

Broadband quantum noise reduction in Virgo using frequency dependent squeezing

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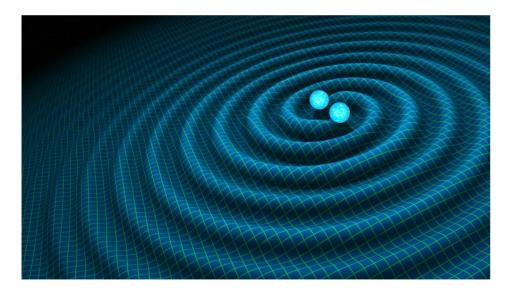


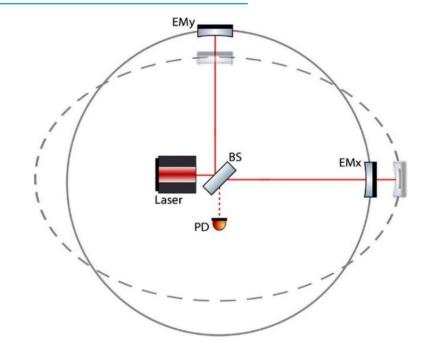
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 Gravitational waves are 'ripples' in the fabric of space-time caused by some of the most violent and energetic processes in the Universe.

 $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu} \longrightarrow g_{\mu\nu}$

small perturbation of the metric tensor





• **Detection principle:** Michelson interferometer measures the difference in phase associated to the passing gravitational wave (GW)

$$\delta \phi_{\rm GW} = \frac{4\pi}{\lambda} \delta L_{\rm GW} \longrightarrow h = \frac{2 \, \delta L_{\rm GW}}{L}$$



Quantum noise in GW detectors

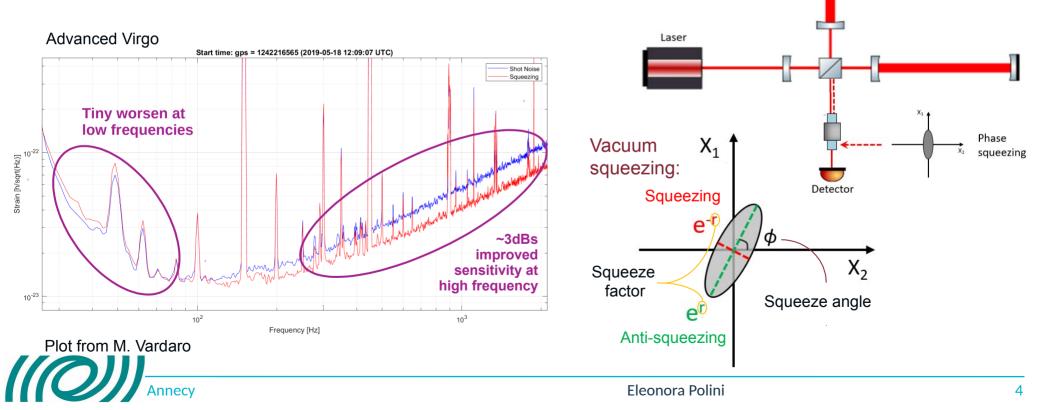
Introduction: • *Quantum noise (QN)* limits the sensitivity of GW detectors QN are vacuum fluctuations entering interferometer's output port Interferometer 10-22 Coating Brownian thermal noise ubstrate thermoelastic nois uspension thermal noise Laser eismic noise strain sensitivity [1/VHz] Quantum noise Total noise 10⁻²³ Amplitude quadrature 10^{3} 10^{2} frequency [Hz] X2 Sensitivity curve of aLIGO Phase Suspended Laser quadrature Photodiode $\hat{\mathcal{H}} = \hbar\omega \left(\hat{a}^{\dagger} \hat{a} + \frac{1}{2} \right) \longrightarrow \hat{\mathcal{H}} = \hbar\omega \left(\hat{X}_1^2 + \hat{X}_2^2 \right)$ Radiation pressure noise Shot noise $\square O$



Interferometer

• First step:

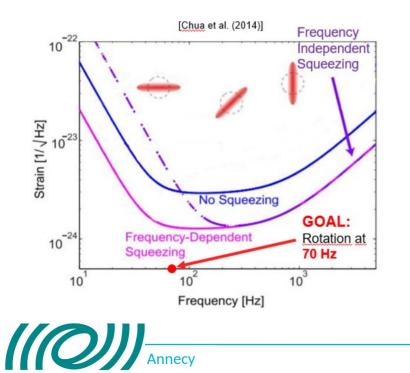
- Injecting squeezed vacuum states from the output port to improve sensitivity, run O3
- Implemented in AdVirgo and aLIGO

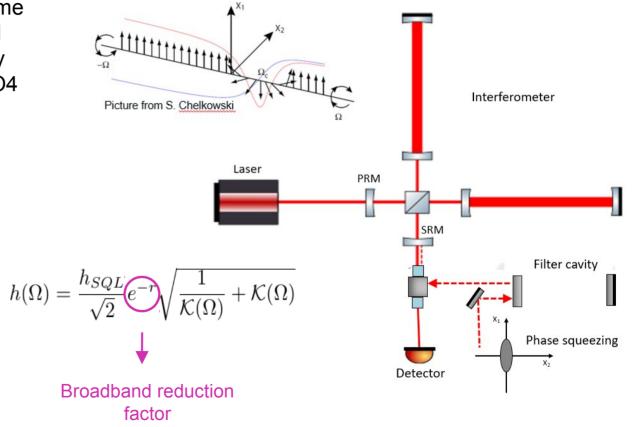




• Next step:

- Vacuum squeezed state angle become frequency dependent when reflected by a detuned Fabry-Perot filter cavity
- Implementation in GW detectors in O4

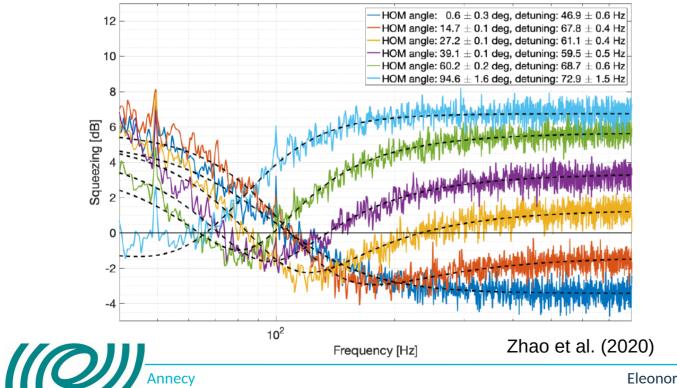


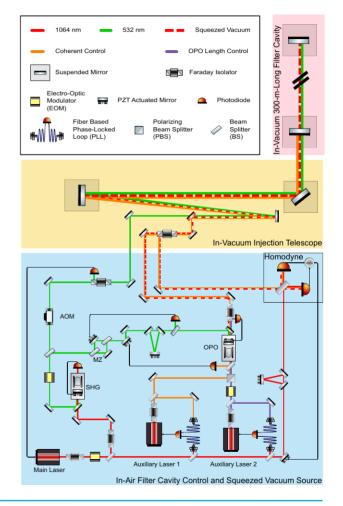




Frequency dependent squeezing (FDS) demonstration

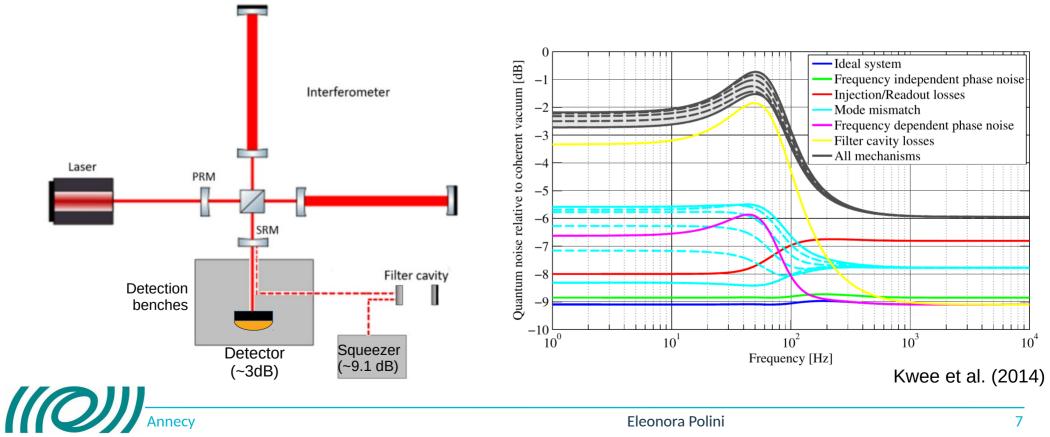
R&D experiment at NAOJ, Tokyo, Japan, was the first demonstration (2020) of a frequency dependent squeezed vacuum source, realized with a 300 m suspended filter cavity. The squeezing rotation takes place in the frequency region (~ 100 Hz) needed to reduce the quantum noise in the whole spectrum of advanced GW detectors.



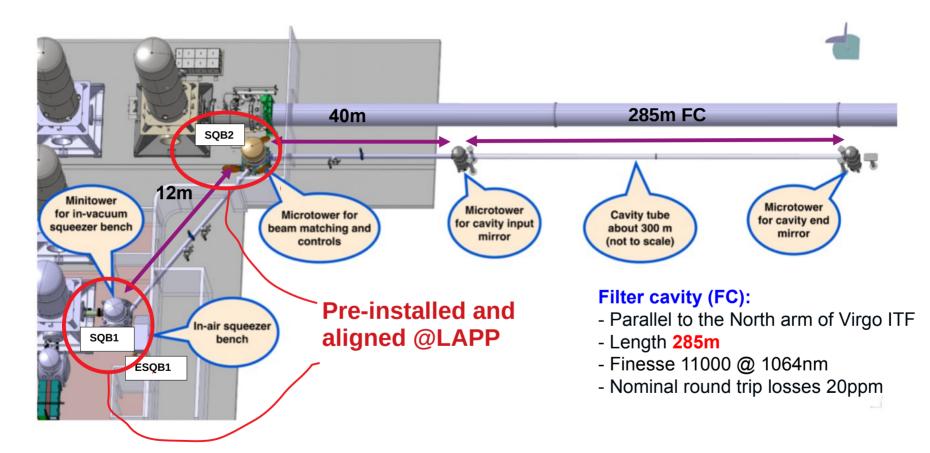




Decoherence (optical losses + mode mismatch) and degradation (phase noise due to phase lock errors + stray light + cavity length fluctuations) mechanisms limit the experimentally achievable QN reduction.

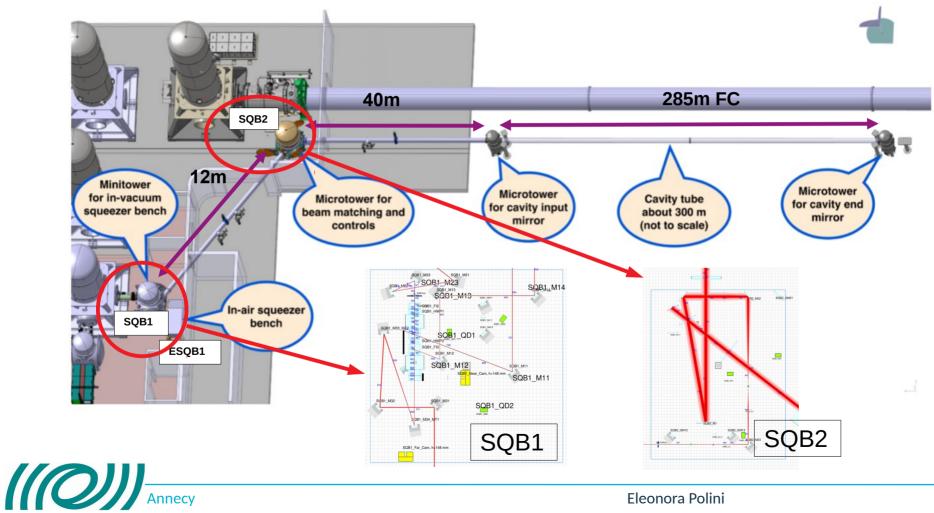








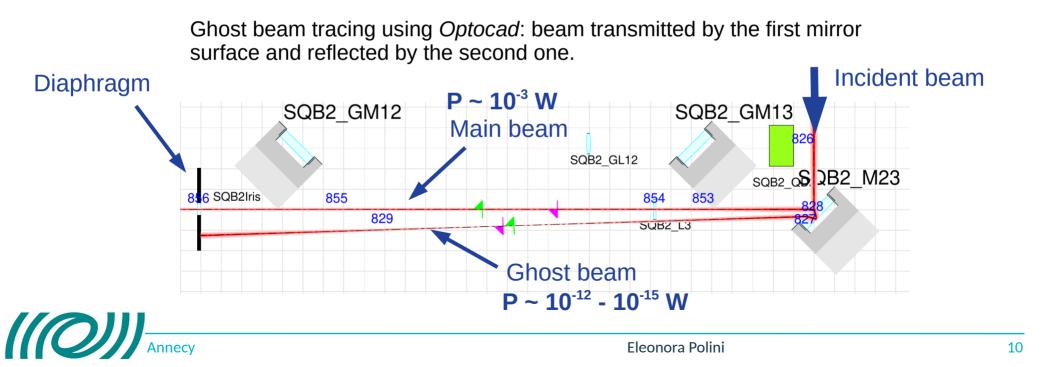
FDS implementation in AdV+





Ghost beams: secondary beams generated by not perfect mirror coatings:

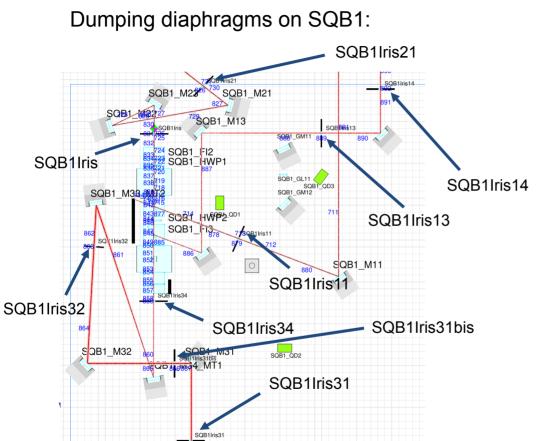
1) we want to <u>dump</u> these beams to **avoid scattered light** on squeezing benches (it has been proved that the squeezing sensitivity enhancement is affected by stray light: <u>Virgo Logbook #48337</u>, <u>#44990</u>);



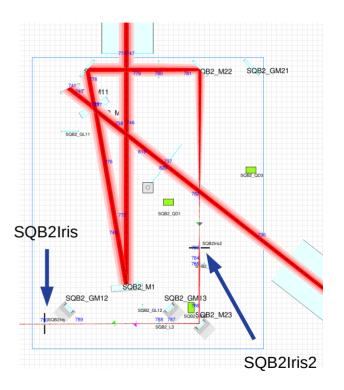


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Annecy

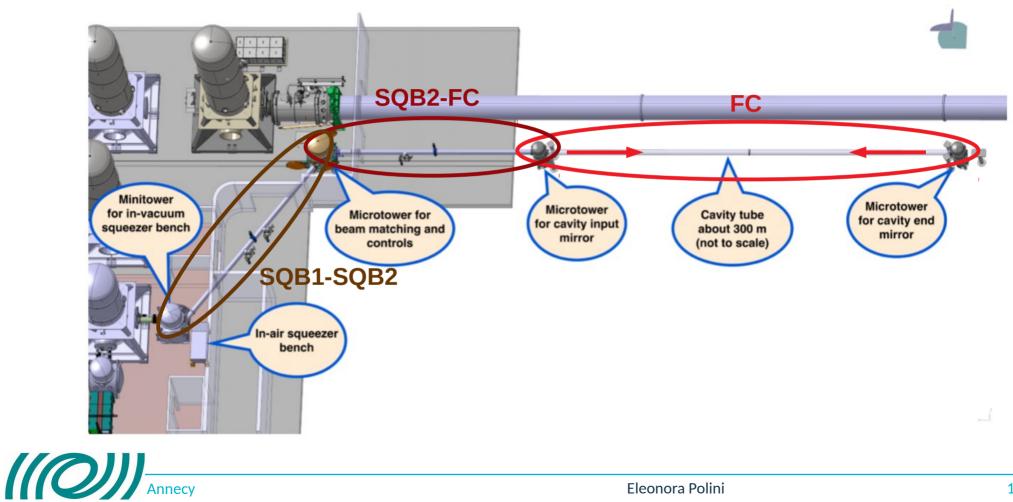


Dumping diaphragms on SQB2:



AdV TDS: VIR-0549B-20



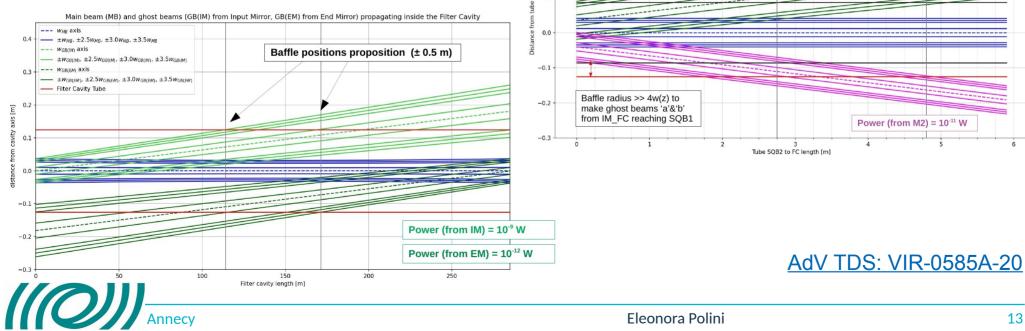




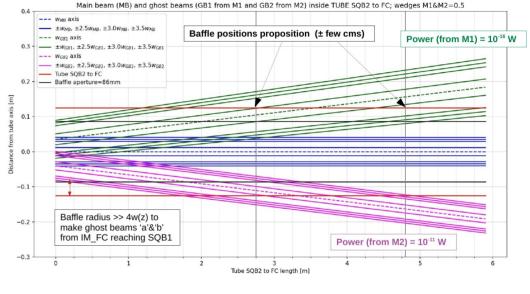
Baffles inside FC and linking tubes

Baffle

Baffles inside FC:



Baffles inside linking tube SQB2-FC:





2) we need to chose an optimal size of diaphragm/baffle apertures to limit losses on the main beam.

Power losses

 10^{-1}

 10^{-3}

 10^{-5}

 10^{-7}

 10^{-9}

 10^{-11}

 10^{-13}

 10^{-15}

Power Losses Normalized

Since the HR coating has 3ppm power losses, we want < 1ppm losses, up to 1 w(z) displacement.

d=1.0w(z)

d = 1.5w(z)

nnecv

2

Aperture radius / beam radius

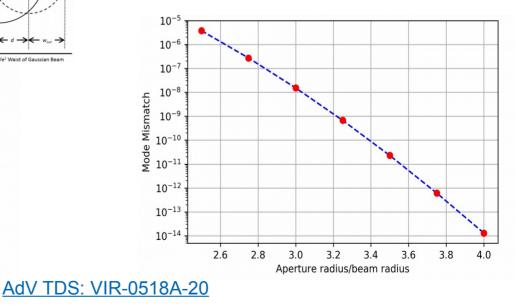
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Gaussian Beam Centered at Gaussian Beam shifted center of circular aperture distance. d. relative to center of circular aperture Outline of 1/e² Waist of Gaussian Bear Circumference of Circular Aperture d=0w(z)d = 0.5w(z)

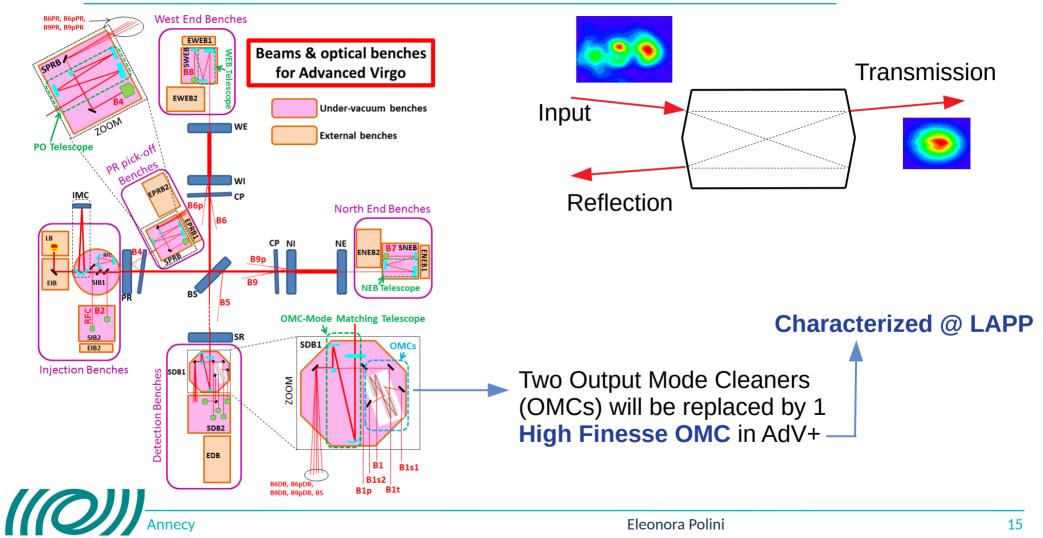
Mode mismatch losses

The total MM requirement for FDS is 1%. We can conclude that mode mismatch due to diffraction is **negligible**.



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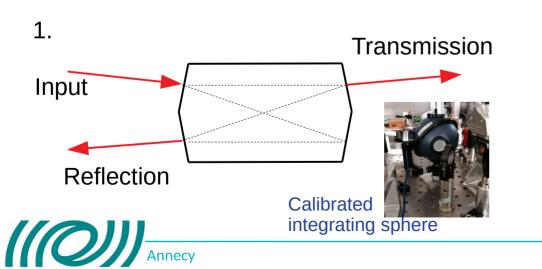


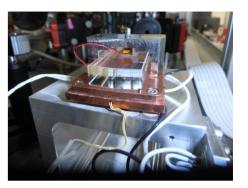


High Finesse OMC characterization

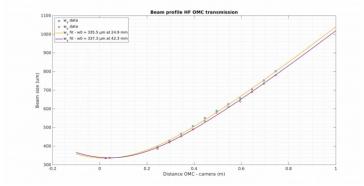
We are characterizing the OMC at LAPP:

- 1. Losses = 3.5% ± 0.3 % (during O3 Losses = 5%) Losses = Input power – Transmitted power – Reflected power
- 2. Radius of Curvature (ρ) of the spherical surface: $\rho \sim 1700$ mm
- 3. High reflective coating residual transmission: few ppm





2.
$$w_0 = \sqrt{\frac{\lambda}{n\pi}\sqrt{2L_{geo}(\rho - 2L_{geo})}}$$



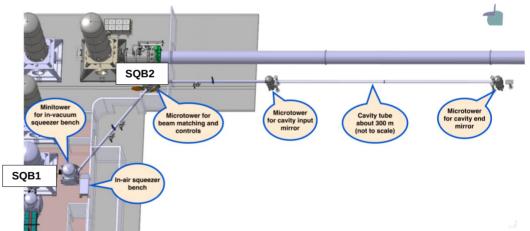


Conclusions:

- Ghost beams study on Frequency Dependent Squeezing system
- Design of diaphragms and baffles
- Construction of diaphragms and baffles
- Mechanical installation of FDS system at Virgo
- Characterization of OMC at LAPP

Next steps:

- Pre-installation of SQB1 and SQB2 at LAPP (December 2020)
- Commissioning of the whole FDS system on site (beginning 2021)
- Commissioning of OMC on site (November 2020)
- X Observing run O4 starts in 2022







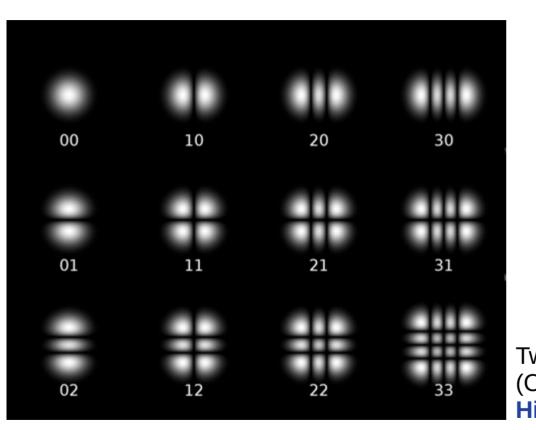


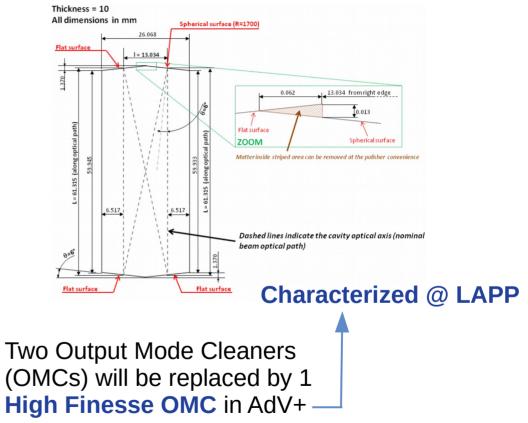


EXTRA SLIDES



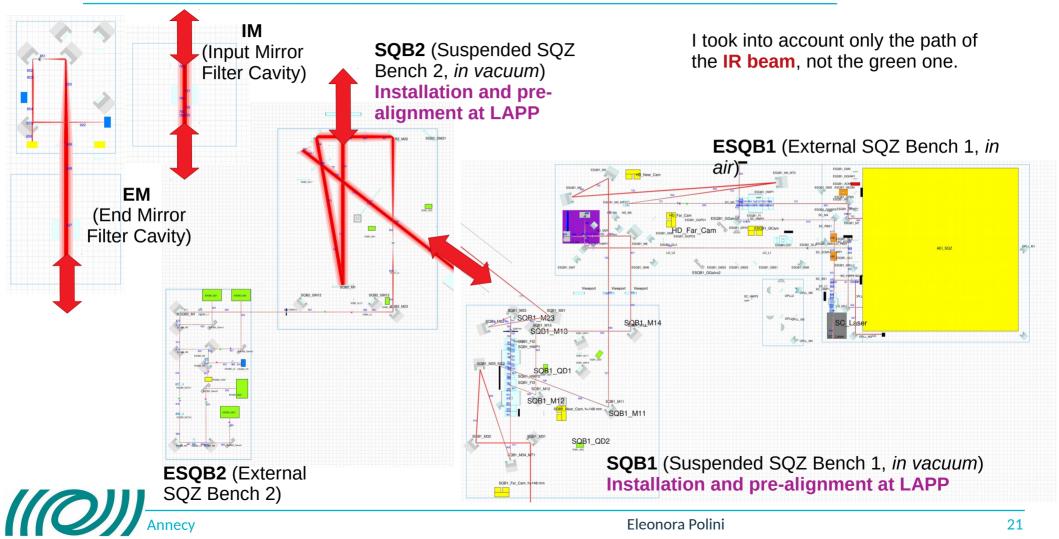




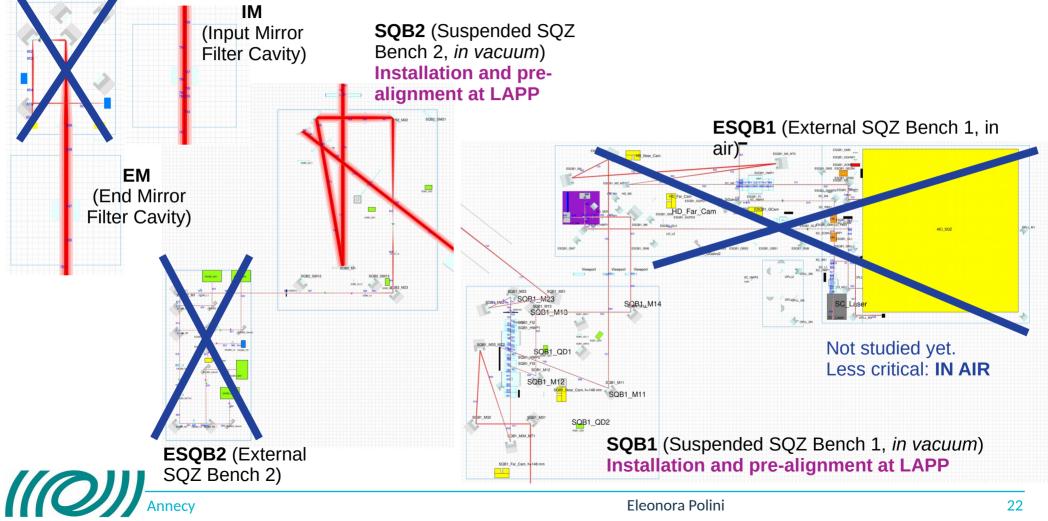










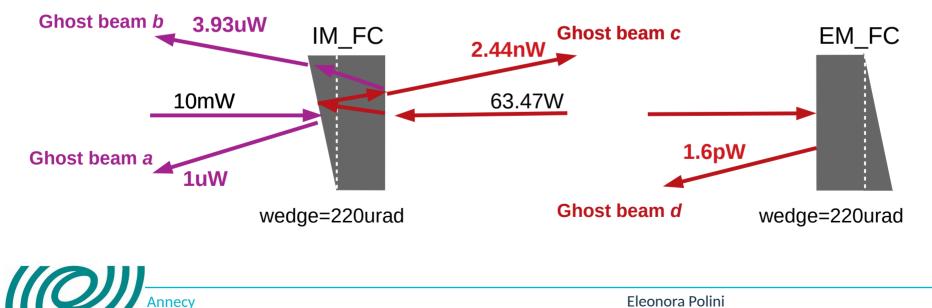




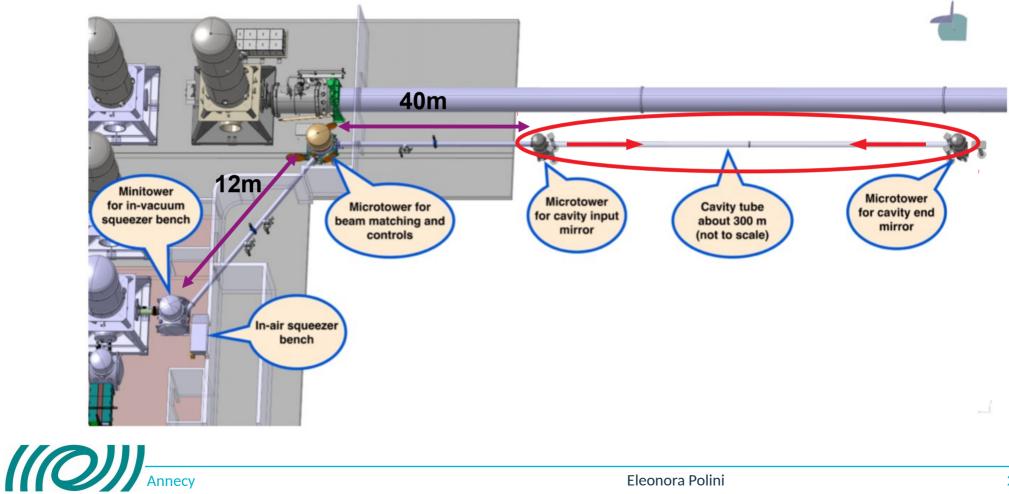
I studied the ghost beams propagation inside different tubes (VIR-0585A-20): Filter Cavity (both Input and End Mirrors), Tube between SOB2 and FC, Tube between SOB2 and SOB1. We are designing the **baffles** to install inside these three tubes.

Filter Cavity (VIR-0473A-20, VIR-0584A-20):

- These ghost beams are way **more powerful** than the others, so we need to dump them properly;
- Ghost beams 'a' & 'b' are dumped on diaphragms on SQB1;
- Ghost beams 'c' & 'd' are symmetrical and dumped on baffles inside FC.



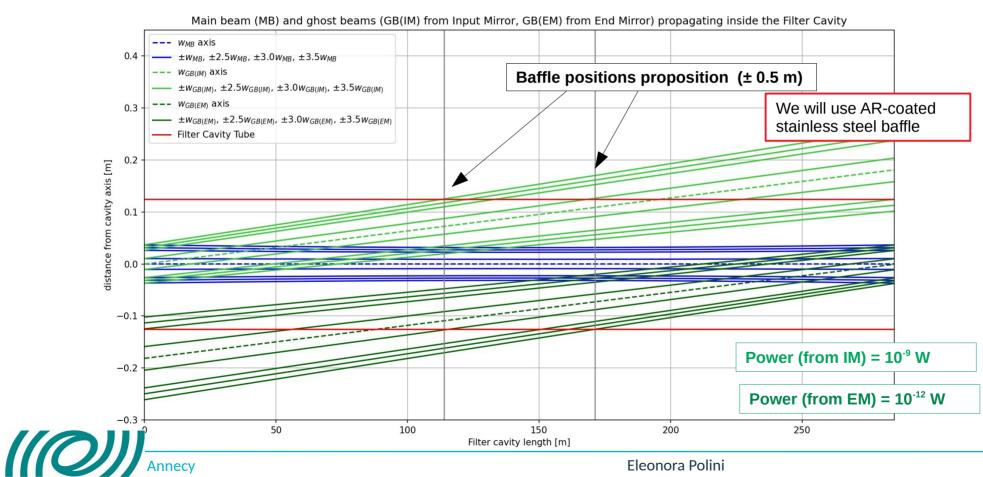


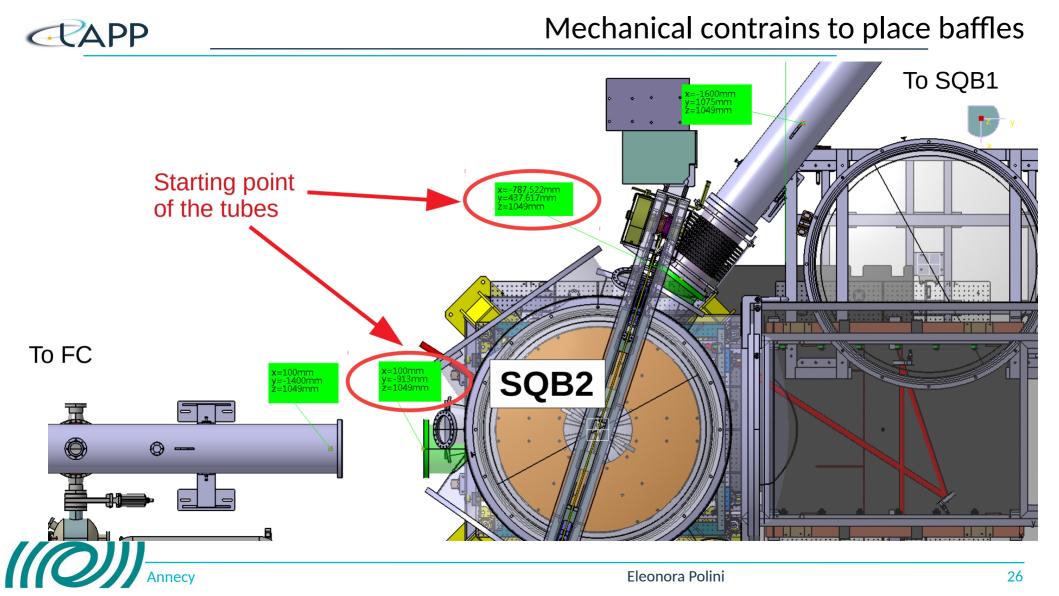




Discussion with FLT to install baffles at different locations inside the tube: **IN PROGRESS**.

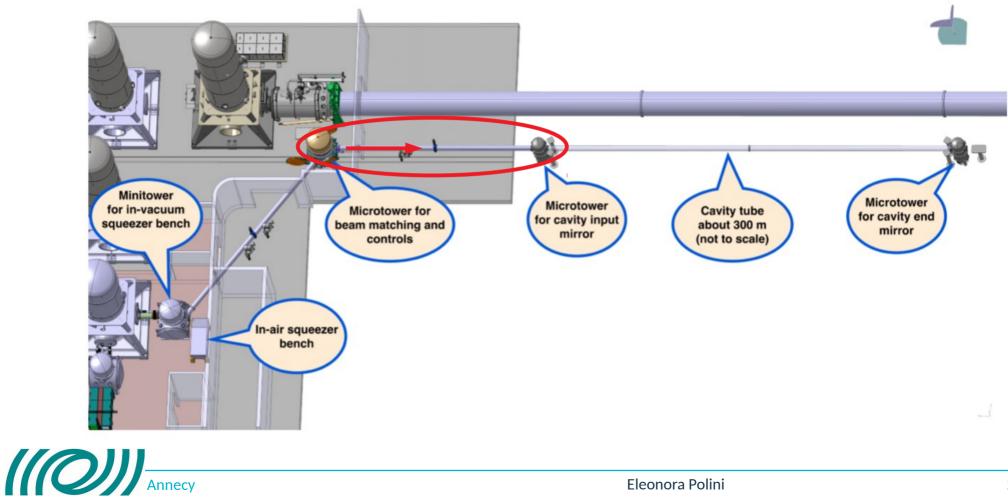
Constraints: available mechanical positions inside FC, do not cut the Ghost Beam in two with the baffle, minimize stray light etc.





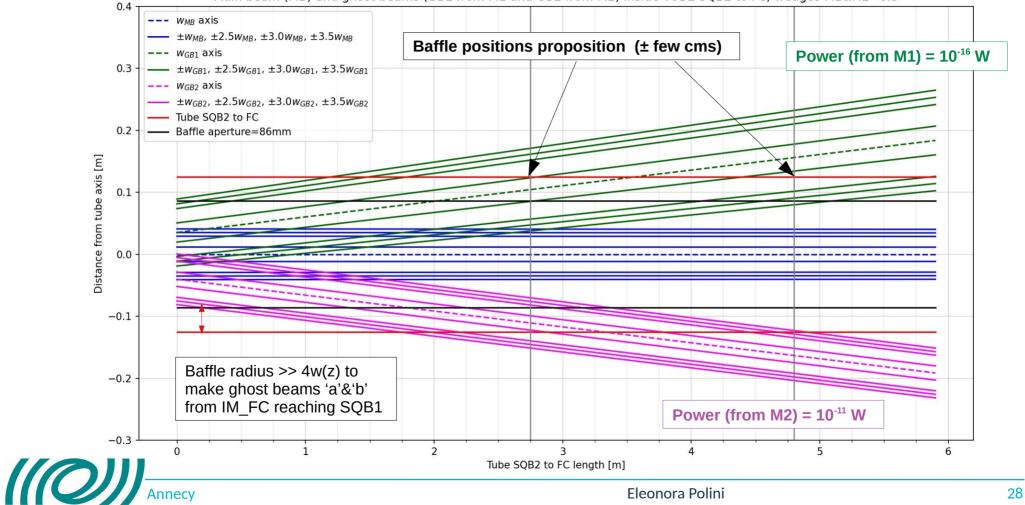


Propagation inside the tubes : **SQB2 to FC**

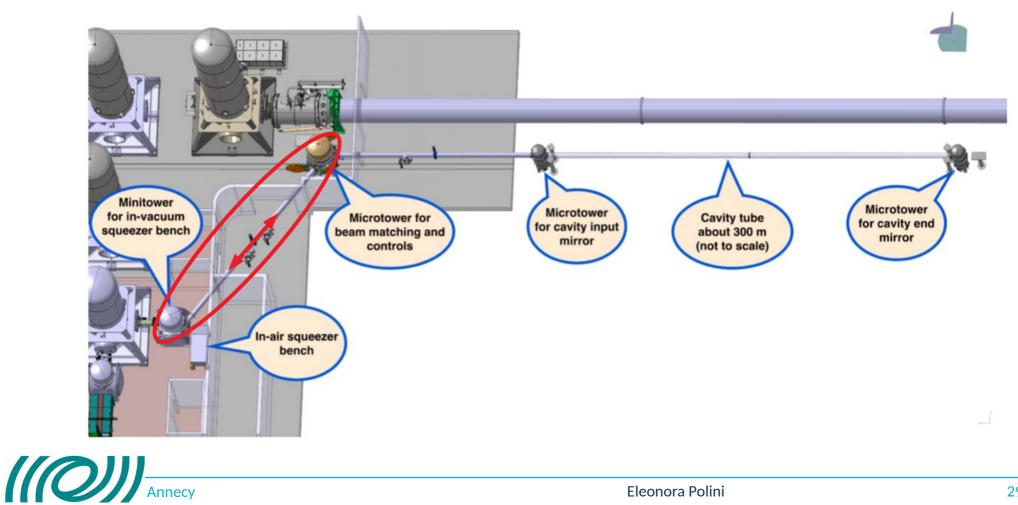




Main beam (MB) and ghost beams (GB1 from M1 and GB2 from M2) inside TUBE SQB2 to FC; wedges M1&M2=0.5

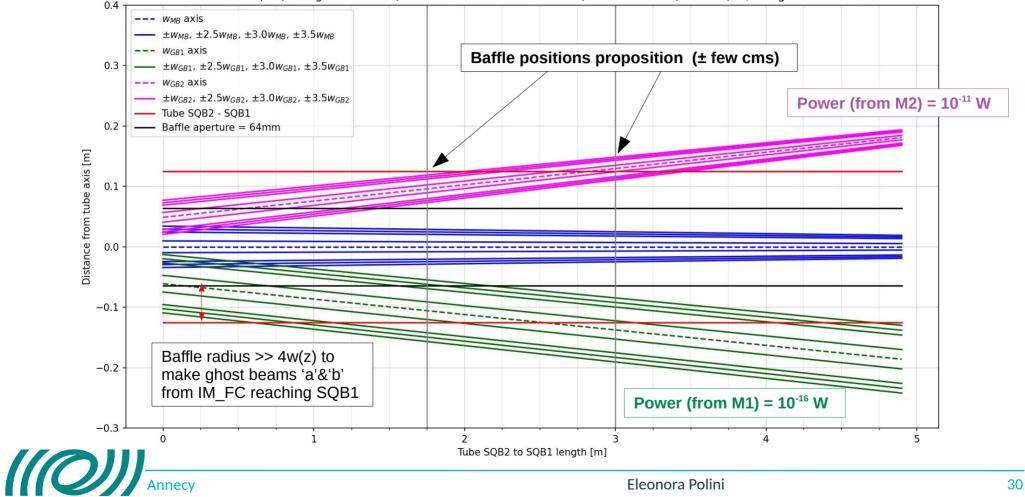








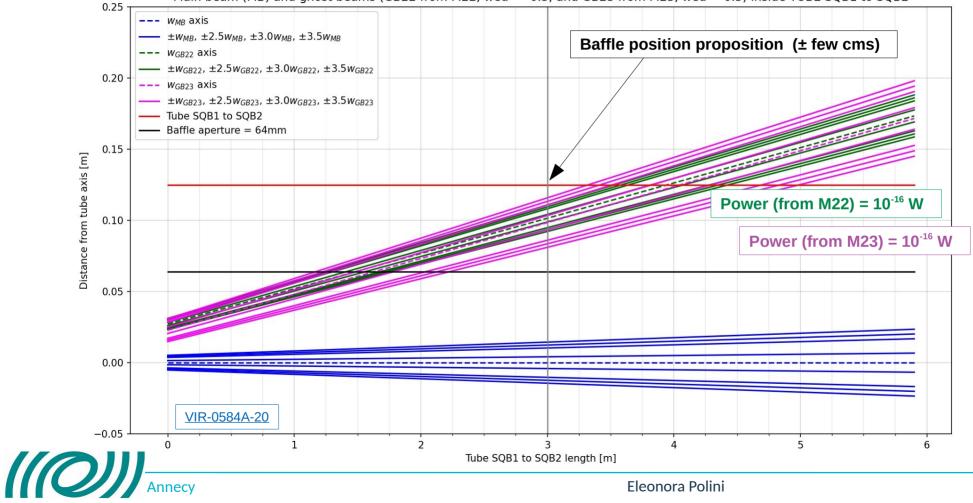
Main beam (MB) and ghost beams (GB1 from M1 and GB2 from M2) inside TUBE SQB2 to SQB1; wedges M1&M2=0.5





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Main beam (MB) and ghost beams (GB22 from M22, wed = -0.5, and GB23 from M23, wed = 0.5) inside TUBE SQB1 to SQB2



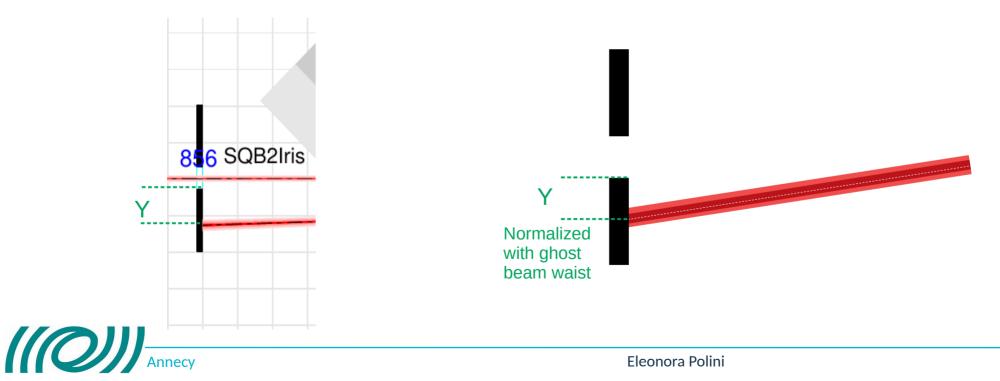


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To understand if the ghost beams are stopped by diaphragms and baffles properly, we introduce a quantity Y defined as:

Y = (distance ghost and main beams – aperture radius)/ghost beam radius.

The <u>fraction of ghost beam power</u> passing through the aperture can be computed using the corrected error function Erfc(Y) = 1 - Erf(Y).





 \square

I took 'standard apertures' in order to ease the production and installation processes. With a = 4, 5.5, 6.5 mm we don't change the performances of diaphragms on benches. Concerning baffles, we need to consider stray light effect at the aperture edges.

