

# **PHQMD: Flow and Clustering**

Michael Winn Supervisor : Prof. J. Aichelin

Subatech, CNRS, University of Nantes IMT Atlantique



#### PHQMD

PHQMD(Parton Hadron Quantum Molecular Dynamics)<sup>1</sup>is a simulation code that is based on PHSD(Parton Hadron String Dynamics)<sup>2</sup>(quasi particle description of QGP of PHSD)but uses N-body dynamics for propagation

#### **Quantum Molecular Dynamics**

Using n-body theory to track nucleon by nucleon the collision, and calculate the mutual interactions between all of the nucleons.

- IQMD<sup>3</sup>(Limited to lower energies 1.5 GeV/A)
- UrQMD<sup>4</sup>(Relativistic, used for coalescence, no potential)

<sup>&</sup>lt;sup>1</sup>J. Aichelin and, E. Bratkovskaya et al. *Phys. Rev. C 101, 044905* 

<sup>&</sup>lt;sup>2</sup>W. Cassing and E. BratkovskayaPhys.Rev.C 78,034919

<sup>&</sup>lt;sup>3</sup>C. Hartnack et al. *Eur. Phys. J.A1:151-169,1998* 

<sup>&</sup>lt;sup>4</sup>S. A. Bass et al.*Nucl.Phys.41:225-370* 

#### Why is it needed ?

Current models are not well adapted to address cluster formation, and current QMD models are not usable at relativistic energies

#### Challenges

Production of clusters with binding energy of  $B_E \approx 8[MeV]$  in environment of the fireball at  $E \approx 100[MeV]$ 

Hyper clusters production  $^{56}\ensuremath{\text{is}}$  an area of renewed interest, and which is not currently addressed

Transition form low energy hadron dominated interactions and high energy quark and gluon dominate reactions

Transition regime is characterised by a finite chemical potential thus a finite baryonic density

<sup>&</sup>lt;sup>5</sup>C. Rappold et al. P.L.B. Volume 747 Pages 129-134

<sup>&</sup>lt;sup>6</sup>J. Adam et al*Phys. Rev. C 93, 024917* 

#### Future experimental data

New energy range from 2 GeV/A - 100 GeV/A, used to investigate the first order phase transition from hadronic to QGP matter, and degrees of freedom of hadronic matter (strangeness) .

- FAIR
- NICA

#### Experimental results so far

- 0.6 [GeV/A] ALADIN <sup>7</sup>
- 1.23 [GeV/A] HADES <sup>8</sup>
- 1.5 [GeV/A] FOPI <sup>9</sup>

<sup>&</sup>lt;sup>7</sup>A. Schüttauf et al N.P.A Volume 607 Pages 457-486

<sup>&</sup>lt;sup>8</sup>J. Adamczewski-Musch et al *Nucl.Phys. A982* 

<sup>&</sup>lt;sup>9</sup>W. Reisdorf et al. *Nuclear Experiment 1112.3180* 

Generalised Ritz variational principle

$$\delta \int_{t_1}^{t_2} \mathrm{dt} \langle \Psi(t) | i \frac{\mathrm{d}}{\mathrm{d}t} - H | \Psi(t) \rangle = 0 \tag{1}$$

Gaussian test wave function which yields the Wigner density:

$$f(\vec{r_i}, \vec{p_i}, \vec{r_{i0}}, \vec{p_{i0}}, t) = \frac{1}{\pi^3 \hbar^3} \exp\left[-\frac{2}{L}(\vec{r_i} - \vec{r_{i0}})^2\right] \exp\left[-\frac{2}{L}(\vec{p_i} - \vec{p_{i0}})^2\right]$$
(2)

Classical type equations of motion for expectation value of the Hamiltonian

$$\dot{r} = \frac{\partial \langle H \rangle}{\partial p} \qquad \dot{p} = -\frac{\partial \langle H \rangle}{\partial r}$$
(3)

Hamiltonian composed of Kinetic term and Two body potentials

$$\langle H \rangle = \sum_{i}^{N} \langle H_{i} \rangle = \sum_{i}^{N} \left( \langle T_{i} \rangle + \sum_{i \neq j}^{N} \langle V_{i,j} \rangle \right)$$
(4)

Potentials

$$\langle V_{i,j} \rangle = \langle V_c \rangle \left( \vec{r} \right) + \langle V_s \rangle \left( \rho \right) + \left( \langle V_{SM} \rangle \left( \Delta \vec{p}^2 \right) \right)$$
(5)

Interaction density

$$\rho_{int}(i,t) = C \sum_{j,j\neq i}^{N} \exp\left[\frac{1}{L}(\vec{r_i} - \vec{r_j})^2\right]$$
(6)



Equations of state parametrised for soft and hard nuclear matter compressibility



<sup>1</sup>Nuclear Physics A640

Optical potential fit to expermimental  $data^{10}$ 



Initial Wood-Saxon position distribution for the three E.o.S.





Conservation of the total system energyEvolution of the average radial nuclei distance, and the average density of the nuclei

### Density profile mid-collision



Density profile of the three equations of state for a collision at E = 0.6 [AGeV] for semi-peripheral impact b = 7 [fm] at T = 10 [fm/c]

#### SACA<sup>11</sup> and MST<sup>12</sup>

To form the cluster we use two methods

- MST : simple form of spanning tree to create clusters based on distance from other nucleons
- SACA : complex simulated annealing of all possible cluster patterns of the nucleons to find the lowest sum of binding energies of all clusters

All the following results are for the SACA method

 <sup>&</sup>lt;sup>10</sup>R. K. Puri, C. Hartnack and J. Aichelin*Phys. Rev. C 54, R28(R)* <sup>11</sup> J. Aichelin, Phys. Rept.202, 233

### FOPI results for E = 1.5 [GeV/A] Multiplicity



Fragment multiplicity

$$\label{eq:zero} \begin{split} Z{=}1 \text{ isotopes multiplicities for central} \\ \text{collisions ( } b < 2 \text{ [fm])} \end{split}$$

 $Z_{bound2}$  the sum of the charges all fragments with  $Z \ge 2$ , and  $Z_{MAX}$  the charge of the largest cluster



Rise and Fall curve for all three equations of state

Zmax curve for all three equations of state

## FOPI results for E = 1.5 [GeV/A] Flow



Direct flow for proton and deuteron with a threshold on transverse momentum

Direct proton flow for 4 different impact parameter classes for all protons with a threshold on transverse momentum

### HADES results for E = 1.23 [GeV/A] Flow



Direct proton flow as a function of rapidity for three  $p_T$  classes compared with the HADES data

Elliptical proton flow as a function of rapidity for three  $p_T$  classes compared with the HADES data

- PHQMD quite nicely replicate heavy fragments spectra and reasonable agreement for Z=1 fragments
- Flow reproduction with similar behaviour for Hard E.o.S. and Soft momentum E.o.S. reinforcing the need for momentum dependence
- Some issues that are being fixed for small fragments at Mid Rapidity