

# Gluon shadowing and $J/\psi$ suppression in p A collisions: uncertainties and constraints

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- **Motivations**
  - final state interaction vs. gluon nPDF
- **World data analysis**
  - impact of nPDF uncertainties
- **Constraints on gluon nPDF from  $J/\psi$  data**
  - specific analysis of E866/NuSea

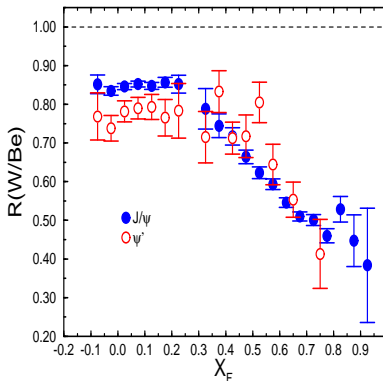
## References

- FA, V.-N. Tram, Eur. Phys. J. C 55 (2008) 449
- V.-N. Tram, FA, Eur. Phys. J. C 61 (2009) 847
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# $J/\psi$ suppression in nuclear collisions

## Experimentally

$$R \equiv \frac{\sigma(p \text{ } A \rightarrow J/\psi \text{ } X)}{A \sigma(p \text{ } p \rightarrow J/\psi \text{ } X)} < 1$$



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Why is  $R \neq 1$ ?

## Two usual suspects

- Inelastic interaction with nuclear matter
- Effects of nuclear parton densities (nPDF)

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## Two usual suspects

- Inelastic interaction with nuclear matter
  - Effects of nuclear parton densities (nPDF)
- Need to know nPDF to understand  $J/\psi$  suppression data
  - Possibility to use  $J/\psi$  suppression data to constrain nPDF

## Aim of this analysis

To determine **quantitatively** the  $J/\psi$  inelastic cross section from a  $\chi^2$  analysis of **all available data** on  $J/\psi$  production on nuclei

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### Data sets

- Inelastic hadroproduction  $h A \rightarrow J/\psi X$
- Inelastic photoproduction  $\gamma A \rightarrow J/\psi X$
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- Inelastic hadroproduction  $h A \rightarrow J/\psi X$
- Inelastic photoproduction  $\gamma A \rightarrow J/\psi X$
- Inelastic leptonproduction  $\gamma^* A \rightarrow J/\psi X$

## Nuclear effects considered

- $J/\psi$  inelastic interaction in nuclear matter
- Nuclear modifications of parton distribution functions



# Production cross section

## Hadroproduction

Colour Evaporation Model is assumed

$$\frac{d\sigma^{\text{no abs}}}{dx_F} \propto \int_{2m_c}^{2m_D} dm \frac{2m}{\sqrt{x_F^2 s + 4m^2}} \left[ f_g^h(x_1) f_g^A(x_2) \hat{\sigma}_{gg \rightarrow c\bar{c}} + \sum_{q=u,d,s} \left\{ f_q^h(x_1) f_q^A(x_2) + f_{\bar{q}}^h(x_1) f_{\bar{q}}^A(x_2) \right\} \hat{\sigma}_{q\bar{q} \rightarrow c\bar{c}} \right]$$

- $f_i^h$ : parton distributions functions in the hadron
- $\hat{\sigma}$ : LO partonic cross sections
- $x_1/x_2$ : projectile / target parton momentum fractions

# Production cross section

## Hadroproduction

Colour Evaporation Model is assumed

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$$\frac{d\sigma^{\text{no abs}}}{dx} \propto f_g^A(x)$$

# Production cross section

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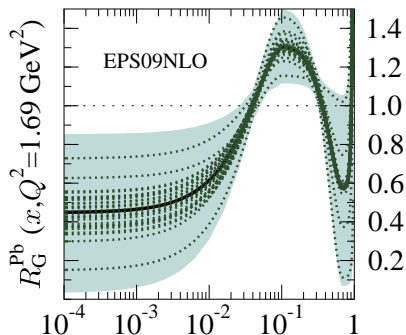
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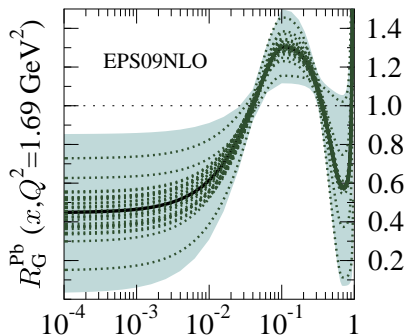
What should be taken for the nuclear densities  $f_i^A$  ?

PDFs in nuclei differ significantly from those in a proton



This should affect the nuclear dependence of  $J/\psi$  production

PDFs in nuclei differ significantly from those in a proton



Issue: large uncertainties

- in the gluon sector
- at small  $x$  and low scales

Several LO/NLO determinations of

$$f_i^A(x, Q^2) / f_i^P(x, Q^2) \quad (i = q, \bar{q}, g)$$

from DIS and DY (and hadron prod. for EPS08) data can be used

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- Eskola–Kolhinen–Salgado EKS98 LO (1998)
- Hirai–Kumano–Miyama HKM LO (2001)
- Eskola–Paukkunen–Salgado EPS08 LO (2008)
- de Florian–Sassot nDS / nDSg NLO (2003)
- Hirai–Kumano–Nagai HKN NLO (2007)
- Eskola–Paukkunen–Salgado EPS09 NLO (2009)
- I. Schienbein et al. nCTEQ NLO (2009)

Differences come from the various data sets and fitting procedures

# Nuclear absorption

Assume factorization between production and rescattering

$$\frac{d\sigma}{dx} = S^{\text{abs}}(A, \sigma_{J/\psi N}) \times \frac{d\sigma^{\text{no abs}}}{dx}$$

$S^{\text{abs}}$  survival probability of  $J/\psi$  in the target nucleus

$$S^{\text{abs}} = \frac{1}{(A-1) \sigma_{J/\psi N}} \int d\mathbf{b} \left( 1 - e^{-(1-1/A) T_A(\mathbf{b}) \sigma_{J/\psi N}} \right)$$

with  $T_A(\mathbf{b})$  thickness function



# Minimization

Cross section determined from the usual  $\chi^2$  minimization

$$\chi^2(\hat{\sigma}_{J/\psi N}) = \min \left[ \chi^2(\sigma_{J/\psi N}) \right]$$

with sophisticated definition of  $\chi^2$  to account for correlated errors

$1\sigma$  error  $\delta\hat{\sigma}_{J/\psi N}$  defined as

$$\Delta\chi^2 \equiv \chi^2(\hat{\sigma}_{J/\psi N} \pm \delta\hat{\sigma}_{J/\psi N}) - \chi^2_{\min} = 1$$

No longer true when correlated errors not properly taken into account !

[ Stump et al. 2002 ]

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Rescaling of the 1 $\sigma$  error (PDG prescription)

$$\delta\bar{\sigma}_{J/\psi N} = S \times \delta\sigma_{J/\psi N}$$

$$S \equiv \left[ \chi^2 / (n - 1) \right]^{1/2} \quad \text{if} \quad \chi^2/\text{ndf} > 1$$

- Data sets

- SPS: NA3, NA38, NA50
- FNAL: E537, E866, E672
- RHIC: PHENIX
- DIS : NMC

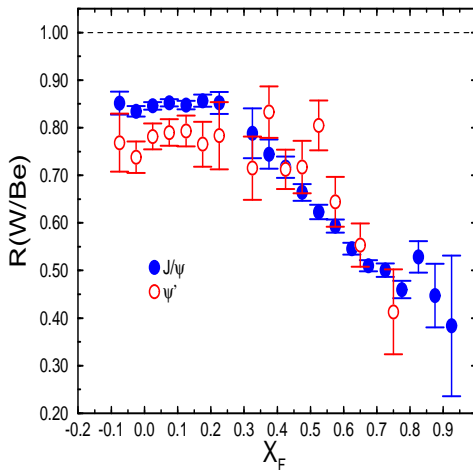
- $J/\psi$  – N energy range  $\sqrt{s_{J/\psi N}} \simeq 6.5 - 25$  GeV

- Kinematic cuts

$J/\psi$  suppression not really understood at large  $x_F$  [→]

$x_F \leq 0.3$  cut imposed to all hadroproduction data

- E772 data rejected because of a  $x_F - p_\perp$  correlation bias
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[ E866/NuSea 1999 ]

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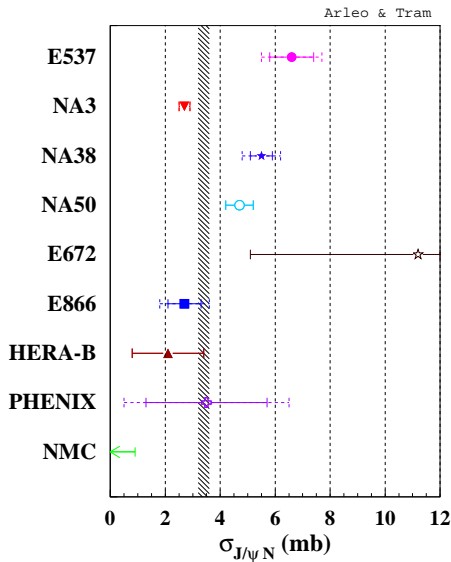
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# Results



## Main observations

- Clear **tension** between

- NA3, E866, HERA-B
- E537, NA38, NA50

$$\sigma_{J/\psi N} \simeq 1\text{--}3 \text{ mb}$$

$$\sigma_{J/\psi N} \simeq 4\text{--}7 \text{ mb}$$

- PHENIX data **do not constrain much**

- need more d Au statistics
- reduced systematic errors

$$\sigma_{J/\psi N} = 3.5 \pm 3.0 \text{ mb}$$

- Muoproduction data consistent with **no absorption**

$$\sigma_{J/\psi N} \leq 0.9 \text{ mb}$$

# Nuclear parton densities

Similar result using nDS parton densities

$$\sigma_{J/\psi N} = 3.5 \pm 0.2 \text{ mb} \quad (\chi^2/\text{ndf} = 1.4)$$

- Almost identical cross sections
- Same  $\chi^2$  than with proton densities



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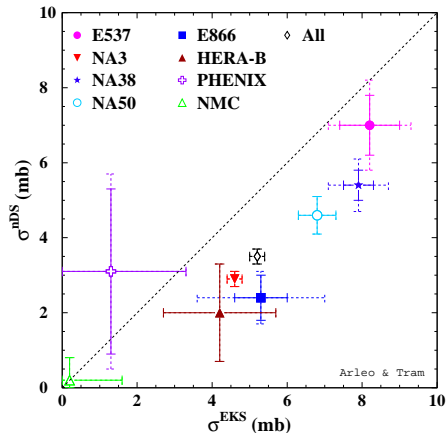
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## ... unlike using EKS98 parton densities

$$\sigma_{J/\psi N} = 5.2 \pm 0.2 \text{ mb} \quad (\chi^2/\text{ndf} = 1.6)$$

because of the stronger antishadowing

# Nuclear parton densities



$$\sigma_{J/\psi N}^{\text{EKS}} \simeq \sigma_{J/\psi N}^{\text{nDS}} + 2 \text{ mb, except at RHIC}$$

# Nuclear parton densities

## Comparing to more nPDF sets

	Proton	nDS	nDSg	EKS98	EPS08	HKM
$\sigma_{J/\psi N}^{\text{nPDF}} \text{ (mb)}$	$3.4 \pm 0.2$	$3.5 \pm 0.2$	$4.0 \pm 0.2$	$5.2 \pm 0.2$	$6.0 \pm 0.2$	$3.6 \pm 0.2$
$\chi^2/\text{ndf}$	1.4	1.4	1.5	1.5	1.7	1.4

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$\chi^2/\text{ndf}$	1.4	1.4	1.5	1.5	1.7	1.4

## Global fit from all experiments and nPDF sets

$$\sigma_{J/\psi N} = 3.5 \pm 0.2 \text{ (stat.)} \pm 2.6 \text{ (syst.) mb}$$

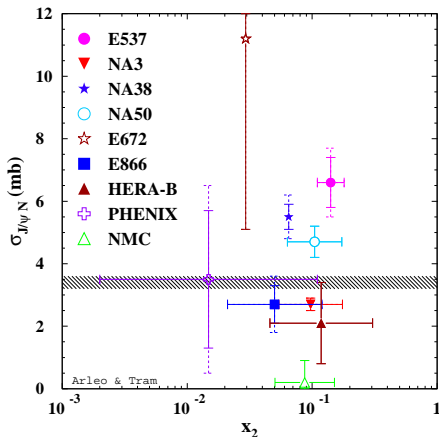
- **Huge spread** from the unknown gluon nPDF
- Extrapolation to **higher energies** (smaller  $x$ ) even more **problematic**

# Energy dependence

Naively,  $\sigma_{J/\psi N}$  should be a function of

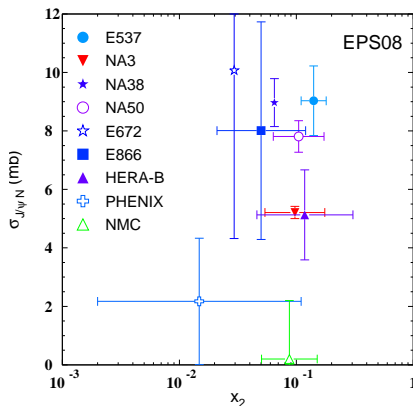
$$\sqrt{s_{J/\psi N}} \simeq m_{J/\psi} \sqrt{\frac{1+x_2}{x_2}}$$

# Energy dependence



- No statistically significant energy dependence of  $\sigma_{J/\psi N}$  found in data

# Energy dependence



• ...yet again this interpretation depends on which nPDF set is used!

# Interlude

Clearly, the study of  $J/\psi$  absorption in nuclei suffers from too large theoretical uncertainties on gluon nPDF – **How to improve this?**



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nPDF sets currently use

- DIS on nuclei
- Drell-Yan in p A collisions
- Scaling violations in DIS on nuclei

valence quarks

sea quarks

gluons

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## A few candidates for the gluon sector

- Jets
- Large  $p_{\perp}$  dileptons
- Heavy-bosons  $W^{\pm}, Z^0$
- Prompt photons

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What about using also heavy-quarkonium production?

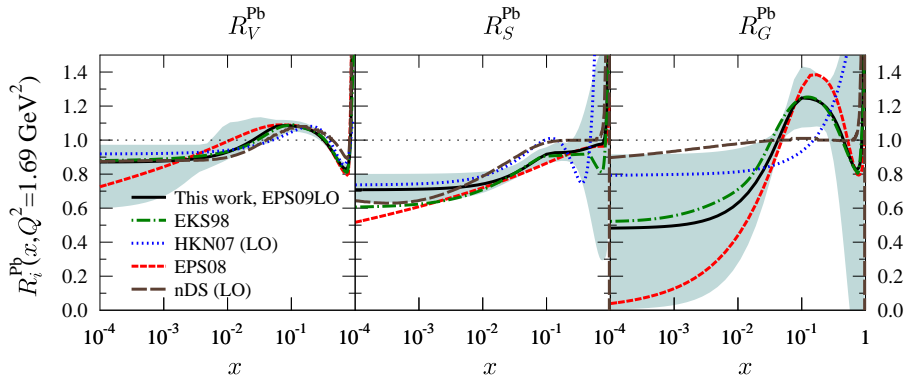
- Surprising at first sight because of the inelastic interaction processes
- Accurate data on large kinematical range might give interesting constraints

# Gluon nPDF from $J/\psi$ data

The behavior of  $G^A/G^P(x)$  **pretty different** from one set to another

- nDS, HKM
  - rather flattish, (anti)shadowing not too pronounced
- nDSg, EKS98, EPS09
  - (anti)shadowing more pronounced
- EPS08
  - very strong gluon (anti)shadowing
  - related to the use of forward hadron production data at RHIC

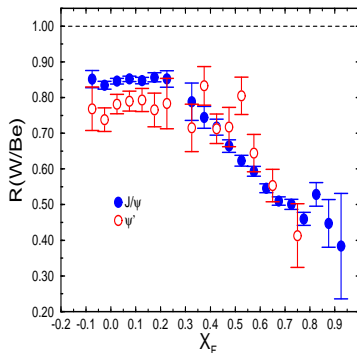
# Gluon nPDF from $J/\psi$ data



- Behavior particularly different in the range  $10^{-2} \lesssim x \lesssim 10^{-1}$
- A fast variation of  $G^A/G^P(x)$  could be seen (in principle) in  $J/\psi$  data

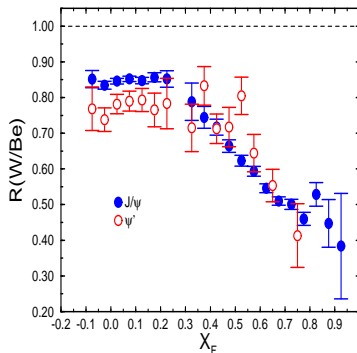
# Gluon shadowing and E866/NuSea data

- $10^{-2} \lesssim x_2 \lesssim 10^{-1}$  corresponds to  $0 \lesssim x_F \lesssim 0.3$  at  $\sqrt{s_{pA}} \simeq 40$  GeV at the  $J/\psi$  mass scale
- Precise E866/NuSea measurements in this range



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Rather flat behavior which tend to rule out a fast variation of  $G^A/G^p(x)$

# A simple analysis

## Fitting E866/NuSea $J/\psi$ data

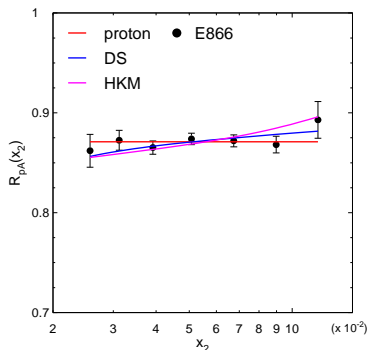
- using all nPDF sets (proton PDF, nDS/nDSg, EKS98, HKM, EPS08)
- inelastic cross section as a **constant free** parameter



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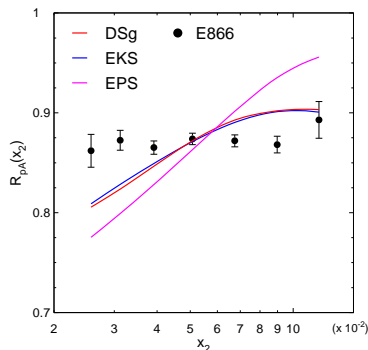


- Behavior consistent with nDS and HKM

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- inelastic cross section as a **constant free** parameter



- Conflict with nDSg and EKS98, worse with EPS08

# A simple analysis

$\chi^2$

nPDF set	proton	DS	DSg	EKS	EPS	HKM
$\chi^2/\text{ndf}$	0.5	0.6	12.0	10.1	35.9	1.2

## Normalizations

	nDS	nDSg	EKS98	EPS08	HKM
$\sigma_{J/\psi N}$ (mb)	2.5	2.7	4.2	4.1	2.5
$S_{\text{abs}}$	0.89	0.88	0.82	0.82	0.89

## Caveat

- Assumption of **constant absorption cross section** can be questioned
- Suppose the absorption weakens with decreasing  $x_2$  (increasing  $x_F$ ) the disagreement would not be as pronounced

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## “Fine tuning issue”

How come a very strong variation of  $G^A/G^p(x)$  be perfectly balanced by nuclear absorption with opposite behavior, resulting into a constant  $J/\psi$  suppression? **Rather unnatural**

- **Systematic study** of  $J/\psi$  suppression in  $\gamma^* A$  and  $p A$  collisions
  - Results somehow spoiled by nPDF uncertainties
- **Constraints on nPDF** from specific  $J/\psi$  data sets
  - E866/NuSea favor nPDF sets with moderate variation in the  $10^{-2} \lesssim x \lesssim 10^{-1}$  range