal{H} = \mathcal{H}_F + \frac{\tau}{\tau_R} \mathcal{H}_H$$

Early times: $\Delta E \sim \frac{1}{\tau}$

Late times: $\Delta E \sim \frac{1}{\tau_R}$

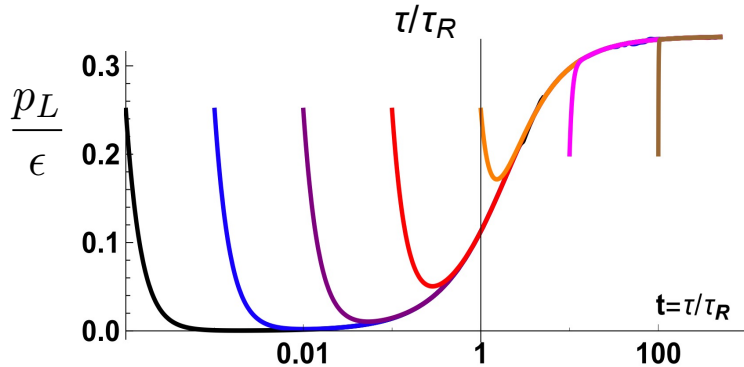
Ground state: far-from-equilibrium slow mode



$$\text{RTA: } \mathcal{H} = \mathcal{H}_F + \frac{\tau}{\tau_R} \mathcal{H}_H$$

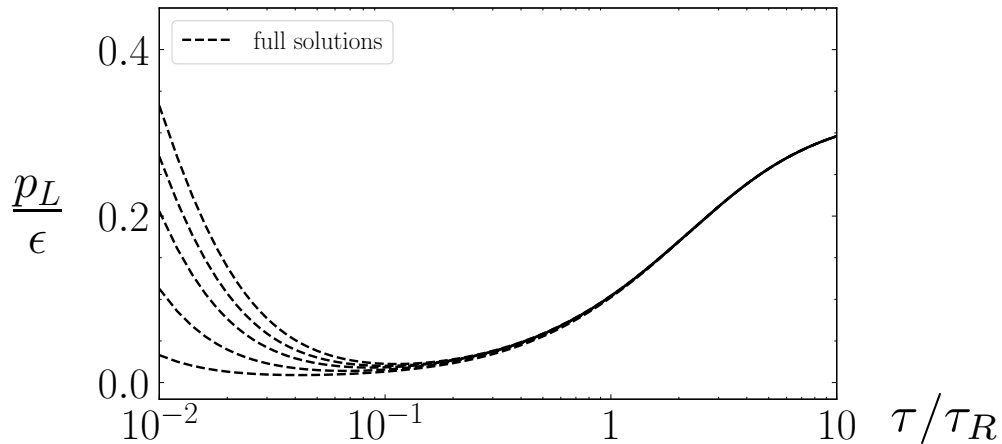
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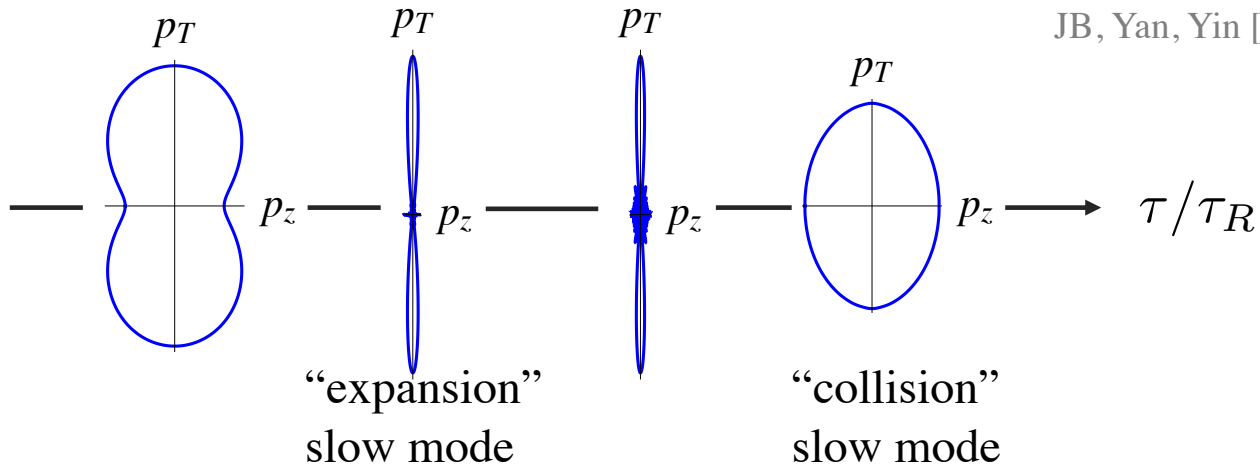


Initial conditions decay to ground state on time scale set by energy gap

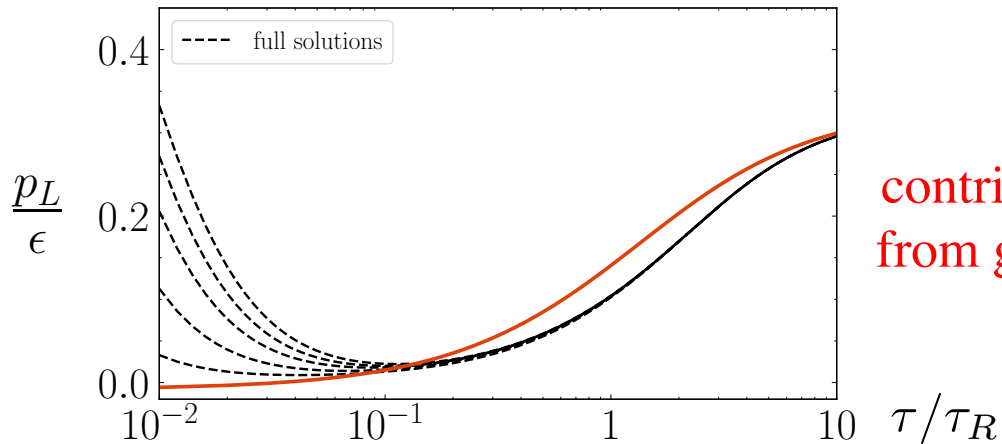
Ground state: far-from-equilibrium slow mode



JB, Yan, Yin [1910.00021]

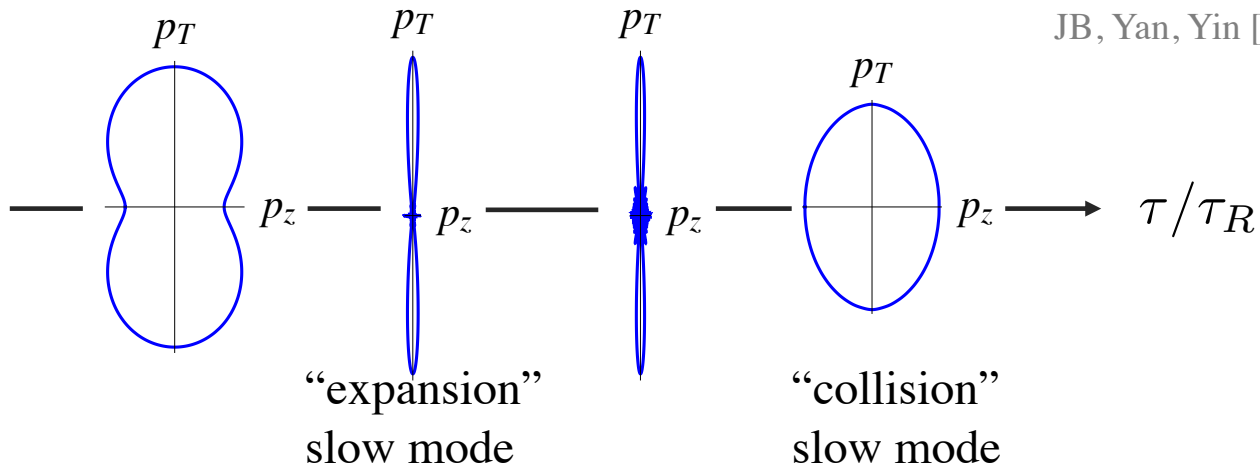


Ground state: far-from-equilibrium slow mode



contribution only
from ground state

JB, Yan, Yin [1910.00021]



A system prepared in its (instantaneous) ground state remains in its (instantaneous) ground state if transitions are suppressed

$$\text{Transition rate} \sim \frac{\partial_{\tau} \log \lambda}{\Delta E_n} \langle 0(\tau) | H | n(\tau) \rangle \quad \lambda = \tau / \tau_R$$

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“Slow quench” adiabaticity

- Hamiltonian evolution slow compared to energy gap
- Small close to hydro limit

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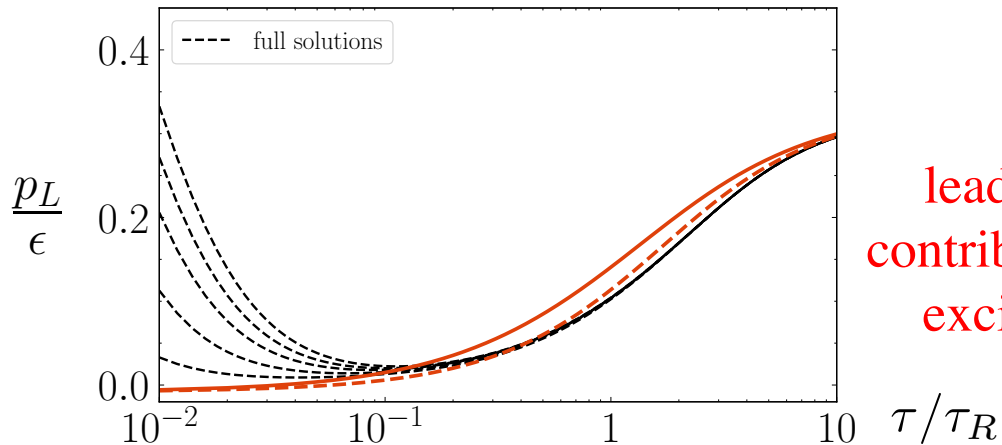
“Slow quench” adiabaticity

- Hamiltonian evolution slow compared to energy gap
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“Fast quench” adiabaticity

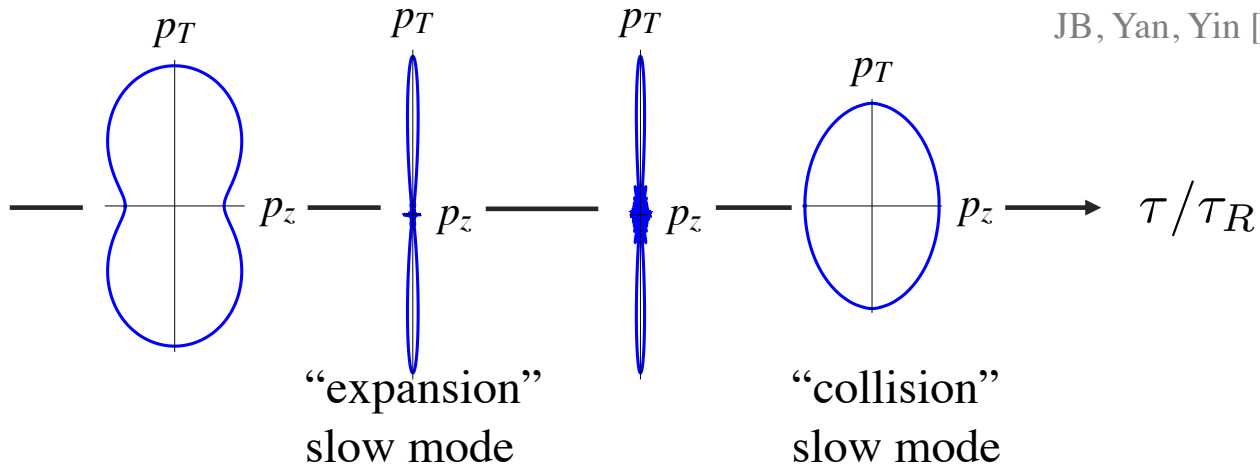
- Matrix element suppressed
- Small at early times because H suppressed by τ

Ground state: far-from-equilibrium slow mode



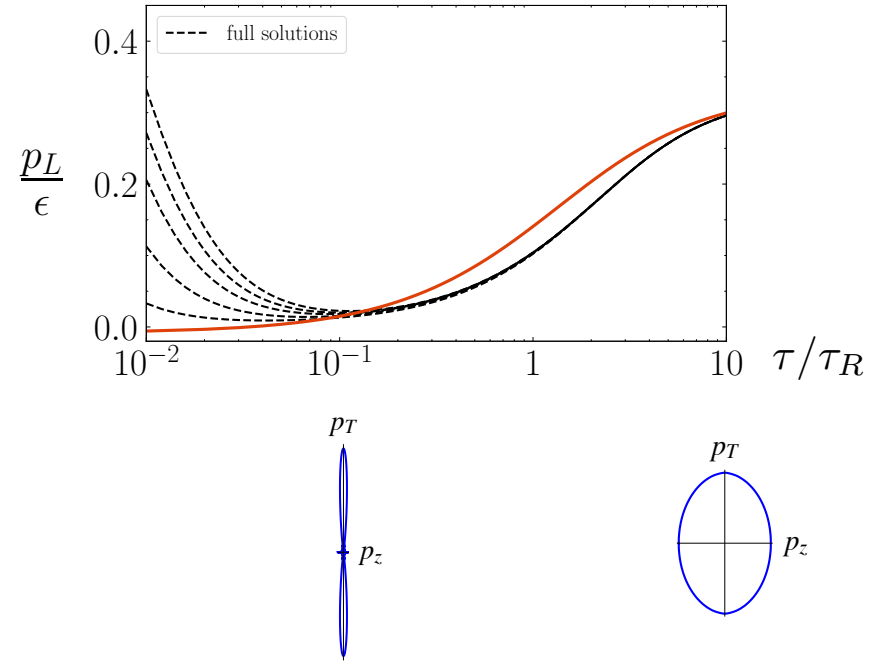
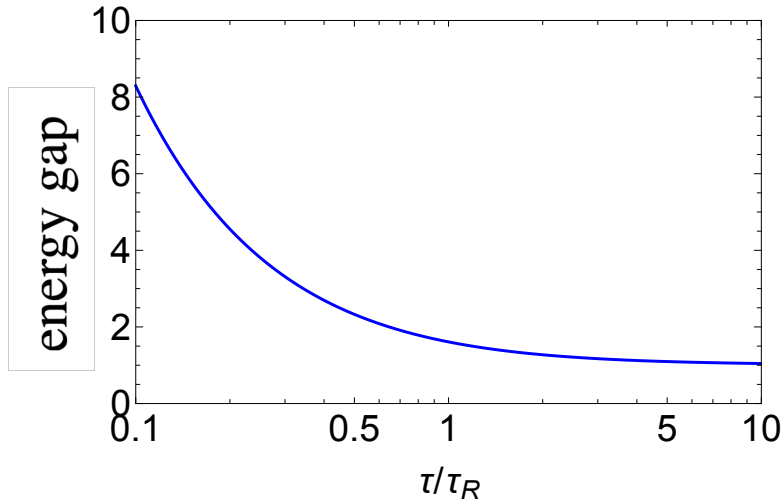
leading order
contributions from
excited states

JB, Yan, Yin [1910.00021]



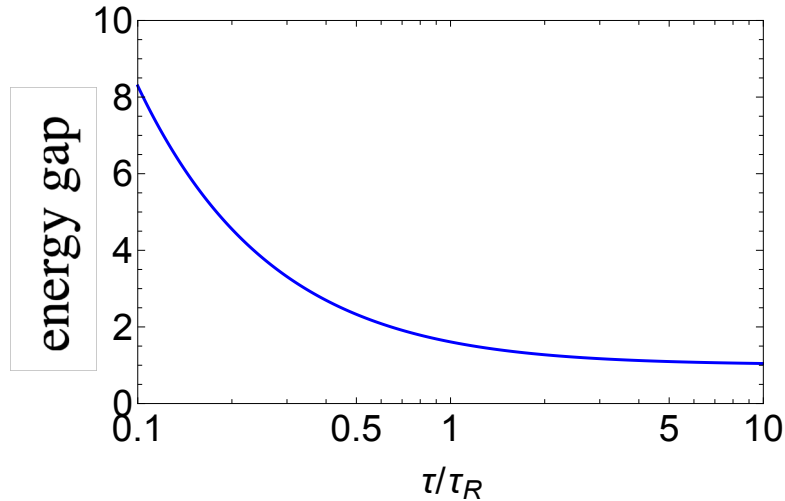
Bjorken expansion

$$P_{2n}(\cos \theta = \frac{p_z}{p})$$



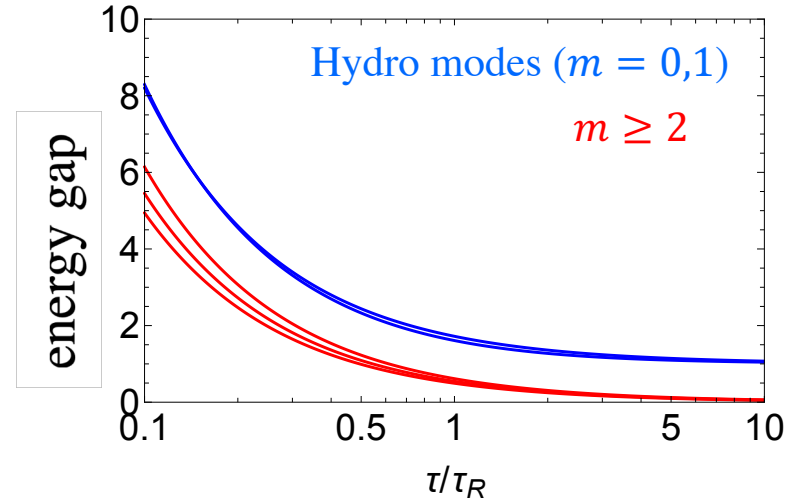
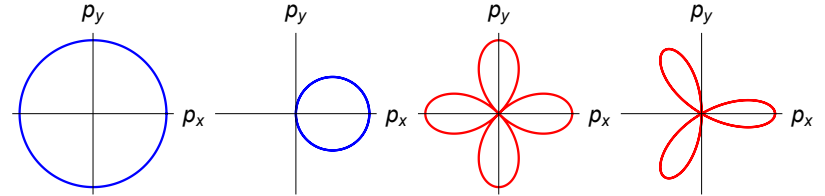
Bjorken expansion

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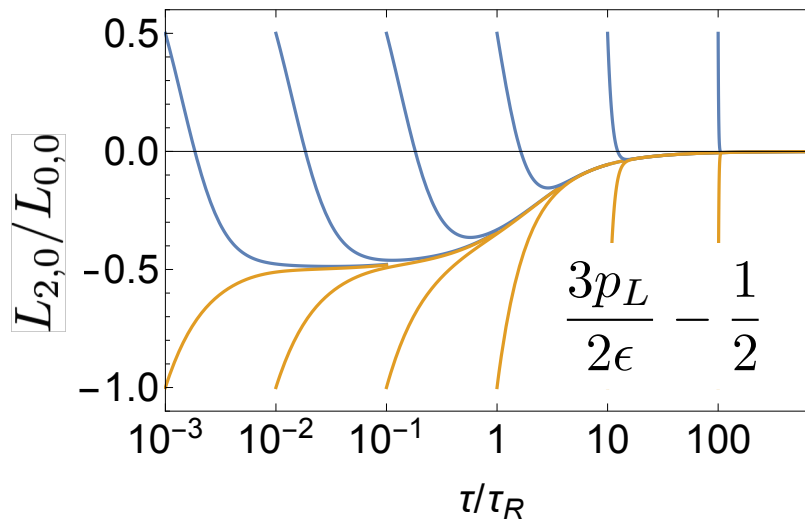
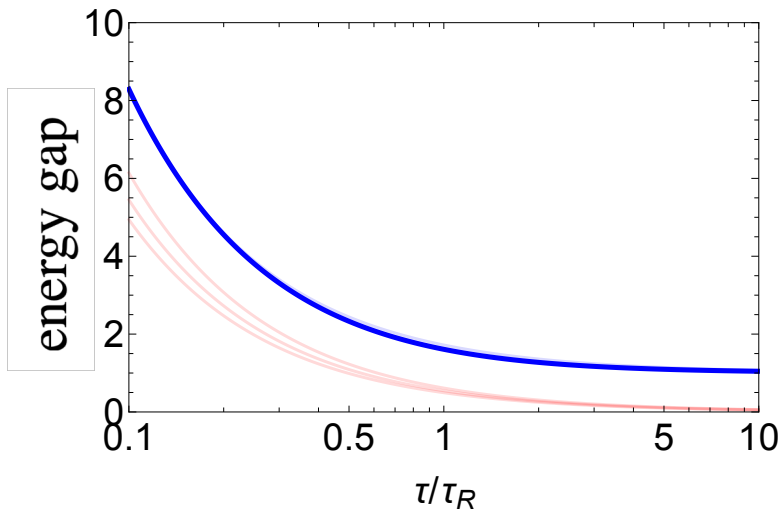


Transverse momentum anisotropy

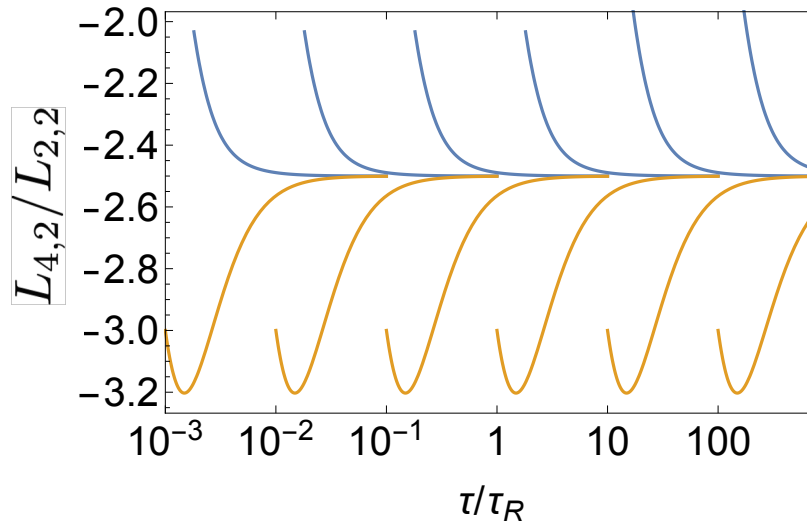
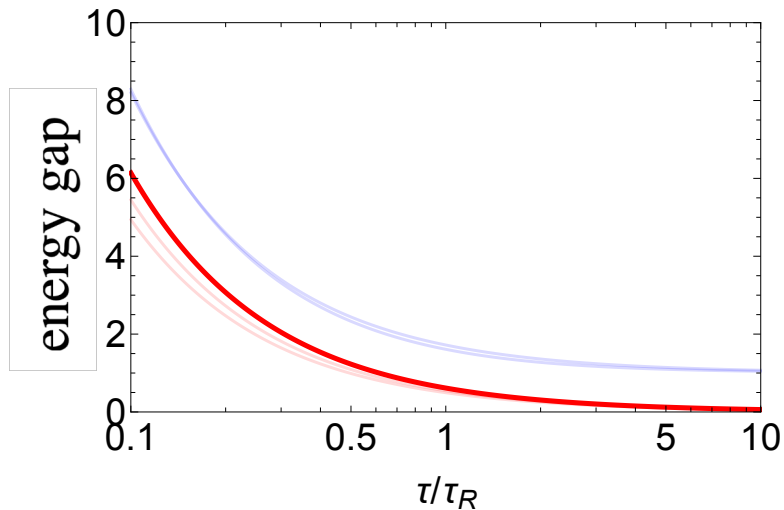
$$Y_l^m(\cos \theta, \phi) = (-1)^m \cos(m\phi) P_l^m(\cos \theta)$$



Beyond Bjorken expansion



Beyond Bjorken expansion



Suggests attractor for non-hydrodynamic modes

Summary

Transition from free-streaming to hydrodynamics described by ground state far-from-equilibrium

New insights for far-from-equilibrium behavior of moments beyond hydrodynamics

