

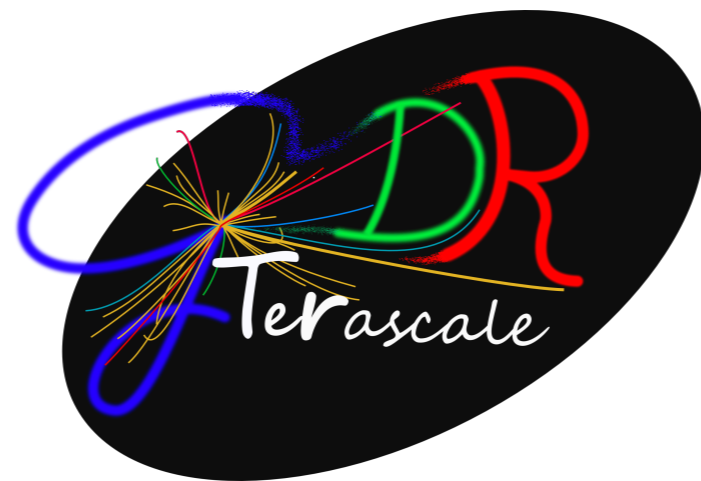


# Primordial Black Holes

*from the early universe*

**Vincent Vennin**

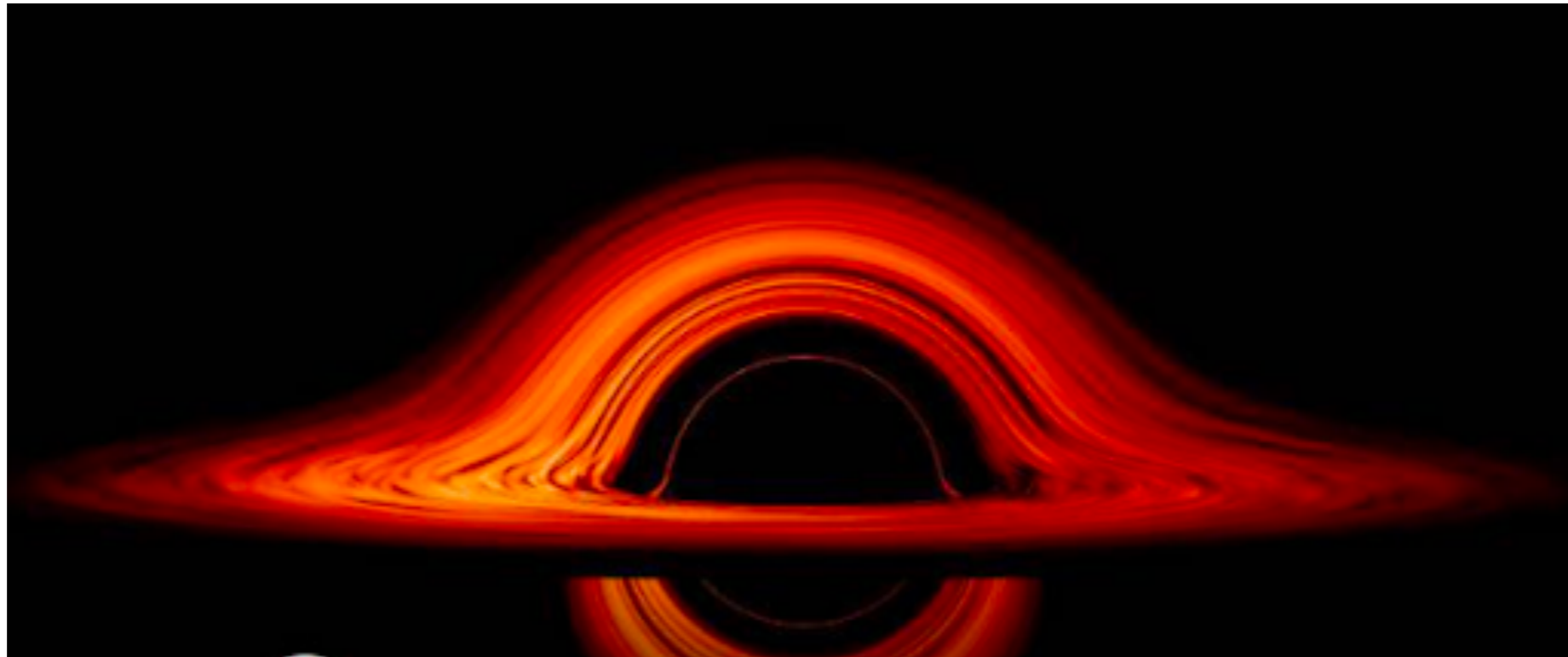
*IRN TERASCALE MEETING*



5 November 2020

# Primordial black holes

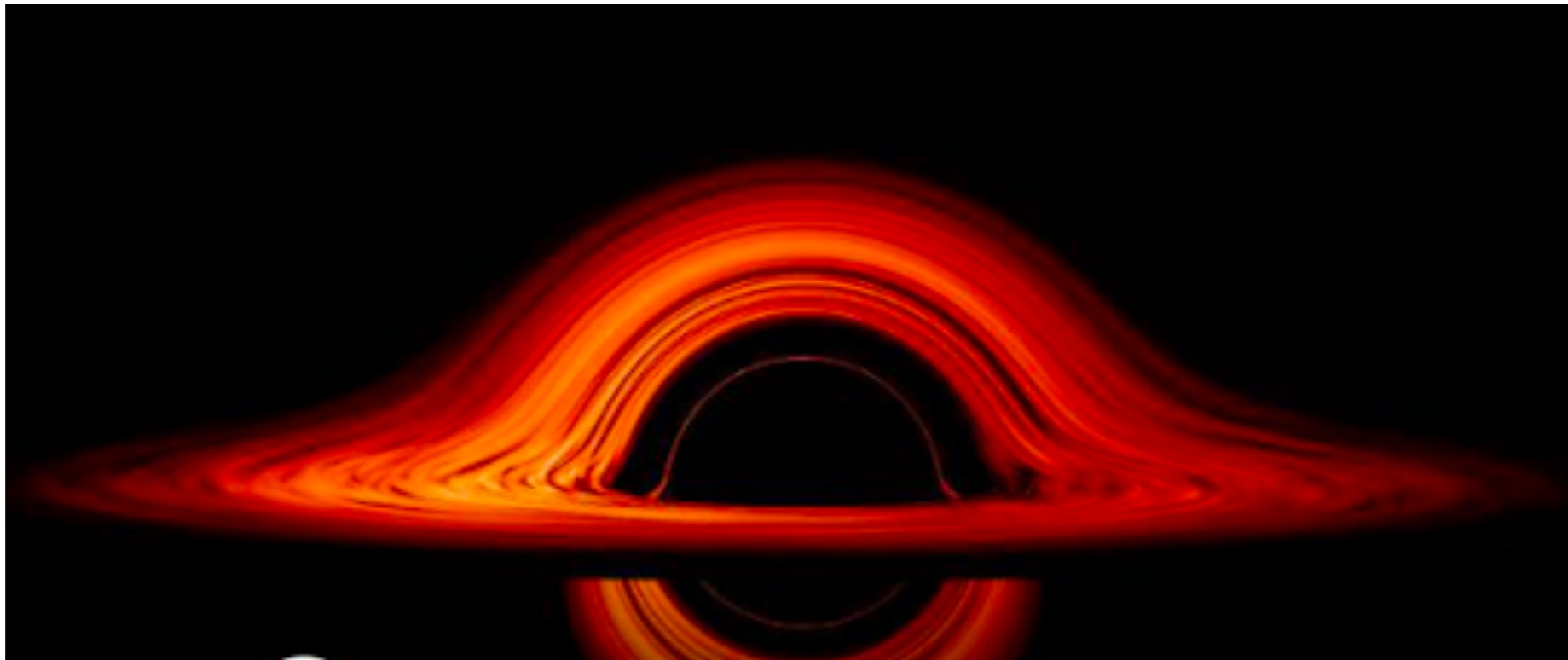
Carr, Hawking 1974



Form from the collapse of large primordial fluctuations as they re-enter the Hubble radius

# Roles in cosmology

- Could constitute part or all of dark matter [Chapline 1975](#)  
 $M = 10^{16} - 10^{17} \text{g}, 10^{20} - 10^{24} \text{g}, 10 - 10^3 M_{\odot}$
- Could provide seeds for supermassive black holes in galactic nuclei  
 $M > 10^3 M_{\odot}$  [Carr, Rees 1984](#)  
[Bean, Magueijo 2002](#)
- Could provide seeds for cosmological structures [Mészáros 1975](#)  
 $M > 10^3 M_{\odot}$  [Afshordi, McDonald, Spergel, 2003](#)
- Could provide progenitors for the LIGO/VIRGO events  
 $M = 10 - 100 M_{\odot}$

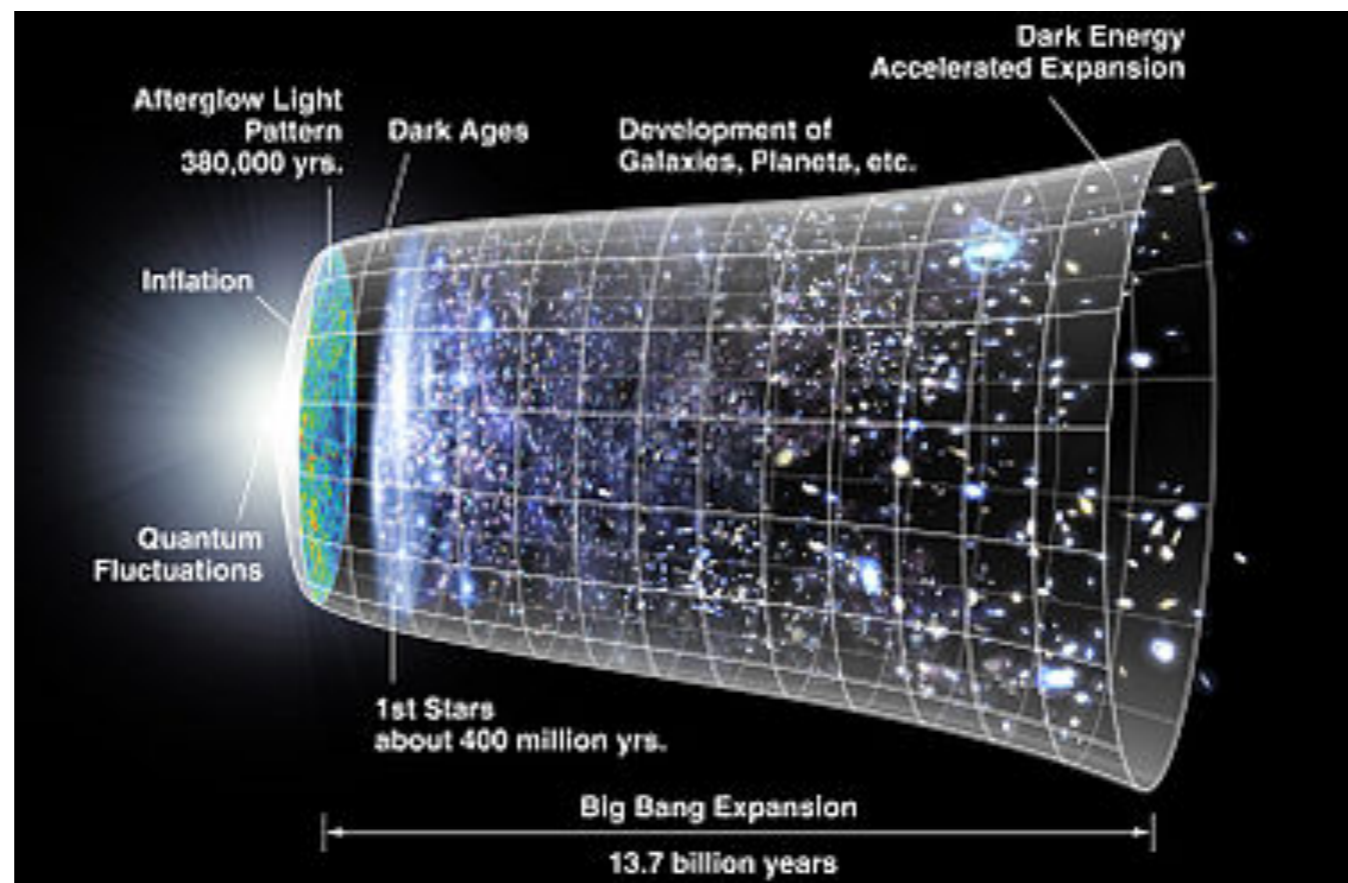


- How can those large primordial fluctuations can be seeded?
- Under what conditions do PBHs form? (How to compute their abundance?)
- How do they evolve?
- How can they be constrained?



# Cosmic Inflation

Inflation is a **high-energy** phase of **accelerated expansion** in the early Universe



Starobinsky

Guth

Linde

Albrecht & Steinhardt

Mukhanov & Chibisov

Guth & Pi

Hawking

Bardeen, Steinhardt & Turner  
(early 80's)

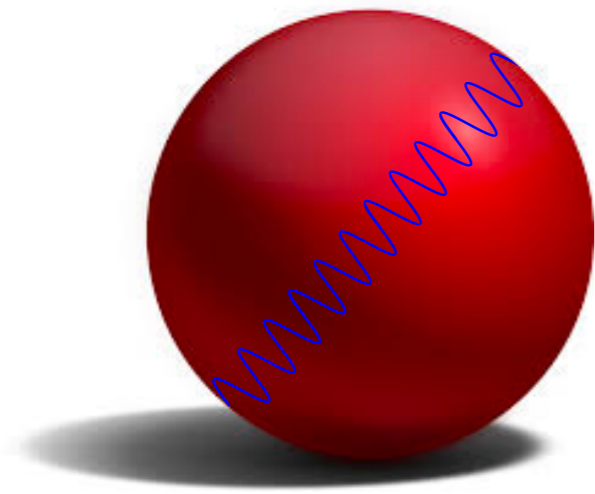
$$ds^2 = -dt^2 + a^2(t) d\vec{x}^2 \quad \text{with} \quad \ddot{a} > 0 \quad \text{and} \quad (10 \text{ MeV})^4 < \rho < (10^{16} \text{ GeV})^4$$

Turns the primordial universe into an ultra-high energy laboratory

# Cosmic Inflation

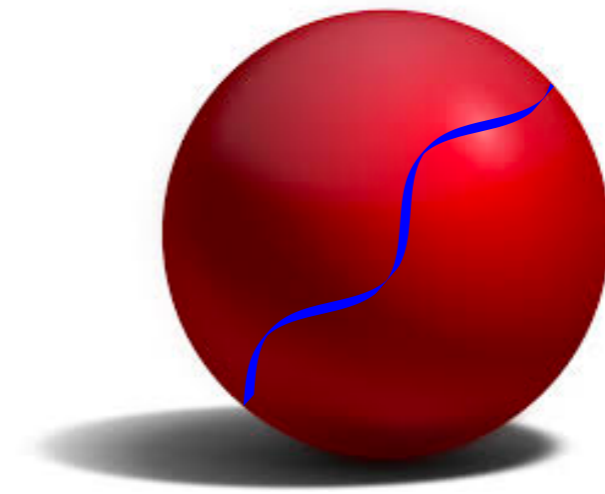
Hubble parameter  $H = \dot{a}/a$

$H^{-1}$  : characteristic time scale, or length scale ( $c = 1$ ), of the expansion



$$\lambda \ll H^{-1}$$

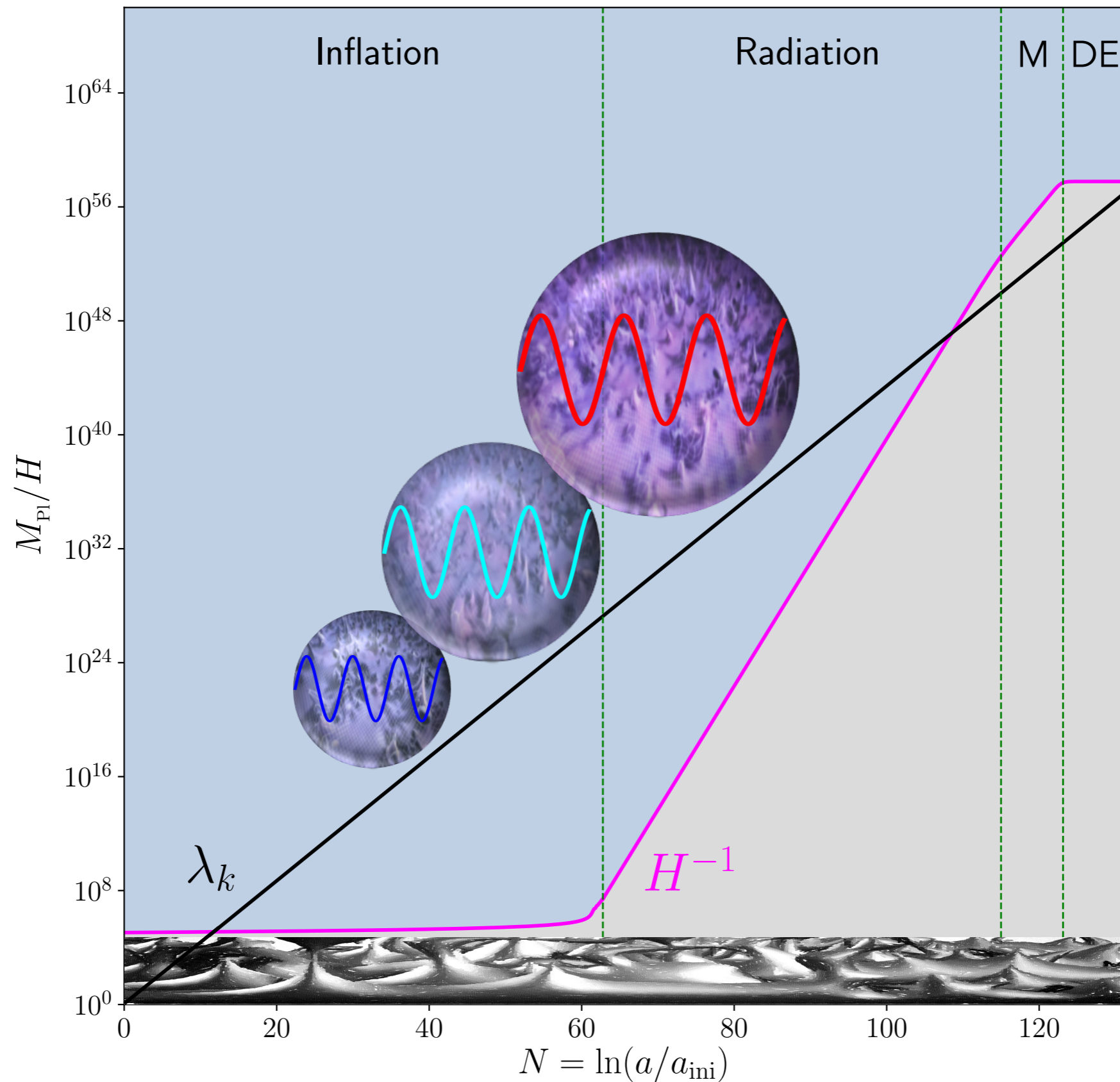
Insensitive to space-time curvature



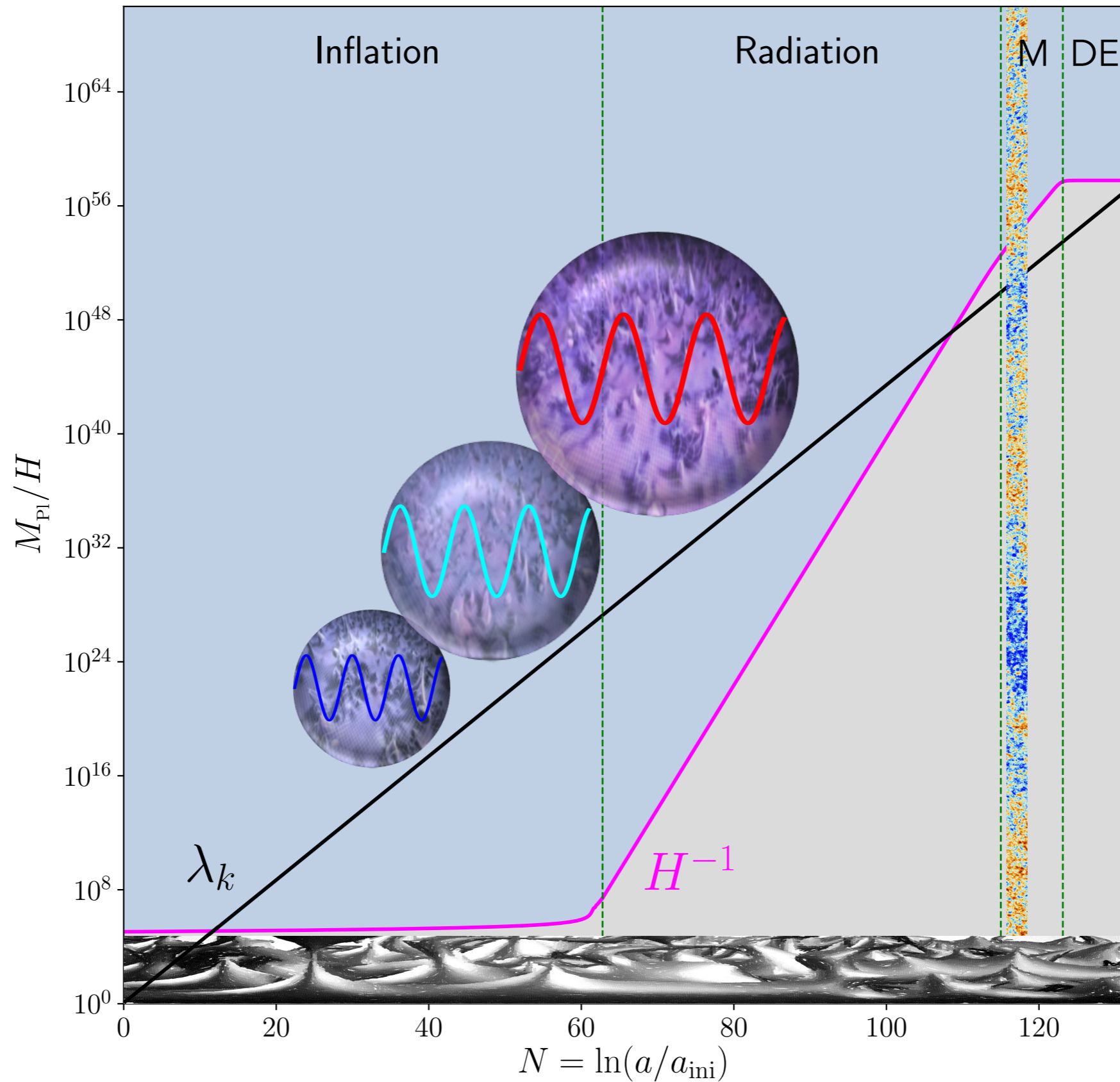
$$\lambda \gtrsim H^{-1}$$

Feels space-time curvature

# Cosmic Inflation

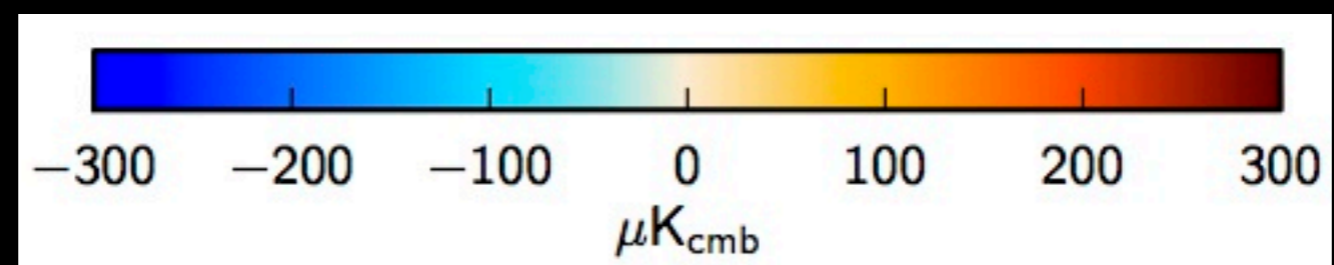
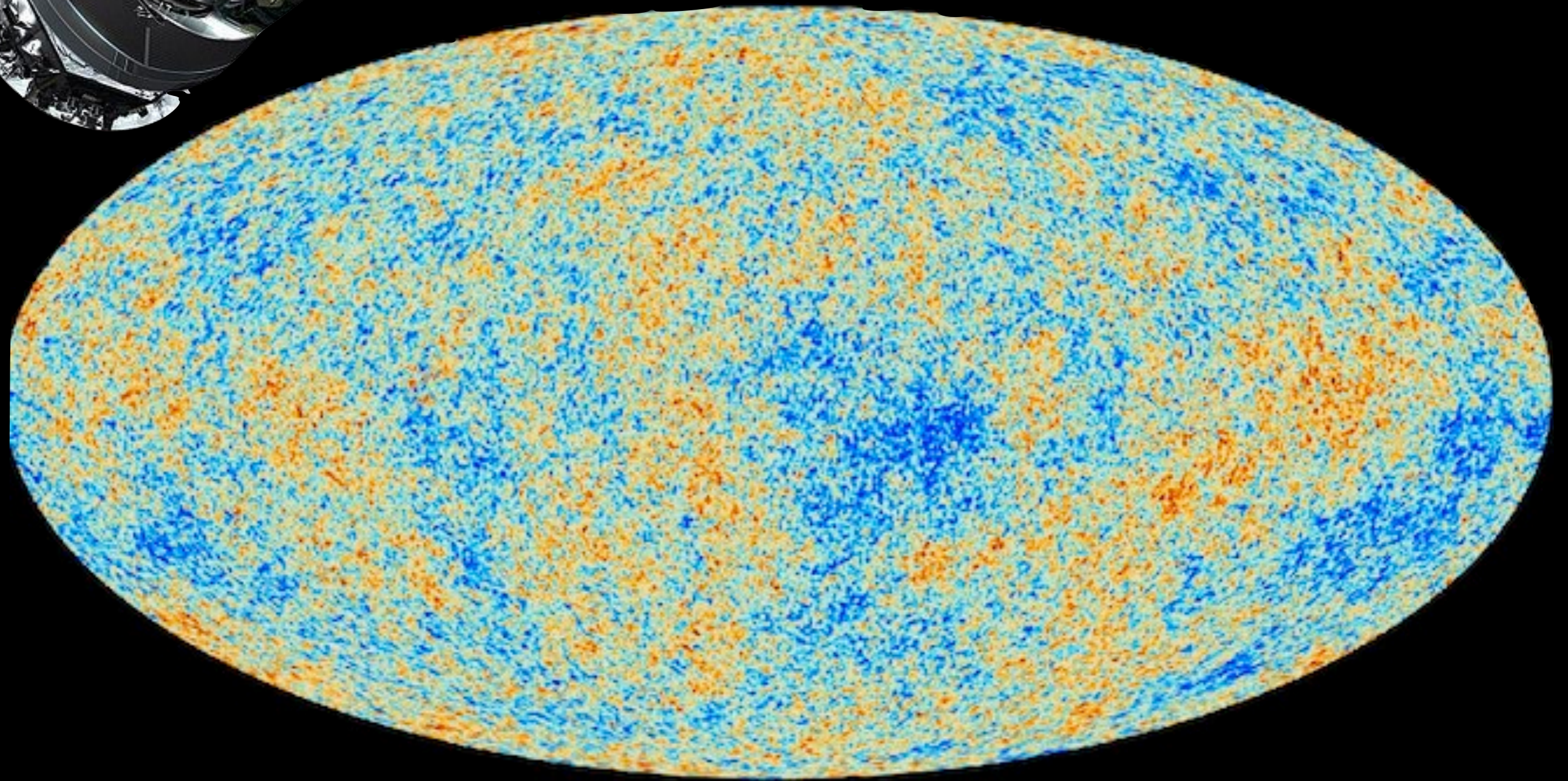
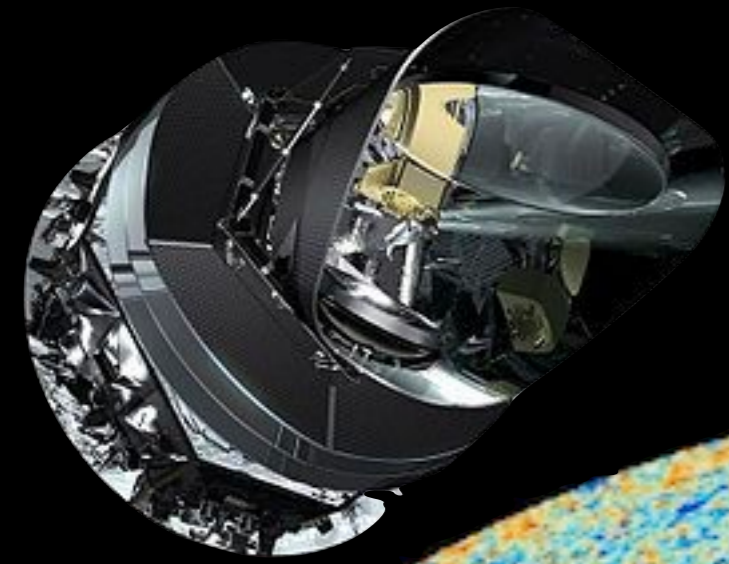


# Cosmic Inflation



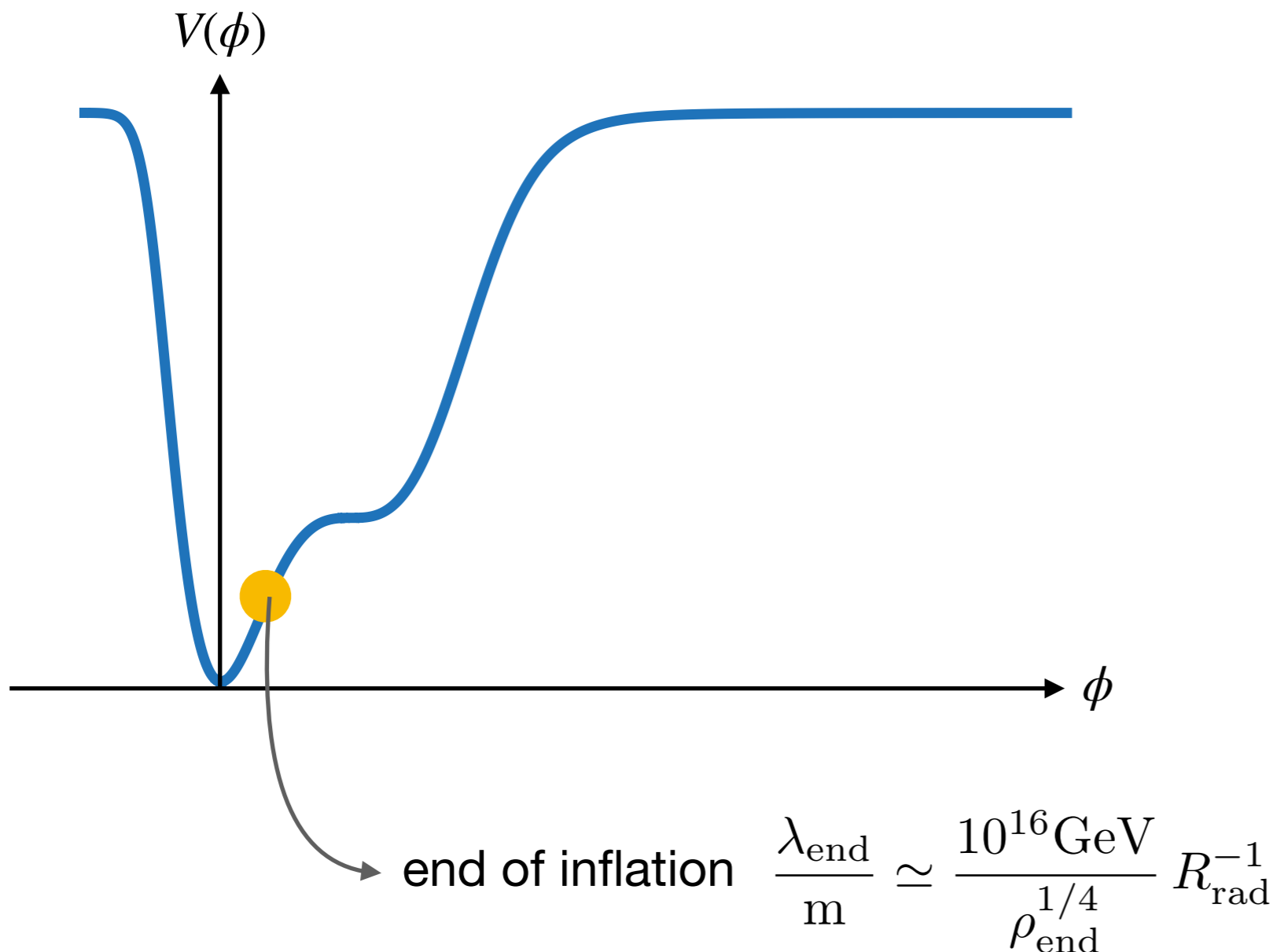


# Planck satellite

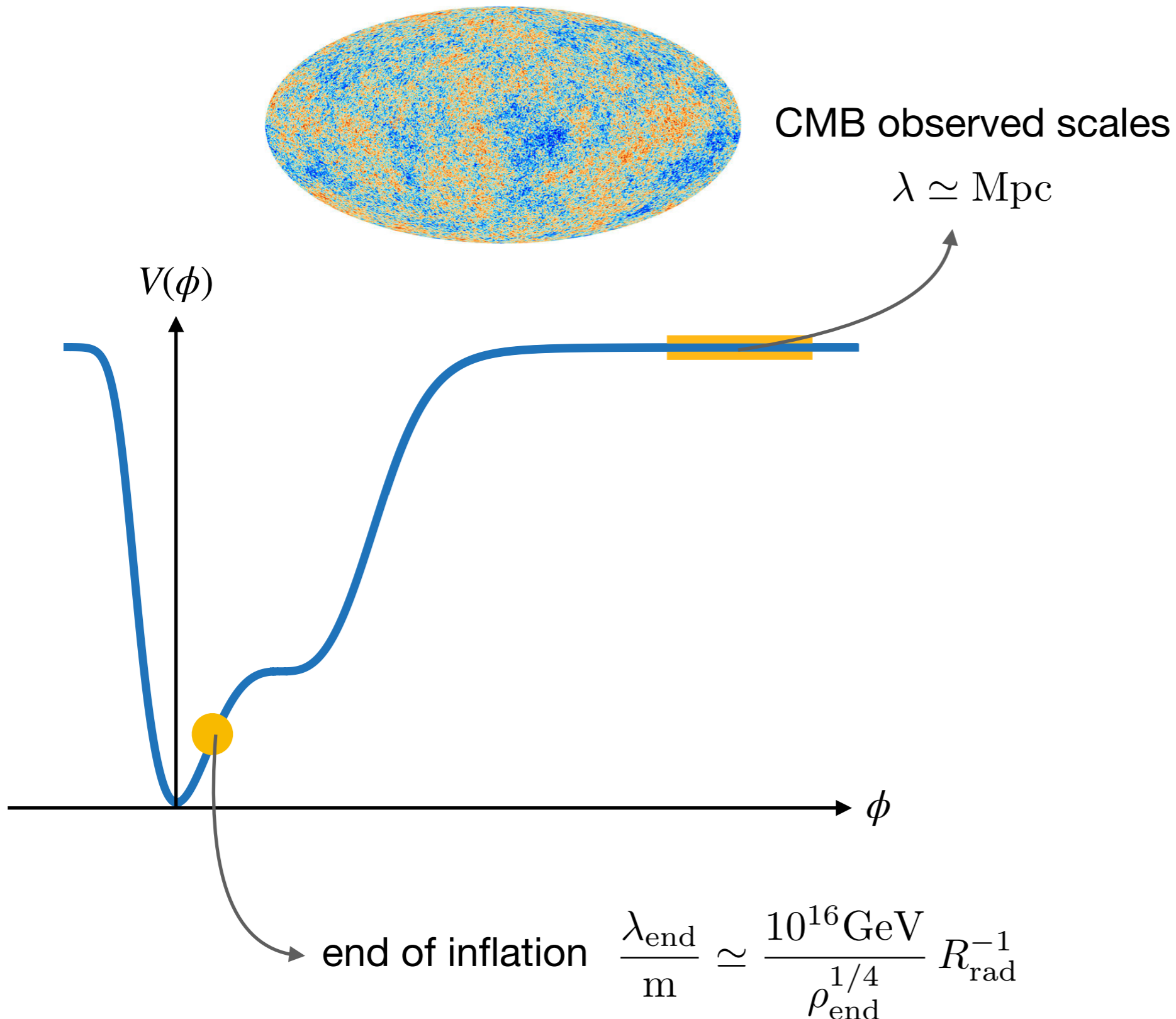




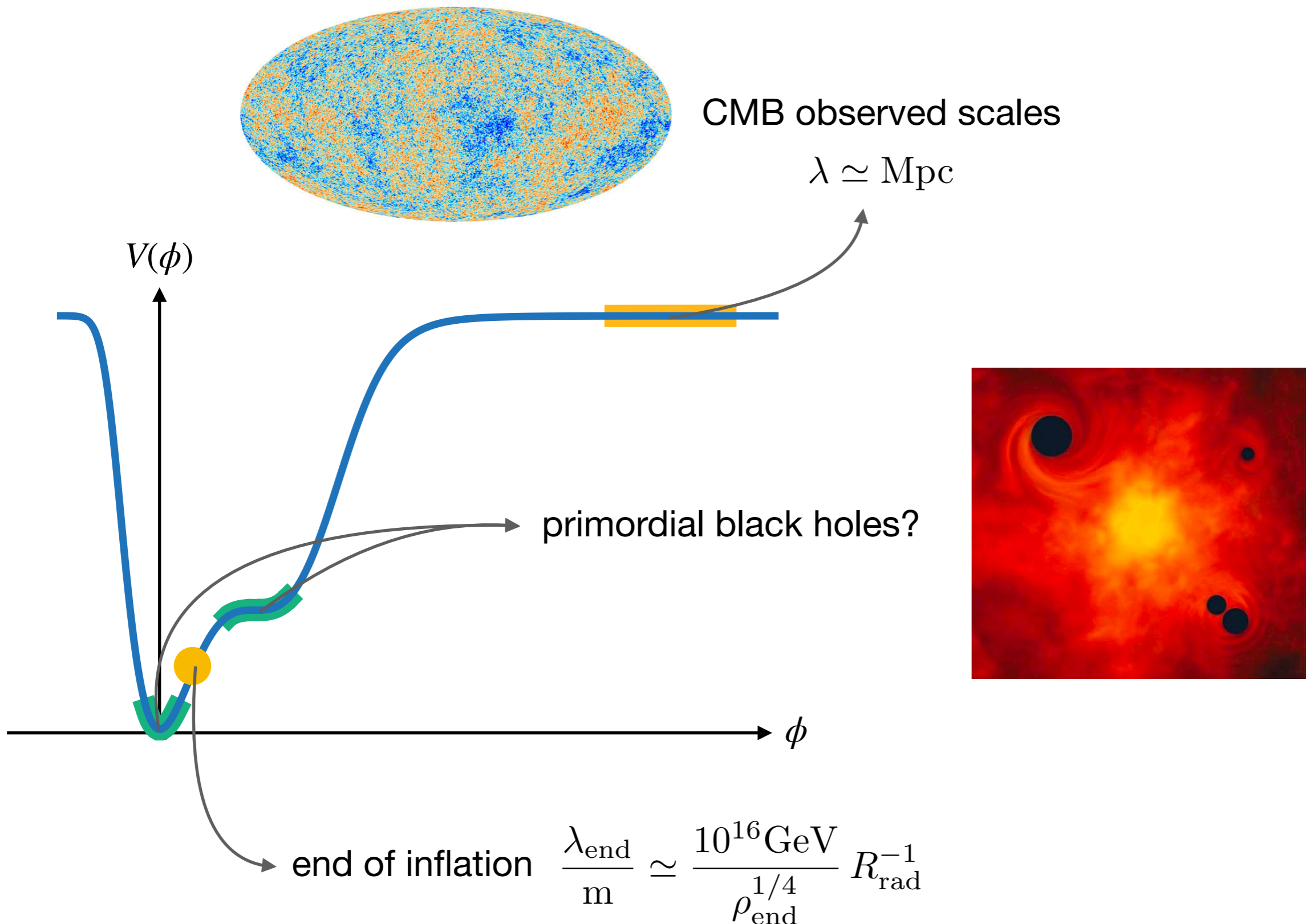
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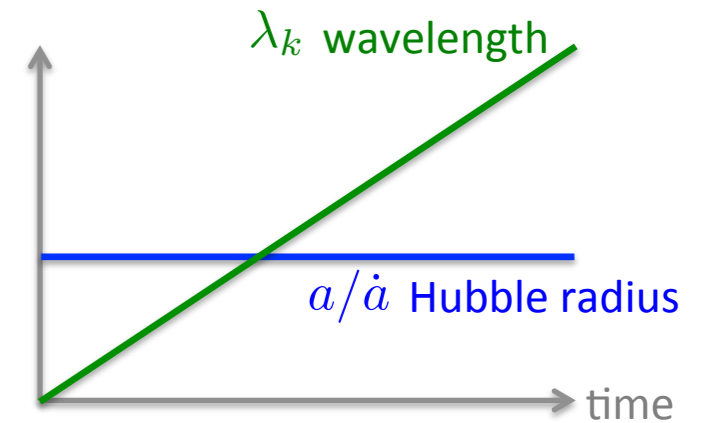
# Case study 1: inflationary enhancement

Dynamics of a scalar field in FLRW cosmology:

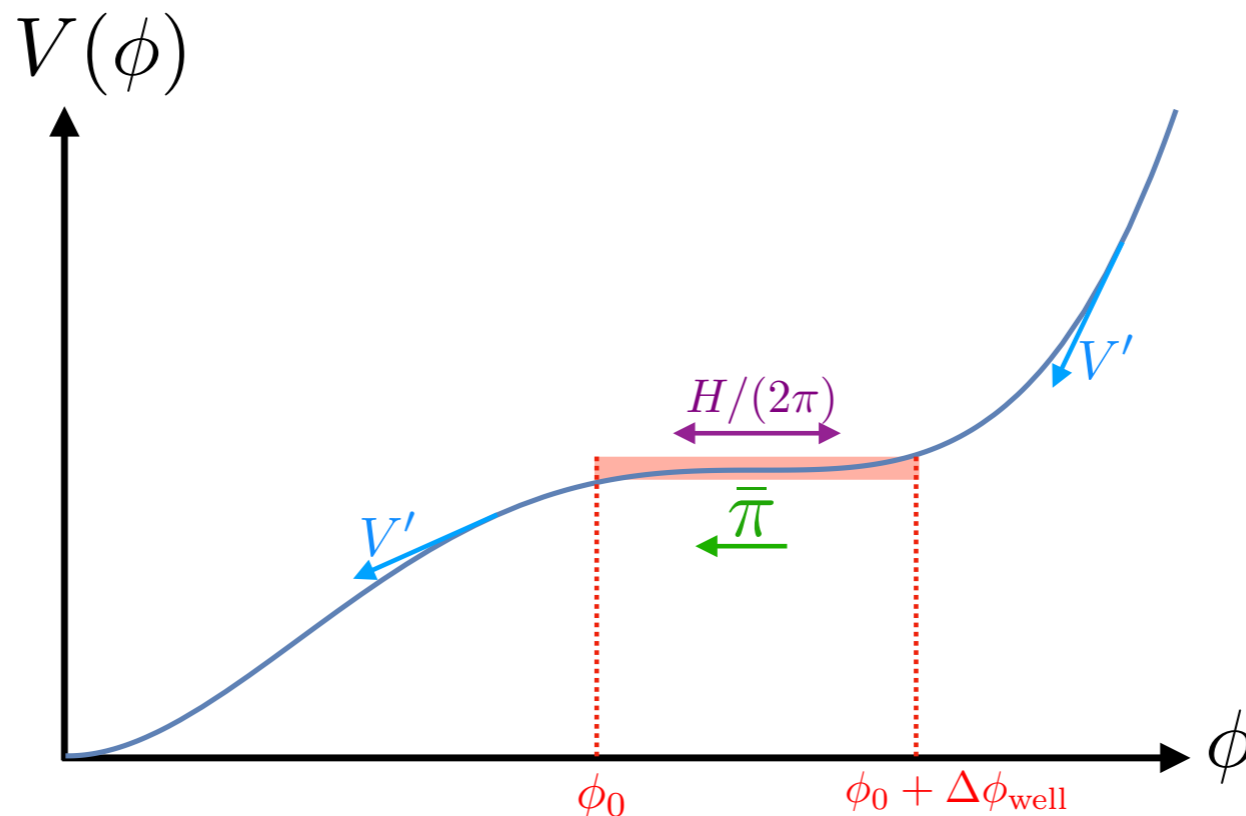
$$\dot{\phi} = \pi + \xi_{\phi}$$

$$\dot{\pi} = -3H\pi - \frac{dV}{d\phi} + \xi_{\pi}$$

friction term  $H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{V(\phi) + \frac{\pi^2}{2}}{3M_{\text{Pl}}^2}$



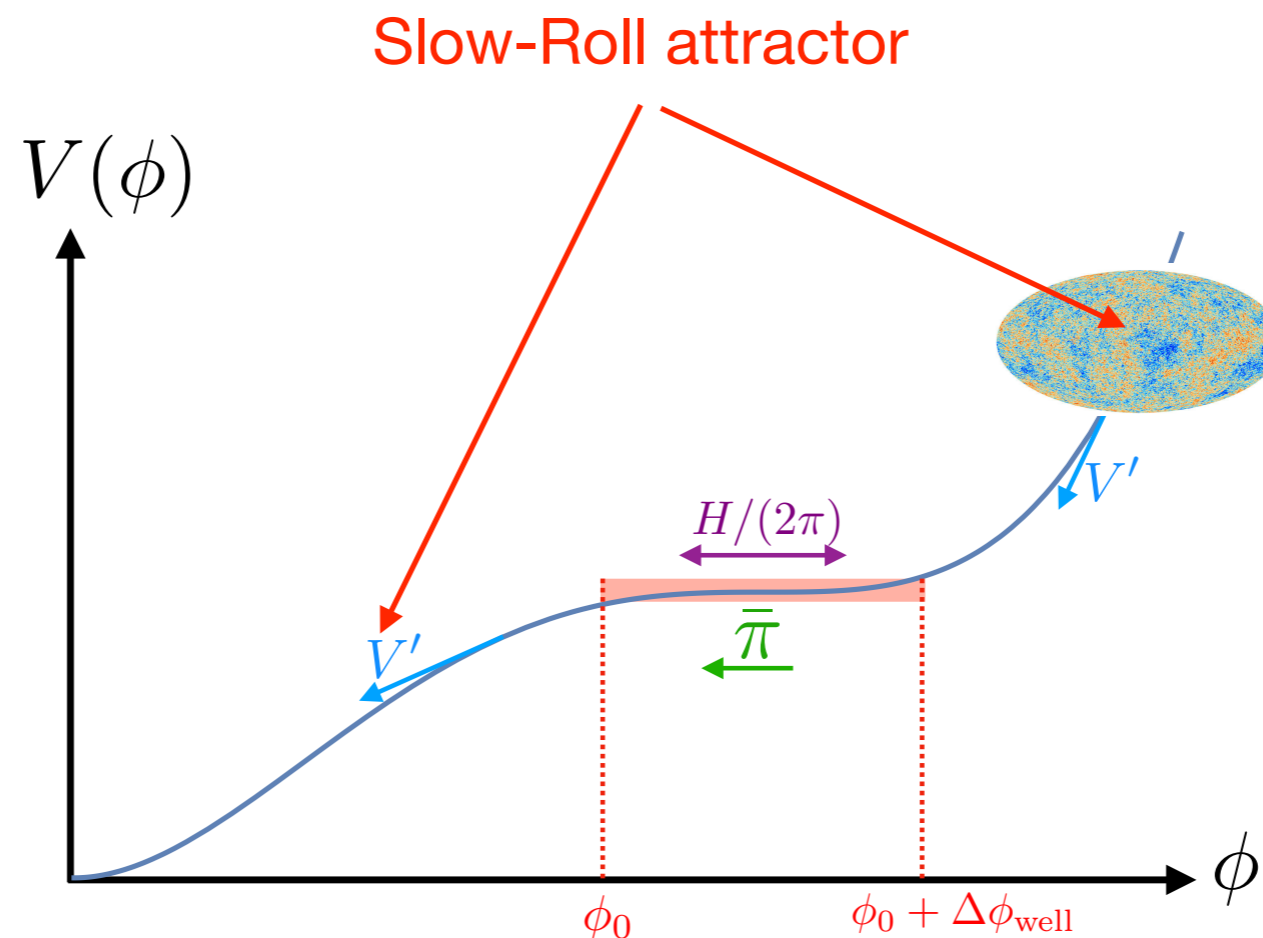
Quantum fluctuations as they cross out the Hubble radius and source the large-scale dynamics



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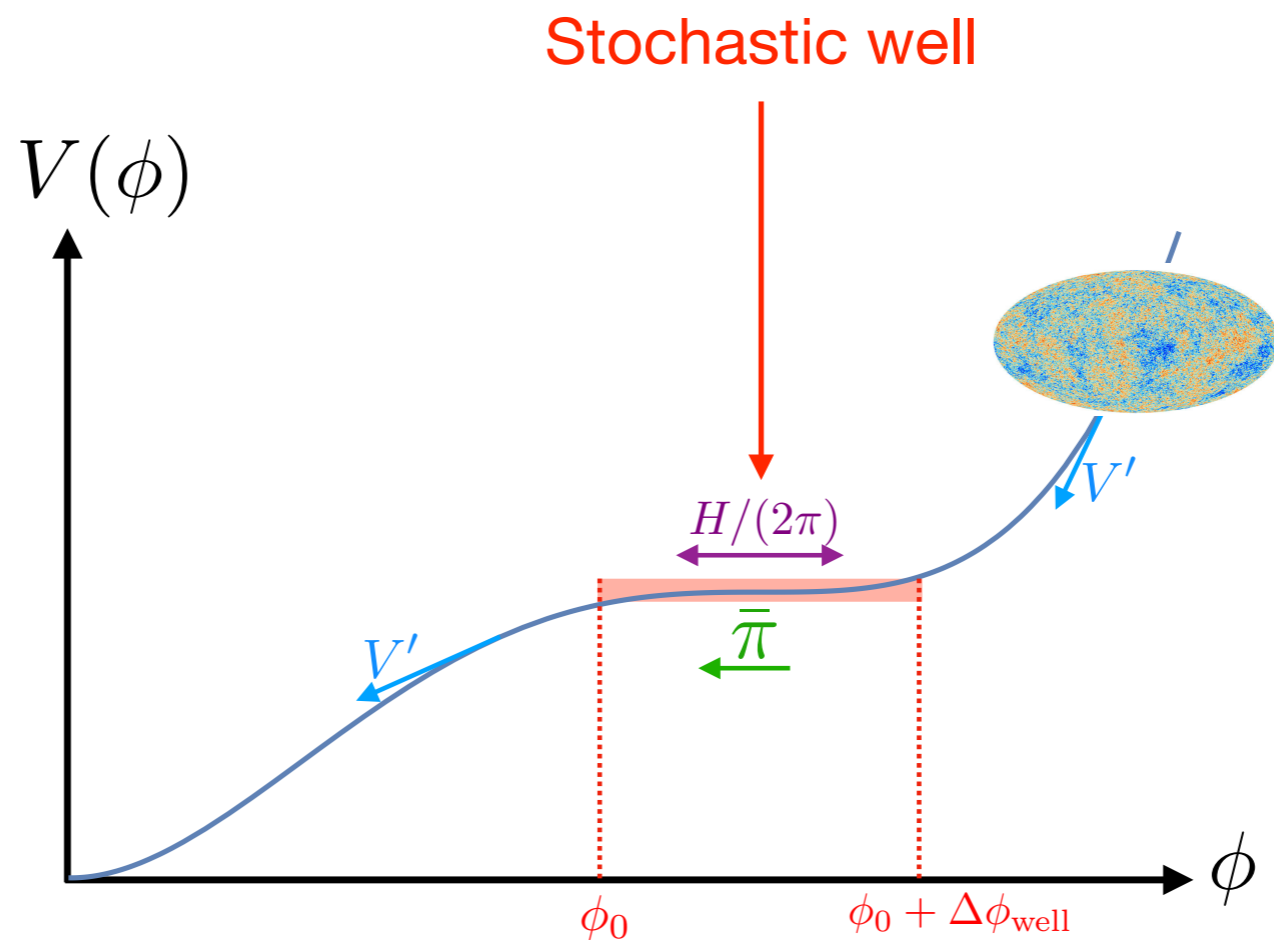


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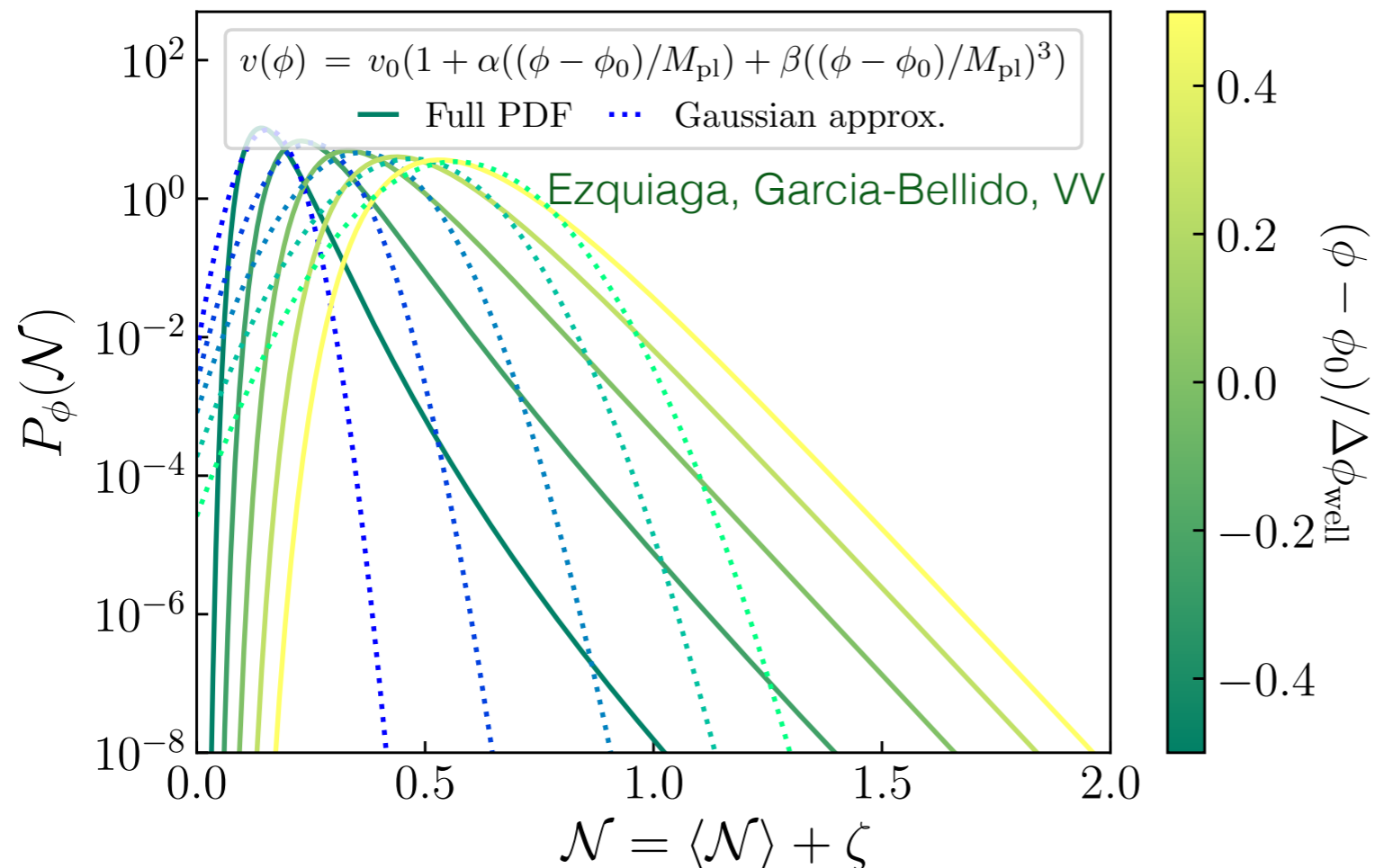
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Need to compute statistics of cosmological perturbations in the presence of non-perturbative stochastic effects

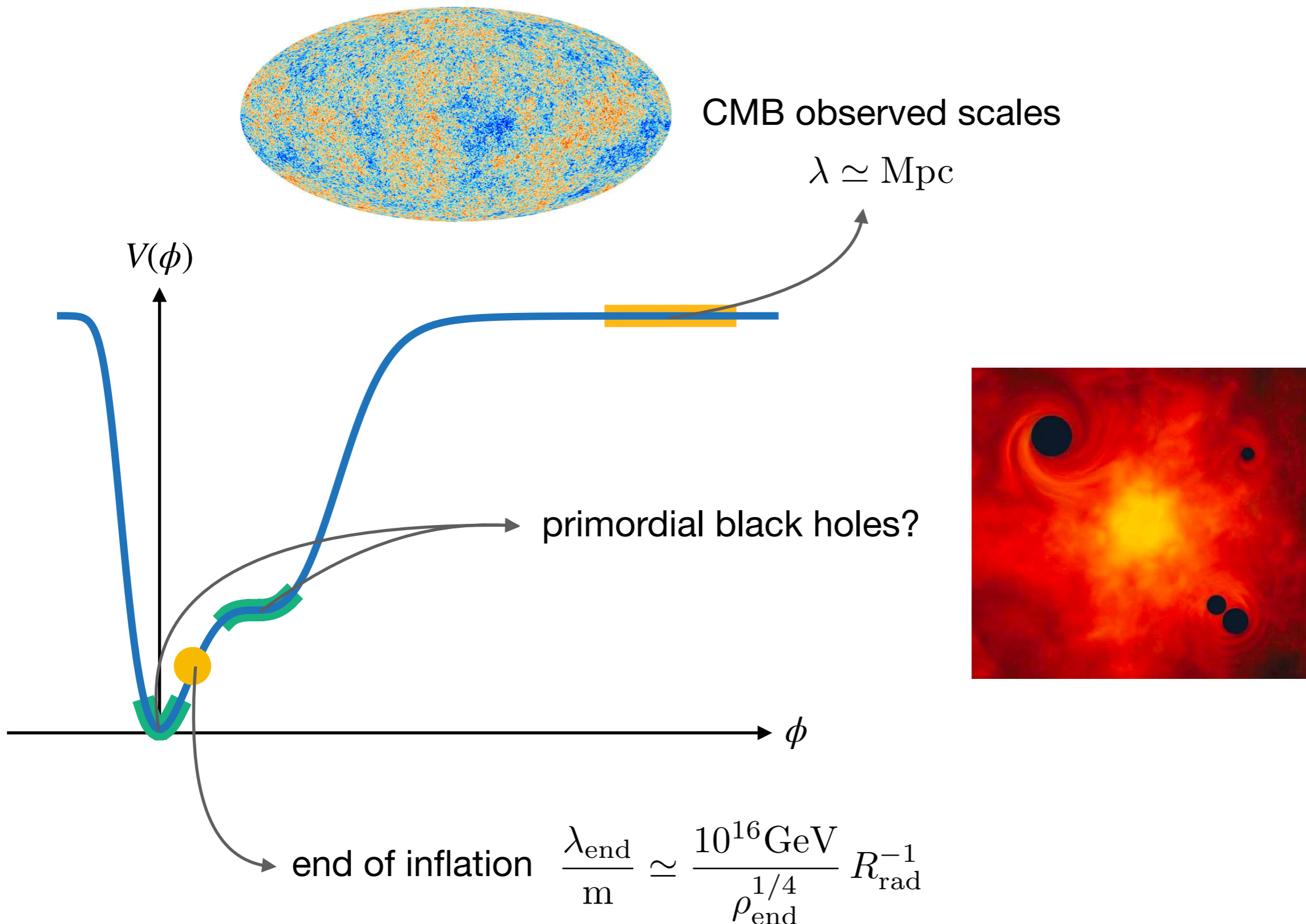


**Stochastic  $\delta N$  formalism**

VV, Starobinsky (2015)  
Pattison, Assadullahi, VV, Wands (2017)

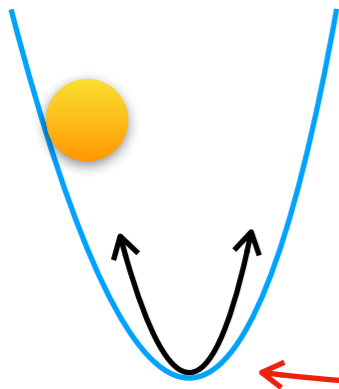


# How can large fluctuations be seeded?



# Case study 2: post-inflationary enhancement

Parametric resonance



$$v_k'' + \left[ k^2 - \frac{(a\sqrt{\epsilon_1})''}{a\sqrt{\epsilon_1}} \right] v_k = 0 \quad \text{with } \epsilon_1 = -\dot{H}/H^2$$

$$\frac{d^2}{dz^2} (\sqrt{a}v_k) + [A_k - 2q \cos(2z)] (\sqrt{a}v_k) = 0$$

with  $z = mt + \pi/4$

$$A_k = 1 + \frac{k^2}{m^2 a^2}, \quad q = \frac{\sqrt{6}}{2} \frac{\phi_{\text{end}}}{M_{\text{Pl}}} \left( \frac{a_{\text{end}}}{a} \right)^{3/2}$$

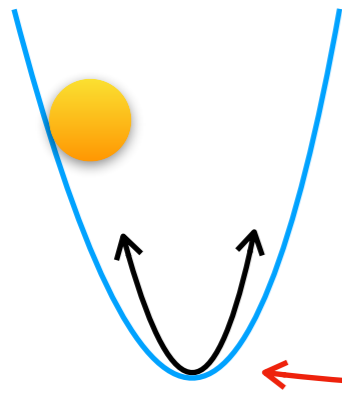
Floquet index:  $\sqrt{a}v_k \propto \exp(\mu_k z)$

Jedamzik, Lemoine, Martin (2010)  
Easter, Flauger, Gilmore (2011)

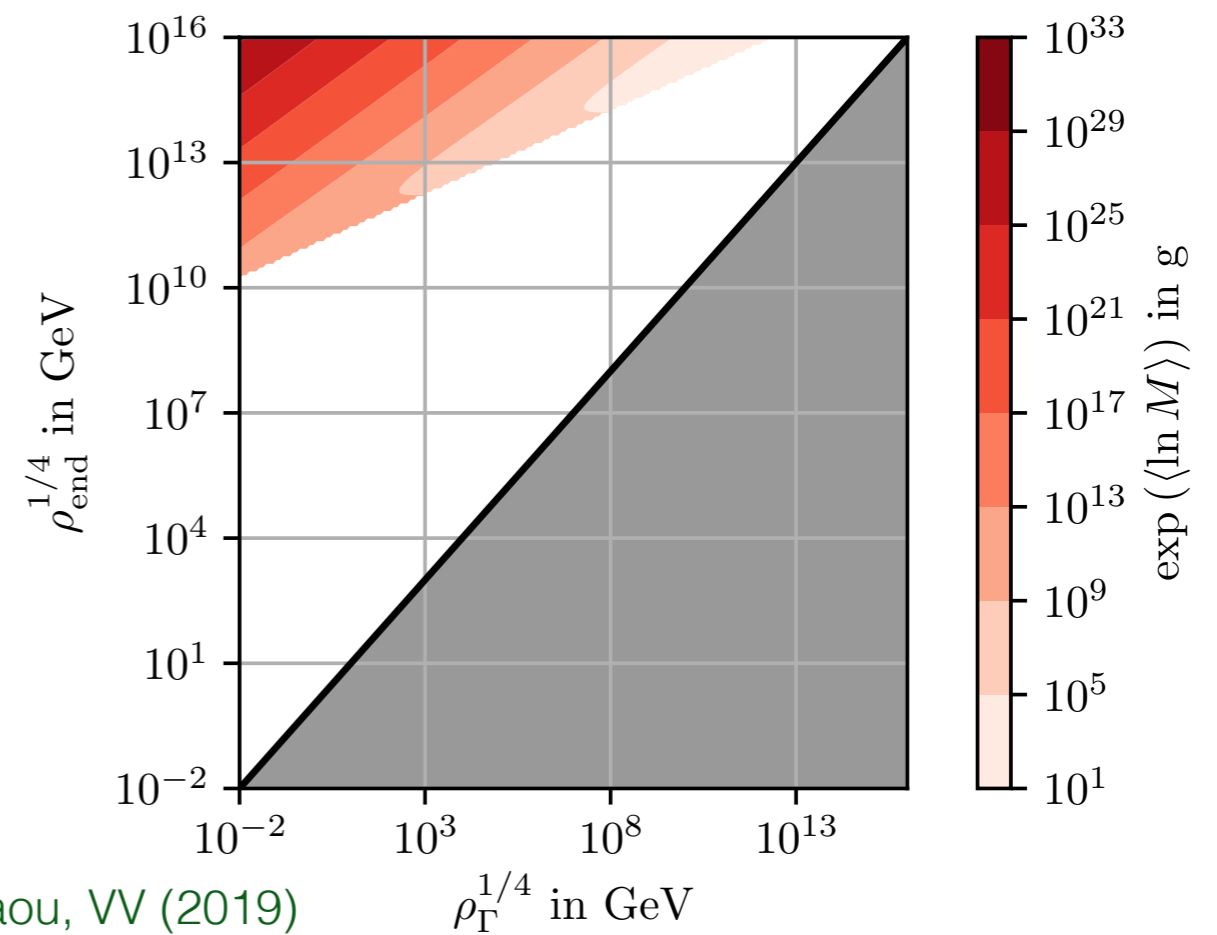
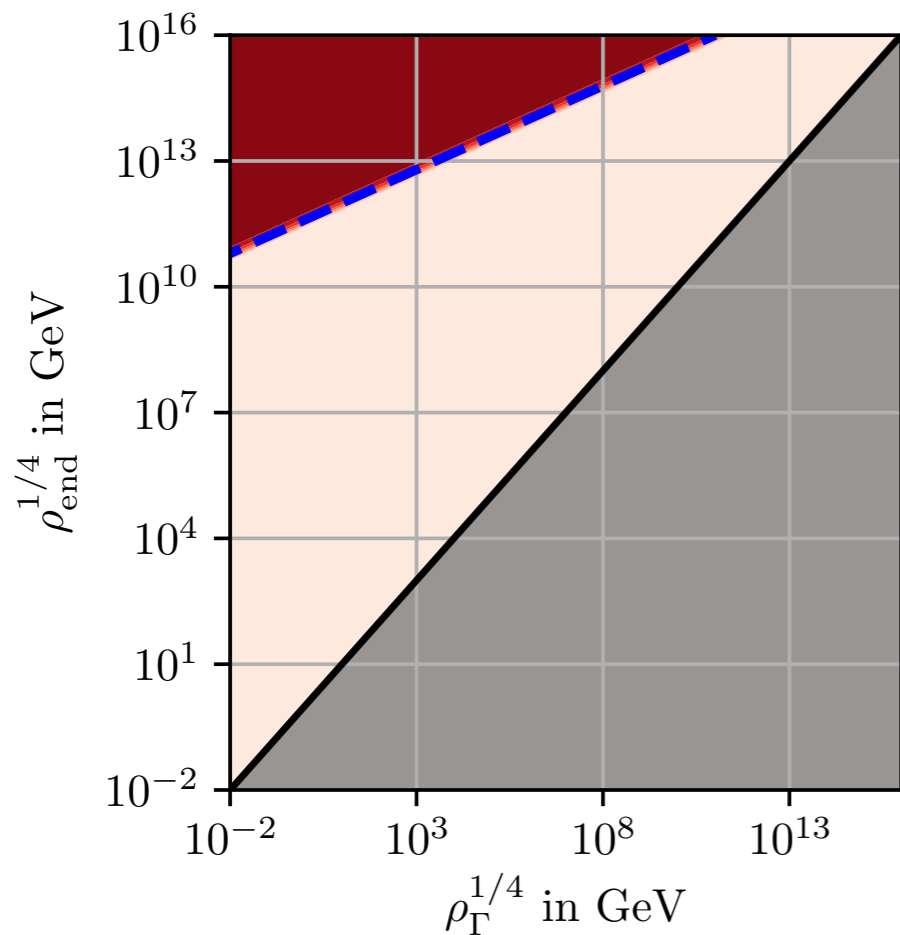


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Martin, Papanikolaou, VV (2019)  
 Martin, Papanikolaou, Pinol, Vennin (2020)  
 Auclair, VV (in prep)

Thank you for your attention!