
CP violation and determination of the “squashed” (b,s) unitarity triangle at FCC-ee

R. Aleksan (CEA/IRFU), L. Oliver (IJCLab), E.Perez (CERN)

3rd Workshop FCC France, Nov 30 – Dec 2, 2021, Annecy

- CP violation and determination of the bs "flat" unitarity triangle at FCC-ee, <https://arxiv.org/abs/2107.02002>
- Study of CP violation in B^\pm decays to $D^0 (\bar{D}^0) K^\pm$ at FCC-ee <https://arxiv.org/abs/2107.05311>

Unitarity triangles

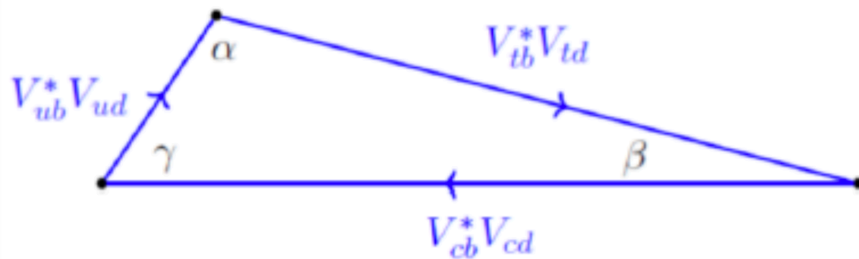
Six triangular relations from unitarity of V_{CKM} .
Among them :

$$V_{CKM} = \begin{bmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{bmatrix}$$

The (usual) “(b, d) triangle” :

$$V_{ub}^* V_{ud} + V_{cb}^* V_{cd} + V_{tb}^* V_{td} = 0$$

$$(\lambda^3, \lambda^3, \lambda^3)$$

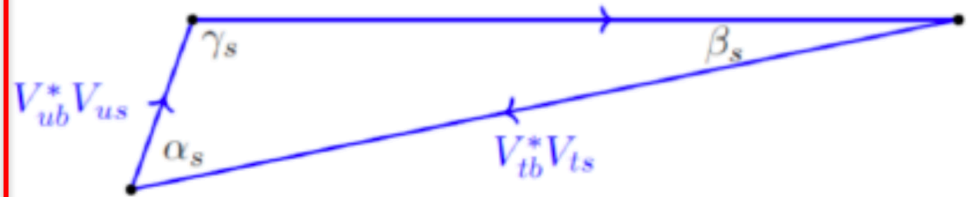


All angles are (quite) large.
Extensively studied experimentally

The “(b, s) triangle” :

$$V_{ub}^* V_{us} + V_{cb}^* V_{cs} + V_{tb}^* V_{ts} = 0$$

$$(\lambda^4, \lambda^2, \lambda^2)$$



$$(\alpha_s, \beta_s, \gamma_s) \sim (67^\circ, 1^\circ, 111^\circ)$$

Can be measured directly at FCC-ee

The “(d, s) triangle”: even more squashed
(SM: Relations between the angles of these triangles)

- Study expected precision at FCC-ee on $(\alpha_s, \beta_s, \gamma_s)$
- Requirements on the detectors from these measurements

Experimental sensitivities and detector response

- Expected sensitivities given for 10^{11} produced \bar{B}_s and $3.9 \cdot 10^{11}$ produced B^+ , corresponding to 150 ab^{-1} at FCC-ee at the Z peak.

- Modelisation of the detector response :

$$\text{Acceptance : } |\cos \theta| < 0.95$$

$$\text{Track } p_T \text{ resolution : } \frac{\sigma(p_T)}{p_T^2} = 2. \times 10^{-5} \oplus \frac{1.2 \times 10^{-3}}{p_T \sin \theta}$$

- Smearing of the momenta and angles of particles in the decays of interest

$$\text{Track } \phi, \theta \text{ resolution : } \sigma(\phi, \theta) \text{ } \mu\text{rad} = 18 \oplus \frac{1.5 \times 10^3}{p_T \sqrt[3]{\sin \theta}}$$

- Parametrisation based on typical performance of a light tracker at a future ee detector
- Excellent EM calo resolution
- Used for most results shown here

$$\text{Vertex resolution : } \sigma(d_{\text{Im}}) \text{ } \mu\text{m} = 1.8 \oplus \frac{5.4 \times 10^1}{p_T \sqrt{\sin \theta}}$$

$$\text{Vertex resolution : } \langle \sigma(d_{\text{Im}}) \rangle \text{ bachelor K in } D_s K$$

$$\langle \sigma(d_{\text{Im}}) \rangle \simeq 10 \text{ } \mu\text{m}$$

$$\text{Calorimeter resolution : } \frac{\sigma(E)}{E} = \frac{3 \times 10^{-2}}{\sqrt{E}} \oplus 5 \times 10^{-3}$$

- Common SW : Full MC events + response of the IDEA detector with DELPHES
 - Detailed description of tracks, accounting for multiple scattering
 - Genuine vertex fitting

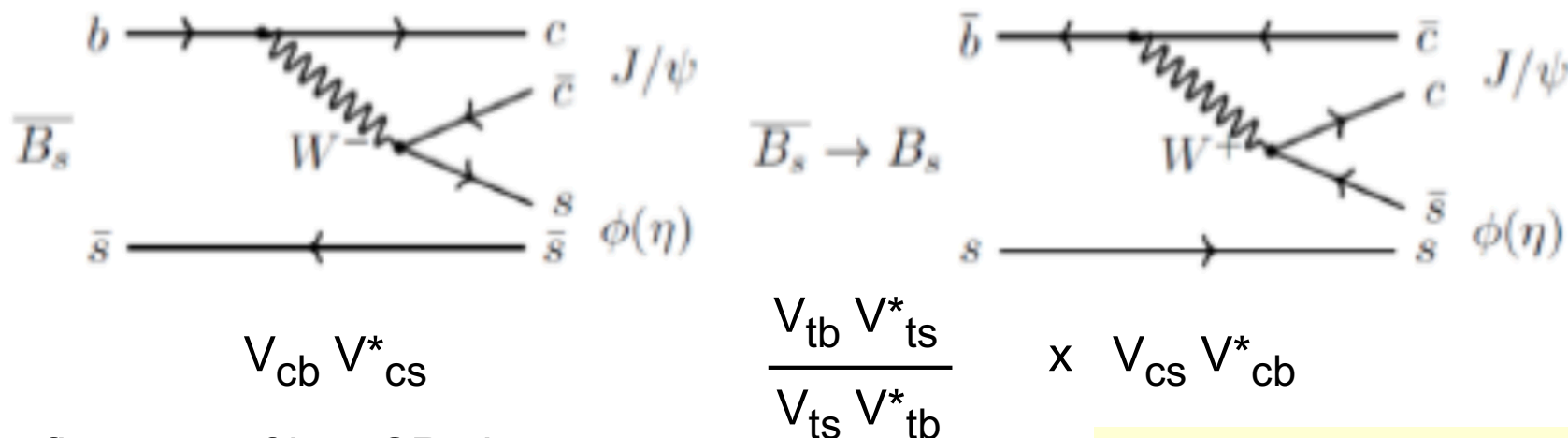
Measurement of β_s

$$\beta_s = \arg \left(-\frac{V_{tb}^* V_{ts}}{V_{cb}^* V_{cs}} \right)$$

β_s : very small in the SM (1 degree), and known precisely.
Hence can set strong constraints on New Physics.

- **Golden channel:** $B_s \rightarrow J/\psi \phi$
 - Measure CP violation in the interference between B_s mixing and $b \rightarrow ccs$ decay
- Analogous of $B_d \rightarrow J/\psi K_s$ from which $\sin(2\beta)$ is extracted
- Already largely used at LHC, but low precision so far:

$$\text{PDG: } \beta_s = (0.60 \pm 0.89)^\circ$$



When final state f is a CP eigenstate:

$$\Phi_{\text{CKM}} (J/\psi \phi) = 2 \beta_s (+\pi)$$

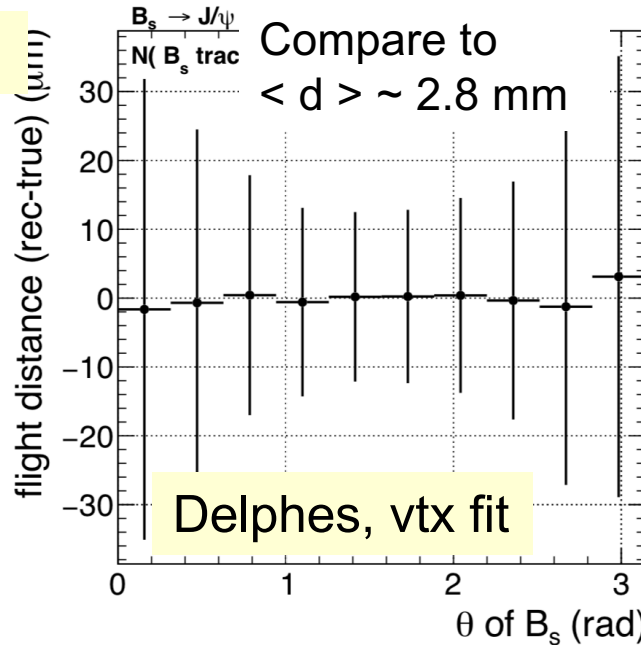
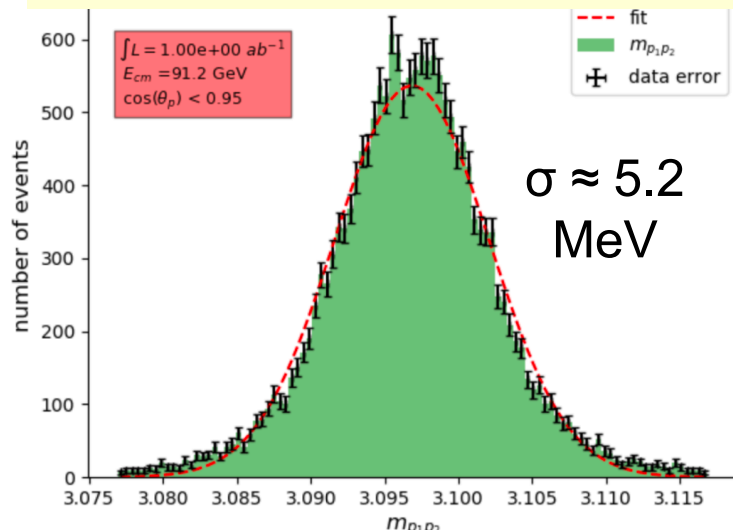
$$\Gamma \left(\overline{B}_s(t) \rightarrow f \right) \sim e^{-\Gamma t} \left[1 \mp \eta_{f,\text{CP}} (1-2\omega) \sin \Phi_{\text{CKM}} \sin(\Delta m_s t) \right]$$

ω = mistag rate, (well) measured independently, see later

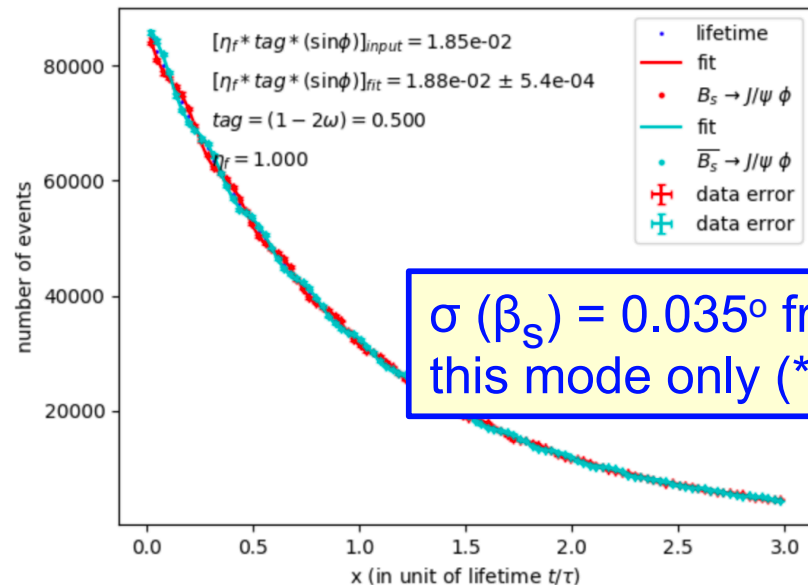
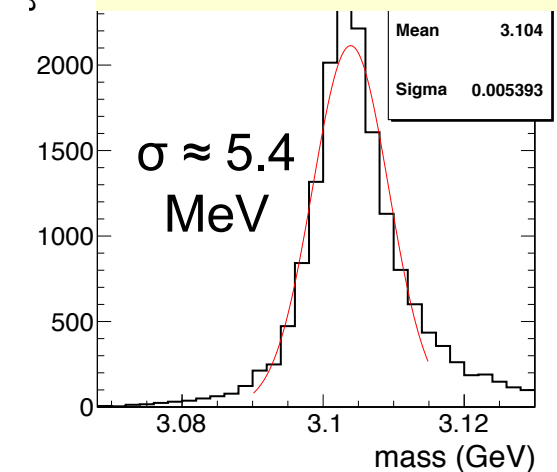
Measurement of β_s

Final state: two muons from J/ψ + two kaons from Φ

Parametrized response, $m(J/\psi)$



Delphes, $m(J/\psi)$ after vtx reco



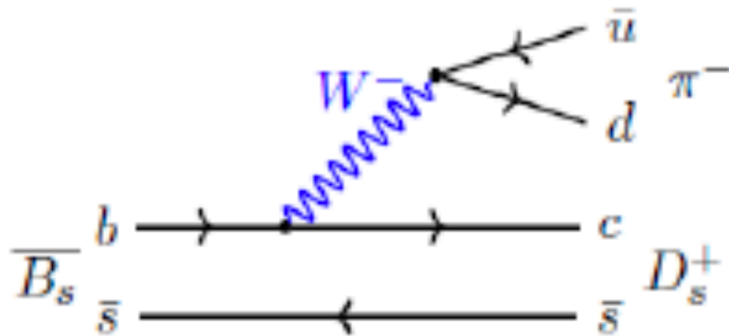
Statistics: 3 M B_s decays, large enough to measure precisely the small modulation.

β_s will be known very precisely !

- O(25x) better than current average
- O(7 x) better than expected from upgrades of LHCb/BelleII (late 2030s)

(*) $J/\psi \Phi$: need an angular analysis as different polarisations of VV have a different CP

Measurement of the mistag rate: $\bar{B}_s \rightarrow D_s^+ \pi^-$

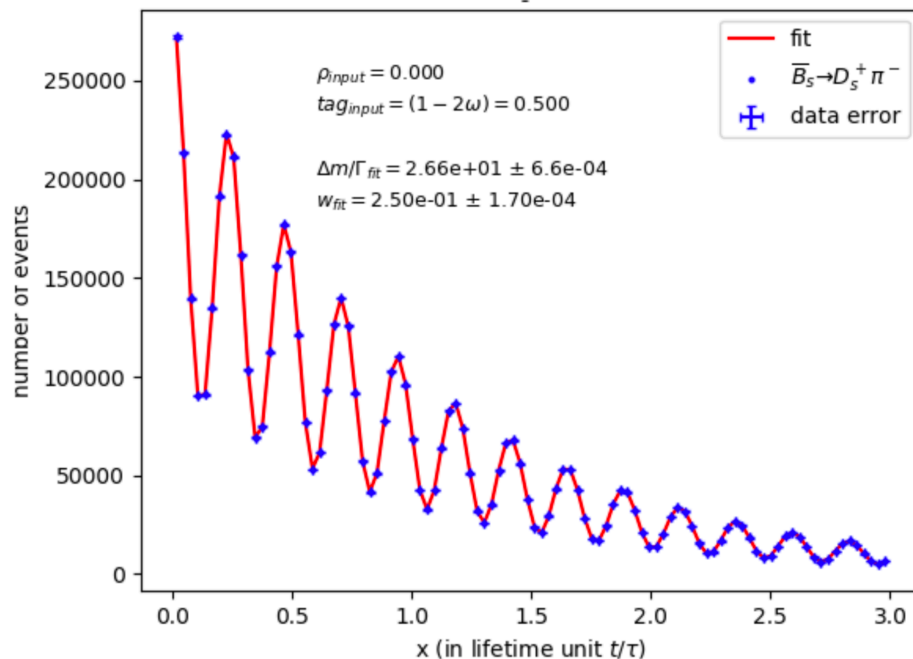


- Relatively large BR, 0.3 %
 - Expect 14 M such decays
- No diagram for $B_s \rightarrow D_s^+ \pi^-$, i.e. flavour-specific decay and **no CP violation** in this mode.

Very convenient for probing the B_s tagging :

$$\Gamma(\bar{B}_s(t) \rightarrow D_s^+ \pi^-) \sim e^{-\Gamma t} \left[(1 - \omega) \cos^2 \Delta m t / 2 + \omega \sin^2 \Delta m t / 2 \right]$$

$$\Gamma(B_s(t) \rightarrow D_s^+ \pi^-) \sim e^{-\Gamma t} \left[\omega \cos^2 \Delta m t / 2 + (1 - \omega) \sin^2 \Delta m t / 2 \right]$$

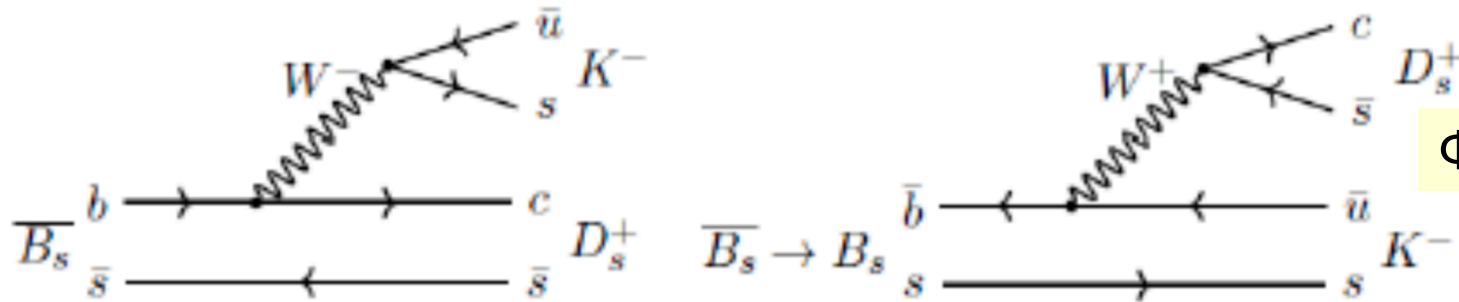


Very precise determination of the mistag rate,
 $\sigma(\omega) \sim 1.3 \cdot 10^{-4}$

And an improved determination of Δm_s ,
 $\Delta m_s / \Gamma_s$ to $5 \cdot 10^{-4}$
 i.e. O(100x) better than current PDG.

Measurement of α_s : B_s to $D_s K$

$$\alpha_s = \arg \left(-\frac{V_{ub}^* V_{us}}{V_{tb}^* V_{ts}} \right)$$



$$\Phi_{\text{CKM}} = \pi - (\alpha_s - \beta_s)$$

- Again CPV from the interference of mixing and decay
- But now have four time-dependent rates

$$\bar{B}_s \rightarrow D_s^+ K^-, \bar{B}_s \rightarrow D_s^- K^+, B_s \rightarrow D_s^+ K^-, B_s \rightarrow D_s^- K^+$$

$$\begin{aligned} \Gamma(B_s(t) \rightarrow f) = & |\langle f | B_s \rangle|^2 e^{-\Gamma t} \left\{ [1 - \omega(1 - \rho^2)] \cos^2 \frac{\Delta m t}{2} \right. \\ & + [\rho^2 + \omega(1 - \rho^2)] \sin^2 \frac{\Delta m t}{2} \\ & \left. - (1 - 2\omega)\rho \sin \phi_{CP}^+ \sin \Delta m t \right\} \end{aligned}$$

$$\begin{aligned} \Gamma(\bar{B}_s(t) \rightarrow \bar{f}) = & |\langle \bar{f} | \bar{B}_s \rangle|^2 e^{-\Gamma t} \left\{ [1 - \omega(1 - \rho^2)] \cos^2 \frac{\Delta m t}{2} \right. \\ & + [\rho^2 + \omega(1 - \rho^2)] \sin^2 \frac{\Delta m t}{2} \\ & \left. + (1 - 2\omega)\rho \sin \phi_{CP}^- \sin \Delta m t \right\} \end{aligned}$$

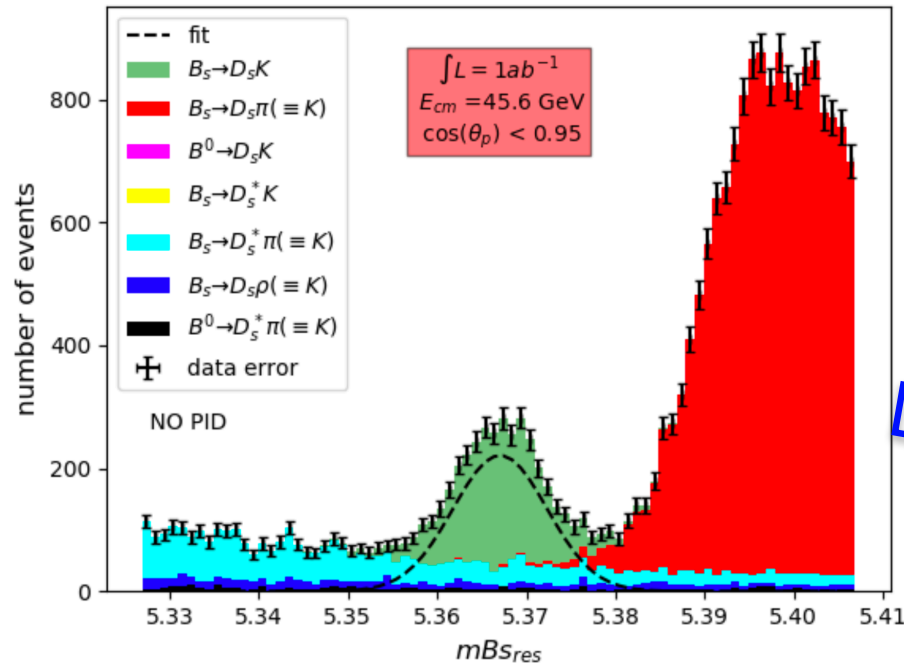
(+ two other equations not shown)

$$\Phi_{CP}^{\pm} = \Phi_{\text{CKM}} \pm \delta$$

δ = strong phase difference
 ρ = ratio of |amplitudes|

Hadronic parameters ρ and δ can be determined from the data together with Φ_{CKM} (2-fold ambiguity) , hence negligible theoretical uncertainties

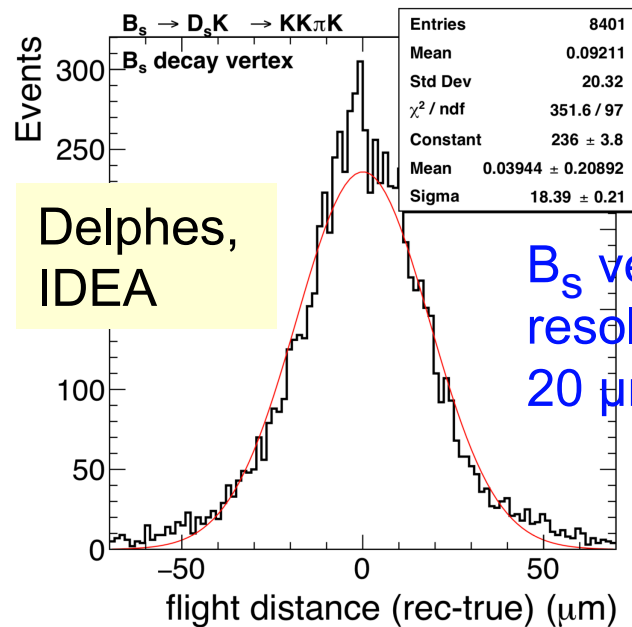
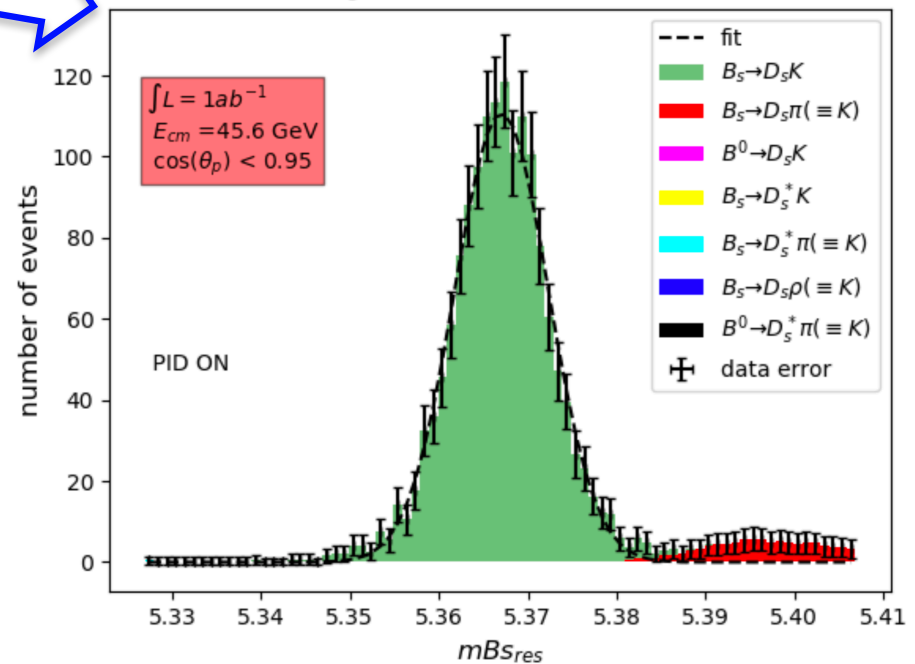
B_s to $D_s K$: signal reconstruction when $D_s \rightarrow \Phi(KK) \pi$



150 ab^{-1} : $O(1 \text{ M})$ $B_s + \bar{B}_s$ such decays.

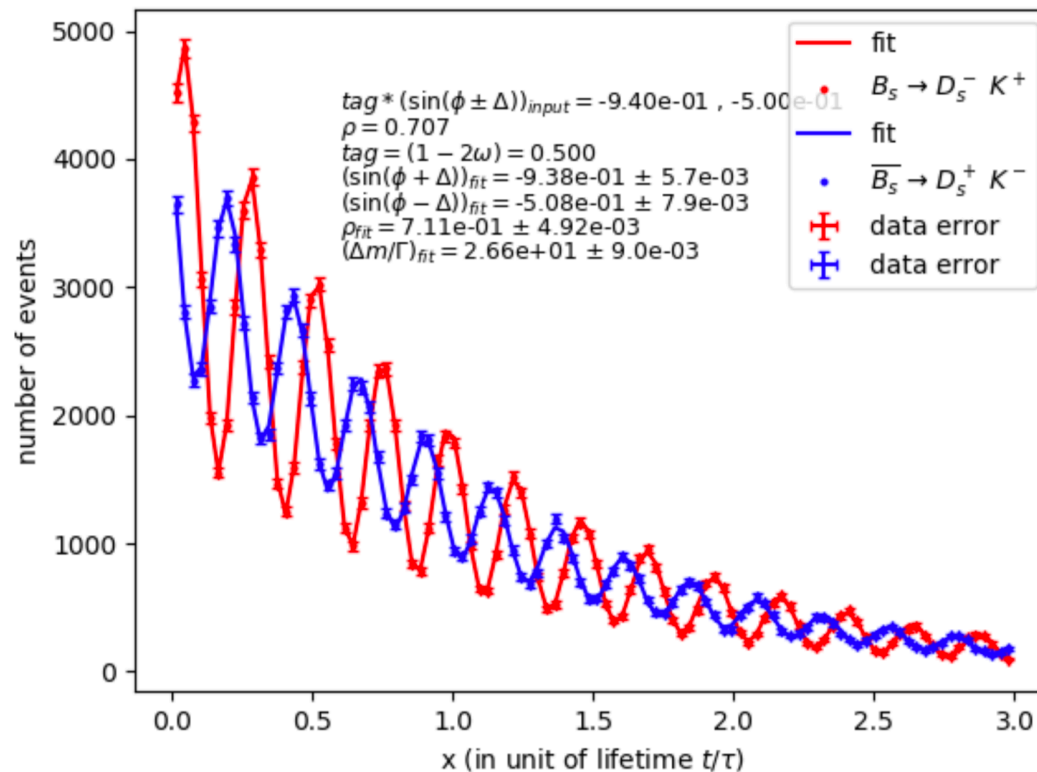
Background from $D_s \pi$ already suppressed thanks to the mass resolution.

Modest PID



Full analysis with Delphes MC samples has started (Marco Scodeggio, Ferrara)

Sensitivity on α_S (or on γ)



Statistical precision on Φ_{CKM} (on α_S)
with 150 ab^{-1} : 0.4 degrees

Resolution on Bs vertex = 20 μm
i.e. O(40) x better than the wave-length of the oscillation. Hence vertexing performance as given with IDEA good enough for this measurement.

Algebra of the args($V_{\alpha i} V_{\beta j}^* / V_{\mu k} V_{\nu l}^*$) :

$\alpha_S = \gamma - \beta_S$ + tiny angle from the (d,s) triangle

Hence this measurements is also a measurement of γ .

Current:

PDG: $\gamma = (71.1^{+4.6}_{-5.3})^\circ$

- i.e. 10x improvement w.r.t. current
- LHCb/BelleII (late 2030s) : similar sensitivity expected - PRD 102, 056023 (2020)

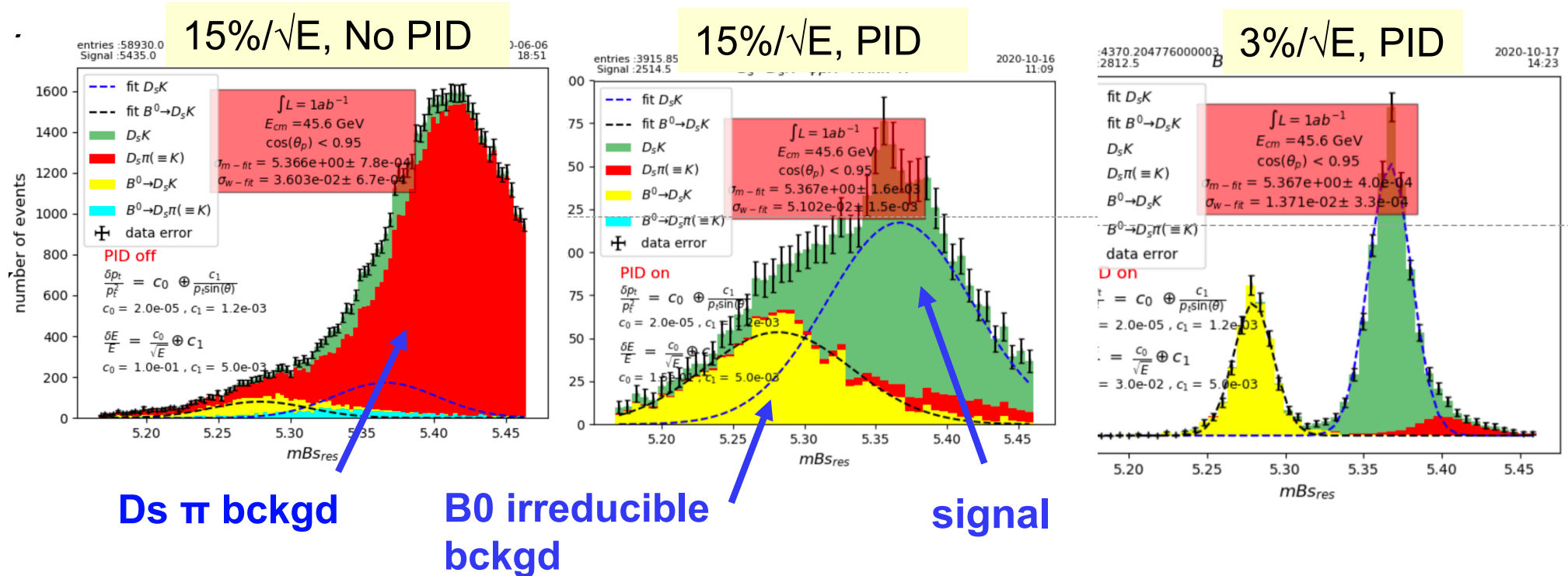
B_s to $D_s K$: inclusion of other modes

Resolution could be further improved by including other modes, that add a **photon** (e.g. $B_s \rightarrow D_s^* K$) or a π^0 (e.g. $B_s \rightarrow D_s K^*$) to the final state.

Most promising boost in statistics: $B_s \rightarrow D_s K$ with $D_s \rightarrow \Phi \rho \rightarrow K^+ K^- \pi \pi^0$.
Could increase the statistics by a **factor of 3**.

Control of backgrounds: exquisite EM resolution ($< 5\%/\sqrt{E}$) + PID (dEdx and ToF at 2m with 20 ns resolution used here)

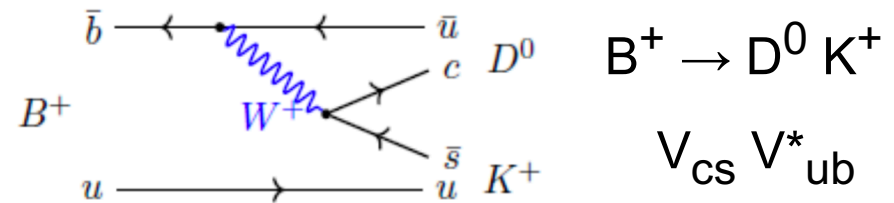
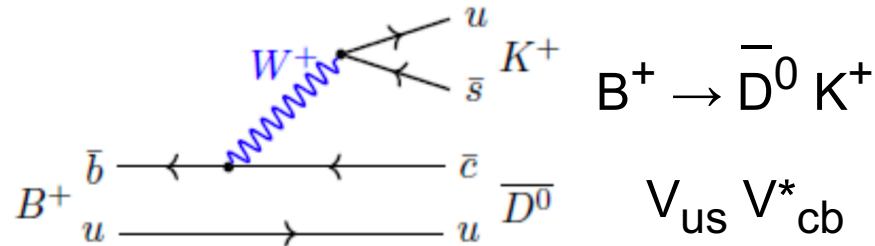
Roy Aleksan, FCC workshop, Nov 2020



Measurement of γ_s : B^+ to $D^0 K^+$

$$\gamma_s = \arg \left(-\frac{V_{cb}^* V_{cs}}{V_{ub}^* V_{us}} \right)$$

Direct CP violation in decays of B^+ to D^0 (\bar{D}^0) K^+ : well-known method to measure the γ angle of the “usual” UT. Can be applied too to measure γ_s .



With a final state f that is accessible to both D^0 and \bar{D}^0 : interference, and CPV.

$$D^0 (\bar{D}^0) \rightarrow K^+ K^- (\eta_{CP} = 1) \text{ or } K_S \pi^0 (\eta_{CP} = -1) : \Phi_{CKM} = \pi + \gamma_s$$

$\Gamma (B^+ \rightarrow f_{(D)} K^+) \neq \Gamma (B^- \rightarrow f_{(D)} K^-)$. Asymmetry \mathcal{A}_{CP} given by :

$$\frac{\pm 2\mathcal{R} \sin \Delta \sin \gamma_s}{1 + \mathcal{R}^2 \mp 2\mathcal{R} \cos \Delta \cos \gamma_s}$$

$$\mathcal{R}^2 = \frac{Br(B^+ \rightarrow D^0 K^+)}{Br(B^+ \rightarrow \bar{D}^0 K^+)}$$

\mathcal{R} already known to 5%,
 can be much improved
 with D^0 semi-leptonic
 decays

Δ = strong phase difference. PDG: $-130^\circ \pm 5^\circ$

Combination of $\mathcal{A}_{CP}^+ (K^+ K^-)$ and $\mathcal{A}_{CP}^- (K_S \pi^0)$ gives
 Δ and γ_s (8-fold ambiguity)

Expected sensitivities

$$\begin{aligned} \text{BR} (B^+ \rightarrow D^0 K^+) &\sim 3.6 \cdot 10^{-4} \\ \text{BR} (B^+ \rightarrow D^0 K^+) &\sim 3.6 \cdot 10^{-6} \\ \text{BR} (D^0 \rightarrow K^+ K^-) &\sim 4.1 \cdot 10^{-3} \\ \text{BR} (D^0 \rightarrow K_S \pi^0) &\sim 1.2 \cdot 10^{-2} \end{aligned}$$

$$\begin{array}{lll} \overline{D^0} K^+ & \overline{D^0} \rightarrow K^+ K^- & \sim 5.8 \cdot 10^5 \\ D^0 K^+ & D^0 \rightarrow K^+ K^- & \sim 5.7 \cdot 10^3 \\ \overline{D^0} K^+ & \overline{D^0} \rightarrow K_S \pi^0 & \sim 1.2 \cdot 10^6 \\ D^0 K^+ & D^0 \rightarrow K_S \pi^0 & \sim 1.2 \cdot 10^4 \end{array}$$

(indicative # of B⁺ decays)

Asymmetries are sizable. E.g. with $\Delta = -130^\circ$ and $\gamma_S = 108^\circ$:

$$\mathcal{A}_{\text{CP}}^+ (K^+ K^-) \approx -15\% \quad \text{and} \quad \mathcal{A}_{\text{CP}}^- (K_S \pi^0) \approx 14\%$$

with **expected statistical uncertainties of $\sim 0.1\%$** (absolute, accounting for approx. acceptance and efficiencies), which corresponds to $\sigma(\gamma_S)$ of 2.8°

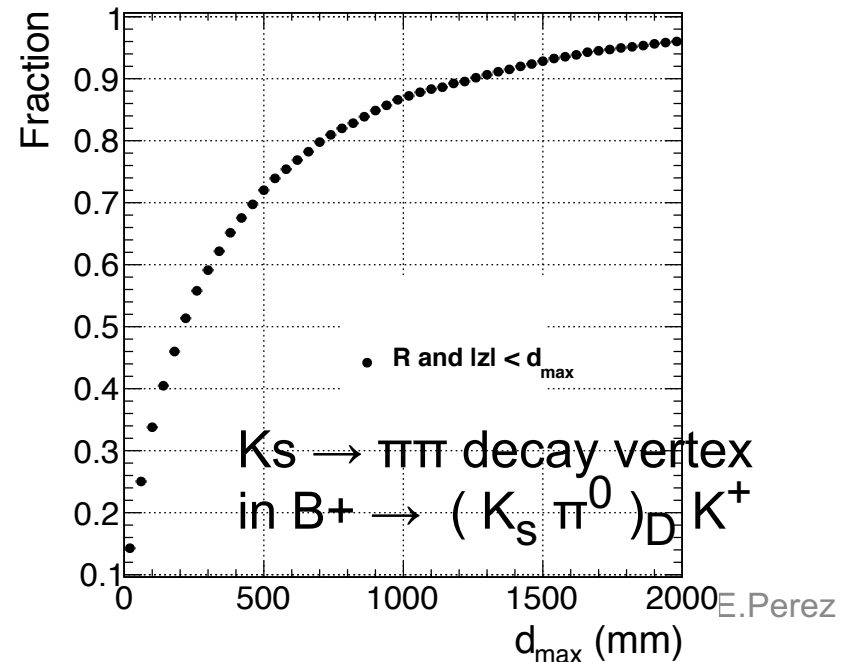
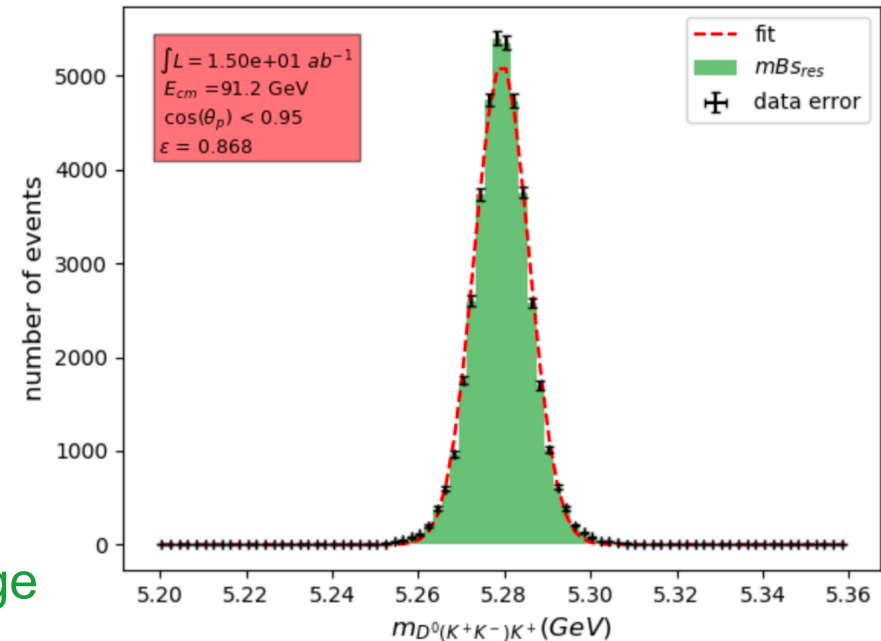
(uncertainty on γ_S depends on the value of Δ – ranges between $< 1^\circ$ to a few deg.)

Possible improvements with additional modes, e.g. $D \rightarrow K_S \eta$, $B^+ \rightarrow D K^{*+}$

Measurement of γ_S to $1^\circ - 2^\circ$ within reach.

Signal reconstruction

- $B^+ \rightarrow (K^+ K^-)_D K^+$: should be quite easy thanks to excellent mass resolution
 $\sigma \sim 6$ MeV on the B^+ mass
- $B^+ \rightarrow (K_S \pi^0)_D K^+$: much more challenging
 - Displaced pion tracks from K_S decay : Up to $O(1m)$ from the IP. Demands a **large enough tracker**
 - Worse mass resolutions :
 - π^0 : naïve σ worsens to 12 MeV even with **exquisite EM resolution** of $3\%/\sqrt{E}$
 - Finite resolution on K_S vertex will degrade this further: IDEA likely much better than Si tracker
 - Hence **more background** comes in: **Requires K/π separation in a wide p range 1 – 30 GeV**



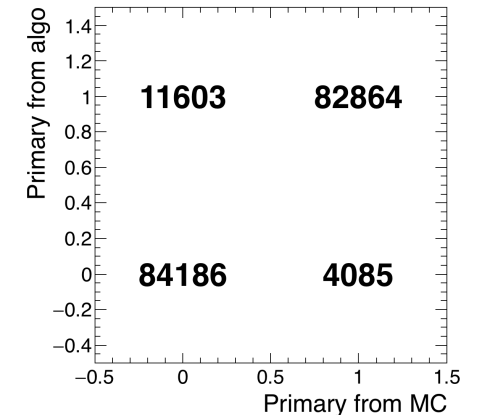
Reconstruction of Ks decays (WIP)

0th version of Ks reconstruction algorithm developed in FCCAnalyses, based on DELPHES samples. Known caveat: over-optimistic resolution on parameters of displaced tracks, will be fixed in the next version (F. Bedeschi).

- Identify the non-primary tracks

- Fit a primary vertex with all tracks
- Remove the track with the highest chi2 if this chi2 is $>$ some cut (25)
- Run the fit again, iterate

Probability
correct
assignment:
 $\sim 90\%$



- Using the non-primary tracks: find Ks candidates from all 2-track combinations

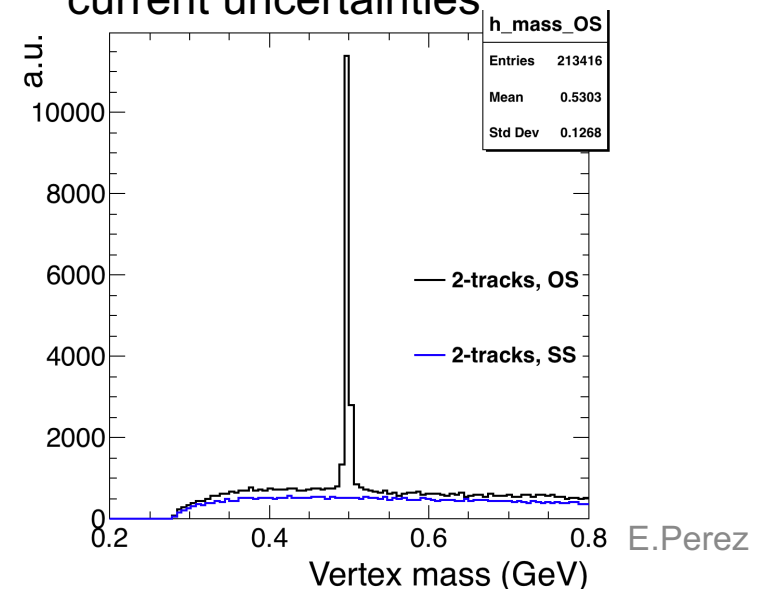
- Fit the 2 tracks to a common vertex
- Propagate the tracks to this vertex
- Build the vertex mass, using $m = \text{pion}$ mass for each leg

Loose chi2 cut, mass window, opposite-charge pairs: good efficiency and purity..

To be quantified with:

- upgraded version of DELPHES
- the FullSim tracking of the CLD detector

Mass resolution ~ 3 MeV with current uncertainties



Summary

- Precise direct measurement of the three angles of the “squashed” (b,s) unitarity triangle possible at FCC-ee :

β_s : to 0.035° or better (0.035%) via $B_s \rightarrow J/\psi \Phi$
- 25x better than the current precision
 α_s : to 0.4° or better (0.5%) via $B_s \rightarrow D_s K$
 γ_s : to 1° (1%) via $B^+ \rightarrow D^0 K^+$

- Simple relation between the phases measured in these three processes :

$$-\Phi (D_s K) + \Phi (J/\psi \Phi) + \Phi (D^0 K) = 0$$

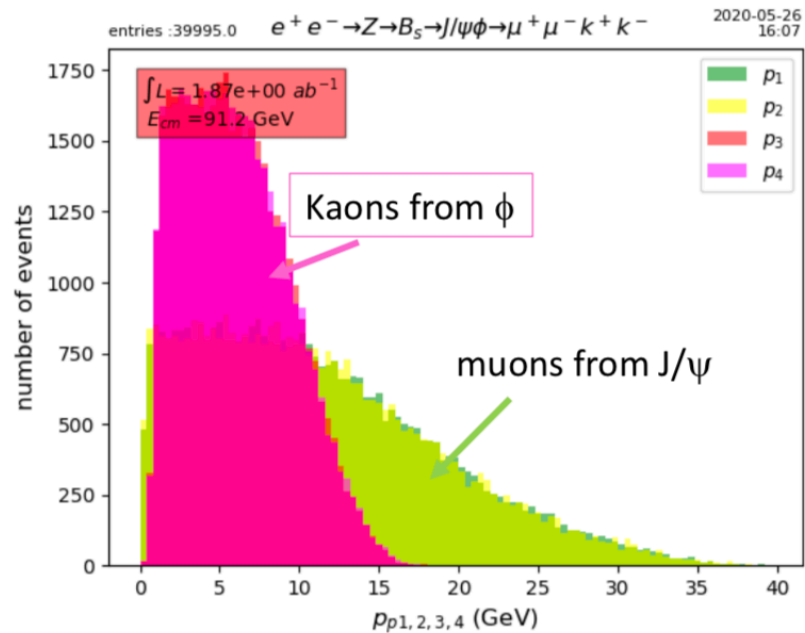
should hold in the Standard Model.

- These processes also provide good benchmarks for detector performance:
 - Excellent tracking performance (mass resolutions)
 - Excellent EM resolution (modes with neutrals)
 - K/Pi separation in a wide p range
 - Ks reconstruction (crucial for many flavour analyses)

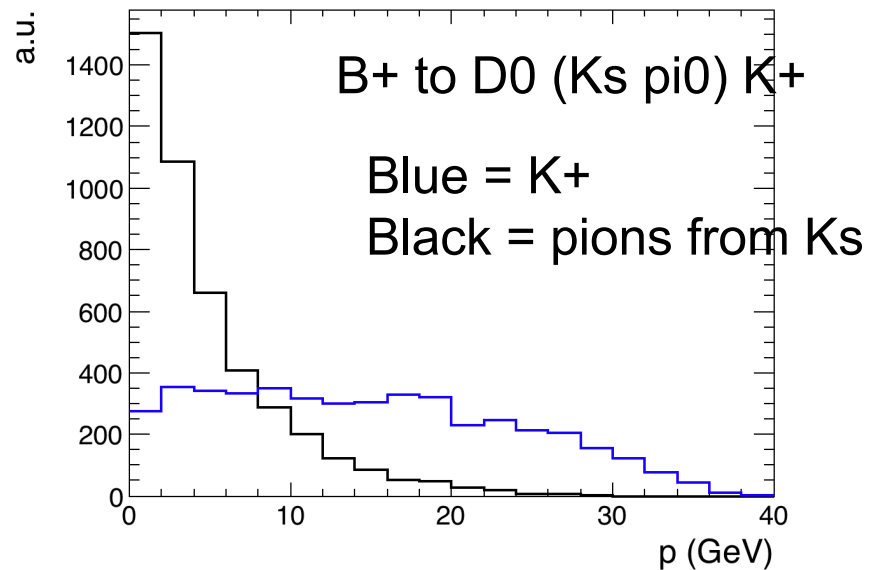
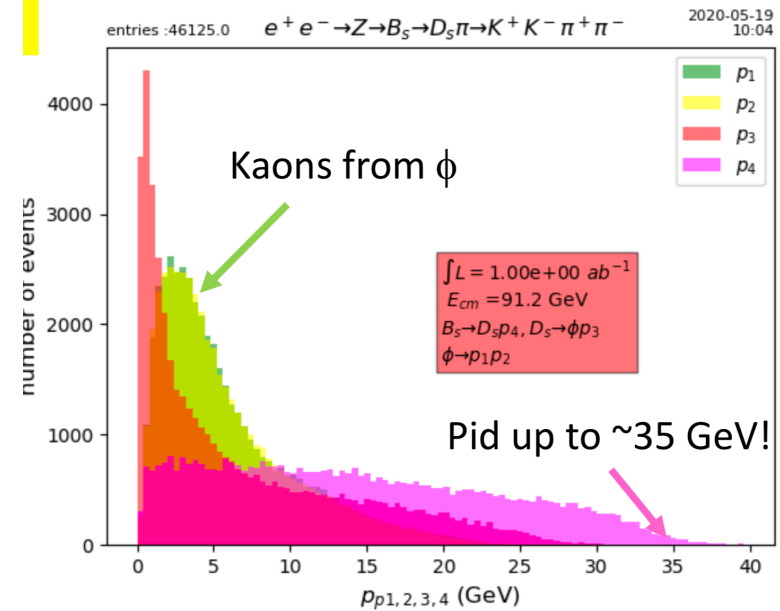
Backup

Kinematic distributions

Bs to Jpsi Phi

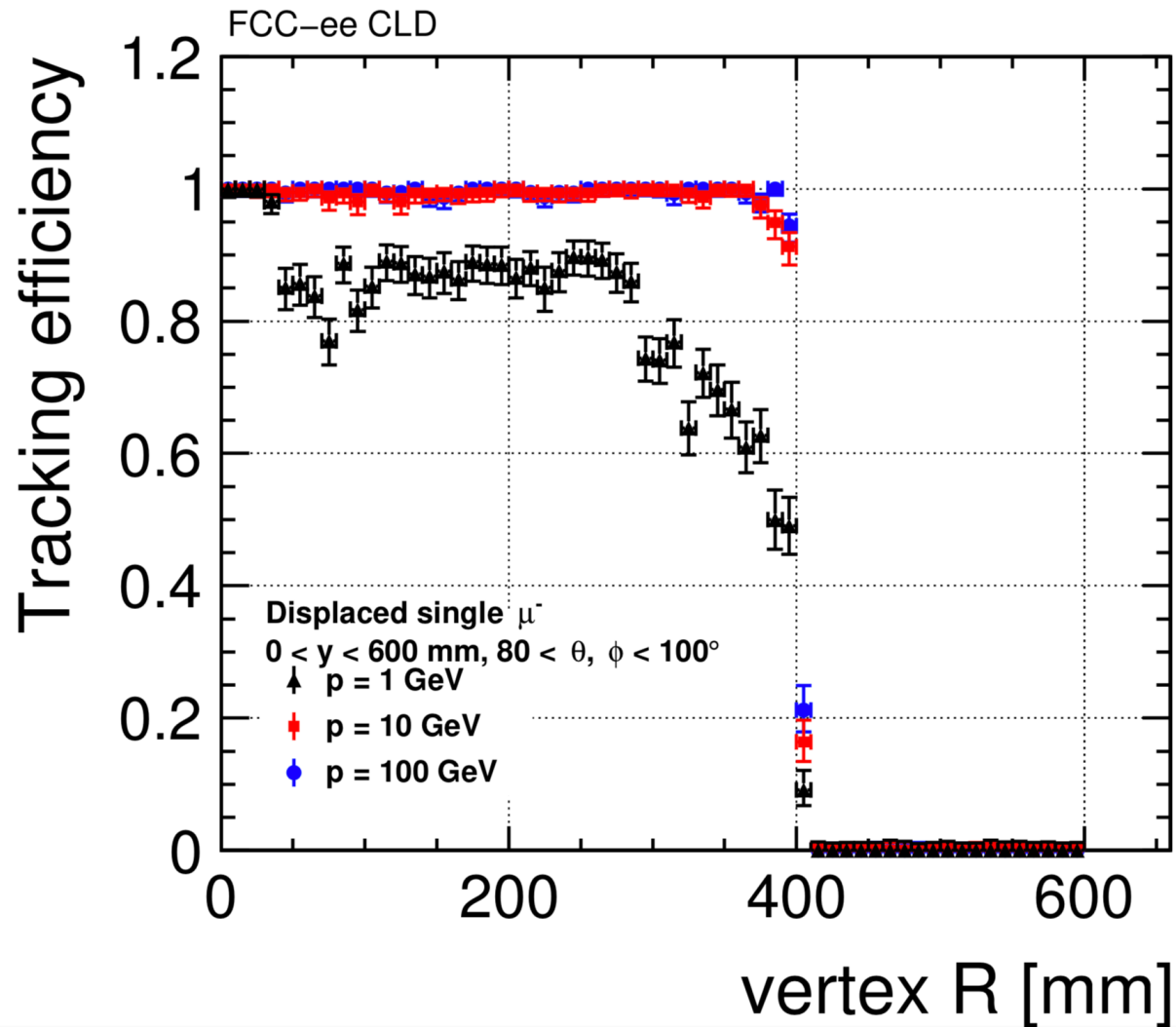


Bs to Ds K



Reconstruction of displaced tracks in CLD

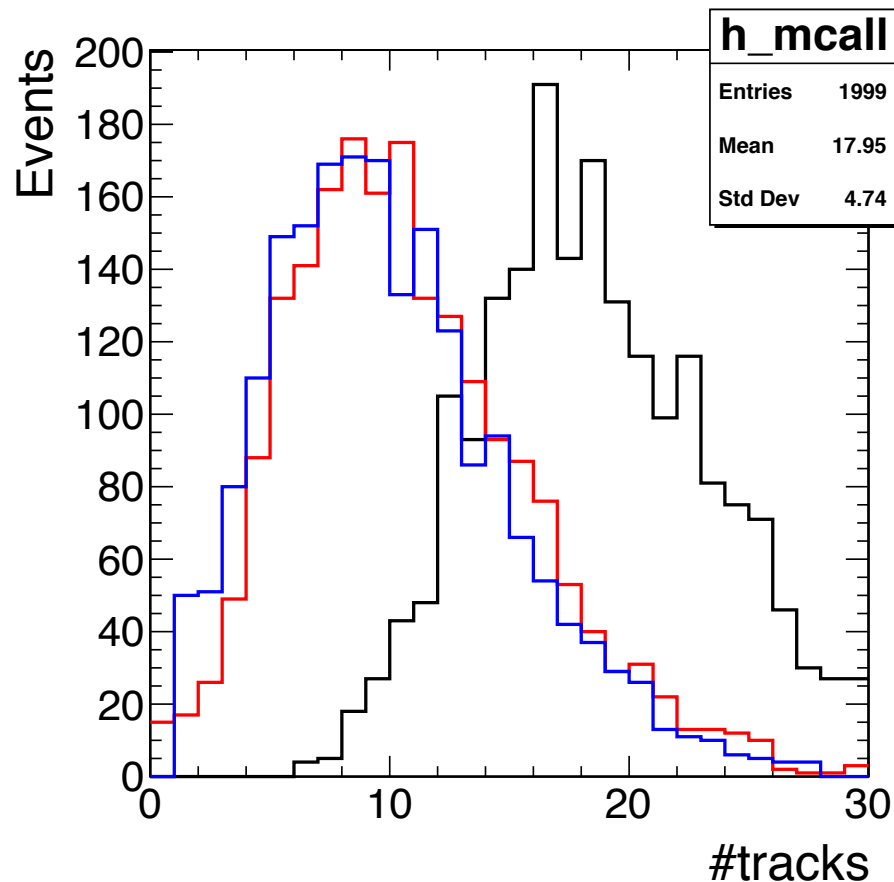
Full simulation results. Cf CLD paper, <https://arxiv.org/abs/1911.12230>



Selection of secondary tracks

In view of a Ks reconstruction: need to select secondary tracks

- Fit a primary vertex with all tracks
- Remove the track with the highest χ^2 if this χ^2 is $>$ some cut (25)
- Run the fit again, iterate



Black = all tracks

Red = tracks that are MC-matched with primary particles

Blue = the reco'd primary tracks with this procedure.

Conclusion: decent selection of non-primary tracks