

# Interplay between astrophysical and laboratory probes of MeV-scale ALPs

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in collaboration with:

Felix Kahlhoefer - arXiv: 2004.01193

May 11th, 2020

# Outline

- Motivation
- Theoretical Framework
- Experimental and Observational Constraints
- Results

# Motivation

- Extensions of the SM? Heavy new particles!
- Despite all experimental efforts so far, no (conclusive) evidence for heavy particles at TeV scale
- Maybe missed some physics along the way?
  - renewed attention for new, much lighter particles
  - remain to be discovered, if they have small interactions
- Popular candidate: Axion-like particle

# Goldstone bosons

- Theoretical motivation: Goldstone's theorem
  - Spontaneous breaking of an (approximate) global symmetry gives rise to (Pseudo-)Goldstone bosons
    - naturally light (or even massless)
    - interactions are suppressed by scale of symmetry breaking
- common explanation for small mass and tiny interactions
- Example:  $U(1)_{\text{PQ}}$  symmetry  $\longleftrightarrow$  QCD axion

# Axion-like particles

- Unlike QCD axion employ model-independent approach
  - Couplings are free parameters
  - Treat mass as an independent parameter

⇒ Axion-like particles (ALPs)

# Axion-like particles

- Unlike QCD axion employ model-independent approach
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    - Treat mass as an independent parameter
- ⇒ Axion-like particles (ALPs)
- ALP couplings therefore include in general
    - SM fermions
    - SM Higgs
    - Gluons
    - Electroweak gauge bosons

# General ALP EFT

- Write down general EFT

[H. Georgi et al., Phys.Lett. 169B (1986), M. Bauer et al., arXiv: 1708.00443]

$$\begin{aligned} \mathcal{L}_{\text{eff}} = & \frac{1}{2}(\partial_\mu a)(\partial^\mu a) - \frac{m_{a,0}^2}{2}a^2 + \frac{\partial^\mu a}{\Lambda} \sum_F \bar{\Psi}_F C_F \gamma_\mu \Psi_F \\ & + g_s^2 C_{GG} \frac{a}{\Lambda} G_{\mu\nu}^a \tilde{G}^{\mu\nu,a} + g^2 C_{WW} \frac{a}{\Lambda} W_{\mu\nu}^a \tilde{W}^{\mu\nu,a} + g'^2 C_{BB} \frac{a}{\Lambda} B_{\mu\nu} \tilde{B}^{\mu\nu} \end{aligned}$$

- $\Psi_F$ : chiral fermion fields
- $X_{\mu\nu}^a$ :  $SU(3)_C$ ,  $SU(2)_L$  and  $U(1)_Y$  field strength tensors

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# Relations between couplings

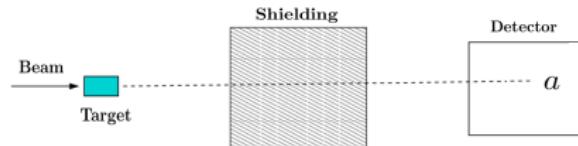
- EWSB generates couplings to  $W$  bosons, photons etc.
  - $W$  boson  $\leftrightarrow C_{WW}$
  - Photon coupling  $C_{\gamma\gamma} = C_{WW} + C_{BB}$
  - $Z\gamma a$  coupling  $C_{\gamma Z} = c_W^2 C_{WW} - s_W^2 C_{BB}$
- But also gluonic contribution ( $m_a \lesssim \Lambda_{\text{QCD}}$ )
  - $C_{\gamma\gamma}^{\text{eff}} \approx C_{WW} + C_{BB} - 1.92 C_{GG} - \frac{1}{3} \frac{m_a^2}{m_\pi^2 - m_a^2} C_{GG}$

[M. Bauer et al., arXiv: 1708.00443]

⇒ Strong correlation between interactions

# Laboratory probes

- Which experiments do we consider?
  - Fixed-target experiments:



[S.N. Gninenko, arXiv: 1204.3583]

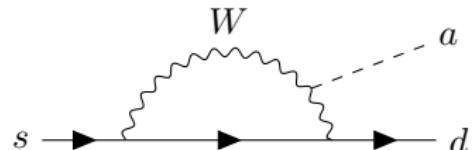
- L3 at LEP performed search for  $Z \rightarrow \gamma + \text{inv.}$  at Z resonance

# Experimental and Observational Constraints

- Rare Meson decays:

- Experimental constraints by E787/E949 on  $K^+ \rightarrow \pi^+ a$
- NA62 focusses on  $K^+ \rightarrow \pi^+ \nu \bar{\nu}$   
→ Assume improvement on bound by an order of magnitude

[R. Fantechi et al., arXiv: 1407.8213]

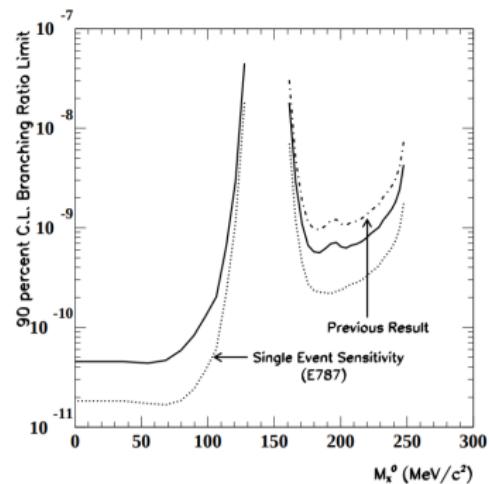


- Ensure invisible decay in experiment

$$\text{BR}(K^+ \rightarrow \pi^+ a)^{\text{exp}} = \text{BR}(K^+ \rightarrow \pi^+ a)^{\text{theo}} e^{-L_{\text{det}}/L_{\text{alp}}}$$

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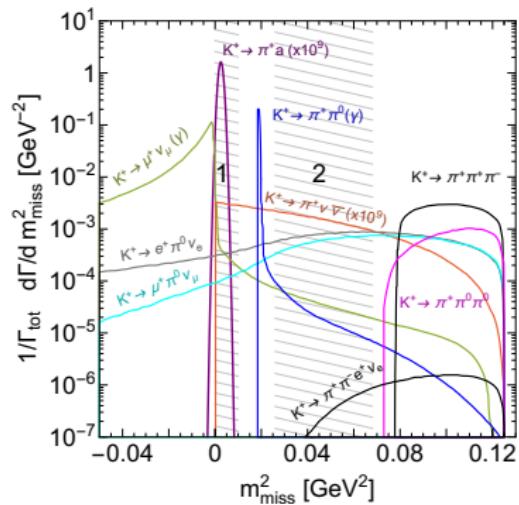
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[S. Adler et al., E787, arXiv: 0403034]

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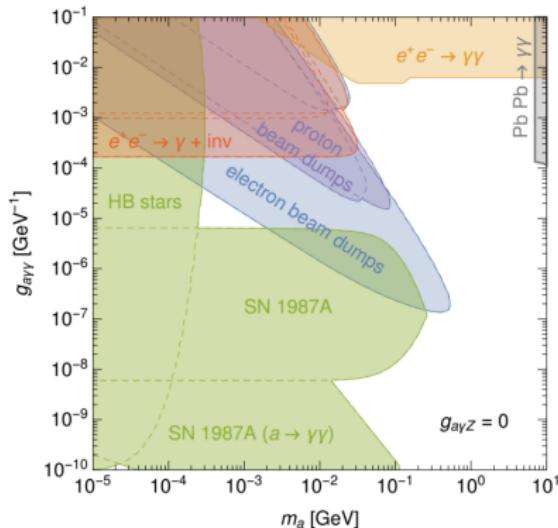
[E. C. Gil et al., NA62, arXiv: 1703.08501]

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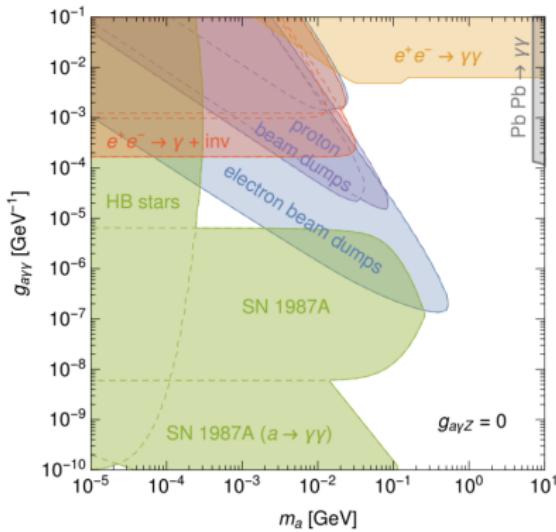
- Electron beam dump (E137)  
 strongest fixed-target bound
- HB star dominant for  $m_a \lesssim 10 \text{ keV}$
- SN 1987A (“cooling” bound)
  - Formation of proto-neutron star
  - Most of energy liberated by neutrino emission



[M. Dolan et al., arXiv: 1709.00009]

# Experimental and Observational Constraints

- SN bound rather simplistic
  - No temperature and density profiles
  - No radius and energy dependence of ALP trapping
    - improve calculation
    - consistent approach:  $C_{\gamma\gamma}$  &  $C_{GG}$
- Below cooling bound: missing ALP burst



[M. Dolan et al., arXiv: 1709.00009]

# SN Bound

- Apply Raffelt criterion for SN bound:  $L_{\text{ALP}} > L_\nu \not\propto$
- ALP Luminosity reads schematically

[J.H. Chang et al., arXiv: 1611.03864]

$$L_{\text{ALP}} \sim \int_{r \leq R_\nu} dV Q e^{-\tau}$$

with  $Q$  as volume emission rate and optical depth  $\tau$

$$Q \sim \int \frac{d^3 k_a}{(2\pi)^3} E_a \Gamma_{\text{prod}}, \quad \tau \sim \int_r^{R_{\text{far}}} dr' \Gamma_{\text{abs}}$$

# ALP Luminosity Formula

- Luminosity formula is given by ( $E_a \equiv \omega$ )

$$L_{\text{ALP}} = \int_{r \leq R_\nu} dV \left( \int \frac{d^3 \mathbf{p}_a}{(2\pi)^3} \omega \Gamma_a e^{-\omega/T} \beta e^{-\tau} + \int \frac{2d^3 \mathbf{p}_\gamma}{(2\pi)^3} \frac{\omega \Gamma_{\gamma \rightarrow a}}{e^{\omega/T} - 1} \beta e^{-\tau} \right)$$

↑  
Nucleon-Scattering:  $N + N \rightarrow N + N + a$

↑  
Primakoff:  
 $\gamma + p \rightarrow a + p$

with  $\beta = \sqrt{1 - m_a^2/\omega^2}$

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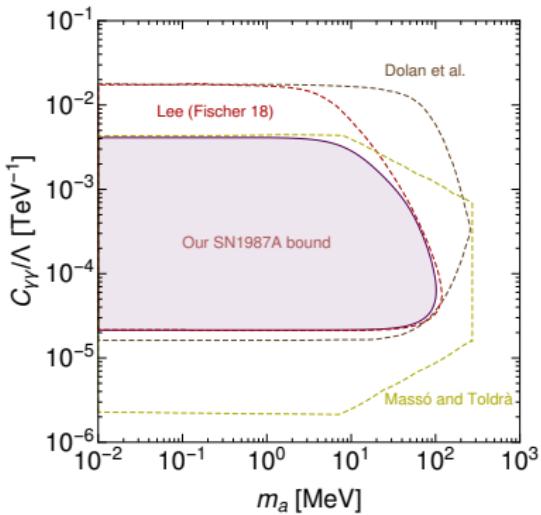
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with  $\beta = \sqrt{1 - m_a^2/\omega^2}$  and the optical depth  $\tau$  reading

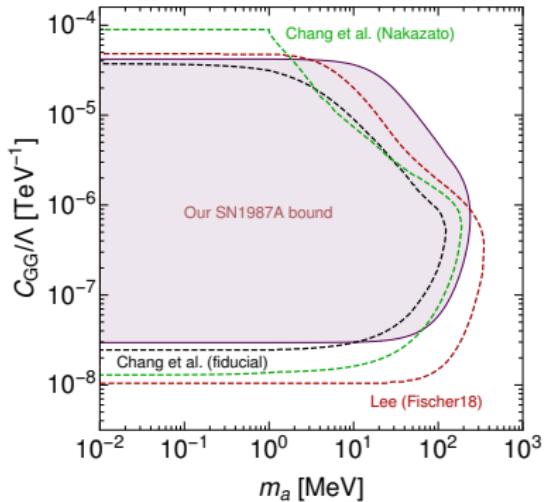
$$\tau = \beta^{-1} \int_r^{R_{\text{far}}} dr' \left( 2\Gamma_{\gamma \rightarrow a} + \gamma^{-1} \Gamma_{a\gamma\gamma} + \Gamma_a \right)$$

# Comparison of SN bounds

## Photon coupling



## Gluon coupling



# Reminder about couplings

- Lagrangian before EWSB

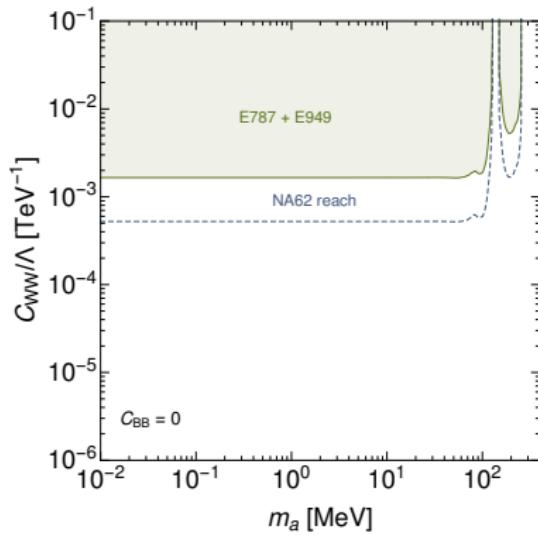
$$\mathcal{L}_{\text{eff}} = g_s^2 C_{GG} \frac{a}{\Lambda} G_{\mu\nu}^a \tilde{G}^{\mu\nu,a} + g^2 C_{WW} \frac{a}{\Lambda} W_{\mu\nu}^a \tilde{W}^{\mu\nu,a} + g'^2 C_{BB} \frac{a}{\Lambda} B_{\mu\nu} \tilde{B}^{\mu\nu}$$

- After EWSB

- W boson  $\leftrightarrow C_{WW}$
- $C_{\gamma\gamma}^{\text{eff}} \approx C_{WW} + C_{BB} - 1.92 C_{GG} - \frac{1}{3} \frac{m_a^2}{m_\pi^2 - m_a^2} C_{GG}$
- $Z\gamma a$  coupling  $C_{\gamma Z} = c_W^2 C_{WW} - s_W^2 C_{BB}$

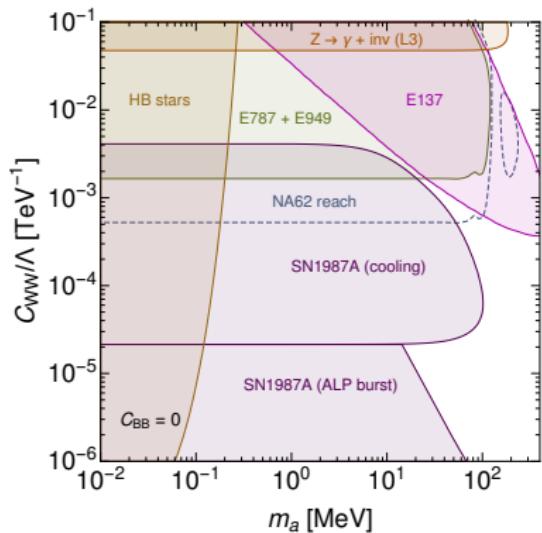
# Couplings to electroweak gauge bosons only

- Consider sizeable  $C_{WW}$



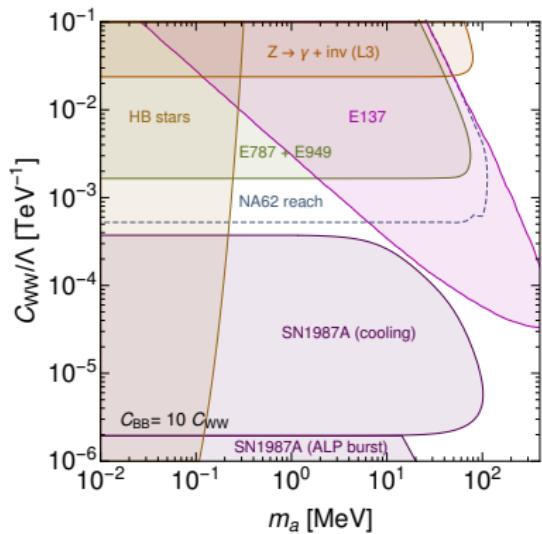
# Couplings to electroweak gauge bosons only

- Consider sizeable  $C_{WW}$
- Have  $C_{\gamma\gamma}^{\text{eff}} \approx C_{WW}$   
 $\rightarrow \text{BR}(a \rightarrow \gamma\gamma) \approx 1$
- How to open up more parameter space?
- Reminder:  
 $C_{\gamma\gamma} = C_{WW} + C_{BB}$



# Couplings to electroweak gauge bosons only

- Consider  $C_{BB} = 10C_{WW}$
- $C_{\gamma\gamma} = C_{WW} + C_{BB}$ 
  - Shifts bounds dependent on  $C_{\gamma\gamma}$
  - In particular SN bounds are shifted downwards
- Further Options?



# Couplings to gluons

- Turn on sizeable gluon coupling

[G. Alonso-Alvarez et al., arXiv: 1811.05466]

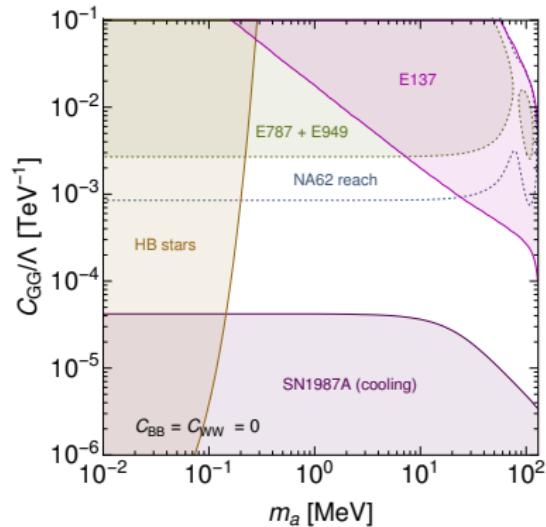
- Contributes to photon coupling

$$C_{\gamma\gamma}^{\text{eff}} \approx -1.92 C_{GG} - \frac{1}{3} \frac{m_a^2}{m_\pi^2 - m_a^2} C_{GG}$$

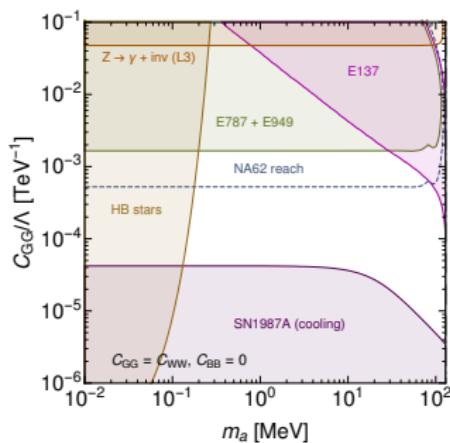
- $K^+ \rightarrow \pi^+ a$  via  $\pi^0, \eta, \eta'$  mixing

[D.S.M. Alves and N. Weiner, arXiv: 1710.03764]

→ How to circumvent the uncertainties?

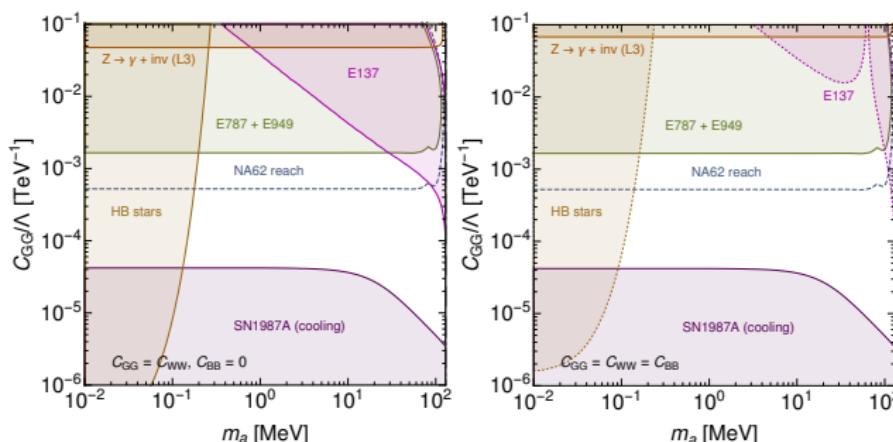


# Couplings to all Standard Model gauge bosons



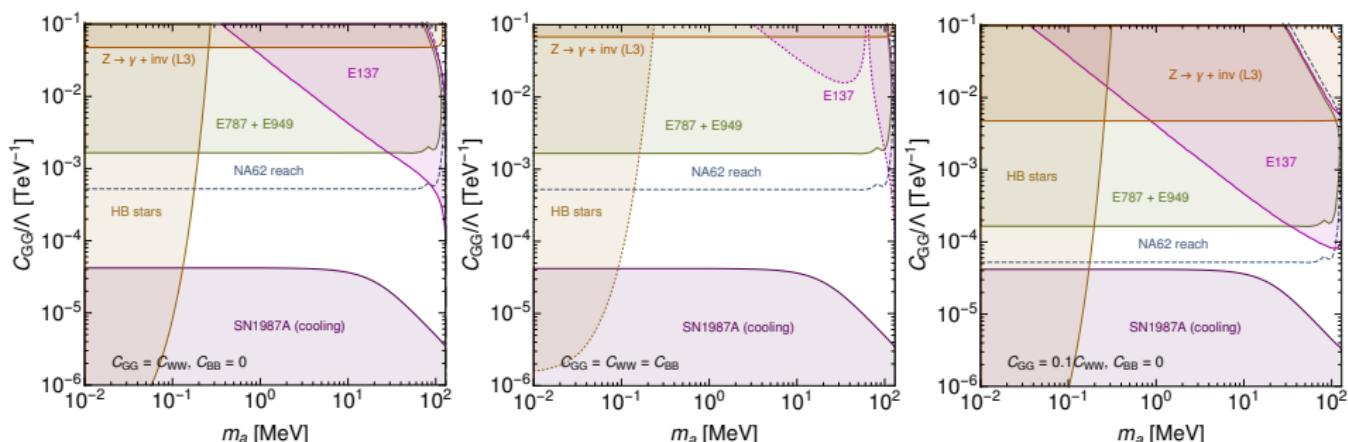
- $C_{GG} = C_{WW}$
- $K^+ \rightarrow \pi^+ a$  via  $C_{WW}$

# Couplings to all Standard Model gauge bosons



- $C_{GG} = C_{WW}$
- $K^+ \rightarrow \pi^+ a$  via  $C_{WW}$
- $C_{GG} = C_{WW} = C_{BB}$
- Cancellation:  
 $C_{\gamma\gamma}^{\text{eff}} \approx 0.08 C_{GG}$

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- $C_{GG} = C_{WW}$
- $K^+ \rightarrow \pi^+ a$  via  $C_{WW}$
- $C_{GG} = C_{WW} = C_{BB}$
- Cancellation:  
 $C_{\gamma\gamma}^{\text{eff}} \approx 0.08 C_{GG}$
- $C_{GG} = 0.1 C_{WW}$
- Lab-Experiment reaches region of SN1987A

# Conclusion

- Studied MeV-scale ALPs in EFT framework
  - Particular focus on  $K^+ \rightarrow \pi^+ a$
- Consistent approach to calculation of SN cooling bound
  - Bound on  $C_{\gamma\gamma}$  more conservative, while on  $C_{GG}$  in agreement with literature

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- Studied MeV-scale ALPs in EFT framework
  - Particular focus on  $K^+ \rightarrow \pi^+ a$
- Consistent approach to calculation of SN cooling bound
  - Bound on  $C_{\gamma\gamma}$  more conservative, while on  $C_{GG}$  in agreement with literature
- Overall ALP phenomenology depends on the couplings that are generated
  - Improved this approach by including all relevant operators

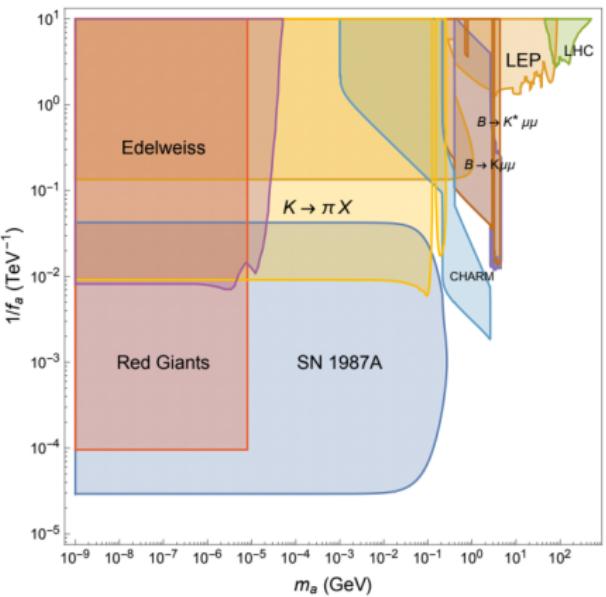
**Thank you for your attention!**

Questions?

# Backup

# Photophobic ALP

- Set  $C_{WW} = -C_{BB}$ 
  - $\rightarrow C_{\gamma\gamma}|_{\text{tree}} = 0$
  - $\rightarrow$  “photophobic ALP”
  - $\rightarrow$  nearly no parameter space open
- Note  $1/f_a \sim C_{WW}$



[N. Craig et al., arXiv: 1805.06538]

# Octet Enhancement

- Hierarchy encountered in Kaon decays

$$\Rightarrow \Gamma(K^0 \rightarrow \pi^+ \pi^-), \Gamma(K^0 \rightarrow \pi^0 \pi^0) \gg \Gamma(K^+ \rightarrow \pi^+ \pi^0)$$

- The very same operators in ChPT lead to an enhanced mixing contribution of  $\eta$  and  $\eta'$

$$i\mathcal{M}(K^+ \rightarrow \pi^+ a) \approx \theta_{a\pi} i\mathcal{M}(K^+ \rightarrow \pi^+ \pi^0) + \theta_{a\eta} i\mathcal{M}(K^+ \rightarrow \pi^+ \eta) + \theta_{a\eta'} i\mathcal{M}(K^+ \rightarrow \pi^+ \eta')$$

[Bardeen et al., Nucl. Phys. B279 (1987)]

# Mixing angles

- The mixing effects can be written as

$$\pi^0 \rightarrow \pi^0 + \theta_{a\pi} a \approx \pi^0 + \epsilon \frac{K_{a\pi} m_a^2}{m_a^2 - m_\pi^2} a$$

$$\eta \rightarrow \eta + \theta_{a\eta} a \approx \eta + \epsilon \frac{K_{a\eta} m_a^2 + m_{a\eta}^2}{m_a^2 - m_\eta^2} a$$

$$\eta' \rightarrow \eta' + \theta_{a\eta'} a \approx \eta' + \epsilon \frac{K_{a\eta'} m_a^2 + m_{a\eta'}^2}{m_a^2 - m_{\eta'}^2} a$$

[D. Aloni et al., arXiv: 1811.03474]

# General Formula for Volume Emission Rate

- The energy loss rate  $Q$  (energy per volume and unit time) is defined by

$$Q = \int \frac{d^3 p_a}{2\omega_a(2\pi)^3} \omega_a \left( \prod_i \int \frac{d^3 p_i}{2\omega_i(2\pi)^3} f_i(\omega_i) \right) \left( \prod_f \int \frac{d^3 p_f}{2\omega_f(2\pi)^3} [1 \pm f_f(\omega_f)] \right)$$
$$\times S \sum_{\text{spins/pol.}} |M|^2 (2\pi)^4 \delta^{(4)} \left( \sum_i p_i^\mu - \sum_f p_f^\mu - p_a^\mu \right)$$

[G. Raffelt, Phys. Rept. 198 (1990)]

# General Formula for Optical Depth

- The optical depth is given by

$$\tau = \beta^{-1} \int dr \Gamma_{\text{abs}},$$

[G. Raffelt, Phys. Rept. 198 (1990)]

where the absorption rate reads

$$\begin{aligned}\Gamma_{\text{abs}} &= \frac{1}{2\omega_a} \left( \prod_i \int \frac{d^3 \mathbf{p}_i}{2\omega_i (2\pi)^3} f_i(\omega_i) \right) \left( \prod_f \int \frac{d^3 \mathbf{p}_f}{2\omega_f (2\pi)^3} [1 \pm f_f(\omega_f)] \right) \\ &\times S \sum_{\text{spins/pol.}} |M|^2 (2\pi)^4 \delta^{(4)} \left( \sum_i p_i^\mu - \sum_f p_f^\mu - p_a^\mu \right)\end{aligned}$$

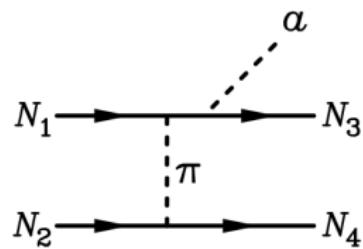
[H. Weldon, Phys. Rev. D28 (1983)]

# Primakoff and Bremsstrahlung processes

## Primakoff

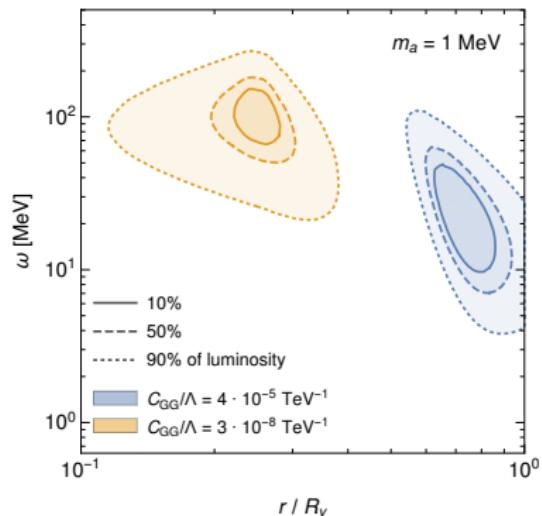
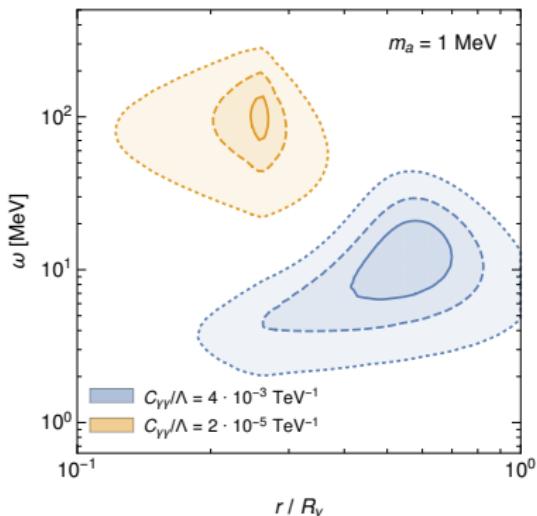


## Bremsstrahlung (OPE)

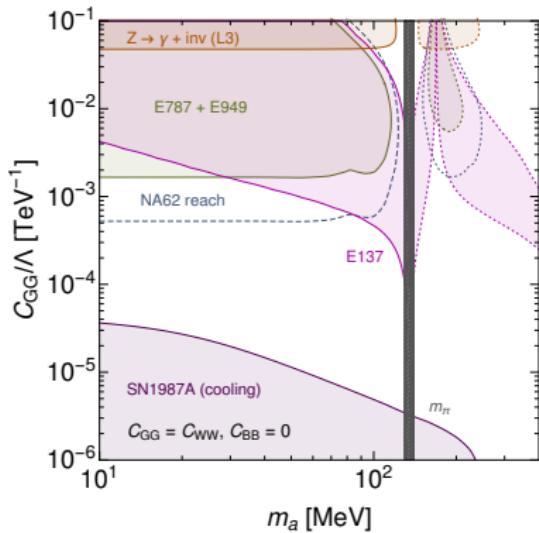
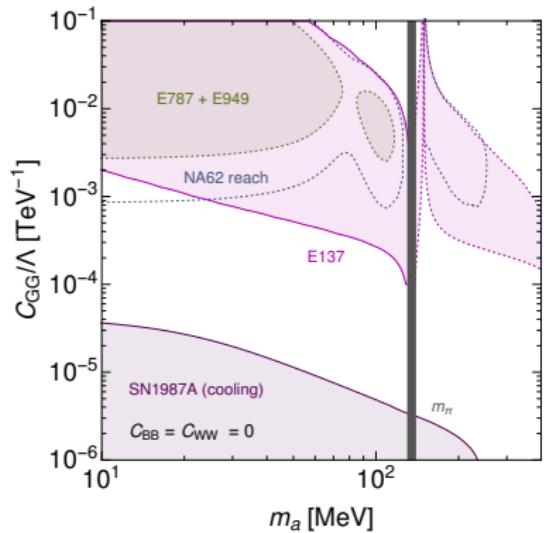


# ALP Luminosity Distribution

What are typical energies for ALPs produced in SN core?



# NA62: Signal Region 2 (1/2)



# NA62: Signal Region 2 (2/2)

