

# Effective Operators in the MSSM

Guillaume Drieu La Rochelle, LAPTH

# MSSM Part I : a promising candidate

- A powerful but insufficient Standard Model

# MSSM Part I : a promising candidate

➤ A powerful but insufficient Standard Model

- Arbitrariness of the parameters

Running of the coupling constants

No Unification

# MSSM Part I : a promising candidate

## ➤ A powerful but insufficient Standard Model

- Arbitrariness of the parameters

Running of the coupling constants

- Quadratic loop divergences

Naturalness issue

No Unification

Hierarchy

# MSSM Part I : a promising candidate

## ➤ A powerful but insufficient Standard Model

- Arbitrariness of the parameters

Running of the coupling constants

- Quadratic loop divergences

Naturalness issue

No Unification

Hierarchy

## ➤ The SUSY solution

MSSM

# MSSM Part I : a promising candidate

## ➤ A powerful but insufficient Standard Model

- Arbitrariness of the parameters

Running of the coupling constants

- Quadratic loop divergences

Naturalness issue

No Unification

Hierarchy

## ➤ The SUSY solution

MSSM

- Supersymmetric

✦ Free of quadratic divergences

✦ Unification of the coupling constants

- Minimal  Most natural extension



# MSSM Part II : The disillusion

- Betrayed by the experiments
  - An **elusive Higgs**

# MSSM Part II : The disillusion

## ➤ Betrayed by the experiments

- An **elusive Higgs**

$$M_{h,H}^2 = \frac{1}{2} \left[ M_{A_0}^2 + M_Z^2 \pm \sqrt{(M_{A_0}^2 + M_Z^2)^2 - 4M_{A_0}^2 M_Z^2 c_{2\beta}^2} \right]$$

$$M_h^2 + M_H^2 = M_{A_0}^2 + M_Z^2$$

$$M_h < M_Z$$



# MSSM Part II : The disillusion

## ➤ Betrayed by the experiments

- An **elusive Higgs**

$$M_{h,H}^2 = \frac{1}{2} \left[ M_{A_0}^2 + M_Z^2 \pm \sqrt{(M_{A_0}^2 + M_Z^2)^2 - 4M_{A_0}^2 M_Z^2 c_{2\beta}^2} \right]$$

$$M_h^2 + M_H^2 = M_{A_0}^2 + M_Z^2$$

Radiative Corrections, but ...

$$M_h < M_Z$$

?!

$$M_h > 114.4 \text{ GeV}$$

# MSSM Part II : The disillusion

## ➤ Betrayed by the experiments

- An **elusive Higgs**

$$M_{h,H}^2 = \frac{1}{2} \left[ M_{A_0}^2 + M_Z^2 \pm \sqrt{(M_{A_0}^2 + M_Z^2)^2 - 4M_{A_0}^2 M_Z^2 c_{2\beta}^2} \right]$$

$$M_h^2 + M_H^2 = M_{A_0}^2 + M_Z^2$$

Radiative Corrections, but ...

$$M_h < M_Z$$

?!

- The **fine-tuning** issue

- High  $t_\beta$
- High stop mass for Radiative Corrections

$$M_h > 114.4 \text{ GeV}$$

- Collider search  $\longrightarrow$  Narrow more and more the susy scenario

# MSSM Part II : The disillusion

## ➤ Betrayed by the experiments

- An **elusive Higgs**

$$M_{h,H}^2 = \frac{1}{2} \left[ M_{A_0}^2 + M_Z^2 \pm \sqrt{(M_{A_0}^2 + M_Z^2)^2 - 4M_{A_0}^2 M_Z^2 c_{2\beta}^2} \right]$$

$$M_h^2 + M_H^2 = M_{A_0}^2 + M_Z^2$$

Radiative Corrections, but ...

$$M_h < M_Z$$

?!

- The **fine-tuning** issue

- High  $t_\beta$
- High stop mass for Radiative Corrections



$$M_h > 114.4 \text{ GeV}$$

- Collider search  $\longrightarrow$  Narrow more and more the susy scenario

## ➤ Far from complete

- Mechanism for the **susy breaking**
- The **mu problem**
- The **little hierarchy problem**

# Effective Field Theory

- Why?  Include high energy physics contribution without the explicit knowing of that high energy theory.
- How?  Global theory at low energy  $\sim$  Low energy lagrangian + Effective Operators



# Effective Field Theory

- Why?  $\longrightarrow$  Include high energy physics contribution without the explicit knowing of that high energy theory.
- How?  $\longrightarrow$  Global theory at low energy  $\sim$  Low energy lagrangian + Effective Operators

## EFT How-To

1. Choose a **scale** high/low energy.
2. Choose a **field theory for low energy** physics (renormalizable).
3. Add some **effective operators**.  
(M,v,f) scales of effective theory
4. Infer the value of  $c_k$  :
  - from measurements
  - knowing of the underlying theory.

# Effective Field Theory

- Why?  Include high energy physics contribution without the explicit knowing of that high energy theory.
- How?  Global theory at low energy  $\sim$  Low energy lagrangian + Effective Operators

## EFT How-To

1. Choose a **scale** high/low energy.
2. Choose a **field theory for low energy** physics (renormalizable).
3. Add some **effective operators**.  
(M,v,f) scales of effective theory
4. Infer the value of  $c_k$  :
  - from measurements
  - knowing of the underlying theory.

$$f^4 \frac{c_k}{M^{d_k}} \mathcal{O}_k \left( \frac{\phi}{v} \right)$$

Compatible with **symmetries** of the low-energy theory

# Supersymmetric EFT

- Rescuing the MSSM : New physics at the TeV scale
  - Done in many Beyond the MSSM models  
NMSSM, phenomenological MSSM, etc..



# Supersymmetric EFT

- Rescuing the MSSM : New physics at the TeV scale
  - Done in many Beyond the MSSM models  
NMSSM, phenomenological MSSM, etc..
- Model independant approach
  - Include *all effective operators* allowed by the symmetries

Gauge  
Susy ...

See later

# Supersymmetric EFT

- Rescuing the MSSM : New physics at the TeV scale
  - Done in many Beyond the MSSM models  
NMSSM, phenomenological MSSM, etc..
- Model independant approach
  - Include *all effective operators* allowed by the symmetries
  - Equivalent to taking a specific BMSSM model  
& **integrating out** heavy particles
  - Infer the consequences of high energy physics

Gauge  
Susy ...

See later

# Supersymmetric EFT

- Rescuing the MSSM : New physics at the TeV scale
  - Done in many Beyond the MSSM models  
NMSSM, phenomenological MSSM, etc..
- Model independant approach
  - Include *all effective operators* allowed by the symmetries
  - Equivalent to taking a specific BMSSM model  
& **integrating out** heavy particles
  - Infer the consequences of high energy physics

Gauge  
Susy ...

See later

- Which corrections to expect
- Each effective operator acts on several observables → **Correlations**

# Choosing our Effective Operators

## ➤ Constraints

$$\blacksquare \quad f^4 \frac{c_k}{M^{d_k}} \mathcal{O}_k \left( \frac{\phi}{v} \right) \quad \boxed{c_k M^{d'_k} \mathcal{O}_k(\phi)} \quad M \sim 10 \text{ Tev}$$

Recap : The MSSM has only dimension up to 4 operators.

# Choosing our Effective Operators

## ➤ Constraints

- $f^4 \frac{c_k}{M^{d_k}} \mathcal{O}_k \left( \frac{\phi}{v} \right)$        $c_k M^{d'_k} \mathcal{O}_k(\phi)$        $M \sim 10 \text{ Tev}$

Recap : The MSSM has only dimension up to 4 operators.

- Susy/Gauge Invariance      Superfield formulation
- Focus on the Higgs sector      (possible extension to other sectors)

# Choosing our Effective Operators

## ➤ Constraints

- $f^4 \frac{c_k}{M^{d_k}} \mathcal{O}_k \left( \frac{\phi}{v} \right)$   $c_k M^{d'_k} \mathcal{O}_k(\phi)$   $M \sim 10 \text{ Tev}$

Recap : The MSSM has only dimension up to 4 operators.

- Susy/Gauge Invariance Superfield formulation
- Focus on the Higgs sector (possible extension to other sectors)

## ➤ Effective Operators

See Antoniadis et al. arxiv/0910.1100

$$\mathcal{L} = \int d^2\theta d^2\bar{\theta} K + \int d^2\theta W + \int d^2\bar{\theta} \bar{W}$$

# Choosing our Effective Operators

## ➤ Constraints

- $$f^4 \frac{c_k}{M^{d_k}} \mathcal{O}_k \left( \frac{\phi}{v} \right) \quad \boxed{c_k M^{d'_k} \mathcal{O}_k(\phi)} \quad M \sim 10 \text{ Tev}$$

Recap : The MSSM has only dimension up to 4 operators.

- Susy/Gauge Invariance Superfield formulation
  - Focus on the Higgs sector (possible extension to other sectors)

## ➤ Effective Operators

See Antoniadis et al. arxiv/0910.1100

$$\mathcal{L} = \int d^2\theta d^2\bar{\theta} K + \int d^2\theta W + \int d^2\bar{\theta} \bar{W}$$

Kähler Potential

Superpotential

Standard

Standard

Effectif

Effectif



# Choosing our Effective Operators

## ➤ Constraints

- $$f^4 \frac{c_k}{M^{d_k}} \mathcal{O}_k \left( \frac{\phi}{v} \right) \quad \boxed{c_k M^{d'_k} \mathcal{O}_k(\phi)} \quad M \sim 10 \text{ Tev}$$

Recap : The MSSM has only dimension up to 4 operators.

- Susy/Gauge Invariance Superfield formulation
- Focus on the Higgs sector (possible extension to other sectors)

## ➤ Effective Operators

See Antoniadis et al. arxiv/0910.1100

$$\mathcal{L} = \int d^2\theta d^2\bar{\theta} K + \int d^2\theta W + \int d^2\bar{\theta} \bar{W}$$

Kähler Potential

Superpotential

Standard  $H_1^\dagger e^{2gV_1} H_1 + H_2^\dagger e^{2gV_2} H_2$

Standard  $\mu H_1.H_2 + \text{Yukawa}$

Effectif  $\frac{\mathbf{a}_1}{M^2} (H_1^\dagger e^{2gV_1} H_1)^2$

Effectif  $\frac{\zeta_1}{M} (H_1.H_2)^2 + \dots$

+  $\frac{\mathbf{a}_2}{M^2} (H_2^\dagger e^{2gV_2} H_2)^2$

Dimension 6  
...

Dimension 5

# Choosing our Effective Operators

## ➤ Constraints

$$\blacksquare f^4 \frac{c_k}{M^{d_k}} \mathcal{O}_k \left( \frac{\phi}{v} \right) \quad \boxed{c_k M^{d'_k} \mathcal{O}_k(\phi)} \quad M \sim 10 \text{ Tev}$$

Recap : The MSSM has only dimension up to 4 operators.

- Susy/Gauge Invariance Superfield formulation
- Focus on the Higgs sector (possible extension to other sectors)

## ➤ Effective Operators

See Antoniadis et al. arxiv/0910.1100

$$\mathcal{L} = \int d^2\theta d^2\bar{\theta} K + \int d^2\theta W + \int d^2\bar{\theta} \bar{W}$$

Kähler Potential

Superpotential

Standard	$H_1^\dagger e^{2gV_1} H_1 + H_2^\dagger e^{2gV_2} H_2$
Effectif	$\frac{\mathbf{a}_1}{M^2} (H_1^\dagger e^{2gV_1} H_1)^2$
	$+ \frac{\mathbf{a}_2}{M^2} (H_2^\dagger e^{2gV_2} H_2)^2 + \dots$

Standard	$\mu H_1 \cdot H_2 + \text{Yukawa}$
Effectif	$\frac{\zeta_1}{M} (H_1 \cdot H_2)^2 + \dots$

# Base of Effective Operators

Dimension 5

$$\frac{\zeta_1}{M} (H_1 \cdot H_2)^2$$

Dimension 6

$$\frac{\mathbf{a}_i}{M^2} (H_i^\dagger e^{2gV_i} H_i)^2$$

$$\frac{\mathbf{a}_3}{M^2} (H_1^\dagger e^{2gV_1} H_1)(H_2^\dagger e^{2gV_2} H_2)$$

$$\frac{\mathbf{a}_4}{M^2} (H_1 \cdot H_2)(H_1^\dagger \cdot H_2^\dagger)$$

$$\frac{\mathbf{a}_5}{M^2} (H_1^\dagger e^{2gV_1} H_1)(H_1 \cdot H_2) + \text{h.c.}$$

$$\frac{\mathbf{a}_6}{M^2} (H_2^\dagger e^{2gV_2} H_2)(H_1 \cdot H_2) + \text{h.c.}$$

# Base of Effective Operators

## Dimension 5

$$\frac{\zeta_1}{M} (H_1 \cdot H_2)^2$$

Strong action on masses

## Dimension 6

$$\frac{\mathbf{a}_i}{M^2} (H_i^\dagger e^{2gV_i} H_i)^2$$

Action on standard couplings

$$\frac{\mathbf{a}_3}{M^2} (H_1^\dagger e^{2gV_1} H_1)(H_2^\dagger e^{2gV_2} H_2)$$

$$\frac{\mathbf{a}_4}{M^2} (H_1 \cdot H_2)(H_1^\dagger \cdot H_2^\dagger)$$

Higgs/Neutralinos masses

$$\frac{\mathbf{a}_5}{M^2} (H_1^\dagger e^{2gV_1} H_1)(H_1 \cdot H_2) + \text{h.c.}$$

$$\frac{\mathbf{a}_6}{M^2} (H_2^\dagger e^{2gV_2} H_2)(H_1 \cdot H_2) + \text{h.c.}$$

# Base of Effective Operators

## Dimension 5

$$\frac{\zeta_1}{M} (H_1 \cdot H_2)^2$$

Strong action on  
masses

High versus Low

$t_\beta$

## Dimension 6

$$\frac{\mathbf{a}_i}{M^2} (H_i^\dagger e^{2gV_i} H_i)^2$$

Action on standard  
couplings

$$\frac{\mathbf{a}_3}{M^2} (H_1^\dagger e^{2gV_1} H_1)(H_2^\dagger e^{2gV_2} H_2)$$

$$\frac{\mathbf{a}_4}{M^2} (H_1 \cdot H_2)(H_1^\dagger \cdot H_2^\dagger)$$

Higgs/Neutralinos  
masses

$$\frac{\mathbf{a}_5}{M^2} (H_1^\dagger e^{2gV_1} H_1)(H_1 \cdot H_2) + \text{h.c.}$$

$$\frac{\mathbf{a}_6}{M^2} (H_2^\dagger e^{2gV_2} H_2)(H_1 \cdot H_2) + \text{h.c.}$$

# Susy Breaking Terms

## ➤ Constraints

- $f^4 \frac{c_k}{M^{d_k}} \mathcal{O}_k \left( \frac{\phi}{v} \right)$       $c_k M^{d'_k} \mathcal{O}_k(\phi)$       $M \sim 10 \text{ Tev}$
- ~~Susy/Gauge Invariance~~     Spurion formulation
- Focus on the Higgs sector

# Susy Breaking Terms

## ➤ Constraints

- $f^4 \frac{c_k}{M^{d_k}} \mathcal{O}_k \left( \frac{\phi}{v} \right)$   $c_k M^{d'_k} \mathcal{O}_k(\phi)$   $M \sim 10 \text{ TeV}$
- ~~Susy/Gauge Invariance~~ Spurion formulation
- Focus on the Higgs sector

## ➤ Spurions

- $\mathbf{a}_1 = a_{10} + a_{11}m_0\theta^2 + \bar{a}_{11}m_0\bar{\theta}^2 + a_{12}m_0\theta^2\bar{\theta}^2$   $m_0 \sim 1 \text{ TeV}$
- $\zeta_1 = \zeta_{10} + \zeta_{10}m_0\theta^2$



# Susy Breaking Terms


## ➤ Constraints

- $f^4 \frac{c_k}{M^{d_k}} \mathcal{O}_k \left( \frac{\phi}{v} \right)$   $c_k M^{d'_k} \mathcal{O}_k(\phi)$   $M \sim 10 \text{ TeV}$
- ~~Susy/Gauge Invariance~~ Spurion formulation
- Focus on the Higgs sector


## ➤ Spurions

- $\mathbf{a}_1 = a_{10} + a_{11}m_0\theta^2 + \bar{a}_{11}m_0\bar{\theta}^2 + a_{12}m_0\theta^2\bar{\theta}^2$   $m_0 \sim 1 \text{ TeV}$
- $\zeta_1 = \zeta_{10} + \zeta_{10}m_0\theta^2$

Example  $\int d^2\theta (1 + Bm_0\theta^2) \mu H_1.H_2 = \int d^2\theta \mu H_1.H_2 + B\mu m_0 h_1.h_2$



Spurion



Standard  
Operator

# Susy Breaking Terms

## ➤ Constraints


- $f^4 \frac{c_k}{M^{d_k}} \mathcal{O}_k \left( \frac{\phi}{v} \right)$ 
 $c_k M^{d'_k} \mathcal{O}_k(\phi)$ 
 $M \sim 10 \text{ TeV}$
- ~~Susy/Gauge Invariance~~
Spurion formulation
- Focus on the Higgs sector

## ➤ Spurions


- $\mathbf{a}_1 = a_{10} + a_{11}m_0\theta^2 + \bar{a}_{11}m_0\bar{\theta}^2 + a_{12}m_0\theta^2\bar{\theta}^2$ 
 $m_0 \sim 1 \text{ TeV}$
- $\zeta_1 = \zeta_{10} + \zeta_{10}m_0\theta^2$

Example


$$\int d^2\theta (1 + Bm_0\theta^2) \mu H_1.H_2 = \int d^2\theta \mu H_1.H_2 + B\mu m_0 h_1.h_2$$




Spurion



Standard Operator



Superpotential term



Scalar higgs mixing

# Theoretical Expectations

➤ Scalar Higgs potential

- D and F terms

$$\mathbf{H}_1 = H_1 + \sqrt{2}\theta\tilde{H}_1 + \theta^2 F_{H_1}$$

➤ Scalar Higgs – Gauge bosons interactions

- $\mathcal{L} = (DH_1)^\dagger DH_1$

➤ Standard fermions masses and couplings

➤ Higgsinos masses

➤ Higgsinos gauginos mixing

# Theoretical Expectations

## ➤ Scalar Higgs potential

- D and F terms

$$\mathbf{H}_1 = H_1 + \sqrt{2}\theta\tilde{H}_1 + \theta^2 F_{H_1}$$

## ➤ Scalar Higgs – Gauge bosons interactions

- $\mathcal{L} = (DH_1)^\dagger DH_1$

## ➤ Standard fermions masses and couplings

## ➤ Higgsinos masses

## ➤ Higgsinos gauginos mixing

## Standard MSSM

$$-\frac{\partial\bar{W}}{\partial\bar{H}_1} \frac{\partial W}{\partial H_1}$$

## Effective Operators

$$-\frac{\partial\bar{W}}{\partial\bar{H}_i} \left( \frac{\partial^2 K}{\partial\bar{H}_i \partial H_j} \right)^{-1} \frac{\partial W}{\partial H_j}$$

# Theoretical Expectations

## ➤ Scalar Higgs potential

- D and F terms

$$\mathbf{H}_1 = H_1 + \sqrt{2}\theta\tilde{H}_1 + \theta^2 F_{H_1}$$

## ➤ Scalar Higgs – Gauge bosons interactions

- $\mathcal{L} = (DH_1)^\dagger DH_1$

## ➤ Standard fermions masses and couplings

## ➤ Higgsinos masses

## ➤ Higgsinos gauginos mixing

## Standard MSSM

$$-\frac{\partial\bar{W}}{\partial\bar{H}_1}\frac{\partial W}{\partial H_1}$$

$$\frac{1}{2}\partial_\mu\phi\partial^\mu\bar{\phi} + i\partial_\mu\phi gv^\mu\bar{\phi} \\ + \text{h.c.} + \bar{\phi}(gv)^2\phi$$

...

## Effective Operators

$$-\frac{\partial\bar{W}}{\partial\bar{H}_i}\left(\frac{\partial^2 K}{\partial\bar{H}_i\partial H_j}\right)^{-1}\frac{\partial W}{\partial H_j}$$

$$\frac{1}{2}K_{i\bar{j}}\partial_\mu\phi^i\partial^\mu\bar{\phi}^{\bar{j}} + iK_{in}\partial_\mu\phi^i(gv)^{n\mu} \\ + K_{nm}(gv)^{n\mu}(gv)_\mu^m + K_n(gv)^{n2}$$

...

# Theoretical Expectations

## ➤ Scalar Higgs potential

- D and F terms

$$\mathbf{H}_1 = H_1 + \sqrt{2}\theta\tilde{H}_1 + \theta^2 F_{H_1}$$

## ➤ Scalar Higgs – Gauge bosons interactions

- $\mathcal{L} = (DH_1)^\dagger DH_1$

## ➤ Standard fermions masses and couplings

## ➤ Higgsinos masses

## ➤ Higgsinos gauginos mixing

## Standard MSSM

$$-\frac{\partial\bar{W}}{\partial\bar{H}_1}\frac{\partial W}{\partial H_1}$$

$$\frac{1}{2}\partial_\mu\phi\partial^\mu\bar{\phi} + i\partial_\mu\phi gv^\mu\bar{\phi} + \text{h.c.} + \bar{\phi}(gv)^2\phi$$

...

## Effective Operators

$$-\frac{\partial\bar{W}}{\partial\bar{H}_i}\left(\frac{\partial^2 K}{\partial\bar{H}_i\partial H_j}\right)^{-1}\frac{\partial W}{\partial H_j}$$

$$\frac{1}{2}K_{i\bar{j}}\partial_\mu\phi^i\partial^\mu\bar{\phi}^{\bar{j}} + iK_{in}\partial_\mu\phi^i(gv)^{n\mu} + K_{nm}(gv)^{n\mu}(gv)_\mu^m + K_n(gv)^{n2}$$

...

An automatic code :

- LanHEP
- Mathematica

# Expected Consequences

➤ Theory  $\longleftrightarrow$  Data

- Renormalisation/Reparametrisation    Input Set  $\{M_Z, M_W, M_{A_0}, e, t_\beta\}$



# Expected Consequences

## ➤ Theory ↔ Data

- Renormalisation/Reparametrisation    Input Set  $\{M_Z, M_W, M_{A_0}, e, t_\beta\}$
- In the effective expansion : parameter  $\frac{1}{M}$
- In the loop expansion

The two effects add up at order 1

⇒ Consider only the  $\frac{1}{M}$  expansion

$$m_h = m_{h0} + \delta_{1\text{-loop}} m_h + \delta_{\text{eff}} m_h$$

# Expected Consequences

## ➤ Theory $\longleftrightarrow$ Data

- Renormalisation/Reparametrisation    Input Set  $\{M_Z, M_W, M_{A_0}, e, t_\beta\}$
- In the effective expansion : parameter  $\frac{1}{M}$
- In the loop expansion

The two effects add up at order 1

$\implies$  Consider only the  $\frac{1}{M}$  expansion

$$m_h = m_{h0} + \delta_{1\text{-loop}} m_h + \delta_{\text{eff}} m_h$$

## ➤ Effects

- Standard coupling  $Z \bar{f} f$  changes

$$g_{Z,L} \longrightarrow g_{Z,L}(\text{SM}) + \delta g_{Z,L}$$

# Expected Consequences

## ➤ Theory $\longleftrightarrow$ Data

- Renormalisation/Reparametrisation    Input Set  $\{M_Z, M_W, M_{A_0}, e, t_\beta\}$
- In the effective expansion : parameter  $\frac{1}{M}$
- In the loop expansion

The two effects add up at order 1

$\implies$  Consider only the  $\frac{1}{M}$  expansion

$$m_h = m_{h0} + \delta_{1\text{-loop}} m_h + \delta_{\text{eff}} m_h$$

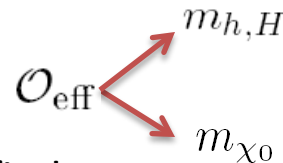
## ➤ Effects

- Standard coupling  $Z \bar{f} f$  changes

$$g_{Z,L} \longrightarrow g_{Z,L}(\text{SM}) + \delta g_{Z,L}$$

- Higgs & Neutralinos masses

Correlations



Sum rules modified

# Expected Consequences

## ➤ Theory ↔ Data

- Renormalisation/Reparametrisation    Input Set  $\{M_Z, M_W, M_{A_0}, e, t_\beta\}$
- In the effective expansion : parameter  $\frac{1}{M}$
- In the loop expansion

The two effects add up at order 1

⇒ Consider only the  $\frac{1}{M}$  expansion

$$m_h = m_{h0} + \delta_{1\text{-loop}} m_h + \delta_{\text{eff}} m_h$$

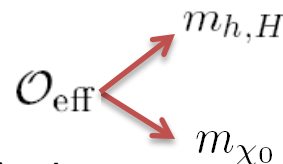
## ➤ Effects

- Standard coupling  $Z \bar{f} f$  changes

$$g_{Z,L} \longrightarrow g_{Z,L}(\text{SM}) + \delta g_{Z,L}$$

- Higgs & Neutralinos masses

Correlations



Sum rules modified

The effect of each operator propagates  
to all the Higgs sector !

# Expected Consequences

## ➤ Theory ↔ Data

- Renormalisation/Reparametrisation    Input Set  $\{M_Z, M_W, M_{A_0}, e, t_\beta\}$
- In the effective expansion : parameter  $\frac{1}{M}$
- In the loop expansion

The two effects add up at order 1

⇒ Consider only the  $\frac{1}{M}$  expansion

$$m_h = m_{h0} + \delta_{1\text{-loop}} m_h + \delta_{\text{eff}} m_h$$

## ➤ Effects

- Standard coupling  $Z \bar{f} f$  changes

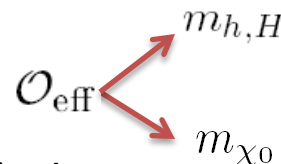
$$g_{Z,L} \longrightarrow g_{Z,L}(\text{SM}) + \delta g_{Z,L}$$



Constraints  
LEP, ...

- Higgs & Neutralinos masses

Correlations



Sum rules modified

The effect of each operator propagates  
to all the Higgs sector !

# Expected Consequences

## ➤ Theory ↔ Data

- Renormalisation/Reparametrisation    Input Set  $\{M_Z, M_W, M_{A_0}, e, t_\beta\}$
- In the effective expansion : parameter  $\frac{1}{M}$
- In the loop expansion

The two effects add up at order 1

⇒ Consider only the  $\frac{1}{M}$  expansion

$$m_h = m_{h0} + \delta_{1\text{-loop}} m_h + \delta_{\text{eff}} m_h$$

## ➤ Effects

- Standard coupling  $Z \bar{f} f$  changes

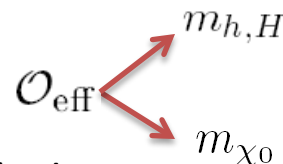
$$g_{Z,L} \longrightarrow g_{Z,L}(\text{SM}) + \delta g_{Z,L}$$



Constraints  
LEP, ...

- Higgs & Neutralinos masses

Correlations



Predictions  
Dark matter searches  
Collider experiments

Sum rules modified

The effect of each operator propagates  
to all the Higgs sector !

# Constraints Coming from standard measurements

- Coupling between Z and standard fermions

$$g_{Z\bar{f}f_{L/R}} = (1 + \kappa_{L/R}) g_{Z\bar{f}f_{L/R}}(\text{SM})$$

# Constraints Coming from standard measurements

➤ Coupling between Z and standard fermions

$$g_{Z\bar{f}f_{L/R}} = (1 + \kappa_{L/R}) g_{Z\bar{f}f_{L/R}}(\text{SM})$$

$$\kappa_{left} = -(a_{10} - t_{\beta}^2 a_{30} + t_{\beta}^4 a_{20}) \frac{2}{e^2(1 + t_{\beta}^2)^2} \frac{M_W^2}{M^2}$$

$$\kappa_{right} = (a_{10} - t_{\beta}^2 a_{30} + t_{\beta}^4 a_{20}) \frac{2(4c_W^2 - 1)}{e^2(1 + 2c_W^1)(1 + t_{\beta}^2)^2} \frac{M_W^2}{M^2}$$



# Constraints Coming from standard measurements

## ➤ Coupling between Z and standard fermions

$$g_{Z\bar{f}f_{L/R}} = (1 + \kappa_{L/R}) g_{Z\bar{f}f_{L/R}}(\text{SM})$$

$$\kappa_{left} = -(a_{10} - t_{\beta}^2 a_{30} + t_{\beta}^4 a_{20}) \frac{2}{e^2(1 + t_{\beta}^2)^2} \frac{M_W^2}{M^2}$$

$$\kappa_{right} = (a_{10} - t_{\beta}^2 a_{30} + t_{\beta}^4 a_{20}) \frac{2(4c_W^2 - 1)}{e^2(1 + 2c_W^1)(1 + t_{\beta}^2)^2} \frac{M_W^2}{M^2}$$

## ➤ Forward Backward asymmetries

Observable	Order0	$a_{10}$	$a_{20}$	$a_{30}$	Exp. Value
$A_{FB}(\mu)$	0.0181	0.000033	0.00052	-0.00013	$0.0169 \pm 0.0013$
$A_{FB}(b)$	0.1053	0.0001	0.0017	-0.0004	$0.0995 \pm 0.0017$
$A_{FB}(c)$	0.0753	0.00008	0.0012	-0.0003	$0.0713 \pm 0.0036$

# Constraints Coming from standard measurements

- Coupling between Z and standard fermions

$$g_{Z\bar{f}f_{L/R}} = (1 + \kappa_{L/R}) g_{Z\bar{f}f_{L/R}}(\text{SM})$$

$$\kappa_{left} = -(a_{10} - t_\beta^2 a_{30} + t_\beta^4 a_{20}) \frac{2}{e^2(1 + t_\beta^2)^2} \frac{M_W^2}{M^2}$$

$$\kappa_{right} = (a_{10} - t_\beta^2 a_{30} + t_\beta^4 a_{20}) \frac{2(4c_W^2 - 1)}{e^2(1 + 2c_W^1)(1 + t_\beta^2)^2} \frac{M_W^2}{M^2}$$

- Forward Backward asymmetries

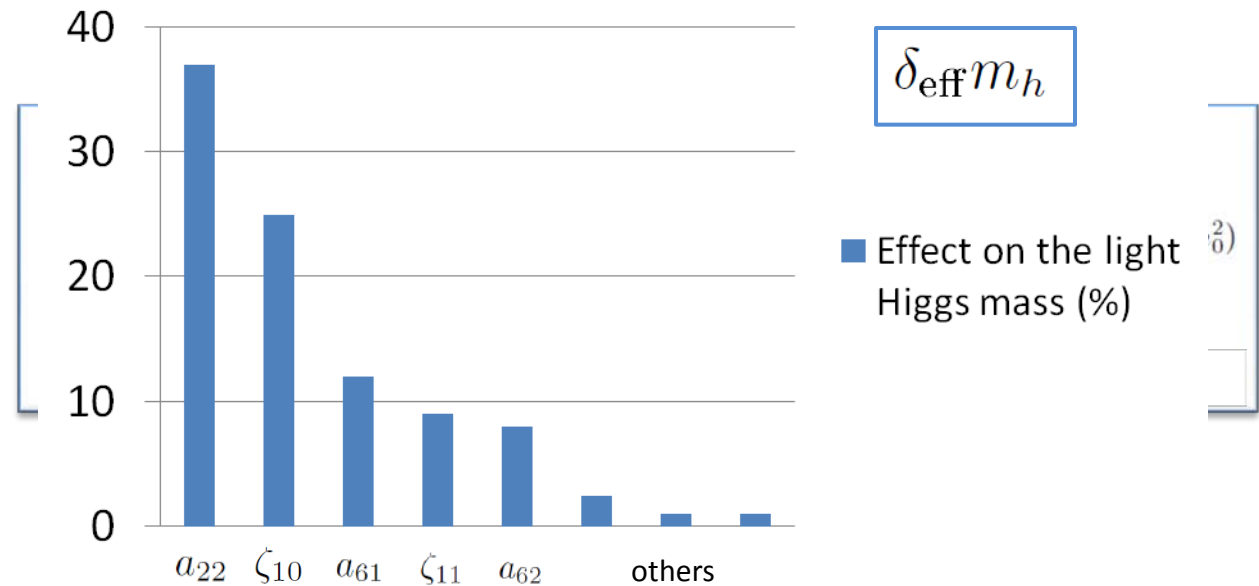
Observable	Order0	$a_{10}$	$a_{20}$	$a_{30}$	Exp. Value
$A_{FB}(\mu)$	0.0181	0.000033	0.00052	-0.00013	$0.0169 \pm 0.0013$
$A_{FB}(b)$	0.1053	0.0001	0.0017	-0.0004	$0.0995 \pm 0.0017$
$A_{FB}(c)$	0.0753	0.00008	0.0012	-0.0003	$0.0713 \pm 0.0036$

- Constraints quite loose  $a_{i0} \sim 1$

# Higgs/Neutralinos Masses

➤ Operators appearing :

	Main Operators	Effects
Neutralino mass	$a_{61}, \zeta_{10}$	A few GeV
Light Higgs mass	$\zeta_{10}, \zeta_{11}, a_{22}, a_{61}, a_{62}$	10-20 GeV



# Higgs/Neutralinos Masses

➤ Operators appearing :

	Main Operators	Effects
Neutralino mass	$a_{61}, \zeta_{10}$	A few GeV
Light Higgs mass	$\zeta_{10}, \zeta_{11}, a_{22}, a_{61}, a_{62}$	10-20 GeV

# Higgs/Neutralinos Masses

➤ Operators appearing :

	Main Operators	Effects
Neutralino mass	$a_{61}, \zeta_{10}$	A few GeV
Light Higgs mass	$\zeta_{10}, \zeta_{11}, a_{22}, a_{61}, a_{62}$	10-20 GeV

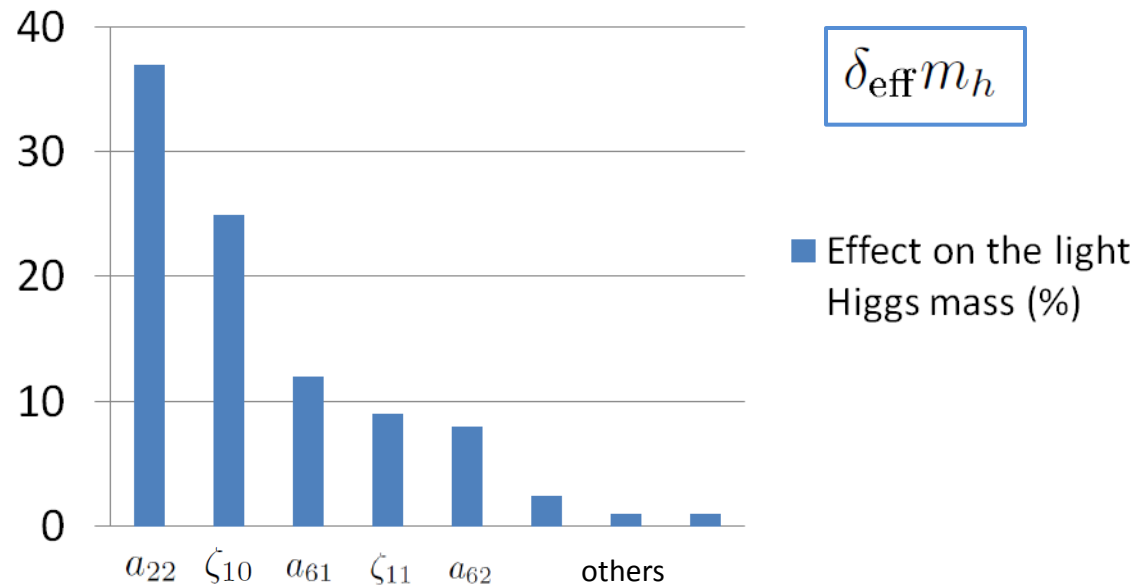
Correction to the light Higgs mass (from Antoniadis et al.)

$$\begin{aligned}
 \delta' m_h^2 = & -2v^2 \left[ \alpha_{22} m_0^2 + (\alpha_{30} + \alpha_{40}) \mu_0^2 + 2\alpha_{61} m_0 \mu_0 - \alpha_{20} m_Z^2 \right] - \frac{(2\zeta_{10} \mu_0)^2 v^4}{m_A^2 - m_Z^2} \\
 & + \frac{v^2}{\tan \beta} \left[ \frac{1}{(m_A^2 - m_Z^2)} \left( 4m_A^2 \left( (2\alpha_{21} + \alpha_{31} + \alpha_{41} + 2\alpha_{81}) m_0 \mu_0 + (2\alpha_{50} + \alpha_{60}) \mu_0^2 + \alpha_{62} m_0^2 \right) \right. \right. \\
 & \left. \left. - (2\alpha_{60} - 3\alpha_{70}) m_A^2 m_Z^2 - (2\alpha_{60} + \alpha_{70}) m_Z^4 \right) + \frac{8(m_A^2 + m_Z^2) (\mu_0 m_0 \zeta_{10} \zeta_{11}) v^2}{(m_A^2 - m_Z^2)^2} \right] \\
 & + \mathcal{O}(1/\tan^2 \beta)
 \end{aligned}$$

# Higgs/Neutralinos Masses

➤ Operators appearing :

	Main Operators	Effects
Neutralino mass	$a_{61}, \zeta_{10}$	A few GeV
Light Higgs mass	$\zeta_{10}, \zeta_{11}, a_{22}, a_{61}, a_{62}$	10-20 GeV





# Higgs/Neutralinos Correlations

➤ Sensitive Operators :

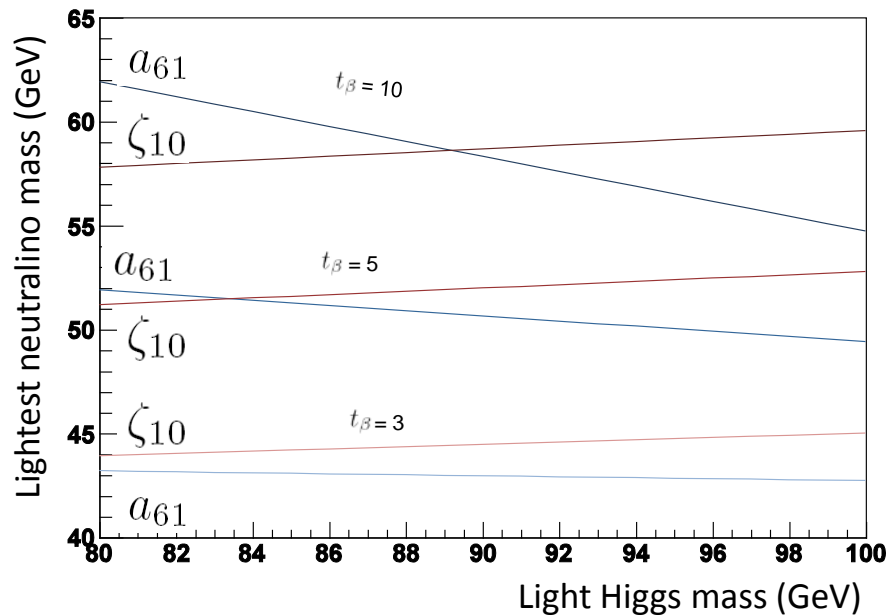
Most sensitive are :  $\zeta_{10}$   $a_{61}$

# Higgs/Neutralinos Correlations

## ➤ Sensitive Operators :

Most sensitive are :  $\zeta_{10}$   $a_{61}$

Higgs and neutralinos  
correlation for  $\zeta_{10}$   $a_{61}$





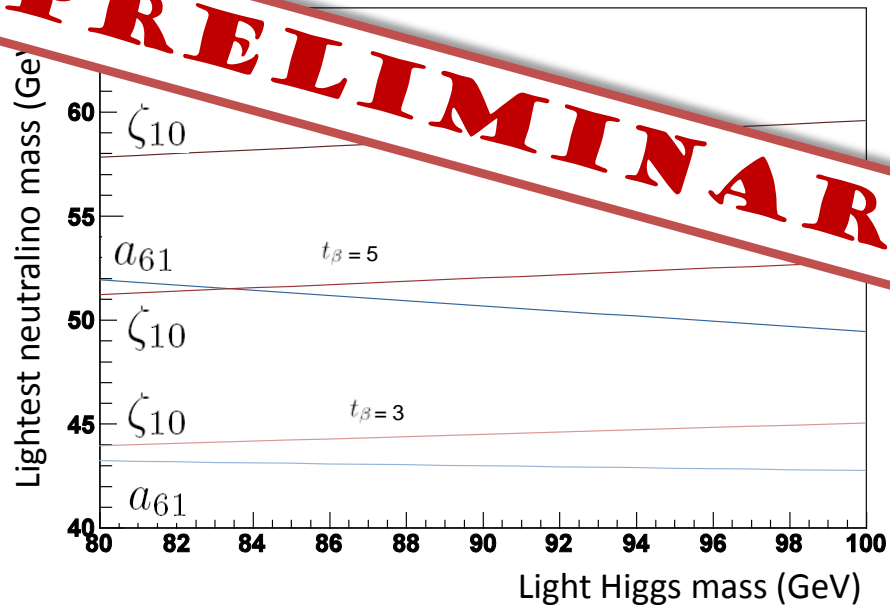
# Higgs/Neutralinos Correlations

➤ Sensitive Operators :

Most sensitive are :  $\zeta_{10}$   $a_{61}$

Higgs and neutralinos  
correlation for  $\zeta_{10}$   $a_{61}$

**PRELIMINARY**



# Analogy with radiatives corrections



# Analogy with radiatives corrections

➤ Similar results :

- Mass sum rules get modified
- Corrections depend on the physical input
- Corrections to different observables are correlated

# Analogy with radiatives corrections

- Similar results :
    - Mass sum rules get modified
    - Corrections depend on the physical input
    - Corrections to different observables are correlated
  
  - Can we include one in the other?
    - Use effective couplings to account for loops
    - Check whether a loop induces the same **correlations** than an effective operator
- ⇒ Fasten one-loop calculations

# Outlook

- Higgs phenomenology
- Effects on Dark Matter and colliders experiments
- Join effective operators with one loop calculations
- Infer what kind of physics leads to what kind of operators