

Université de Mons

Υ production in dAu and AuAu collisions at RHIC

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In collaboration with E. G. Ferreiro, F. Fleuret, J. P. Lansberg and A.

Rakotozafindrabe

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↑ production in dAu and AuAu collisions

Outline

- Introduction and motivations
- 2 Experimental situation
- On the kinematics of Υ production
- The Glauber Monte Carlo
- 5 Results for dAu collisions
- 6 Results for AuAu collisions
- Conclusions and perspectives

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- Comparison of three differents shadowing parametrisations

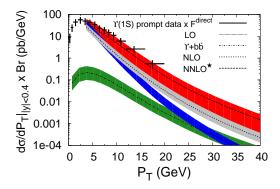
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- Glauber Monte Carlo model to simulate dAu and AuAu collisions at RHIC
- Two main production schemes $(1 \rightarrow 2, 2 \rightarrow 2)$
- Comparison of three differents shadowing parametrisations
- Three absorption cross sections

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Experimental situation

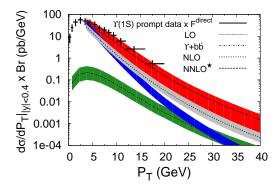
Results at 1.8 TeV



CSM describes well the data at NNLO*

Experimental situation

Results at 1.8 TeV

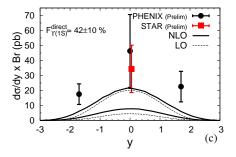


- CSM describes well the data at NNLO*
- However LO CSM is sufficient to describe low pT data

Experimental situation

S. J. Brodsky and J. P. Lansberg, arXiv:0908.0754 [hep-ph].

Results at 200 GeV



- Upper dashed line, $m_{\Upsilon} = 4.5$ GeV,
- Lower dashed line, $m_{\Upsilon} = 5.0$ GeV,

 $m_{\Upsilon} = 4.5$ GeV, taken in the following plots

If $\mathcal{F}_{g}^{A}(x, \vec{r}, z, \mu_{f})$ gives the distribution of a gluon of mom. fract. x at a position \vec{r}, z in a nucleus A, the differential cross-section reads:

$$\frac{d\sigma_{AB}}{dy \, dP_T \, d\vec{b}} =$$

 $2 \rightarrow 1$ kinematics with intrinsic $\textit{p}_{\mathcal{T}}$

 $2 \rightarrow 2$ kinematics with extrinsic $\textit{p}_{\mathcal{T}}$

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 $\mathbf{2}
ightarrow \mathbf{1}$ kinematics with intrinsic p_T

 $\int d\vec{r}_A dz_A dz_B$ $\times \mathcal{F}_g^A(x_1^0, \vec{r}_A, z_A, \mu_f) \mathcal{F}_g^B(x_2^0, \vec{r}_B, z_B, \mu_f)$ $\times \sigma_{gg}^{\text{Intr.}}(x_1^0, x_2^0)$ $\times S_A(\vec{r}_A, z_A) S_B(\vec{r}_B, z_B)$ $2 \rightarrow 2$ kinematics with extrinsic p_{T}

$$\int dx_1 dx_2 \int d\vec{r}_A dz_A dz_B \times \mathcal{F}_g^A(x_1, \vec{r}_A, z_A, \mu_f) \mathcal{F}_g^B(x_2, \vec{r}_B, z_B, \mu_f) \times 2\hat{s} P_T \frac{d\sigma_{gg \rightarrow \Upsilon + g}}{d\hat{t}} \delta(\hat{s} - \hat{t} - \hat{u} - M^2) \times S_A(\vec{r}, z_A) S_B(\vec{r}_B, z_B)$$

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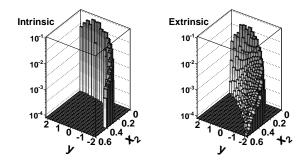
 $\int d\vec{r}_A dz_A dz_B$ $\times \mathcal{F}_g^A(x_1^0, \vec{r}_A, z_A, \mu_f) \mathcal{F}_g^B(x_2^0, \vec{r}_B, z_B, \mu_f)$ $\times \sigma_{gg}^{\text{Intr.}}(x_1^0, x_2^0)$ $\times S_A(\vec{r}_A, z_A) S_B(\vec{r}_B, z_B)$

$$x_{1,2} = \frac{m_T}{\sqrt{s_{NN}}} \exp\left(\pm y\right) \equiv x_{1,2}^0(y, P_T)$$

 $\mathbf{2} \rightarrow \mathbf{2}$ kinematics with extrinsic p_T

$$\int dx_1 dx_2 \int d\vec{r}_A dz_A dz_B \times \mathcal{F}_g^A(x_1, \vec{r}_A, z_A, \mu_f) \mathcal{F}_g^B(x_2, \vec{r}_B, z_B, \mu_f) \times 2\hat{s} P_T \frac{d\sigma_{gg \rightarrow \Upsilon + g}}{d\hat{t}} \delta(\hat{s} - \hat{t} - \hat{u} - M^2) \times S_A(\vec{r}, z_A) S_B(\vec{r}_B, z_B)$$

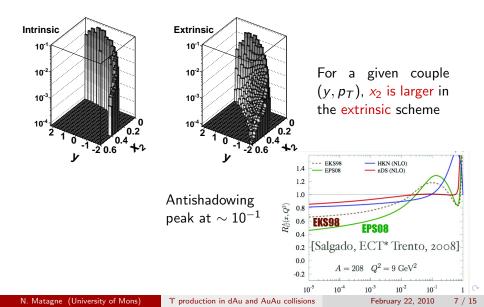
$$\delta(..) \rightarrow x_2 = \frac{x_1 m_T \sqrt{s_{NN}} e^{-y} - M^2}{\sqrt{s_{NN}} (\sqrt{s_{NN}} x_1 - m_T e^y)}$$

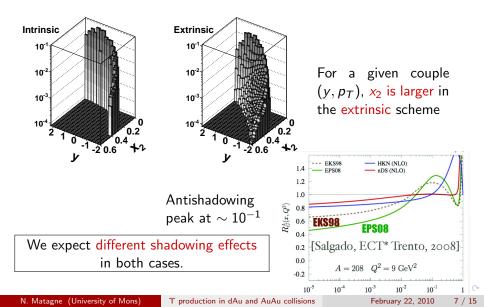


For a given couple (y, p_T) , x_2 is larger in the extrinsic scheme

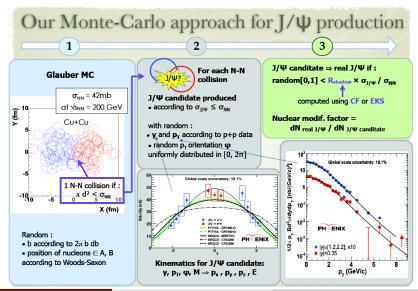
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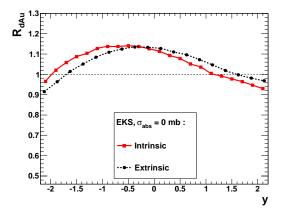
The Glauber Monte Carlo



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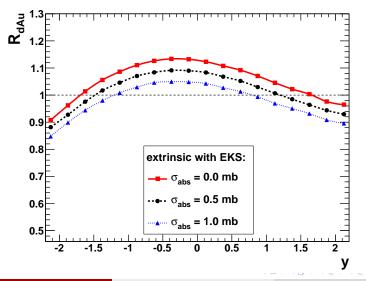
↑ production in dAu and AuAu collisions

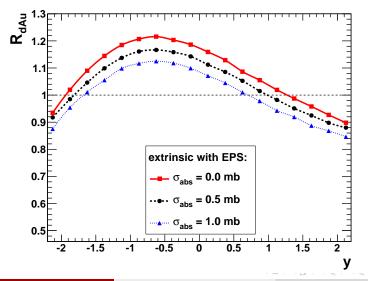
Illustration of the differences between intrinsic and extrinsic p_T

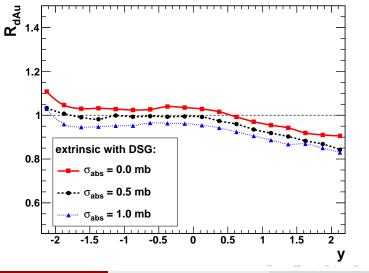


See Andry's talk

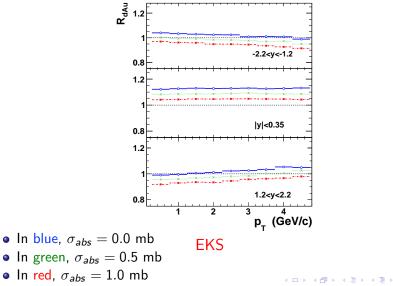
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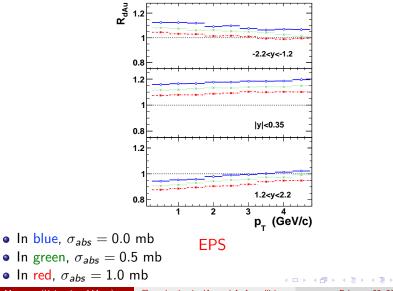


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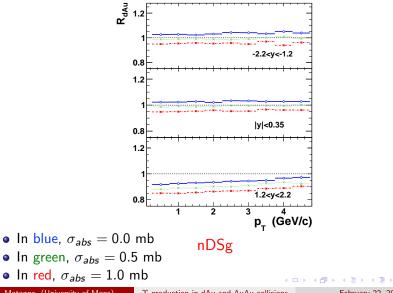
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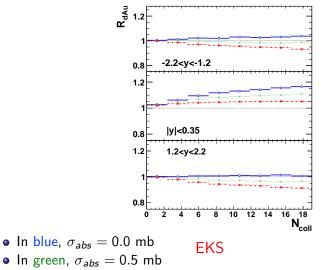
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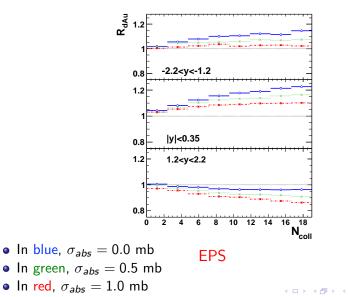
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• In red,
$$\sigma_{abs} = 1.0 \text{ mb}$$

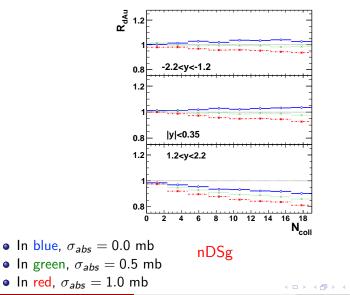
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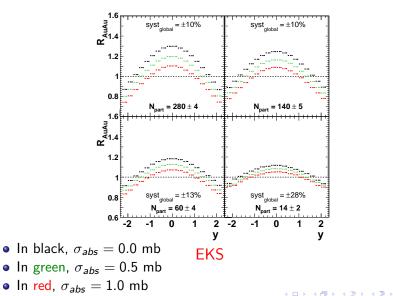
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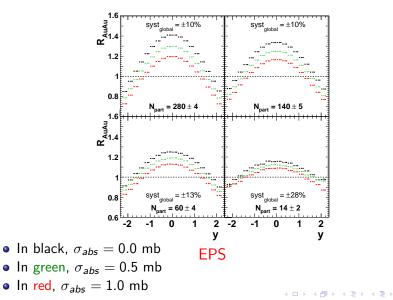


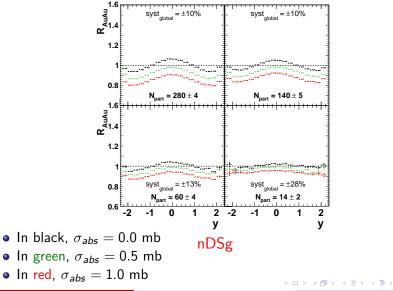
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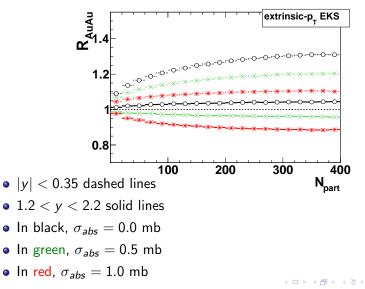


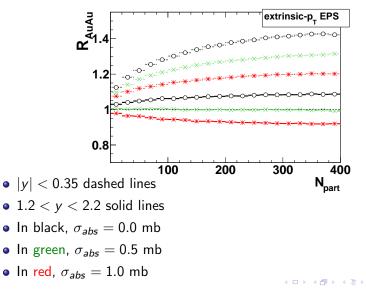
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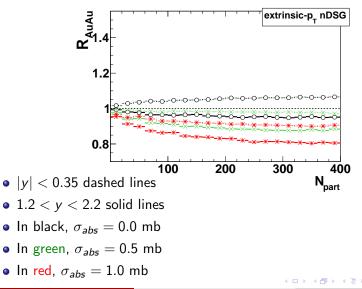




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- Only very forward (backward) region shows $R_{dAU} < 1$ due to shadowing

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- Production of Upsilon sits in the antishadowing region at RHIC energy
- Only very forward (backward) region shows $R_{dAU} < 1$ due to shadowing
- \bullet Within the commonly accepted $\sigma_{\textit{abs}},$ one should expect an excess of Υ
- ... unless there is no antishadowing (see nDSg)

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