Model independent Higgs boson couplings determination at FCC-ee

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Higgs boson couplings at the HL-LHC



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- Model dependence: results depend on assumptions (e.g. on BSM decays) ٠
- ٠ TeV scale, we need to go (sub-)percent level and model-independent measurements

Typical precision: 1.5-4% (mainly couplings to gauge bosons, 3rd generation fermions, muons), limited by systematic uncertainties

To improve our understanding of the Higgs sector further, and search for O(%) deviations induced by BSM physics at the multi-



Higgs boson production at the FCC-ee



"Clean" (good S/B, no pileup) and abundant Higgs boson production

S ~ 10⁶ ٠

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- **S/B (before selection) ~ 1/100-1000**
- Much larger than pp colliders (S/B as low as 10^{-7} for hadronic signatures such as gg \rightarrow H \rightarrow bb) •



Small wrt pp colliders due to small xsection (but large acceptance!); large wrt linear ee colliders due to larger luminosity





Model independence: the ZH Higgs-strahlung process

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Measuring the Higgs couplings without strong assumptions on how it decays (but assume same kind of interactions as in the SM..) can be done by "tagging" Higgs boson production without observing its decay, detecting only the accompanying Z boson \Rightarrow recoil mass



Model independence: the ZH Higgs-strahlung process

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Without looking at the Higgs decay:

$$\sigma(ee \to ZH) \Rightarrow g^2_{HZZ}$$

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Reconstructing H \rightarrow **ZZ:** $\sigma(ee \rightarrow ZH)BR(H \rightarrow ZZ) \propto \frac{g_{HZZ}^4}{\Gamma} \Rightarrow \Gamma$

Reconstructing other Higgs Boson decays H \rightarrow XX:

 $\sigma(ee \to ZH)BR(H \to XX) \propto \frac{g_{HZZ}^2 g_{HXX}^2}{\Gamma} \Rightarrow g_{HZZ}^2$

Looking at "invisible" Higgs decays (large missing energy)

Measuring the Higgs couplings without strong assumptions on how it decays (but assume same kind of interactions as in the SM..) can be done by "tagging" Higgs boson production without observing its decay, detecting only the accompanying Z boson \Rightarrow recoil mass

$$E_{ff}, \vec{p}_{H} + \vec{p}_{ff}) = (\sqrt{s}, \vec{0}) \Rightarrow m_{recoil}^{2} = (\sqrt{s} - E_{ff})^{2} - p_{ff}^{2}$$

$$= s + m_{Z}^{2} - 2E_{ff}\sqrt{s} \approx m_{H}^{2}$$

$$\overset{\times 10^{3}}{10}$$

$$2^{5} \qquad 2^{0} \qquad 2^{1} \qquad 2^{$$



Improving the precision with the VBF process



In combination with measurements @240 GeV, any $\sigma(vvH)BR(H \rightarrow XX)$ measurement @365 GeV leads to a

a new determination of g_{HWW}:

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 $\frac{\sigma(ee \to \nu\nu H)BR(H \to XX)}{\sigma(ee \to ZH)BR(H \to XX)} \propto \frac{g_{HWW}^2}{g_{HZZ}^2}$

a new determination of g_{HXX}:

 $\frac{\sigma(ee \to \nu\nu H)BR(H \to XX)}{\sigma(ee \to ZH)BR(H \to WW)} \propto \frac{g_{HXX}^2}{g_{HZZ}^2}$

a new determination of the total width:

 $\sigma(ee \to \nu\nu H)BI$

$$R(H \to XX) \propto \frac{g_{HWW}^2 g_{HXX}^2}{\Gamma}$$



Analysis strategy

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Analysis strategy first outlined in public document with studies for LEP3 (ee in LEP/LHC tunnel) [LEP3 note]

- Detector performance in initial studies were estimated using CMS detector simulation •
- Multivariate analysis assumed to reduce bkg by x2 wrt simple selection, ~15-20% improvement in sensitivity ٠

Projections updated with a fast simulation of the performance foreseen for an optimised detector design (CLD) [slides]

- all-silicon tracker for reconstruction of charged particle and interaction vertices ٠
- High-granularity calorimeter for particle-flow reconstruction together with the silicon tracker ٠
- Muon system outside of the magnet coil ٠
- \Rightarrow ~10–40% improvement in sensitivity wrt CMS ٠

2 IPs, 5/ab (2.5/ab/IP) @240 GeV, 1.5/ab @365 GeV

Similar analyses also performed and detailed in CEPC CDR, leading to similar expected performance

In the following slides, I will show a few examples of cases that were studied in some detail, before giving the overall picture of the sensitivity to the different couplings

Numbers in CDR are based on ILC/CLIC expected sensitivities, scaled to account for the larger integrated luminosity of FCCee



Measurement of $\sigma(HZ)$ and of $\sigma(HZ)^*BR(H \rightarrow invisible)$

Total ZH cross section:

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- Opposite-charge, same-flavour lepton pair (e, μ)
- m_{II}~m_Z (after FSR correction)
- Kinematic cuts to reject bkg from $ee \rightarrow II(\gamma)$ (p_T^{II}, p_L^{II}, lepton acoplanarity) and $ee \rightarrow ZZ$ (lepton acollinearity)
- Perform S+B fit to mrecoil distribution
 - 0.5% uncertainty on event rate at FCCee

Invisible branching ratio: ZH, $Z \rightarrow II/bb$, $H \rightarrow invisible$

- $Z \rightarrow \parallel$ as before
- Z→bb: require >=2 jets, recluster original jets until 2-jets are obtained. Require at least 1 b-tagged jet, m_{jj}~m_Z, and same jet kinematic cuts as in the Z→ll case
- Veto any extra particles
- Perform S+B fit to m_{recoil} distribution
 - BR<0.3% at 95% C.L. at FCCee







Measurement of hadronic Higgs boson branching ratios

H→bb

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3 channels: 2I ($Z \rightarrow II$), high E_T^{miss} ($Z \rightarrow vv$), 4j ($Z \rightarrow qq$)

- **2I**: same selection as for $\sigma(HZ)$ + at least one b-tagged jets in recoil.
- **High E_T^{miss}**: similar selection as for BR($H \rightarrow inv$) measurement, require $m_{miss} \sim m_Z$. Rescale jet energies such that m_{miss}=m_Z
- 4j: at least 4-jets, reclustered into 4 iteratively recombining the jet pair with the smallest ٠ invariant mass, and recalibrated (no change in direction) from E,p conservation
 - Reject jets without hadrons or <5 particles (reject Z \rightarrow leptons)
 - Reject events with large missing energy: visible mass > 180 GeV. (reject $Z \rightarrow vv$) ٠
 - Veto events in which m_{j1j2} , $m_{j3j4} \sim m_W$ or m_{j1j2} , $m_{j3j4} \sim m_Z$ (suppress WW, ZZ) •
 - Keep events with $m_{j1j2} \sim m_H$, $m_{j3j4} \sim m_Z$. j1 and j2 must be b-tagged
- Fit recoil or dijet mass distribution ٠

0.3% uncertainty on event rate from combination of the 3 channels at FCCee ٠

H→cc, gg

- ٠ confirmed with dedicated full-simulation study and implementation of real flavour-tagging algorithm)



Similar selection as H→bb, use other working points of b-tagging algorithms (different b-jet efficiency and c, g rejection) or dedicated c, g tagging to obtain system of equations relating measured Higgs yields to BR(H→bb), BR(H→cc), BR(H→gg)

GeV

Events/2

Events/1 GeV

2.2% uncertainty on cc, 1.9% on gg event rate at FCCee (sensitivity very dependent on vertex detector design, should be

Expected coupling uncertainties (k framework)

Collider	HL-LHC	FCC-ee			
Luminosity (ab^{-1})	3	5 @ 240 GeV	+1.5 @ 365 GeV	+ HL-LHC	
Years	25	3	+4	_	FCC CDR volume 1. Physics
$\delta\Gamma_{\rm H}/\Gamma_{\rm H}~(\%)$	SM	2.7	1.3	1.1	opportunities
$\delta g_{\mathrm{HZZ}}/g_{\mathrm{HZZ}}$ (%)	1.5	0.2	0.17	0.16	https://doi.org/10.1140/epjc/
$\delta g_{\rm HWW}/g_{\rm HWW}$ (%)	1.7	1.3	0.43	0.40	<u>s10052-019-6904-3</u>
$\delta g_{ m Hbb}/g_{ m Hbb}$ (%)	3.7	1.3	0.61	0.56	
$\delta g_{ m Hcc}/g_{ m Hcc}$ (%)	SM	1.7	1.21	1.18	
$\delta g_{ m Hgg}/g_{ m Hgg}$ (%)	2.5	1.6	1.01	0.90	
$\delta g_{ m H au au}/g_{ m H au au}$ (%)	1.9	1.4	0.74	0.67	
$\delta g_{\mathrm{H}\mu\mu}/g_{\mathrm{H}\mu\mu}$ (%)	4.3	10.1	9.0	3.8	
$\delta g_{ m H\gamma\gamma}/g_{ m H\gamma\gamma}$ (%)	1.8	4.8	3.9	1.3	
$\delta g_{\rm Htt}/g_{\rm Htt}$ (%)	3.4	_	_	3.1	
BR _{EXO} (%)	SM	< 1.2	< 1.0	< 1.0	

- For **FCC-ee₂₄₀**, global fit to recoil cross section and σ^*BR leads to **absolute** couplings measurements with: ٠
 - **0.2% uncertainty** in coupling to **Z** (>10x better than HL-LHC) ٠
 - **<2% uncertainty** in couplings to **W**, τ , **g**, **b**, **c** (1.3, 1.4, 1.5, 2.8, **O(30)** x **better** than HL-LHC) •
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- **KW** improvement directly related to gww measurement in WBF •
- K_{g} , K_{τ} , K_{b} improvement driven by larger stat as well as reduced uncertainty on Γ - Γ_{W} .
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poor constraints on couplings to **µ** and **Zy** (rare processes, low yield) and **none** on coupling to **top** (ttH not accessible)

FCC-ee₃₆₅: more stat as well as WW \rightarrow H production \Rightarrow smaller uncertainties (up to x3), <% for W, g, b, τ (2.5-6x better than HL-LHC)

Improvement for couplings inducing rare decays (small contribution to Γ), κ_{μ} , κ_{ν} , goes down like $\sqrt{\sigma}$ (benefiting only from larger stat)

Combination with HL-LHC leads to 1.3% uncertainty on κ_γ and 3-4% on κ_μ, κ_{t.} (10-40% better than HL-LHC, no hypothesis on Γ)













Measuring first generation Higgs couplings?

- 1st generation couplings very small \Rightarrow **BR(H** \rightarrow ee/dd/uu) too small to be measured ~O(10-9)
- Can only measure κ_e through study of s-channel Higgs production in ee collisions •
- $\sigma_{ee \rightarrow H}=1.64 \text{ fb}, \Gamma_{H}=4.2 \text{ MeV} \Rightarrow \text{need very high luminosity+very small E spread at } \sqrt{S}=m_{H}$ ٠
- At FCC-ee@125 GeV, L=2e36 cm⁻²s⁻¹ \Rightarrow 20 ab⁻¹/yr \Rightarrow ~30k H/yr if no E spread and ISR ٠
 - **ISR + 4.2 MeV energy spread** $\Rightarrow ~80\%$ reduction (290 ab, ~6k H/yr) ٠

Feasibility study w/ Pythia8: generate $ee \rightarrow H \rightarrow X$ and $ee \rightarrow VV$, $\gamma\gamma$, gg, qq (q=t,b,c,l) (<u>slides</u>)

- 10 final states chosen based on S and S/B: bb, gg, ττ, WW, ZZ, γγ ٠
- Selection based on kinematic quantities characterising single particles, pairs, or global event •
- Simplified assumptions on efficiencies and fake rates for b/c/l/ τ tagging algorithms (and $e \leftrightarrow \gamma$) •

Most significant channel: $H \rightarrow WW^* \rightarrow Ivjj$

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E _{i1.i2} < 52,45 GeV	← Kills e⁺e⁻→qq̄		
$m_{w(lv)}^{1} > 12 \text{ GeV/c}^{2}$	← Kills e⁺e⁻→qq	qā:	σ
E _{lepton} > 10 GeV	← Kills e⁺e⁻→q̄q	ττ:	σ
ME > 20 GeV	← Kills e⁺e⁻→qq	\ \ \\\/*·	G.
m _{ME} < 3 GeV/c ²	← Kills e⁺e⁻→ττ	VVVV .	0.
BDT MVA ← Kills	e⁺e⁻→WW* continuum	H(WW*):	σ
(exploits opposite W [±] µ	oolarizations in H decay)		



 $\Rightarrow \sigma(after) = 4 ab$ = 22 pb $\Rightarrow \sigma(after) = 2.6 ab$ 1 pb = 16.3 fb $\Rightarrow \sigma(after) = 2.7$ fb = 23 ab $\Rightarrow \sigma$ (after) = 8 ab

For L_{int}=10 ab⁻¹ S/√B =80/√27000 ≈ 0.5 Significance ≈ 0.5



Measuring first generation Higgs couplings?

Combining all channels together:

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Current monochromatisation options (*): energy spread = 6 or 10 MeV, L_{int}=3/ab or 7/ab per experiment*year

- **κ**_e < 2.2 @95% CL (1 experiment, 1 year) ٠
- Could exclude $\kappa_e = 0$ @2.5 σ w/ 2 experiments after 1 yr ($\kappa_e = 1 \pm 0.4$) ٠
- Could reach >3 σ on κ_e w/ 2 experiments running 3 yrs at Higgs pole (κ_e =1±0.23) •
- Constraints on κ_e are x100 (x30) better than at HL-LHC (FCC-hh) •
- Few years at $\sqrt{s}=m_{H}$ are not in baseline FCC-ee run plan but would be a very interesting add-on!





Conclusion

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- **generation** (κ_e), setting constraints that are order of magnitudes better than those achievable at hadron machines and can reach sensitivity to SM
- Haven't talked about self-coupling as well as EFT interpretation \Rightarrow in next talk by Roberto •

Clean and abundant production of Higgs boson at FCC-ee, in well-known initial state, with accompanying tagging Z boson, and small backgrounds, will lead to model-independent determination, with (sub-)% level uncertainty, of absolute Higgs boson couplings

Additional running at $\sqrt{s}=125$ GeV could lead to a measurement of the coupling to the 1st

More work is needed in the future on detector optimisation, full simulation and analysis optimisation efforts to further improve the sensitivity and consolidate the results







FCCee - uncertainties on the event rates

are displayed in the last columns

\sqrt{s} (GeV)	240		365	
Luminosity (ab^{-1})	5		1.5	
$\delta(\sigma BR)/\sigma BR$ (%)	HZ	$\nu \overline{\nu} H$	HZ	$\nu\overline{\nu}H$
$H \rightarrow any$	± 0.5		± 0.9	
$H \to b \bar{b}$	± 0.3	± 3.1	± 0.5	± 0.9
$H \rightarrow c \bar{c}$	± 2.2		± 6.5	± 10
$H \rightarrow gg$	± 1.9		± 3.5	± 4.5
$\rm H ightarrow W^+W^-$	± 1.2		± 2.6	± 3.0
$H \rightarrow ZZ$	± 4.4		± 12	± 10
$H \to \tau \tau$	± 0.9		± 1.8	± 8
$H \rightarrow \gamma \gamma$	± 9.0		± 18	± 22
$H \to \mu^+ \mu^-$	± 19		± 40	
$H \rightarrow invis.$	< 0.3		< 0.6	

Table 4.1 Relative statistical uncertainty on the measurements of event rates, providing $\sigma_{HZ} \times BR(H \rightarrow XX)$ and $\sigma_{\nu\bar{\nu}H} \times BR(H \rightarrow XX)$, as expected from the FCC-ee data. This is obtained from a fast simulation of the CLD detector and consolidated with extrapolations from full simulations of similar linear-collider detectors (SiD and CLIC). All numbers indicate 68% C.L. intervals, except for the 95% C.L. sensitivity in the last line. The accuracies expected with 5 ab⁻¹ at 240 GeV are given in the middle columns, and those expected with 1.5 ab⁻¹ at $\sqrt{s} = 365$ GeV



Global EFT fit



Figure S.1: One- σ precision reach at the FCC on the effective single Higgs couplings, Higgs selfcoupling, and anomalous triple gauge couplings in the EFT framework. Absolute precision in the EW Compared with LHC and LEP, FCC-ee/eh will improve the measurements of EFT parameters by roughly one order of magnitude. A combination with the LHC measurements provides a marginal improvement for most of the parameters. For $g_{H\gamma\gamma}^{eff}$, $g_{HZ\gamma}^{eff}$ and $g_{H\mu\mu}$, the improvements are more significant, as the small rates and clean signals make the LHC reaches comparable to that of lepton colliders. Other couplings, e.g. geff_{HVV} and g_{Hbb}, are also indirectly improved in the combination. It should be noted that the measurements of the H \rightarrow gg branching fraction only constrain a linear combination of g^{eff}_{Hgg} and g_{Htt}. These two couplings are thus only constrained independently by lepton colliders when ttH production is measured. Therefore, combination with LHC measurements is required for the FCC-ee to constrain these couplings independently. The resulting bound on gHtt is then even substantially better than that set by the LHC alone.

g_{Hff} : modified Yukawa coupling

$$(g_{HXX}^{\text{eff}})^2 \equiv \Gamma_{H \to X} / \Gamma_{H \to X}^{\text{SM}}$$



FCC-ee Higgs Couplings

Unique measurements at highest precision

Collider	HL-LHC	ILC ₂₅₀	CLIC ₃₈₀	FCC-ee			FCC-eh
Luminosity (ab ⁻¹)	3	2	0.5	5 @ 240 GeV	+1.5 @ 365 GeV	+ HL-LHC	2
Years	25	15	8	3	+4	-	20
$\delta\Gamma_{\rm H}/\Gamma_{\rm H}~(\%)$	SM	3.6	4.7	2.7	1.3	1.1	SM
$\delta g_{\mathrm{HZZ}}/g_{\mathrm{HZZ}}$ (%)	1.5	0.30	0.60	0.2	0.17	0.16	0.43
$\delta g_{\rm HWW}/g_{\rm HWW}$ (%)	1.7	1.7	1.0	1.3	0.43	0.40	0.26
$\delta g_{ m Hbb}/g_{ m Hbb}$ (%)	3.7	1.7	2.1	1.3	0.61	0.56	0.74
$\delta g_{ m Hcc}/g_{ m Hcc}$ (%)	SM	2.3	4.4	1.7	1.21	1.18	1.35
$\delta g_{ m Hgg}/g_{ m Hgg}$ (%)	2.5	2.2	2.6	1.6	1.01	0.90	1.17
$\delta g_{\rm H\tau\tau}/g_{\rm H\tau\tau}$ (%)	1.9	1.9	3.1	1.4	0.74	0.67	1.10
$\delta g_{ m H\mu\mu}/g_{ m H\mu\mu}$ (%)	4.3	14.1	n.a.	10.1	9.0	3.8	n.a.
$\delta g_{ m H\gamma\gamma}/g_{ m H\gamma\gamma}$ (%)	1.8	6.4	n.a.	4.8	3.9	1.3	2.3
$\delta g_{\rm Htt}/g_{\rm Htt}$ (%)	3.4	_	-	-	-	3.1	1.7
BR _{EXO} (%)	SM	< 1.8	< 3.0	< 1.2	< 1.0	< 1.0	n.a.

Uncertainties not limited by experimental or theoretical uncertainties. Statistics sets the floor.

Indirect sensitivity to Higgs self-coupling



Coupling sensitivities (k framework) including FCC-hh

• photons, to Z+photon, and to muons) or set more stringent limits on the invisible BR

kappa-0	FCC-ee	FCC-ee/eh/hh						1 0 .			
	240 365			kappa-2 scenario	FCC-ee ₂₄₀	FCC-ee ₃₆₅	FCC-ee/eh/hh	kappa-3 scenario	FCC-ee ₂₄₀	FCC-ee ₃₆₅	FCC-ee/eh/hh
<i>к</i> _W [%]	1.3 0.43	0.14		к _W [%]	1.3	0.44	0.2	к _W [%]	0.88	0.41	0.19
κ _Z [%]	0.20 0.17	0.12		κ _Z [%]	0.21	0.18	0.17	$\kappa_Z[\%]$	0.20	0.17	0.16
к g [%]	1.7 1.0	0.49		$\kappa_g \ [\%]$	1.7	1.0	0.52	$\kappa_{g}[\%]$	1.2	0.9	0.5
κ _γ [%]	4.7 3.9	0.29		κ _γ [%]	4.8	3.9	0.32	κ _γ [%]	1.3	1.3	0.31
$\kappa_{Z\gamma}$ [%]	81* 75*	0.69		$\kappa_{Z\gamma}$ [%]	71.*	66.*	0.71	$\kappa_{Z\gamma}$ [%]	10.*	10.*	0.7
κ_{c} [%]	1.8 1.3	0.95		κ_c [%]	1.8	1.3	0.96	$\kappa_c [\%]$	1.5	1.3	0.96
<i>κ</i> _t [%]	— —	1.0		κ_t [%]	-	-	1.0	$\kappa_t ~[\%]$	3.1	3.1	0.96
к _b [%]	1.3 0.67	0.43		к _b [%]	1.3	0.69	0.48	к _b [%]	1.	0.64	0.48
κ _μ [%]	10 8.9	0.41		κ _μ [%]	10.	8.9	0.43	κ_{μ} [%]	4.	3.9	0.43
$\kappa_{ au}$ [%]	1.4 0.73	0.44		$\kappa_{ au}$ [%]	1.4	0.74	0.49	κ _τ [%]	0.94	0.66	0.46
			΄ (BR _{inv} (<%, 95% CL)	0.22	0.19	0.024	BR _{inv} (<%, 95% CL)	0.22	0.19	0.024
				BR _{unt} (<%, 95% CL)	1.2	1.1	1.0	BR _{unt} (<%, 95% CL)	1.2	1.	1.

With respect to FCC-ee, FCC-hh has high-enough energy to produce ttH (and thus the coupling to the top) as well as much larger cross sections, which are particularly beneficial for measuring the couplings that induce rare decays (couplings to



Peak luminosity vs sqrt(s)





