

Constraining neutrino mass using three-point mean relative velocity statistics

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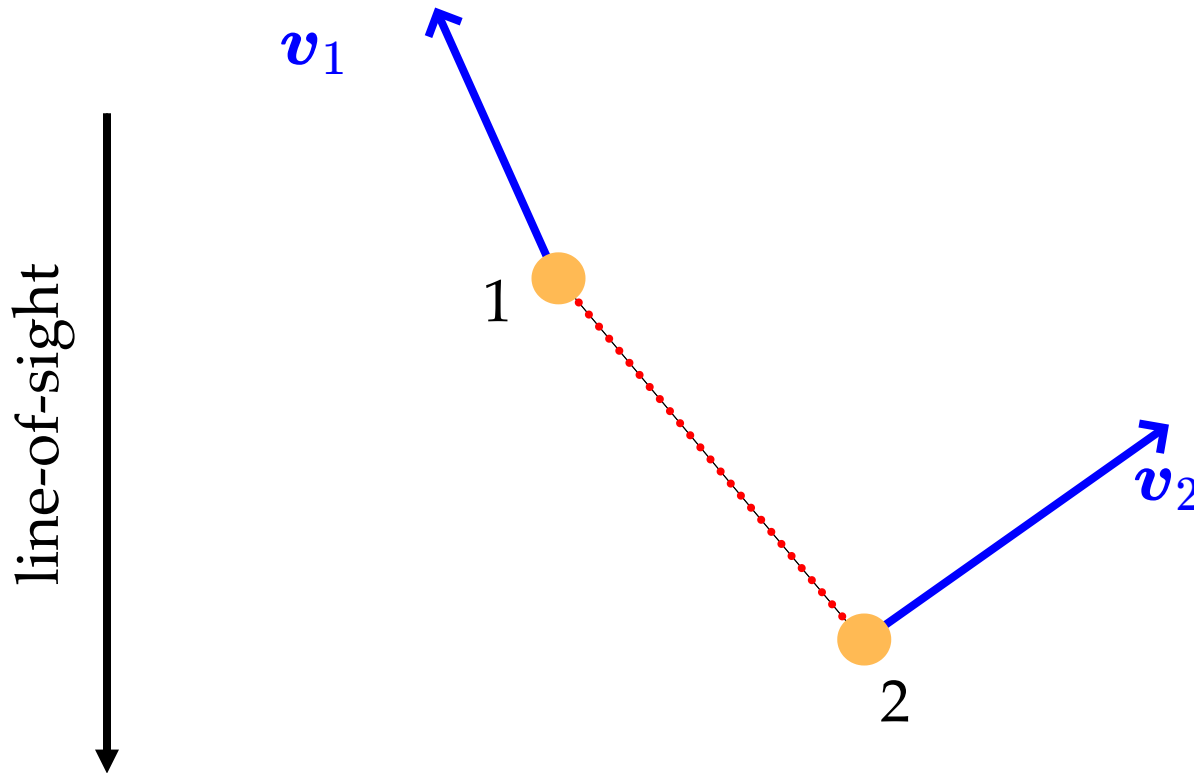
Current constraints on summed neutrino mass

- The Karlsruhe Tritium Neutrino (KATRIN) experiment has reported the first direct detection of sub-eV neutrino mass, with an upper limit on the 'effective neutrino mass' of 0.8 eV. (Aker et al. 2021)
- By combining cosmic microwave background, baryonic acoustic oscillation, and redshift-space galaxy clustering to obtain an upper limit of 0.102 eV. (eBOSS Collaboration 2021)
- Considering extensions of the standard cosmological model, the upper limit becomes less stringent. (e.g. Vagnozzi et al. 2018; Choudhury & Hannestad 2020)



Pairwise velocity

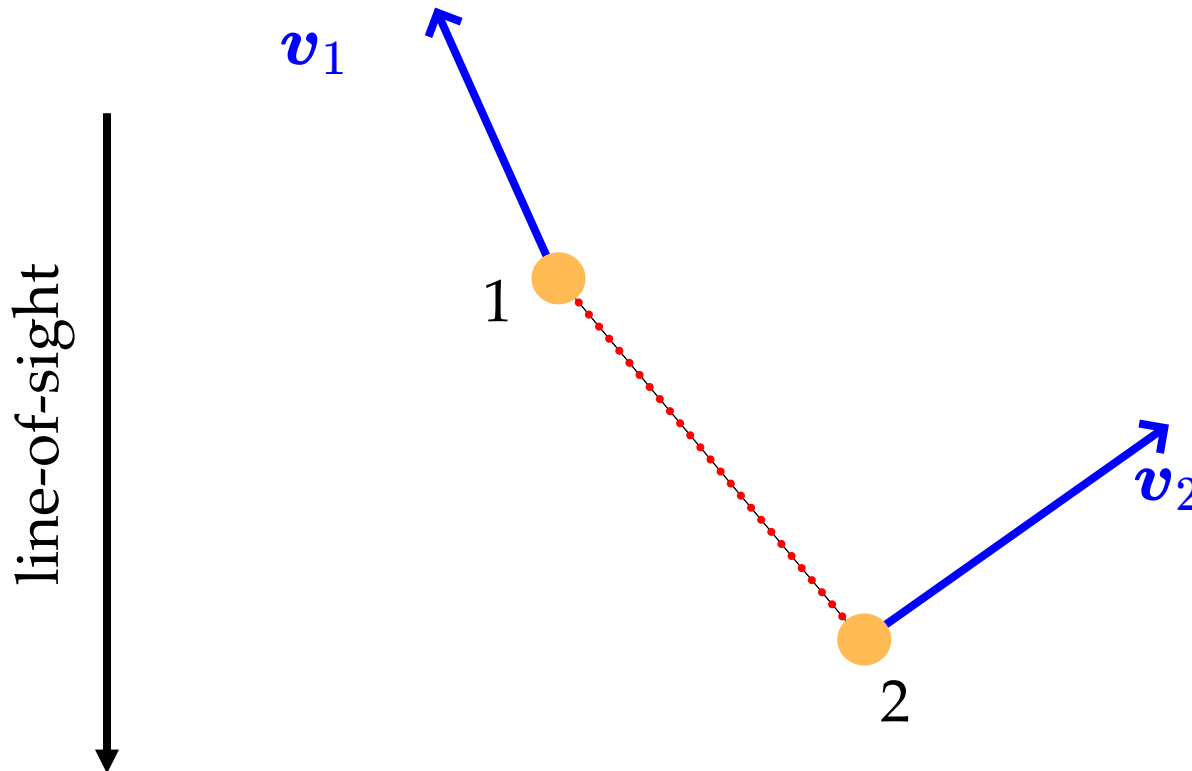
$$w_r(r_{12}) = (\mathbf{v}_2 - \mathbf{v}_1) \cdot \hat{\mathbf{r}}_{12}$$



Pairwise velocity

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$$\langle \mathbf{w}_{12} | \mathbf{r}_{12} \rangle_p = \frac{\langle (1 + \delta_1)(1 + \delta_2)(\mathbf{v}_2 - \mathbf{v}_1) \rangle}{\langle (1 + \delta_1)(1 + \delta_2) \rangle}$$

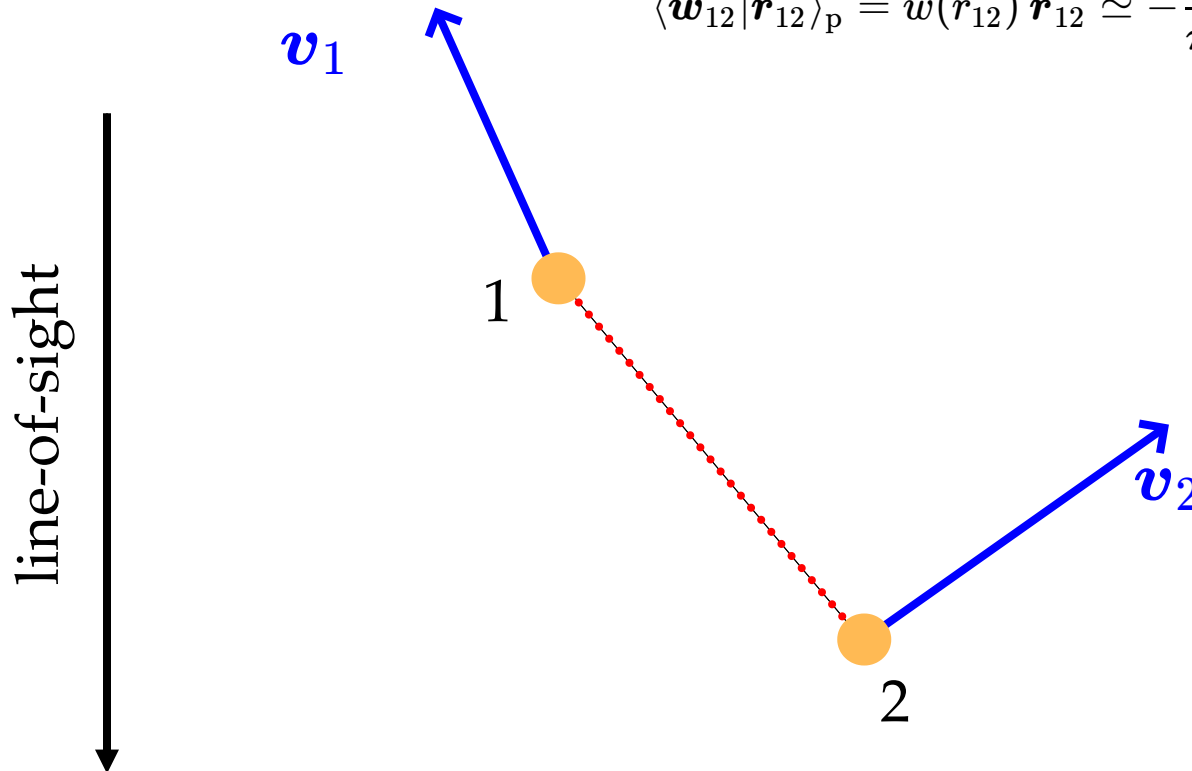


Pairwise velocity

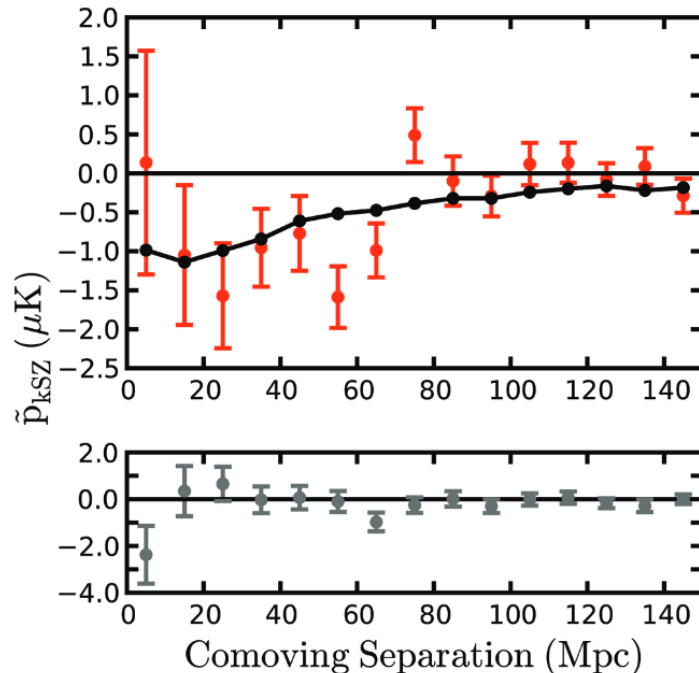
$$w_r(r_{12}) = (\mathbf{v}_2 - \mathbf{v}_1) \cdot \hat{\mathbf{r}}_{12}$$

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$$\langle \mathbf{w}_{12} | \mathbf{r}_{12} \rangle_p = \bar{w}(r_{12}) \hat{\mathbf{r}}_{12} \simeq -\frac{f}{\pi^2} \hat{\mathbf{r}}_{12} \int_0^\infty k j_1(kr_{12}) P(k) dk$$



- The first significant detection of the kinetic Sunyaev-Zeldovich (kSZ) effect was achieved through the mean pairwise mean velocity. (Hand et al. 12)



(Hand et al. 12)

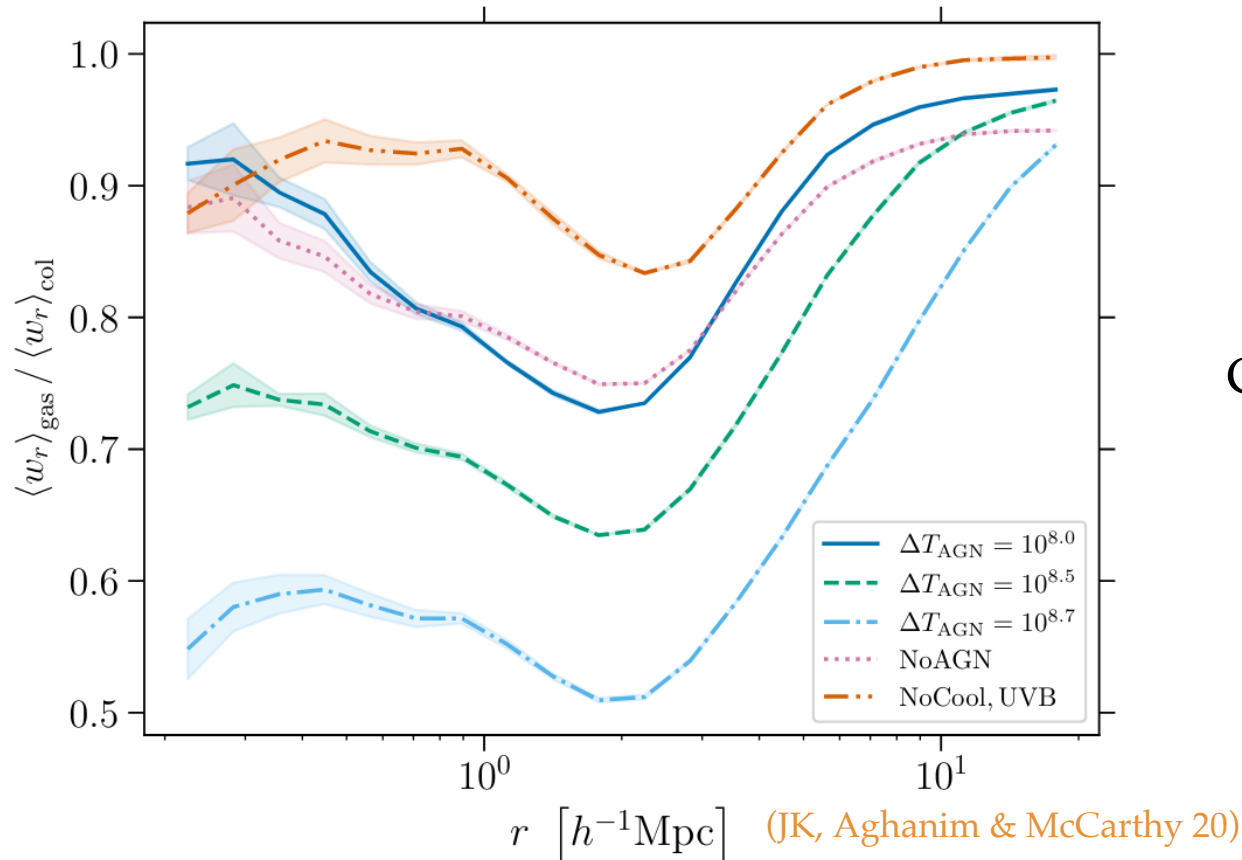
$$\frac{\Delta T^{\text{kSZ}}(r_{12})}{T_{\text{cmb}}} \simeq -\tau \frac{\bar{w}(r_{12})}{c}$$



- The first significant detection of the kinetic Sunyaev-Zeldovich (kSZ) effect was achieved through the mean pairwise mean velocity. (Hand et al. 12)
- The pairwise velocity measurement from kSZ experiments has been shown to be a novel probe to constrain the summed neutrino mass. (Mueller et al. 15)
- The impact of massive neutrinos and its interplay with baryonic physics at non-linear scales. (JK, Aghanim & McCarthy 20)



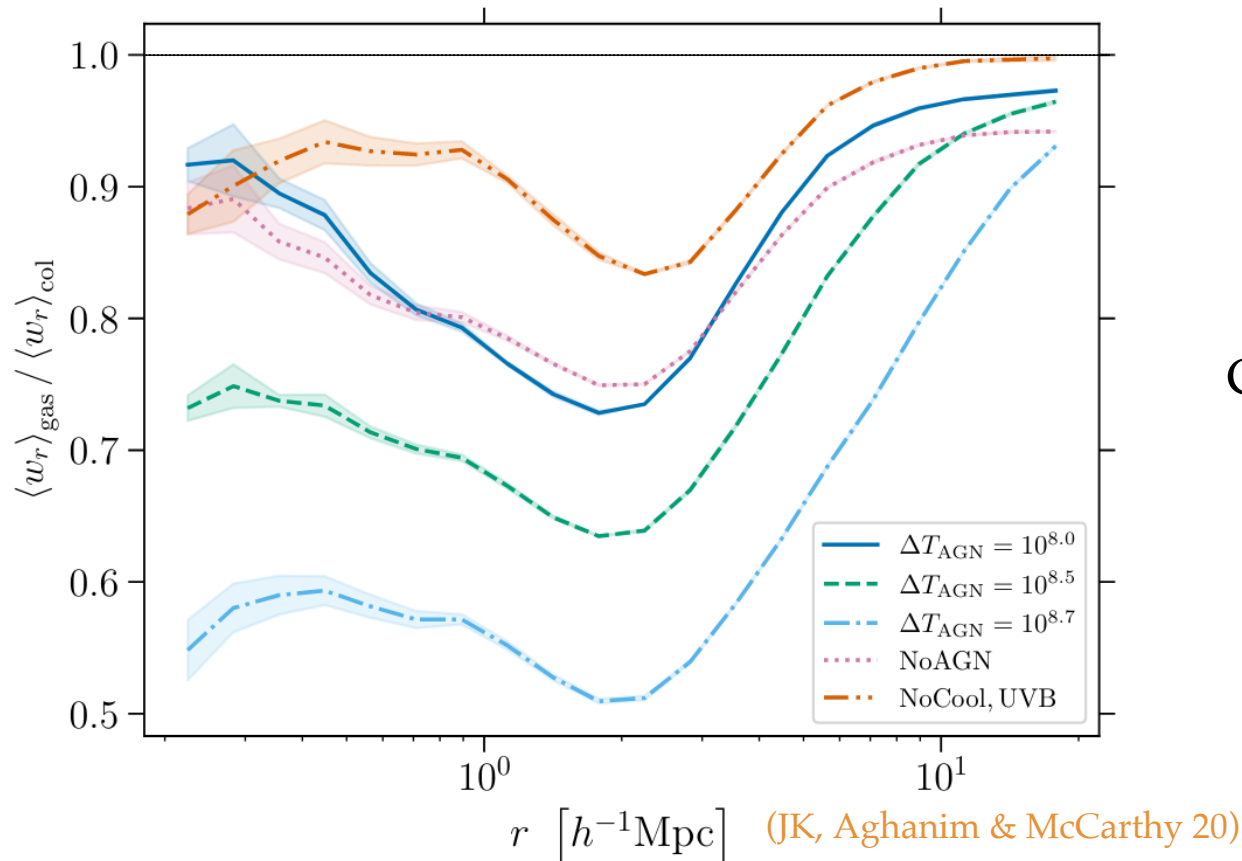
Impact of baryonic physics



Cosmo-OWLS
simulations

- Assumption of gas following the matter is not valid even at scales of roughly 20 Mpc/h.

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Is there any other domain in cosmology where the pairwise velocities are important?

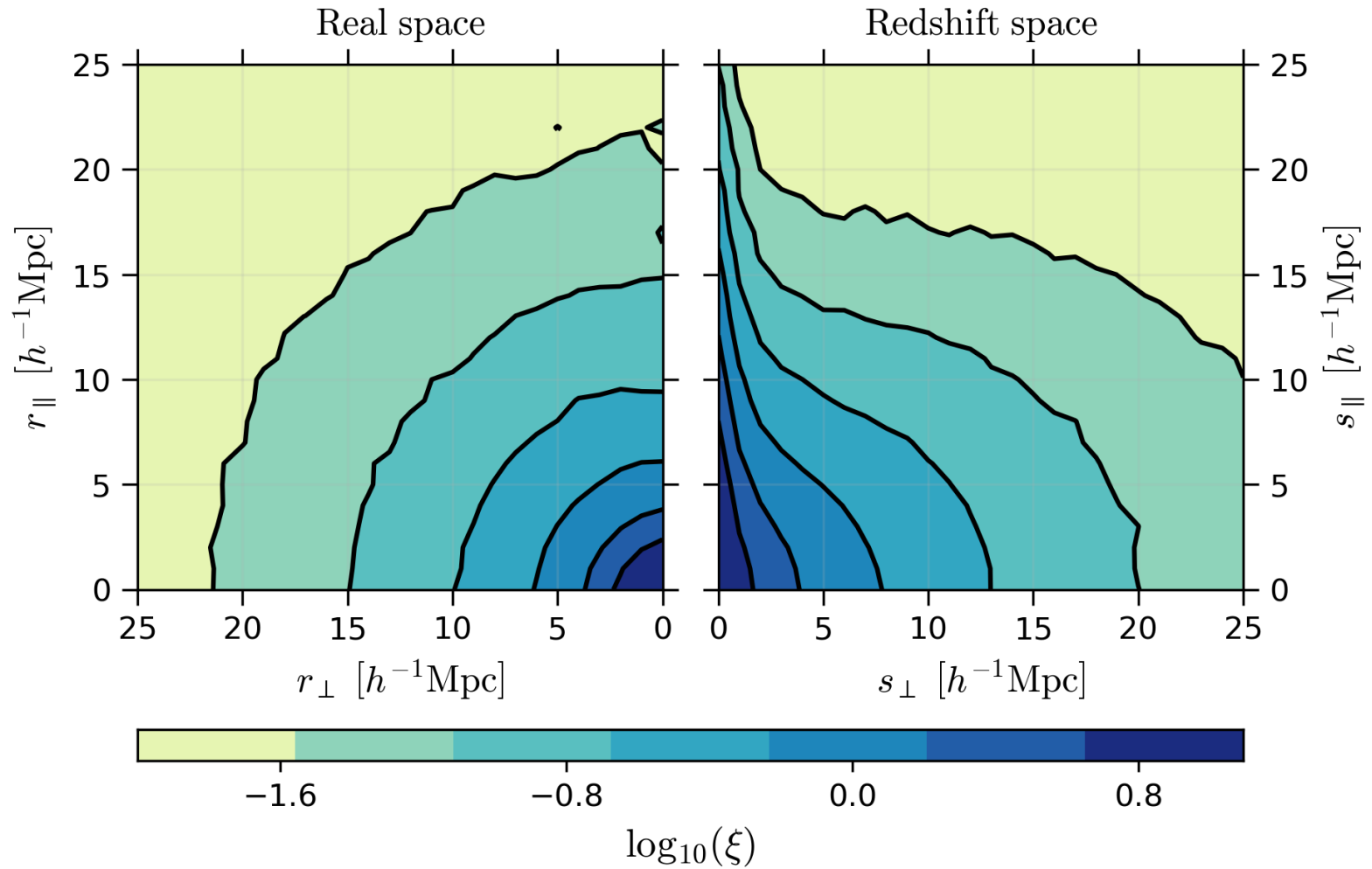


Is there any other domain in cosmology where the pairwise velocities are important?

Yes, in modelling the anisotropic correlation function in redshift space!



Two-point streaming model



Two-point streaming model

$$1 + \xi_s(\mathbf{s}_\perp, \mathbf{s}_\parallel) = \int_{-\infty}^{+\infty} [1 + \xi(r)] \mathcal{P}_{w_\parallel}^{(2)}(w_\parallel | \vec{r}) dr_\parallel$$

(Peebles 80, Fisher 95, Scoccimarro 04,
JK & Porciani 18, Vlah & White 19)

$\xi_s(\mathbf{s}_\perp, \mathbf{s}_\parallel)$: Anisotropic redshift-space correlation function

$\xi_r(r)$: Isotropic real-space correlation function

$\mathcal{P}_{w_\parallel}^{(2)}(w_\parallel | \vec{r})$: Relative pairwise line-of-sight velocity distribution



Can we generalise streaming model to three-point and higher order statistics?



Can we generalise streaming model to three-point and higher order statistics?

Yes we can!



n -point streaming model

- Generalised the streaming model framework to n -point statistics.

$$\mathcal{G}_n = \int \mathcal{F}_n \mathcal{P}_{\mathbf{w}_{\parallel}}^{(n)} d\mathbf{w}_{12\parallel} \dots d\mathbf{w}_{mn\parallel}$$

where,

(JK & Porciani 20)

\mathcal{G}_n : the (anisotropic) n -point full CF in redshift space

\mathcal{F}_n : the (isotropic) n -point full CF in real space

$\mathcal{P}_{\mathbf{w}_{\parallel}}^{(n)}$: the joint probability density of $n-1$ relative line-of-sight peculiar velocity

The equation is exact under distant observer approximation



Full correlation function in three-point

$$\langle \delta_1 \delta_2 \delta_3 \rangle = \text{[diagram showing five terms: three unconnected dots, two pairs of connected dots, and a triangle of three connected dots]}$$

Credits: Bernardeau et al. 01

Thus the three-point streaming model is

$$\begin{aligned} & 1 + \xi_s(s_{12\parallel}, s_{12\perp}) + \xi_s(s_{23\parallel}, s_{23\perp}) + \xi_s(\check{s}_{31\parallel}, s_{31\perp}) + \zeta_s(\mathbf{s}_{12}, \mathbf{s}_{23}) \\ &= \int [1 + \xi(\check{r}_{12}) + \xi(\check{r}_{23}) + \xi(\check{r}_{31}) + \zeta(\check{r}_{12}, \check{r}_{23}, \check{r}_{31})] \\ & \quad \mathcal{P}_{\mathbf{w}_{\parallel}}^{(3)}(s_{12\parallel} - r_{12\parallel}, s_{23\parallel} - r_{23\parallel} | \check{r}_{12}, \check{r}_{23}) dr_{12\parallel} dr_{23\parallel} \end{aligned}$$

(JK & Porciani 20)



Full correlation function in three-point

$$\langle \delta_1 \delta_2 \delta_3 \rangle = \text{[Diagram showing five terms representing different topologies for the three-point correlation function: three isolated points, two points connected by a line, one point connected to two others, and a fully connected triangle.]}$$

Credits: Bernardeau et al. 01

Thus the three-point streaming model is

$$\begin{aligned}
 & 1 + \xi_s(s_{12\parallel}, s_{12\perp}) + \xi_s(s_{23\parallel}, s_{23\perp}) + \xi_s(\check{s}_{31\parallel}, s_{31\perp}) + \zeta_s(\mathbf{s}_{12}, \mathbf{s}_{23}) \\
 &= \int [1 + \xi(\check{r}_{12}) + \xi(\check{r}_{23}) + \xi(\check{r}_{31}) + \zeta(\check{r}_{12}, \check{r}_{23}, \check{r}_{31})] \\
 & \quad \mathcal{P}_{\mathbf{w}_{\parallel}}^{(3)}(s_{12\parallel} - r_{12\parallel}, s_{23\parallel} - r_{23\parallel} | \check{r}_{12}, \check{r}_{23}) dr_{12\parallel} dr_{23\parallel}
 \end{aligned}$$

(JK & Porciani 20)



Full correlation function in three-point

$$\langle \delta_1 \delta_2 \delta_3 \rangle = \text{[Diagram 1]} + \text{[Diagram 2]} + \text{[Diagram 3]} + \text{[Diagram 4]} + \text{[Diagram 5]}$$

The diagram shows five terms representing different topologies for the three-point correlation function. The first term is three unconnected points. The second and third terms are two points connected by a line, with a third point unconnected. The fourth term is two points connected by a line, with a third point connected to one of them. The fifth term is a fully connected triangle.

Credits: Bernardeau et al. 01

Thus the three-point streaming model is

$$1 + \xi_s(s_{12\parallel}, s_{12\perp}) + \xi_s(s_{23\parallel}, s_{23\perp}) + \xi_s(\check{s}_{31\parallel}, s_{31\perp}) + \zeta_s(\mathbf{s}_{12}, \mathbf{s}_{23})$$

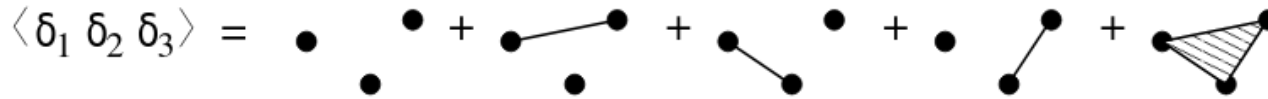
$$= \int [1 + \xi(\check{r}_{12}) + \xi(\check{r}_{23}) + \xi(\check{r}_{31}) + \zeta(\check{r}_{12}, \check{r}_{23}, \check{r}_{31})]$$

$$\mathcal{P}_{w_{\parallel}}^{(3)}(s_{12\parallel} - r_{12\parallel}, s_{23\parallel} - r_{23\parallel} | \check{r}_{12}, \check{r}_{23}) dr_{12\parallel} dr_{23\parallel}$$

(JK & Porciani 20)



Full correlation function in three-point



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Thus the three-point streaming model is

$$\begin{aligned}
 & 1 + \xi_s(s_{12\parallel}, s_{12\perp}) + \xi_s(s_{23\parallel}, s_{23\perp}) + \xi_s(\check{s}_{31\parallel}, s_{31\perp}) + \zeta_s(\mathbf{s}_{12}, \mathbf{s}_{23}) \\
 &= \int [1 + \xi(\check{r}_{12}) + \xi(\check{r}_{23}) + \xi(\check{r}_{31}) + \zeta(\check{r}_{12}, \check{r}_{23}, \check{r}_{31})] \\
 & \quad \mathcal{P}_{w_{\parallel}}^{(3)}(s_{12\parallel} - r_{12\parallel}, s_{23\parallel} - r_{23\parallel} | \check{r}_{12}, \check{r}_{23}) \, dr_{12\parallel} \, dr_{23\parallel}
 \end{aligned}$$

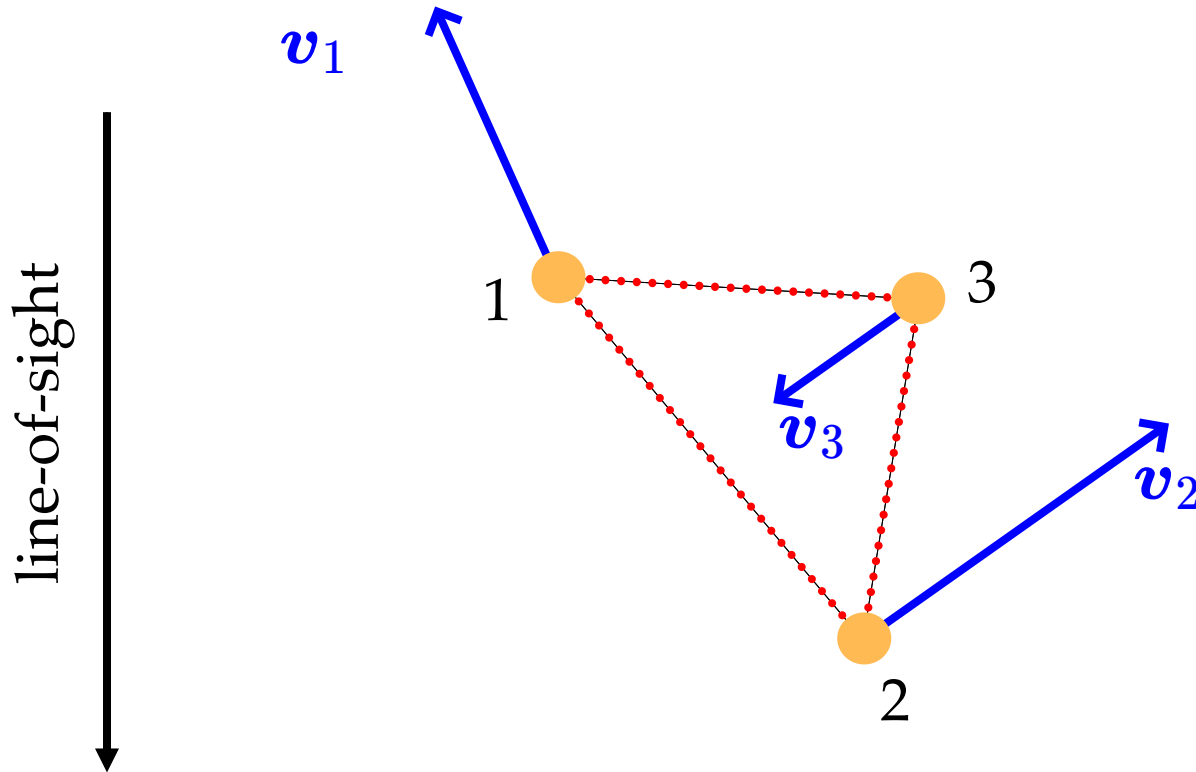
(JK & Porciani 20)



What does $\mathcal{P}_{w_{\parallel}}^{(3)}$ look like?

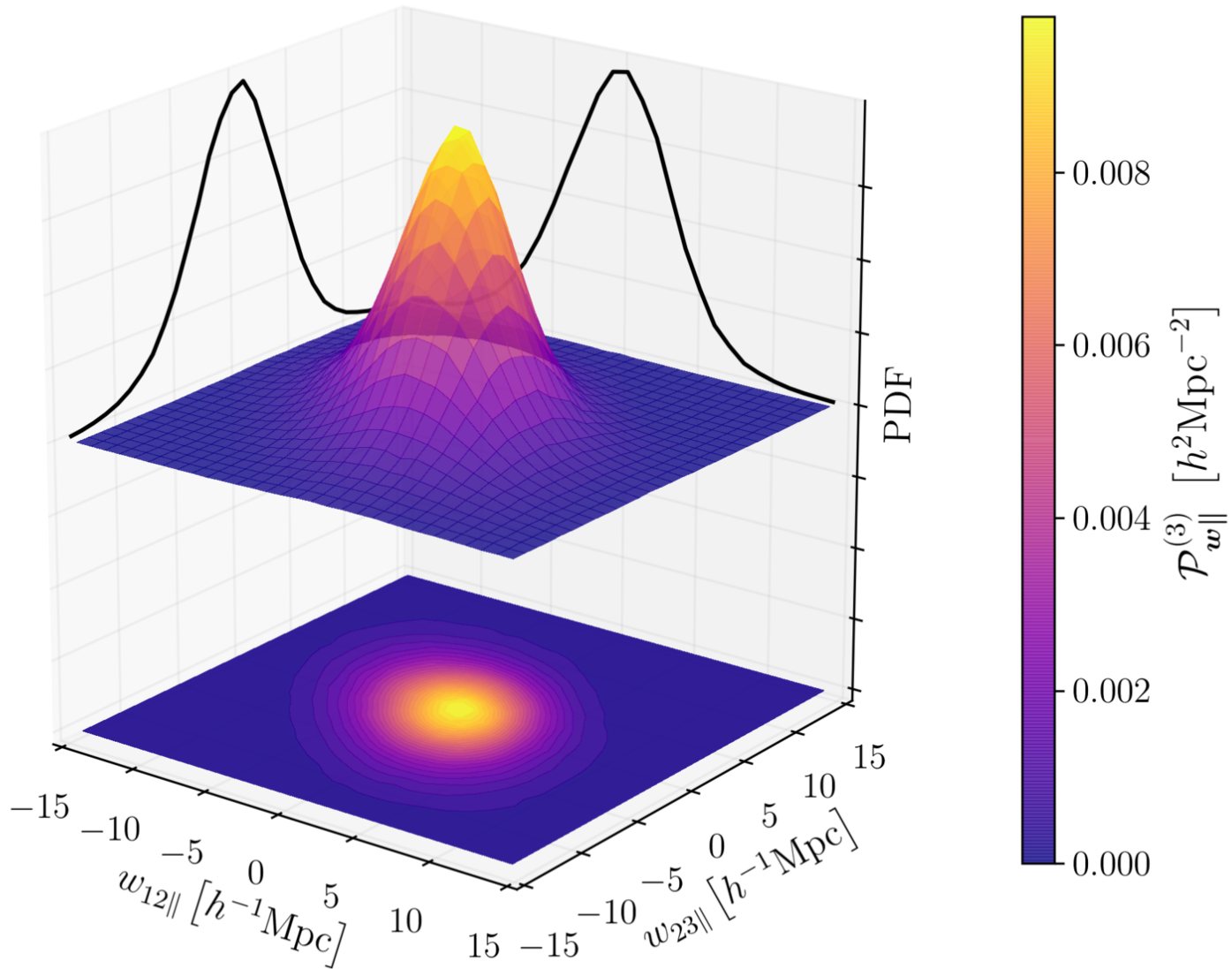


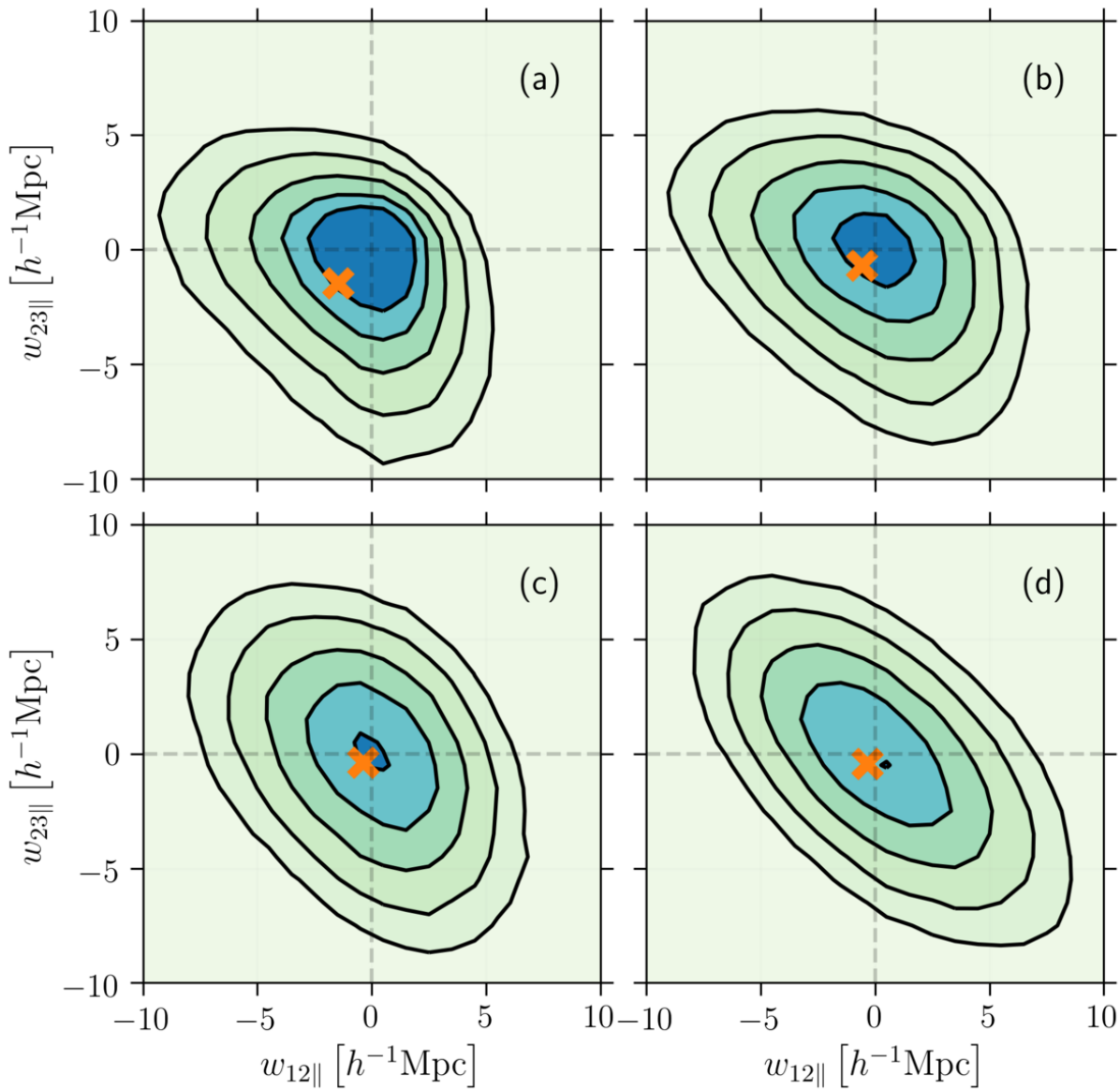
Three-point relative velocity



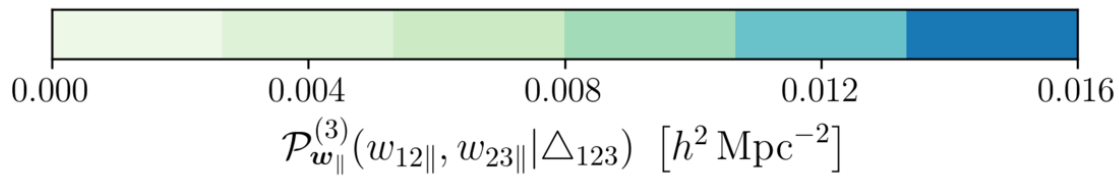
$$R_{12}(r_{12}, r_{23}, r_{31}) = (\mathbf{v}_2 - \mathbf{v}_1) \cdot \hat{\mathbf{r}}_{12}$$







(JK & Porciani 20)



Can we predict the velocity moments?

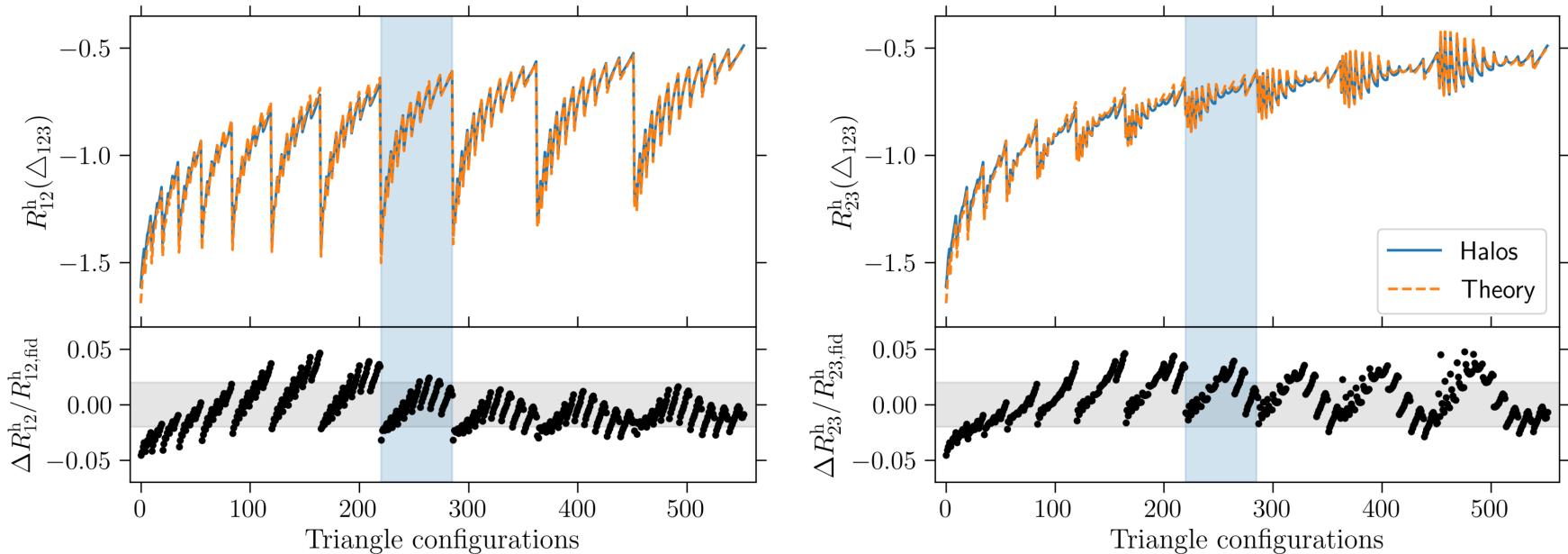


Halos - Mean relative velocity between pairs in a triplet

Triangular configurations: $50 \leq r_{31} \leq r_{23} \leq r_{12} \leq 120 h^{-1} \text{Mpc}$

Minimum halo mass: $M_h > 5 \times 10^{13} h^{-1} M_\odot$

15,000 simulations of the Quijote suite of simulations



$$R_{ij}^h(\Delta_{123}, M_h) = b(M_h) R_{ij}(\Delta_{123})$$

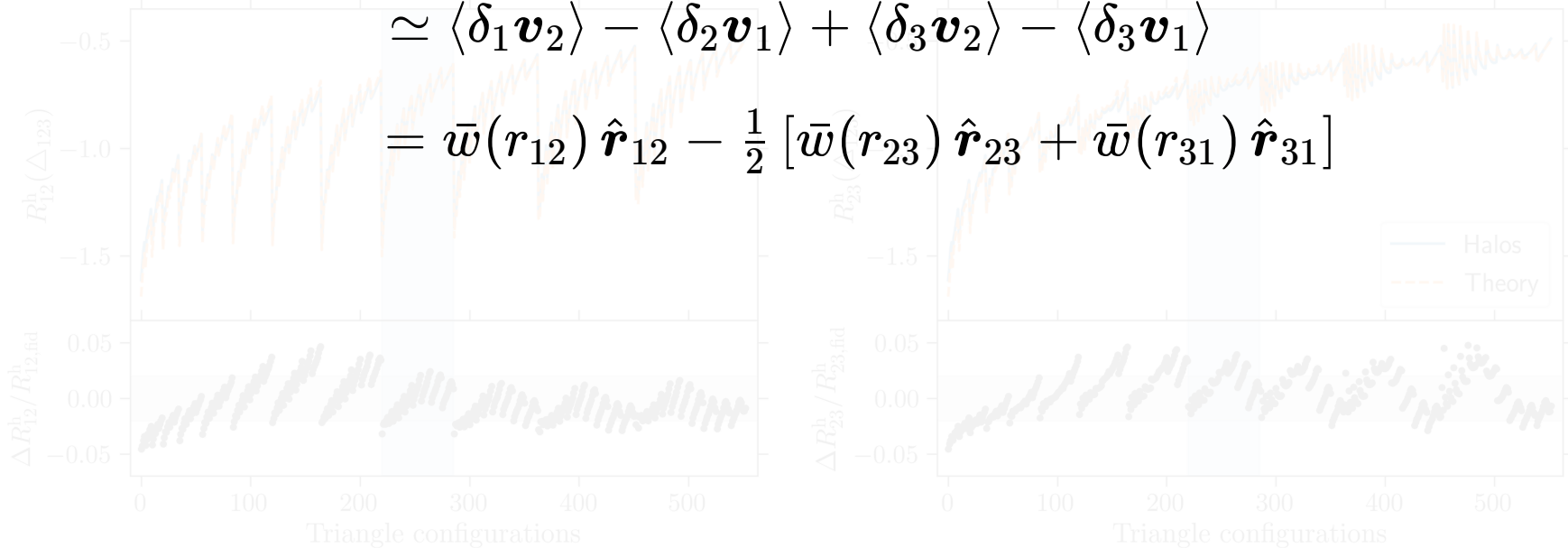
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$$\langle \mathbf{w}_{12} | \Delta_{123} \rangle_t = \frac{\langle (1 + \delta_1)(1 + \delta_2)(1 + \delta_3)(\mathbf{v}_2 - \mathbf{v}_1) \rangle}{\langle (1 + \delta_1)(1 + \delta_2)(1 + \delta_3) \rangle}$$

$$\simeq \langle \delta_1 \mathbf{v}_2 \rangle - \langle \delta_2 \mathbf{v}_1 \rangle + \langle \delta_3 \mathbf{v}_2 \rangle - \langle \delta_3 \mathbf{v}_1 \rangle$$

$$= \bar{w}(r_{12}) \hat{\mathbf{r}}_{12} - \frac{1}{2} [\bar{w}(r_{23}) \hat{\mathbf{r}}_{23} + \bar{w}(r_{31}) \hat{\mathbf{r}}_{31}]$$



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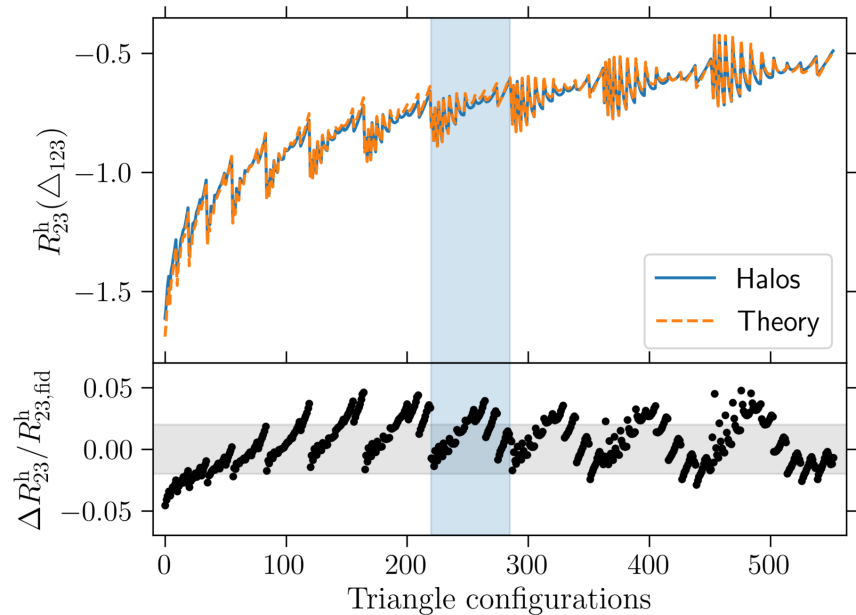
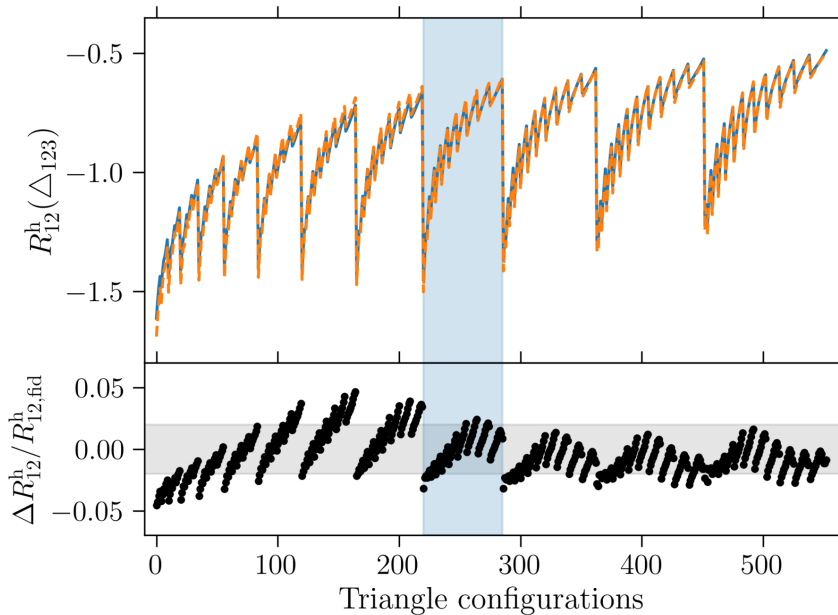


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$$R_{12}(\Delta_{123}) = \bar{w}(r_{12}) - \frac{1}{2} \left[\bar{w}(r_{23}) \cos \chi - \bar{w}(r_{31}) \frac{r_{12} + r_{23} \cos \chi}{\sqrt{r_{12}^2 + r_{23}^2 + 2r_{12}r_{23} \cos \chi}} \right]$$

Can we use the relative velocity statistics to constrain cosmology?



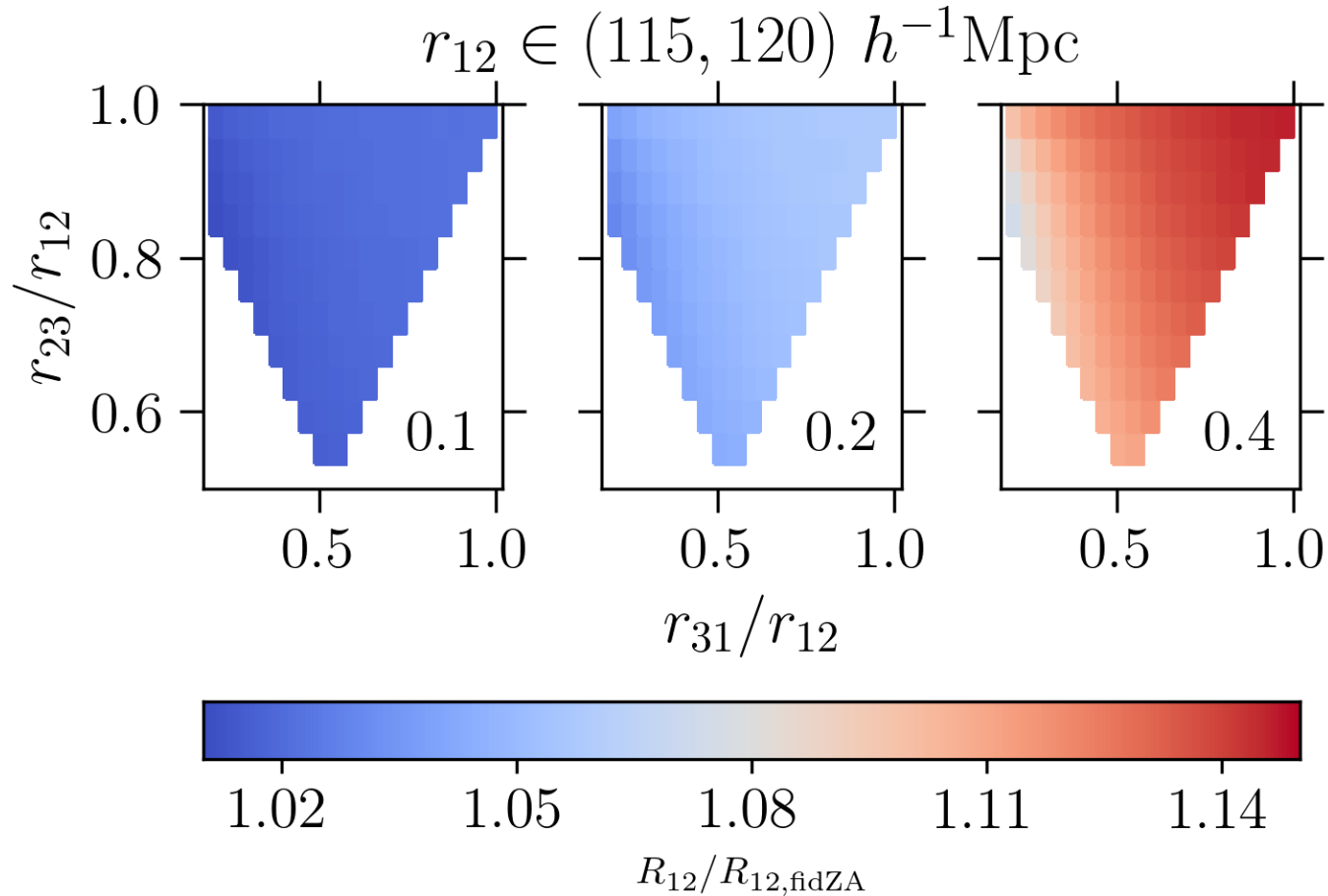
Halos - Mean relative velocity between pairs in a triplet

- Fisher matrix formalism
- Compute the covariance matrix for the relative velocity statistics using 15,000 simulations from the Quijote suite.
- Derivatives also computed directly using the Quijote suite.

-
- The Quijote simulations is a suite of 44,100 full N-body simulations. (Villaescusa-Navarro et al. 20)
 - This work runs parallel to the quantification of the information content from the halo redshift-space bispectrum using the Quijote suite. (Hahn et al. 20)



Effect of massive neutrinos

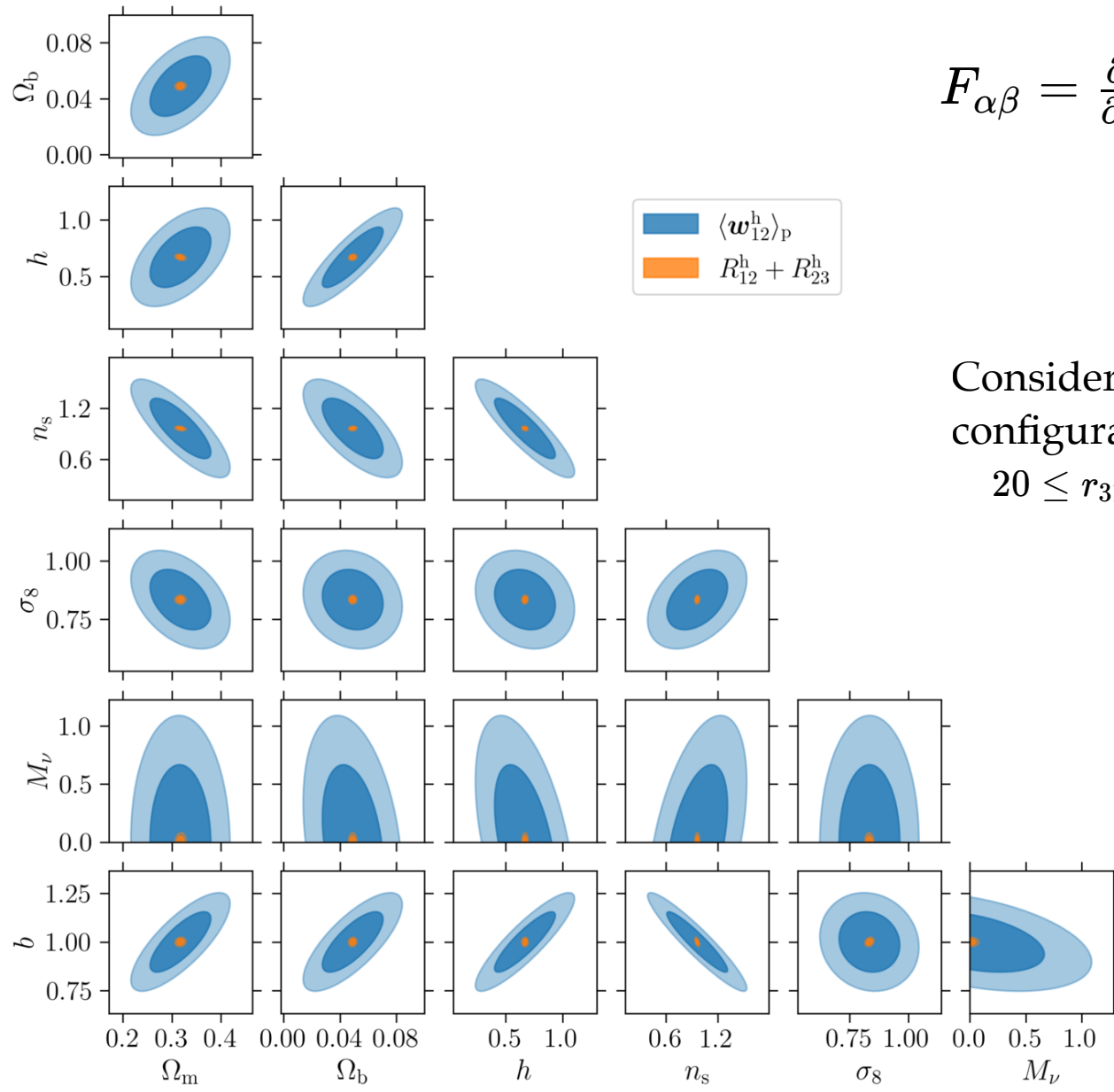


(JK & Aghanim 21)



$$F_{\alpha\beta} = \frac{\partial \mathcal{S}}{\partial \theta_\alpha} \cdot \hat{\mathbf{C}}^{-1} \cdot \frac{\partial \mathcal{S}^\top}{\partial \theta_\beta}$$





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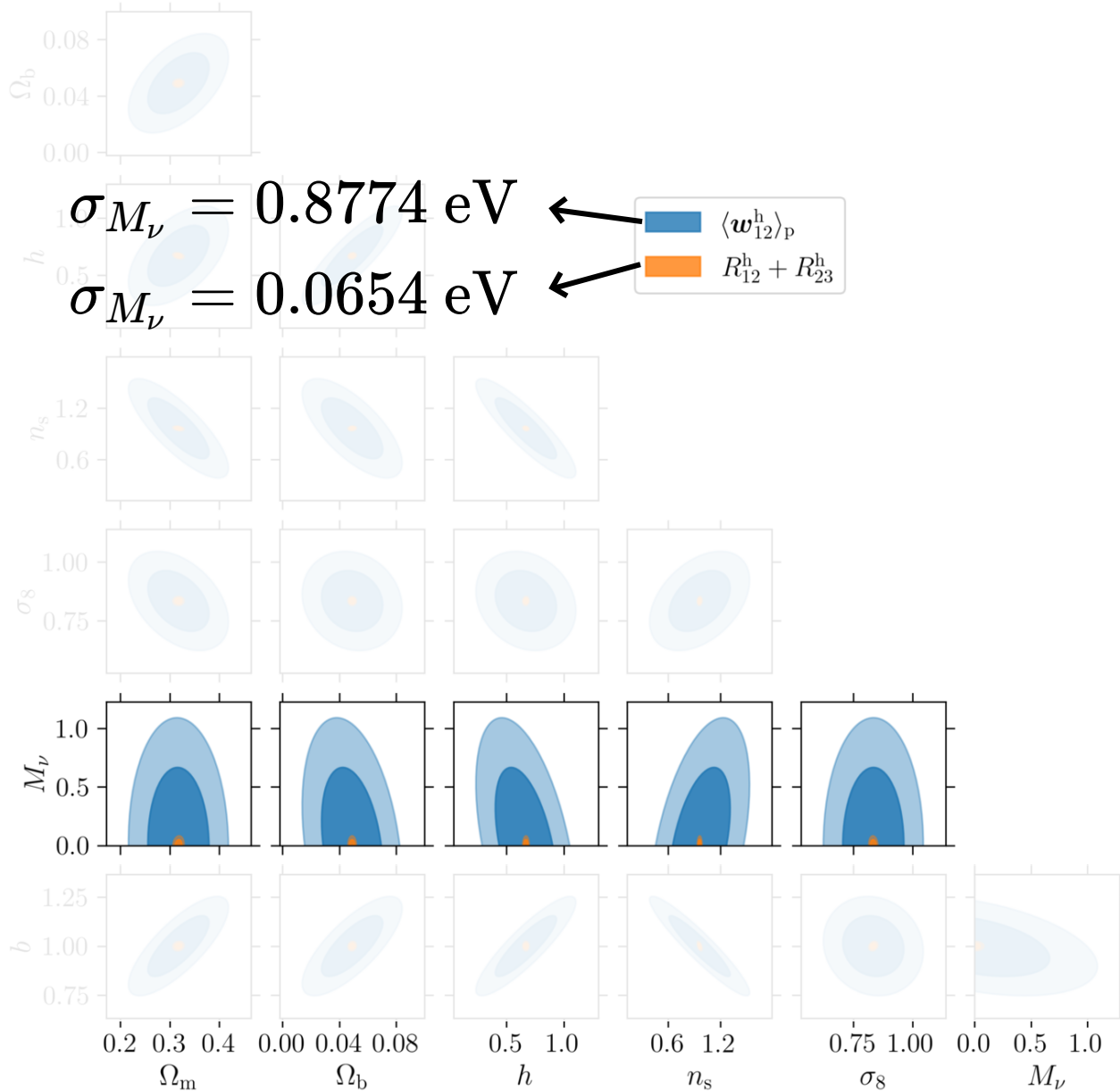
Considering all triangular configurations:

$$20 \leq r_{31} \leq r_{23} \leq r_{12} \leq 120$$



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But is this really competitive with
clustering statistics?



Summary statistics	Matter density	Baryon density	Hubble parameter (h)	Spectral index	Sigma 8	Summed neutrino mass (eV)
mean relative velocity	0.0091	0.0024	0.0226	0.0221	0.0156	0.0655
power spectrum multipoles	2.6	4.8	4.9	5.7	2.3	4.5
bispectrum monopole	1.2	1.7	1.7	1.5	0.9	0.8

(Hahn et al. 20)

* Numbers within the orange box denotes the factor of improvement of 1 sigma constraint from three-point relative velocity with respect to the corresponding clustering statistics.



What about the optical depth degeneracy?



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Joseph Kuruvilla - Ateliers action dark energy 2021



We saw earlier that $\frac{\Delta T^{\text{kSZ}}(r_{12})}{T_{\text{cmb}}} \simeq -\tau \frac{\bar{w}(r_{12})}{c}$ thus $\Delta T^{\text{kSZ}} \propto \tau f \sigma_8^2$



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- One of the attempt to break this degeneracy has been proposed by using fast radio bursts (FRB).

(Madhavacheril et al. 19)

- However I introduce a new estimator which is independent of optical depth, and also sigma 8.

(Kuruvilla, 2021, to be submitted)



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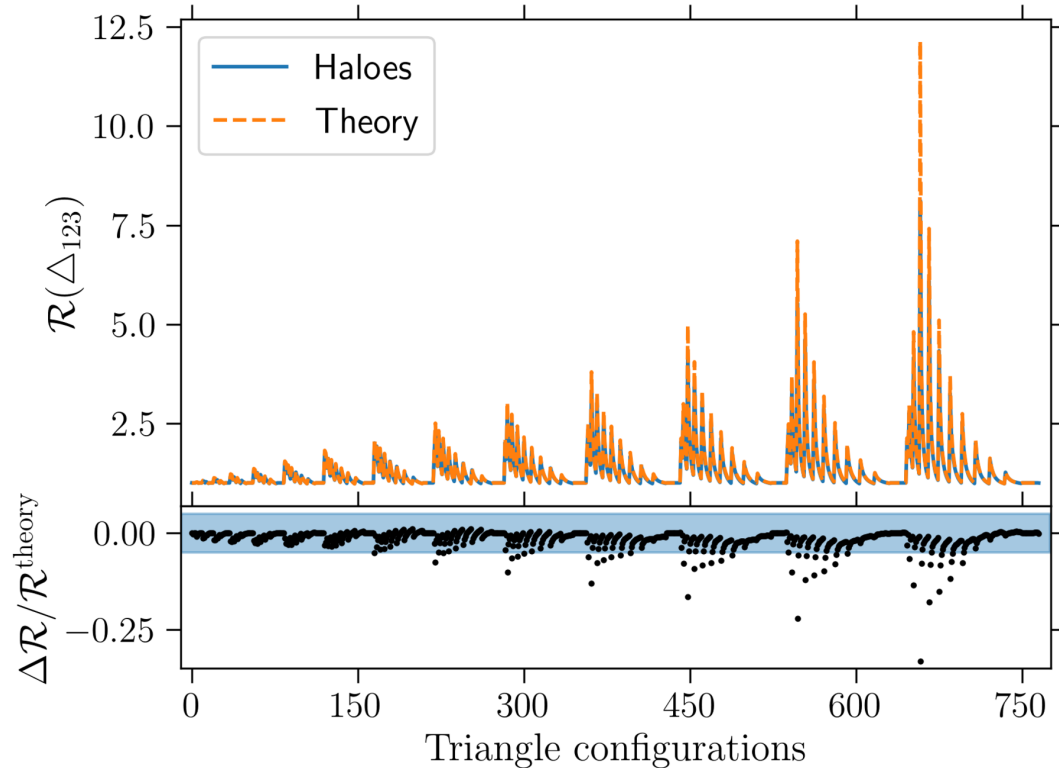
(Kuruvilla, 2021, to be submitted)

$$\mathcal{R}(\Delta_{123}) = \frac{\Delta T_{12}^{\text{kSZ}}(\Delta_{123})}{\Delta T_{23}^{\text{kSZ}}(\Delta_{123})} \equiv \frac{R_{12}^{\text{h}}(\Delta_{123})}{R_{23}^{\text{h}}(\Delta_{123})} = \frac{R_{12}(\Delta_{123})}{R_{23}(\Delta_{123})}$$



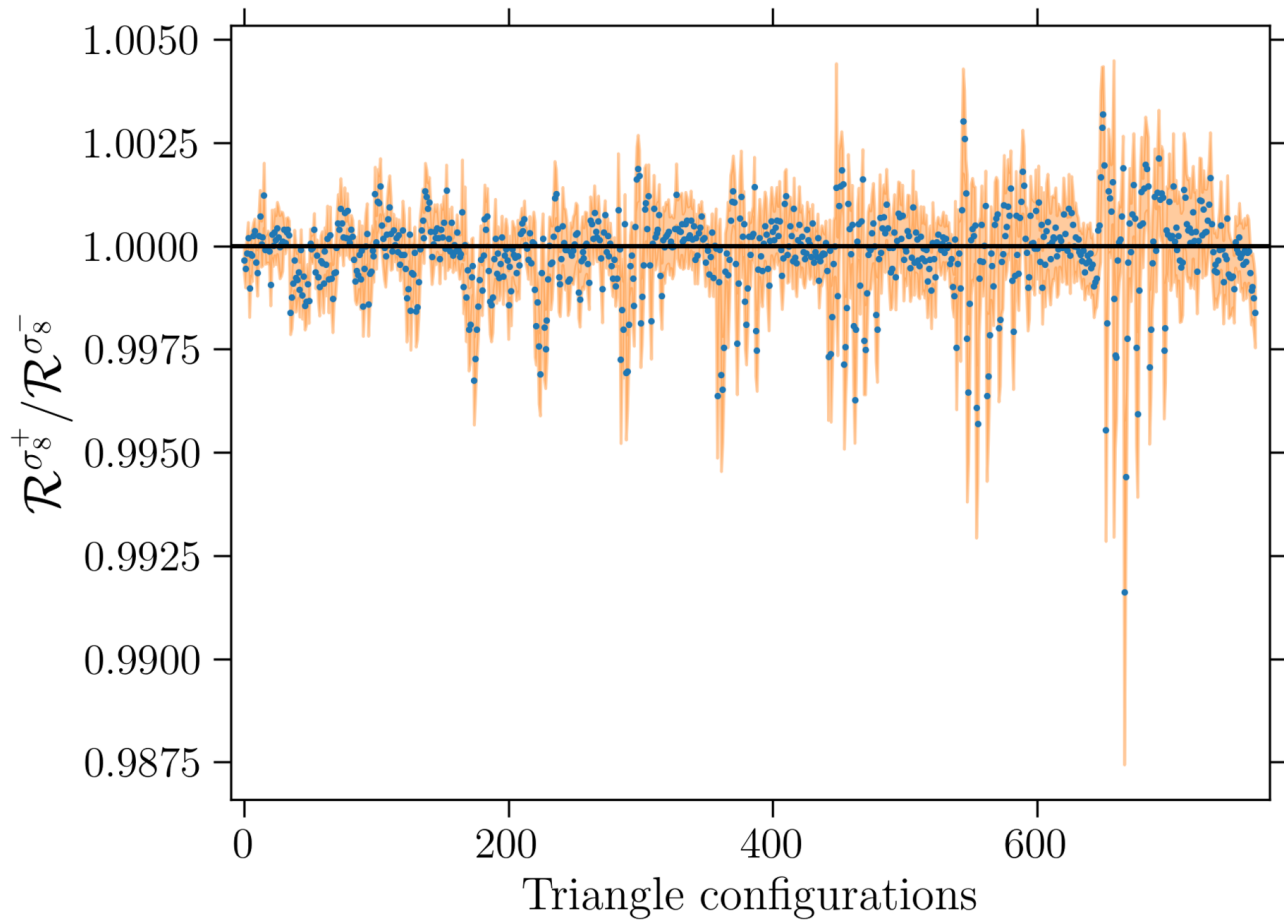
Triangular configurations: $40 \leq r_{31} \leq r_{23} \leq r_{12} \leq 120 h^{-1} \text{Mpc}$

Minimum halo mass: $M_h > 5 \times 10^{13} h^{-1} M_\odot$

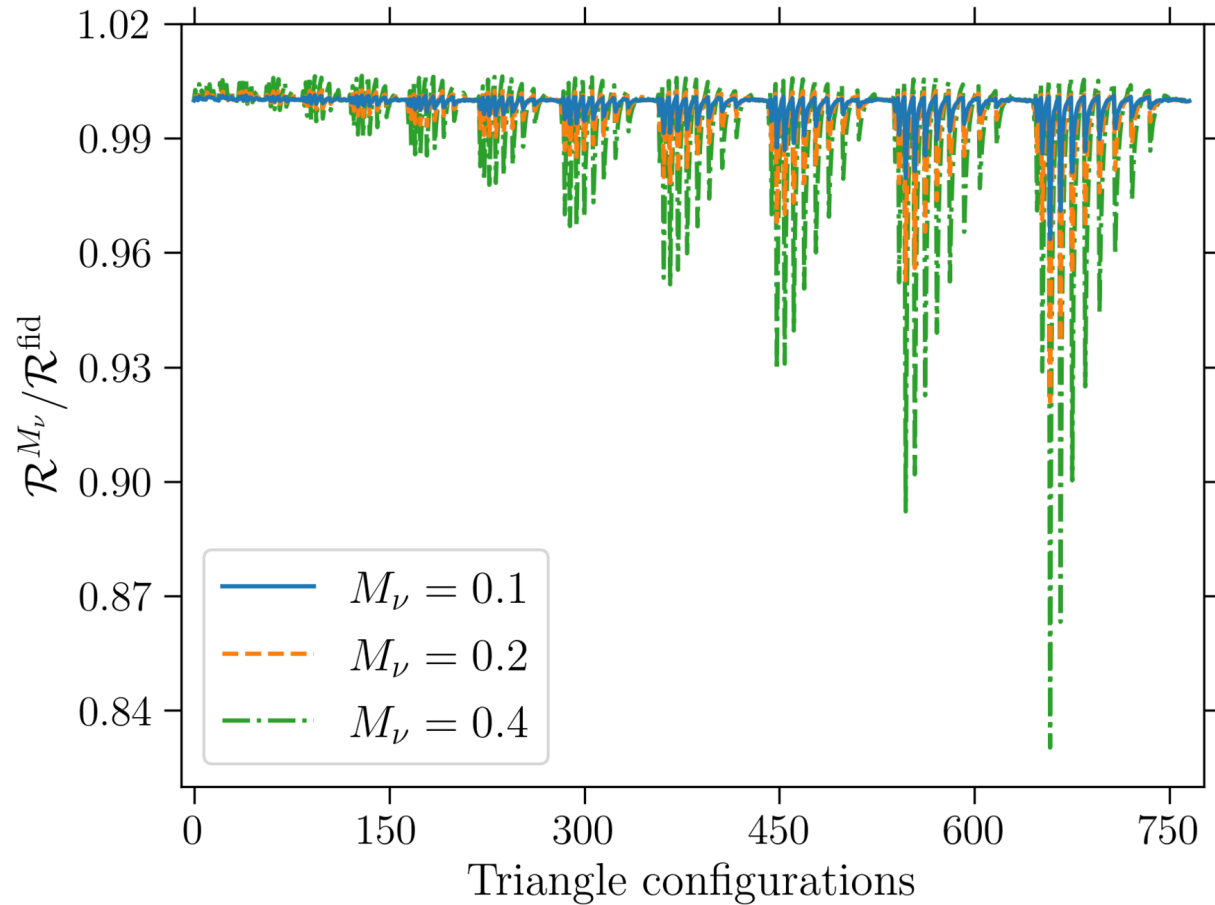


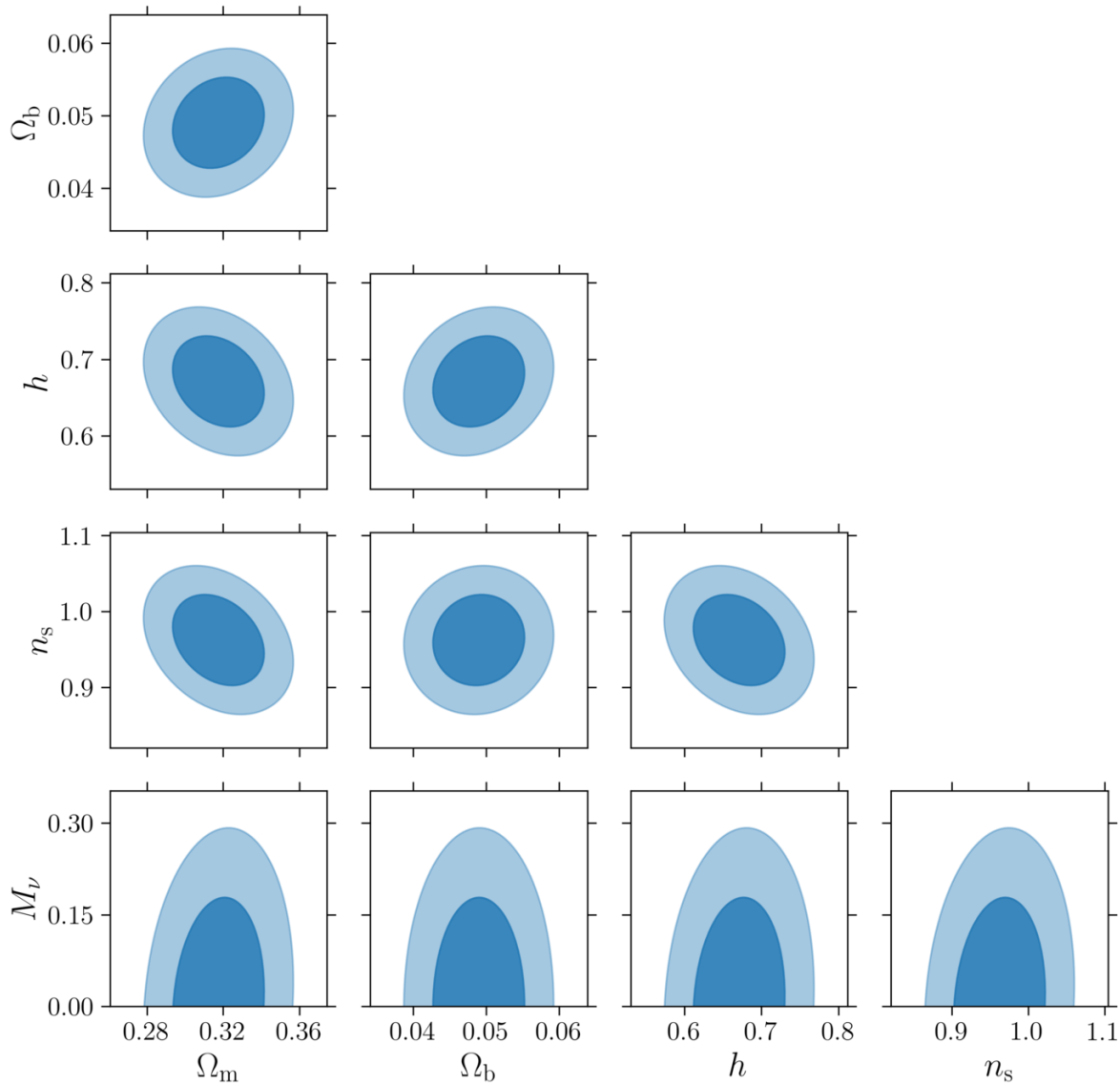
(Kuruvilla, 2021, to be submitted)

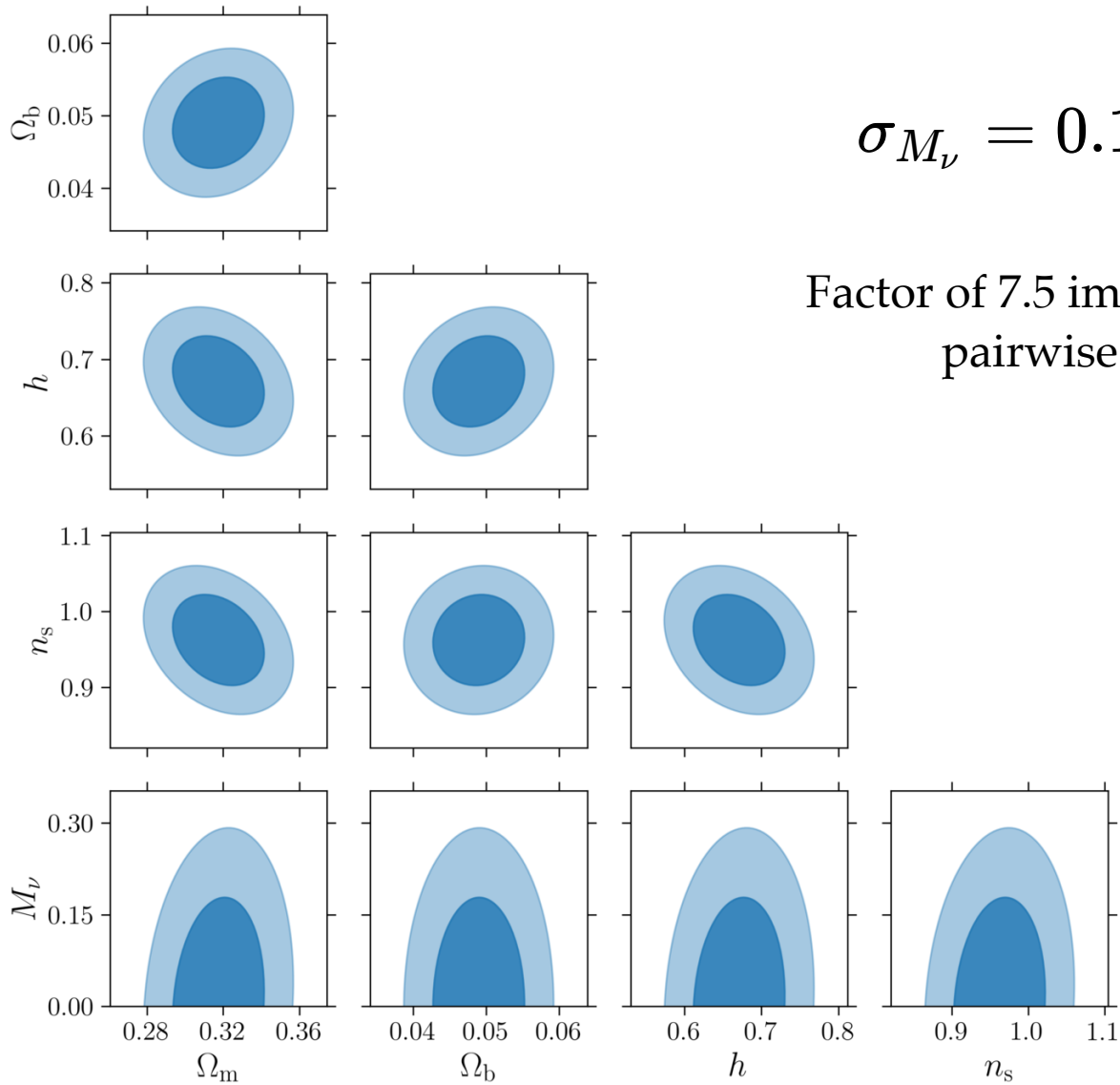




Effect of summed neutrino mass







$$\sigma_{M_\nu} = 0.1175 \text{ eV}$$

Factor of 7.5 improvement from
pairwise velocities



Conclusions

- Cosmological constraints from the mean three-point relative velocity statistics are competitive with those obtained from the bispectrum, while having sizeable improvements with respect to the power spectrum.

(Kuruvilla & Aghanim 2021, submitted to A&A)

- Introduced a new estimator which enables to constrain the summed neutrino mass independent of optical depth and σ_8 .

(Kuruvilla 2021, to be submitted)



Thank you for listening!



