

$t\bar{t}$ Spin Correlations in the di-leptons channel using Run 2 data & EFT interpretation

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Date : 18/03/2021



Objective

1. Spin Observables
2. Effective field theory interpretation

Why Top quark is so special ?

Top is an ideal quark for spin measurements :

- ▶ decays before it can form bound states
- ▶ spin information transferred to daughter particles
- ▶ expect top spin observables to be well predicted by perturbative QCD

$$\text{lifetime} < \frac{\text{QCD timescale}}{\text{spin-flip timescale}} \ll 10^{-25} \text{ s} < 10^{-24} \text{ s} \ll 10^{-21} \text{ s}$$

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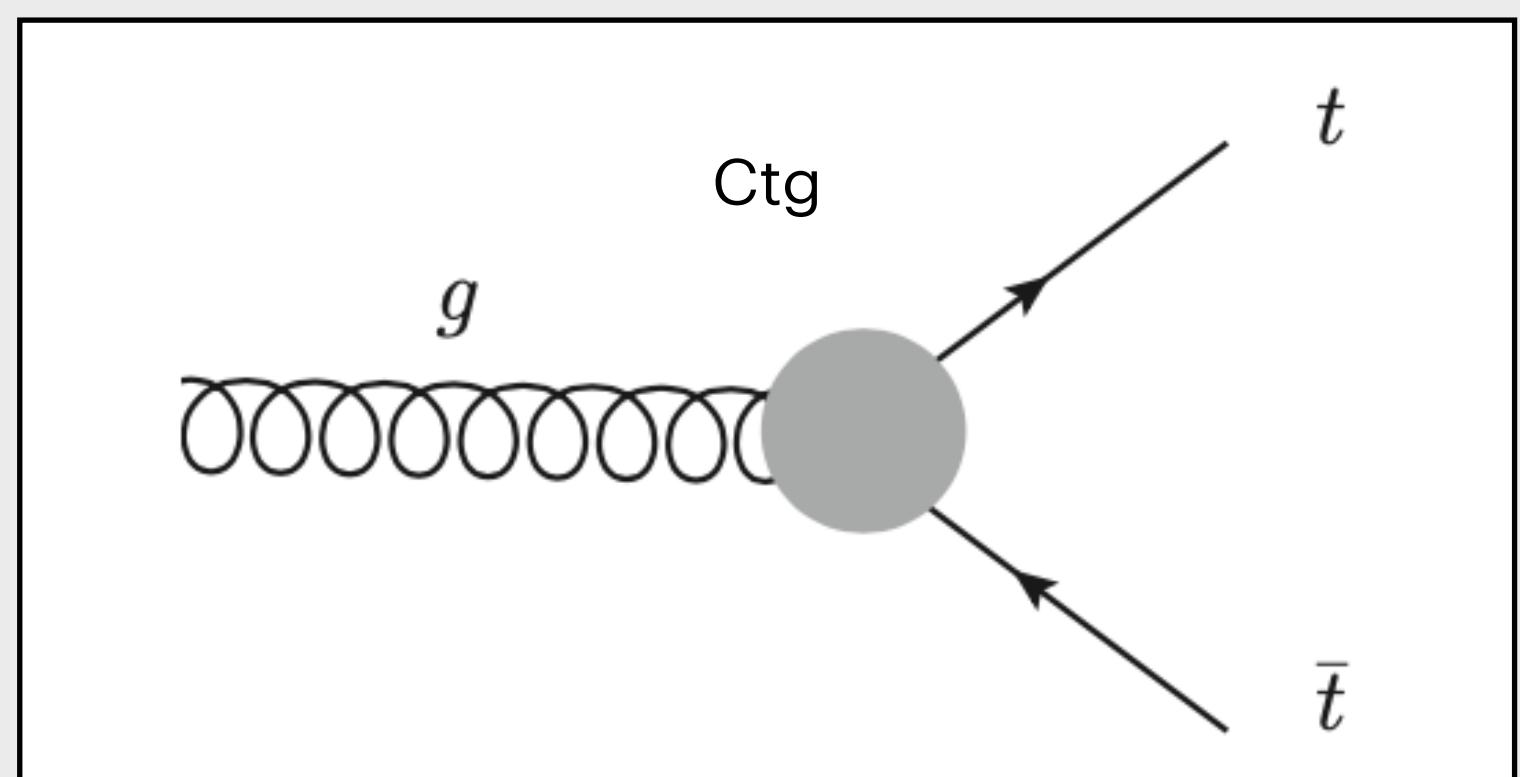
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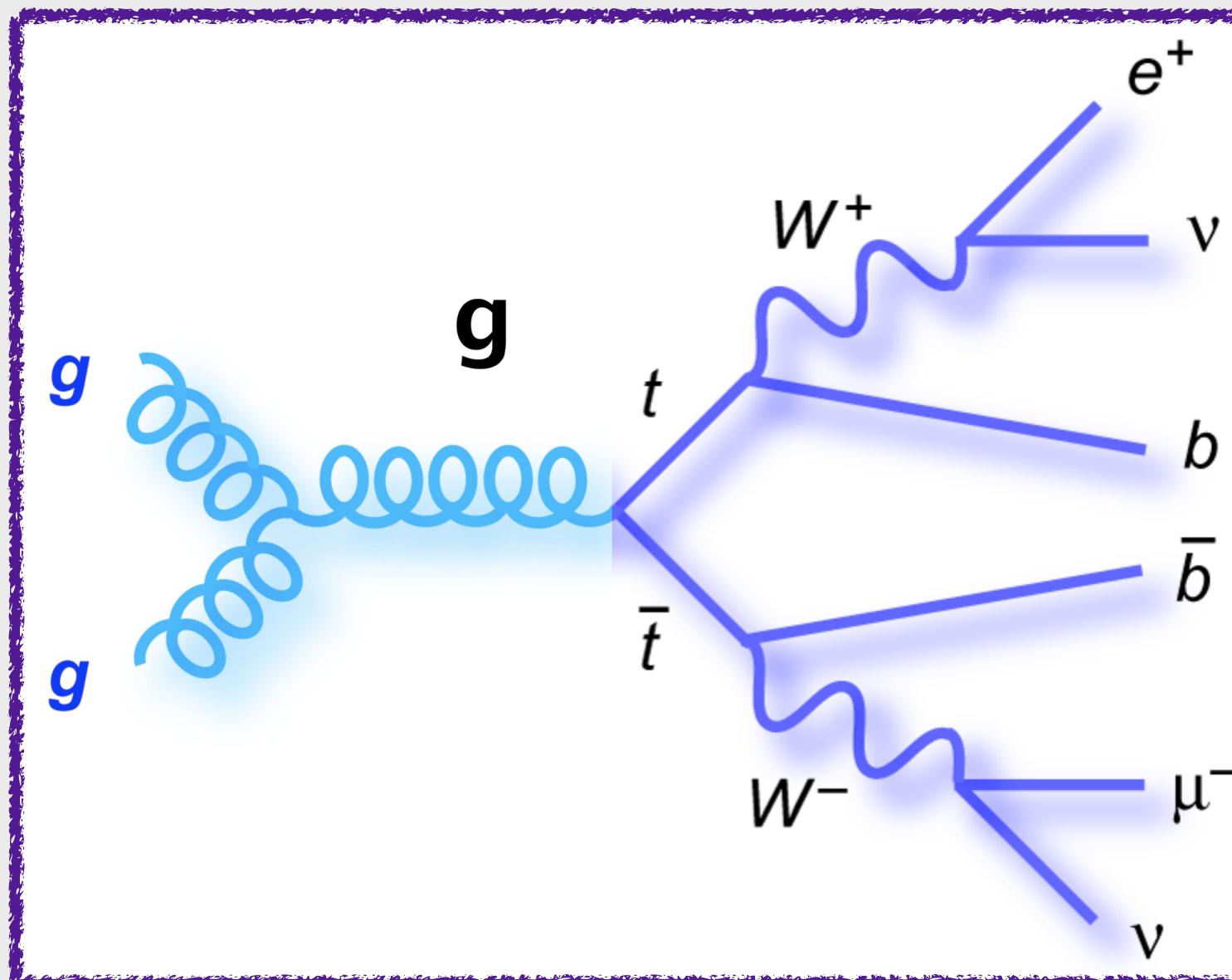
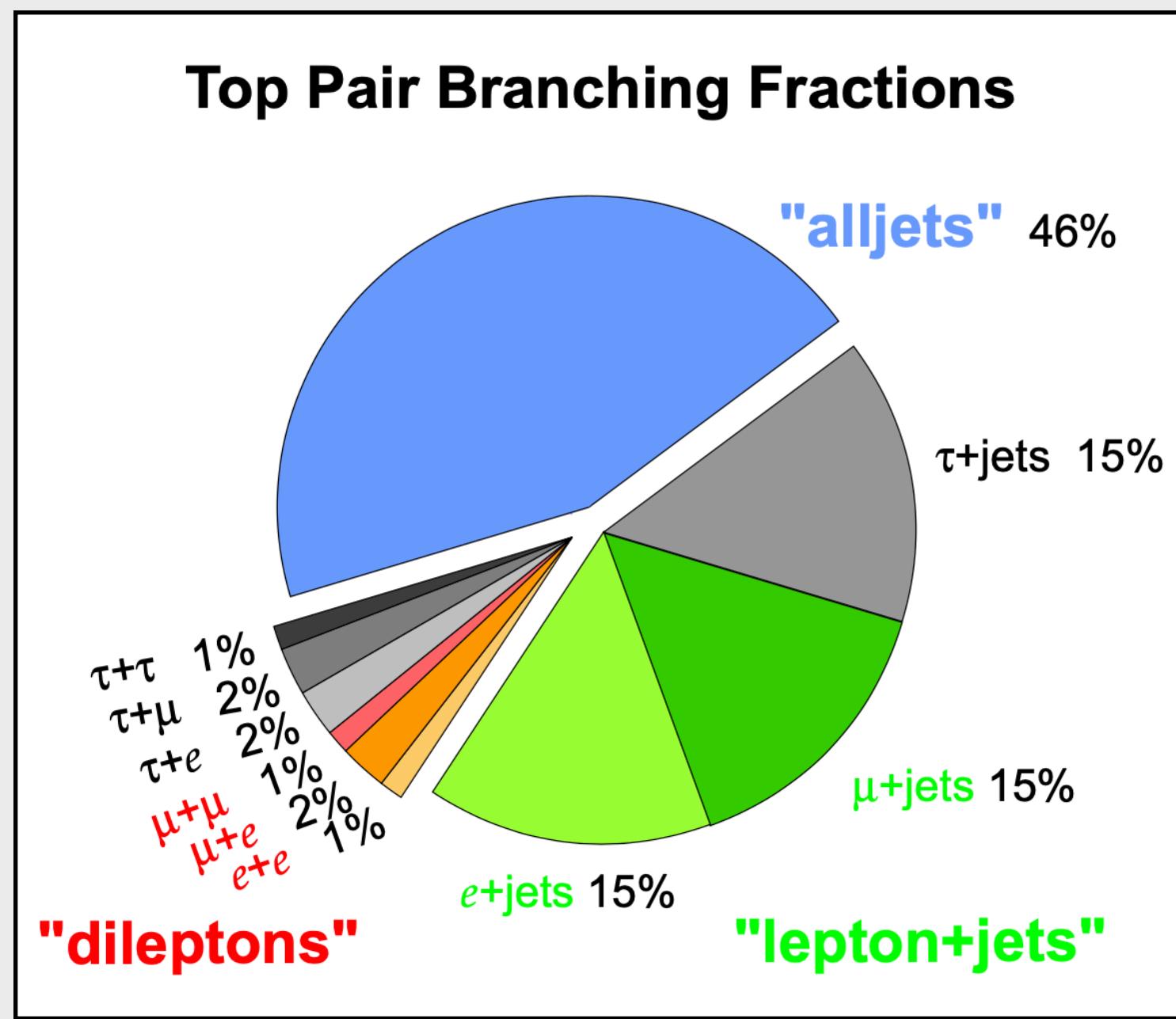
Top spin measurements are a powerful probe of new physics in $t\bar{t}$ production :

- ▶ new mediator would change spin structure
- ▶ sensitive to many dim-6 EFT operators

$$\text{lifetime} < \frac{\text{QCD timescale}}{\text{spin-flip timescale}} \ll 10^{-25} \text{ s} < 10^{-24} \text{ s} \ll 10^{-21} \text{ s}$$



Top quarks pair decay mode



All - hadronic

- ❖ Largest BR
- ❖ Largest QCD background
- ❖ Event fully constrained

Semi - leptonic

- ❖ High BR
- ❖ Medium background
- ❖ Event constrained

Di-leptonic

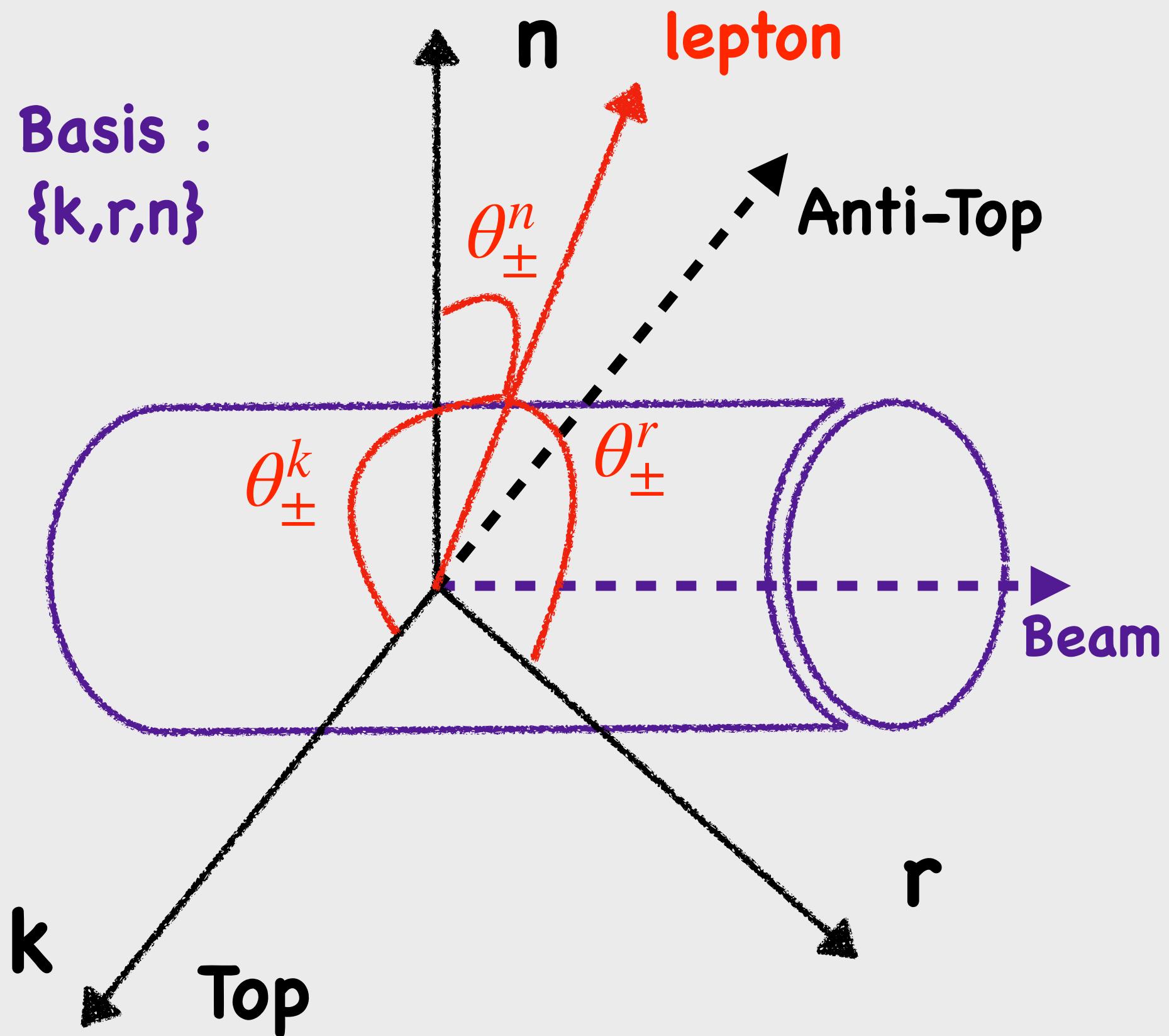
- ❖ Small BR
- ❖ S/B Good
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Spin Observables

Probing the spin observables :

- ✓ The dominant effect of the spin correlations is to correlate the angles of the decay products between the top quark and anti-top quark :

$$\frac{1}{\sigma} \frac{d^2\sigma}{d \cos(\theta_+^a) d \cos(\theta_-^b)} = \frac{1}{4} (1 + B_+^a \cos(\theta_+^a) + B_-^b \cos(\theta_-^b) - C(a, b) \cos(\theta_+^a) \cos(\theta_-^b))$$



The subscript +(-) refers to the top (anti-top) quark

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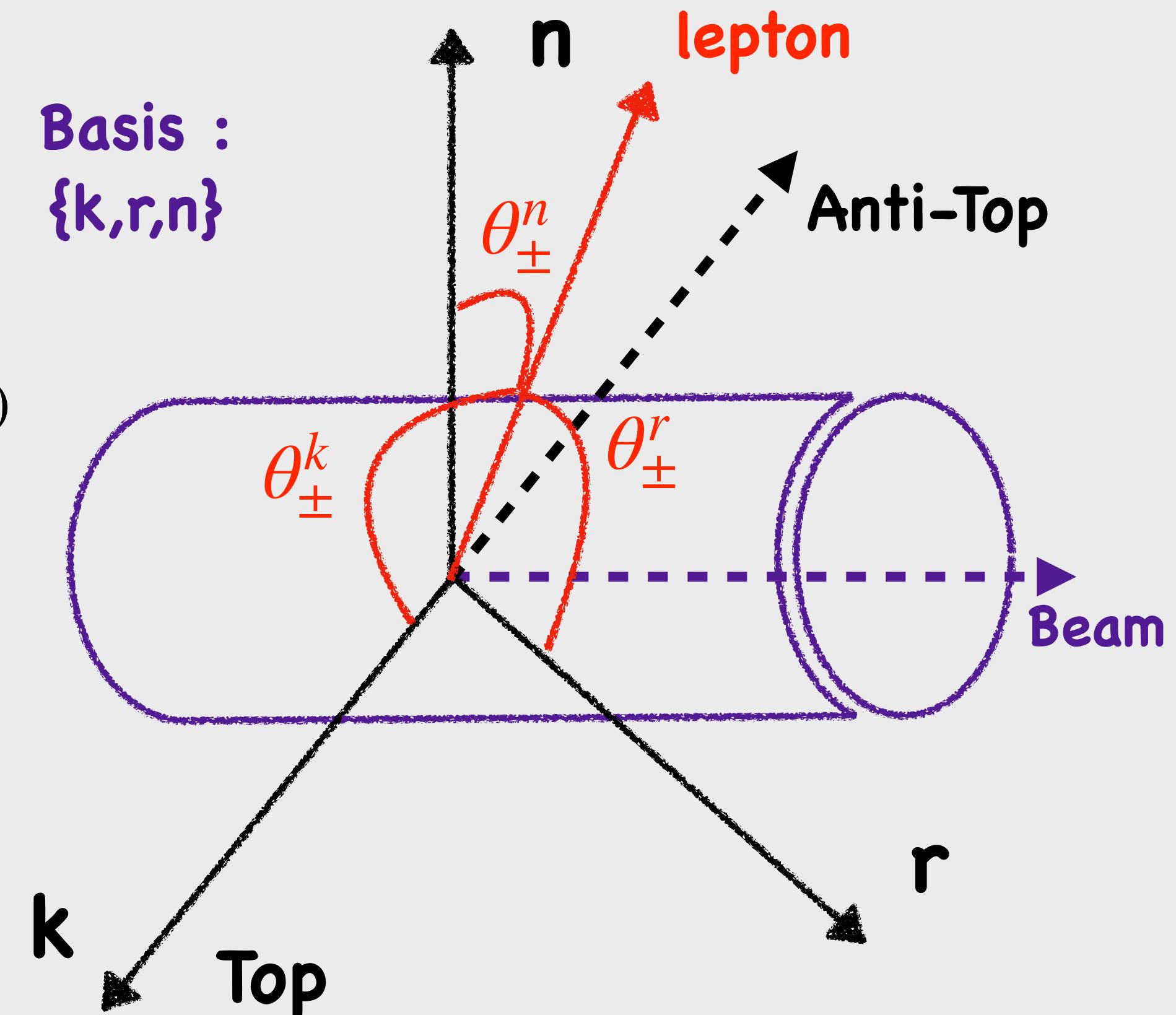
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- B_+^a, B_-^b and $C(a, b)$ are the polarisation and spin correction in quantisation axis a and b where $a, b = (\hat{k}, \hat{n}, \hat{r})$.

$C(a, b) = -9 < \cos(\theta_a^+) \cos(\theta_b^-) > \Rightarrow 9$ correlations

$B^a = 3 < \cos(\theta^a) > \Rightarrow 6$ polarisations



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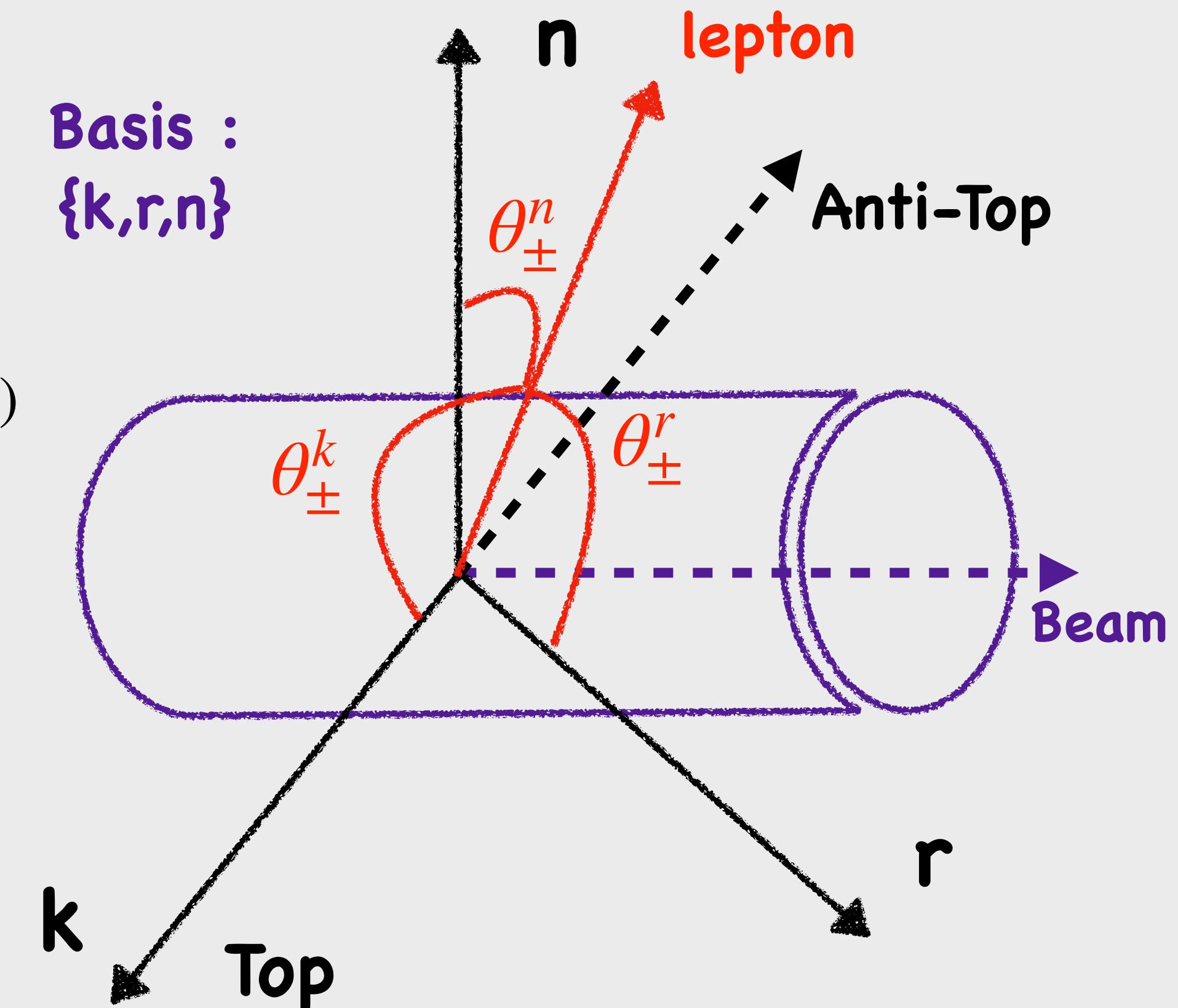
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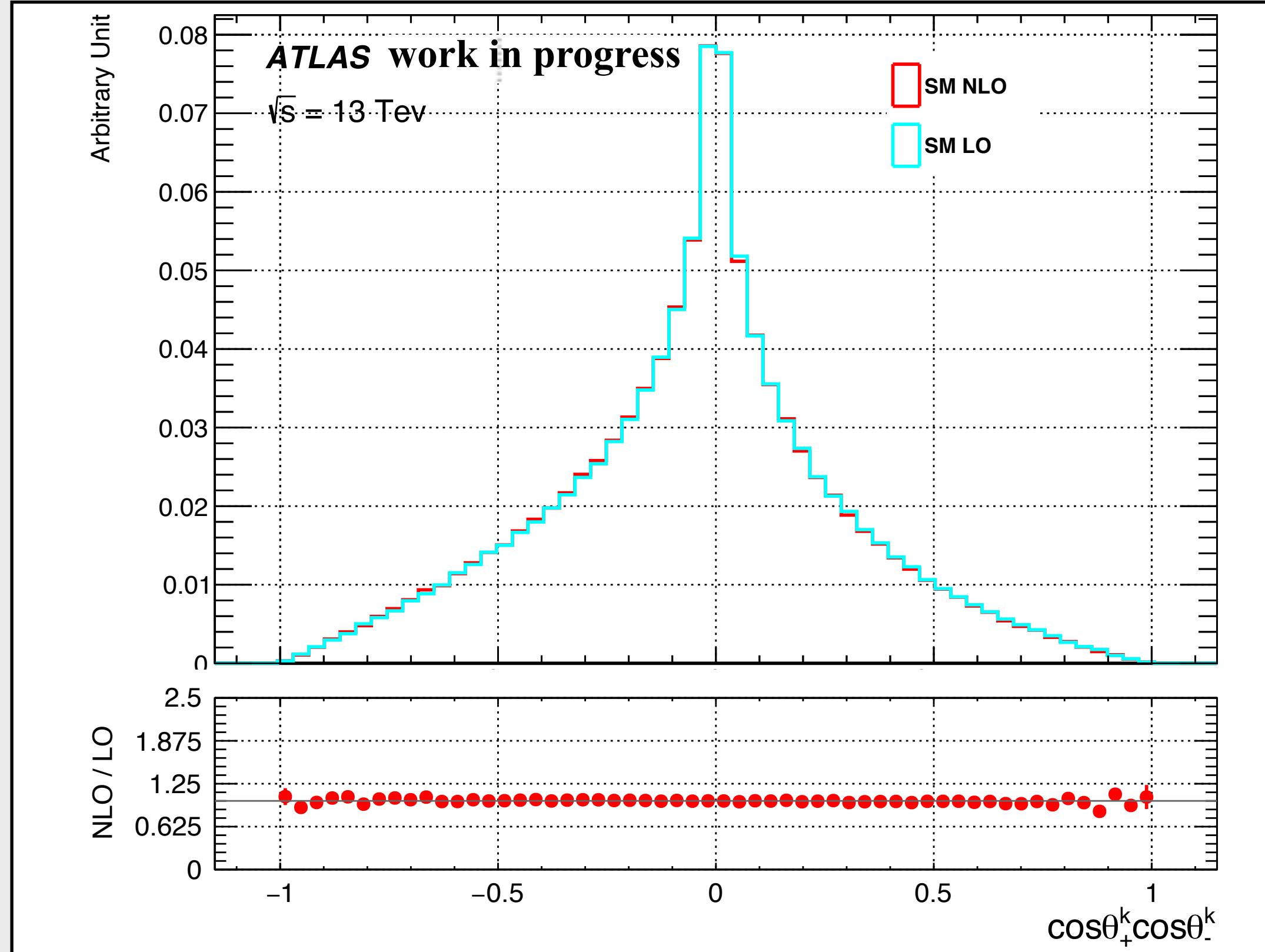
- ✓ these 15 coefficients completely characterise spin dependence of $t\bar{t}$ production and can be measured experimentally.



The subscript $+(-)$ refers to the top (anti-top) quark

Spin Correlations : $C(k,k)$

- ✓ Distributions for the correlation of top spins along k axis (probing diagonal of C (a, b) matrix) : $C(k,k) = -9 <\cos(\theta_k^+) \cos(\theta_k^-)>$



SM NLO : $C(k, k) = 0.366313 +/- 0.0042$
(stat)

SM LO : $C(k, k) = 0.341856 +/- 0.0042$
(stat)

- ✓ Spin correlations along each axis consistent with SM expectations (NLO from [1508.05271](#))

Interpretation

Introduction

- In effective field theory (EFT) language, the Standard Model Lagrangian is the first term in an effective Lagrangian

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_{D>4} \sum_i \frac{c_i^{(D)}}{\Lambda^{D-4}} \mathcal{O}_i^{(D)}$$

- Where Λ generically represents the scale of the new physics. c_i are dimensionless Wilson coefficients.
- EFT implemented in [dim6top model](#) and [SMEFT@NLO model](#)
- These model are implemented inside [MadGraph5_aMC@NLO framework](#) which translate a Lagrangian into a MC sample.

EFT interpretation

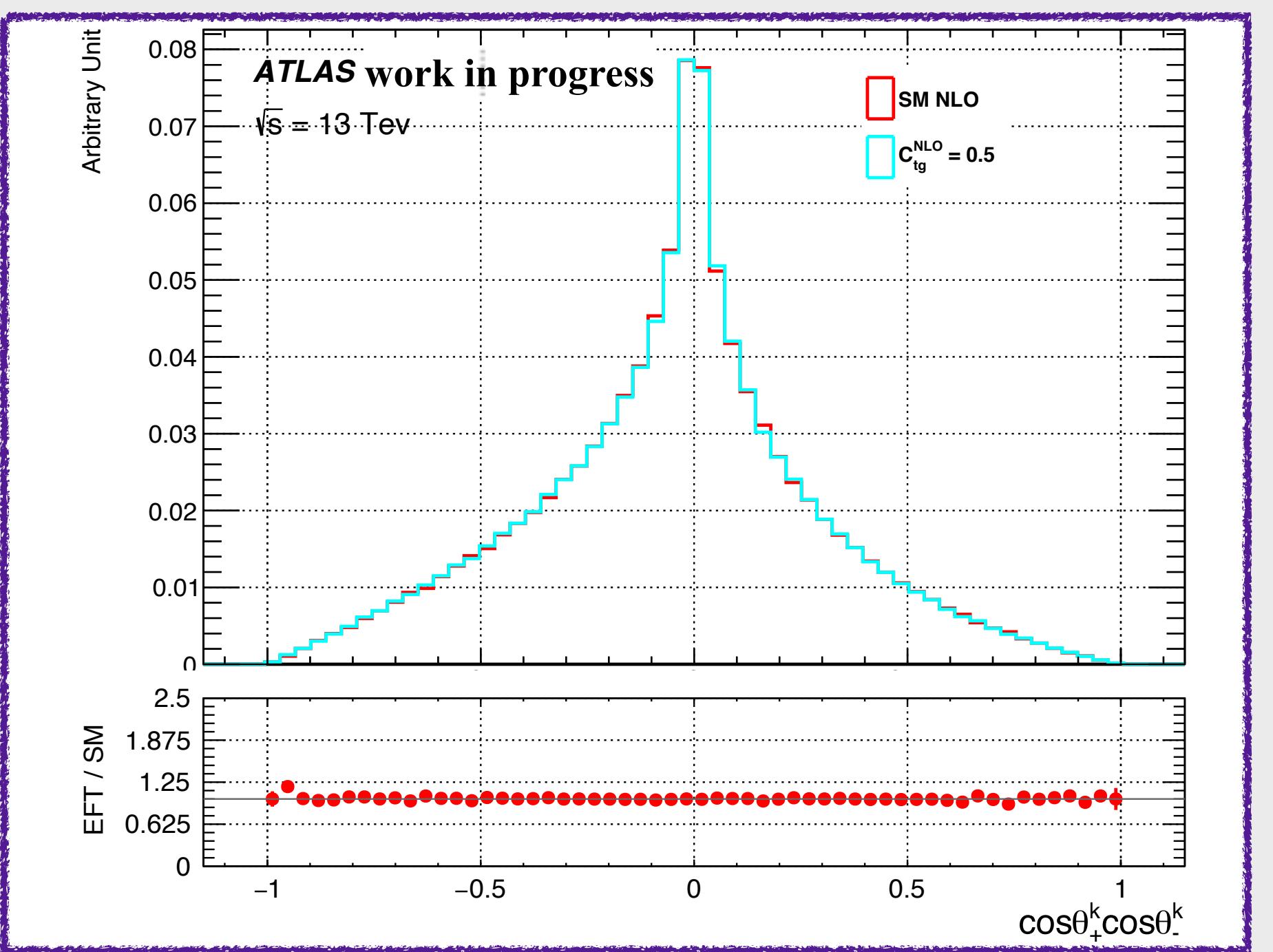
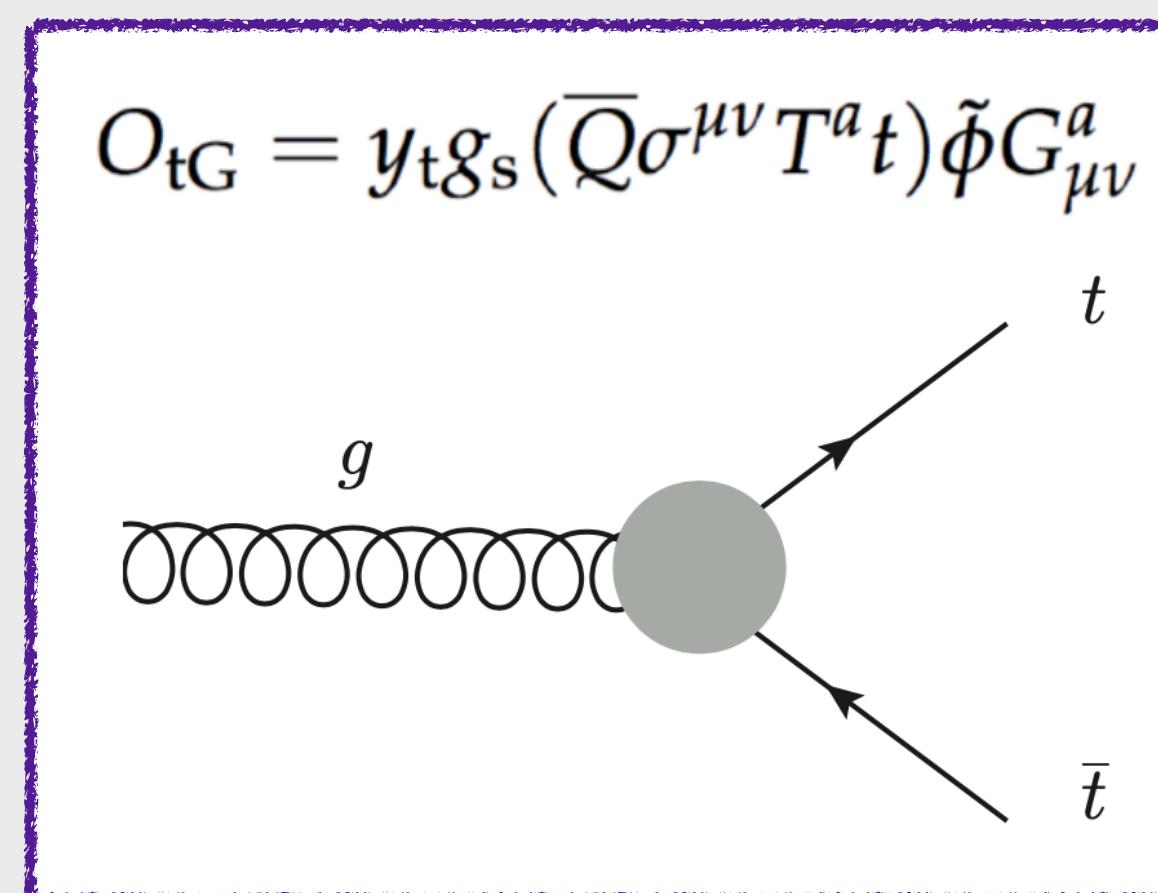
- The measured coefficients probe **most of the lowest-order EFT operators** relevant to LHC $t\bar{t}$ production.

$$O_{tG} = y_t g_s (\bar{Q} \sigma^{\mu\nu} T^a t) \tilde{\phi} G_{\mu\nu}^a$$

A Feynman diagram illustrating the process of top quark-antiquark (tt-bar) production via gluon-gluon fusion. A gluon (g) represented by a wavy line enters from the left and interacts with another gluon at a vertex. This interaction produces a virtual top quark (t) and an antiquark (t-bar). The top quark then decays at a vertex into a standard top quark (t) and an antineutrino (nu-bar). The antiquark (t-bar) also decays at a vertex into a standard antiquark (t-bar) and an electron (e). The final state consists of the standard top quark (t) and the electron (e).

EFT interpretation

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SM NLO : $C(k, k) = 0.366313 +/- 0.0042 \text{ (stat)}$

Ctg NLO : $C(k, k) = 0.375982 +/- 0.0042 \text{ (stat)}$

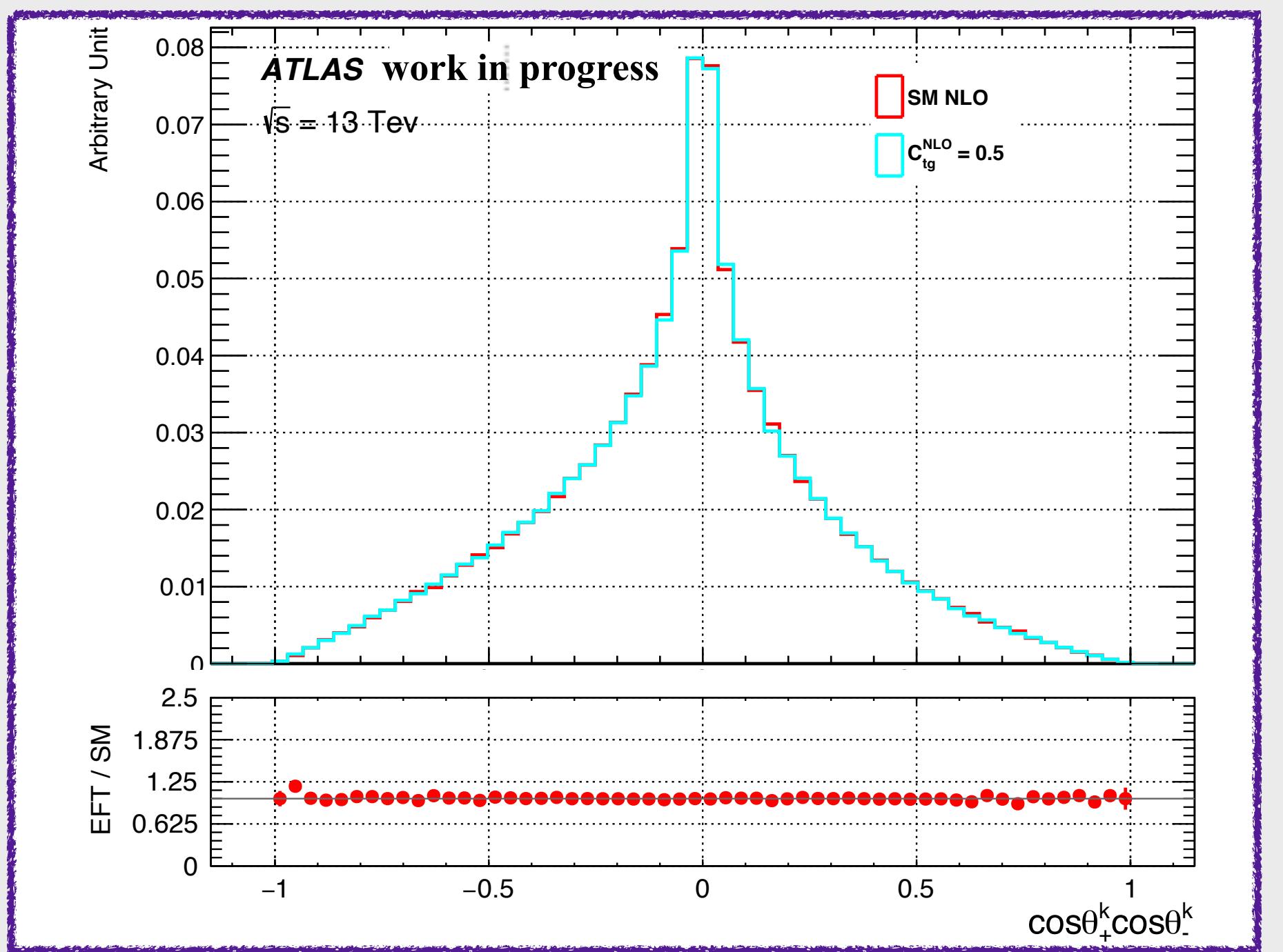
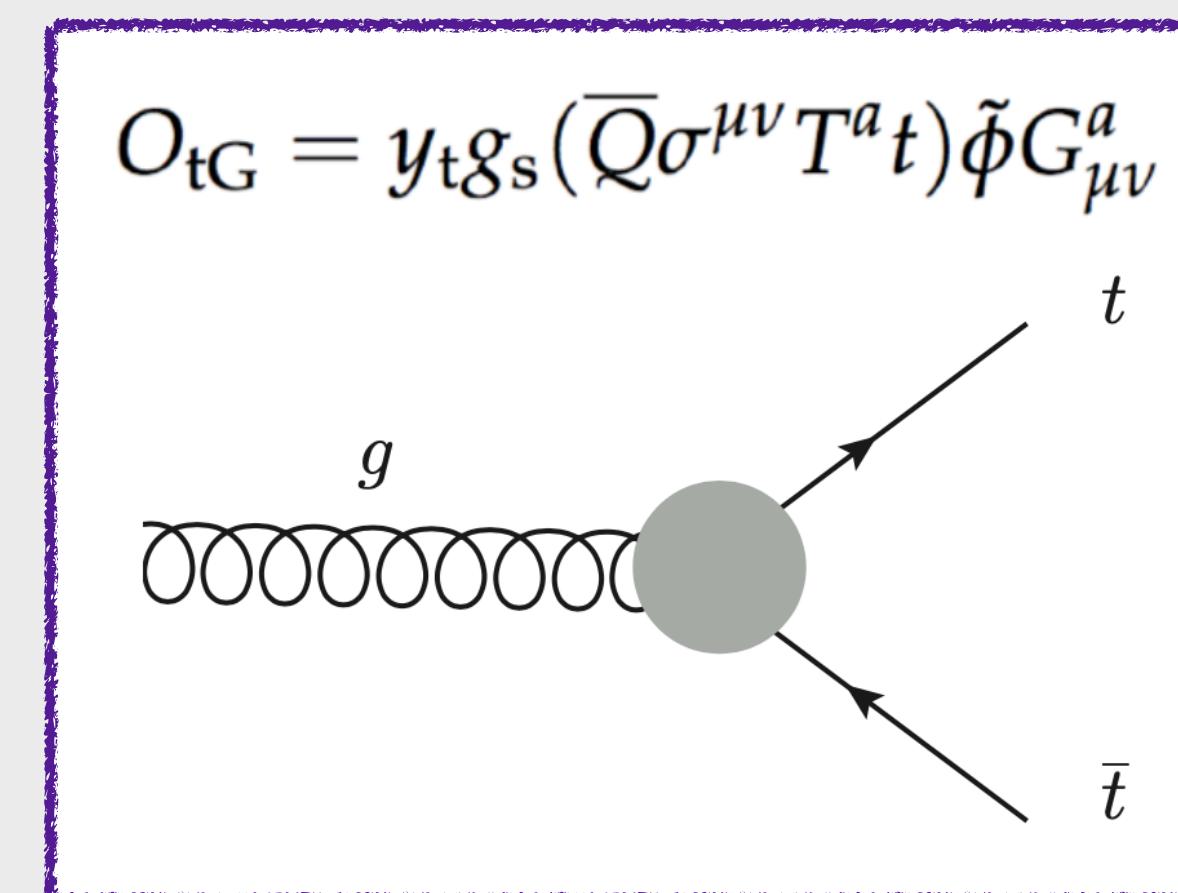
EFT interpretation

- The measured coefficients probe **most of the lowest-order EFT operators relevant to LHC $t\bar{t}$ production.**

Spin correlation :

- The impact of C_{tg} is low.

- For other spin corrections observables, the effect is very low or not observed.



SM NLO : $C(k, k) = 0.366313 +/- 0.0042$ (stat)

C_{tg} NLO : $C(k, k) = 0.375982 +/- 0.0042$ (stat)

Compute α_i and β_i

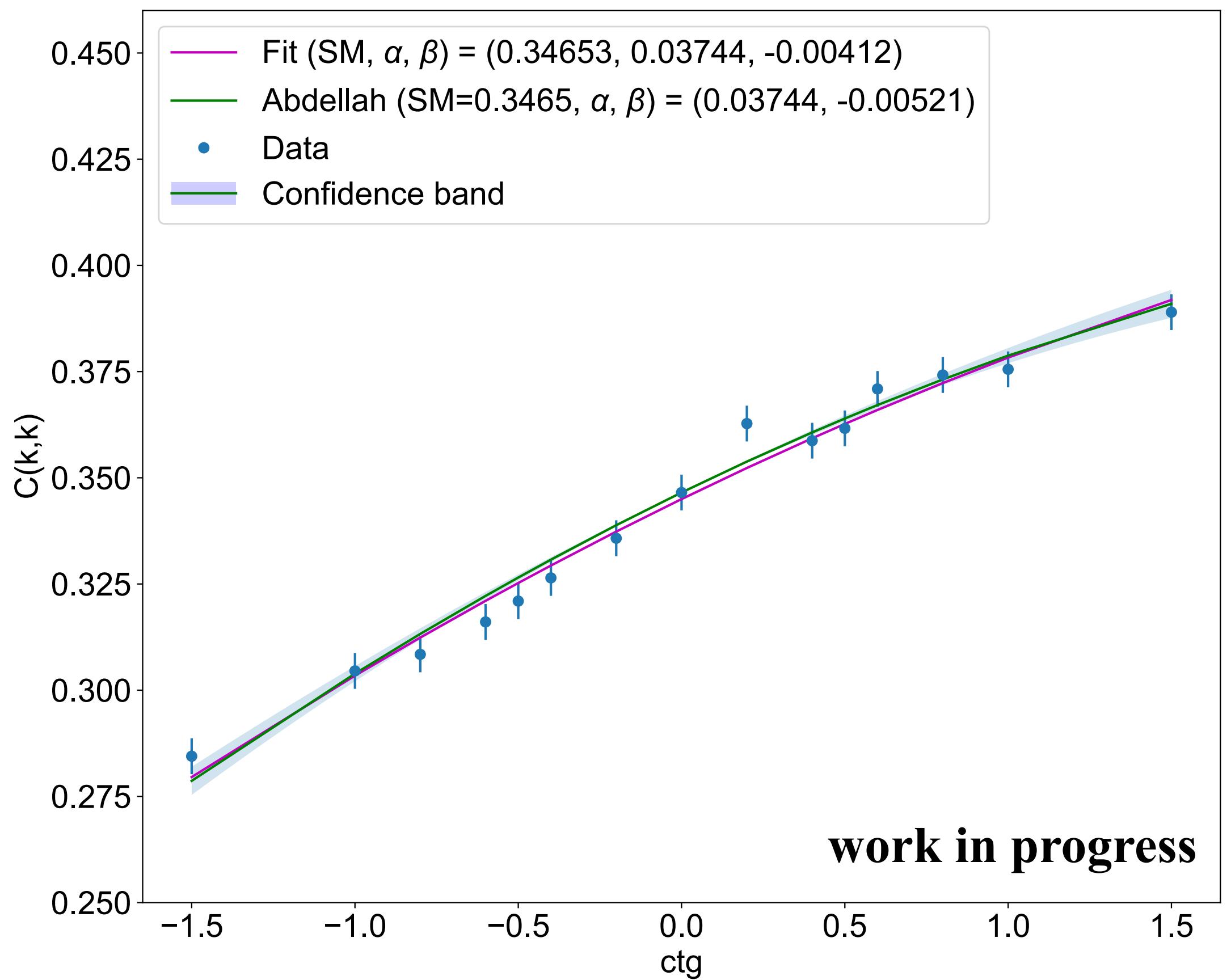
$$C(k, k)_{EFT} \approx C(k, k)_{SM} + \boxed{\alpha_{tg} \frac{c_{tg}}{\Lambda^2}} + \boxed{\beta_{tg} \frac{c_{tg}^2}{\Lambda^4}}$$

Linear Term

quadratic Term

α_i/Λ^2 and β_i/Λ^4 at LO : $C(k,k)$

SMEFT@NLO model

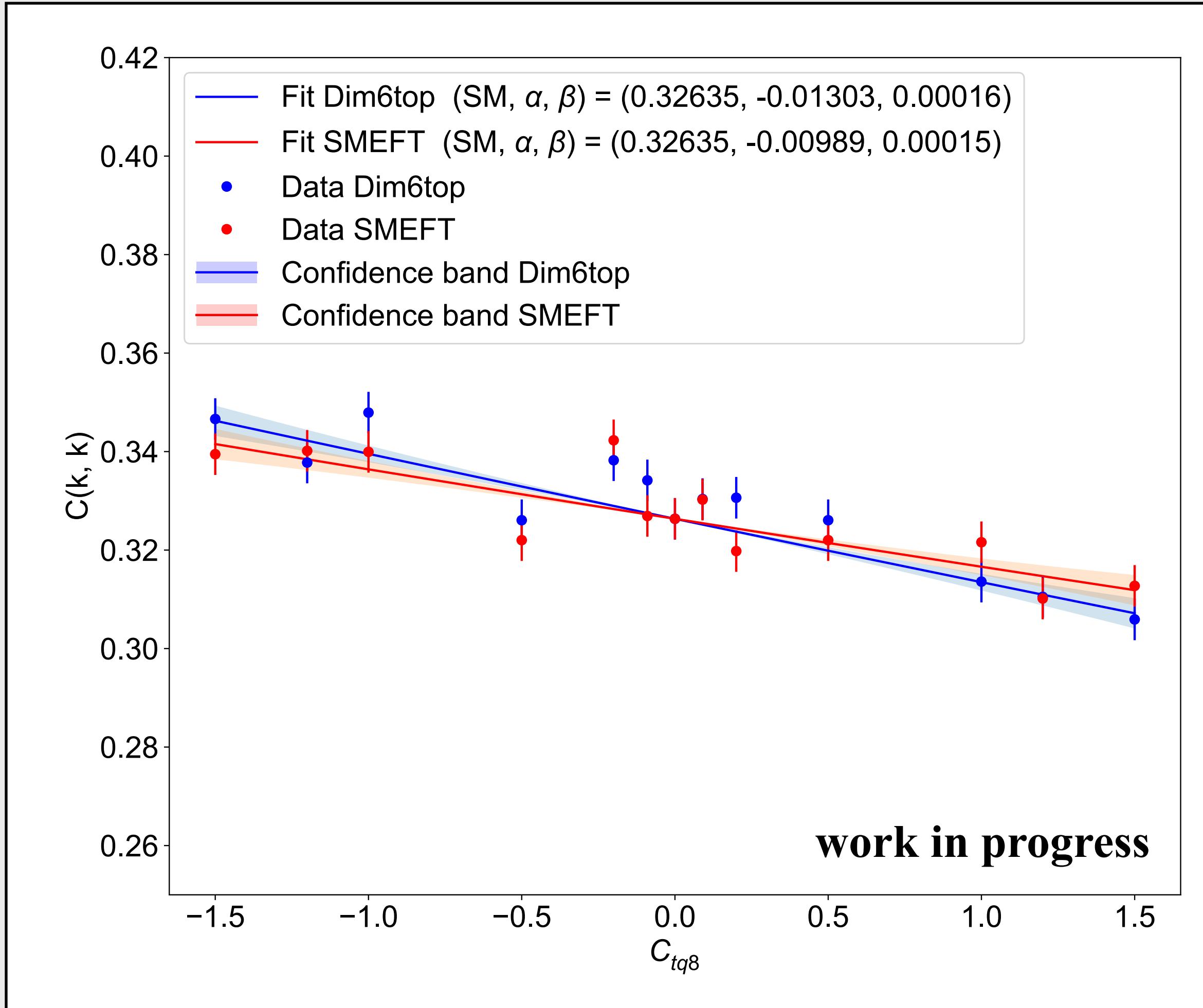


- SMEFT model is used to generate MC sample
- The value of C_{tg} affect the spin correlation

$$C(k, k)_{EFT} \approx C(k, k)_{SM} + \boxed{\alpha_{tg} \frac{c_{tg}}{\Lambda^2}} + \boxed{\beta_{tg} \frac{c_{tg}^2}{\Lambda^4}}$$

Linear Term quadratic Term

Comparaison between SMEFT and Dim6Top Model



SMEFT model and Dim6top model show appx. same value of α_{ctq8} and β_{ctq8} [within the statistical uncertainties.]

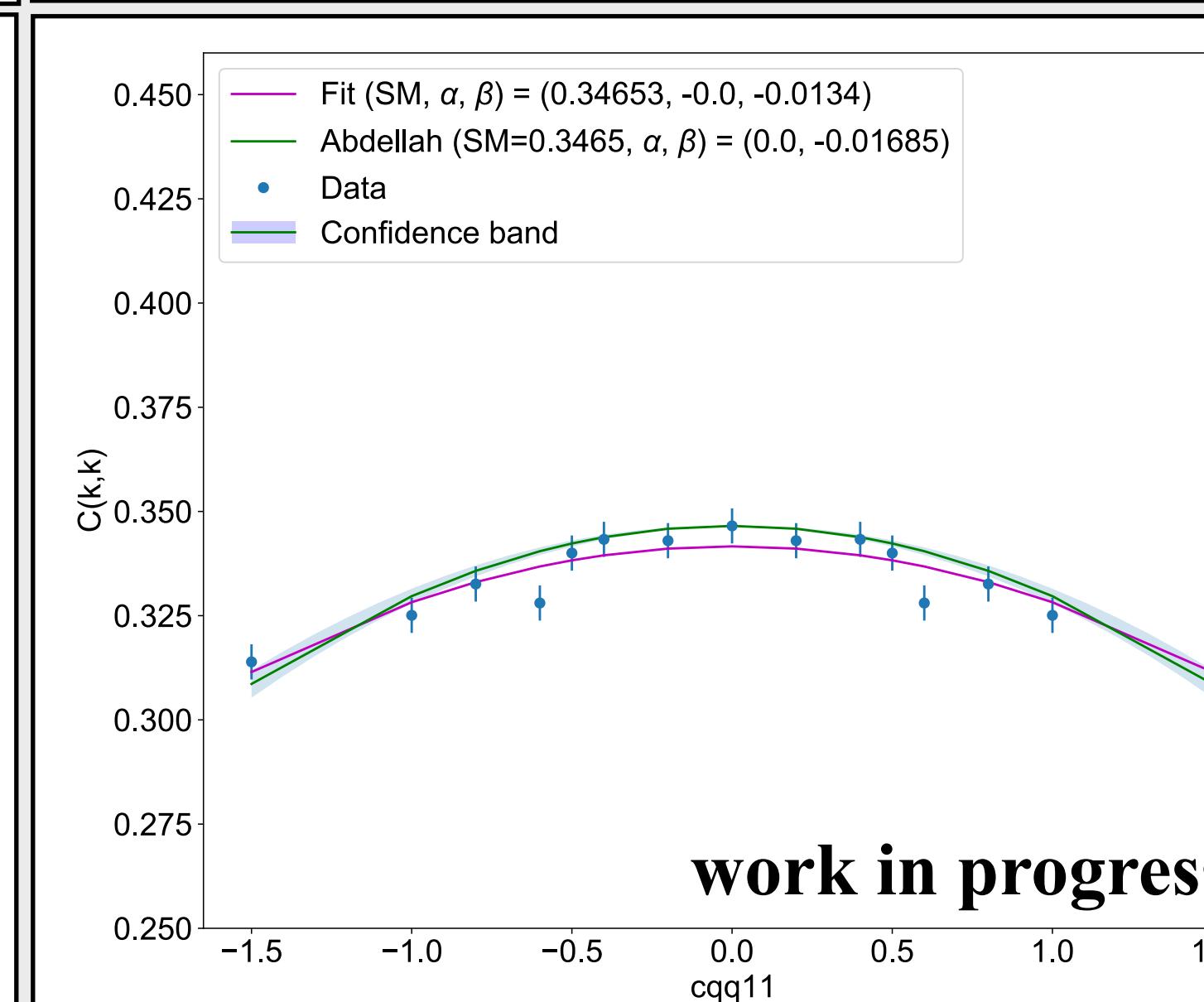
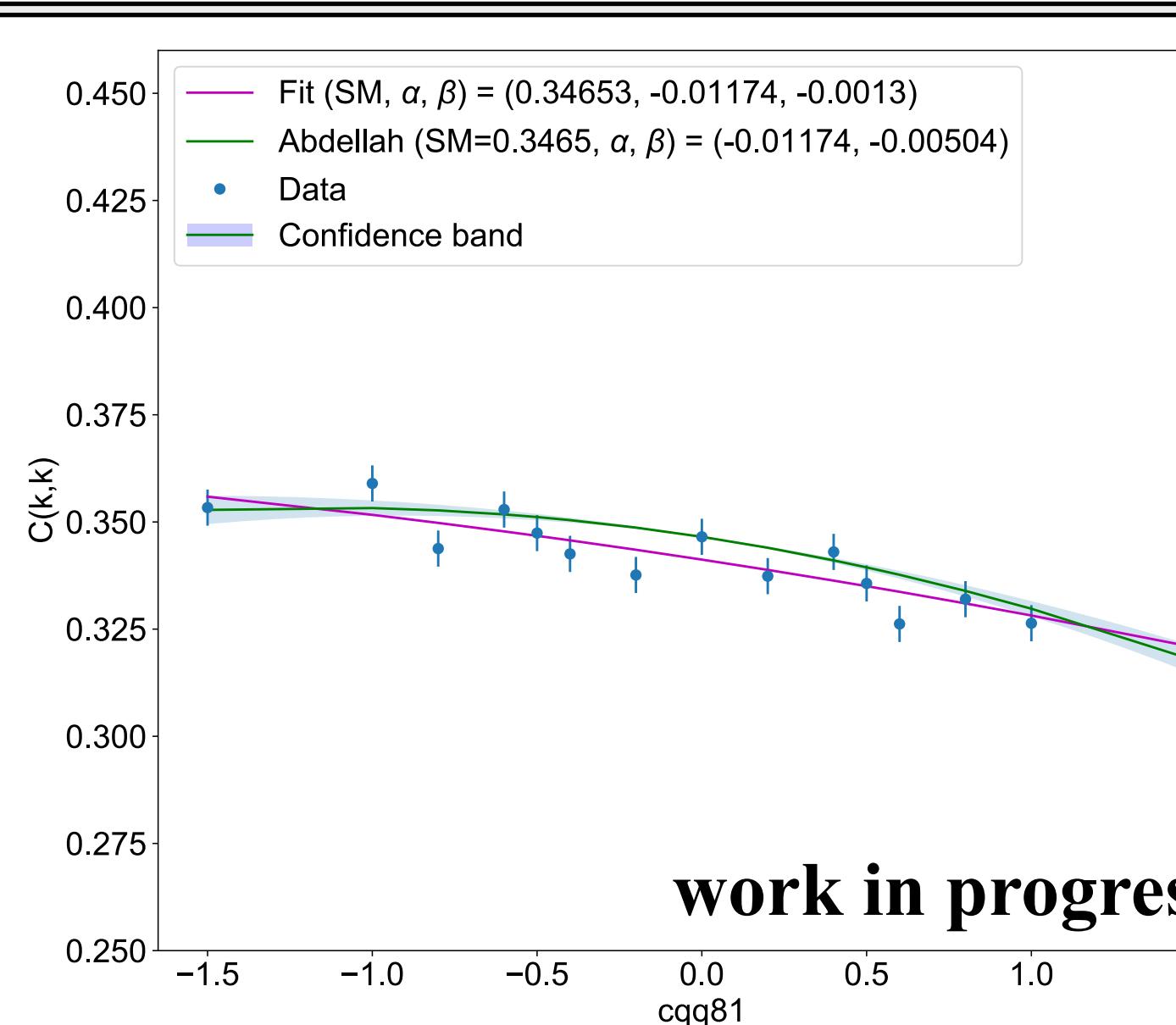
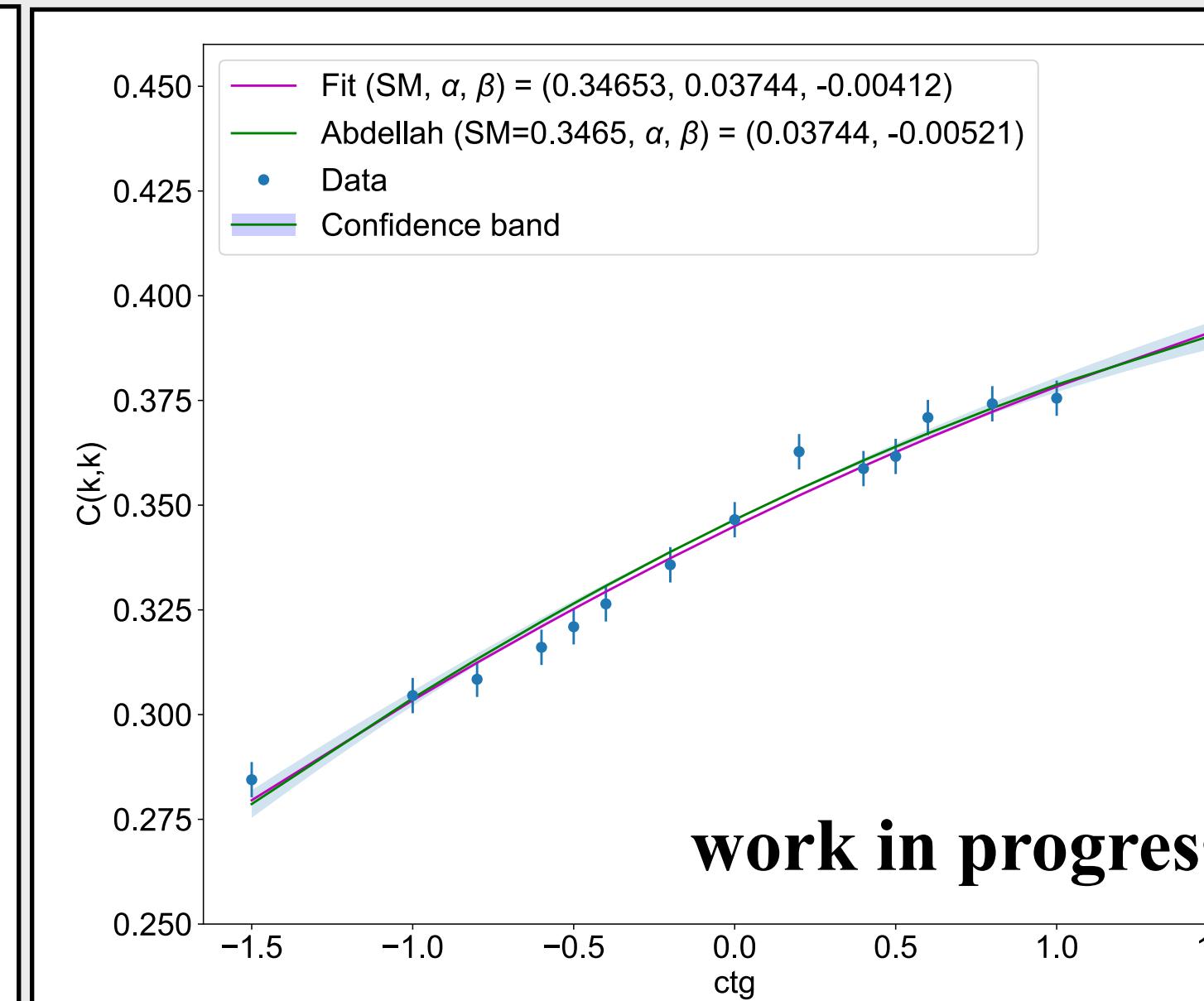
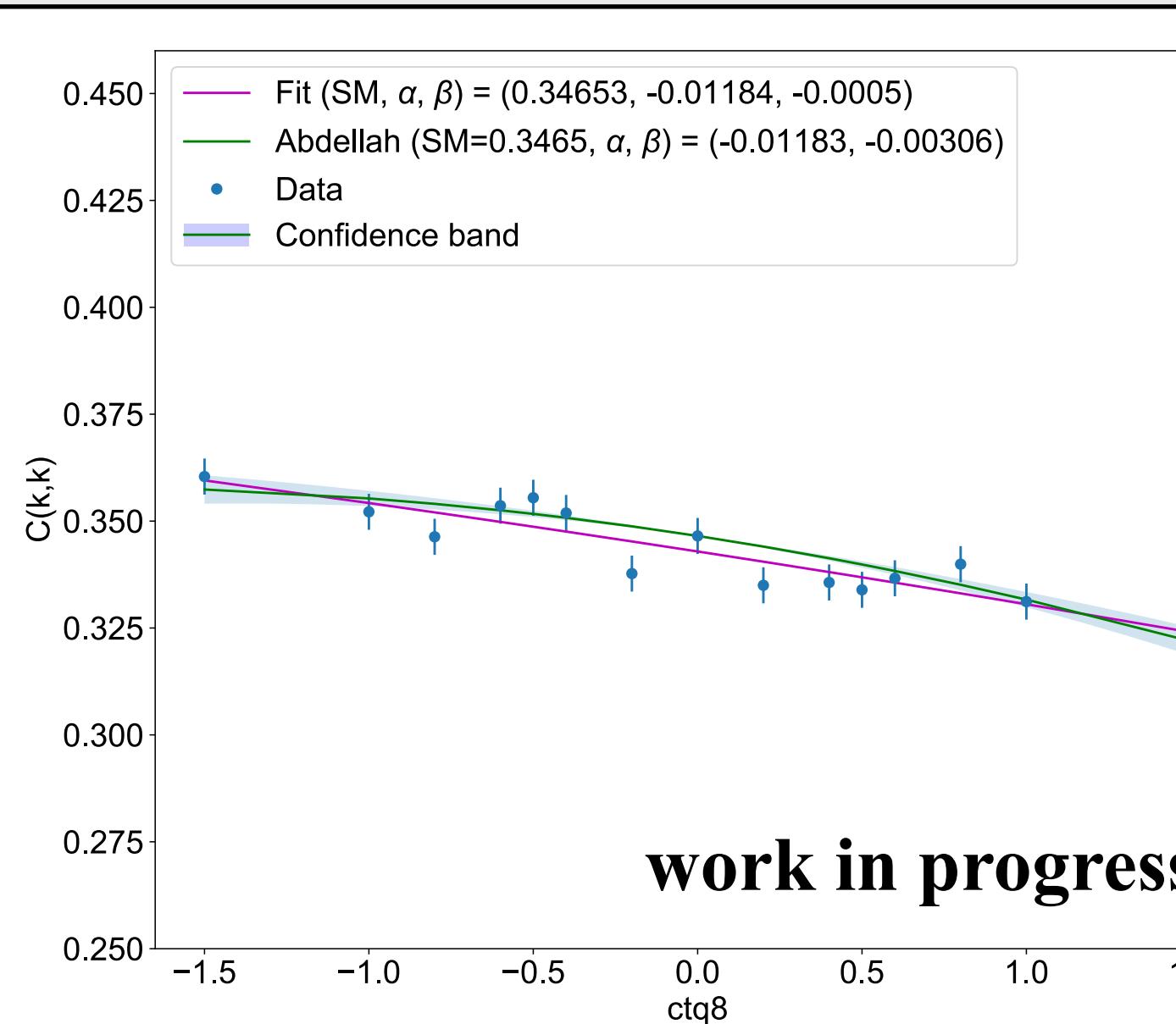
Summary

1. Spin observables is sensitive to a different coefficient of the spin density matrix of $t\bar{t}$ production.
2. Precision top quark spin measurements are a powerful probe of new physics and complementary to other approaches
3. Spin observables are sensitive to many BSM operator which can be use to constrain the Wilson coefficients
4. A Comparaison between SMEFT and Dim6top is shown

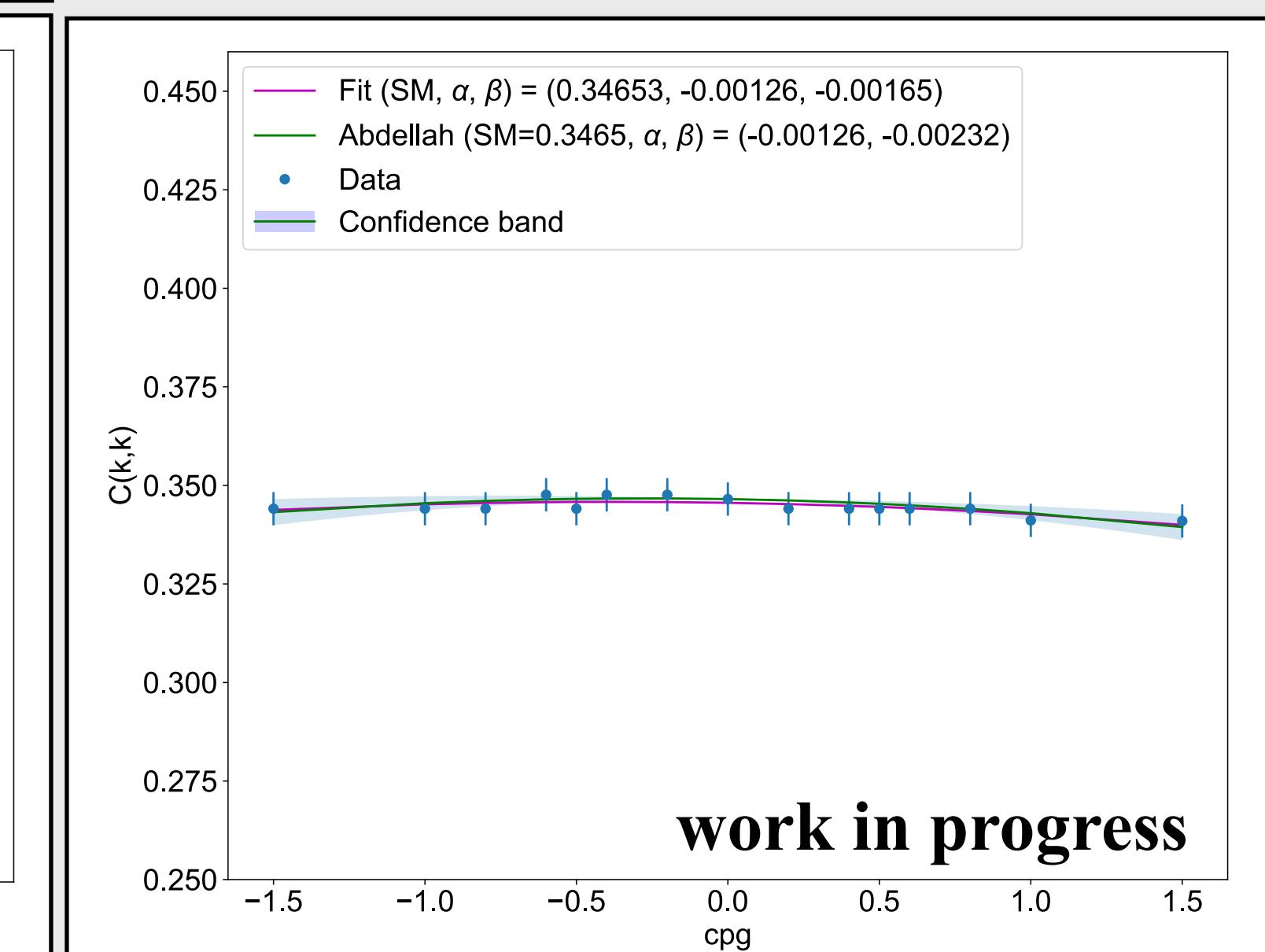
Back up

α_i/Λ^2 and β_i/Λ^4 at LO : $C(k,k)$

C(r,r) and C(n,n) in back-up



- SMEFT model is used to generate MC sample
- The value of WC affect the spin correlation

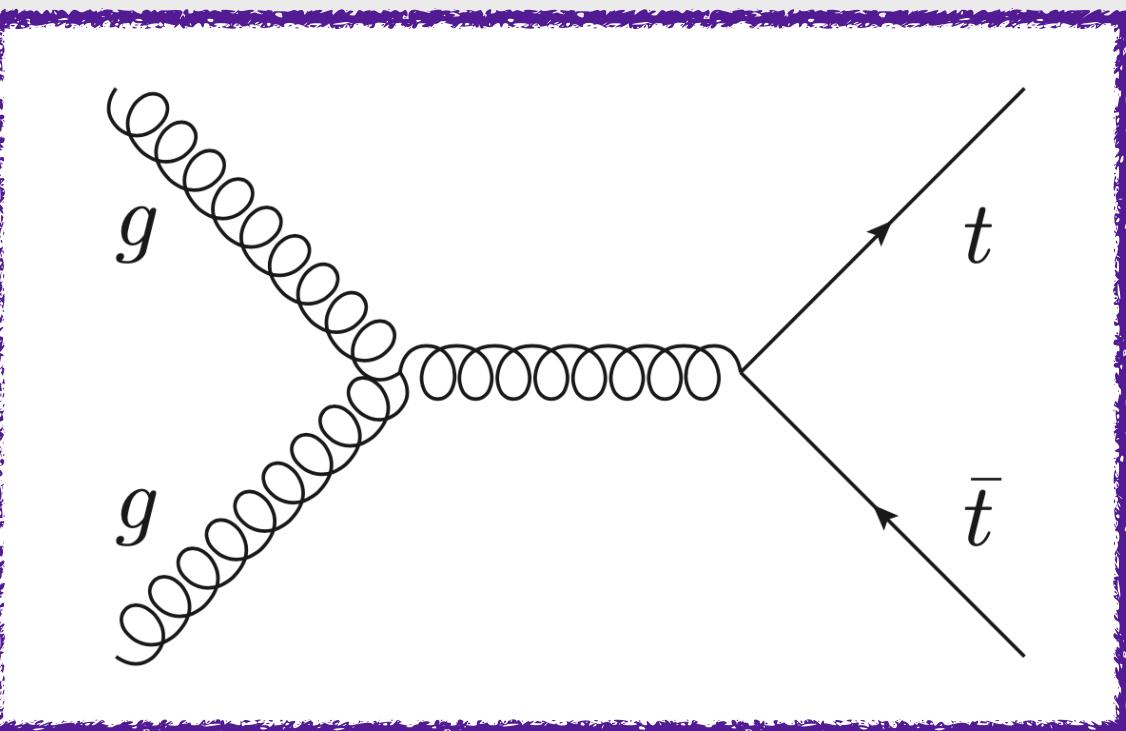


Introduction

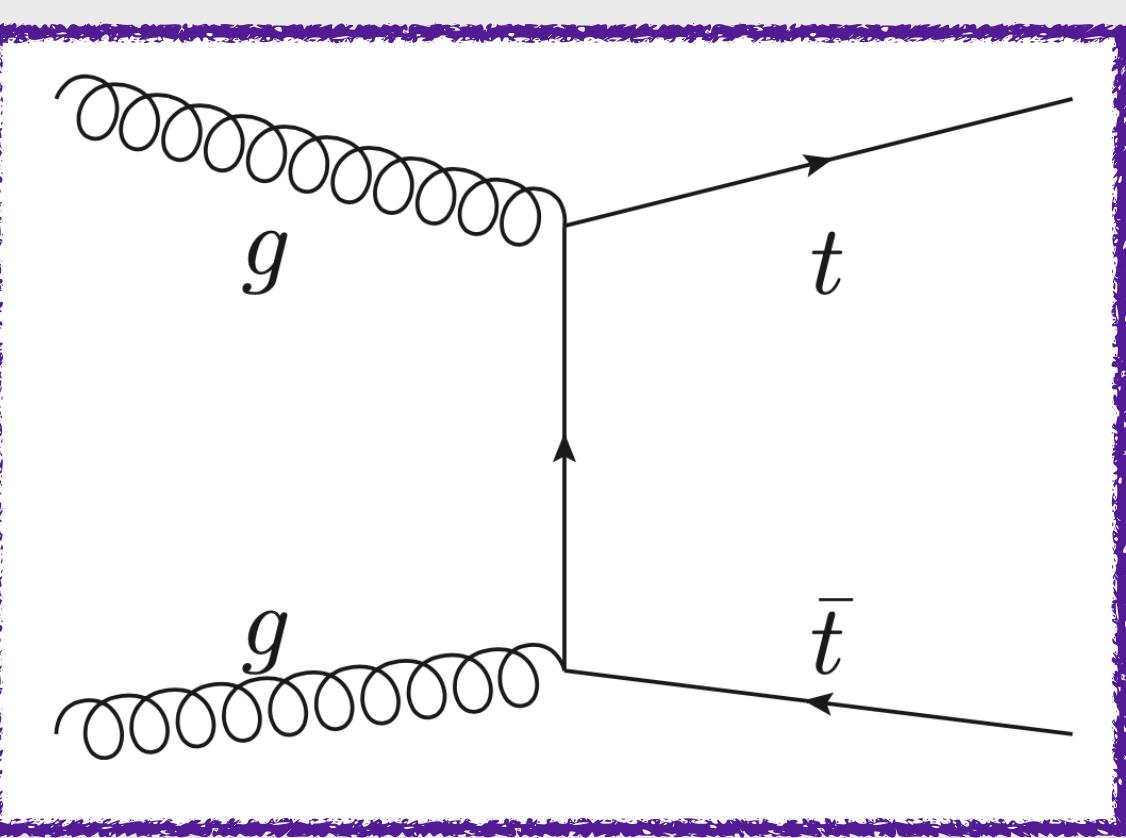
- ✓ Heaviest fundamental particle (known) :
 - ◆ $m_t = 173.34 \pm 0.27(\text{stat}) \pm 0.71(\text{syst})\text{GeV}$ [[link](#)]
- ✓ Short life time 10^{-24} , so it decay before hadronization.
- ✓ LHC is a Top Factory, large paire production.

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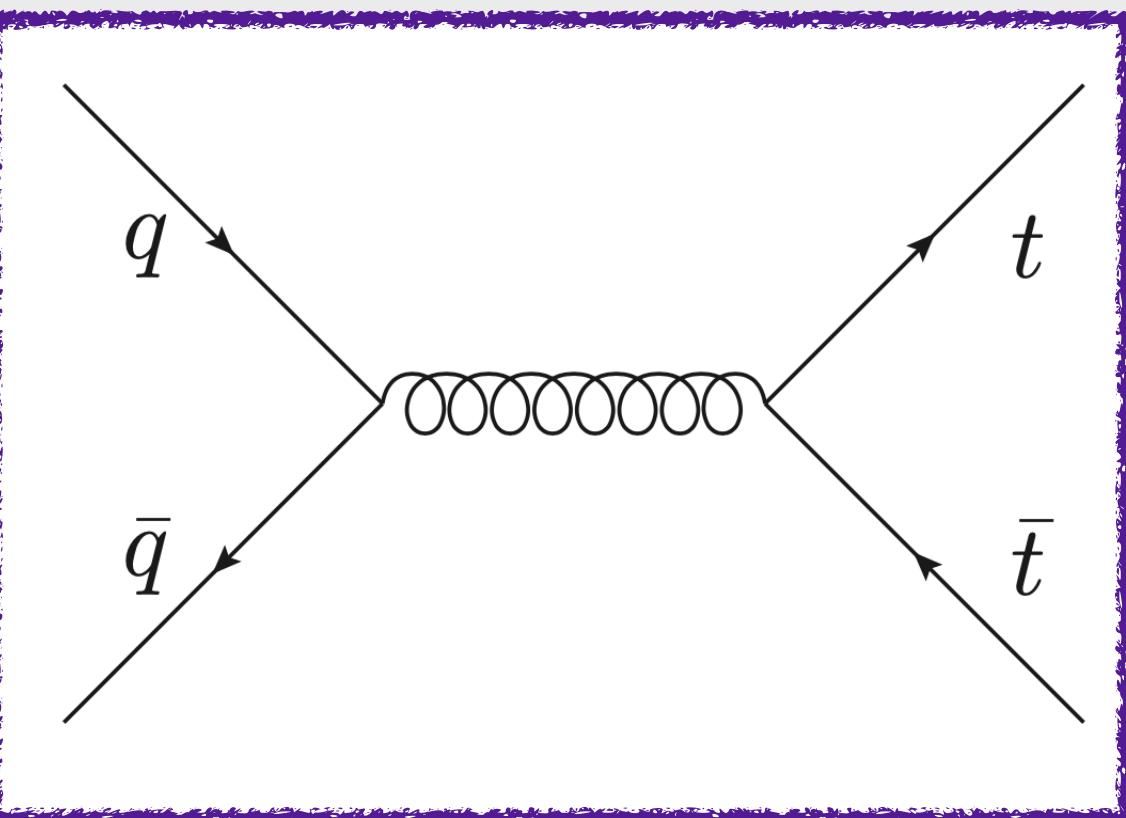
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- For leading-order (LO) : $q\bar{q}$ and gg initiated subprocesses contribute.



90%

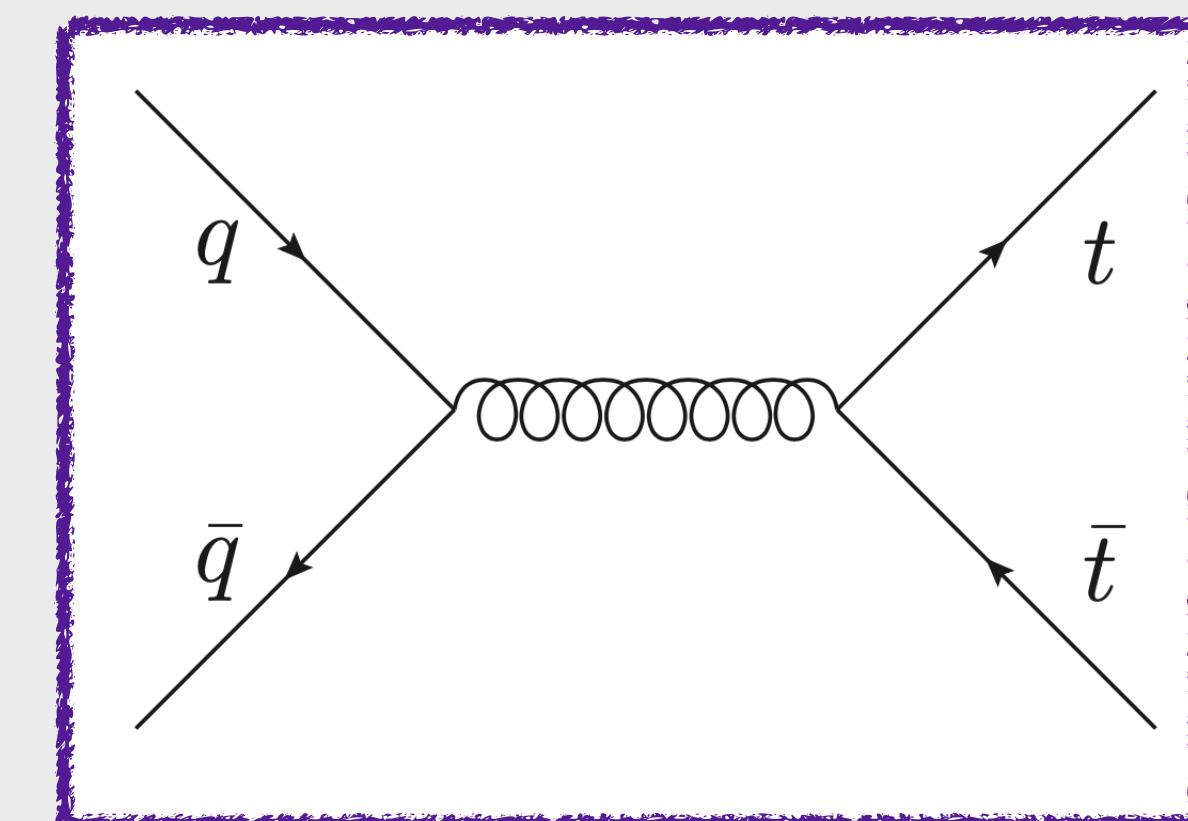
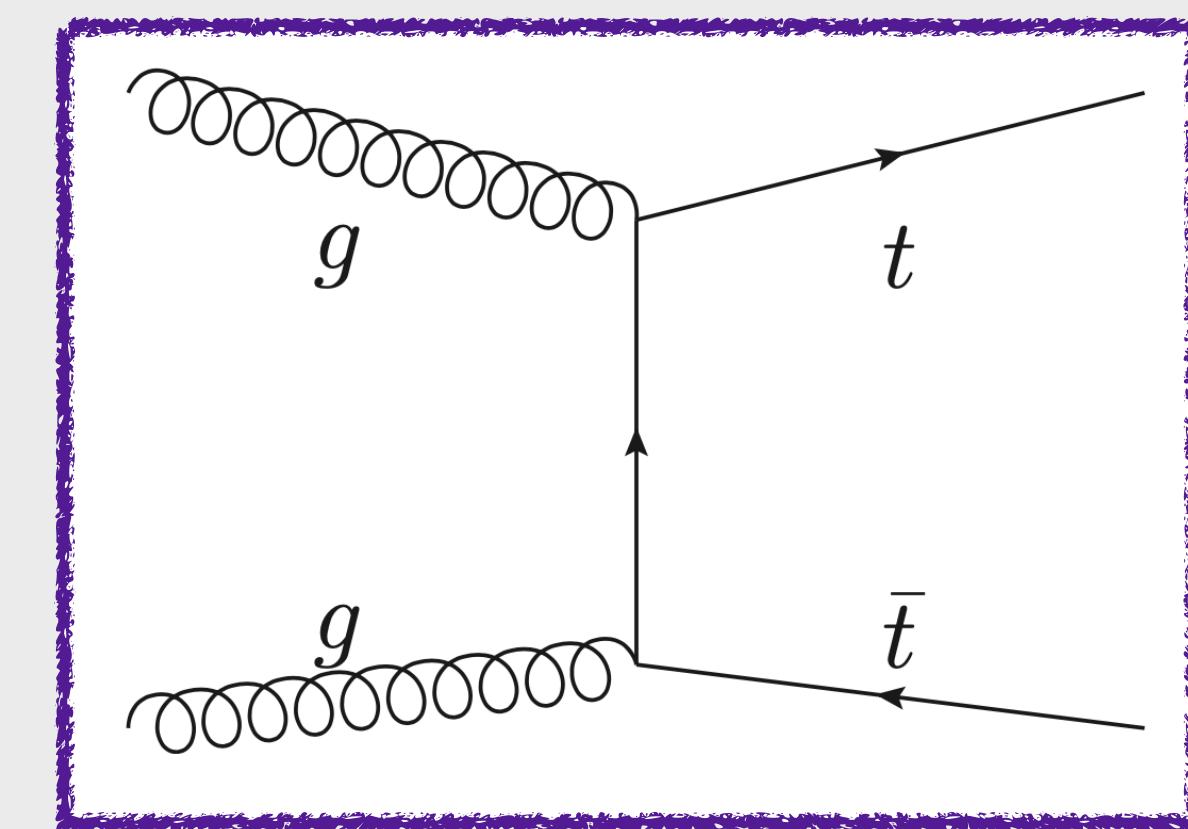
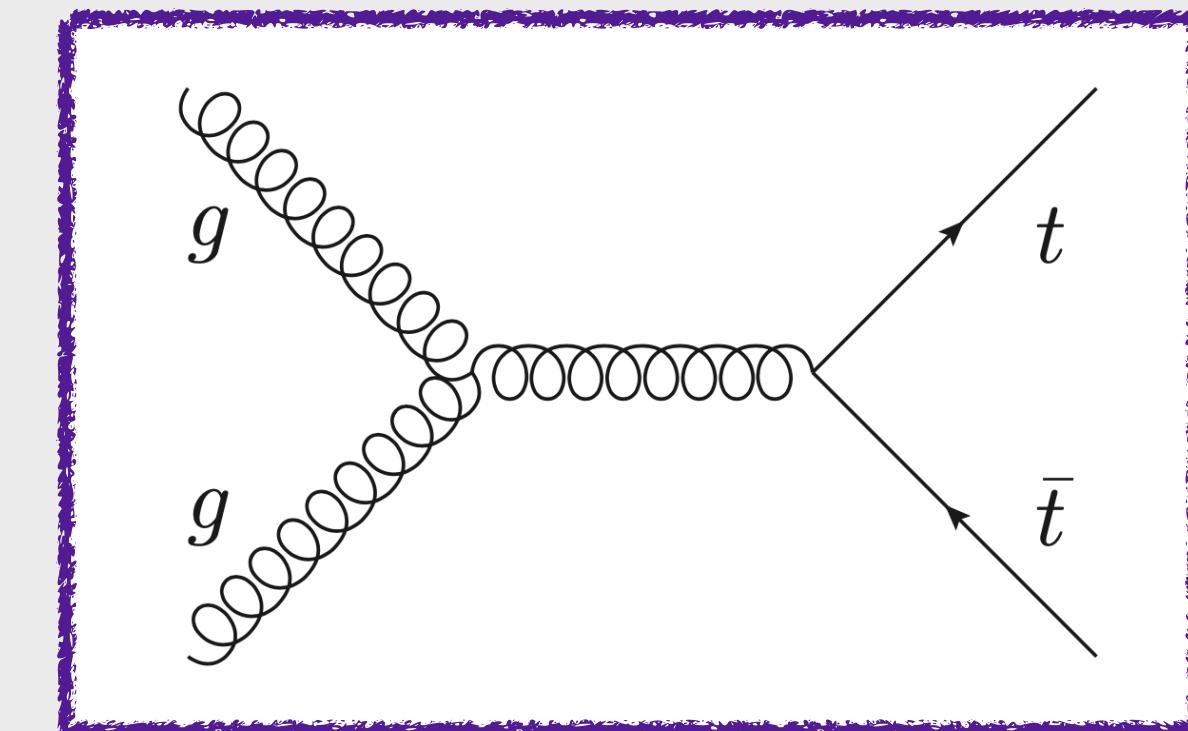
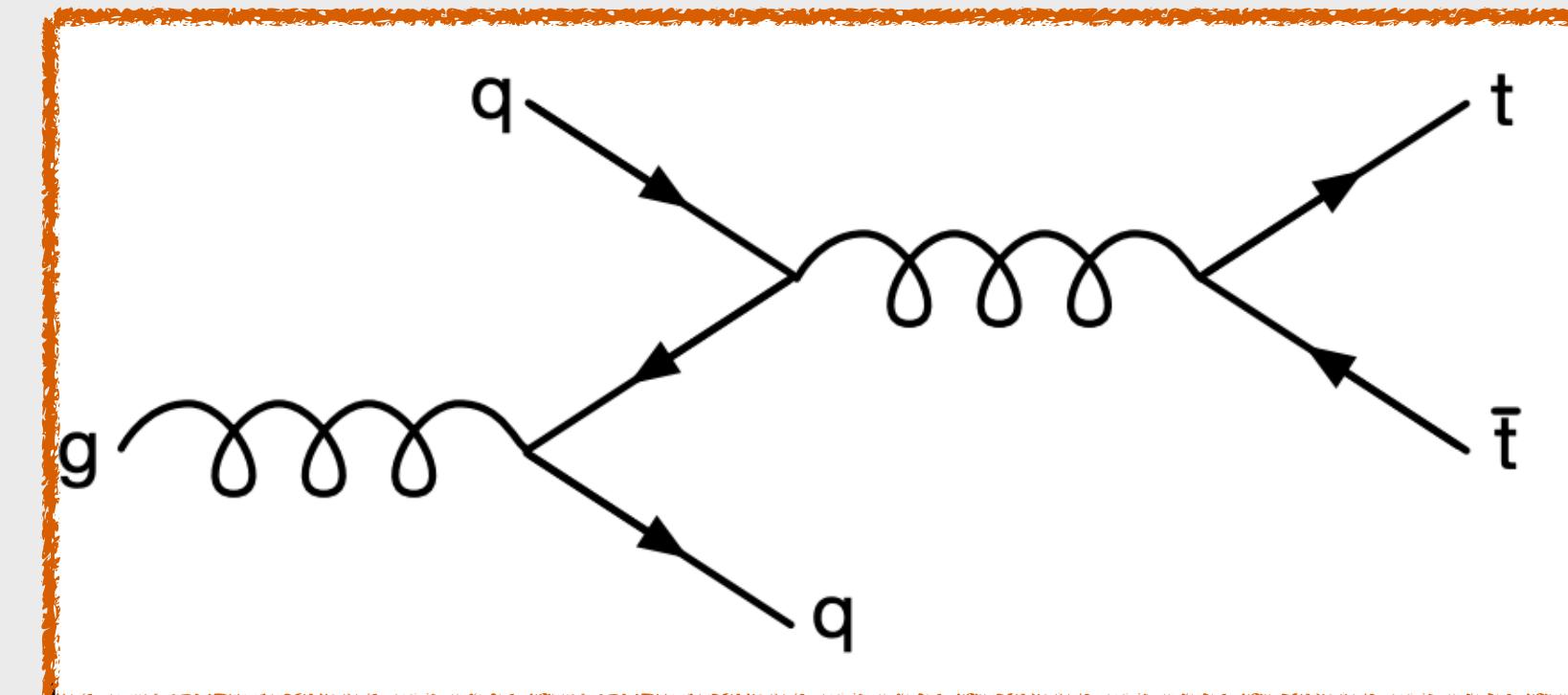
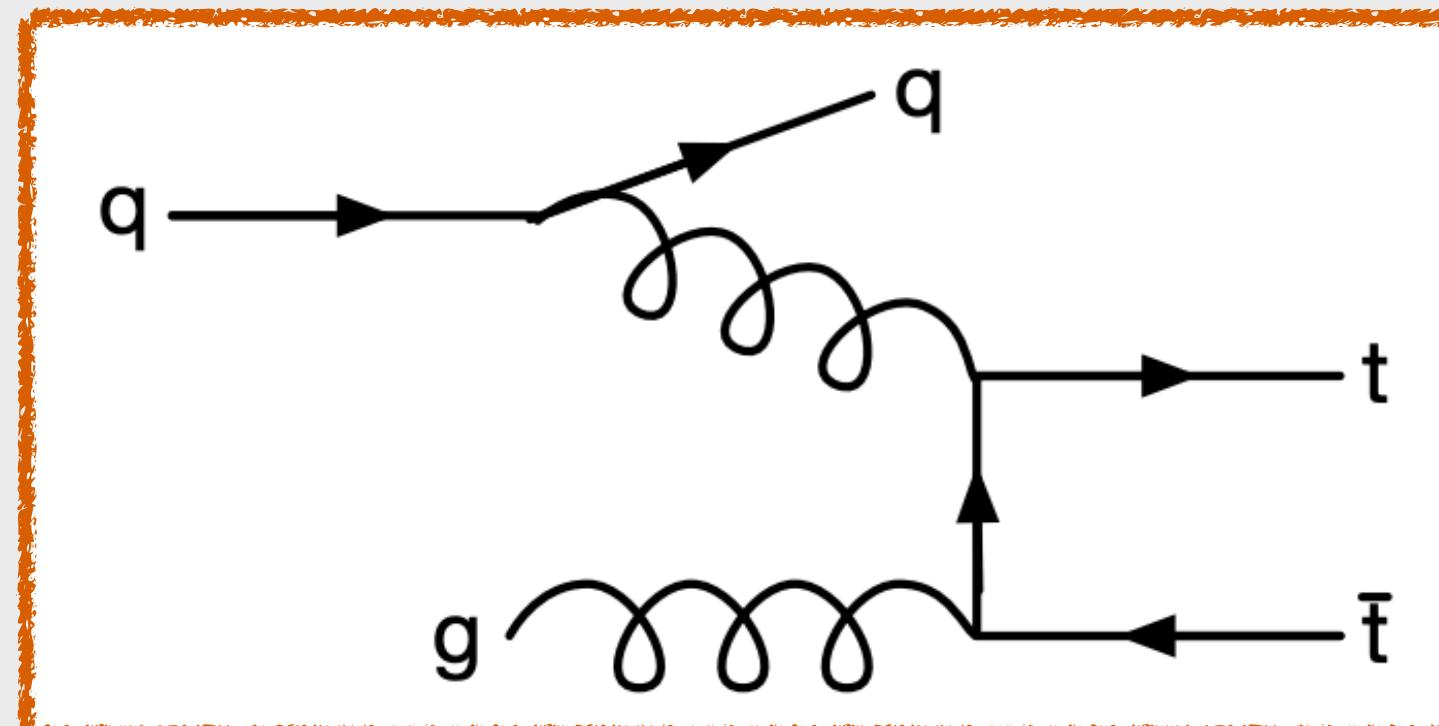


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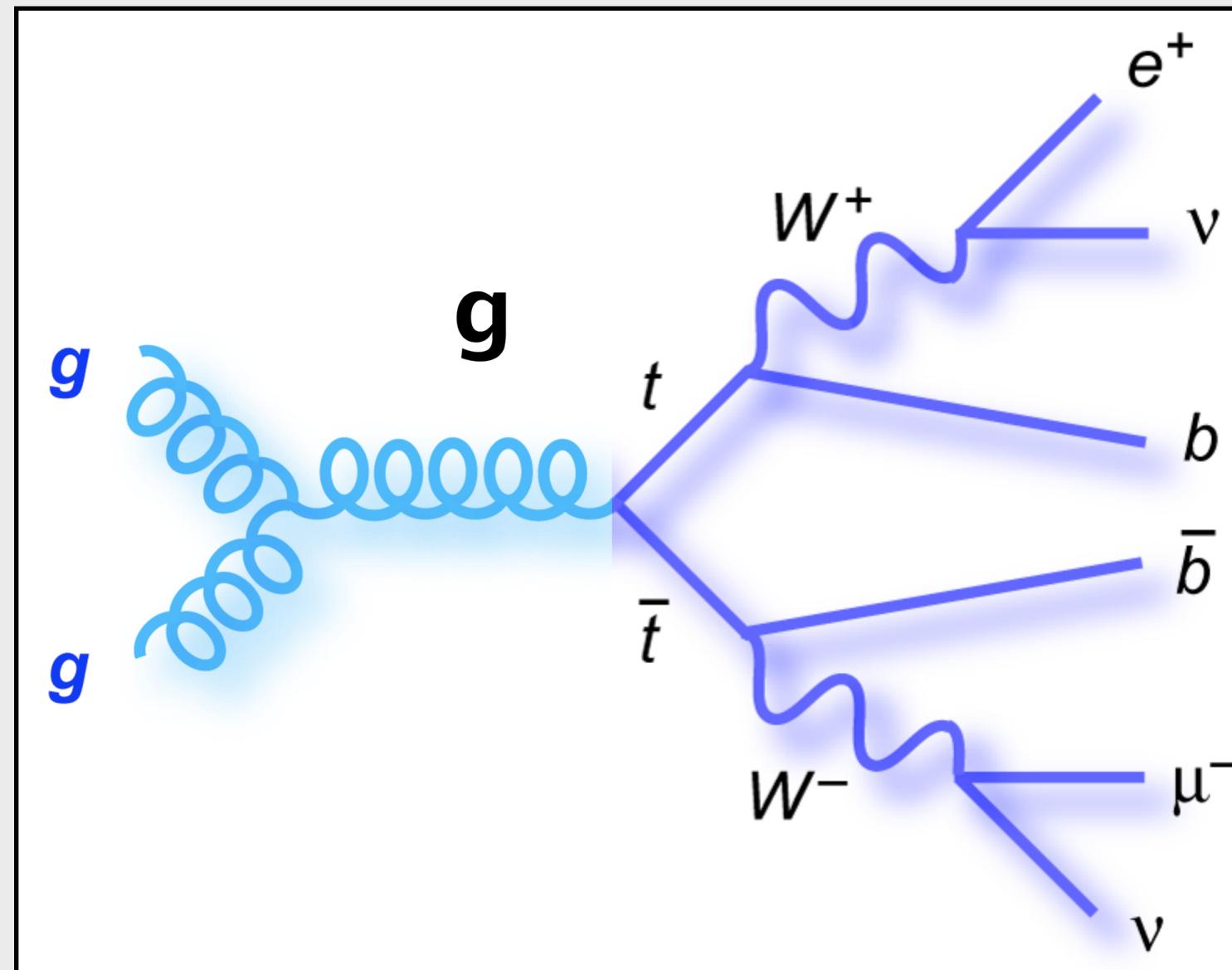
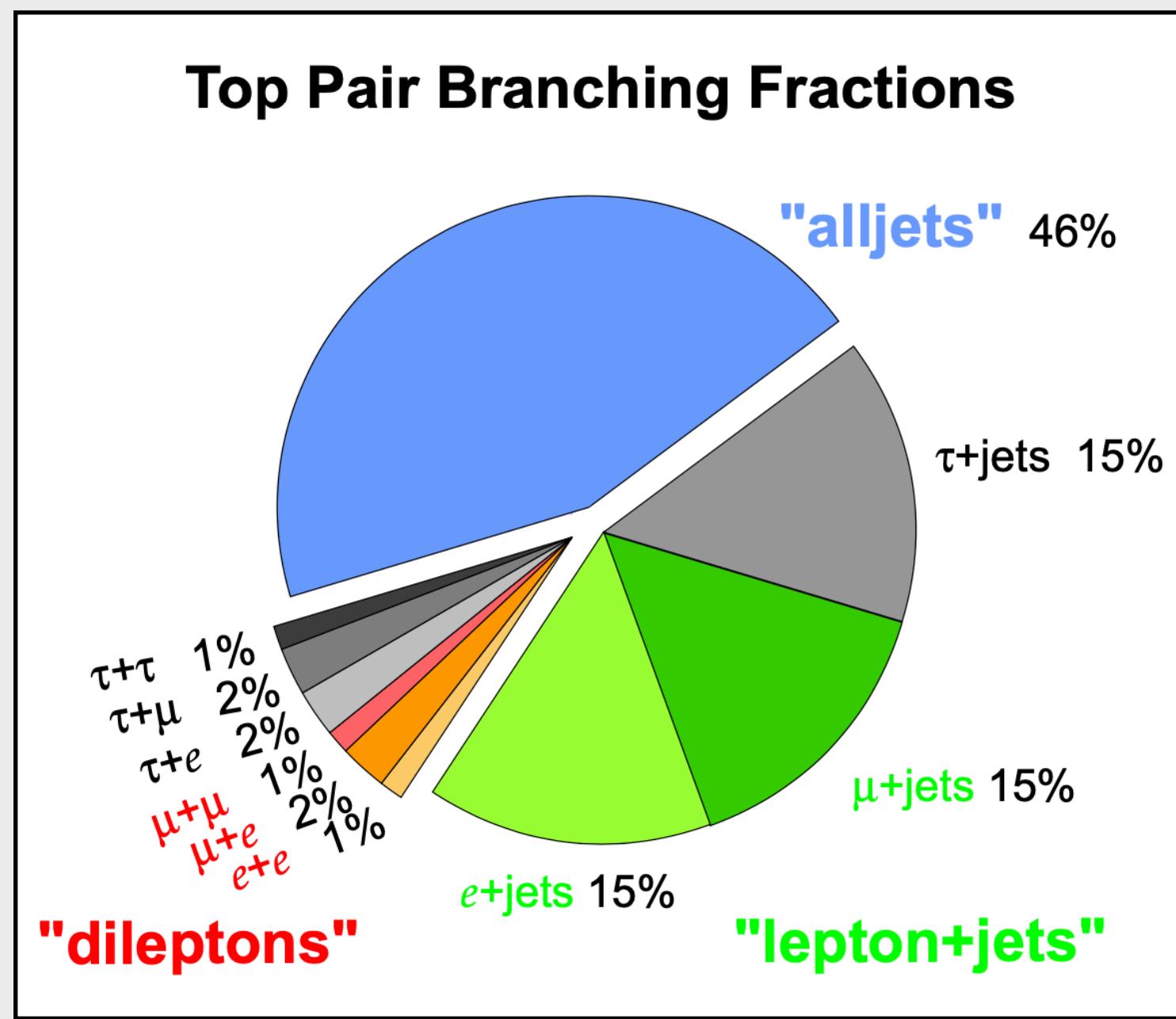
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- LHC is a Top Factory, large pair production.
- For leading-order (LO) : $q\bar{q}$ and gg initiated subprocesses contribute.
- For next-to-leading-order (NLO) : qg initiated subprocess contribute.



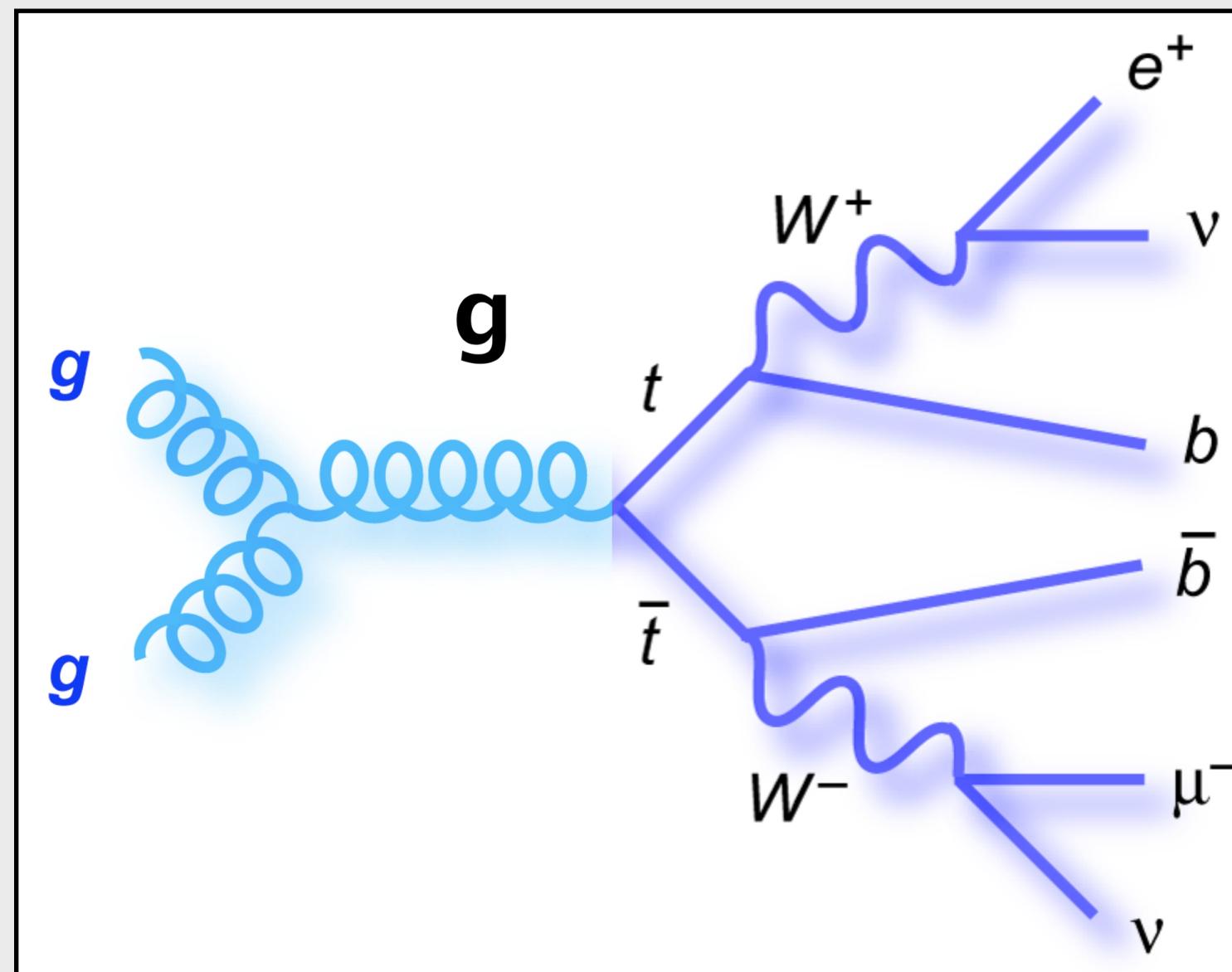
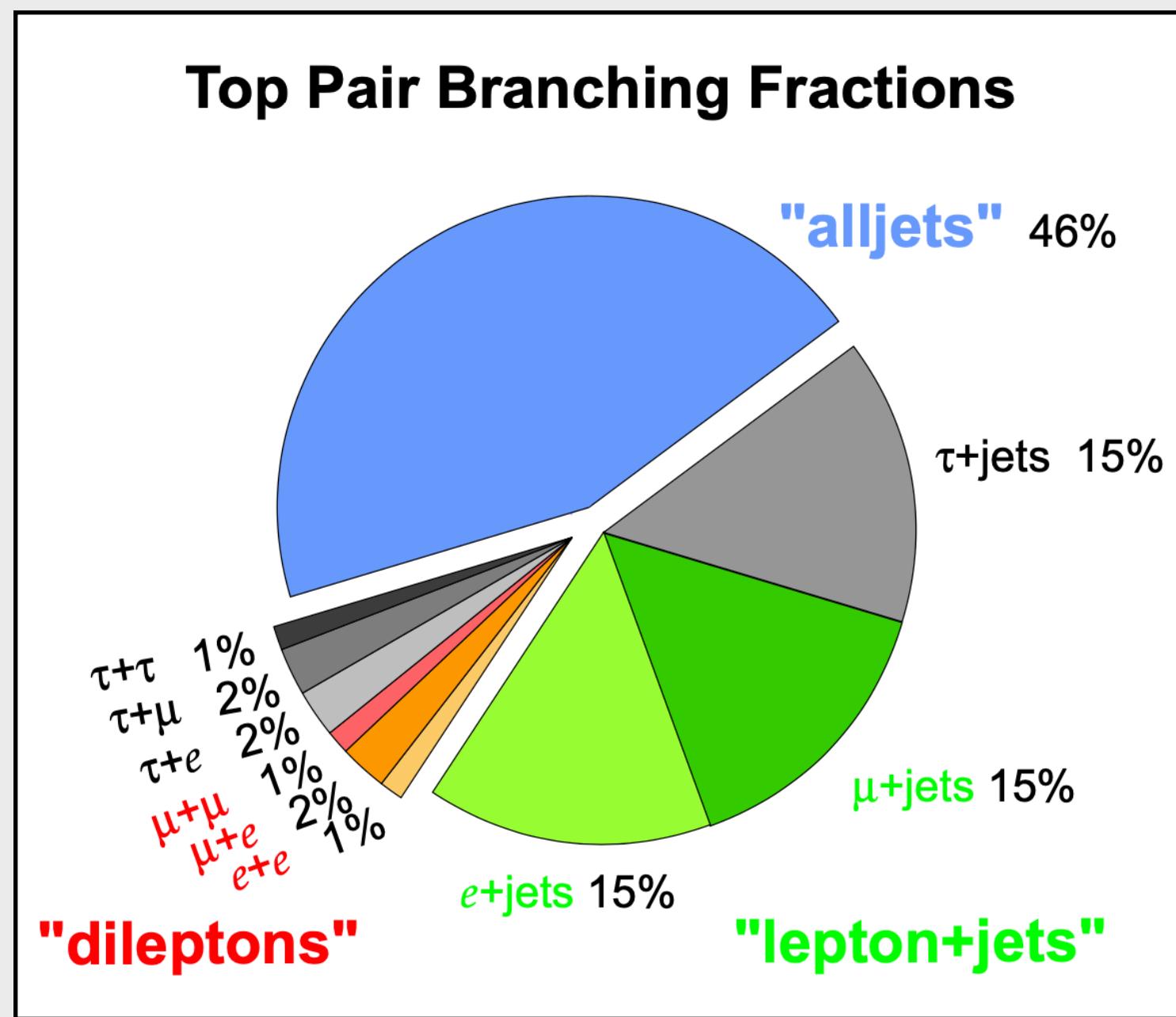
~4%

Top quarks pair decay mode



- All - hadronic
 - ❖ Largest BR
 - ❖ Largest QCD background
 - ❖ Event fully constrained
- Semi - leptonic
 - ❖ High BR
 - ❖ Medium background
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- Di-leptonic
 - ❖ Small BR
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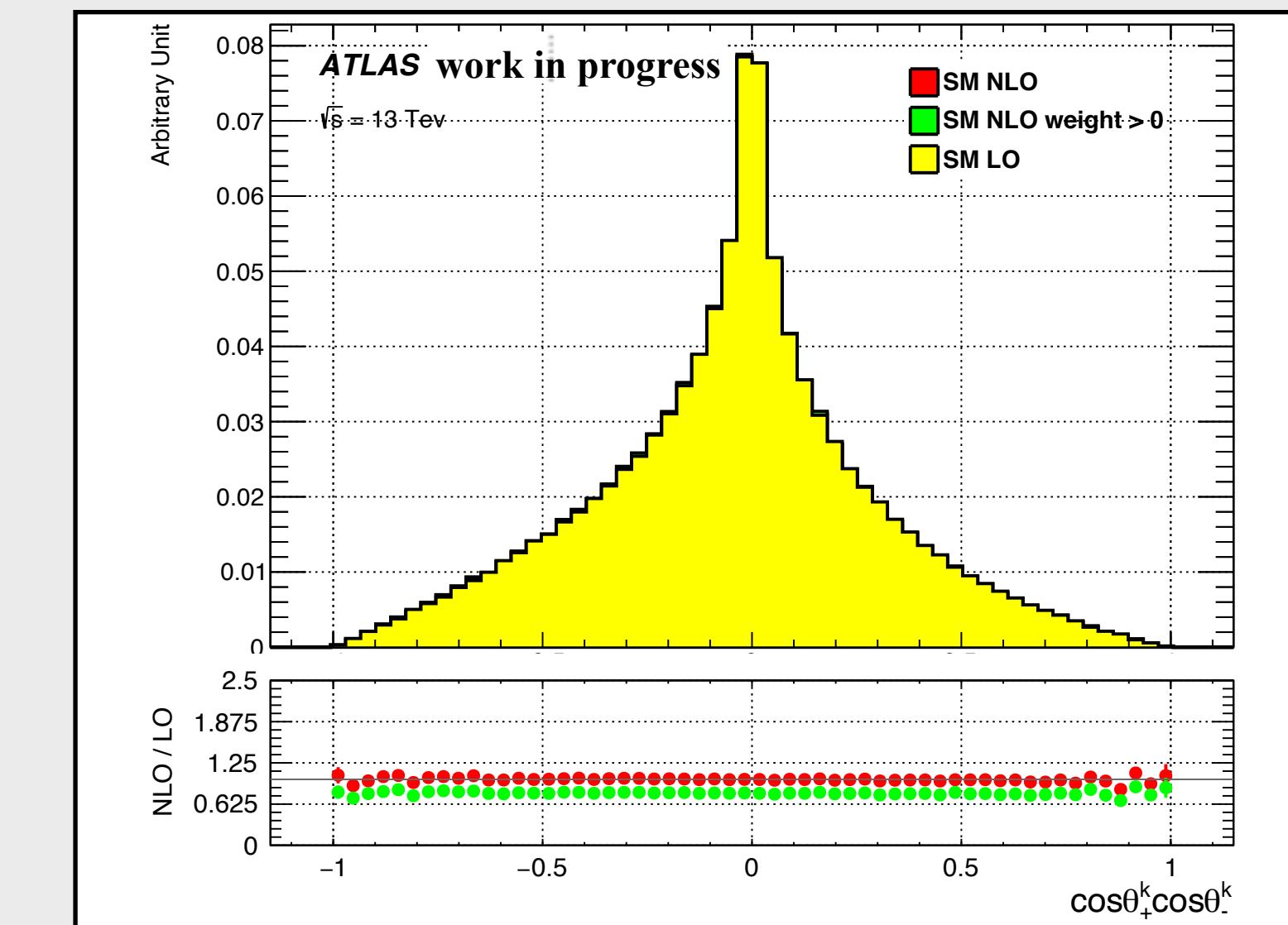
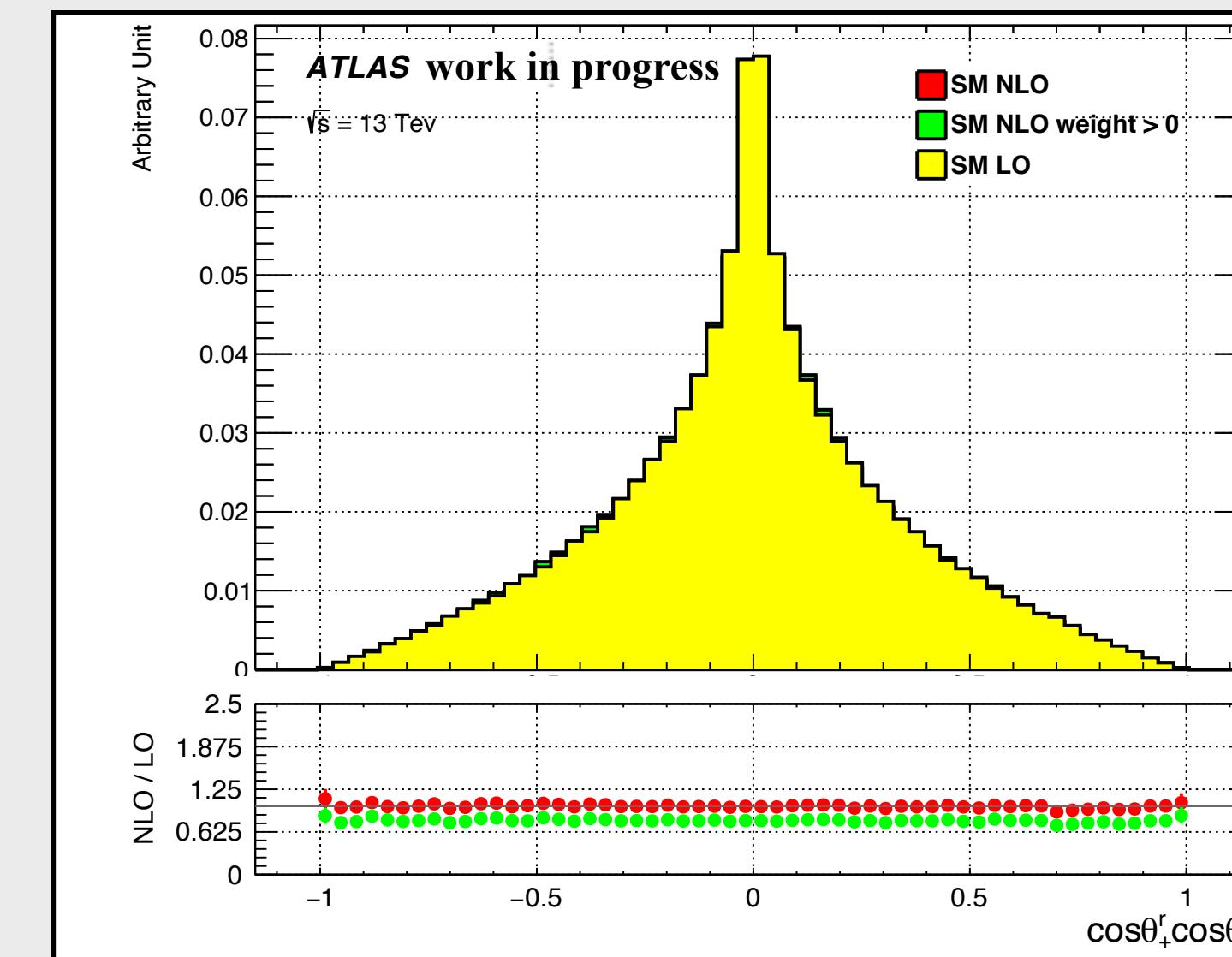
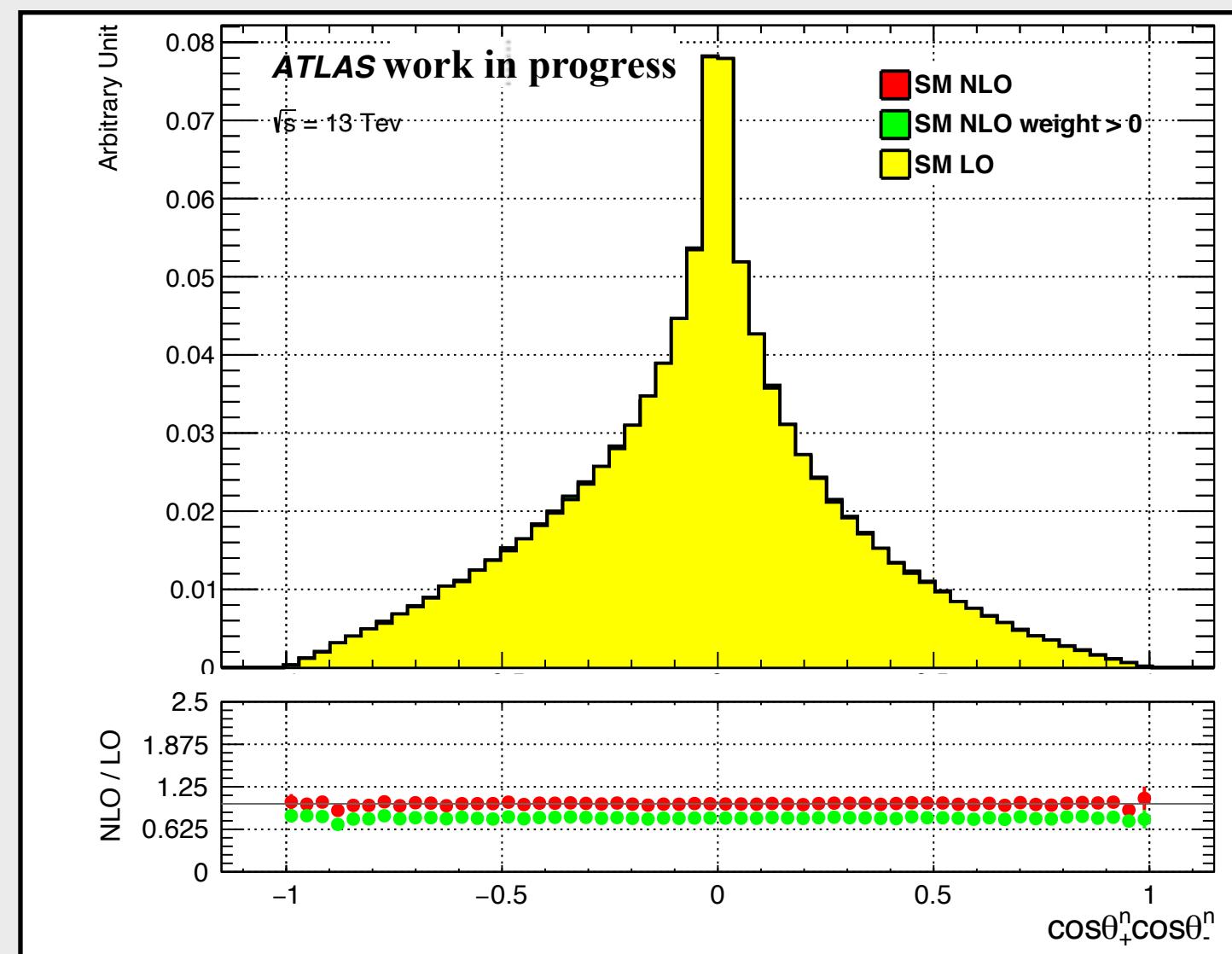
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Spin Correlations

Distributions for the correlation of top spins along each axis (probing diagonal of C (a, b) matrix)



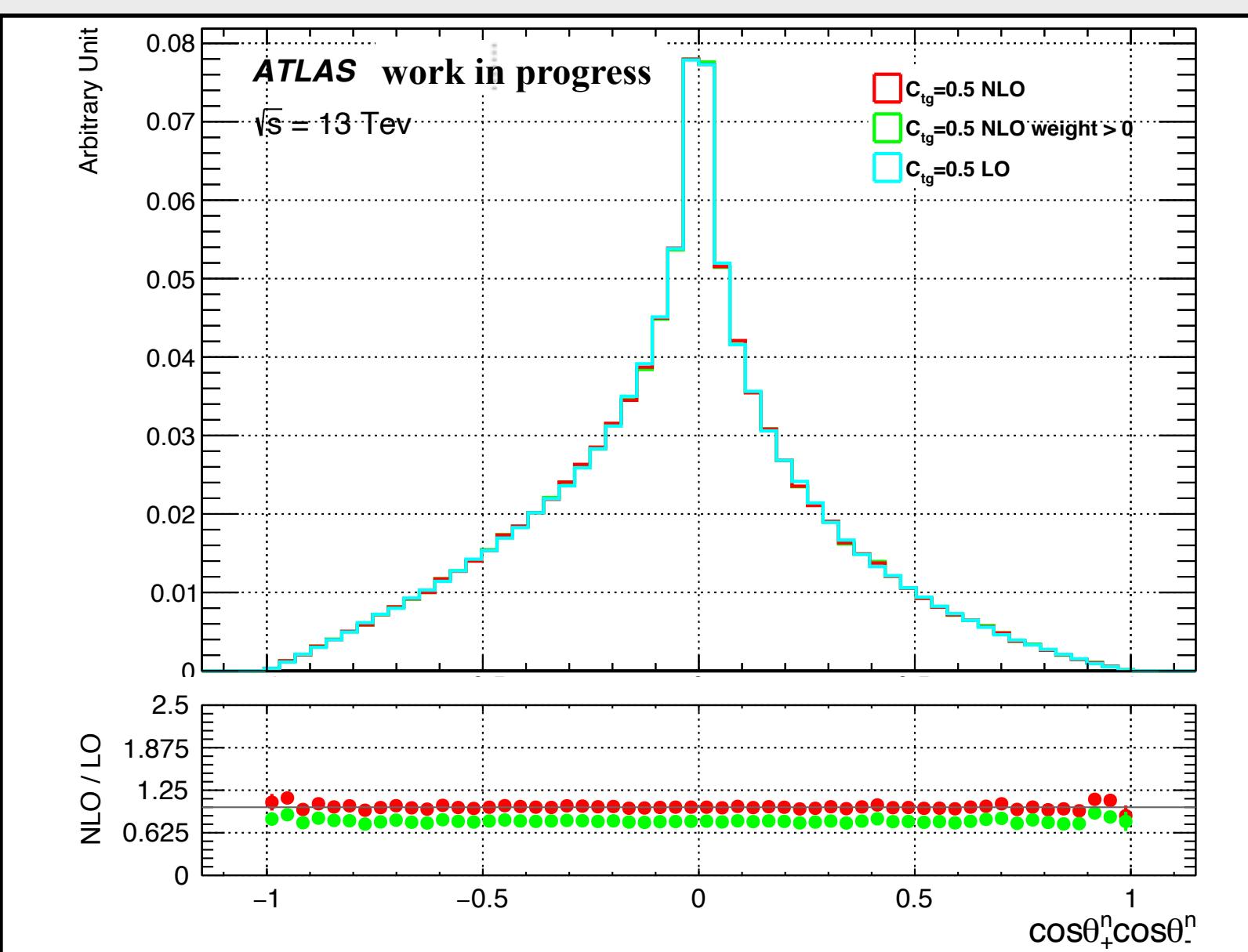
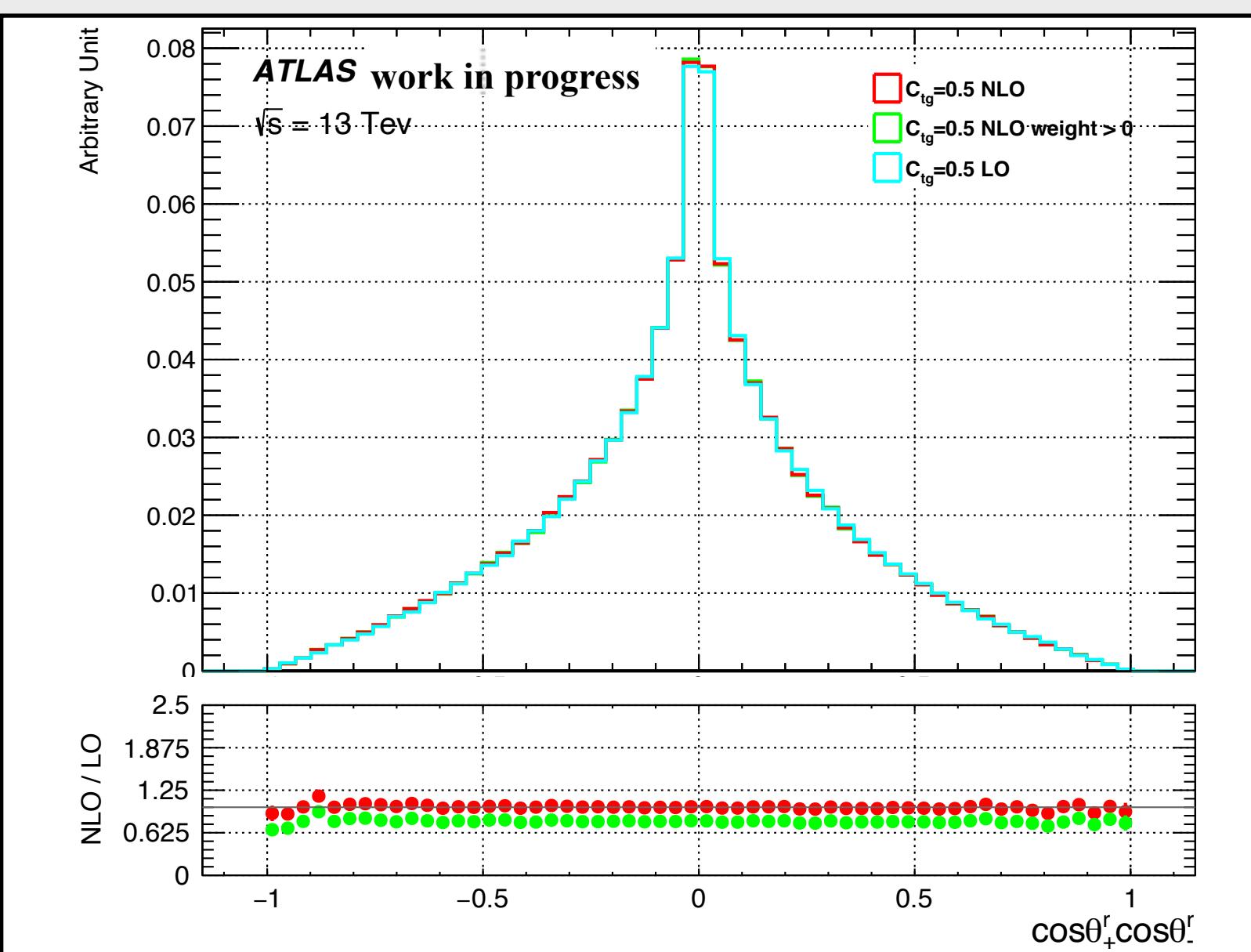
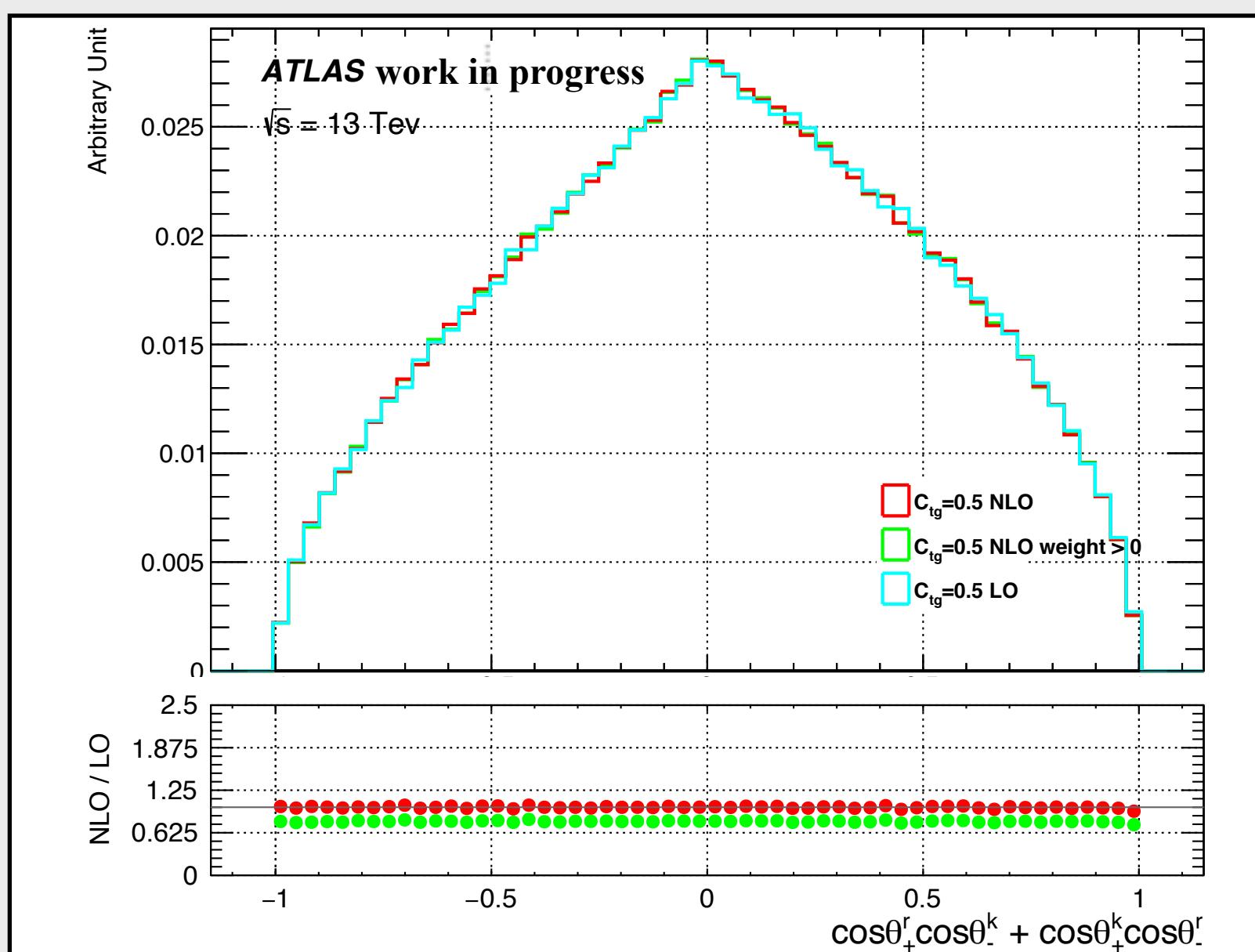
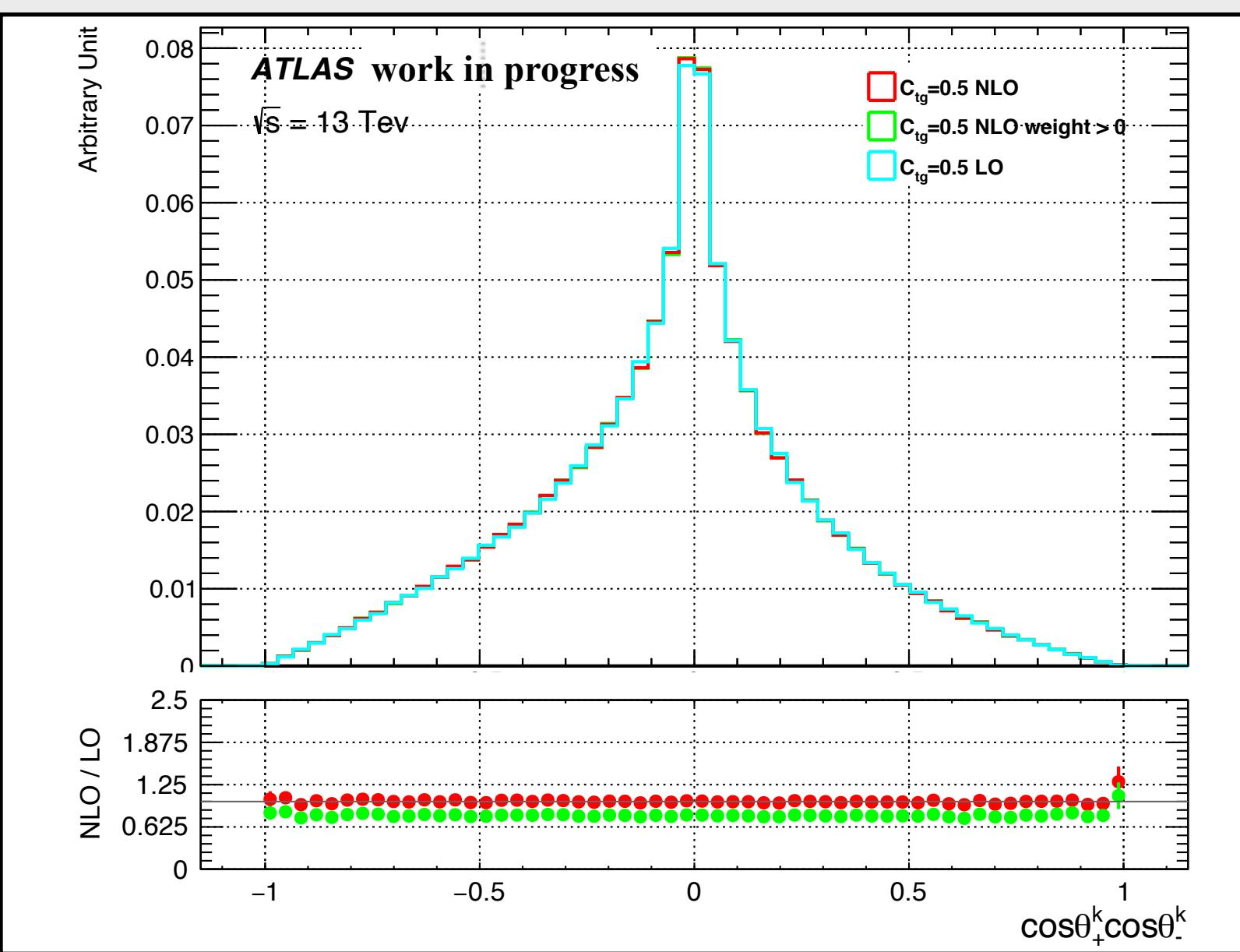
Spin correlations along each axis consistent with SM expectations (NLO from [1508.05271](#))

Spin Correlation Vs Ctg

Spin correlation :

- The impact of C_{tg} at NLO/LO is low, except $C(r,k) + C(k,r)$.

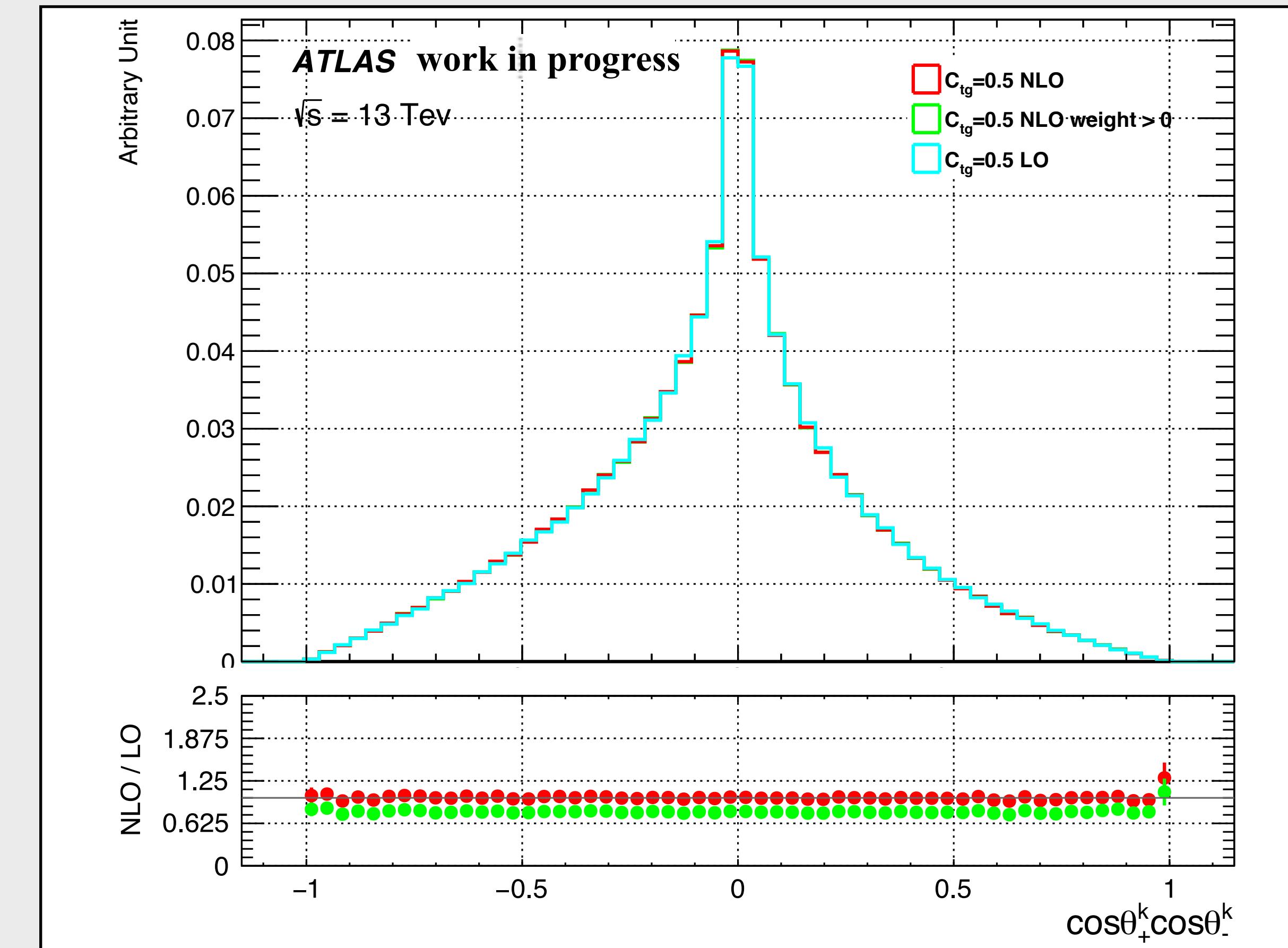
- For other spin corrections observables, the effect is very low or not observed.



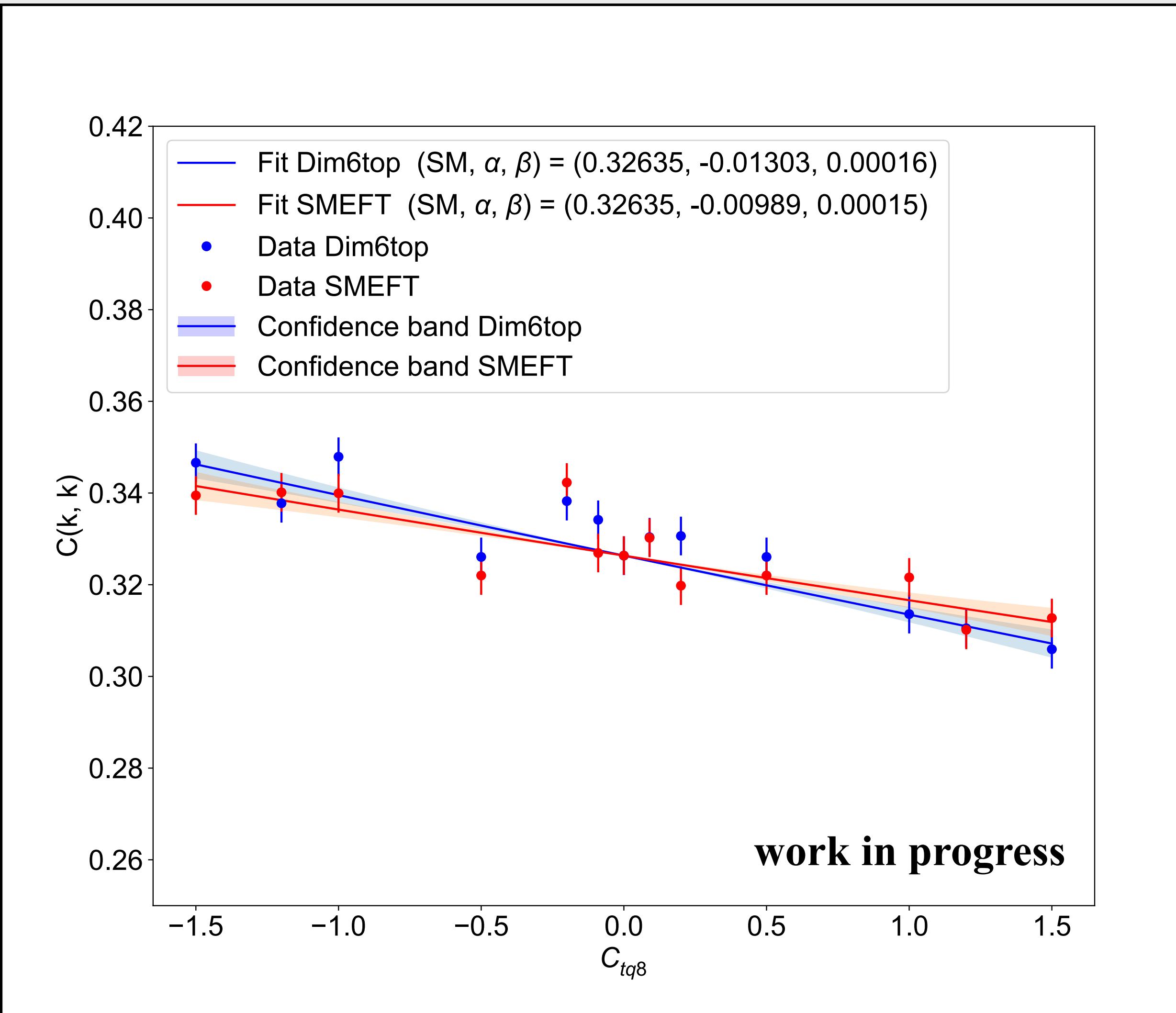
Spin Correlation Vs C_{tg}

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α_i/Λ^2 and β_i/Λ^4 at LO : $C(k,k)$



Dim6Top and SMEFT model are used to generate MC sample

Which Wilson coefficients affects $t\bar{t}$ production the most ?

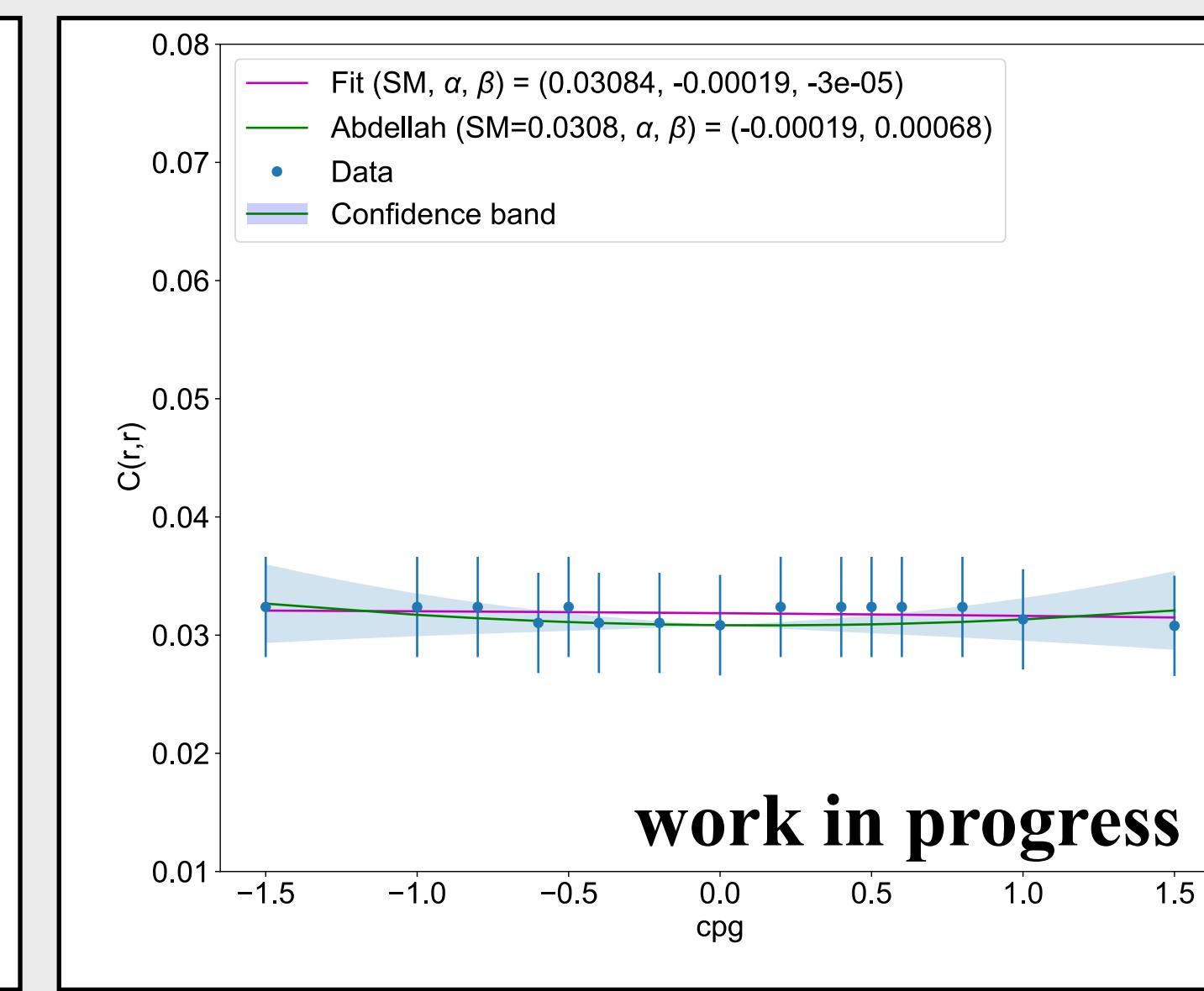
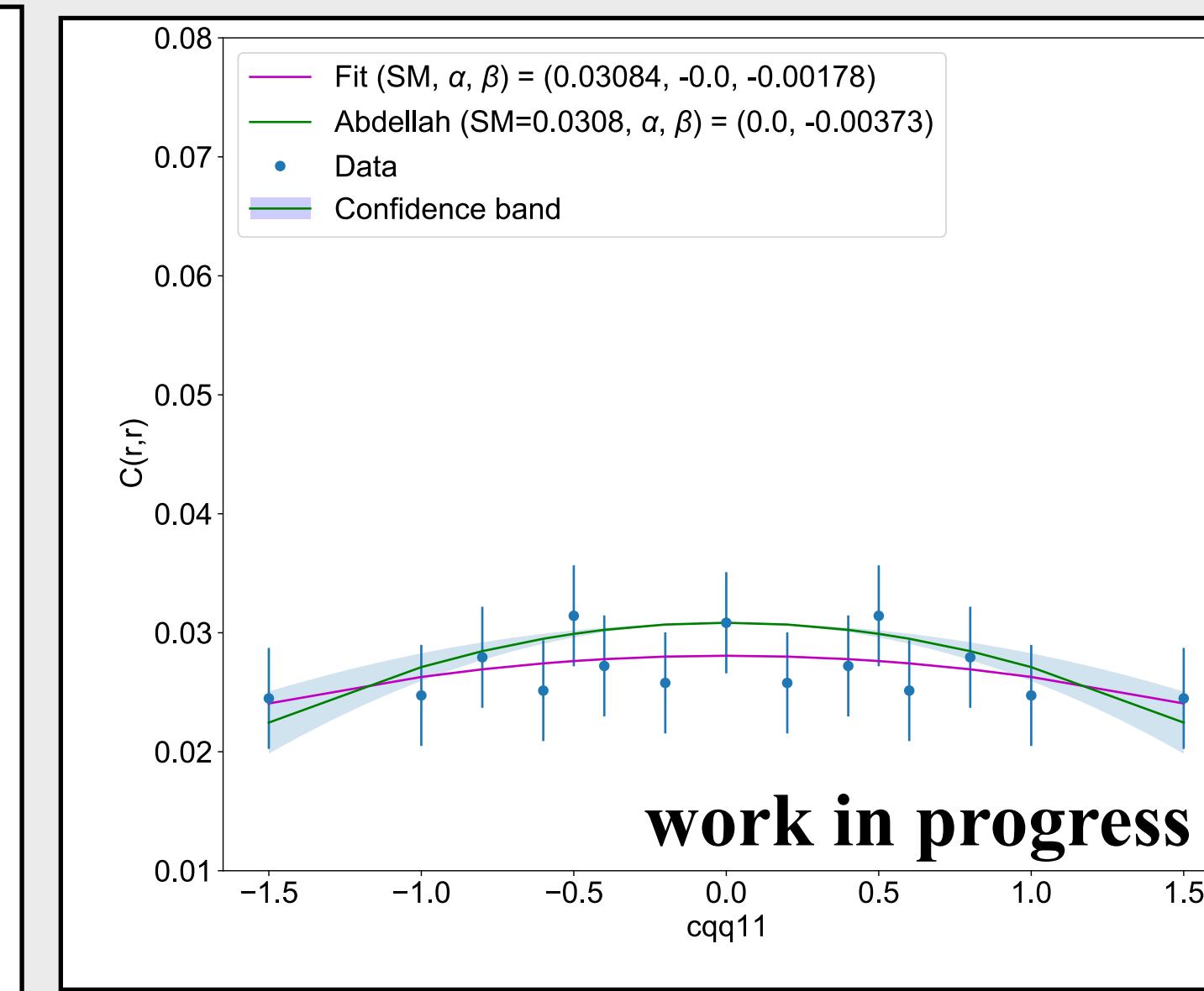
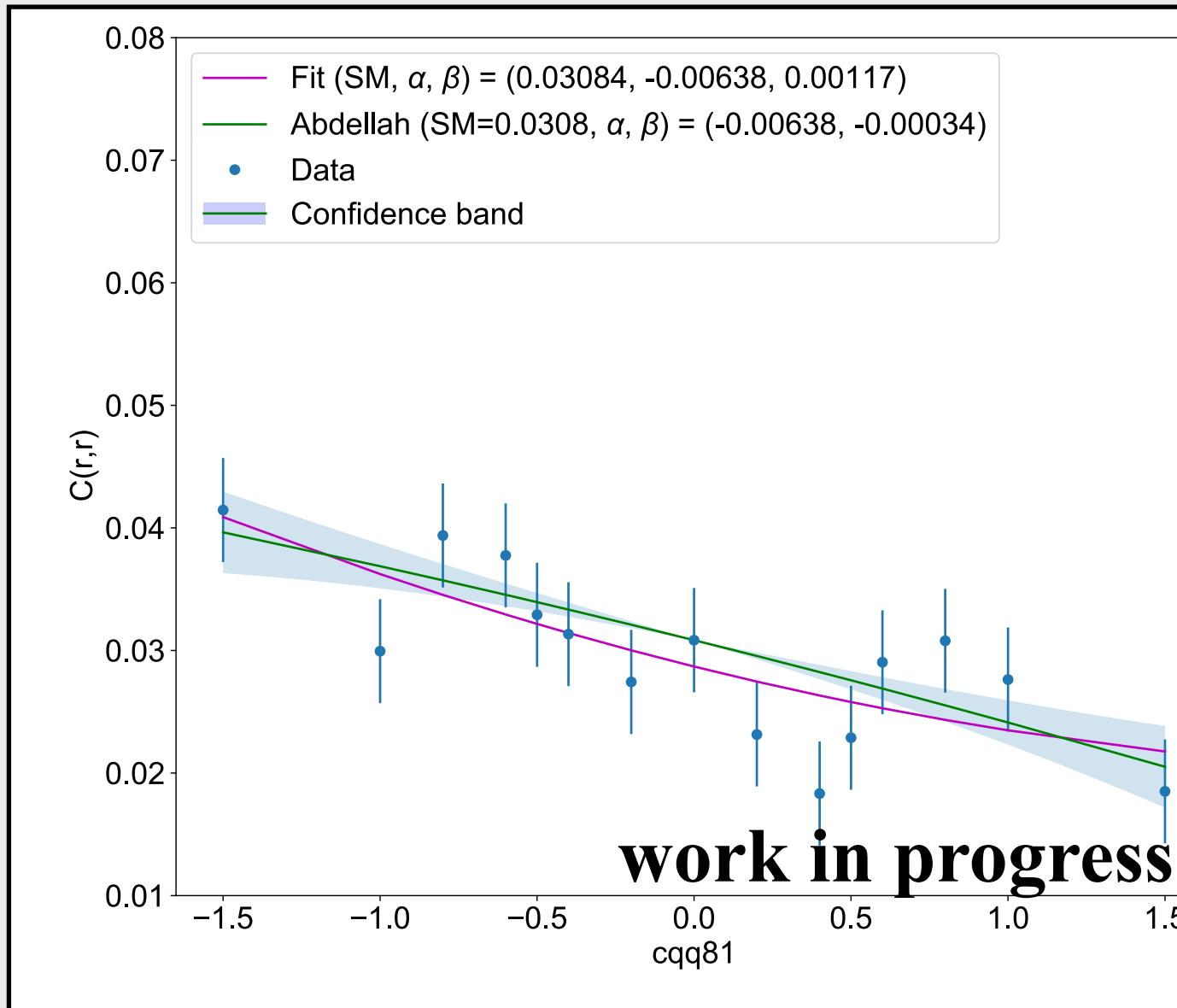
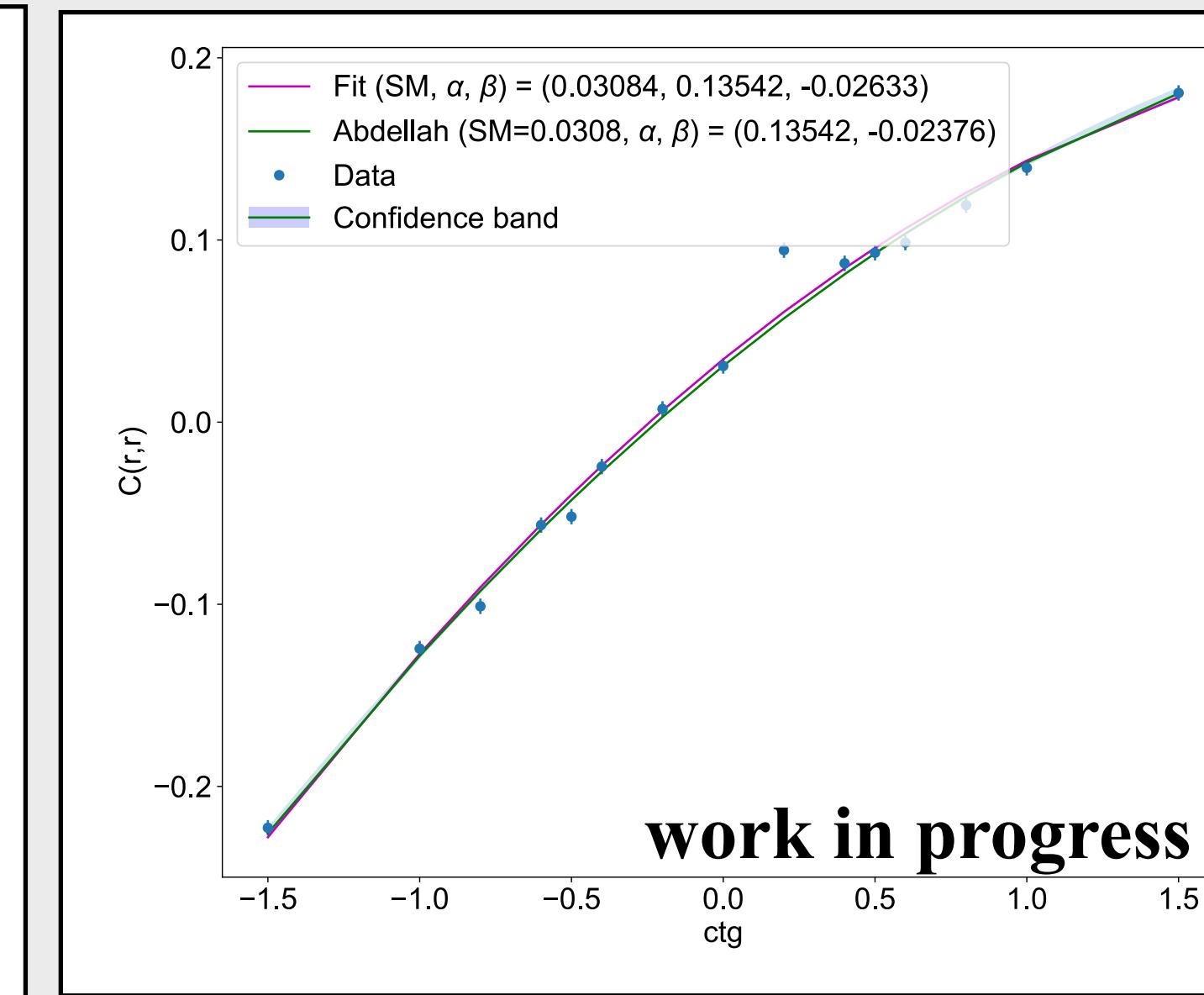
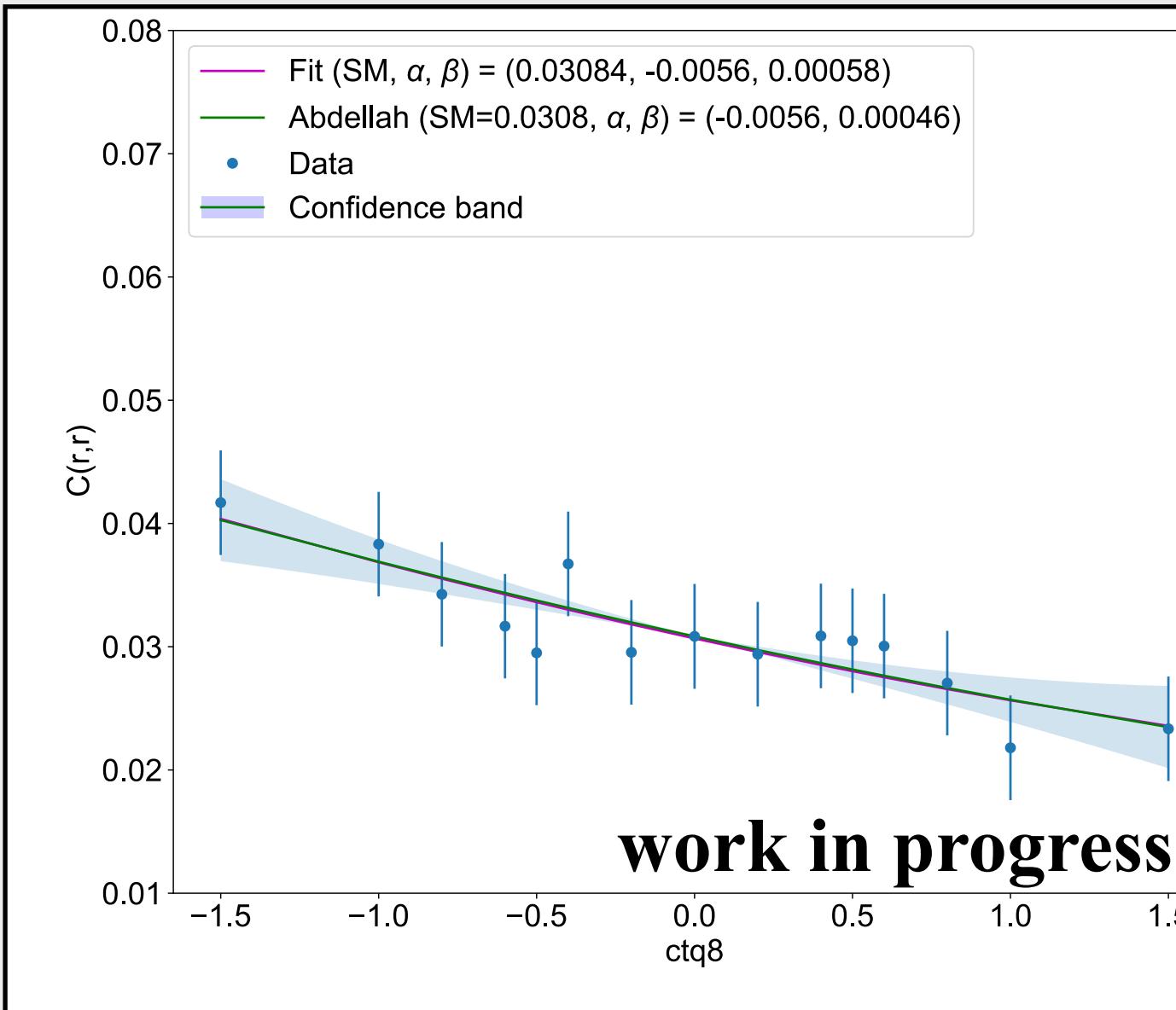
- 18 operator expect to affect $t\bar{t}$ process :

- 4-quark (2-heavy, 2-light) operator
- Heavy quark boson

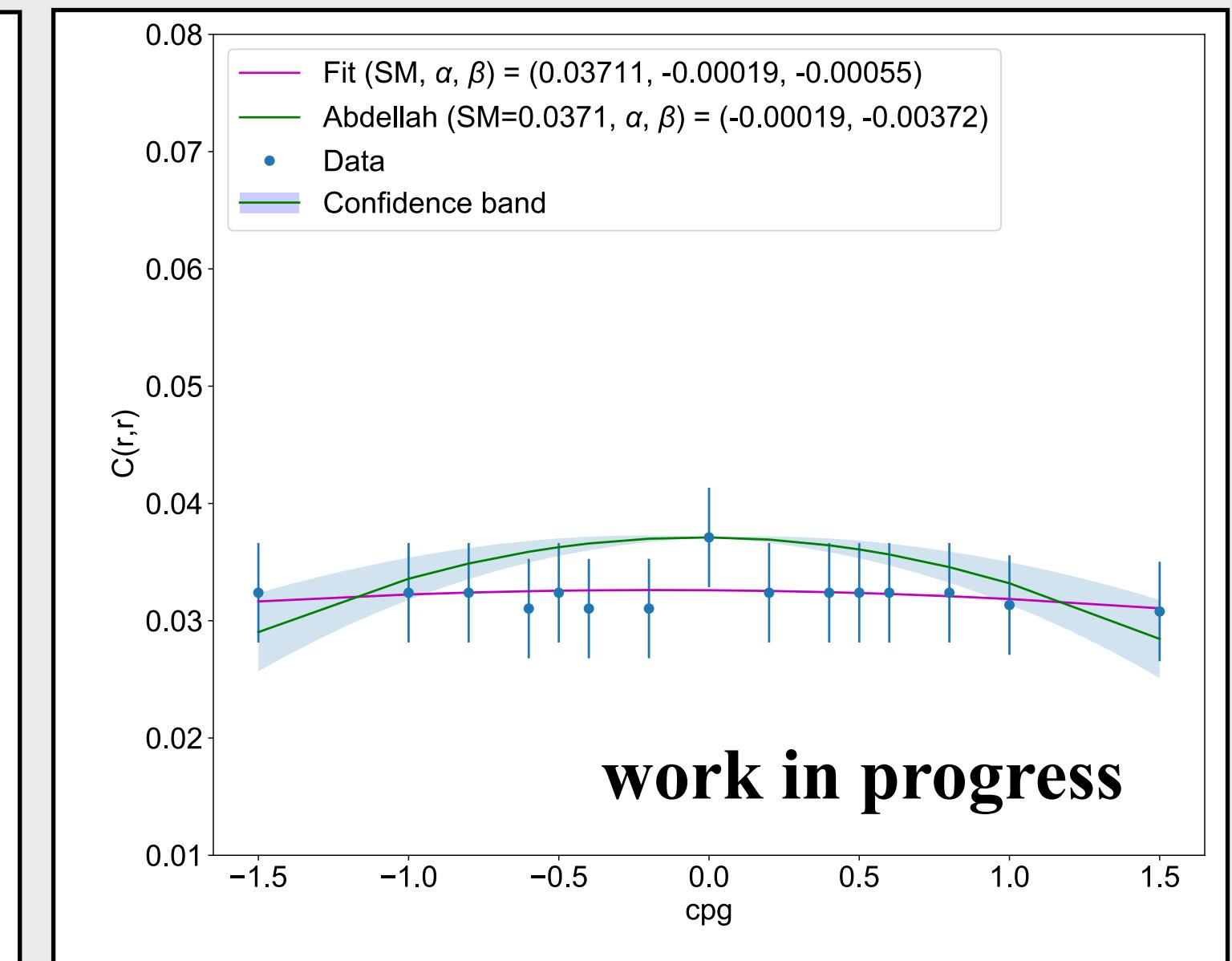
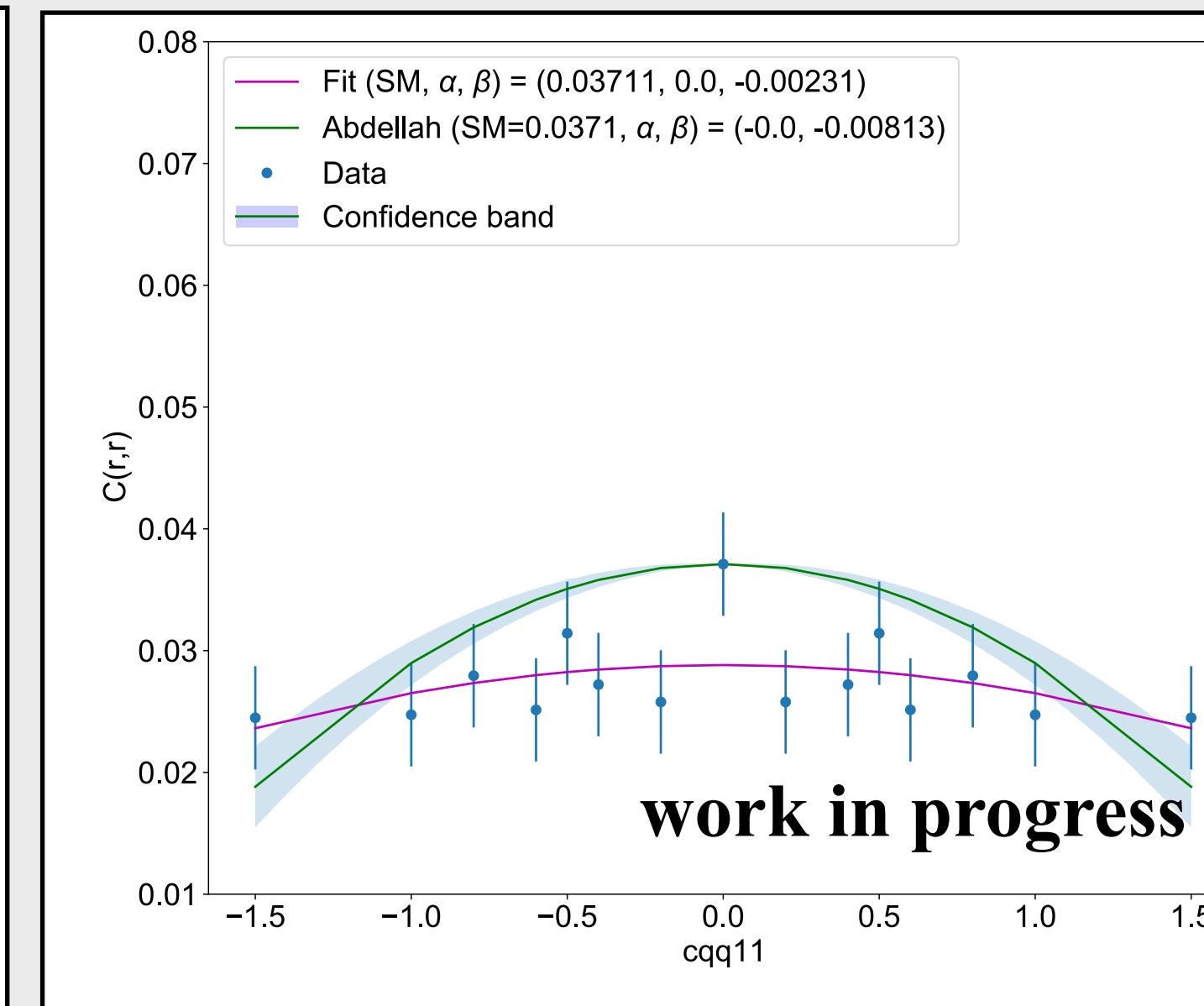
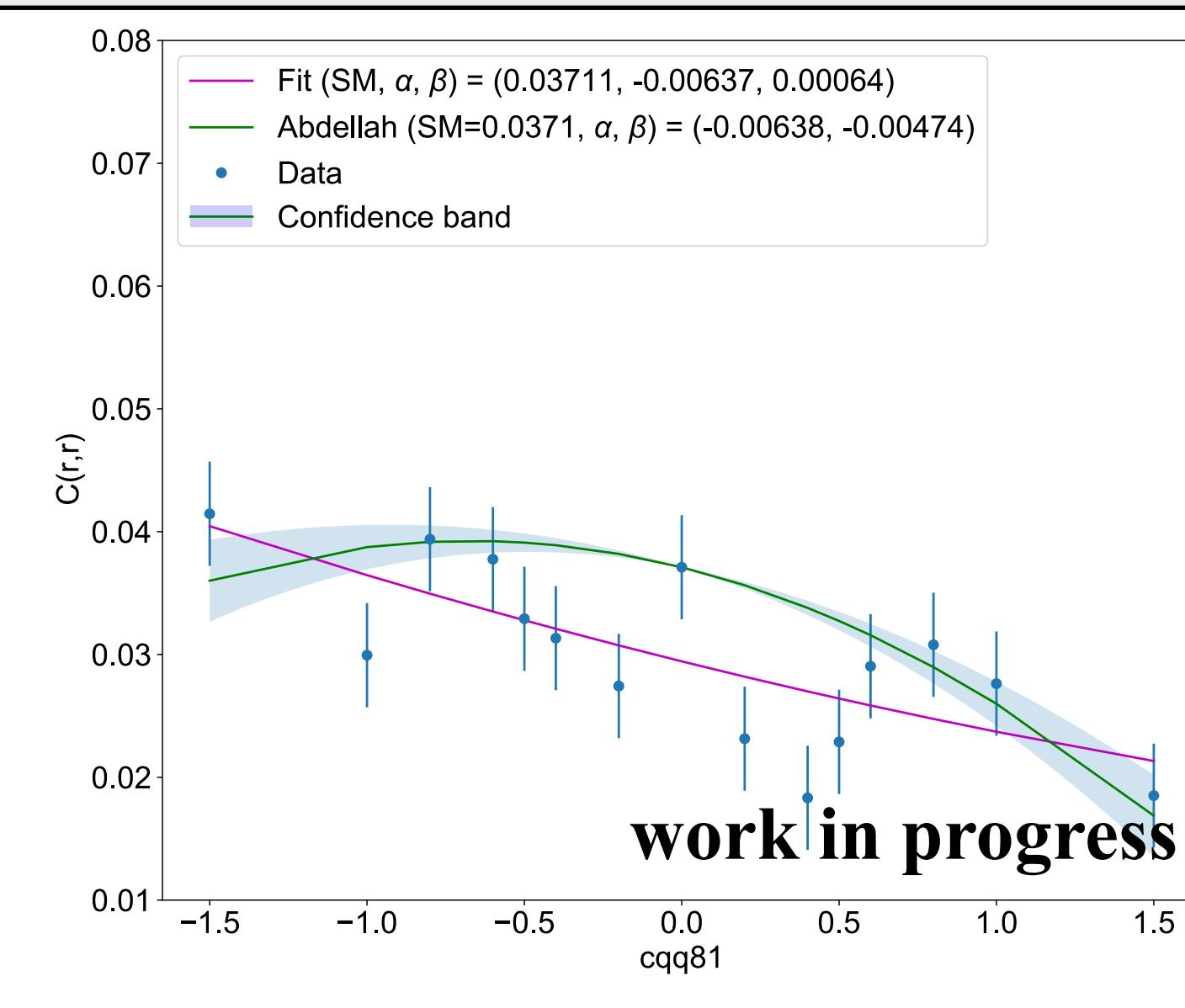
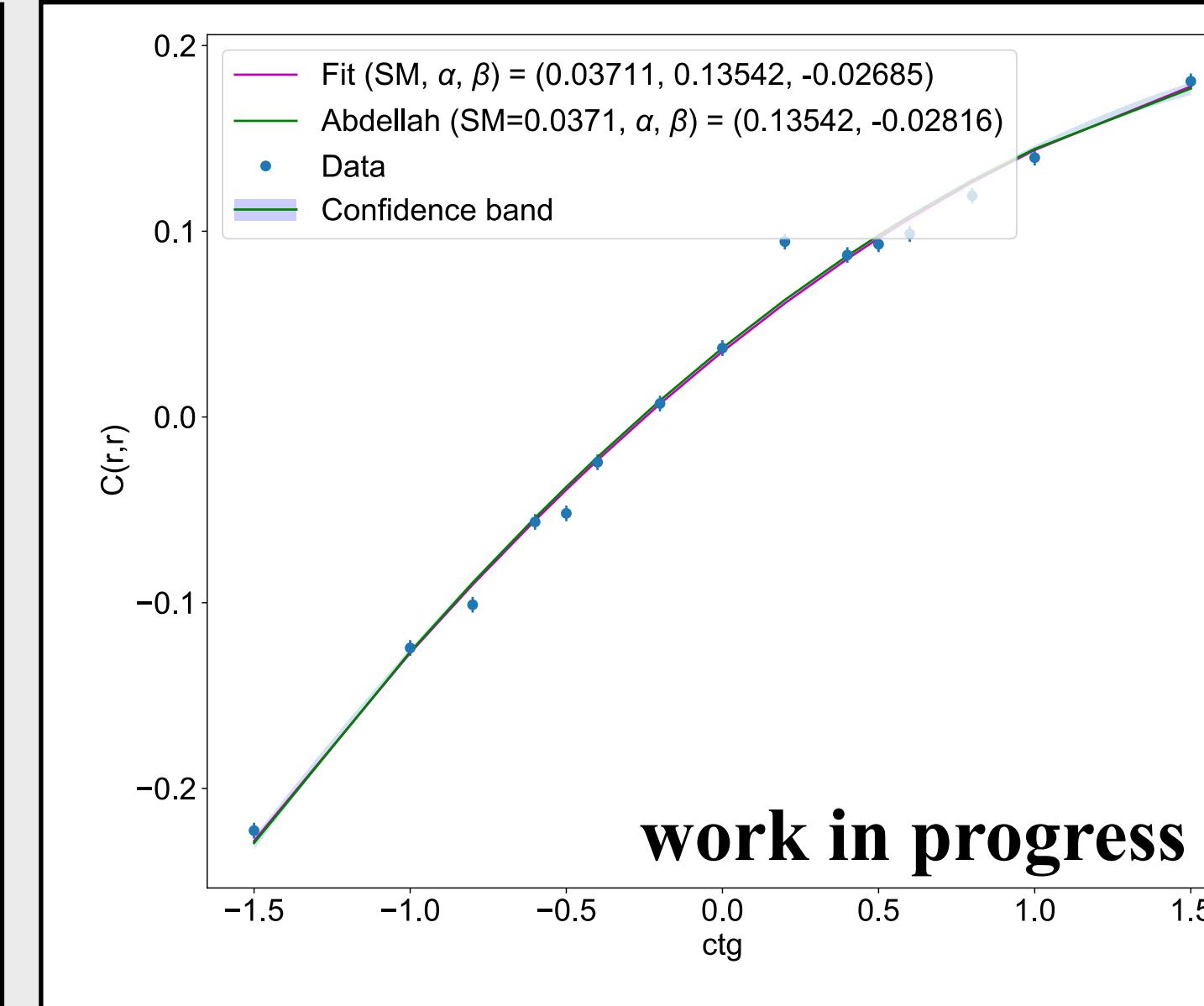
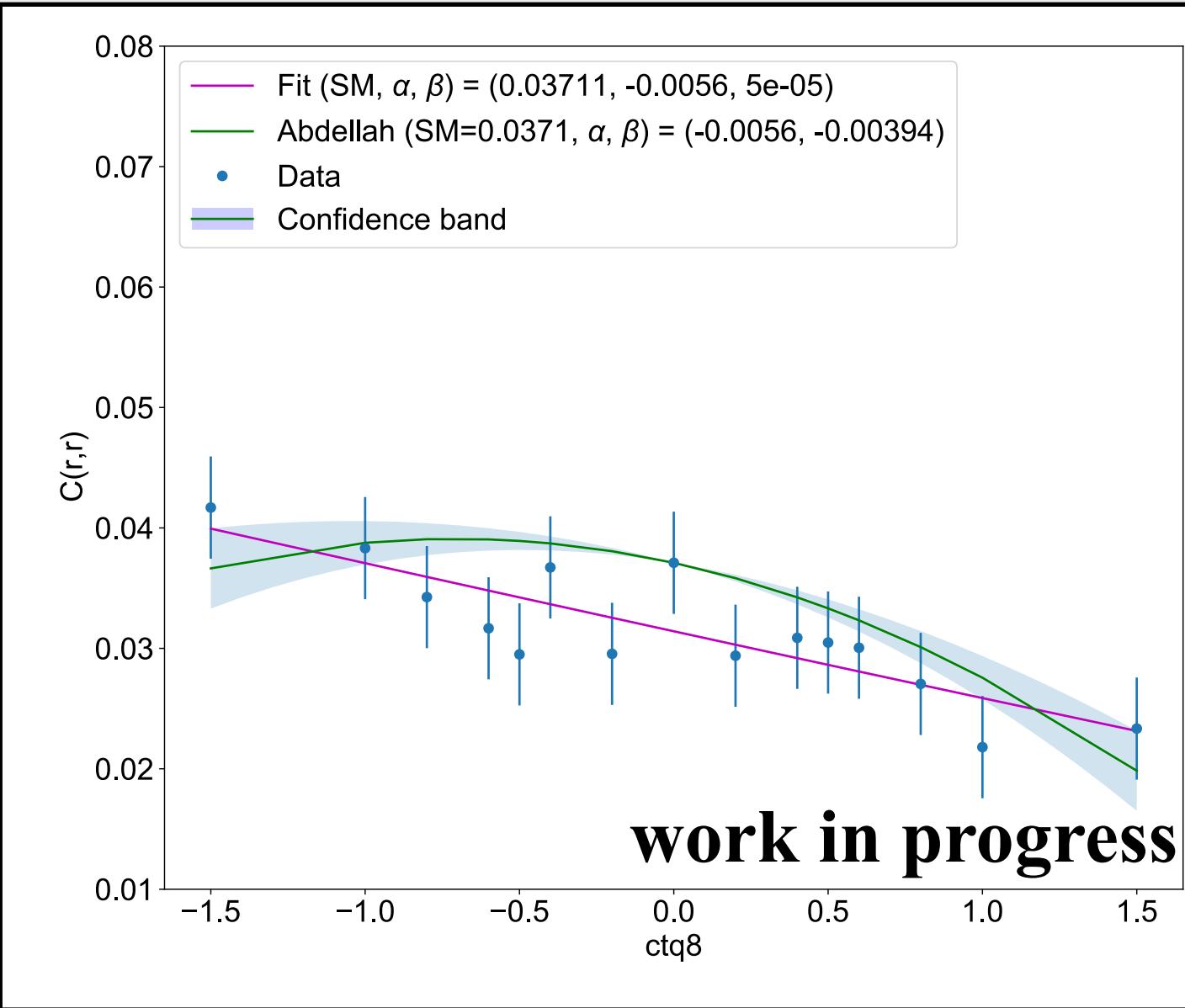
- We Can not prob gluon self-coupling cG in dim6top or SMEFT@NLO

parameter	$t\bar{t}$	single t
$C_{Qq}^{1,8}$	Λ^{-2}	—
$C_{Qq}^{3,8}$	Λ^{-2}	$\Lambda^{-4} [\Lambda^{-2}]$
C_{tu}^8, C_{td}^8	Λ^{-2}	—
$C_{Qq}^{1,1}$	$\Lambda^{-4} [\Lambda^{-2}]$	—
$C_{Qq}^{3,1}$	$\Lambda^{-4} [\Lambda^{-2}]$	Λ^{-2}
C_{tu}^1, C_{td}^1	$\Lambda^{-4} [\Lambda^{-2}]$	—
C_{Qu}^8, C_{Qd}^8	Λ^{-2}	—
C_{tq}^8	Λ^{-2}	—
C_{Qu}^1, C_{Qd}^1	$\Lambda^{-4} [\Lambda^{-2}]$	—
C_{tq}^1	$\Lambda^{-4} [\Lambda^{-2}]$	—
$C_{\phi Q}^-$	—	—
$C_{\phi Q}^3$	—	Λ^{-2}
$C_{\phi t}$	—	—
$C_{\phi tb}$	—	Λ^{-4}
C_{tZ}	—	—
C_{tW}	—	Λ^{-2}
C_{bW}	—	Λ^{-4}
C_{tG}	Λ^{-2}	$[\Lambda^{-2}]$

α_i/Λ^2 and β_i/Λ^4 at LO : $C(r,r)$



α_i/Λ^2 and β_i/Λ^4 at LO : $C(r,r)$



α_i/Λ^2 and β_i/Λ^4 at LO : $C(k,k)$

