

# Nucleons from neutrino interactions

and why  $\nu$  oscillations need  $\nu$  interaction measurements!

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# The precision era of $\nu$ oscillations

## Latest results

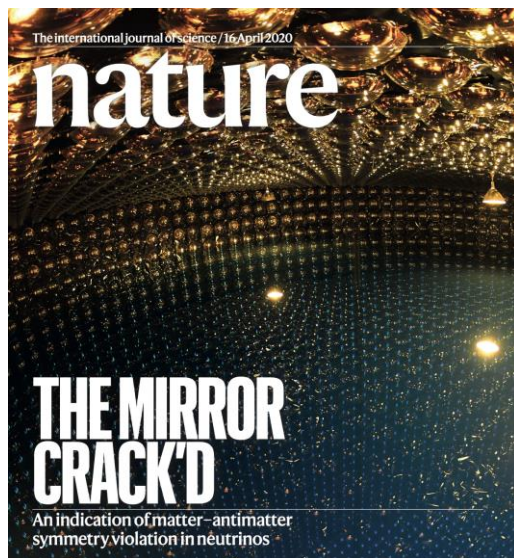
- **Indication of CP violation!**
- Currently largely limited by statistics ... but not for long!



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## Current systematic uncertainties

Source (T2K)	$N(\nu_e)$
Binding Energy	7.1%
<b>Total Syst.</b>	<b>8.8%</b>
Nature 580, 339-344 / arXiv:2101.03779	
Source (NOVA)	$N(\nu_e)$
$\sigma_{\nu N}$ and FSI	7.7%
<b>Total Syst.</b>	<b>9.2%</b>
Phys. Rev. D 98, 032012	

- Current results use  $\sim 100 \nu_e + \bar{\nu}_e$ , expect **1000-2000** for DUNE/HK
- $\sim 2-3\%$  stat. precision on CP asymmetry

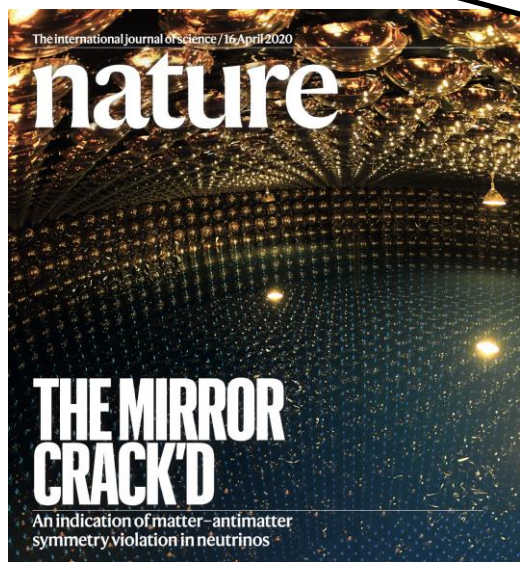
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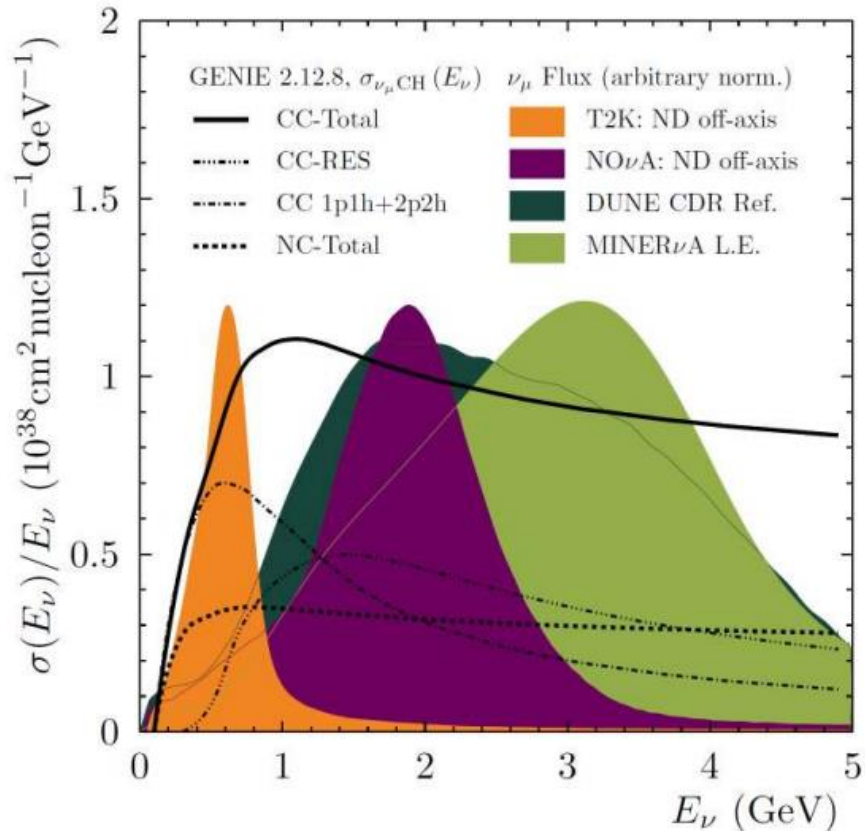
**Neutrino interaction uncertainties must be reduced for DUNE/Hyper-K to succeed**



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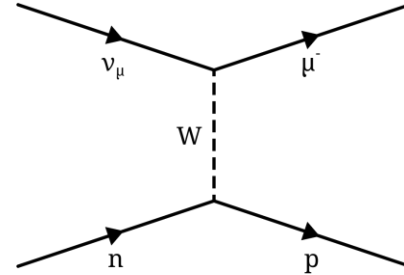
- use expect for DUNE/HK
- 1000-2000
- ~2-3% stat. precision on CP asymmetry

# Neutrino interactions

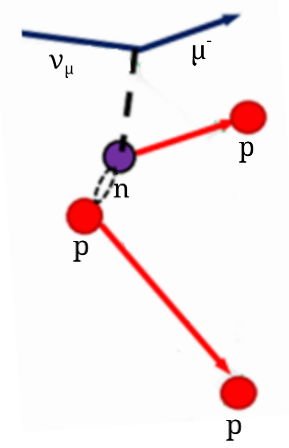


**Key challenge: estimating  $E_\nu$**

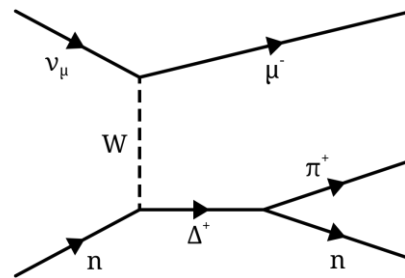
**CCQE (1p1h)**  
(Charged-Current Quasi-Elastic)



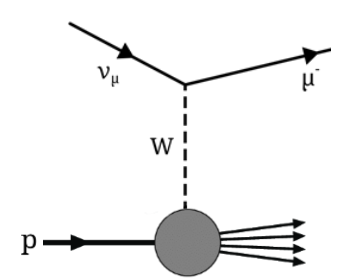
**CC 2p2h**  
(2 particle, 2 hole)



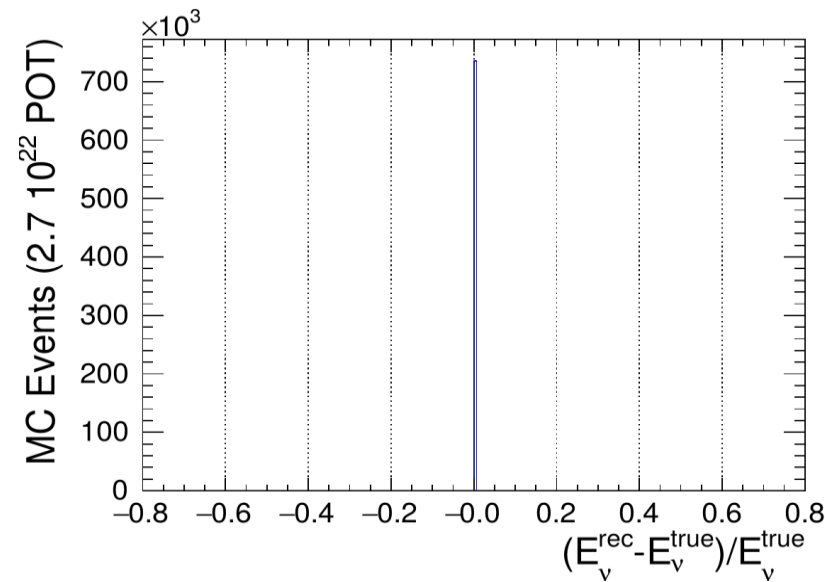
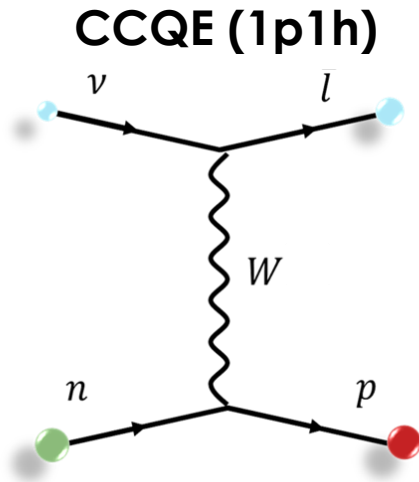
**CCRES**  
(Charged-Current Resonant)



**CCDIS**  
(Deep Inelastic Scattering)



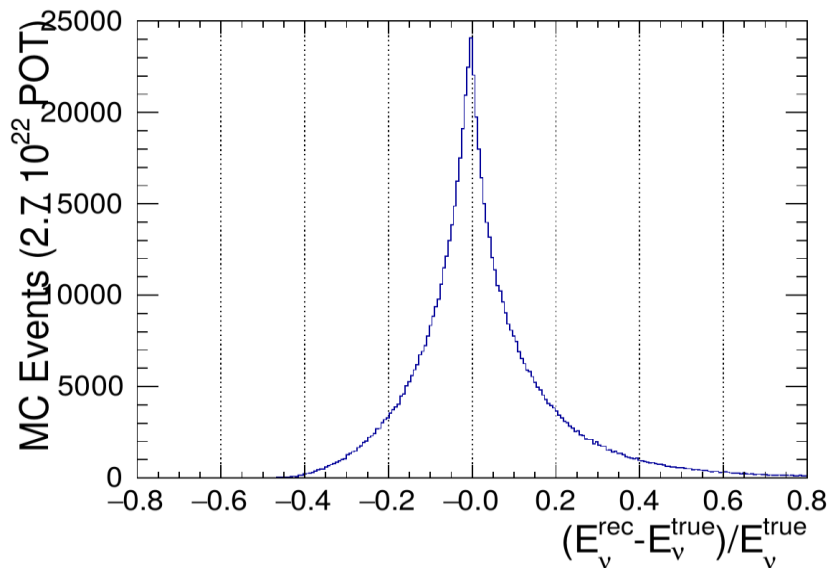
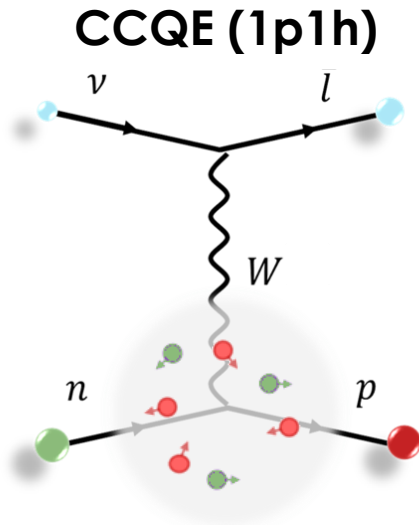
# Nuclear effects and $E_\nu$ (T2K/HK)



$$E_\nu = \frac{m_p^2 - (m_n - E_b)^2 - m_\mu^2 + 2(m_n - E_b)E_\mu}{2(m_n - E_b - E_\mu + p_\mu \cos \theta_\mu)}$$

Proxy for  $E_\nu$  from lepton kinematics is perfect only for **CCQE elastic scattering** off a **stationary nucleon**

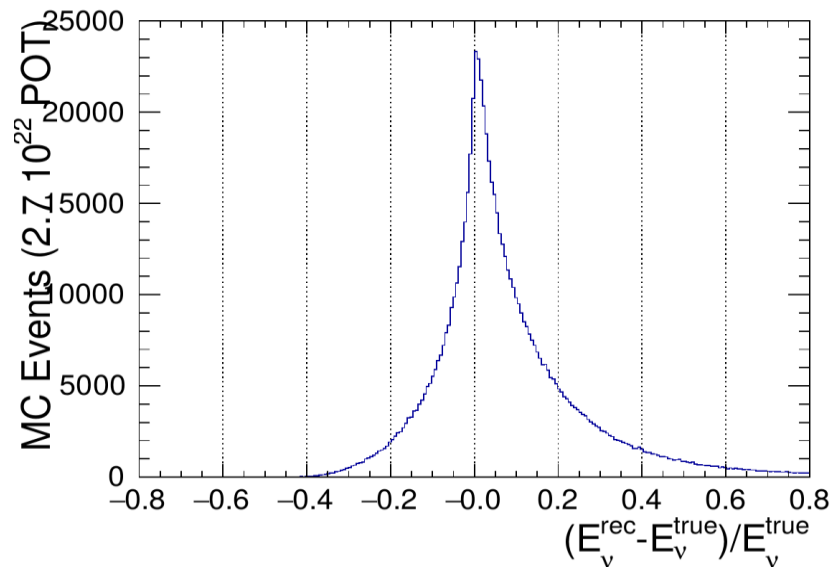
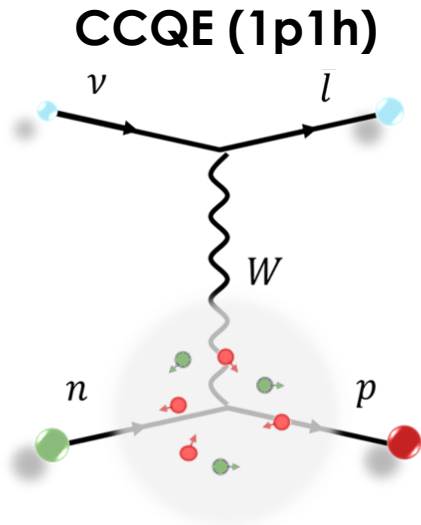
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The motion of the nucleons inside the nucleus (*Fermi motion*) causes a **smearing** on  $E_\nu$

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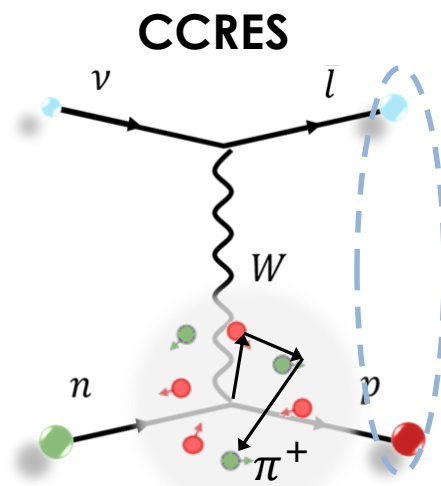
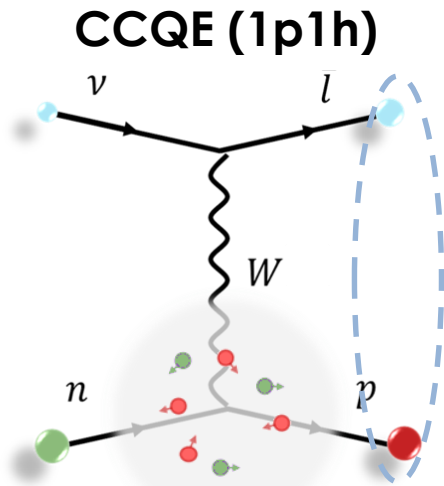
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The motion of the nucleons inside the nucleus (*Fermi motion*) causes a **smearing** on  $E_\nu$

The energy loss in the nucleus (to extract the struck nucleon from its shell) introduces a **bias**

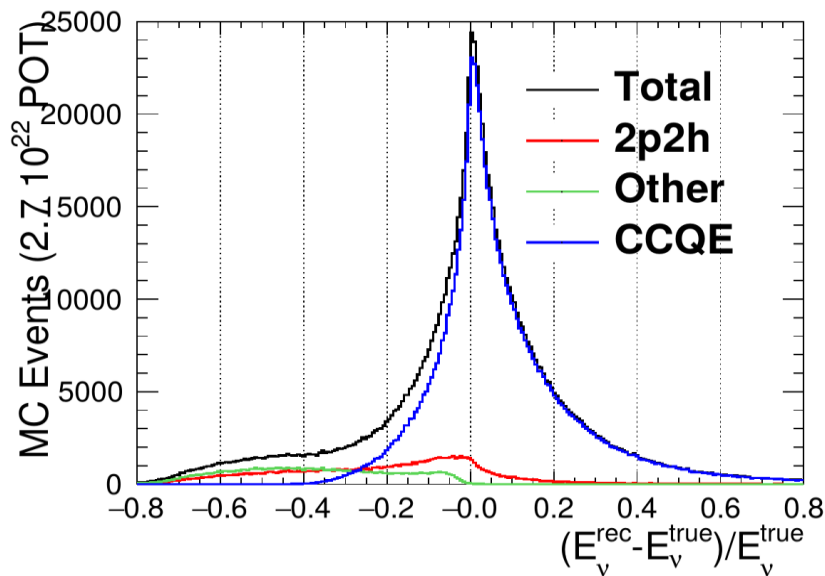
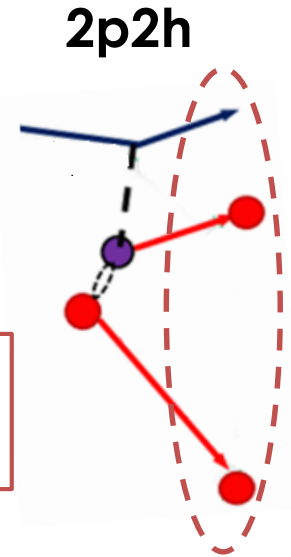


# Nuclear effects and $E_\nu$ (T2K/HK)



**Final state interactions** (FSI) can cause different interaction modes to have the same final state

Interactions off a bound state of two nucleons can result in **2p2h** final states



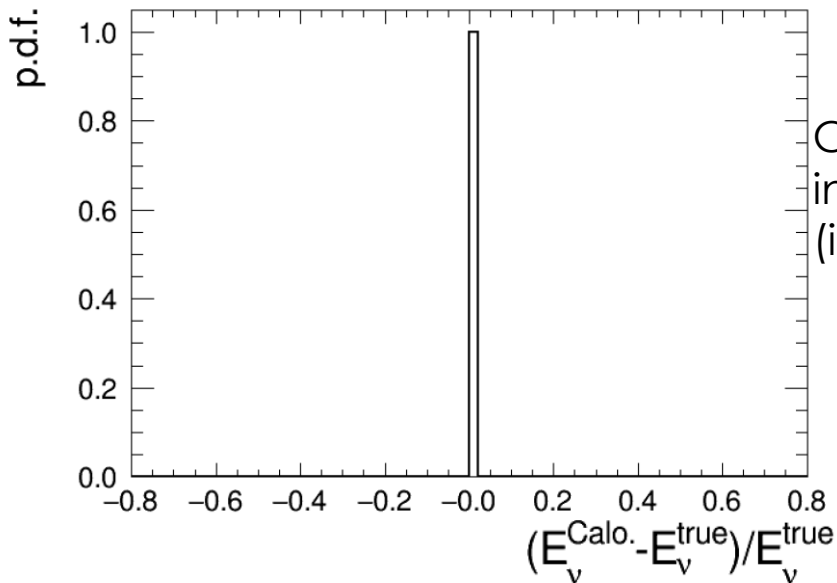
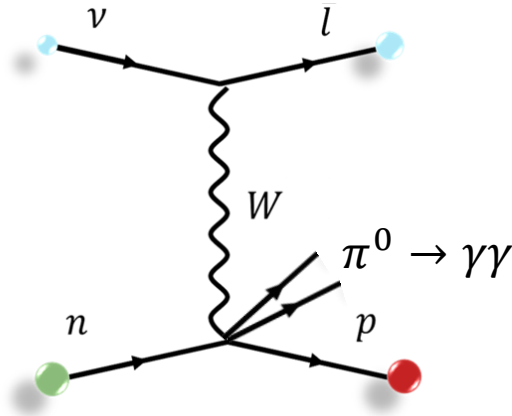
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The motion of the nucleons inside the nucleus (*Fermi motion*) causes a **smearing** on  $E_\nu$

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Not a good proxy for non-CCQE events: 2p2h and CC1π with pion abs. FSI

# Nuclear effects and $E_\nu$ (SBN/DUNE/NOvA)

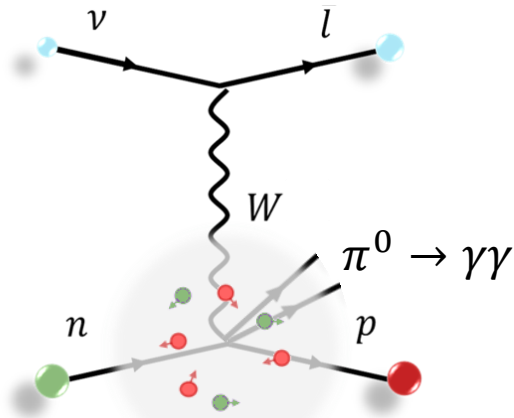


$$E_\nu^{\text{calo}} = E_\ell + E_{\text{had.}} = E_\ell + \Sigma T_p + \Sigma T_{\pi^\pm} + \Sigma E_\gamma$$

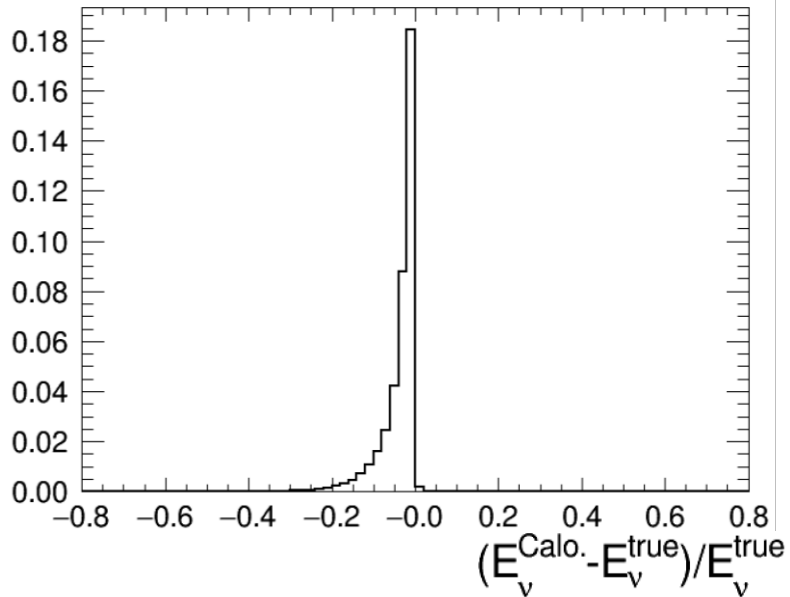
Calculation from calorimetry is perfect only for interactions **without neutrons and charged pions** (ignore heavier mesons here) off a **stationary nucleon**

Usefulness is not restricted to QE-like interactions (no final state pions)

# Nuclear effects and $E_\nu$ (SBN/DUNE/NOvA)



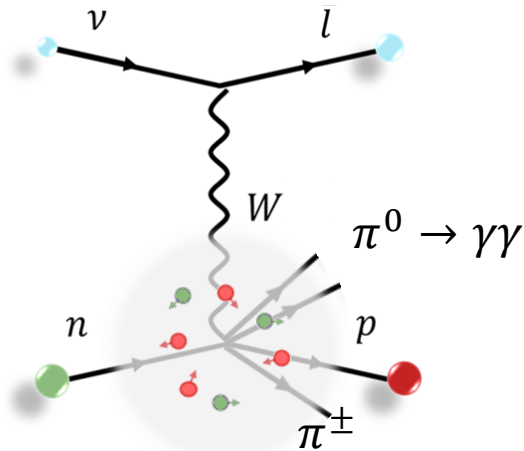
All events without  $n$  or  $\pi^\pm$



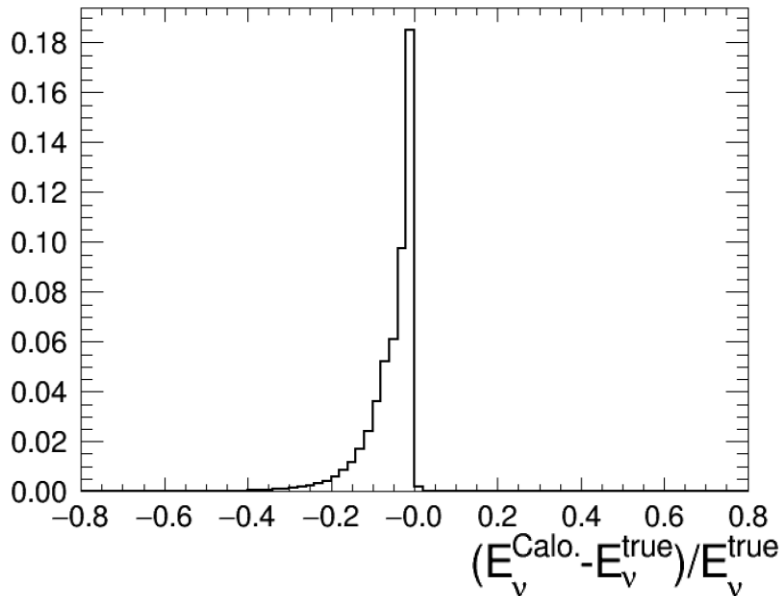
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Impact of initial state effects (Fermi motion and removal energy) smaller than in QE approach

# Nuclear effects and $E_\nu$ (SBN/DUNE/NOvA)



All events without  $n$

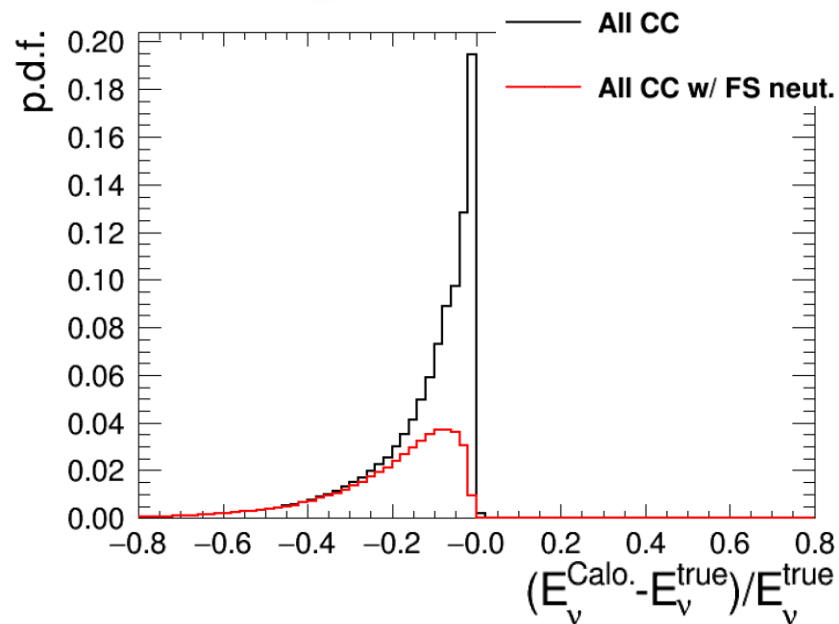
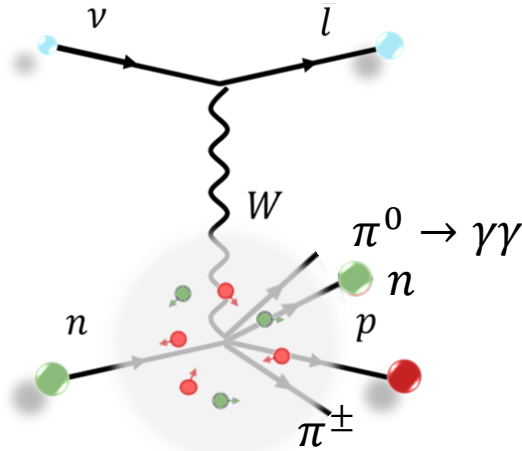


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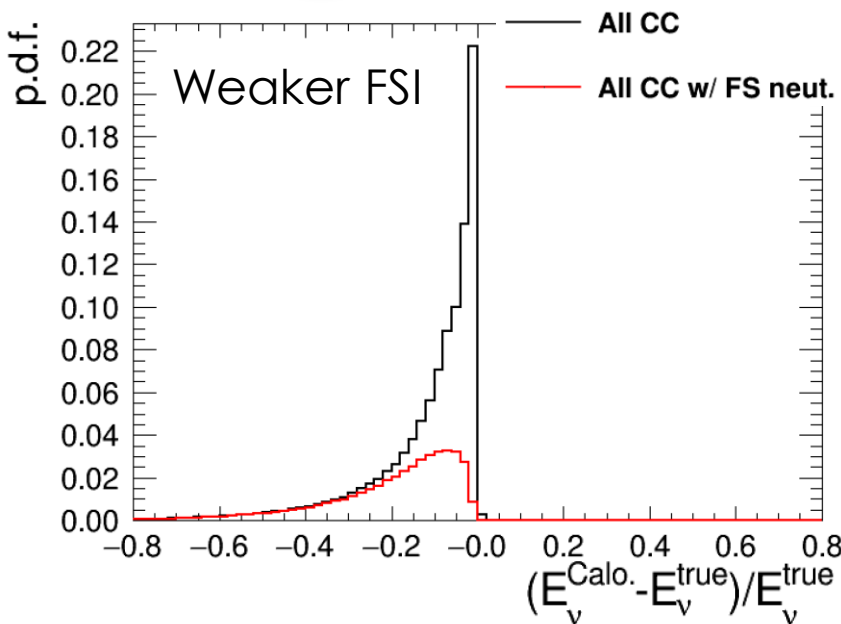
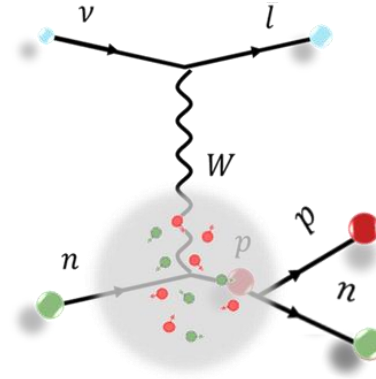
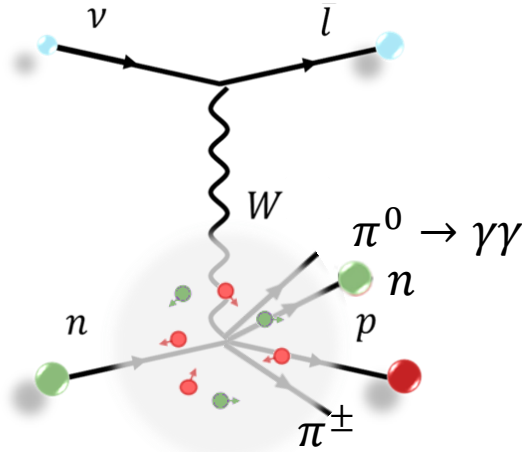
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# Nuclear effects and $E_\nu$ (SBN/DUNE/NOvA)

## Final State Interactions (FSI)



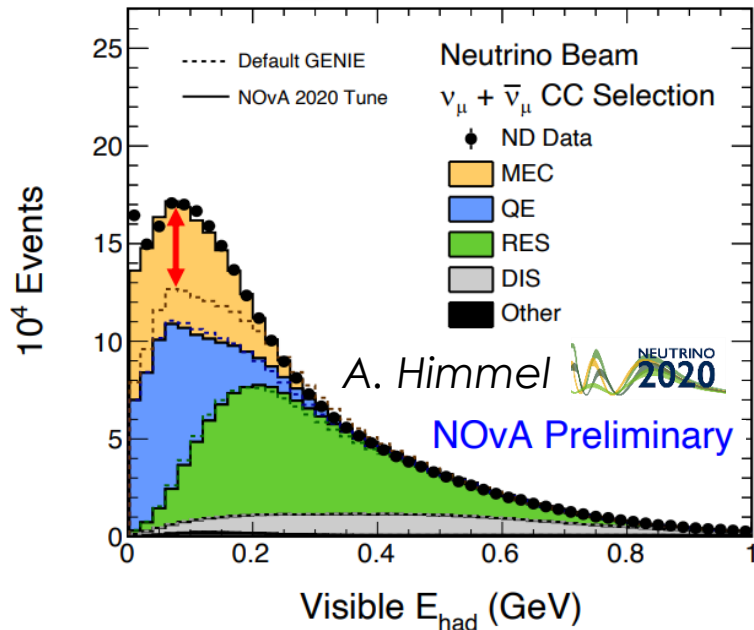
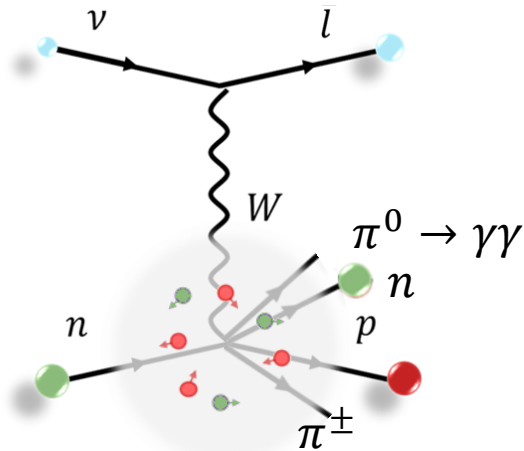
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Fraction of  $E_\nu$  in **Neutrons** is critical – affected importantly by FSI (and more important for  $\bar{\nu}$ )

# Nuclear effects and $E_\nu$ (SBN/DUNE/NOvA)



- Complex interaction topologies make  $E_{had}$  tough to model
- NOvA find strong data/simulation discrepancy at low  $E_{had}$  (before applying a 2p2h modification)
- Covered by generous systematics, but this must be better understood for DUNE

# What we need to know (a non exhaustive list!)

## T2K/HK

("kinematic"  $E_\nu$  proxy)

### Critical

- Nuclear ground state: **Fermi motion** and "**binding energy**"
- **2p2h** and **pion absorption FSI** contributions to  $0\pi$  final states

### Important

- Impact of **nucleon FSI** on  $\sigma(\nu_e)/\sigma(\nu_\mu)$  (see backups for details)
- Differences between interactions on Carbon and Oxygen

## SBN/DUNE/NOvA

("calorimetric"  $E_\nu$  proxy)

### Critical

- Neutron production:
  - **FSI**
  - **2p2h**
  - DIS hadronisation

### Important

- Charged pion multiplicities (e.g. from **FSI**)
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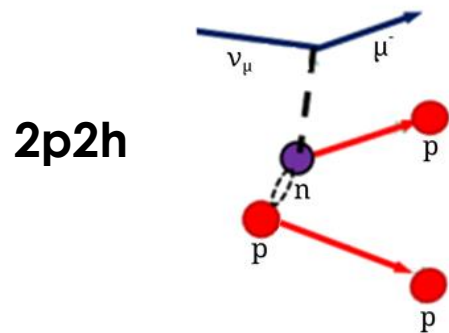
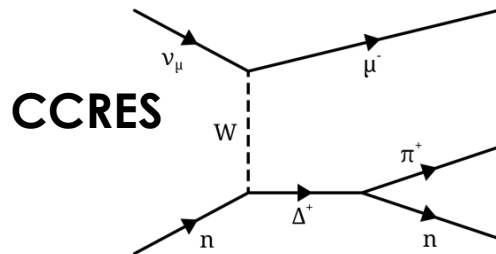
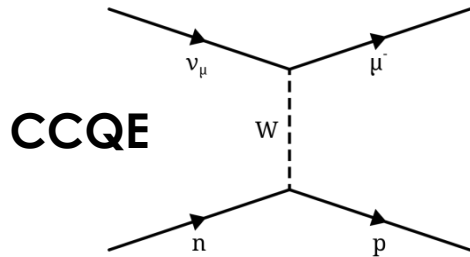
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Neutrino interaction modelling is crucial all upcoming experiments, but different experiments have different priorities: **complimentary approaches!**



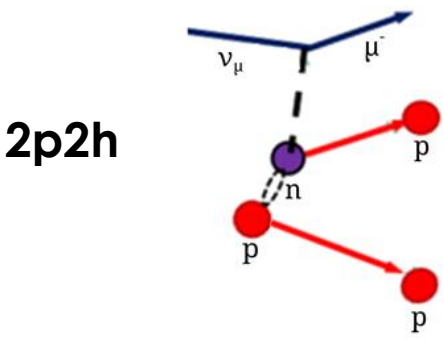
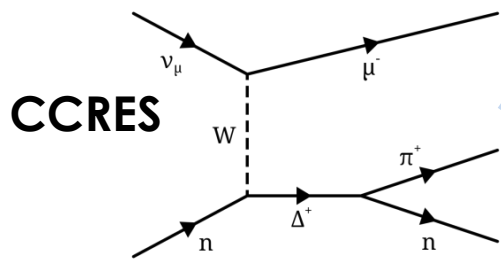
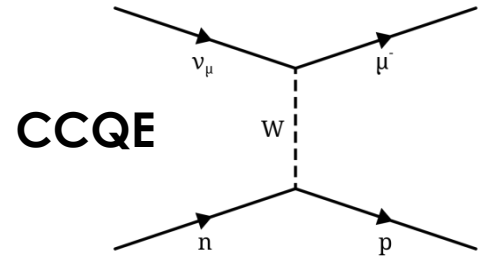
# What can we measure?

## Interaction Modes

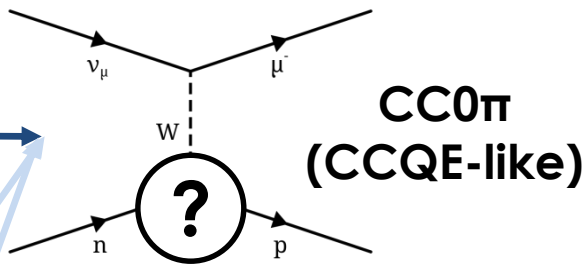


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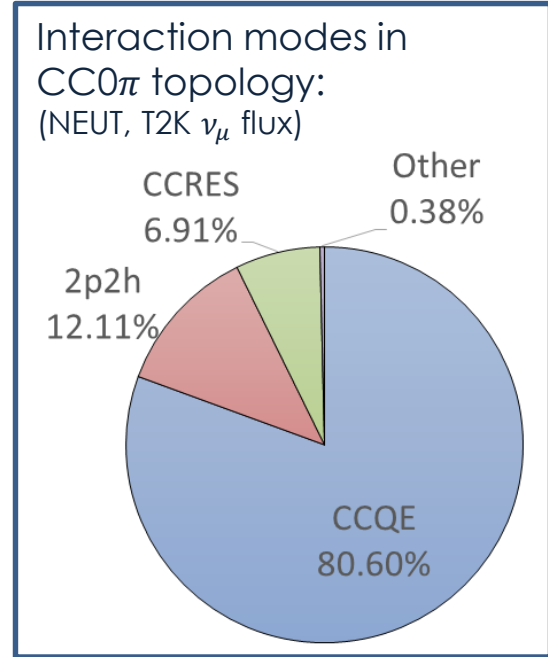
## Interaction Modes



## Interaction Topologies

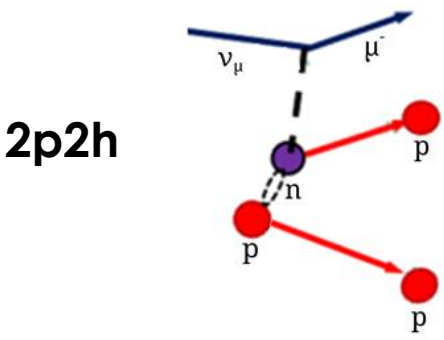
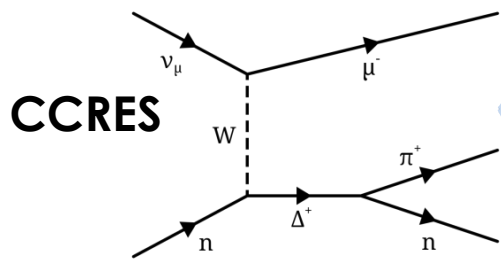
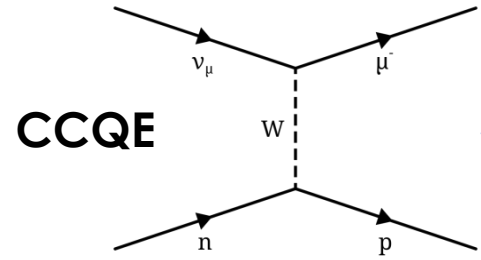


- We measure what our detector actually sees
- But this contains all sorts of physics ...

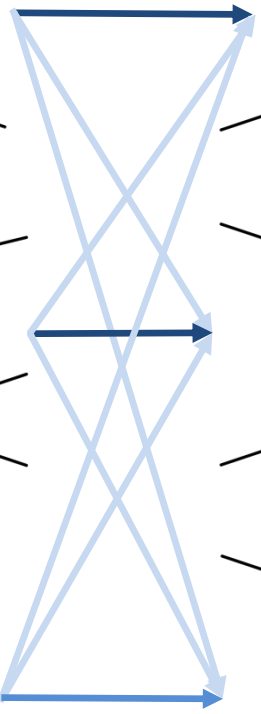
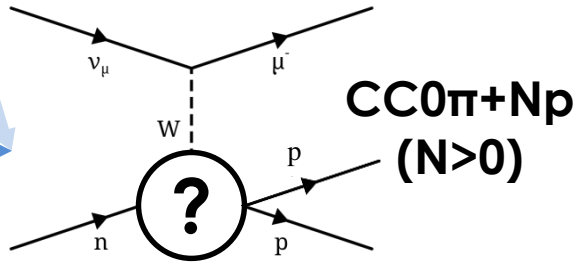
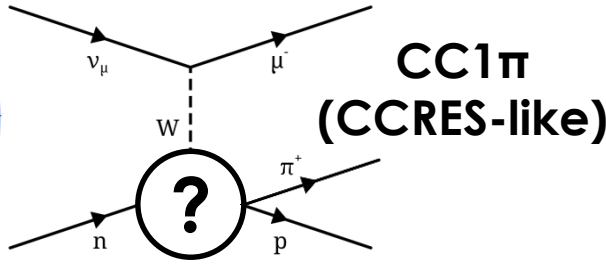
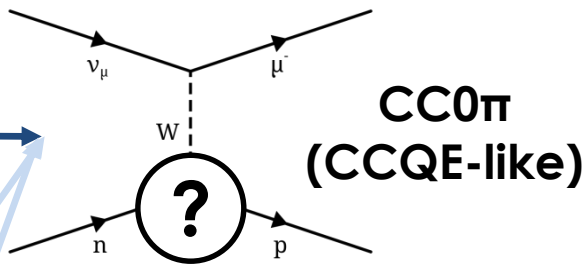


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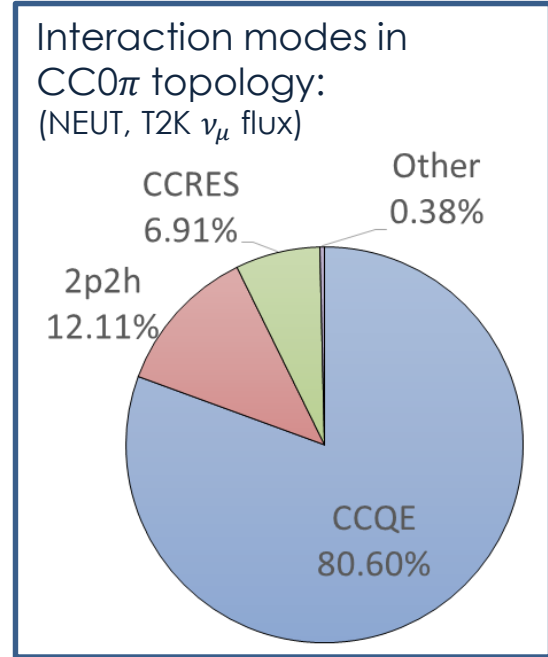
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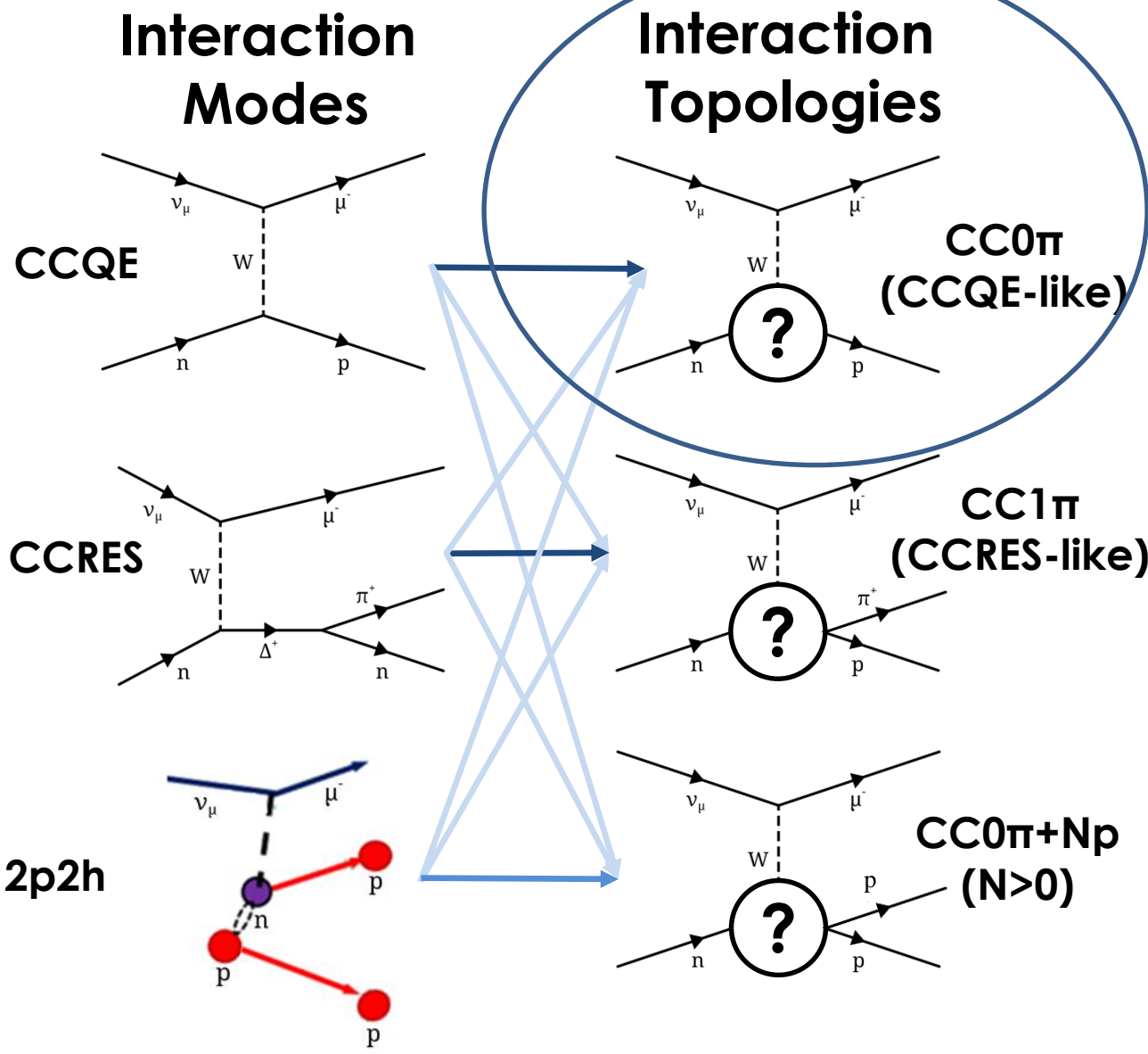
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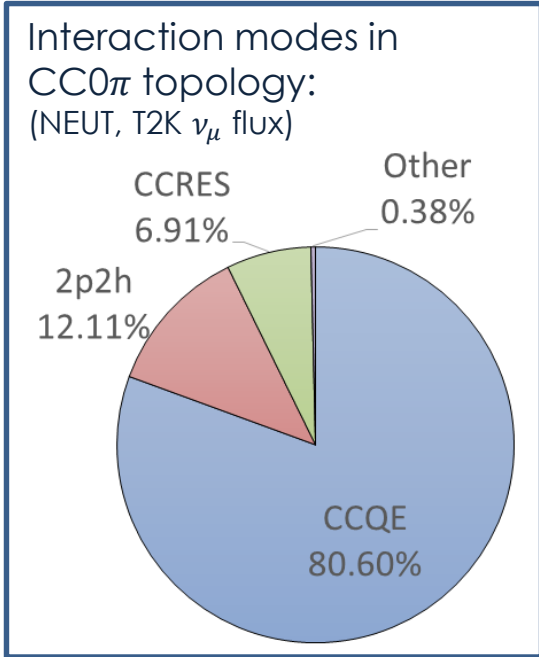
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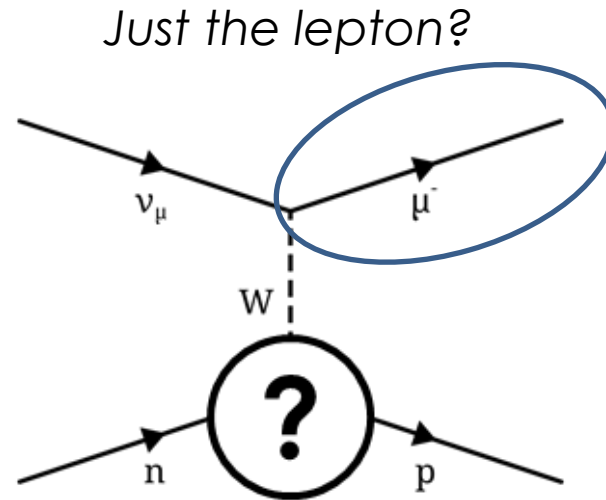
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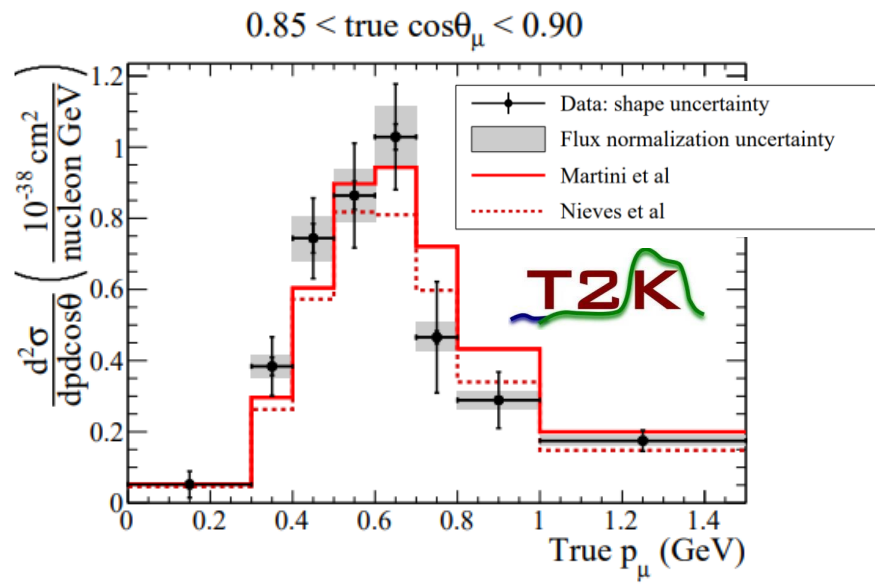
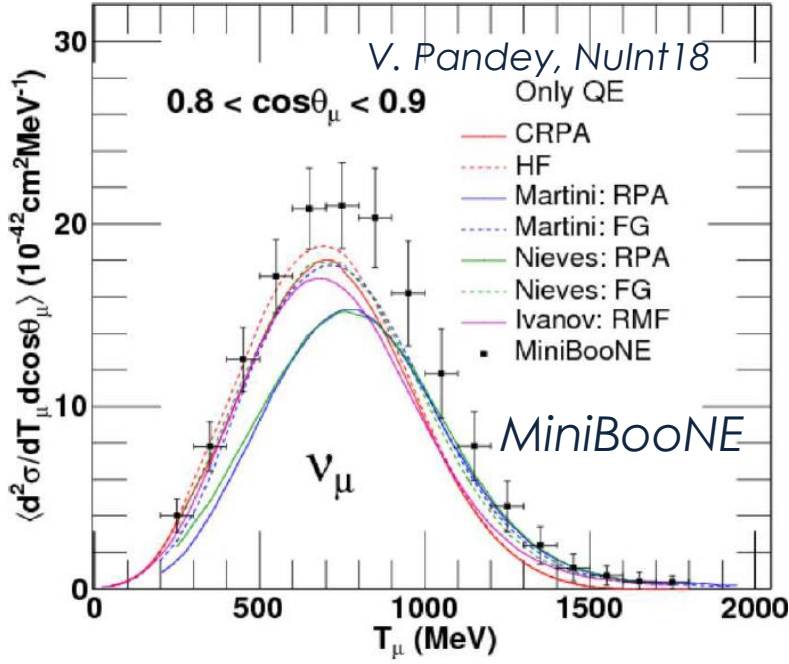
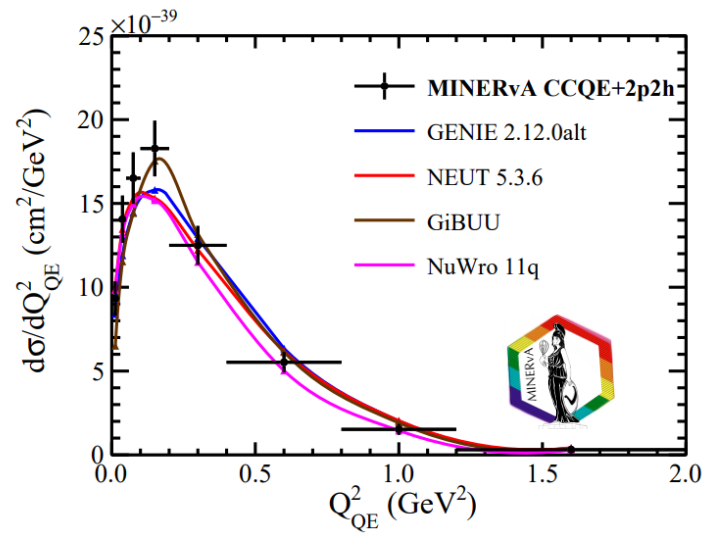


# Which observables?



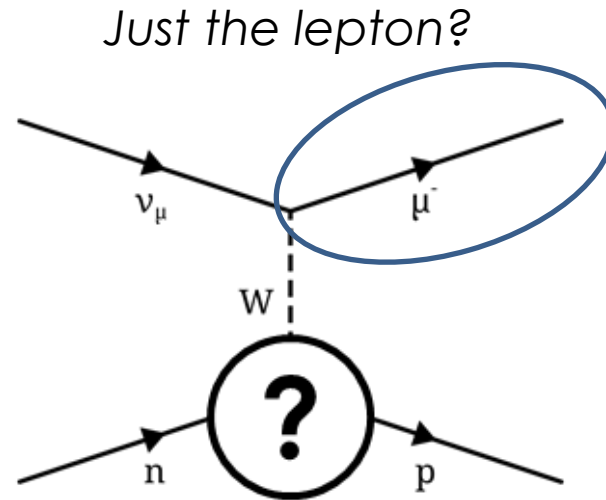
# Which observables?

- Models can broadly describe data
- But the wide-band nature of neutrino beams means the impact of nuclear effects is “smeared out”
- Dramatically different models give very similar predictions!



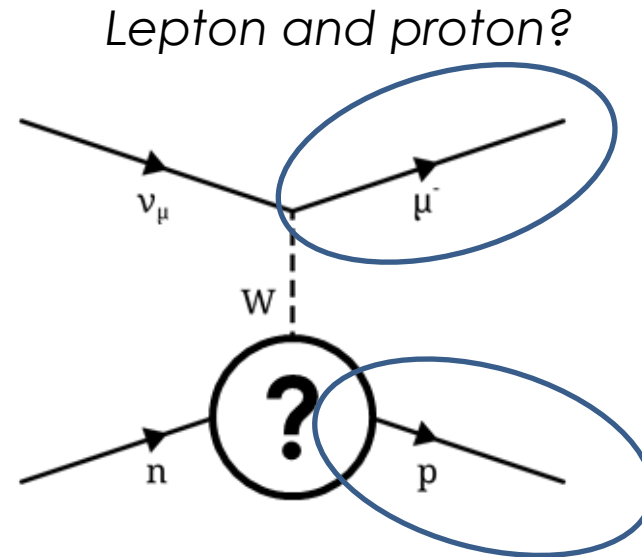


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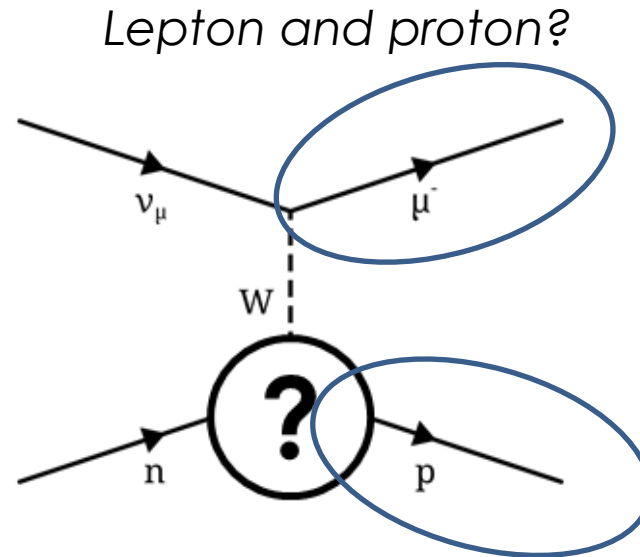


Difficult to untangle nuclear effects due to wide band neutrino fluxes.

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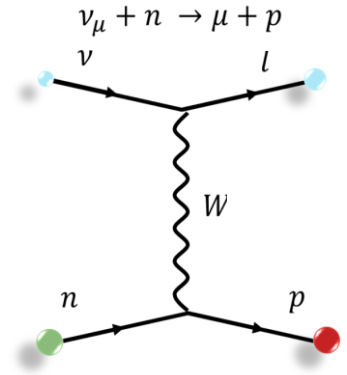
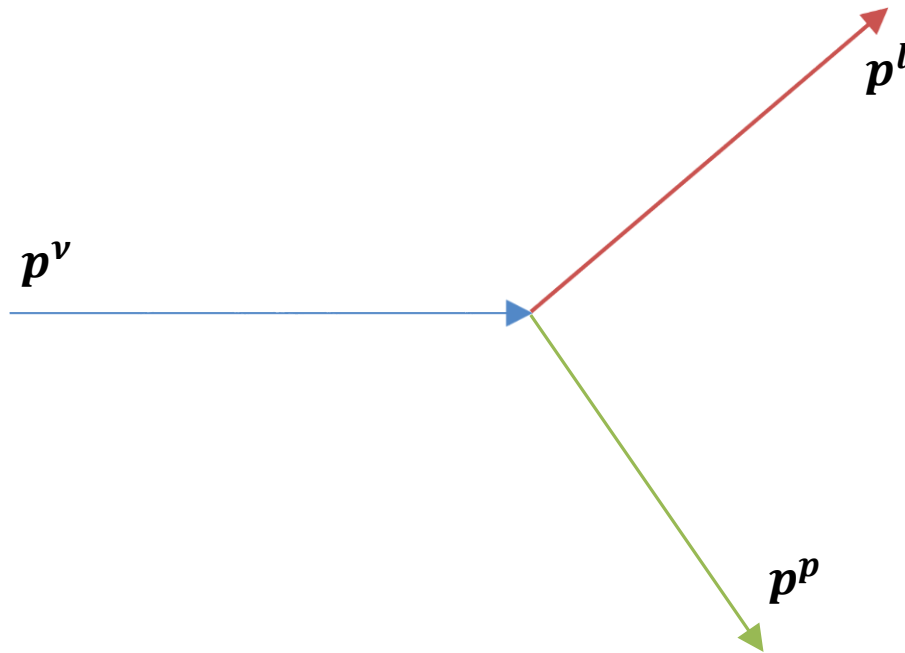


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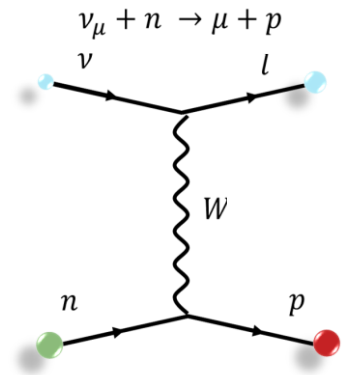
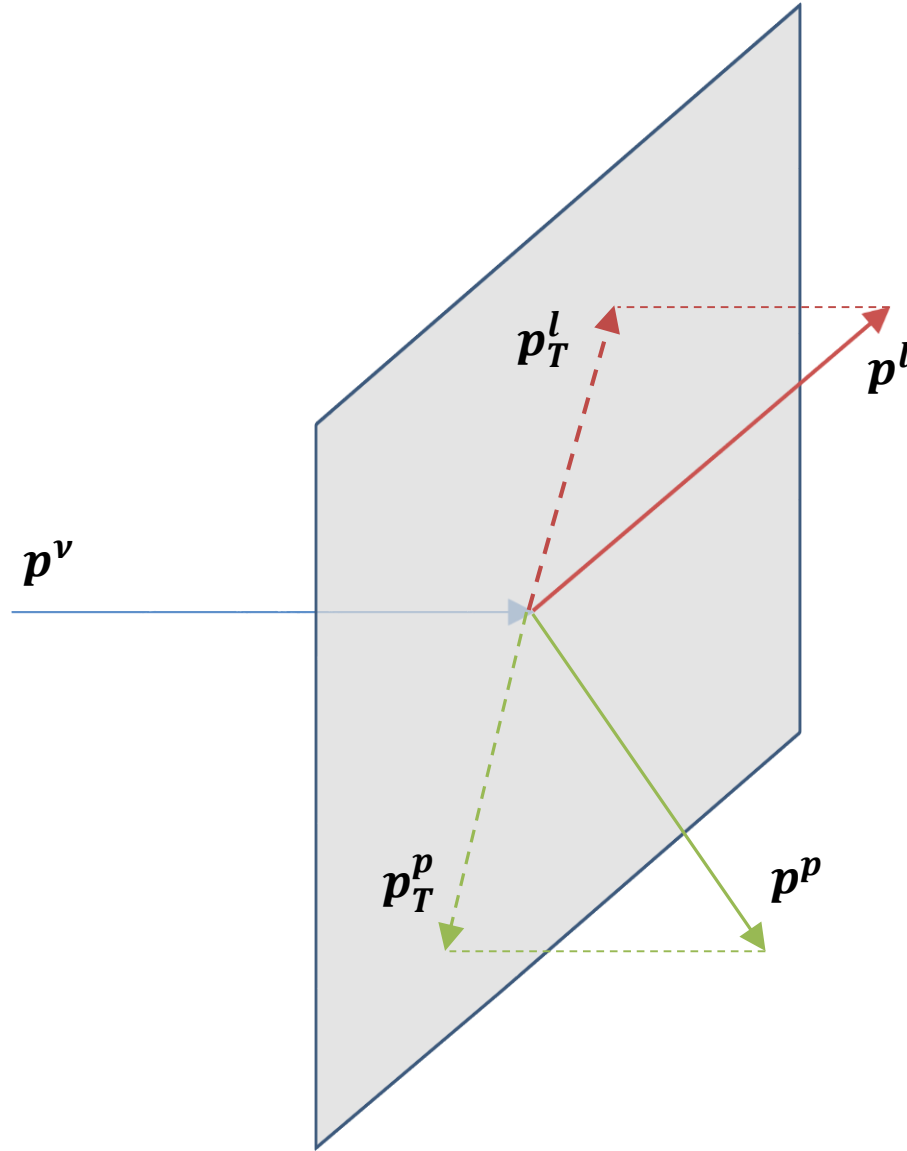
Correlations between the muon and proton kinematics allow us to disentangle nuclear effects from neutrino energy

# Single Transverse Variables



No nuclear Effects

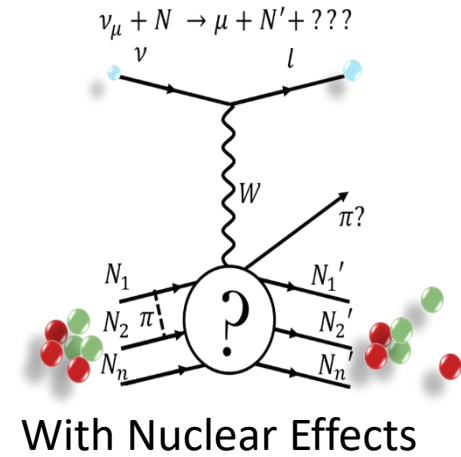
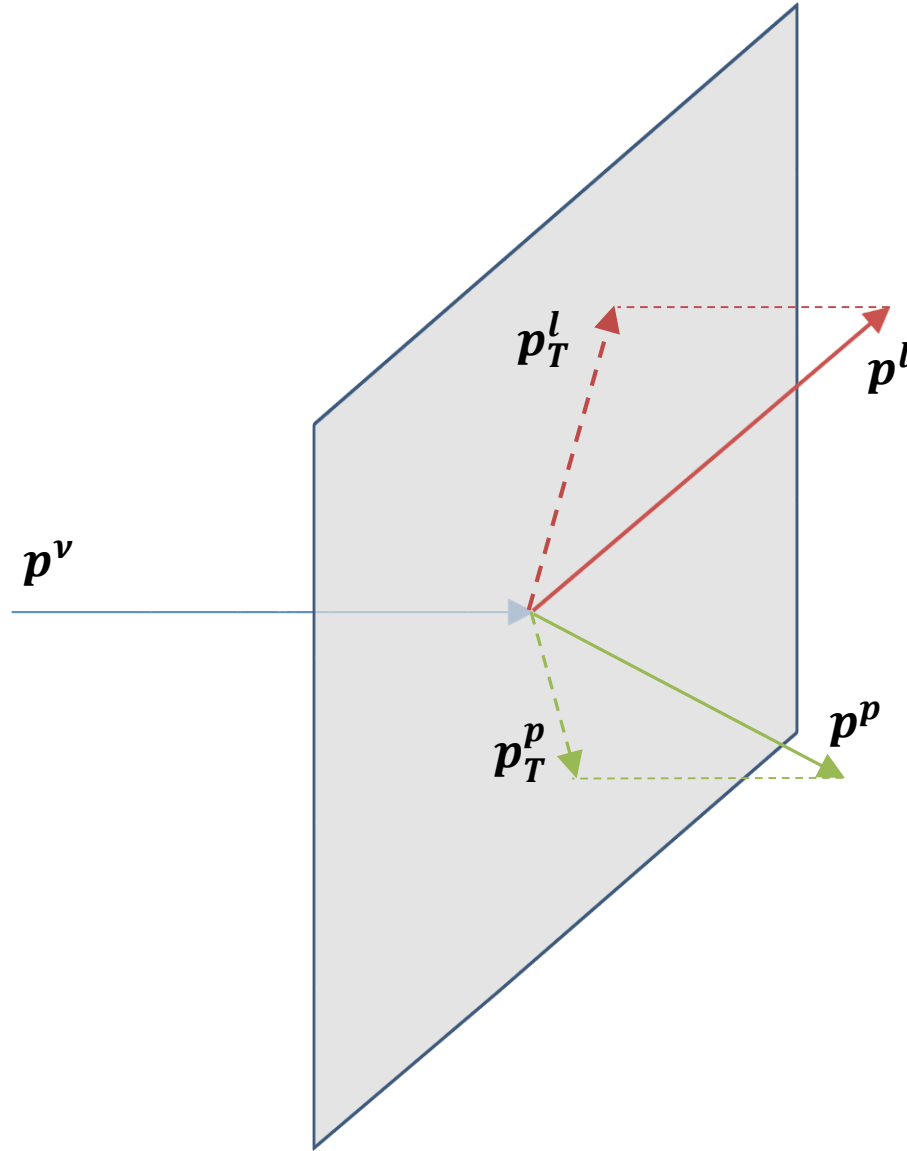
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$$p_T^l = -p_T^p$$

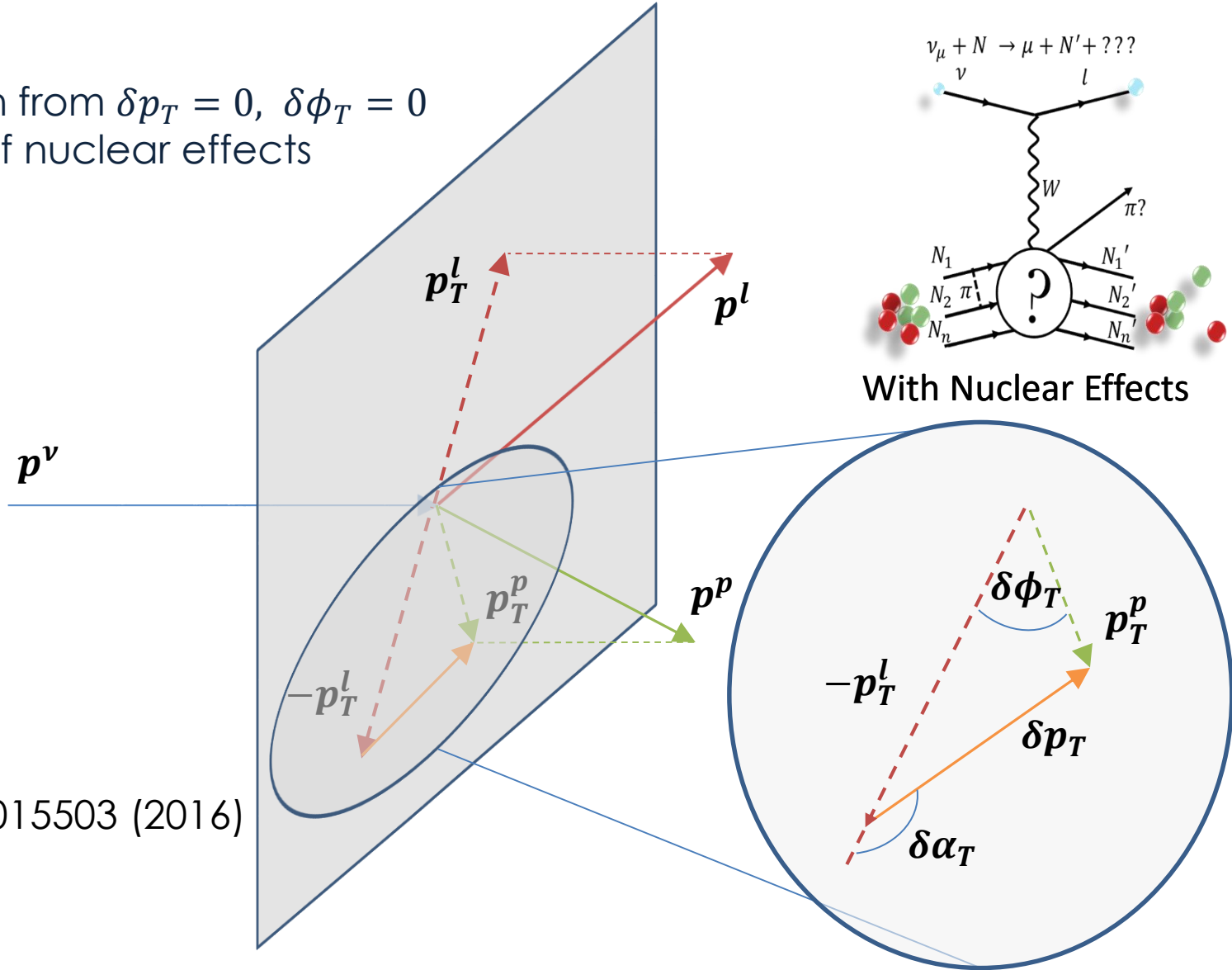
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# Single Transverse Variables

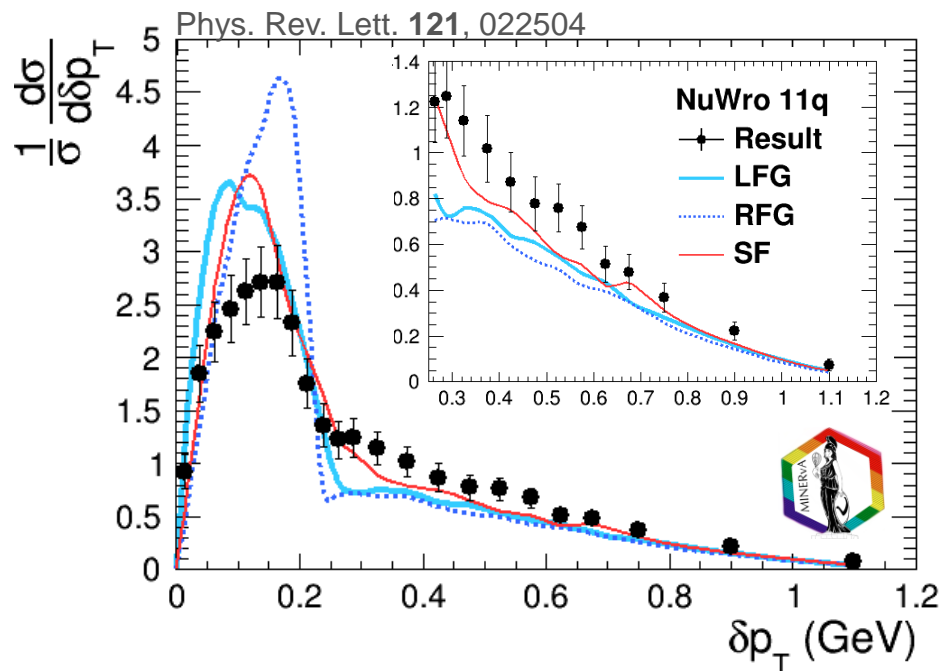
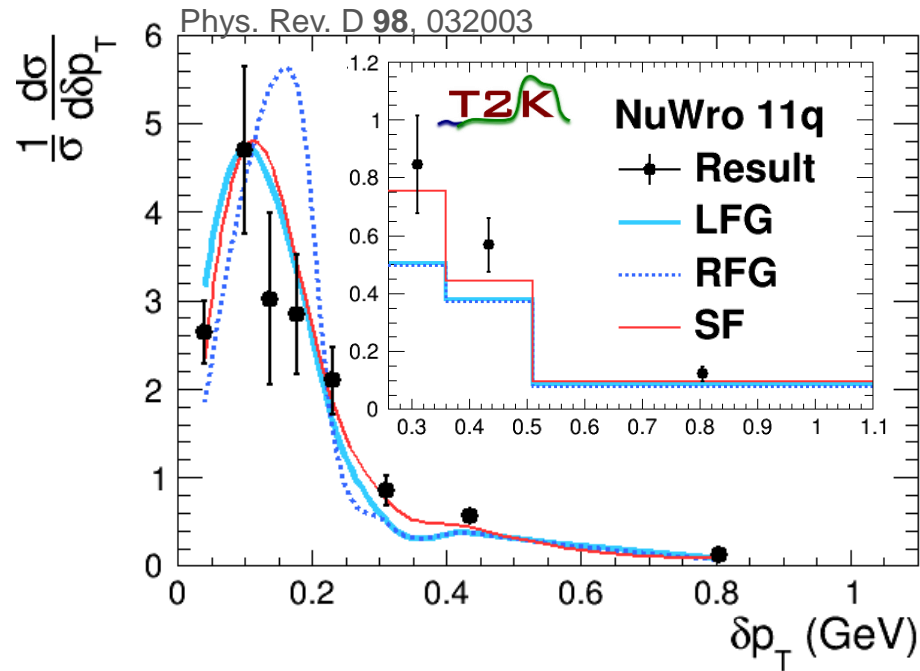
- Any deviation from  $\delta p_T = 0$ ,  $\delta\phi_T = 0$  is indicative of nuclear effects



With Nuclear Effects

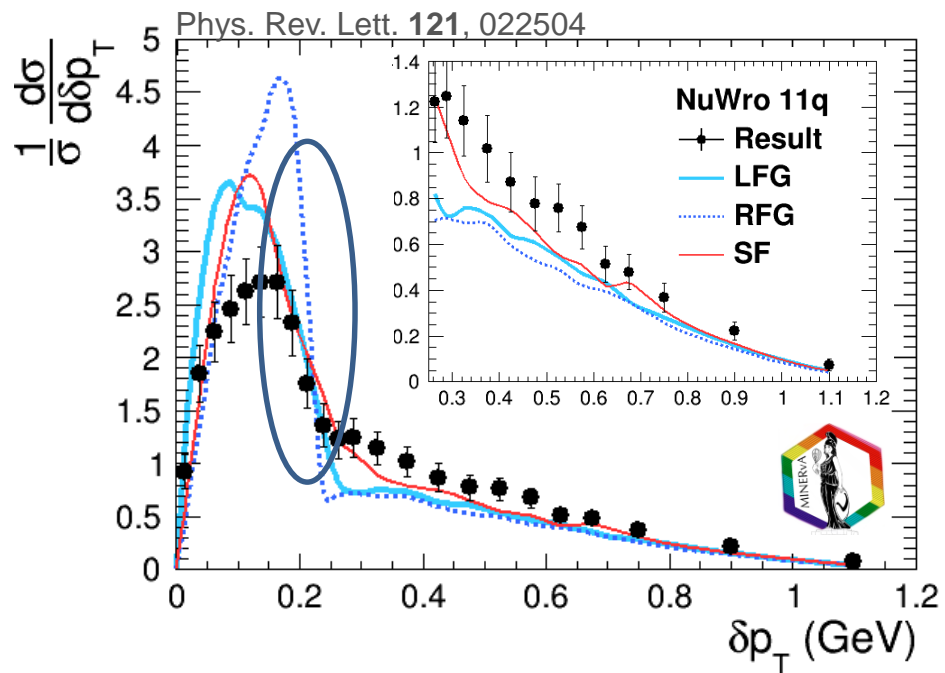
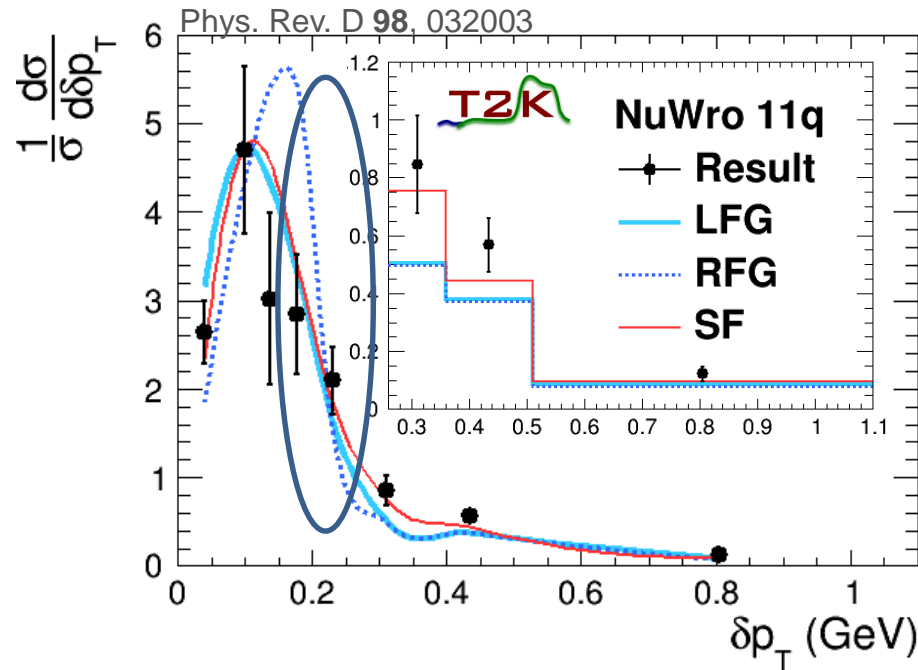
Phys. Rev. C **94**, 015503 (2016)

# Current measurements



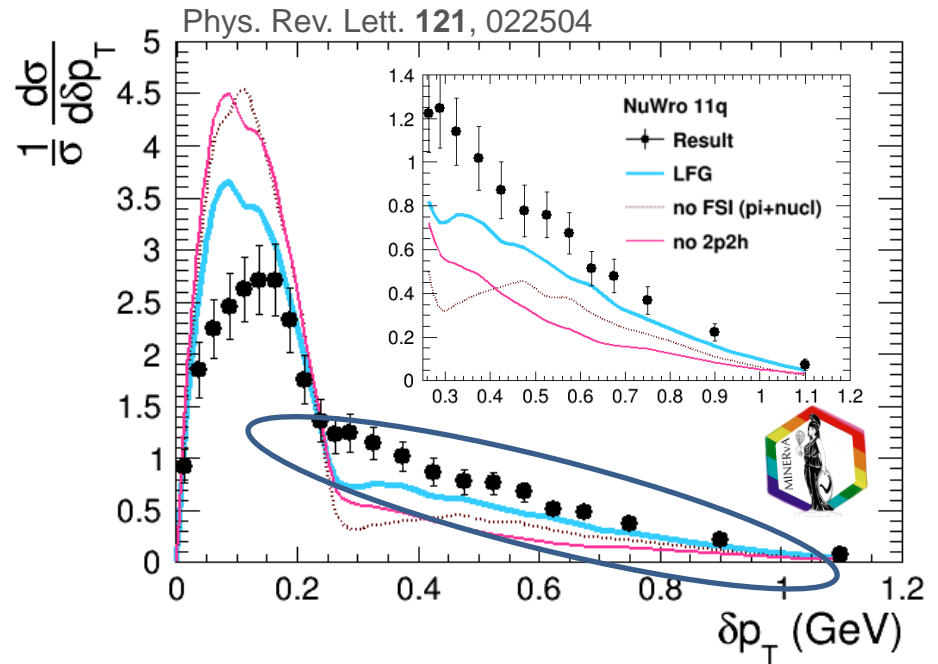
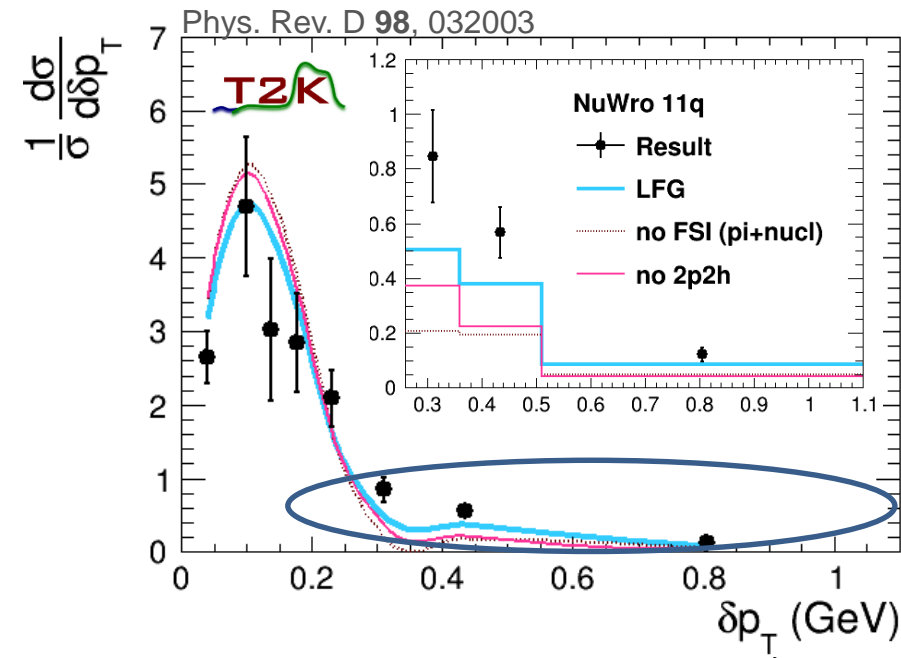


# Current measurements



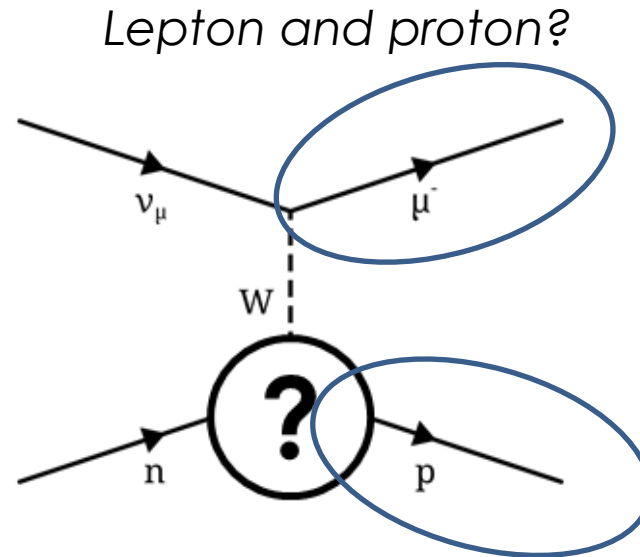
- The bulk of the distribution is sensitive to initial state nuclear effects
- But none of the models are able to provide a complete description of the data ...

# Current measurements



- The tail is more sensitive to 2p2h and FSI effects
- But, again, no model can describe the observations
- Clearly we have work to do in our modelling of nuclear effects!

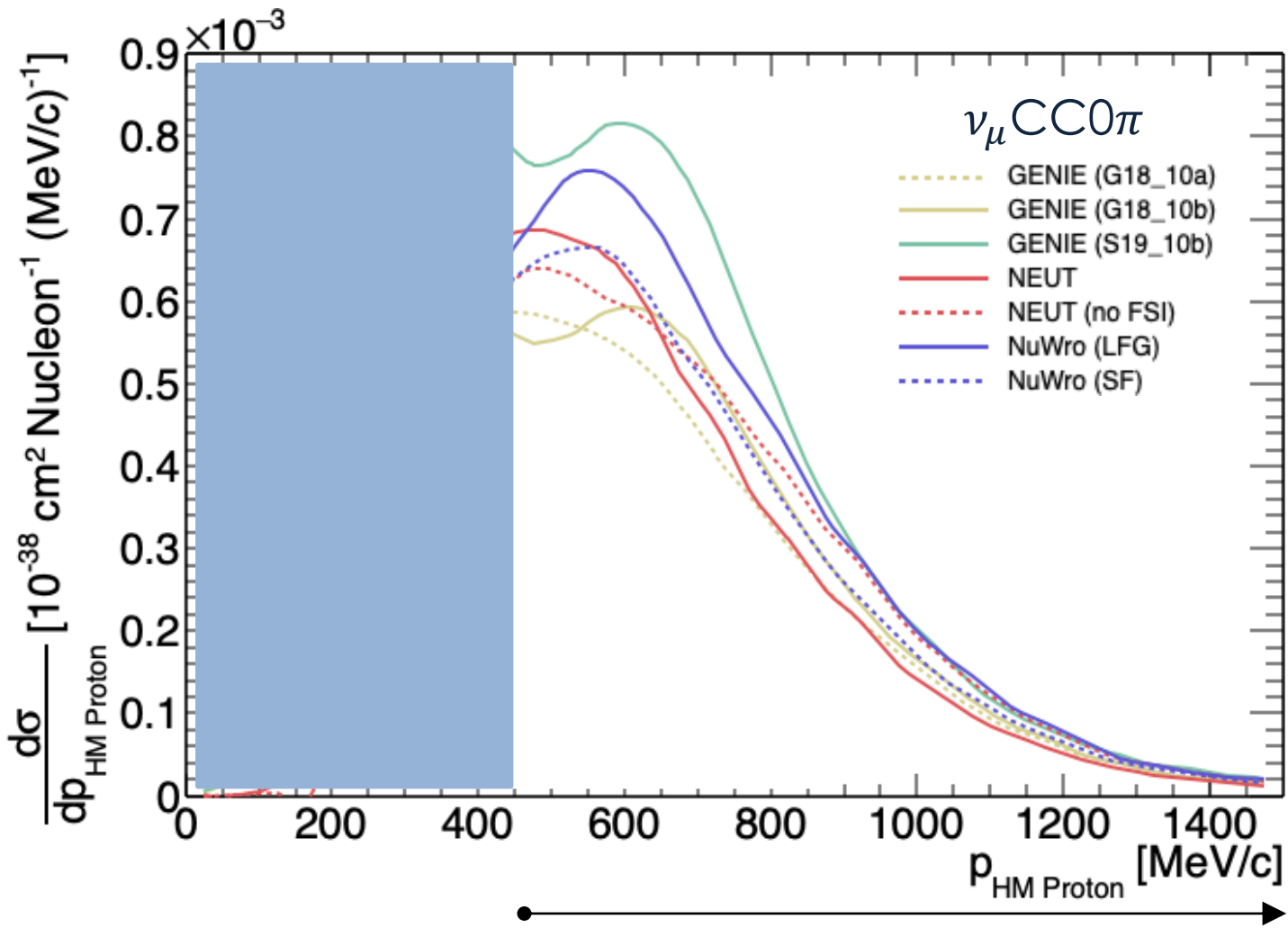
# Which observables?



Correlations between the muon and proton kinematics allow us to disentangle nuclear effects from neutrino energy

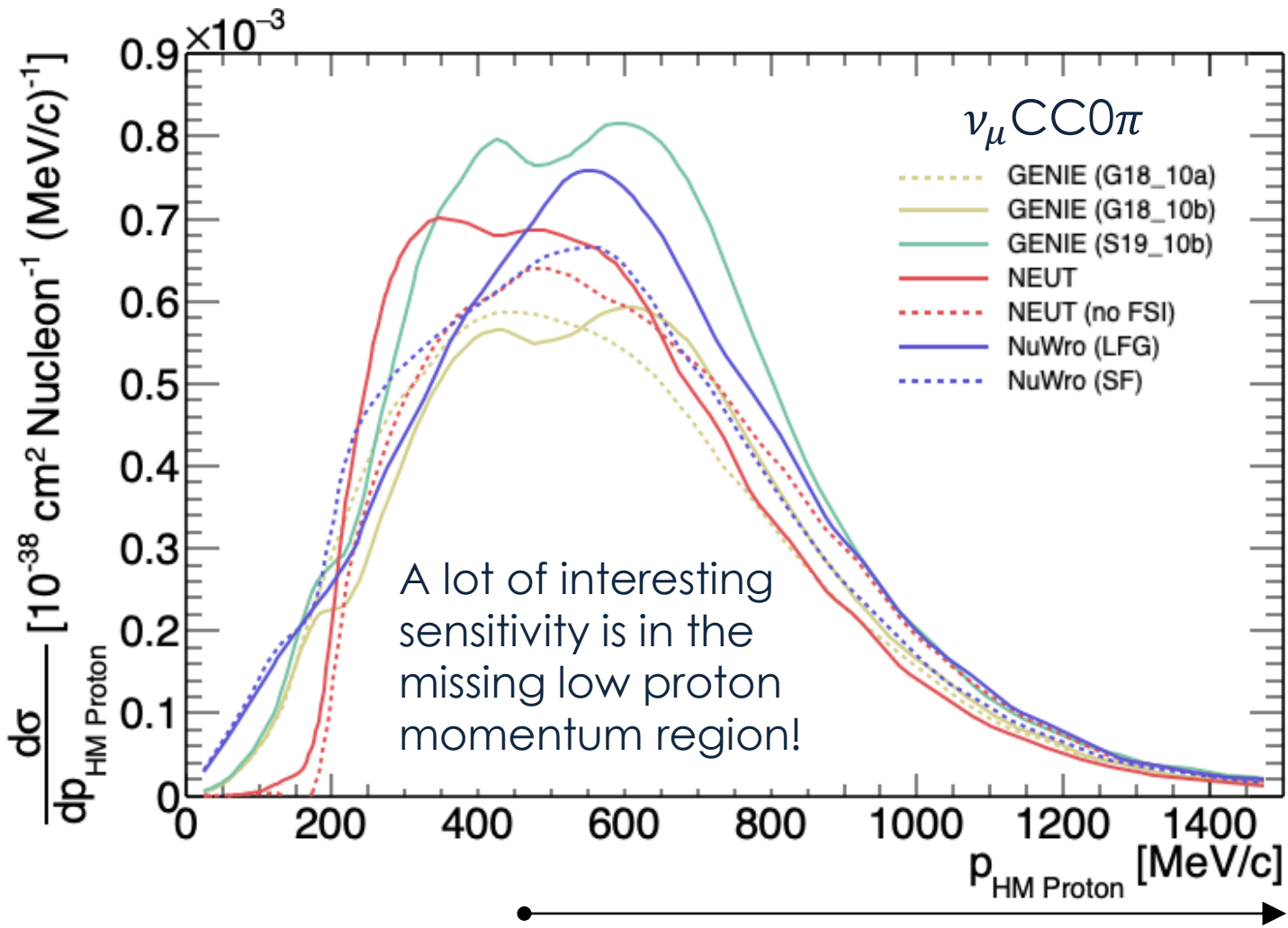
However, limited detector acceptance means low statistics and that we don't see the full story ...

# Proton momentum thresholds

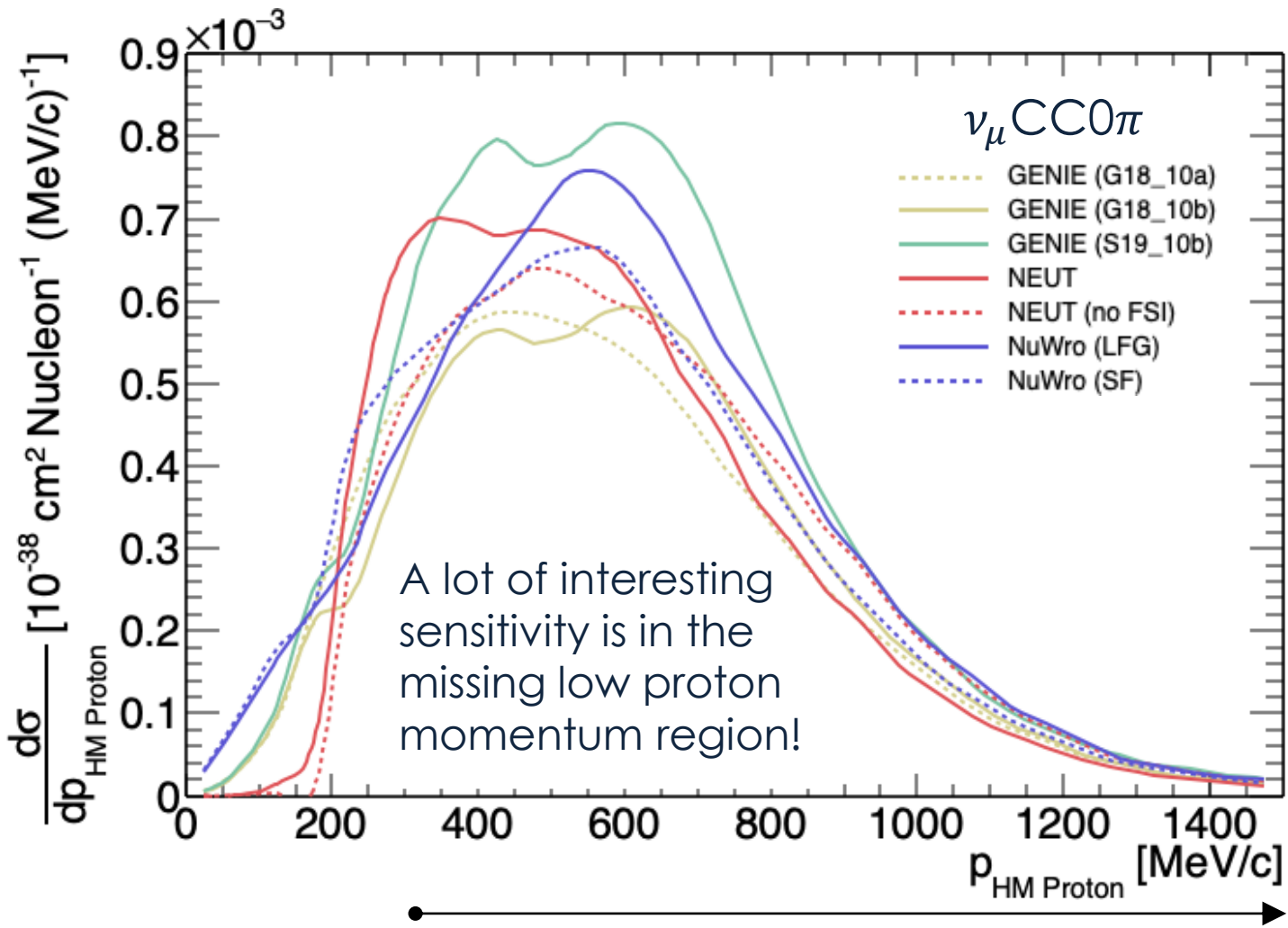


The range of proton momentum current experiments are sensitive to

# Proton momentum thresholds



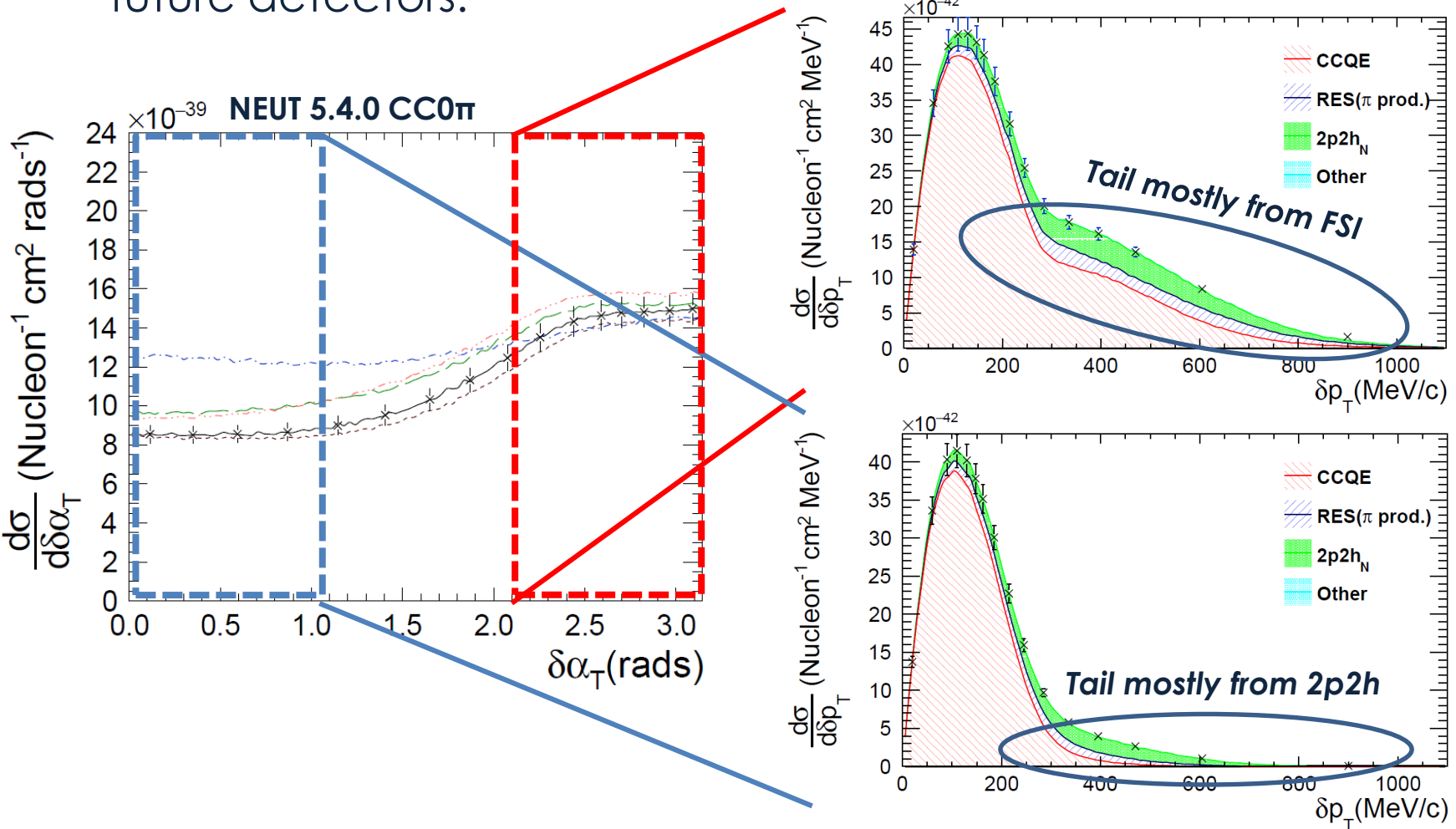
# Proton momentum thresholds



●————→  
 The range of proton momentum ND280 Upgrade and LAr detectors can access

# Enhanced sensitivity

- Measuring  $\delta p_T$  in bins of  $\delta\alpha_T$  may allow excellent separation of nuclear effects - makes use of high statistics / low thresholds from future detectors.



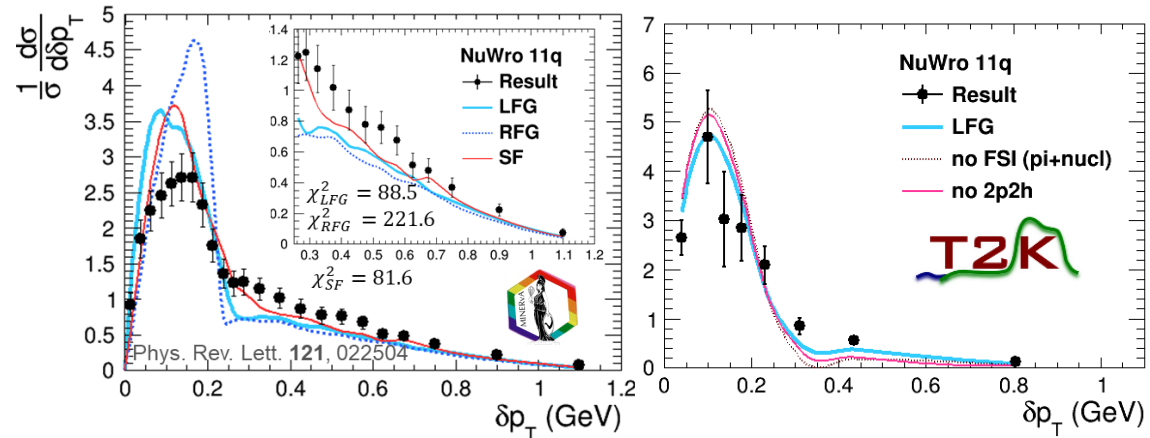
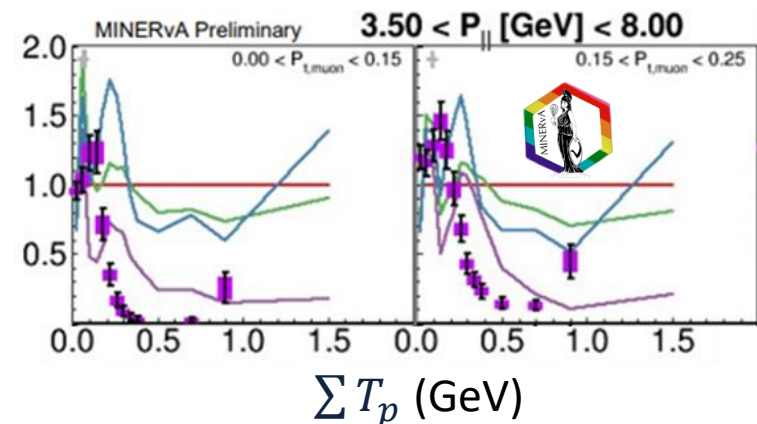
# Summarising current measurements

- ✓ T2K, MINERvA and others have made a wide range of innovative cross-section measurements aimed to target the nuclear physics most pertinent to future oscillation analyses

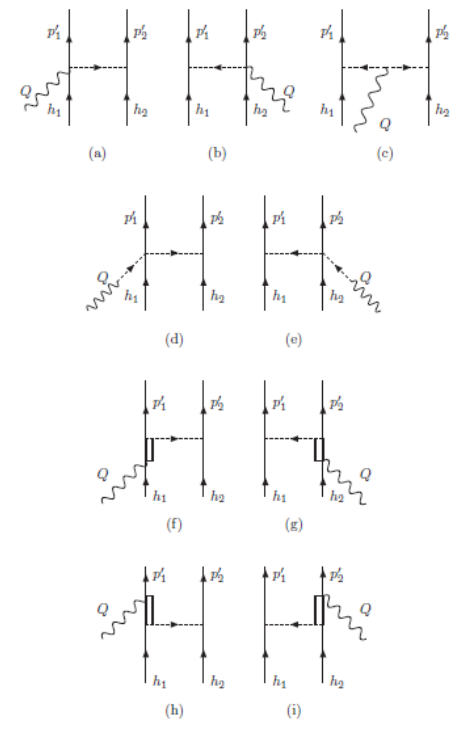
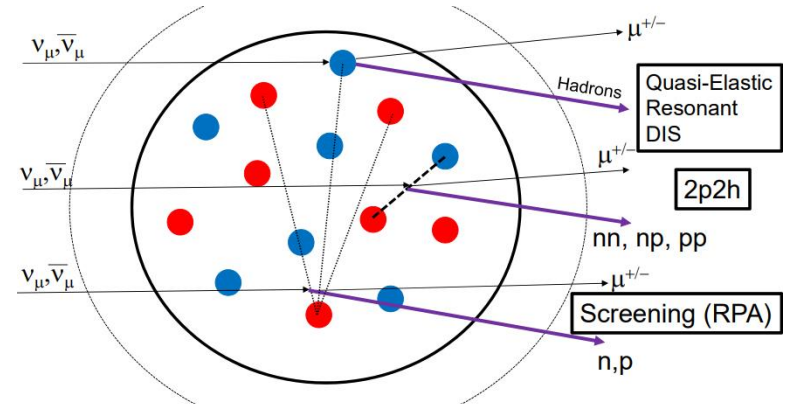
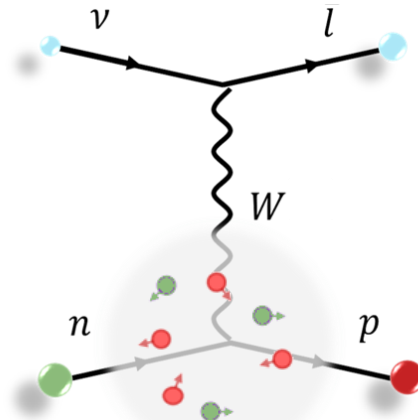
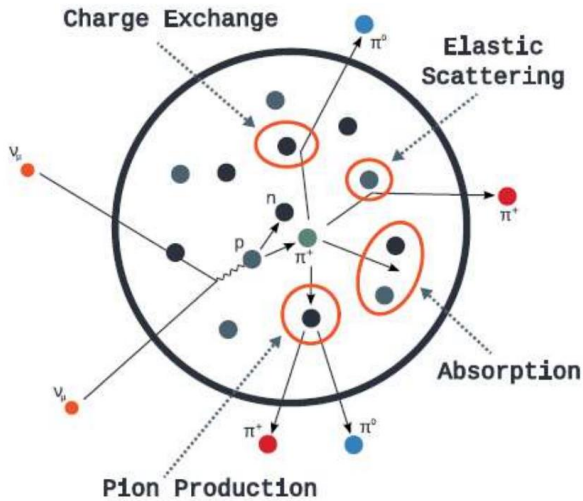
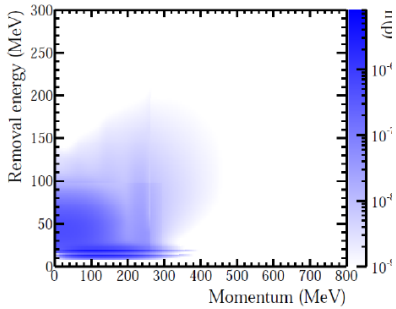
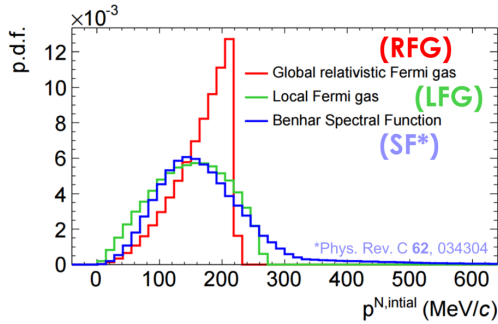


# Summarising current measurements

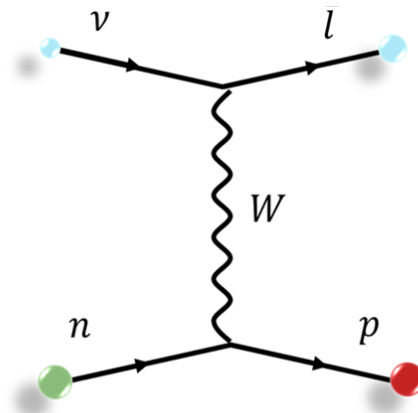
- ✓ T2K, MINERvA and others have made a wide range of innovative cross-section measurements aimed to target the nuclear physics most pertinent to future oscillation analyses
- X None of our current simulations are able describe more than the lepton kinematics ...
- It's critical to continue making new measurements with better detectors!



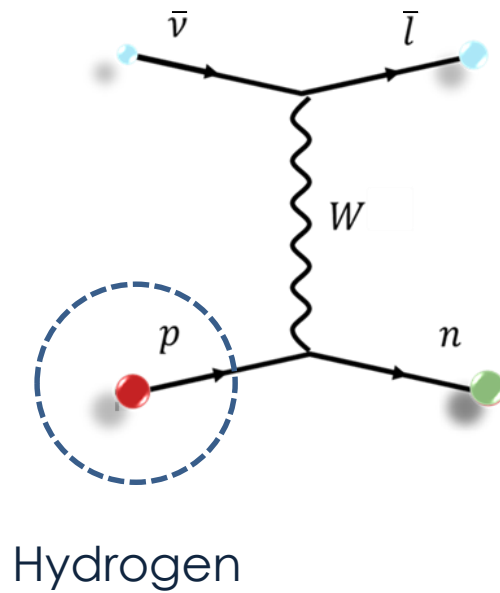
# Nuclear targets are hard ...



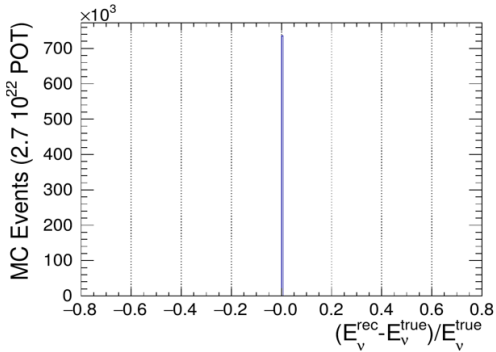
... nucleon targets are not



# ... nucleon targets are not

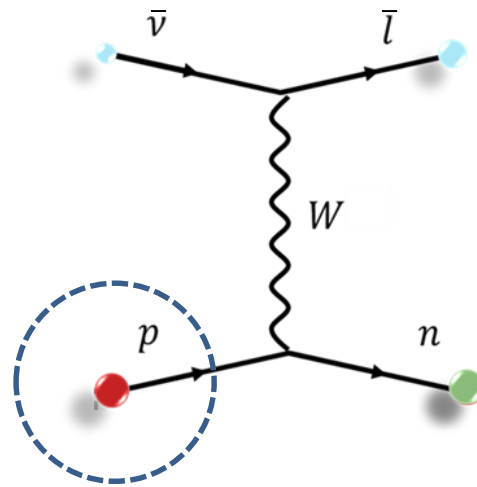


# ... nucleon targets are not



$$E_\nu = \frac{m_p^2 - (m_n - E_b)^2 - m_\mu^2 + 2(m_n - E_b)E_\mu}{2(m_n - E_b - E_\mu + p_\mu \cos \theta_\mu)}$$

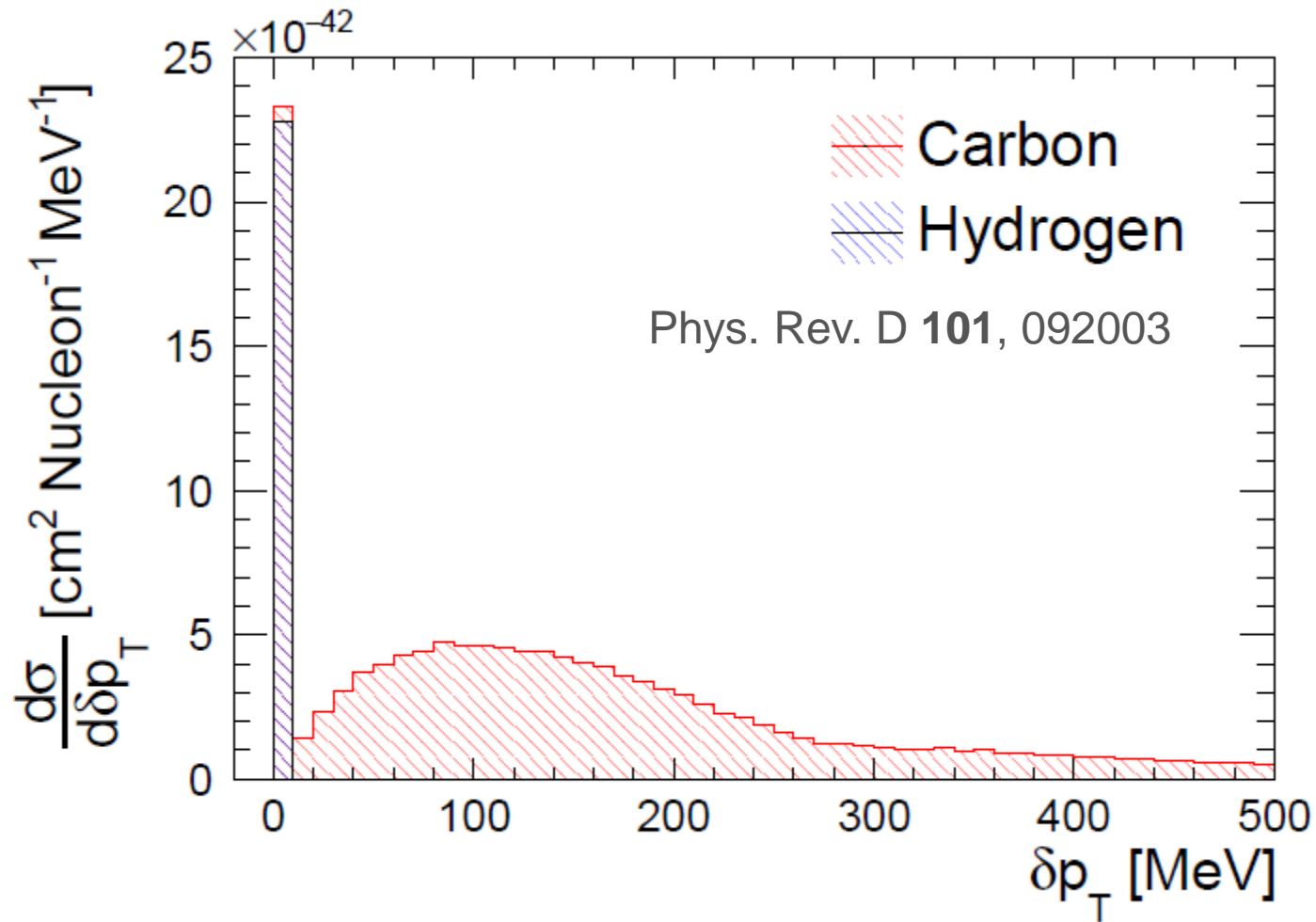
Proxy for  $E_\nu$  from lepton kinematics is perfect only for **CCQE elastic scattering** off a **stationary nucleon**



Hydrogen

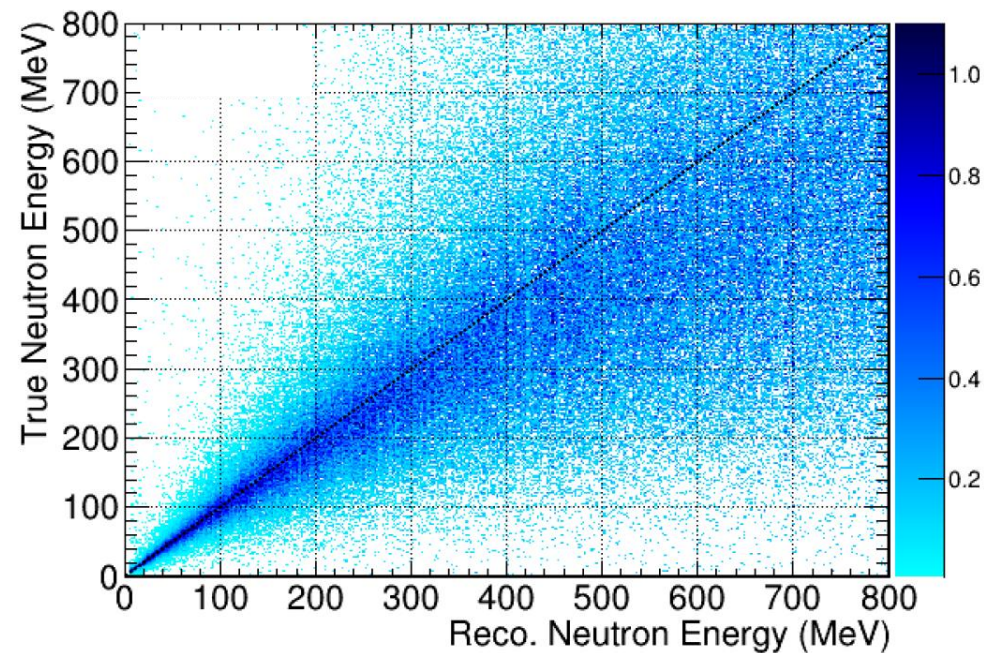
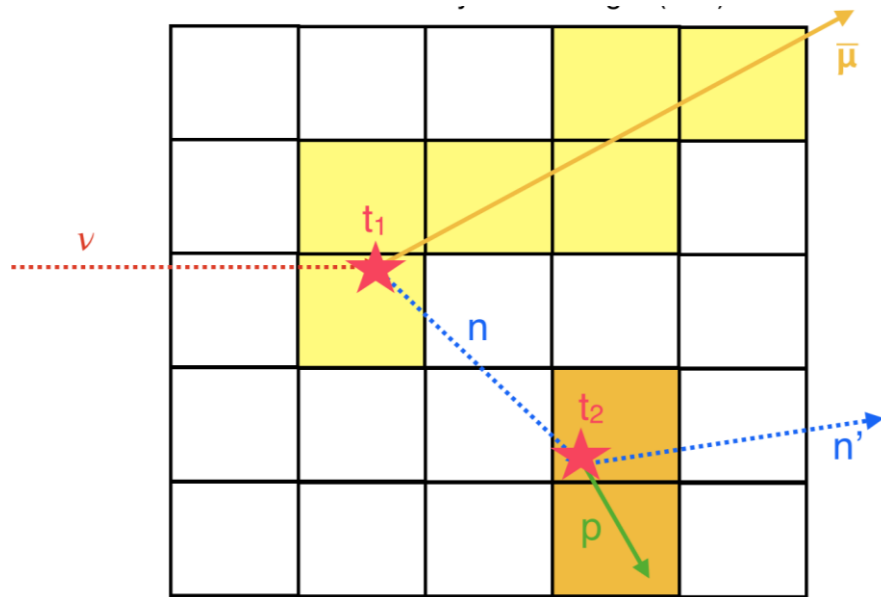


# Separating H in CH (scintillator)



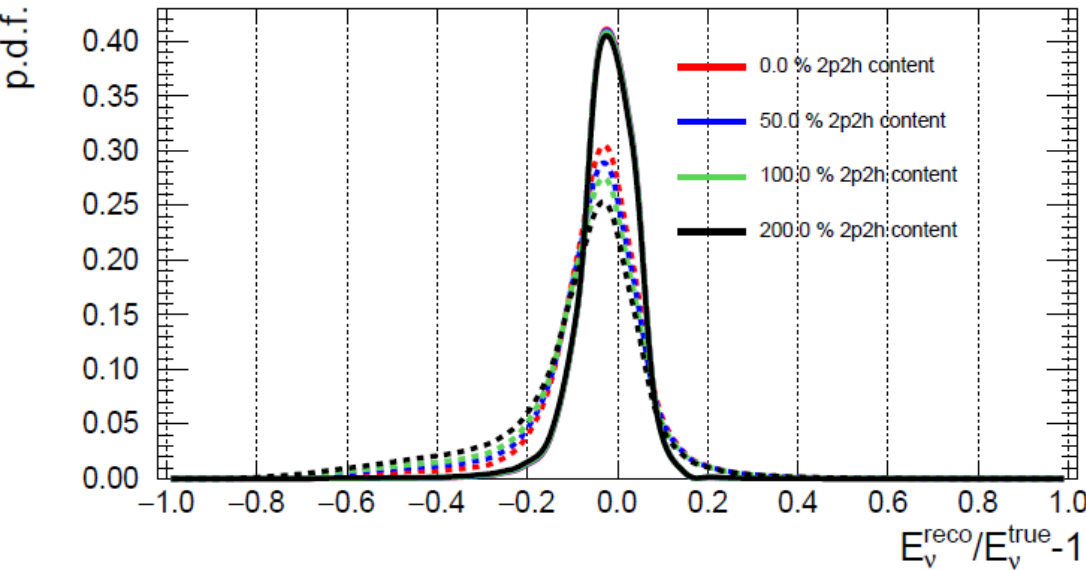
# Neutrons with a 3D scintillator detector

- Can look for neutrons via their re-interaction within a detector
- If the path is long enough ( $>20$  cm) neutron energy is measured using the time of flight with resolution 15-30% (for  $\sim 1$  ns timing resolution)

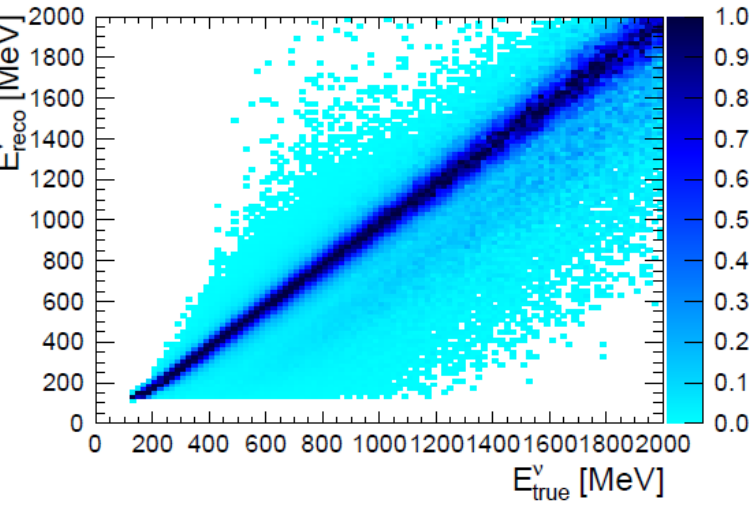


Phys. Rev. D **101**, 092003

# Reconstructing the anti-neutrino flux



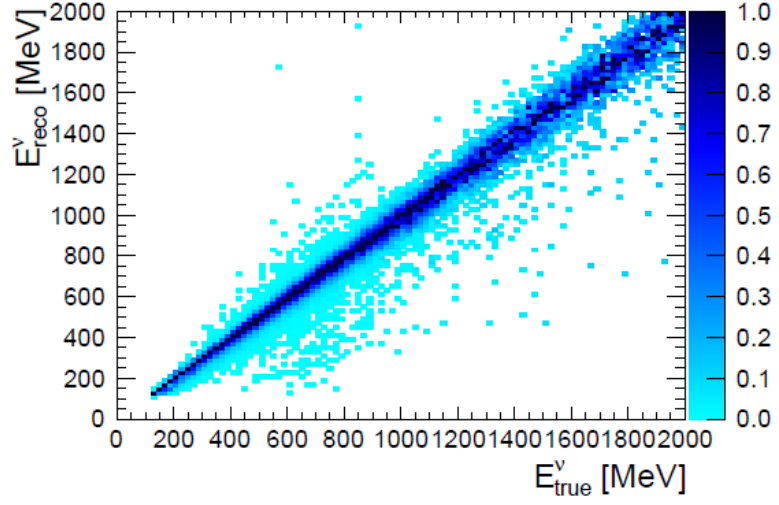
- Improved  $E_\nu$  resolution
  - RMS: 15%  $\rightarrow$  7%
- Insensitive to nuclear effects (even the carbon events are the ones that aren't much affected)
- In-situ anti-neutrino flux measurement!



$\delta p_T > 40 \text{ MeV}$

**➔**

lever arm  $> 10 \text{ cm}$



Phys. Rev. D **101**, 092003



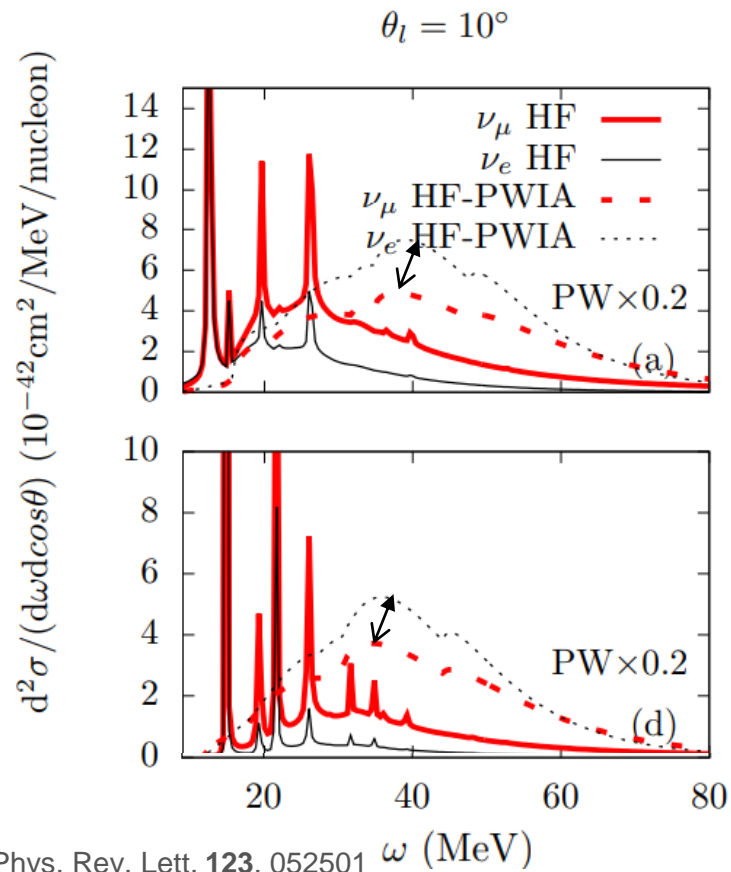
# Summary

- To avoid future oscillation analyses becoming **pre-maturely limited by systematic uncertainties**, *it is essential to better understand nuclear effects in neutrino-nucleus interactions*
- Nuclear effects can only be effectively understood by **measuring outgoing lepton and nucleon kinematics**
- Current measurements are limited by detector capabilities, but already **indicate a poor modelling of the essential processes**
- **Future measurements of nucleons** from neutrino interactions and **continued collaboration with the nuclear theory community** are critical
- Measuring neutrons from neutrino interactions offers the possibility of interesting new measurements

# Backups

# Nuclear effects and $\sigma(\nu_e)/\sigma(\nu_\mu)$

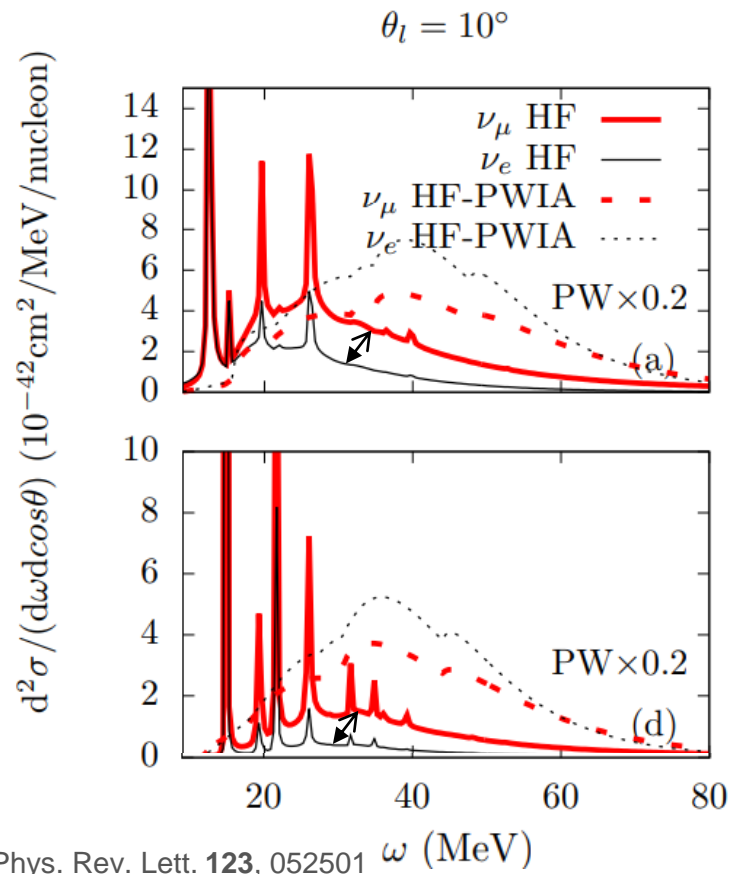
- Ratio of  $\nu_e$  to  $\nu_\mu$  critical for future oscillation analyses
  - Measure  $\nu_\mu$  at ND but need to know about  $\nu_e$  to measure  $\delta_{CP}$
- This is also subject to subtleties in the nuclear physics...



- If the outgoing nucleon exits the nucleus as a “plane wave” (no FSI):  $\sigma(\nu_e) > \sigma(\nu_\mu)$

# Nuclear effects and $\sigma(\nu_e)/\sigma(\nu_\mu)$

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  - Measure  $\nu_\mu$  at ND but need to know about  $\nu_e$  to measure  $\delta_{CP}$
- This is also subject to subtleties in the nuclear physics...



- If the outgoing nucleon exits the nucleus as a “plane wave” (no FSI):  $\sigma(\nu_e) > \sigma(\nu_\mu)$
- If the outgoing nucleon is distorted by the nuclear potential (FSI):  $\sigma(\nu_e) < \sigma(\nu_\mu)$

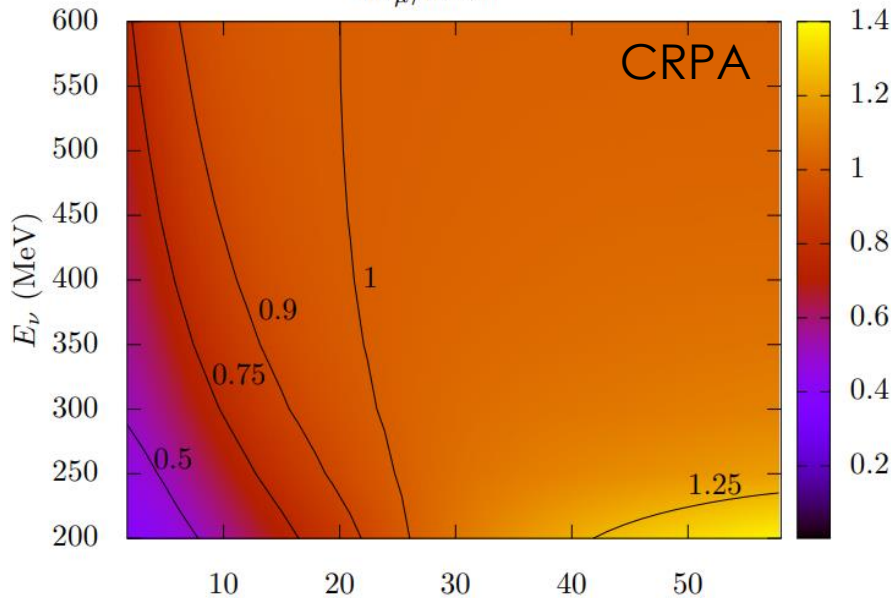
# Nuclear effects and $\frac{d\sigma_e/d\cos\theta}{d\sigma_\mu/d\cos\theta}$

- Different models can predict quite different cross section ratios!
- Important for T2K/HK!

	$E_\nu = 200 \text{ MeV}$		$E_\nu = 600 \text{ MeV}$	
<b>Model</b>	5°	60°	5°	60°
RFG (w/PB)	0.64	1.61	0.97	1.03
SF (full)	1.41	1.92	1.04	1.03
CRPA	~0.5	~1.4	~0.9	~1.0

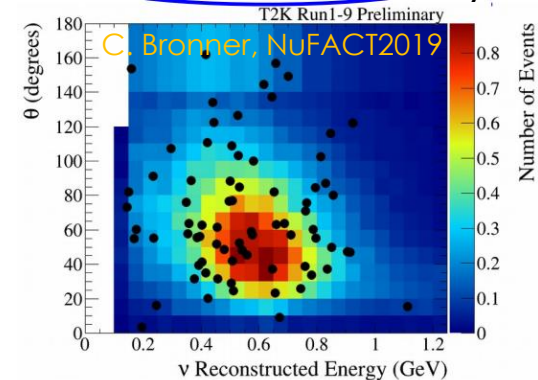
Tabulated from Phys. Rev. C **96**, 035501 and the left figure

$$\frac{d\sigma_e/d\cos\theta}{d\sigma_\mu/d\cos\theta}$$

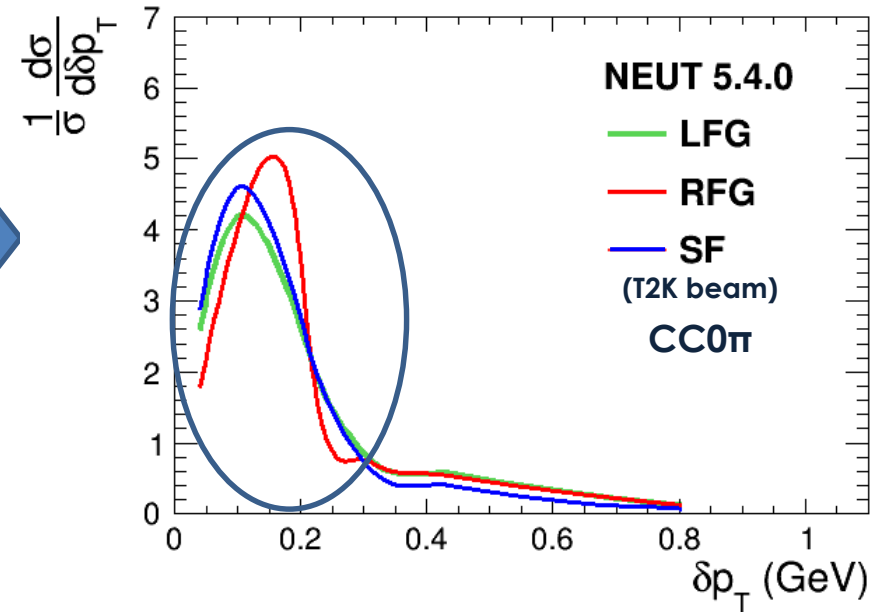
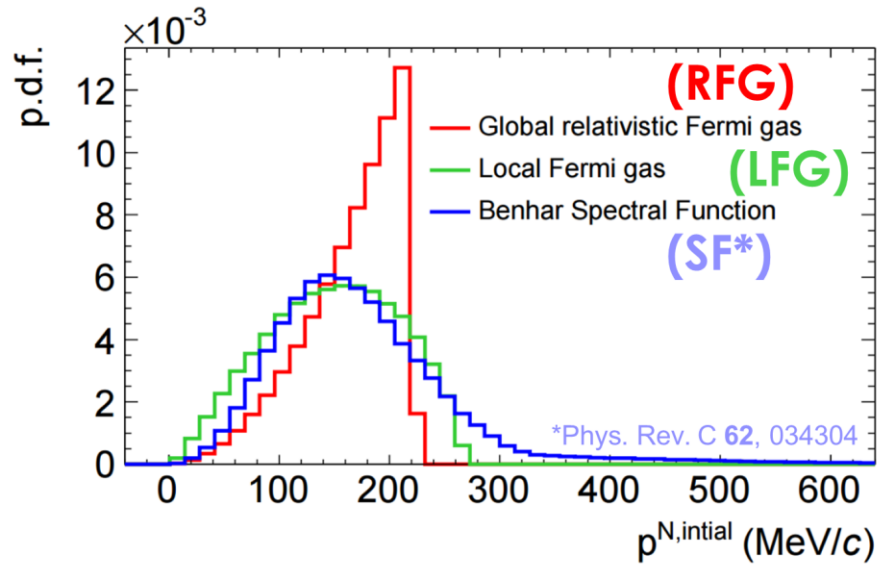


Phys. Rev. Lett. **123**, 052501  $\theta_l$  (degrees)

These differences are predicted in regions that are relevant to T2K/HK oscillation analyses

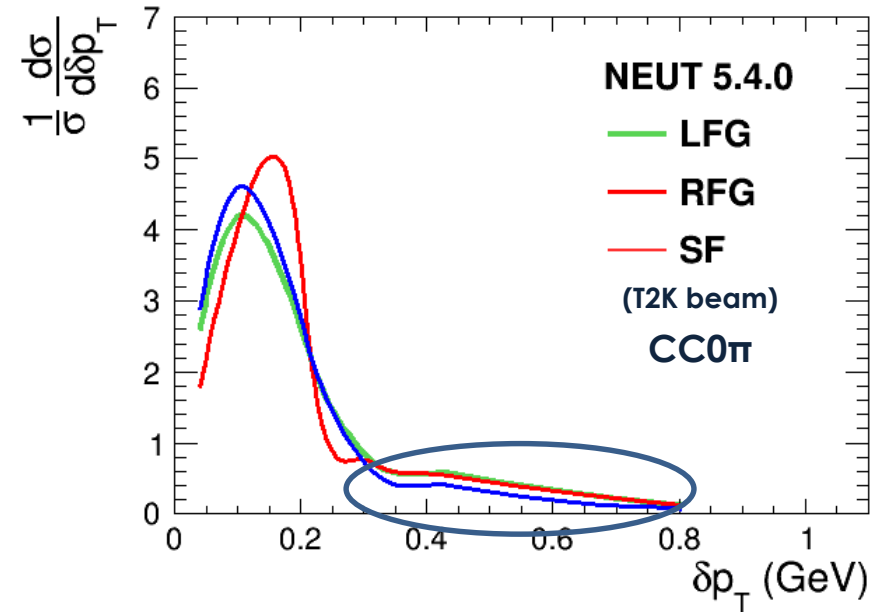
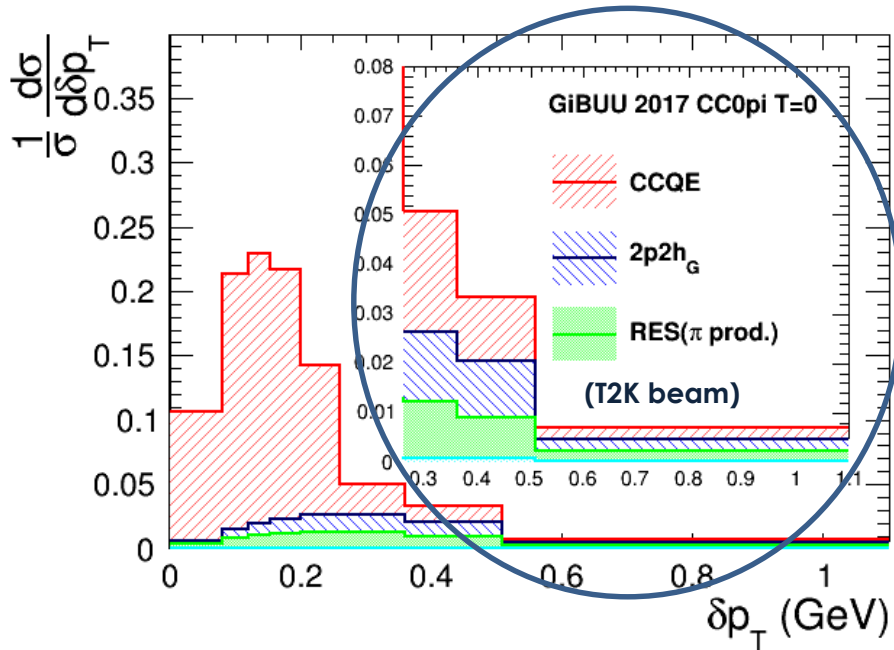


# STV model discrimination - $\delta p_T$



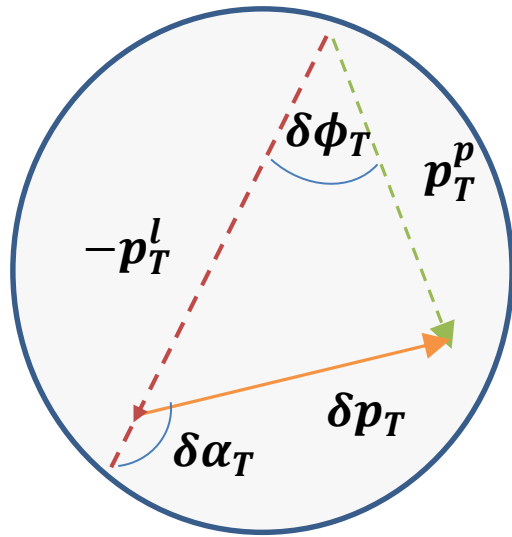
- In the absence of other nuclear effects,  $\delta p_T$  is the transverse projection of the Fermi motion.
- Since this motion is isotropic,  $\delta p_T \rightarrow$  Fermi motion

# STV model discrimination - $\delta p_T$

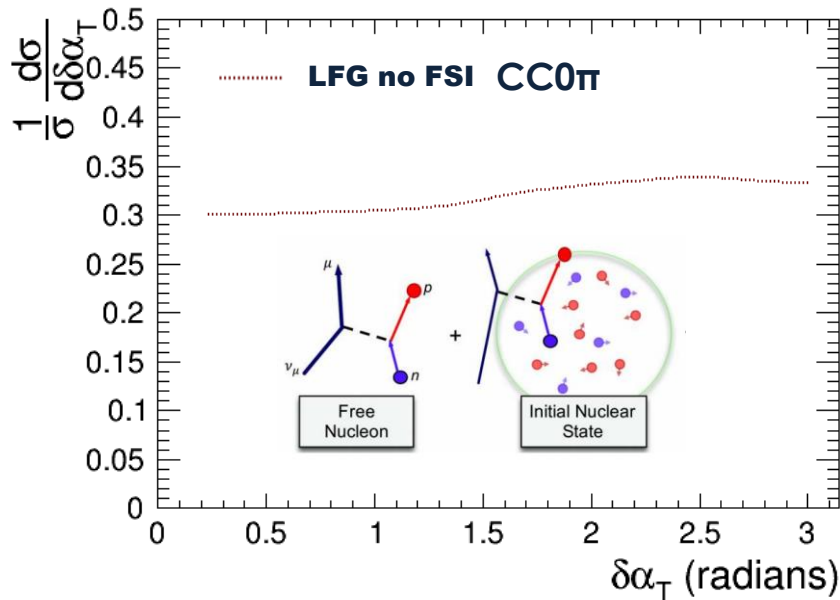


- In the absence of other nuclear effects,  $\delta p_T$  is the transverse projection of the Fermi motion.
- Since this motion is isotropic,  $\delta p_T \rightarrow$  Fermi motion
- Cross section beyond the Fermi surface must come from physics beyond RFG  $\rightarrow$  2p2h, FSI, SRCs ...

# STV model discrimination - $\delta\alpha_T$



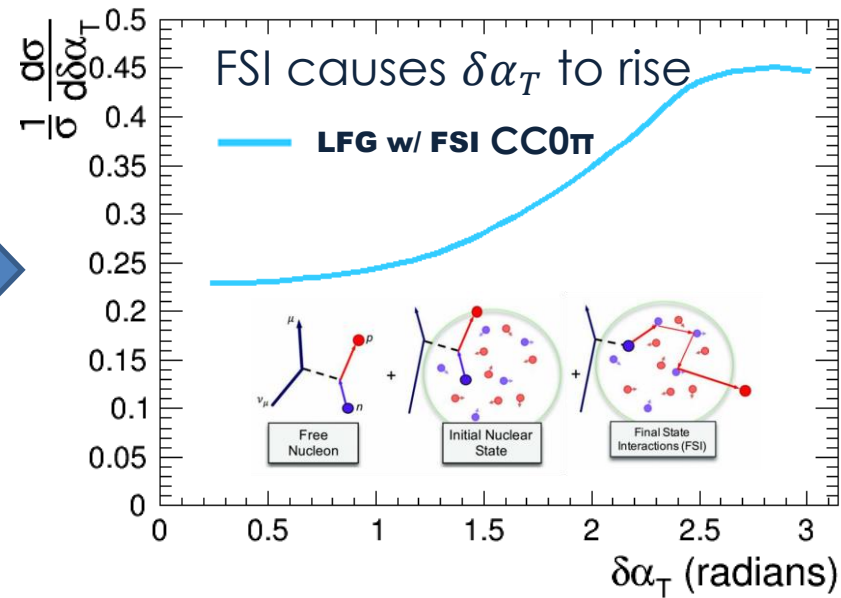
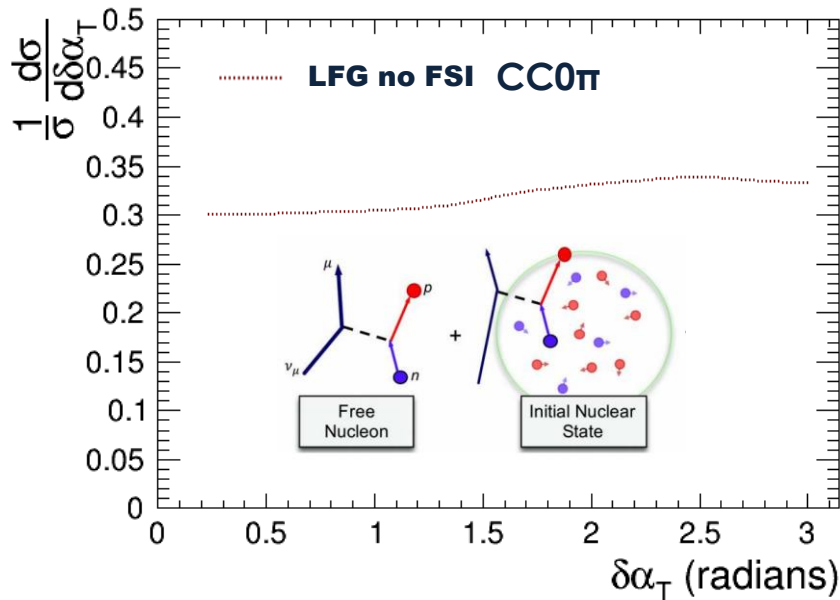
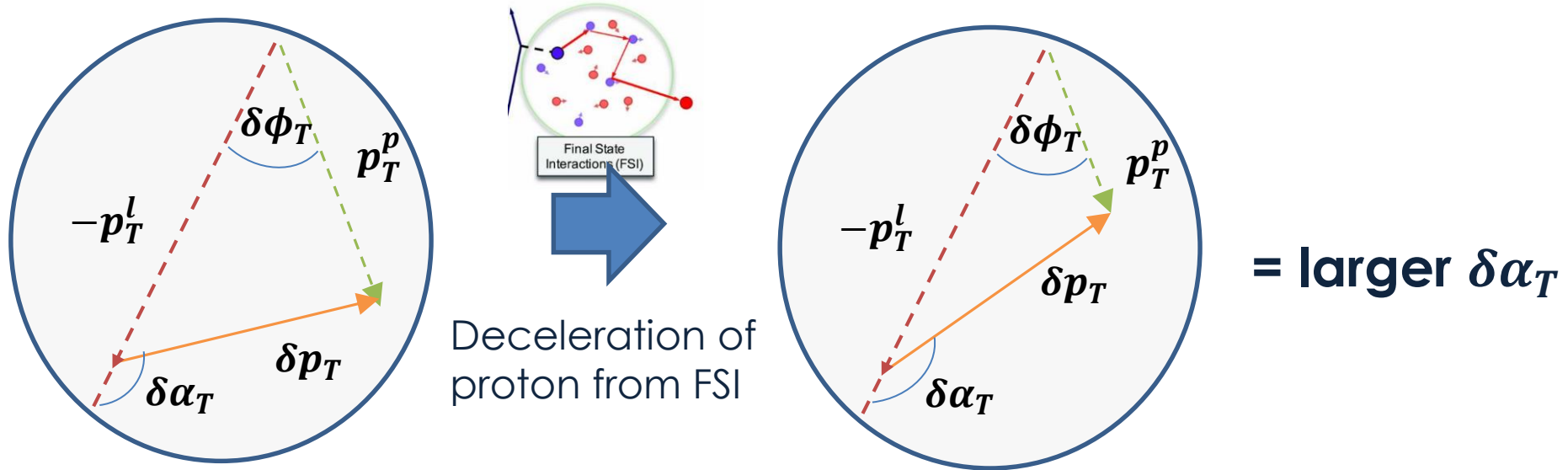
Consider imbalance from only Fermi motion



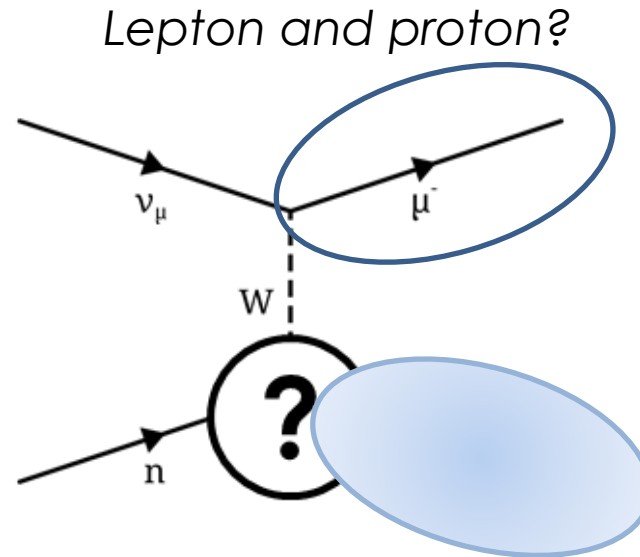
Fermi motion is isotropic so no preferred  $\delta\alpha_T$  direction



# STV model discrimination - $\delta\alpha_T$



# Which observables?



Correlations between the muon and proton kinematics allow us to disentangle nuclear effects from neutrino energy

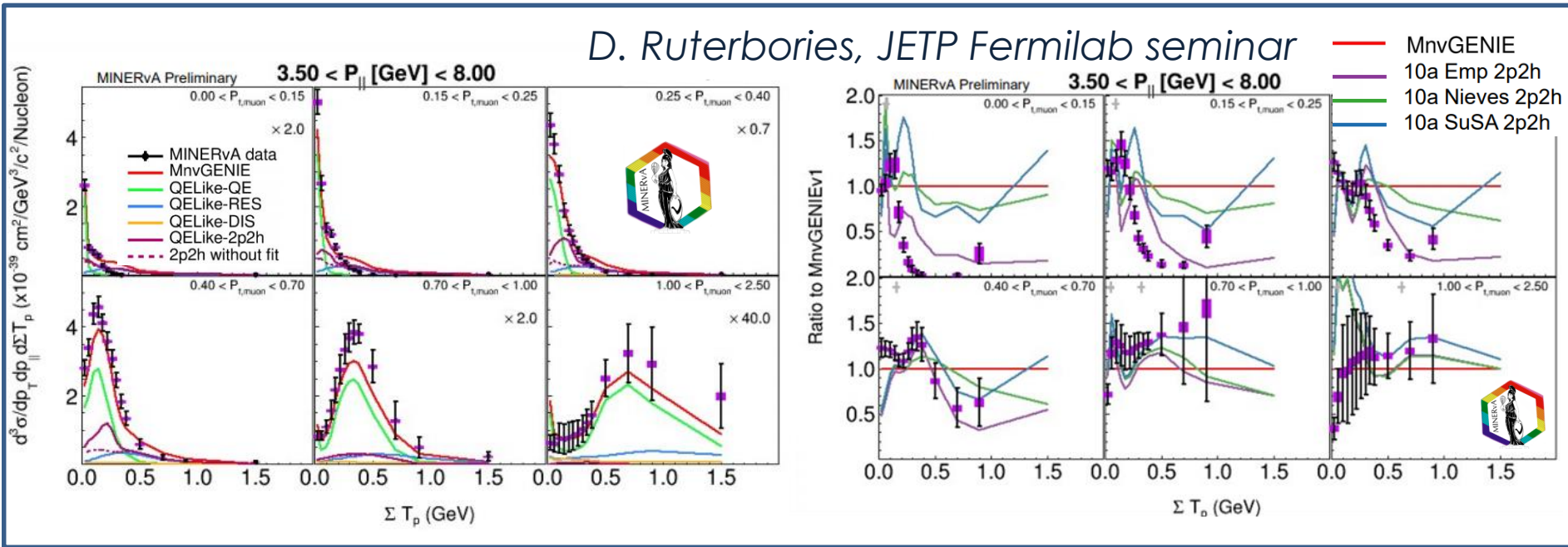
However, limited detector acceptance means low statistics and that we don't see the full story ...

To avoid this, we can rely on calorimetric methods (no need to reconstruct individual protons)

# Calorimetric measurements

- Sum energy deposited in the detector not associated with lepton
  - Reject pions: excess energy is  $\Sigma T_{proton}$
- Measure  $d^3\sigma/dp_T^\mu dp_{||}^\mu d\Sigma T_p$

D. Ruterbories, JETP Fermilab seminar



- **Best model agreement with data:**  $\chi^2 = 5062/238$
- Worrisome for calorimetric reconstruction of  $E_\nu$  (DUNE/SBN/NOvA) ...

End