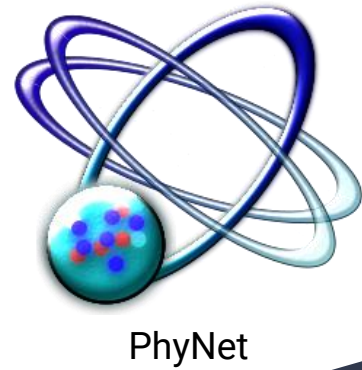
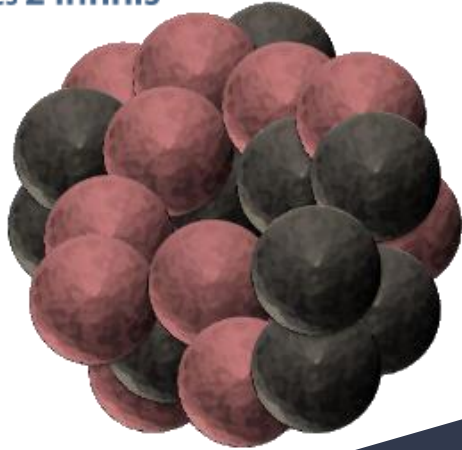


Nuclear structure and α radioactivity



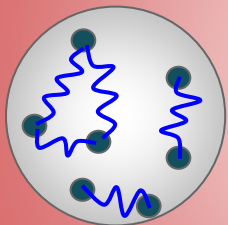
Florian MERCIER



- I. The nuclear many body problem
- II. Building an interaction
- III. Many body interacting problem
- IV. Alpha and cluster radioactivity

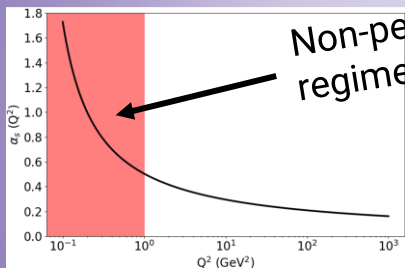
Tackling the nuclear many body problem

Quantum many body interacting problem !



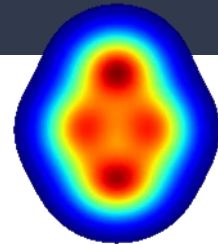
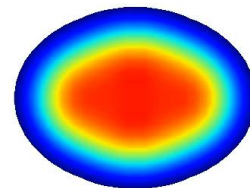
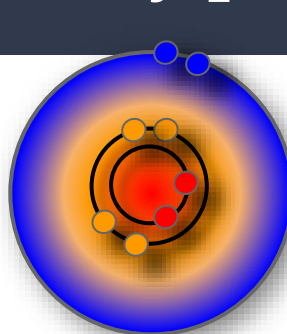
Many (strongly) interacting particles but not enough for statistical approximation
TOOOOO hard to be solved exactly

Interaction coming from (non-perturbative) QCD !



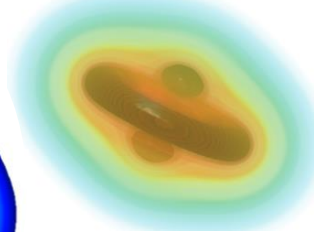
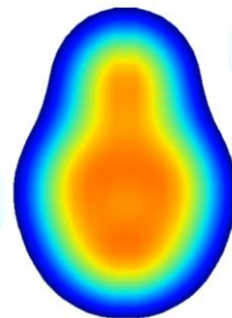
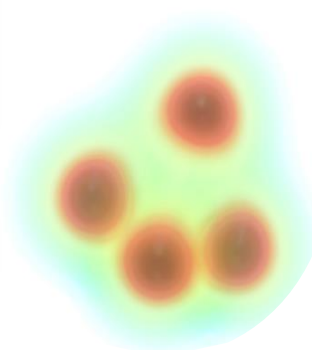
Non-perturbative regime at low energy !

EFT
Relevant degrees of freedom ?
 $2N, 3N, 4N, \dots$ interactions



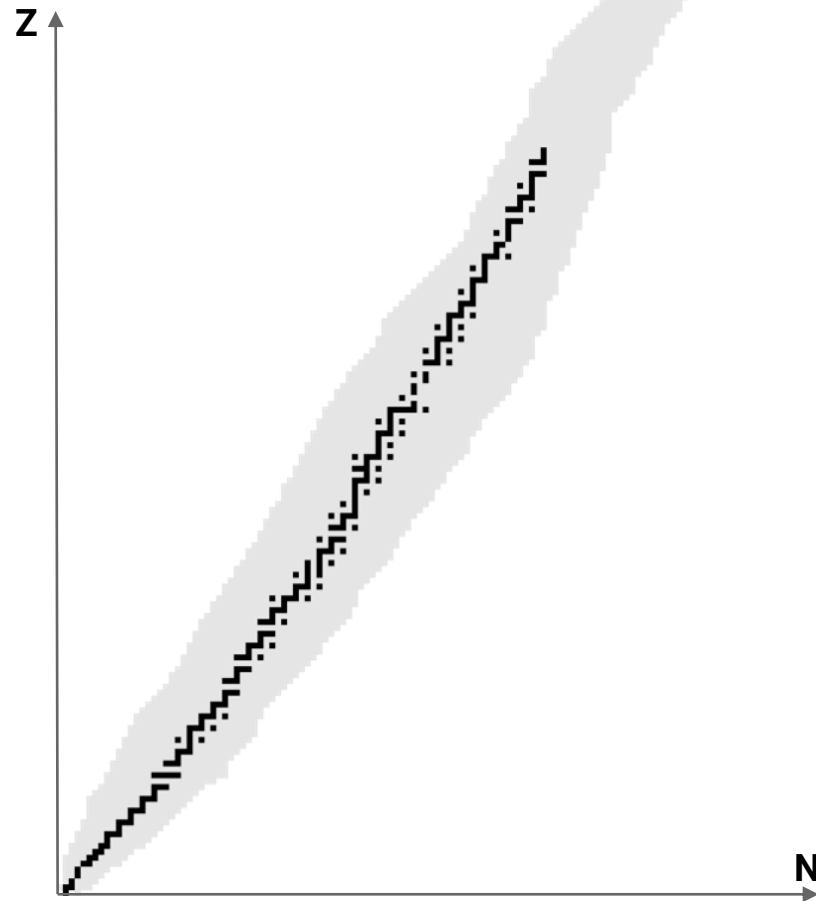
Huge phenomenology !

Many emerging properties : deformation, radioactivity, halo, neutron skin, superfluidity, clustering, excitations, ...



The different methods to deal with nuclear system

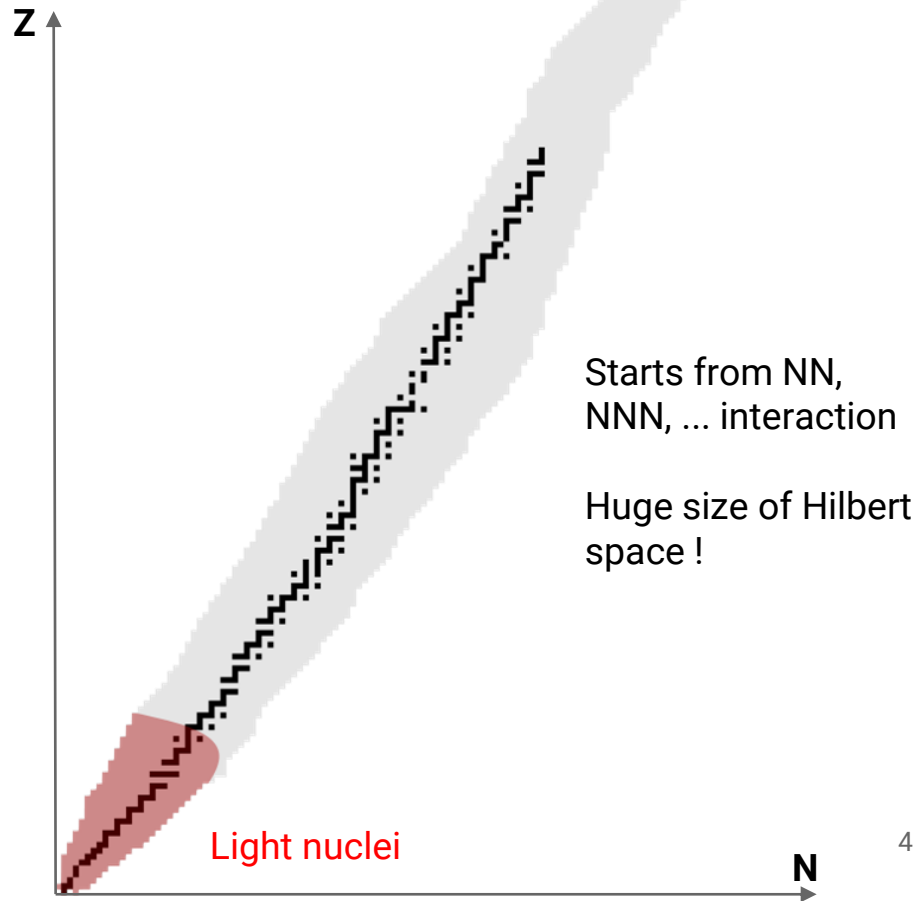
Nuclear landscape



The different methods to deal with nuclear system

- Ab Initio

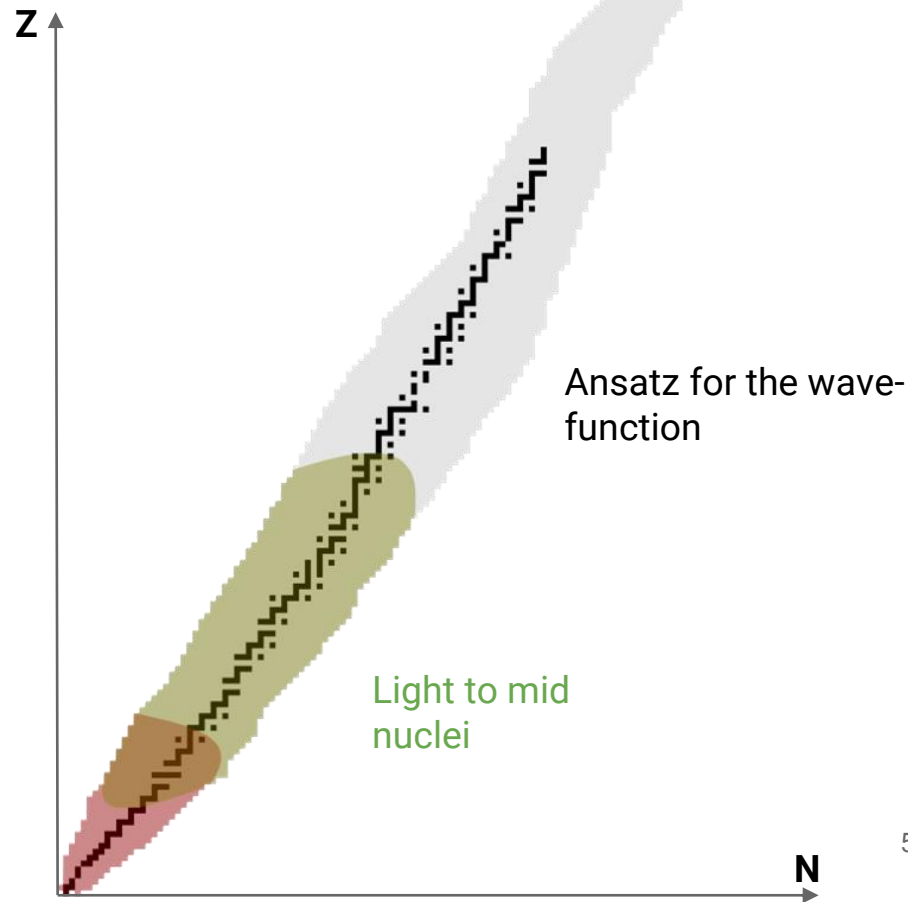
Nuclear landscape



The different methods to deal with nuclear system

- Ab Initio
- Configuration interaction

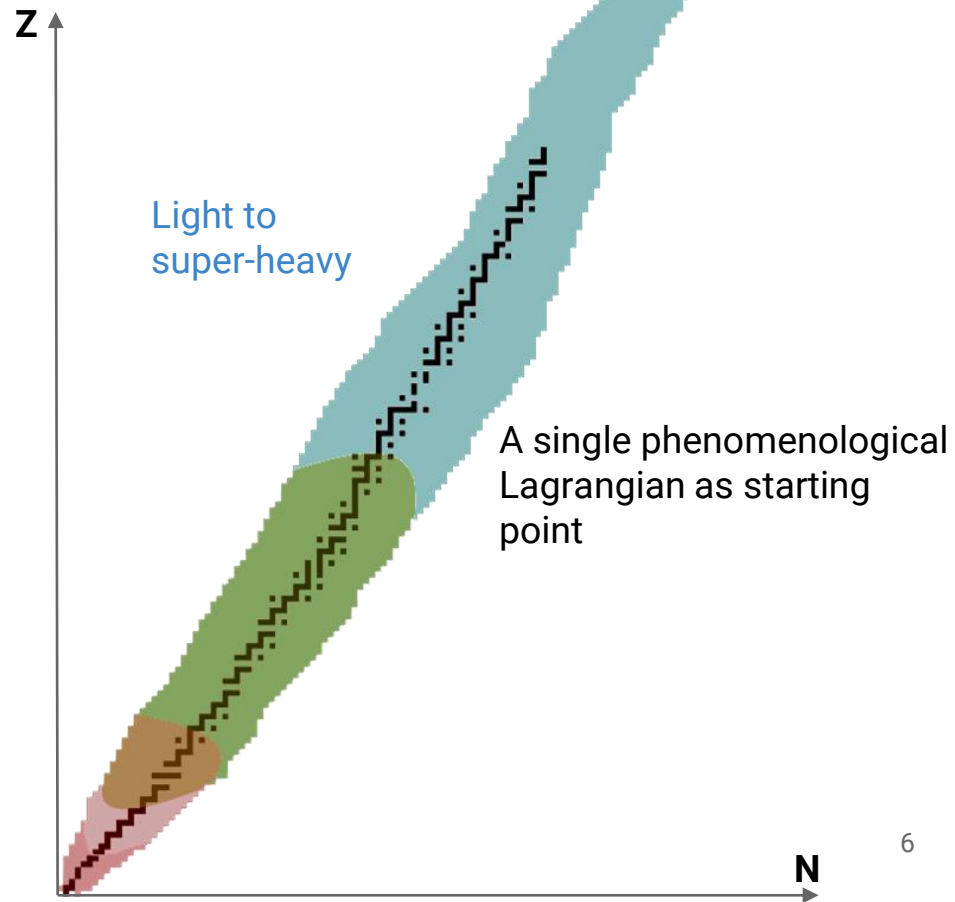
Nuclear landscape



The different methods to deal with nuclear system

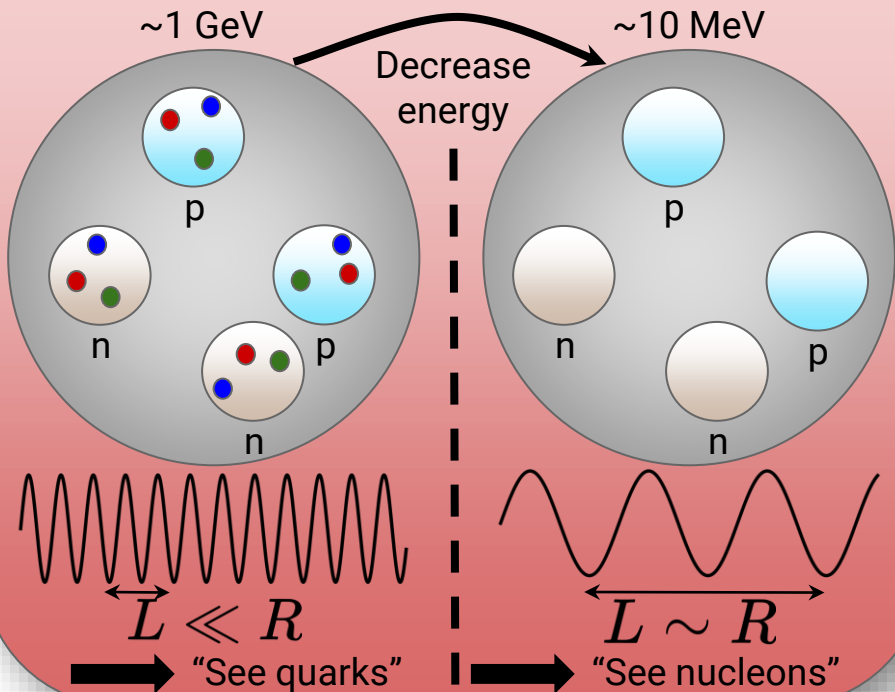
- Ab Initio
- Configuration interaction
- Energy density functional

Nuclear landscape

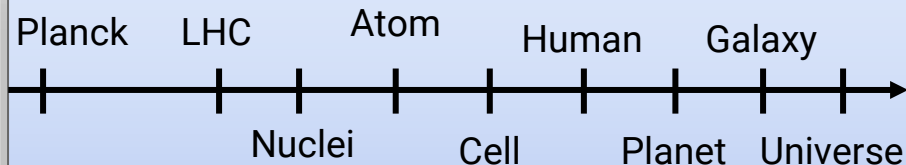


Building an interaction : what is the idea ?

Relevant degrees of freedom ?



Effective field theory

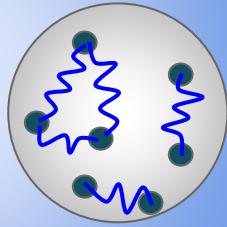



- ➡ Scale separation is the key ingredient
- ➡ Zooming out (decreasing energy, increasing scale) leads to different theories !
- ➡ Allows a scale by scale study in Science !

Many body interacting problem

Strongly **interacting** many body system

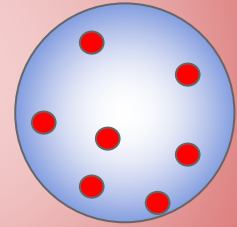
VERY HARD



Mapping

Fermi liquid

One-body problem involving **independent** quasi-particles

Mean-field



Idea :

$$H = \left[\sum_i \frac{p_i^2}{2m} \right] + \left[\sum_{ij} V_{ij} + \sum_{i<j<k} V_{ijk} + \dots \right]$$

$$= \left[\sum_i \frac{p_i^2}{2m} + U_i \right] + \left[\sum_{ij} V_{ij} + \sum_{i<j<k} V_{ijk} + \dots - \sum_i U_i \right]$$

$$= H_0 + \mathcal{V}_{res}$$

Such that $|H_0| \gg |\mathcal{V}_{res}|$

One body term

Find the “best” U_i

As much physics
as possible

Justify perturbative
approach

Hartree-Fock theory

Symmetries

$$H = H_0 + \mathcal{V}_{res} \text{ with a certain } U_i$$

We know that the Hamiltonian H possesses some symmetries

$$[H, N] = 0 \quad \text{Particle number conservation}$$

$$[H, J] = 0 \quad \text{Invariance under rotation}$$

Let us choose the same symmetries for H_0

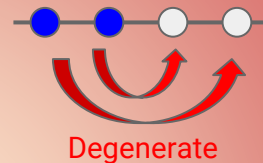
$$[H_0, N] = 0 \quad [H_0, J] = 0$$

Problem

Perturbation theory



Improvement by going at higher order



$$E = \mathcal{E} + \sum_{n>0} \underbrace{\frac{|\langle \phi_n | \mathcal{V}_{res} | \phi_0 \rangle|^2}{\mathcal{E} - \mathcal{E}_n}}_{\gg 1} + \dots$$

For (quasi-) degenerate state

We are missing an essential ingredient in \mathcal{V}_{res} !
Change U_i to change the physical content

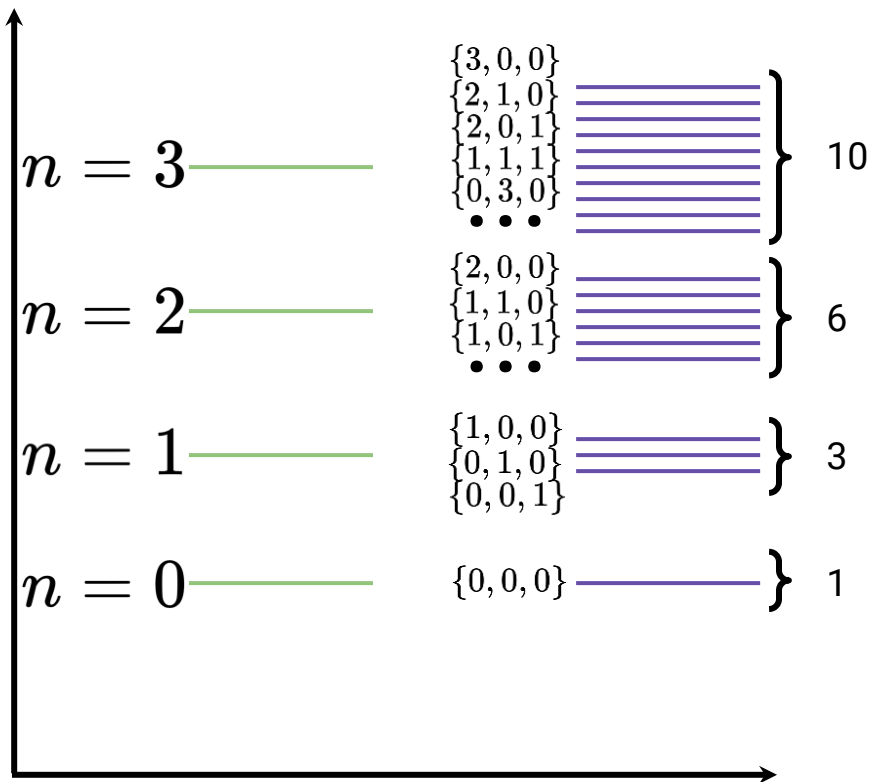
How to raise degeneracies ?

The isotropic (symmetric) harmonic oscillator

$$E_n^{sym} = \hbar\omega(n_x + n_y + n_z + 3/2) \\ = \hbar\omega(n + 3/2)$$

The anisotropic (asymmetric) harmonic oscillator

$$E_{n_x n_y n_z}^{asym} = \hbar(\omega_x n_x + 1/2) + \hbar(\omega_y n_y + 1/2) \\ + \hbar(\omega_z n_z + 1/2)$$



Hartree-Fock-Bogoliubov theory

Symmetries

$$H = H_0 + \mathcal{V}_{res} \text{ with a certain } U_i$$

We know that the Hamiltonian H possesses some symmetries

$$[H, N] = 0 \quad \text{Particle number conservation}$$

$$[H, J] = 0 \quad \text{Invariance under rotation}$$

↓ Let us choose the same symmetries for H_0

$$[H_0, N] = 0 \quad [H_0, J] = 0$$

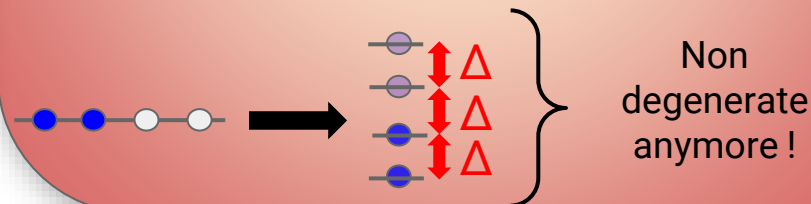
Problem Solution

Assume $[H'_0, N] \neq 0 \longrightarrow$ Particle number not conserved anymore !

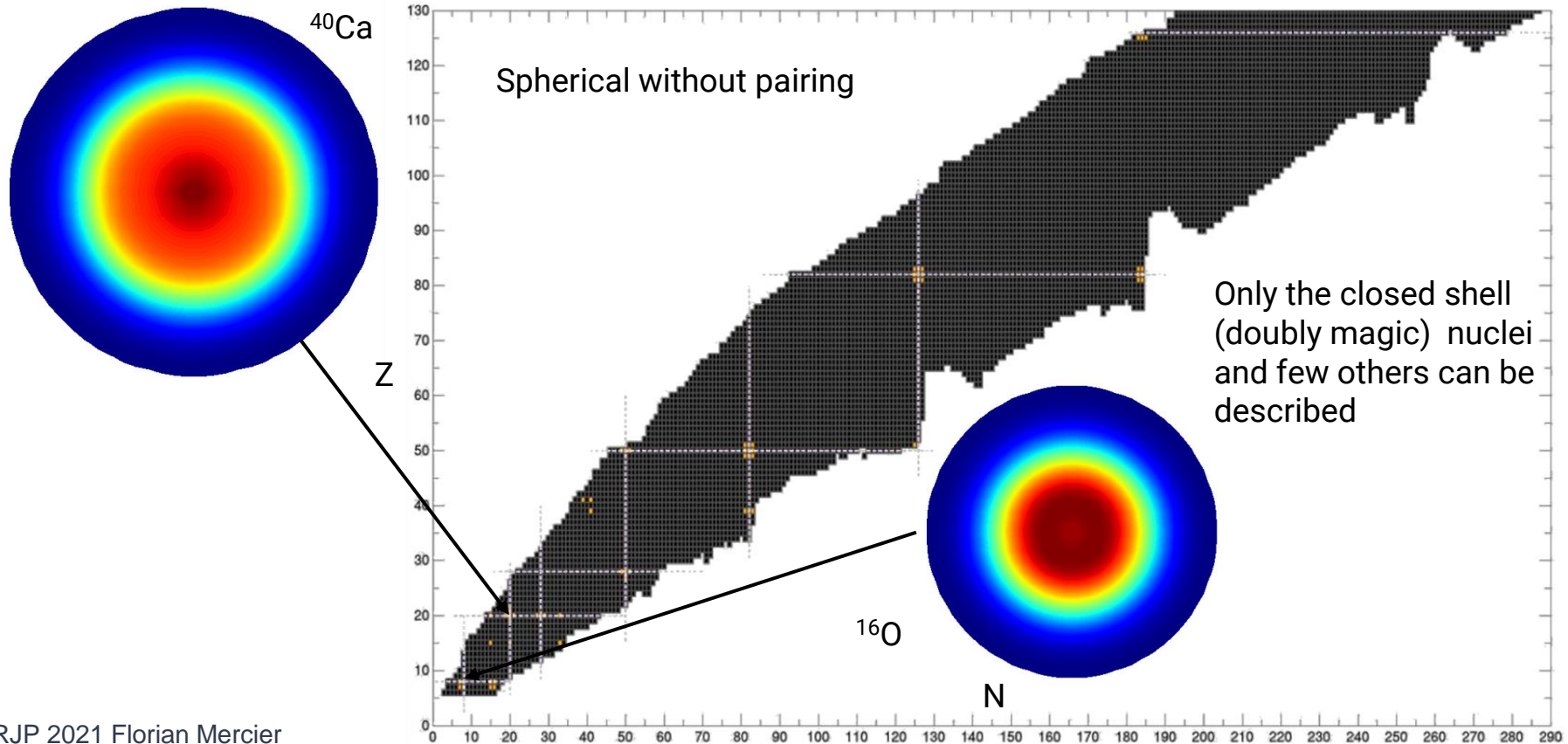
Particles can create pairs which behave as single degree of freedom

As for spin pairing, it leads to a splitting of the energy and a new minimum energy configuration !

$$E = \epsilon \pm \Delta$$



Impact of symmetry breaking : no breaking



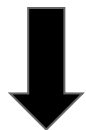
Impact of symmetry breaking : U(1) breaking

What are we missing then ?

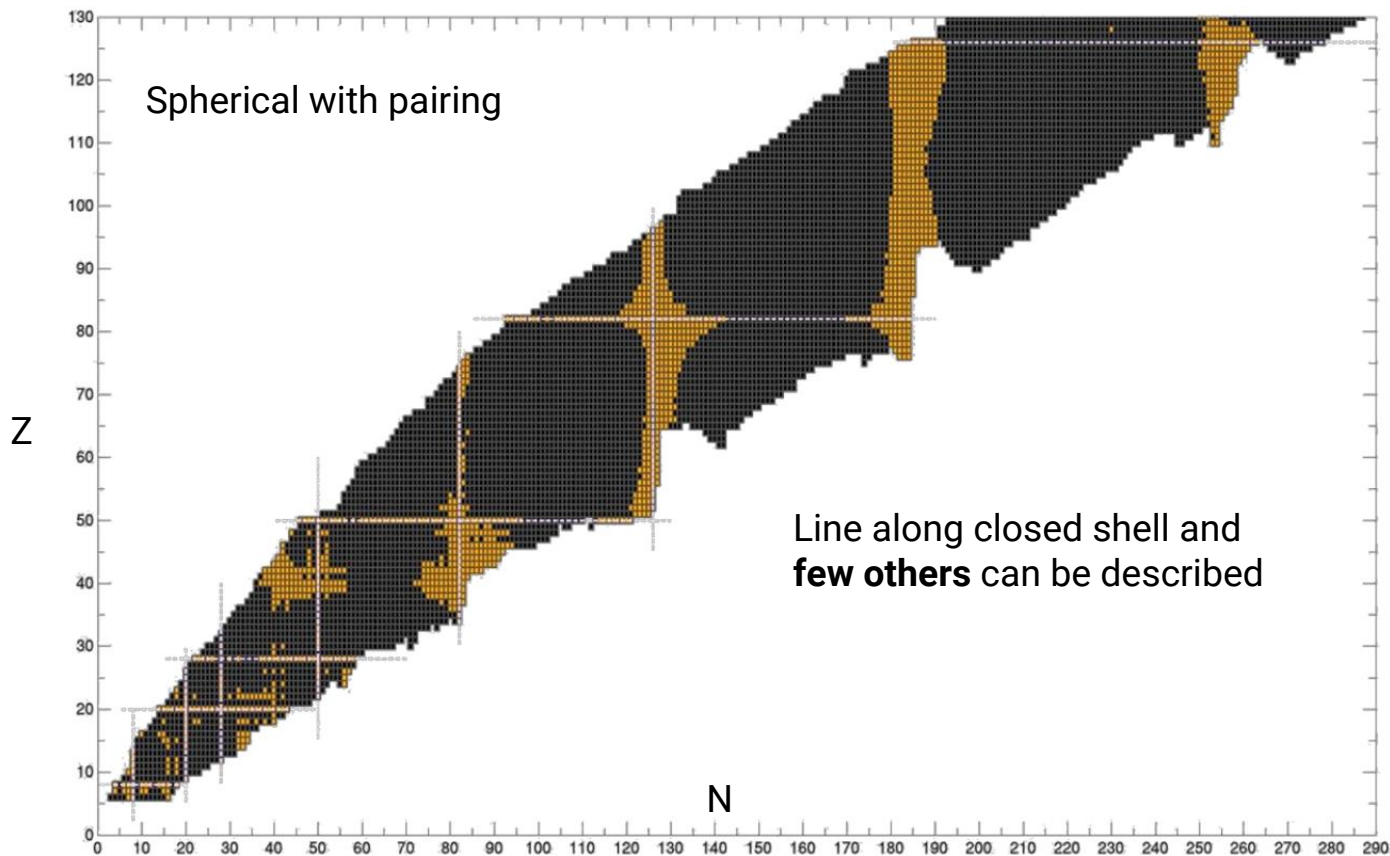


Let us break other symmetries !

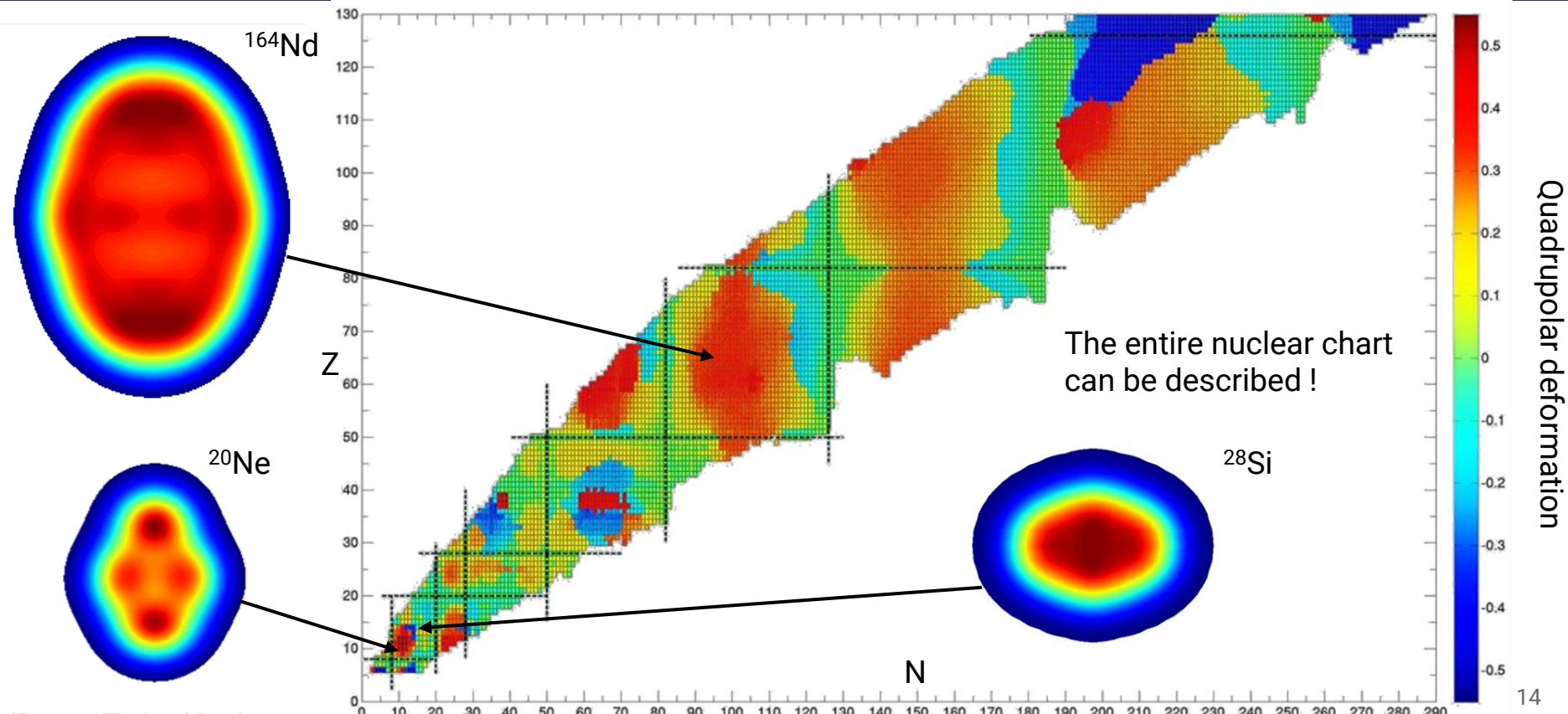
$$[H_0, J] \neq 0$$



Nuclei can now be deformed



Impact of symmetry breaking : U(1) and SO(3)



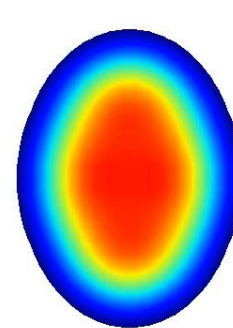
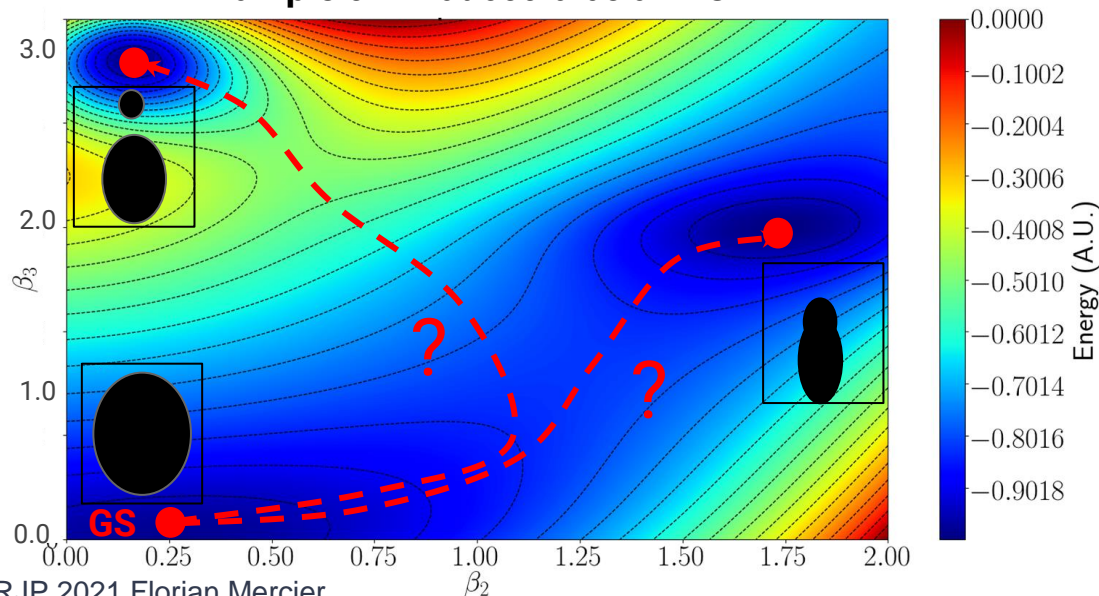
Deformations and Potential Energy Surface

Impose the deformation in
quadrupolar/octupolar
shapes and compute E

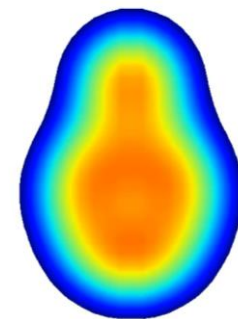


It will give $E(\beta_2, \beta_3) \geq E_{\min}$

Example of what could be a PES



Quadrupolar β_2

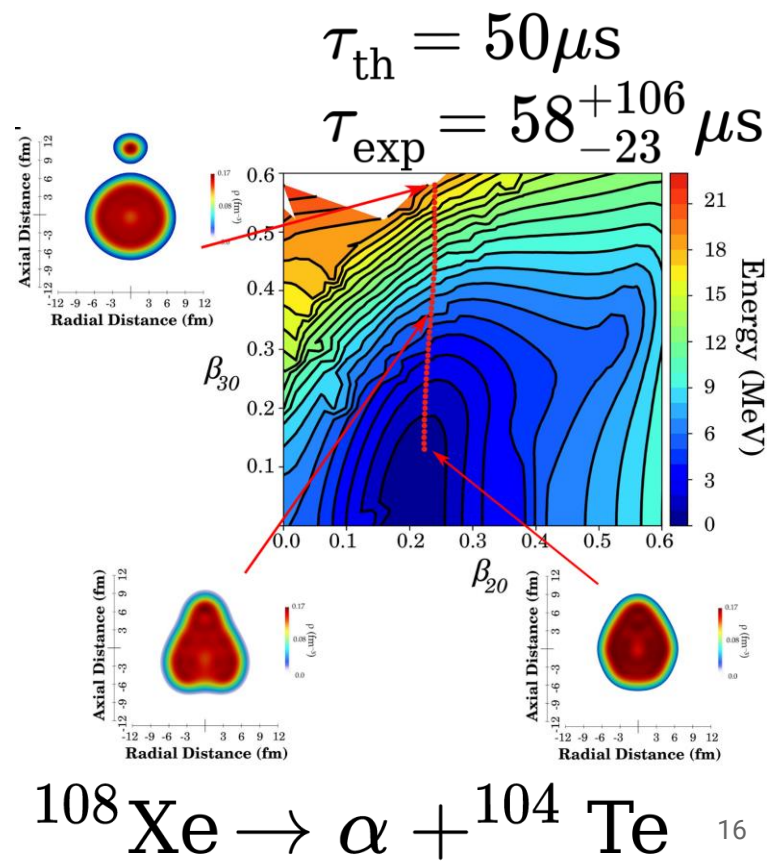
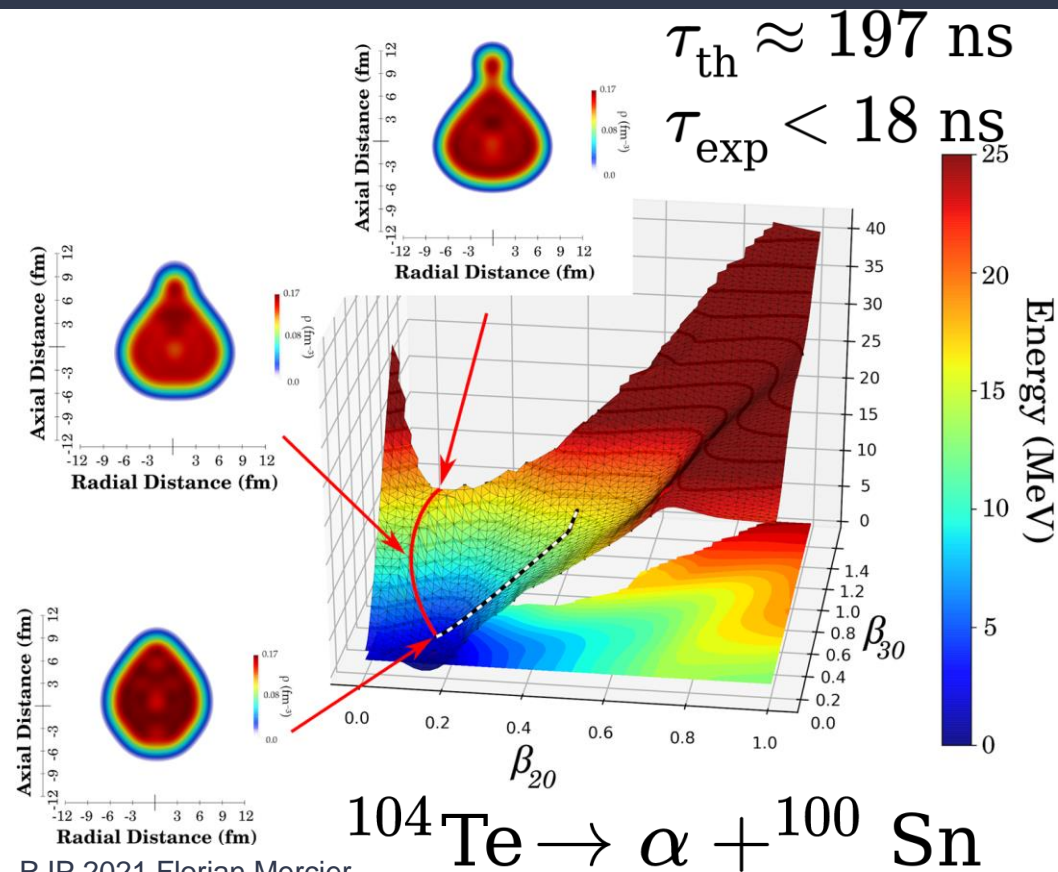


Octupolar β_3

- Can shapes evolve from GS to α emission ?
- Is it possible to compute the associated probability ?

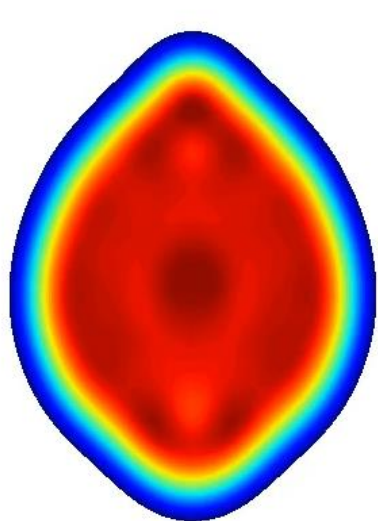
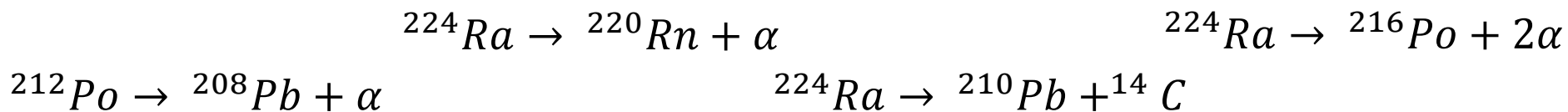
Alpha decay chain $^{108}\text{Xe} \rightarrow ^{104}\text{Te} \rightarrow ^{100}\text{Sn}$

F. Mercier and al., Phys. Rev. C 102, 011301(R) (2020)

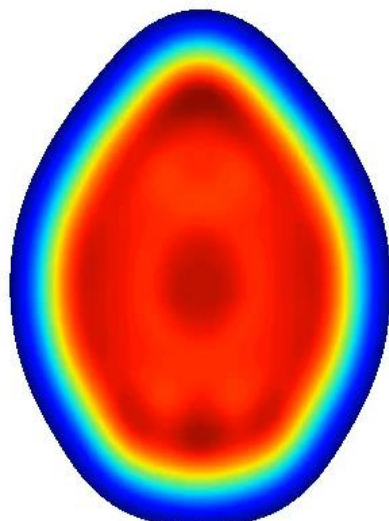


Other results for heavier nuclei

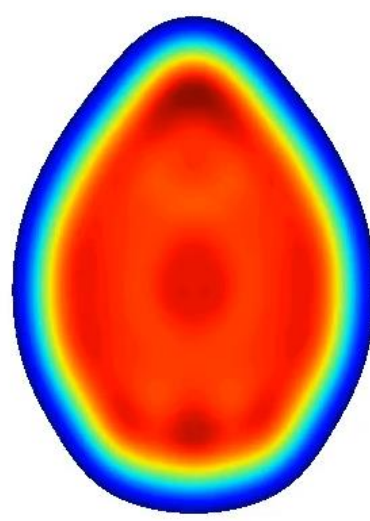
F. Mercier and al., to be published



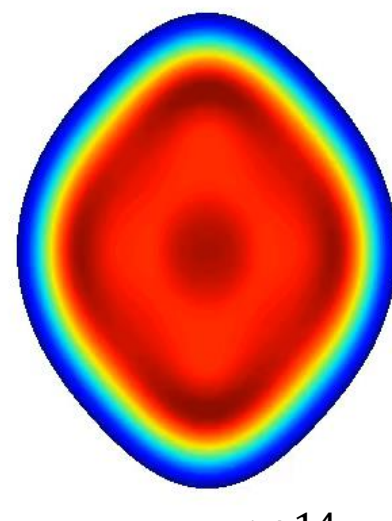
$$\tau_{th} \approx 6 \times 10^{-7} \text{ s}$$
$$\tau_{exp} \approx 3 \times 10^{-7} \text{ s}$$



$$\tau_{th} \approx 8 \times 10^5 \text{ s}$$
$$\tau_{exp} \approx 3 \times 10^5 \text{ s}$$



$$\tau_{th} \approx 3 \times 10^{17} \text{ s}$$
$$\tau_{exp} \approx 7 \times 10^{15} \text{ s}$$



$$\tau_{th} \approx 10^{14} \text{ s}$$
$$\tau_{exp} = ??$$