

Gravitational Wave Follow-up

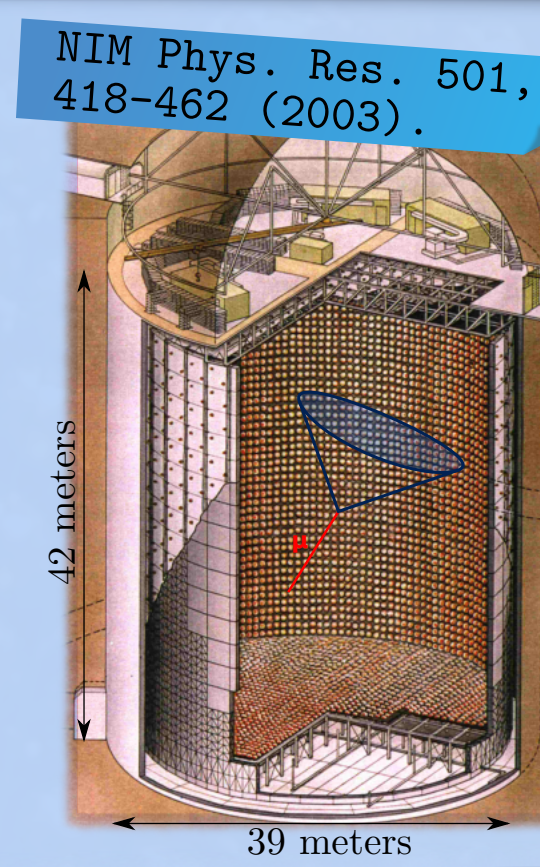
- Still no evidence of common astrophysical source of Gravitational Waves (GW) and neutrinos (IceCube-170922A).
- Such detection would also significantly improve the localization of the source, making EM follow-up observations easier.
- Advanced LIGO and Virgo are sending realtime public alerts since April 2019 (O3 run) with a rate of 1/week.
- They have recently published GWTC-2 catalog covering first half of O3.

Super-Kamiokande

- 50 kton water Cherenkov detector
- Located in Mozumi mine, Gifu-ken, Japan
- Sensitive to MeV-TeV neutrinos
- Running since 1996, with six different periods
- In Summer 2020, Gadolinium was dissolved into the water (allow better neutron tagging)

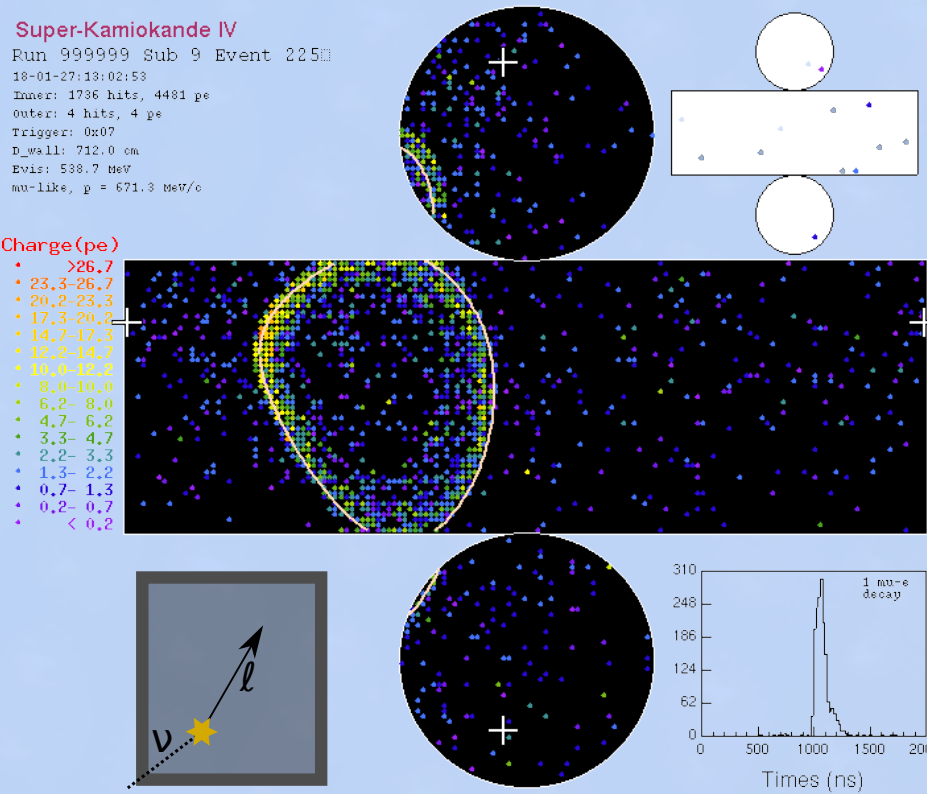
Different phases

Start	End
	SK-I
04/1996	07/2001
	SK-II
10/2002	10/2005
	SK-III
07/2006	09/2008
	SK-IV
09/2008	06/2018
	SK-V
01/2019	07/2020
	SK-VI
07/2020	ongoing



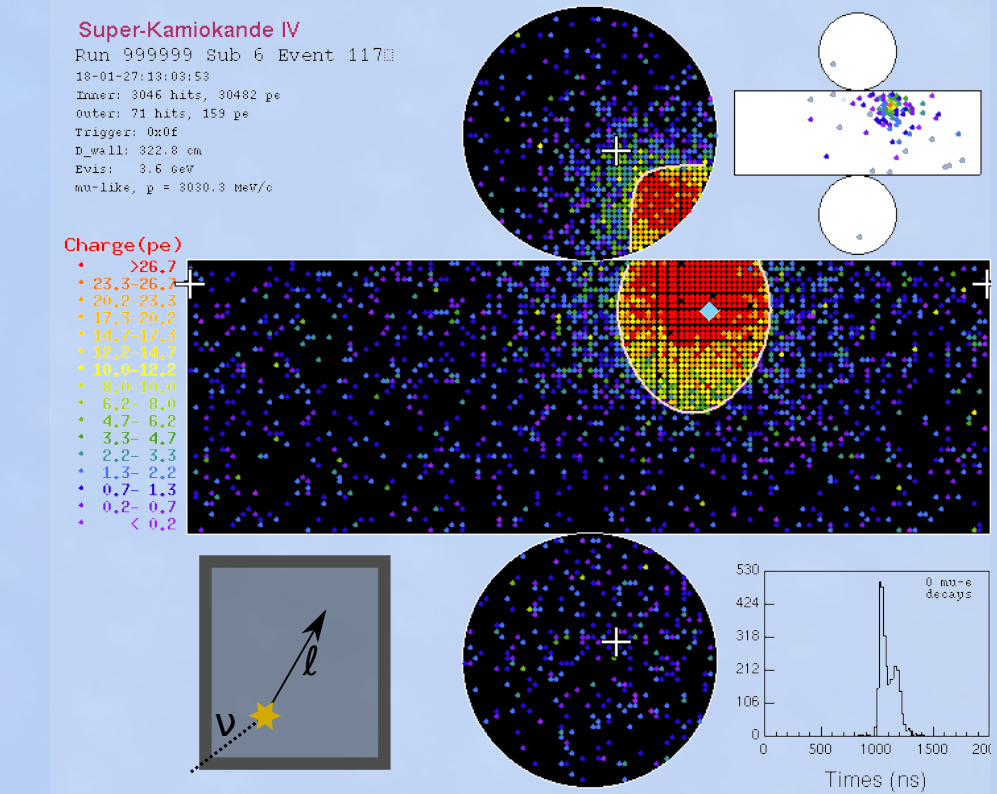
Samples

Fully-Contained (FC)



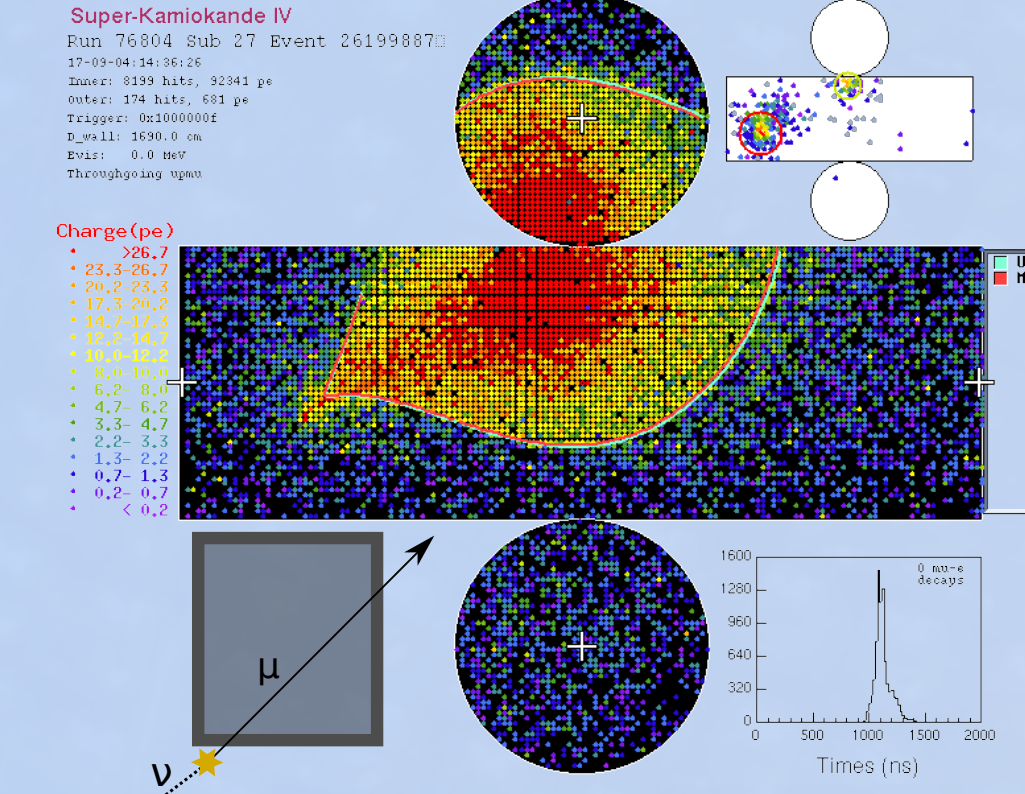
FULL-SKY COVERAGE
Energy range: 0.1-10 GeV
Background in 1000 s: 0.112

Partially-Contained (PC)

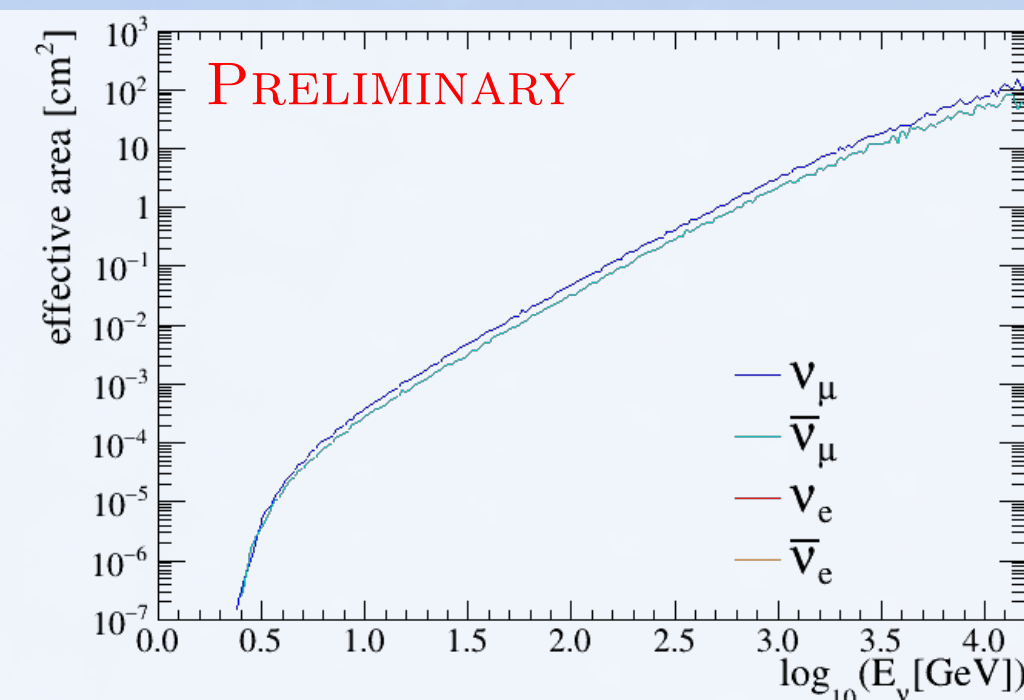
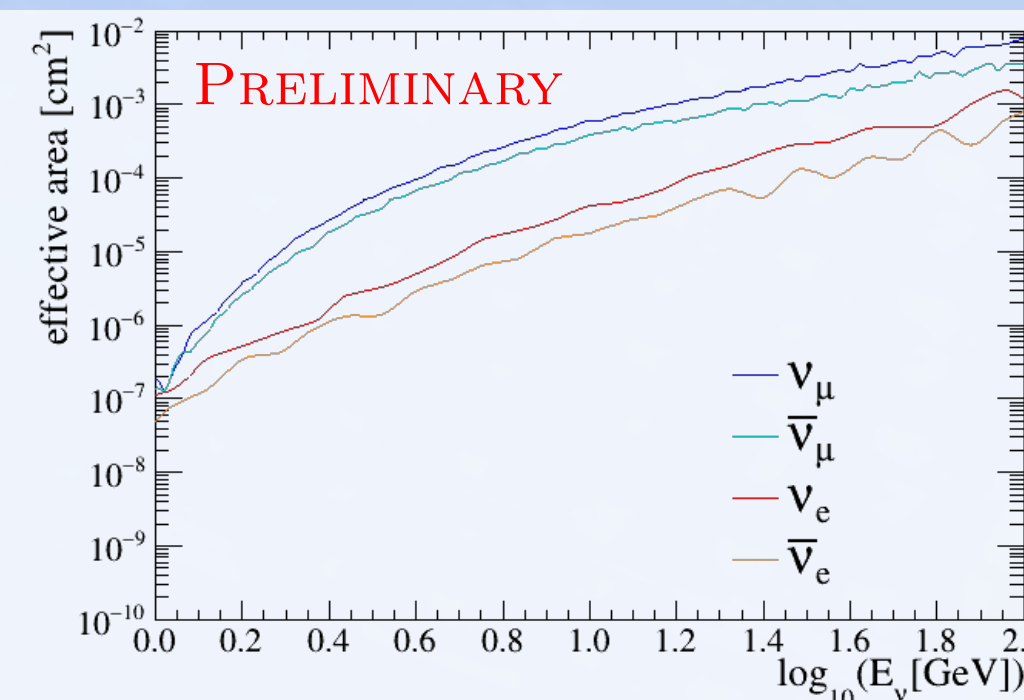
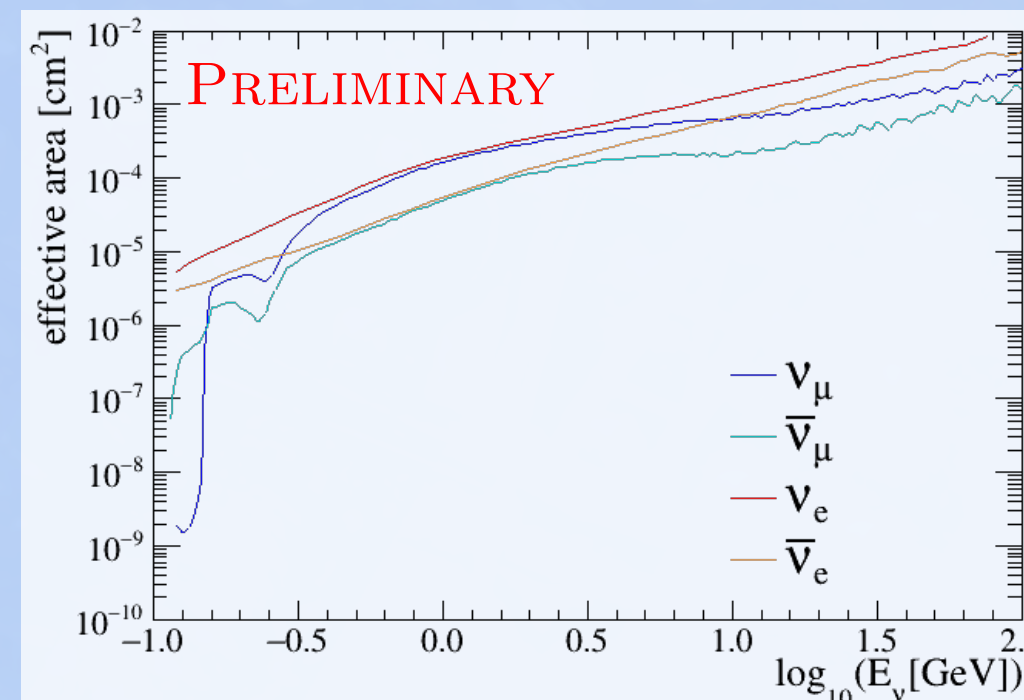


FULL-SKY COVERAGE
Energy range: 0.1-100 GeV
Background in 1000 s: 0.007

Upward-going μ (UPMU)

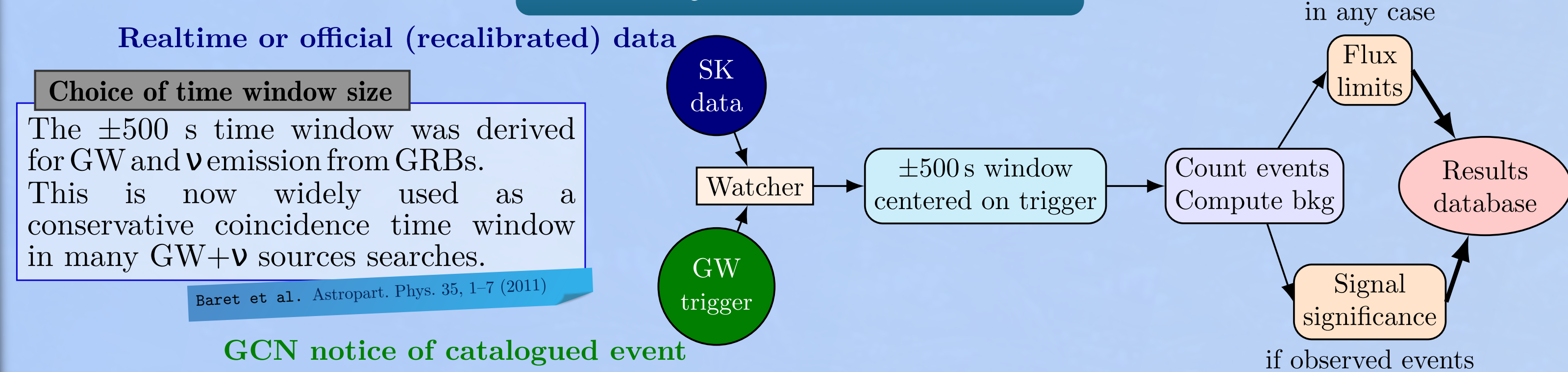


ONLY BELOW THE HORIZON
Energy range: few GeV - TeV
Background in 1000 s: 0.016



For each sample:
$$N_{\text{signal}} = \int A_{\text{eff}}(E, \Omega) \times \frac{dn(E)}{dE} \times dE \quad \& \quad N_{\text{background}} = T \times \text{rate}_{\text{atmospheric background rate}}$$

Analysis method



Sample-by-sample statistics

The neutrino flux from a point-source is:

$$\frac{dn}{dE_\nu} = \phi_0 E_\nu^{-2} \quad \text{assuming } E^{-2} \text{ spectrum}$$

For a given sample s (FC, PC or UPMU):

- expected background B_{exp} is known
- observed number of event N_{obs} is known
- compute Poisson upper limit on signal

If source position Ω is perfectly known:

$$\phi_{0,\text{up}}^{s,f}(\Omega) = \frac{N_{90}}{\int_{E_{\text{min}}}^{E_{\text{max}}} dE_\nu A_{\text{eff}}^{s,f}(E_\nu, \Omega) E_\nu^{-2}} c(\Omega)$$

To take into account wide localisation, we define:

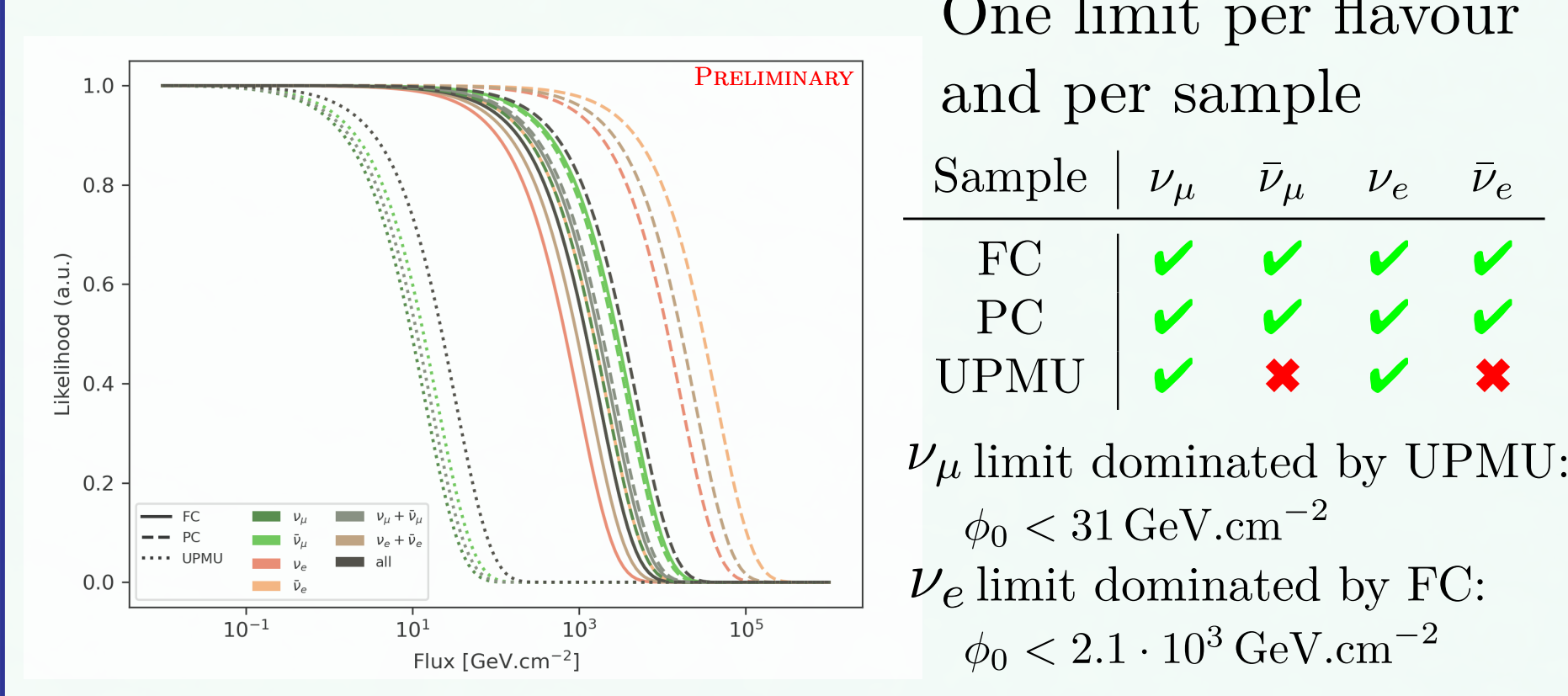
$$\mathcal{L}(\phi_0) = \int \frac{(c(\Omega)\phi_0 + N_{\text{exp}})^{N_{\text{obs}}}}{N_{\text{obs}}!} e^{-c(\Omega)\phi_0 + N_{\text{exp}}} \mathcal{P}_{\text{GW}}(\Omega) d\Omega$$

GW probability distribution

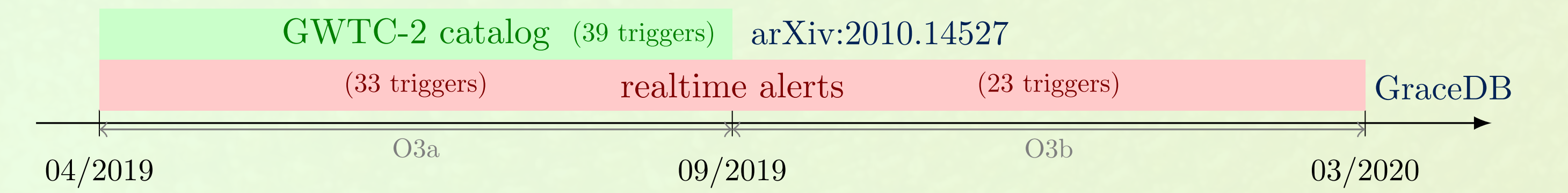
and find Bayesian upper limit (flat prior on flux):

$$\int_0^{\phi_{0,\text{up}}} \mathcal{L}(\phi_0) d\phi_0 = 0.90$$

Example of S190412m (realtime alert)



O3 results

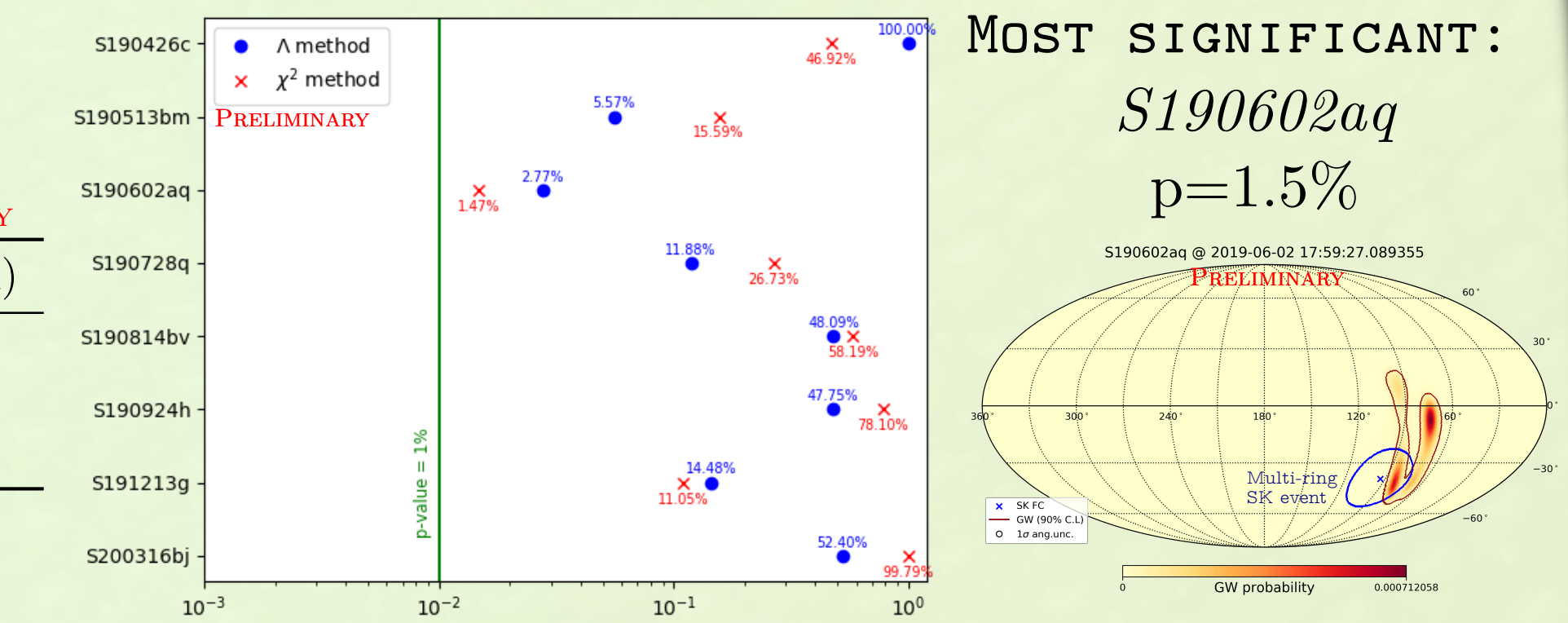


Results of SK GW follow-up of O3 realtime alerts presented at Neutrino 2020

- 46 GWs with SK running
- 8 coincident neutrinos

Sample	Total observed	Total expected
FC	6	5.06
PC	0	0.33
UPMU	2	0.74

NO SIGNIFICANT OBSERVATION



Ongoing analysis of GWTC-2 events (publication soon...)

- Catalog events: more precise sky localisation, distance estimation and properties (masses...) of the source objects.
- This could be used to constrain precisely neutrino emission from the source.
- Distance uncertainties and correlations covered by using this likelihood:

Veske et al. arXiv:2001.00566

$$\mathcal{L}(E_{\text{iso}}; TS_m, \mathcal{V}_{\text{GW}}) = \int \sum_{k=0}^2 \frac{(c'(r, \Omega) E_{\text{iso}})^k}{k!} e^{-c'(r, \Omega) E_{\text{iso}}} \times P_k(TS_m) \mathcal{V}_{\text{GW}}(r, \Omega) r^2 dr d\Omega$$

conversion from E_{iso} to number of events
GW 3D localisation

- It may also be used to combine different GW of the same type to get a stronger constrain:

$$\mathcal{L}(E_{\text{iso}}) = \prod_i \mathcal{L}^{(i)}(E_{\text{iso}}) \quad \text{or} \quad \mathcal{L}(k_\nu) = \prod_i \mathcal{L}^{(i)}(k_\nu C_i)$$

(e.g. total mass)
assuming all objects have exact same emission
assuming emission \propto source characteristic C_i

Combined analysis

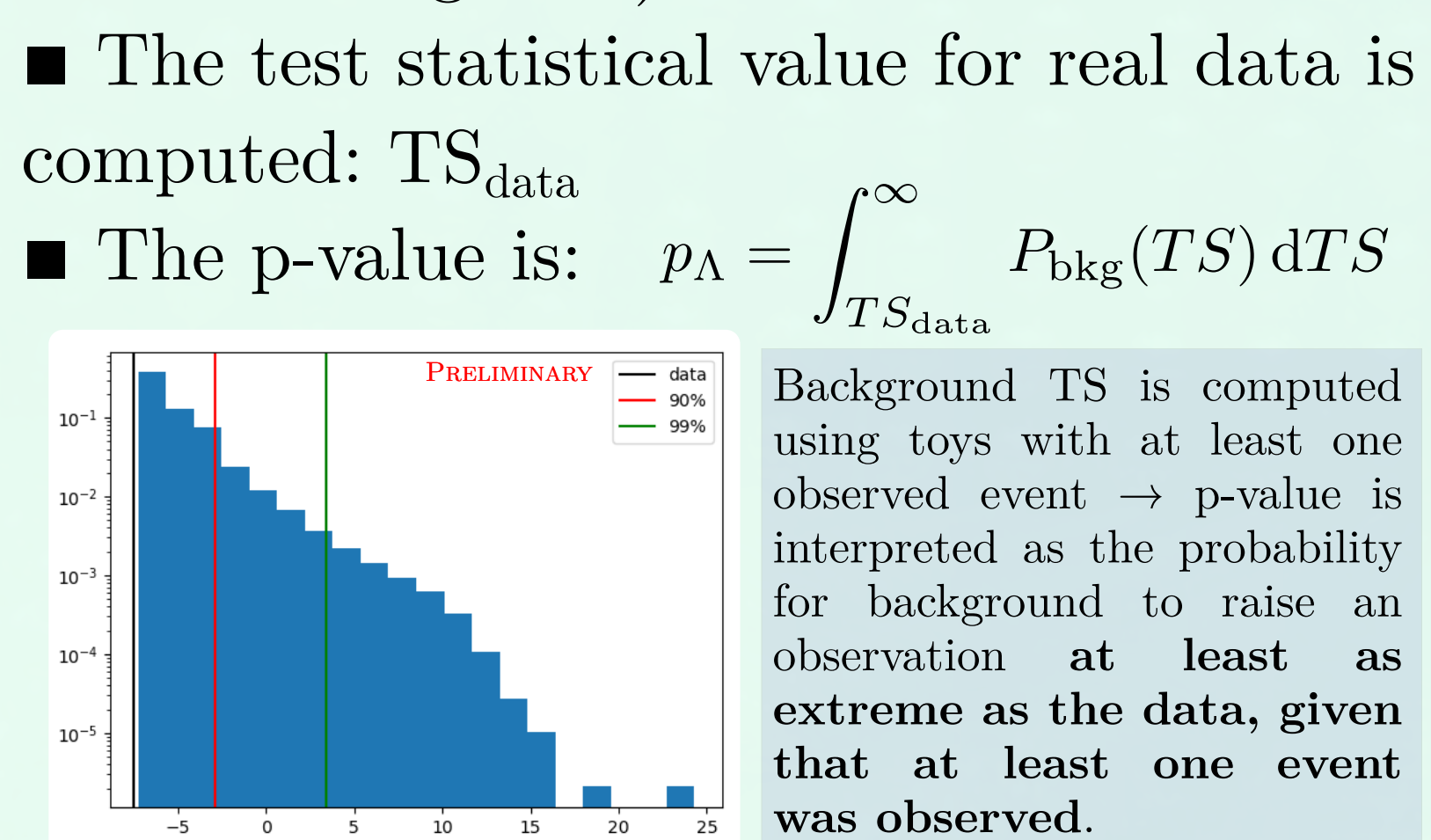
- Define a likelihood and test statistic to separate between signal (point-source) and background:

$$\mathcal{L}^{(s)}(n_S^{(s)}; \gamma; \vec{x}_S) = \frac{e^{-(n_S^{(s)} + n_B^{(s)})} (n_S^{(s)} + n_B^{(s)})^{N(s)}}{N(s)!} \prod_{i=1}^{N(s)} \frac{n_S^{(s)} S^{(s)}(\vec{x}_i, E_i; \vec{x}_S, \gamma) + n_B^{(s)} B^{(s)}(\vec{x}_i, E_i)}{n_S^{(s)} + n_B^{(s)}} \quad \& \quad TS = \max_{\Omega} [\Lambda(\Omega)] \quad \text{with} \quad \Lambda(\Omega) = 2 \sum_s \ln \left[\frac{\mathcal{L}^{(s)}(n_S^{(s)}; \gamma; \vec{x}_S; \Omega)}{\mathcal{L}^{(s)}(n_S^{(s)}; \gamma; \vec{x}_S; 0; \Omega)} \right] + 2 \ln \mathcal{P}_{\text{GW}}(\Omega)$$

values maximising the likelihood
GW probability distribution

Compute signal significance

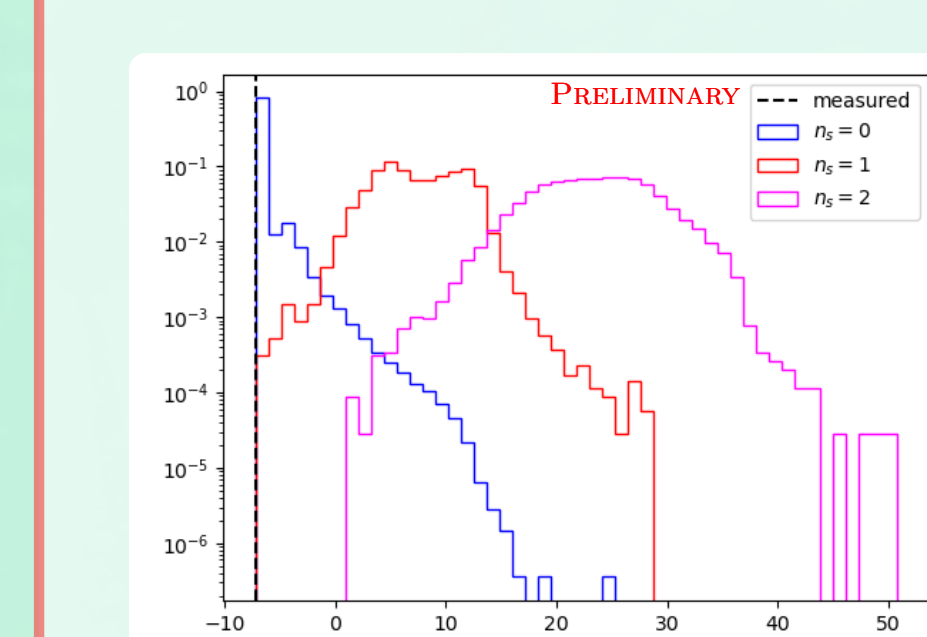
- The background distribution of TS is obtained using time-scrambled neutrinos and expected background rate (FC + PC + UPMU together).
- The test statistical value for real data is computed: TS_{data}
- The p-value is: $p_\Lambda = \int_{TS_{\text{data}}}^{\infty} P_{\text{bkg}}(TS) dTS$



Compute flux limits

Inspired from arXiv:2001.00566

- Compute TS distribution for signal injections (consistent with GW localisation), with $n_{\text{sig}}=0$ (background-only), $n_{\text{sig}}=1$, $n_{\text{sig}}=2$.
- Compute $P_0(TS_{\text{data}})$, $P_1(TS_{\text{data}})$ and $P_2(TS_{\text{data}})$, the associated "data" probabilities
- Compute likelihood: $\mathcal{L}(\phi_0; TS_m, \mathcal{P}_{\text{GW}}) = \int \sum_{k=0}^2 \frac{(c(\Omega)\phi_0)^k}{k!} e^{-c(\Omega)\phi_0} \times P_k(TS_m) \times \mathcal{P}_{\text{GW}}(\Omega) d\Omega$
- Find Bayesian upper limit as in sample-by-sample approach.



> This gives the most constraining results for each flavour individually + one can combine flavours (e.g. $\nu_\mu + \bar{\nu}_\mu$)

Combining flavours require an assumption on distribution of flavours. Used scenario is:

$$\nu_e : \nu_\mu : \nu_\tau = (1 : 2 : 0)_{\text{source}} \rightarrow (1 : 1 : 1)_{\text{Earth}}, \quad \phi_\nu = \phi_{\bar{\nu}}$$