Flavor composition of high-energy astrophysical neutrinos: *Why you should be excited* 

Mauricio Bustamante

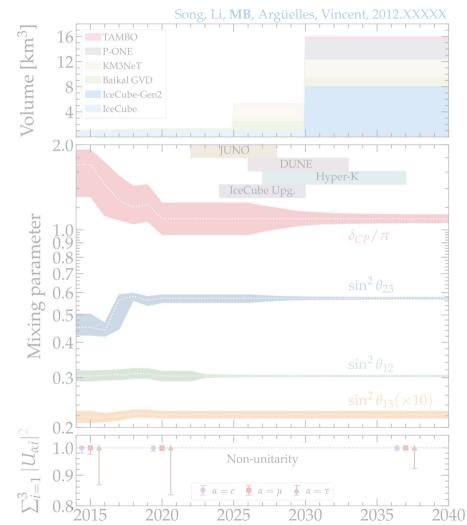
Niels Bohr Institute, University of Copenhagen

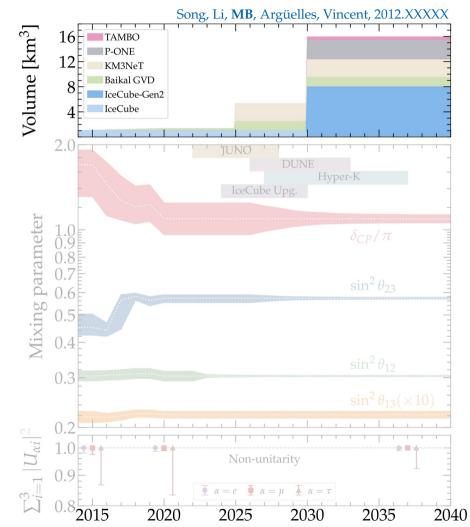
UNIVERSITY OF COPENHAGEN





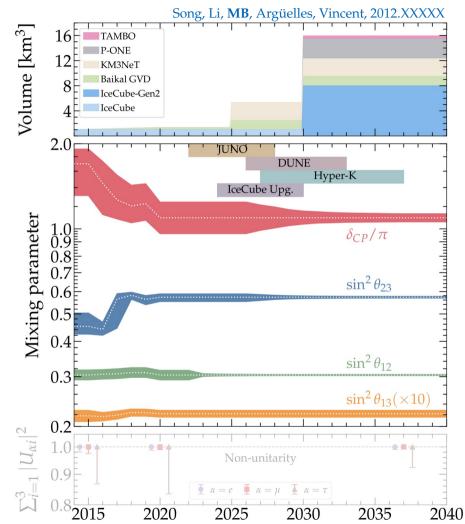






#### *Flavor measurements:*

New neutrino telescopes = more events, better flavor measurement

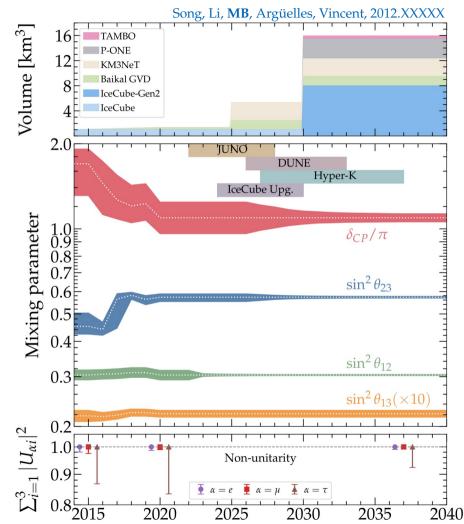


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#### Oscillation physics:

We will know the mixing parameters better (JUNO, DUNE, Hyper-K, IceCube Upgrade)



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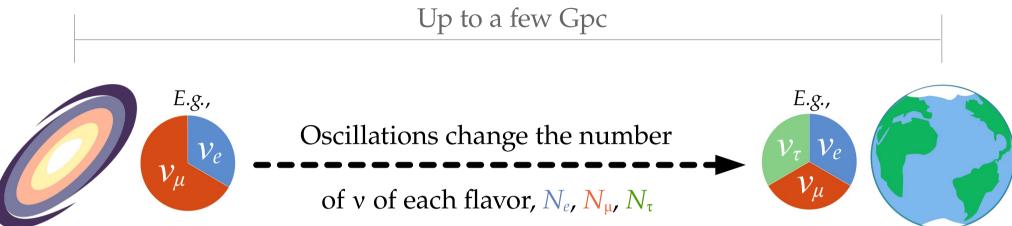
#### Oscillation physics:

We will know the mixing parameters better (JUNO, DUNE, Hyper-K, IceCube Upgrade)

*Test of the oscillation framework:* We will be able to do what we want even if oscillations are non-unitary

#### Astrophysical sources

#### Earth



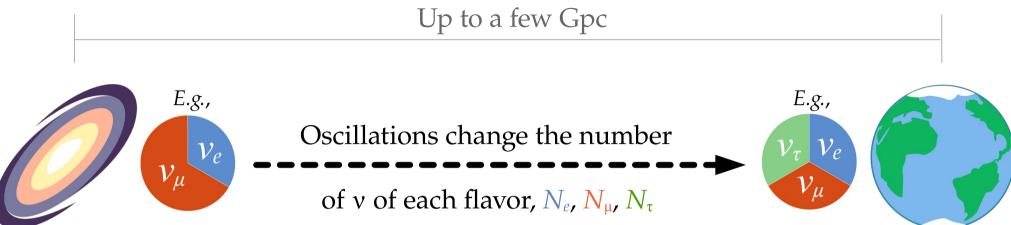
## Different production mechanisms yield different flavor ratios: $(f_{e,S}, f_{\mu,S}, f_{\tau,S}) \equiv (N_{e,S}, N_{\mu,S}, N_{\tau,S})/N_{tot}$

Flavor ratios at Earth ( $\alpha = e, \mu, \tau$ ):

$$f_{\alpha,\oplus} = \sum_{\beta=e,\mu,\tau} P_{\nu_{\beta}\to\nu_{\alpha}} f_{\beta,S}$$

#### Astrophysical sources

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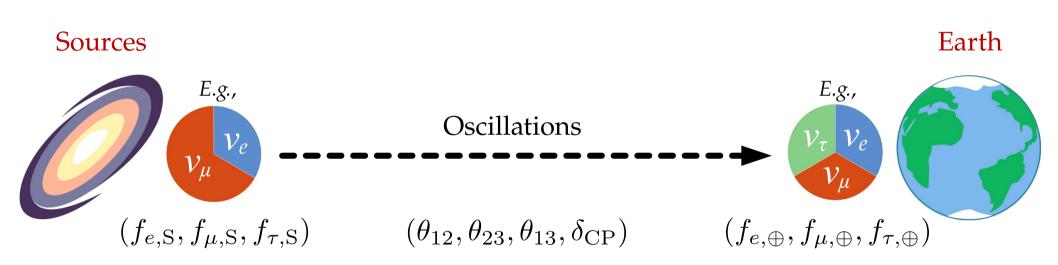


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Flavor ratios at Earth (
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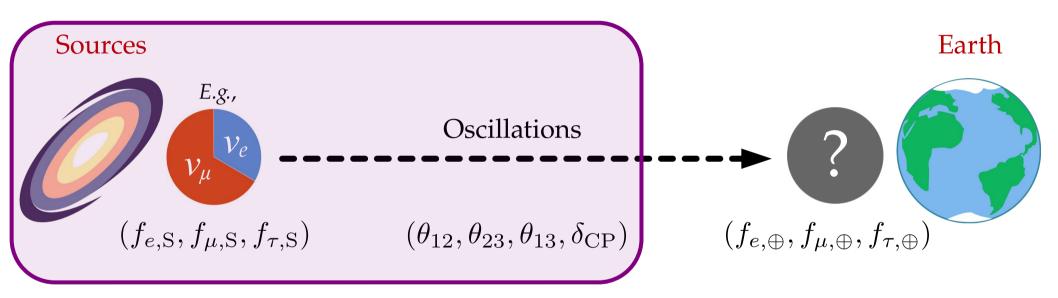
$$f_{\alpha, \oplus} = \sum_{\beta = e, \mu, \tau} P_{\nu_{\beta} \to \nu_{\alpha}} f_{\beta, S}$$
Standard oscillations  
or  
new physics

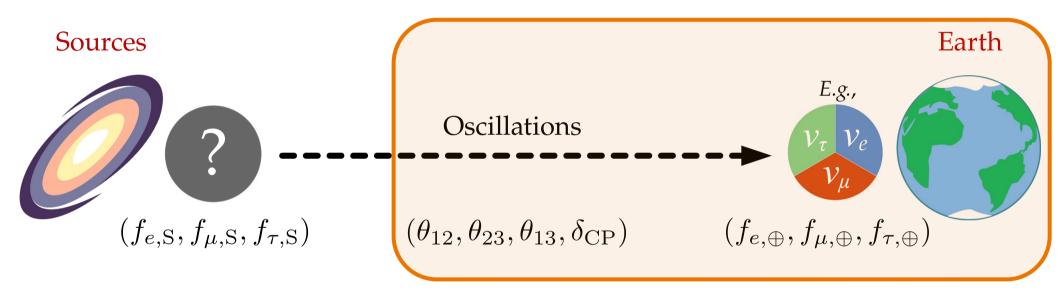
#### *From sources to Earth:* we learn what to expect when measuring $f_{\alpha,\oplus}$



*From Earth to sources:* we let the data teach us about  $f_{\alpha,S}$ 

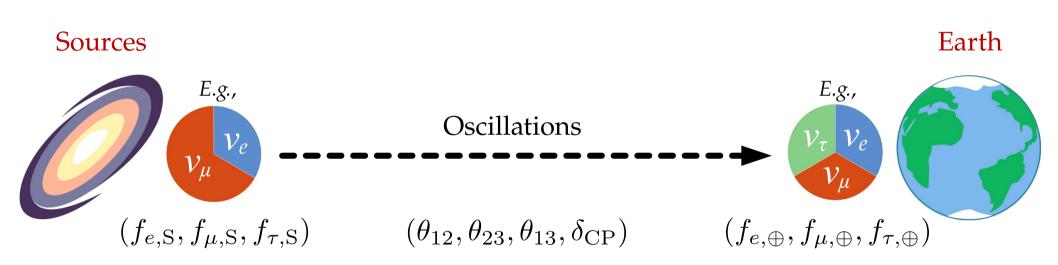
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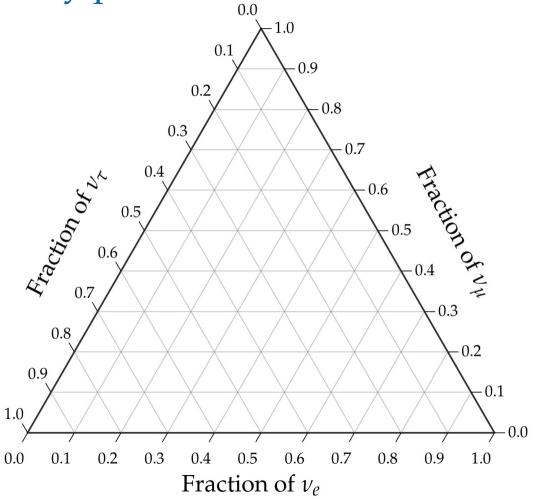


*From Earth to sources:* we let the data teach us about  $f_{\alpha,S}$ 

Assumes underlying unitarity – sum of projections on each axis is 1

How to read it: Follow the tilt of the tick marks

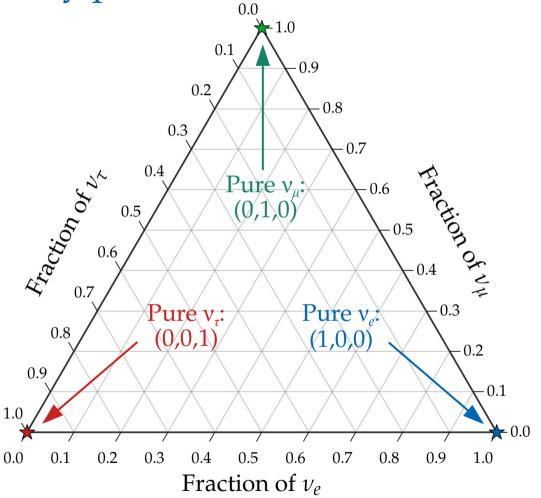
Always in this order:  $(f_e, f_\mu, f_\tau)$ 



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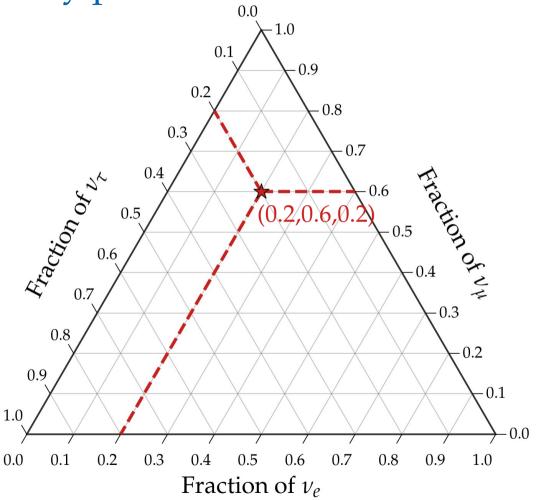
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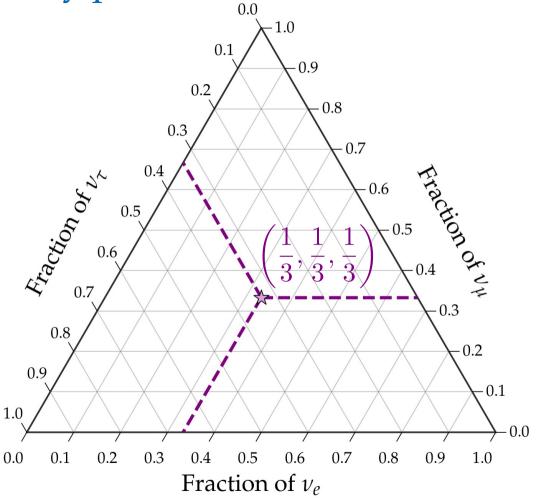
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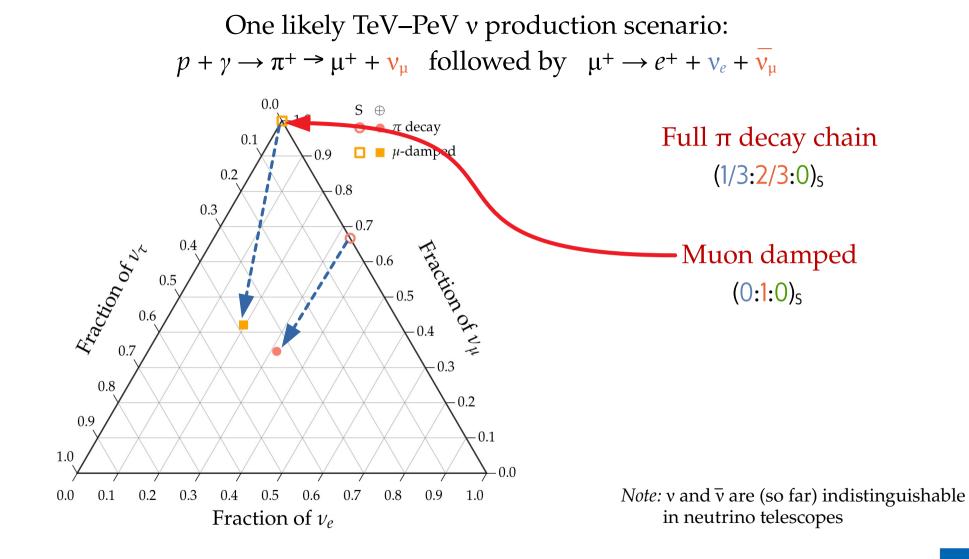
One likely TeV–PeV v production scenario:  $p + \gamma \rightarrow \pi^+ \rightarrow \mu^+ + \nu_{\mu}$  followed by  $\mu^+ \rightarrow e^+ + \nu_e + \overline{\nu_{\mu}}$ 

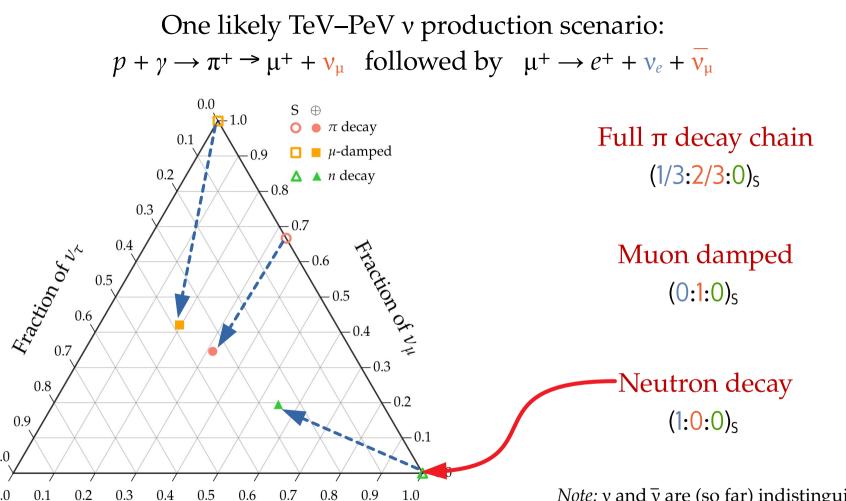
# Full $\pi$ decay chain (1/3:2/3:0)<sub>s</sub>

*Note:* v and  $\overline{v}$  are (so far) indistinguishable in neutrino telescopes

One likely TeV–PeV v production scenario:  $p + \gamma \rightarrow \pi^+ \rightarrow \mu^+ + \nu_{\mu}$  followed by  $\mu^+ \rightarrow e^+ + \nu_e + \overline{\nu_{\mu}}$ 0.0 S O -1.0  $\pi$  decay Full  $\pi$  decay chain 0.1-0.9  $(1/3:2/3:0)_{S}$ 0.2 - 0.8 0.3 -0.7 Fraction of Vr Fraction of NH 0.4- 0.6 0.5 - 0.5 0.6 -0.30.8 -0.2 0.9 -0.1 1.0 -0.0 *Note:* v and  $\overline{v}$  are (so far) indistinguishable 0.0 0.2 0.6 0.7 0.8 0.9 1.0 0.1 0.3 0.40.5 in neutrino telescopes Fraction of  $v_e$ 

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1.0

0.0

0.1

0.2

0.3

0.4

0.6

Fraction of  $\nu_{e}$ 

0.7

0.8

0.9

1.0

*Note:* v and  $\overline{v}$  are (so far) indistinguishable in neutrino telescopes

How does IceCube see TeV–PeV neutrinos?

#### Deep inelastic neutrino-nucleon scattering

Neutral current (NC)Charged current (CC)

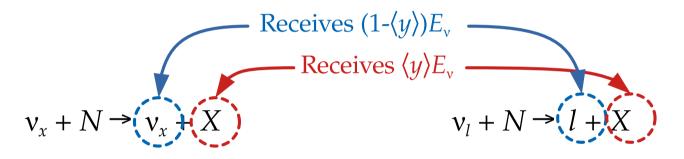
$$v_x + N \rightarrow v_x + X$$

 $v_l + N \rightarrow l + X$ 

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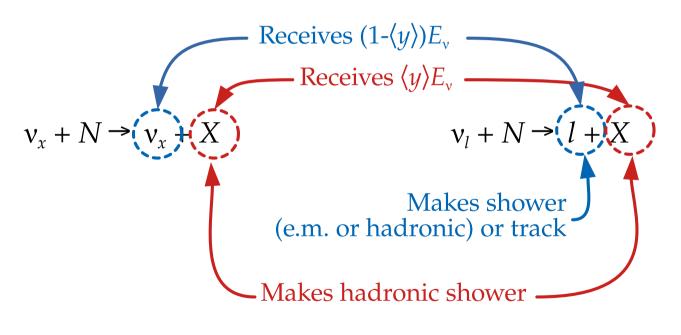


At TeV–PeV, the average inelasticity  $\langle y \rangle = 0.25-0.30$ 

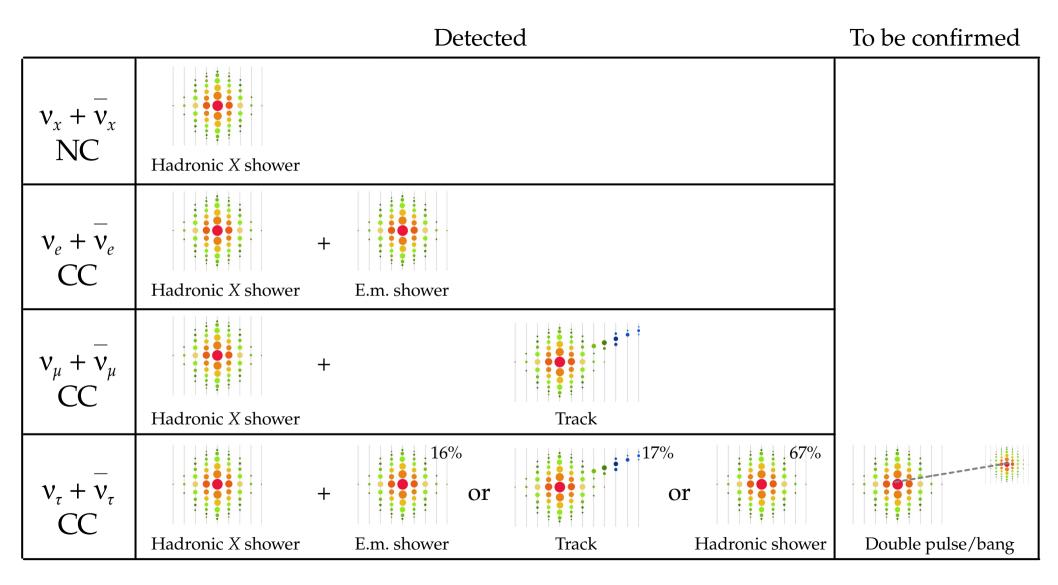
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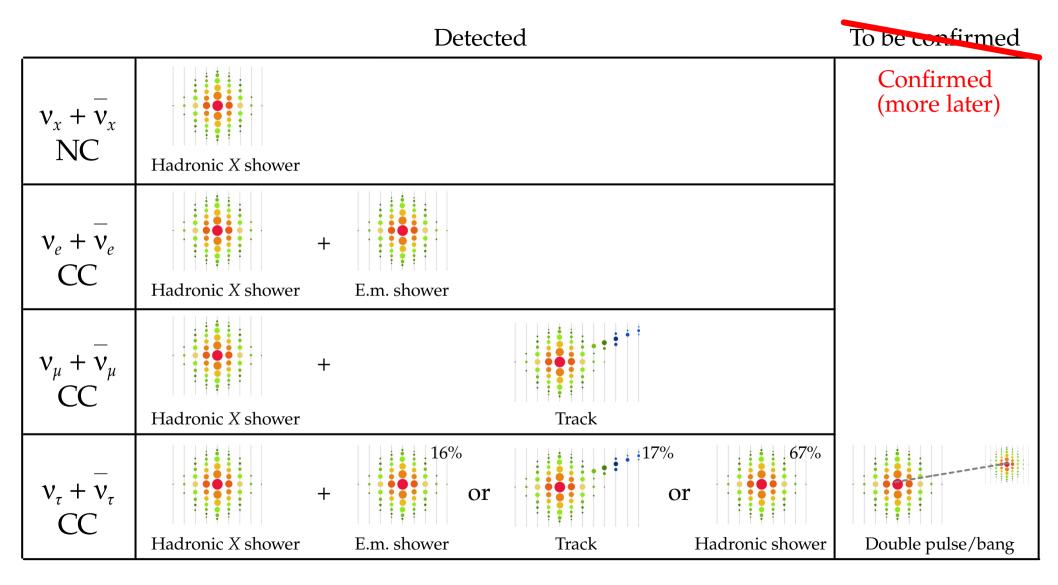
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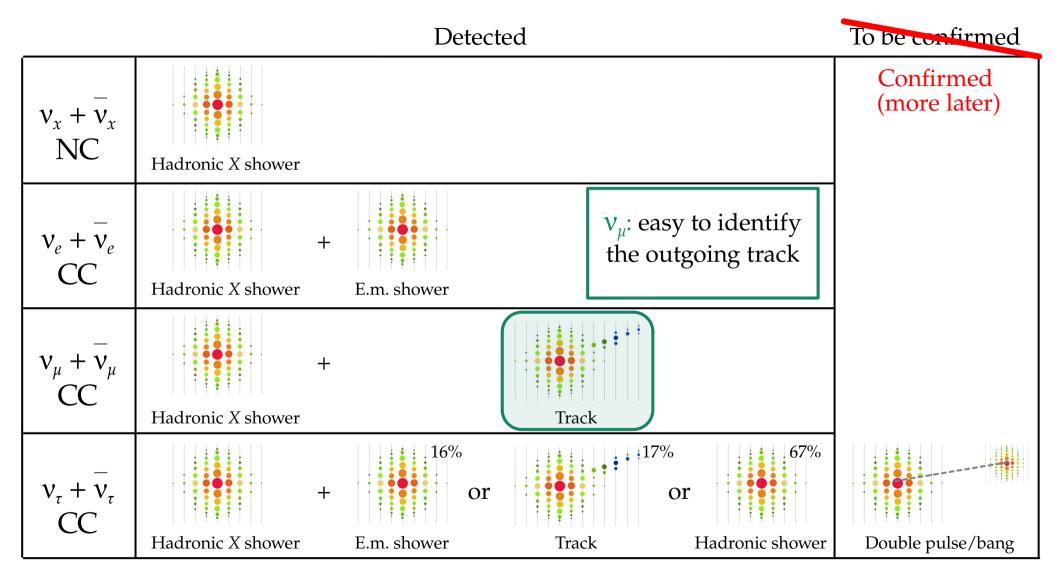
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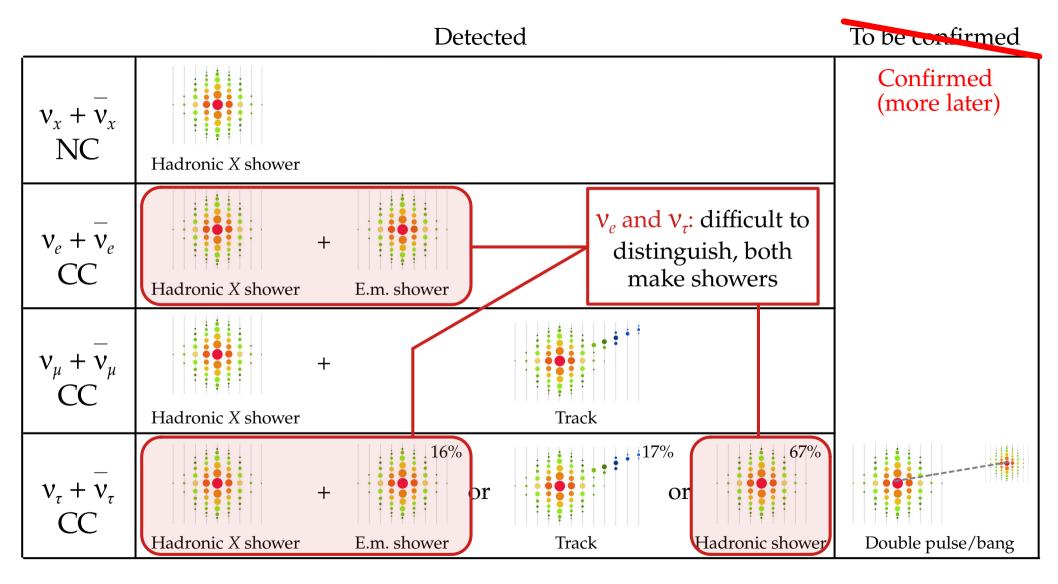


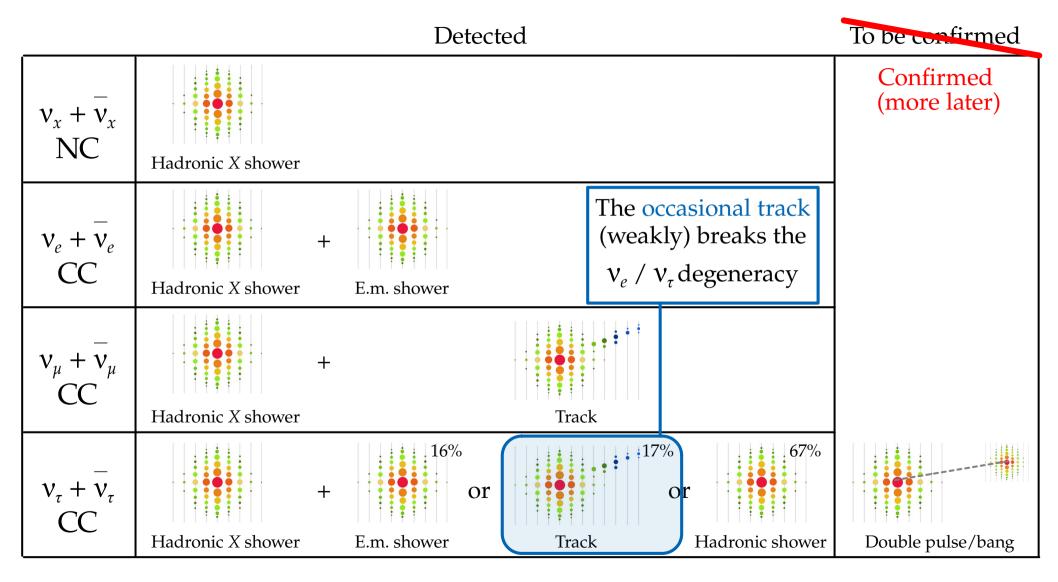
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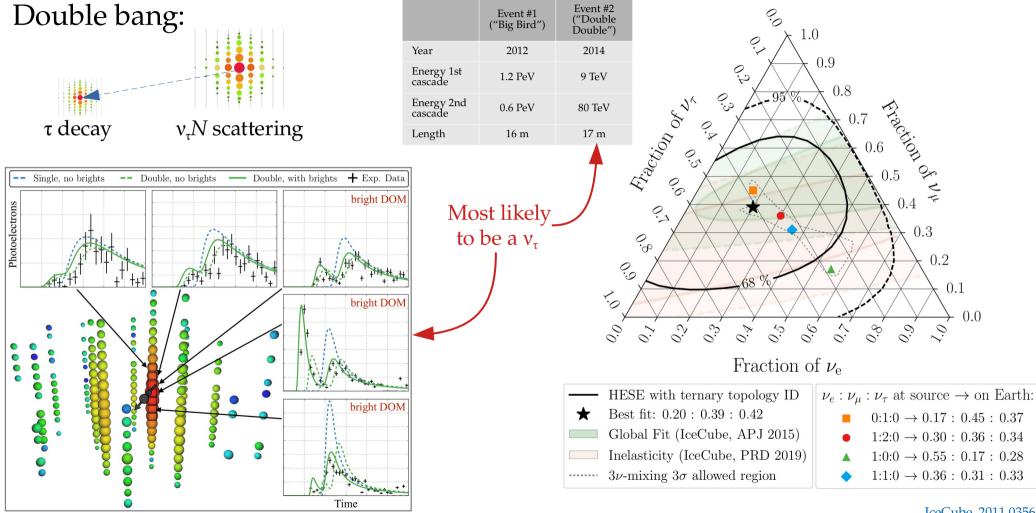






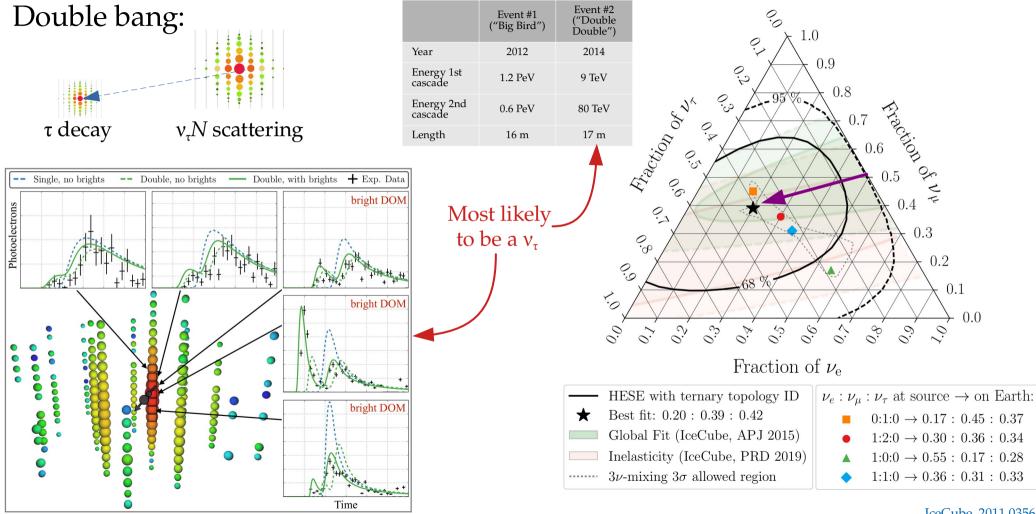


# New (IC 7.5 yr): First identified high-energy astrophysical $v_{\tau}$



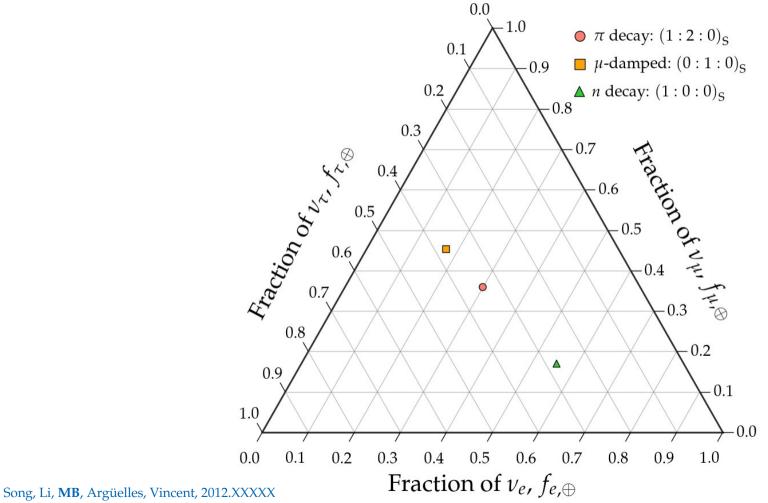
IceCube, 2011.03561

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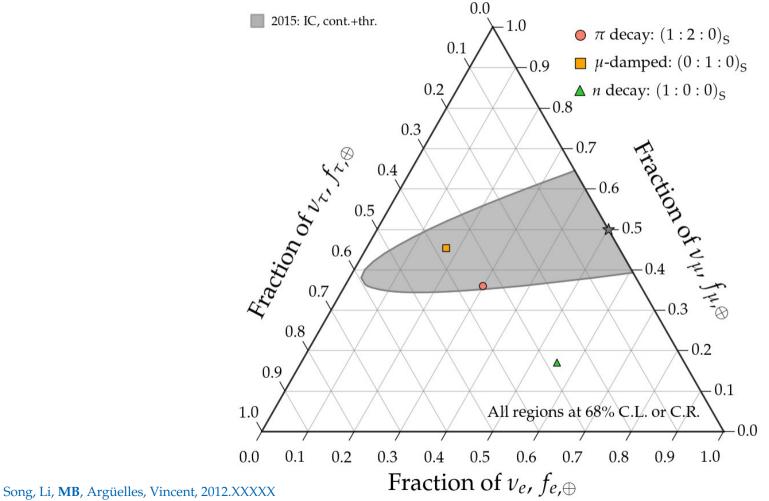


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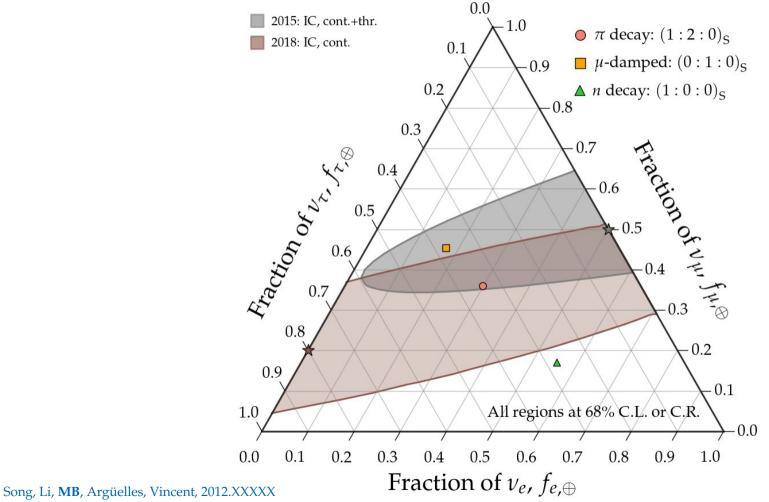
Song, Li, **MB**, Argüelles, Vincent, 2012.XXXXX IceCube, *PRL* 2015, *ApJ* 2015, *PRD* 2019, 2008.04323, 2011.03561



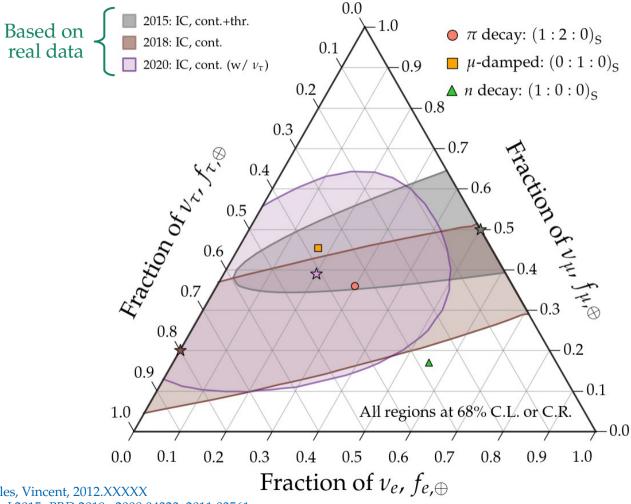
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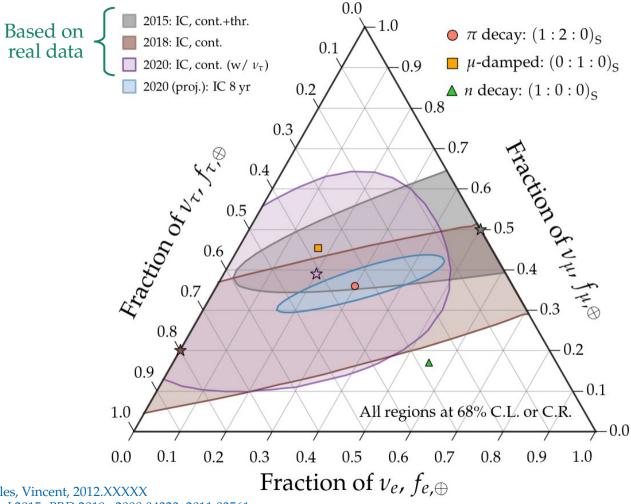
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*Status today:* 

Measurements are compatible with standard expectations (but errors are large!)

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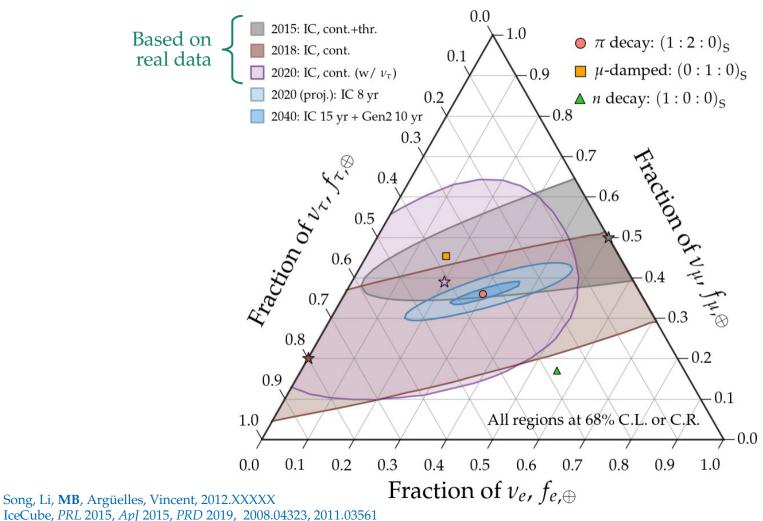
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*Near future* (~2020): × **5 reduction** using 8 yr of IC contained + thru.

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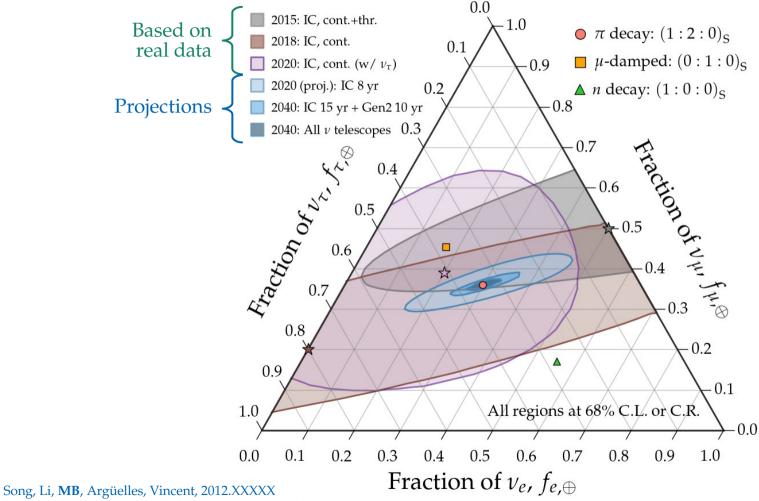
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14



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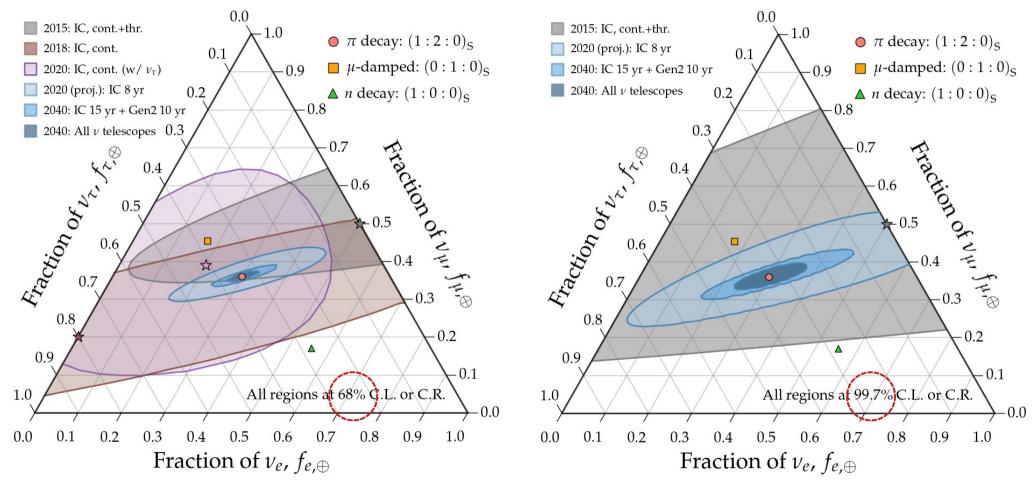
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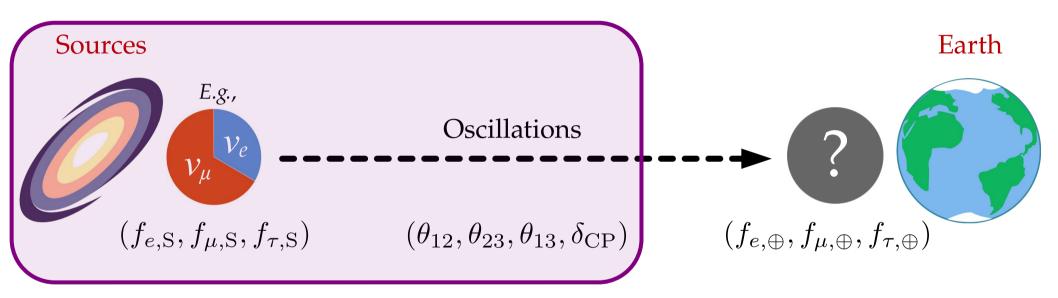
*Coming up* (~2040): **× 10 reduction** using Gen2 and all v telescopes

IceCube, *PRL* 2015, *ApJ* 2015, *PRD* 2019, 2008.04323, 2011.03561



Song, Li, **MB**, Argüelles, Vincent, 2012.XXXXX IceCube, *PRL* 2015, *ApJ* 2015, *PRD* 2019, 2008.04323, 2011.03561

#### *From sources to Earth:* we learn what to expect when measuring $f_{\alpha,\oplus}$



Theoretically palatable flavor regions  $\equiv MB, Beacom, Winter, PRL 2015$ Allowed regions of flavor ratios at Earth derived from oscillations

*Note:* The original palatable regions were frequentist [MB, Beacom, Winter, *PRL* 2015]; the new ones are Bayesian

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Ingredient #1: Flavor ratios at the source,  $(f_{e,S}, f_{\mu,S}, f_{\tau,S})$ 

Fix at one of the benchmarks (pion decay, muon-damped, neutron decay)

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Explore all possible combinations

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0.65

0.55

 $\sin^2 \theta_{23}$ 

0.60

2020: Use  $\chi^2$  profiles from 2.0 the NuFit 5.0 global fit 1.8 (solar + atmospheric 1.6 1.4 + reactor + accelerator) 1.2 Esteban *et al.*, *JHEP* 2020  $\delta_{\rm CP}/\pi$ www.nu-fit.org 1.0 0.8 0.6 0.4 0.2 NuFit 5.0 0.400.45 0.50

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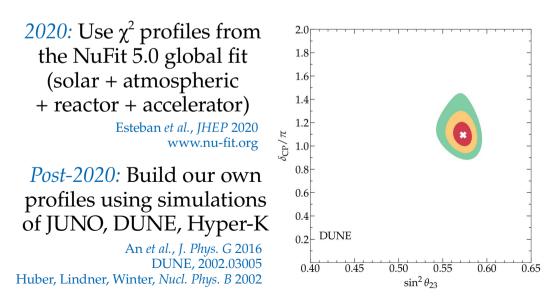
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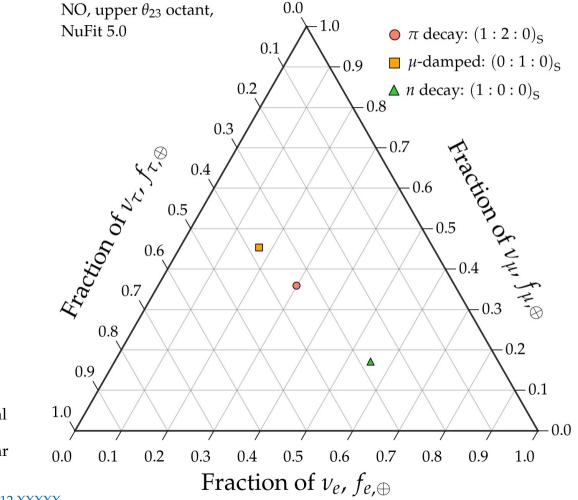
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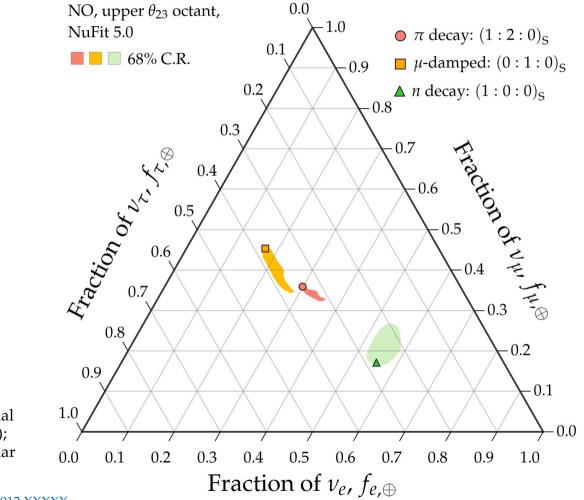
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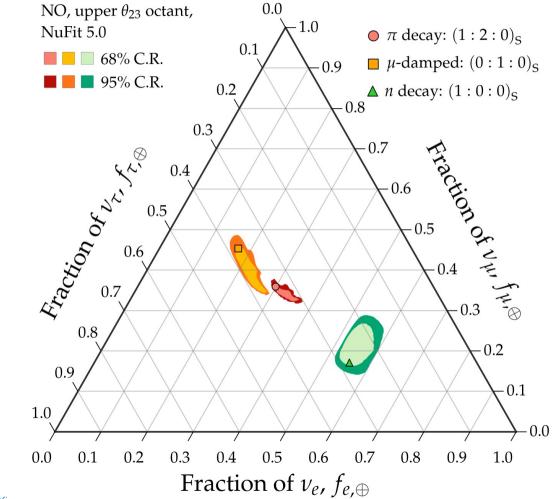




Note: All plots s

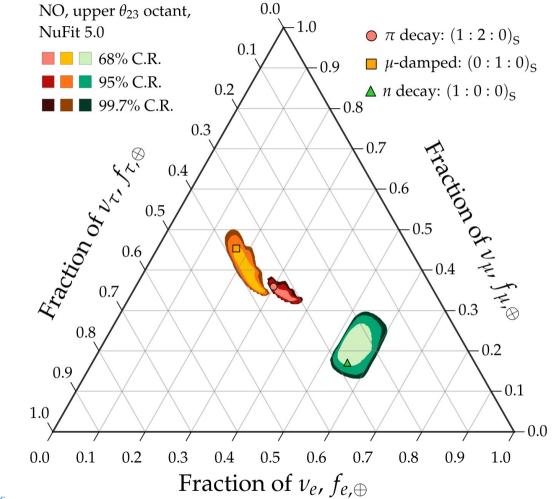


Note:

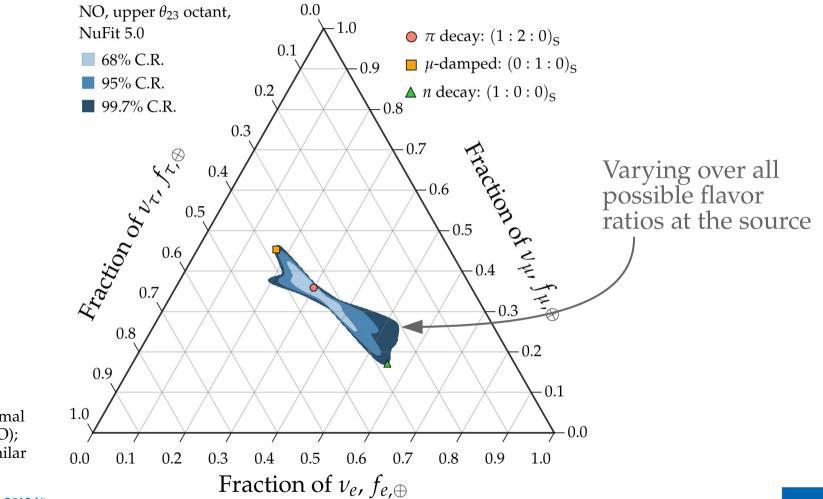


*Note:* All plots shown are for normal

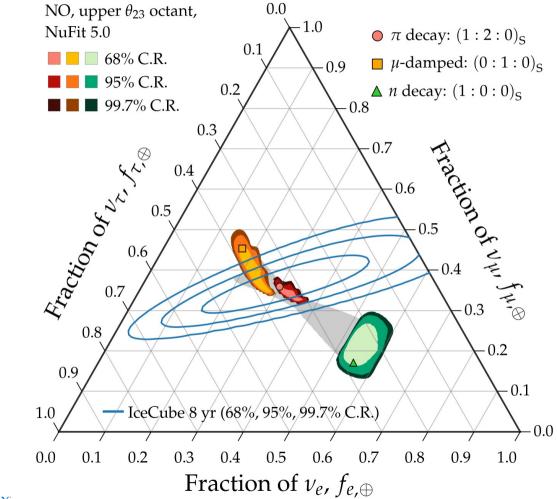
neutrino mass ordering (NO); inverted ordering looks similar



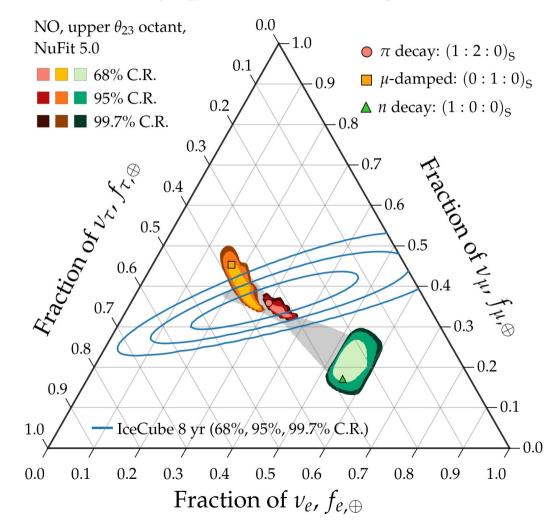
Note:



Note:



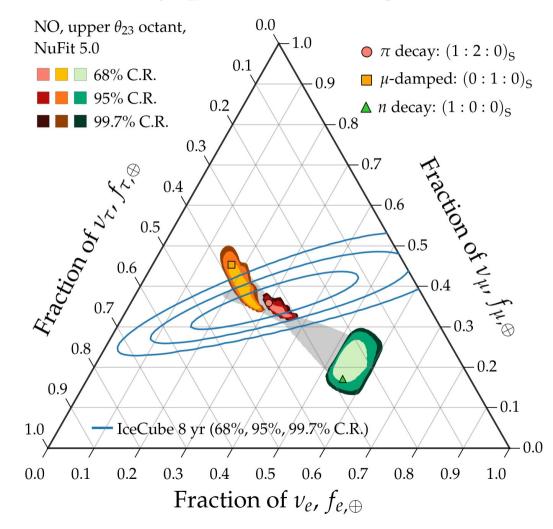
#### Note:



Two limitations:

*Allowed flavor regions overlap* – Insufficient precision in the mixing parameters

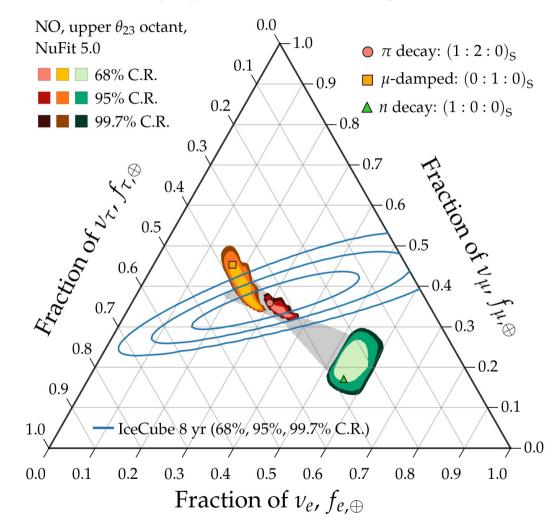
Measurement of flavor ratios – Cannot distinguish between pion-decay and muon-damped benchmarks even at 68% C.R. (1σ)



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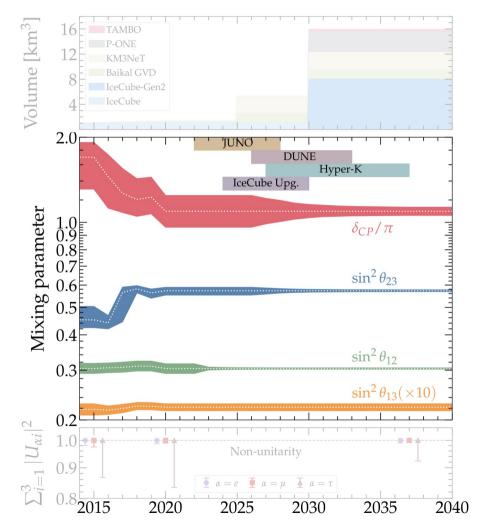
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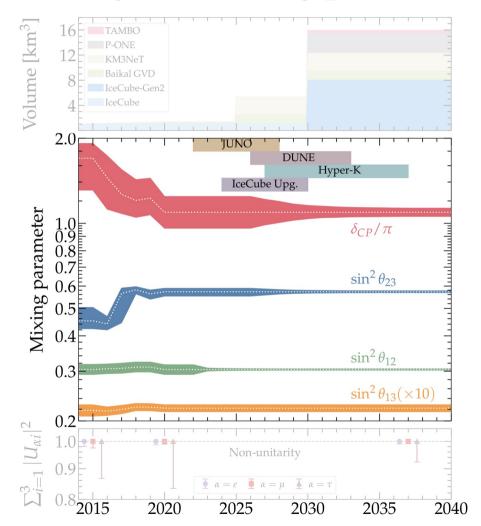
Measurement of flavor ratios – Cannot distinguish between pion-decay and muon-damped benchmarks even at 68% C.R. (1σ) Will be overcome by 2040



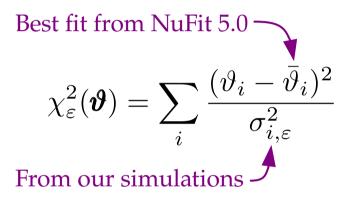
We can compute the oscillation probability more precisely:

$$f_{\alpha,\oplus} = \sum_{\beta=e,\mu,\tau} P_{\beta\alpha} f_{\beta,\mathrm{S}}$$

So we can convert back and forth between source and Earth more precisely

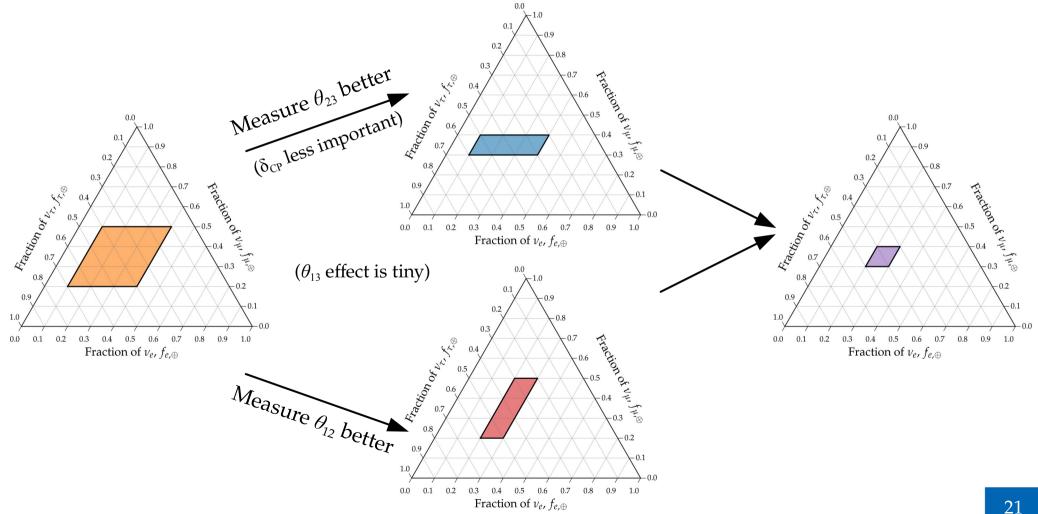


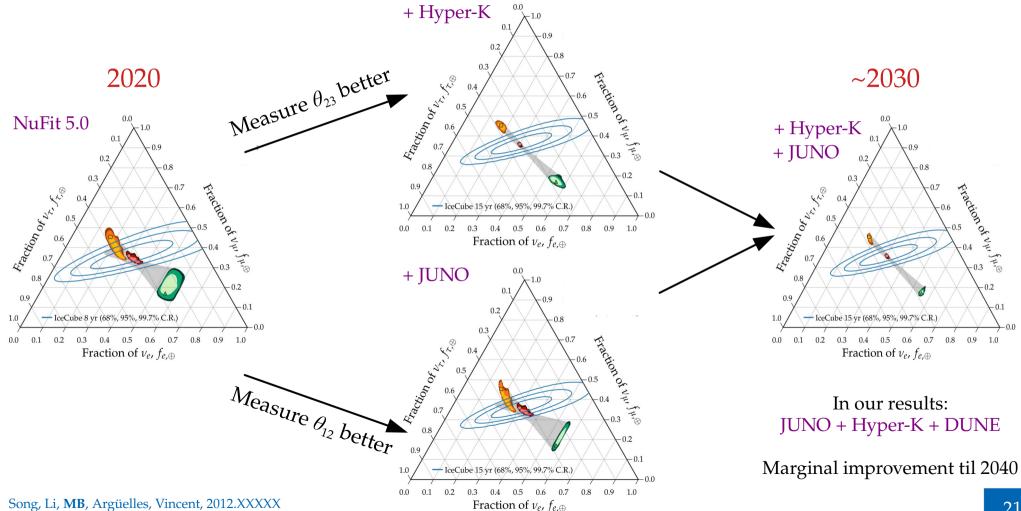
For a future experiment ε = JUNO, DUNE, Hyper-K:



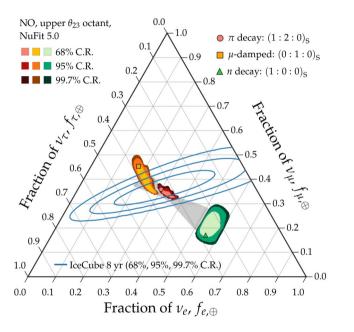
We combine experiments in a likelihood:

$$-2\log \mathcal{L}(\boldsymbol{\theta}) = \sum_{\varepsilon} \chi_{\varepsilon}^2(\boldsymbol{\vartheta})$$



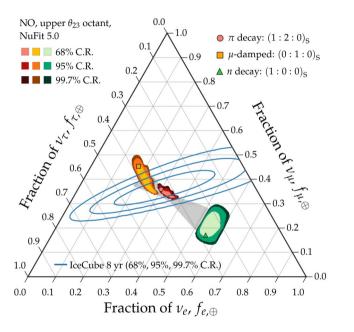


2020



Allowed regions: overlapping Measurement: imprecise

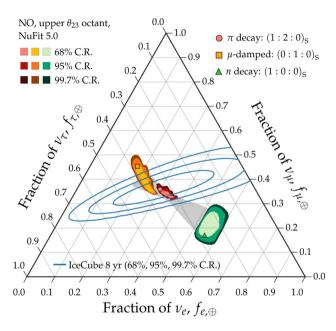
2020



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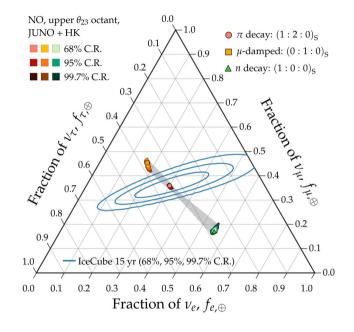
Not ideal

2020



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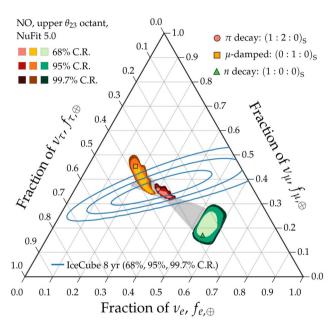
Not ideal



2030

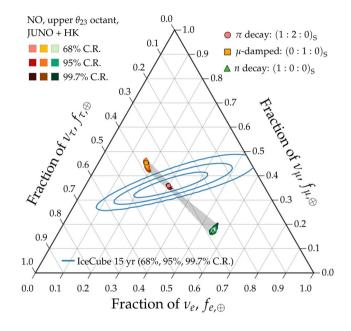
Allowed regions: well separated Measurement: improving

2020



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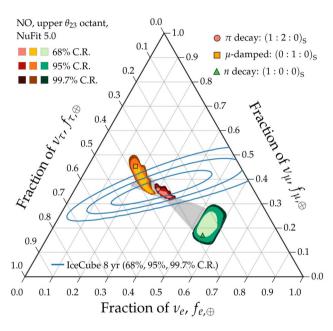


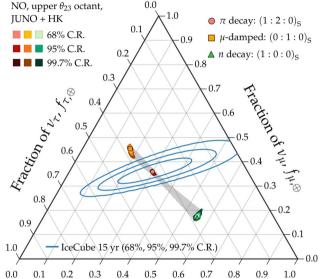
2030

Allowed regions: well separated Measurement: improving

Nice

2020





2030

0.0

Fraction of  $v_e$ ,  $f_{e,\oplus}$ 

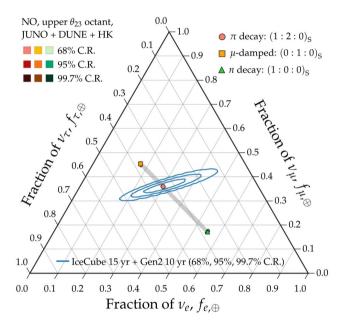
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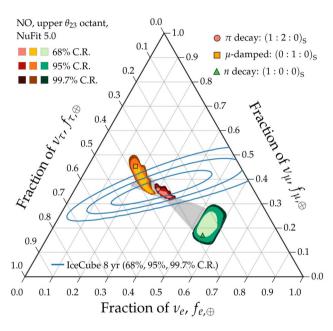
Nice

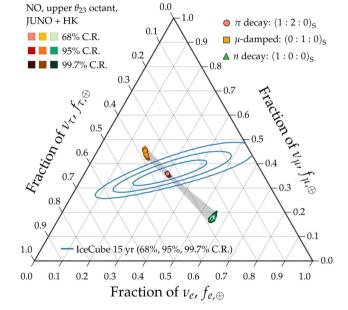
2040



Allowed regions: well separated Measurement: precise

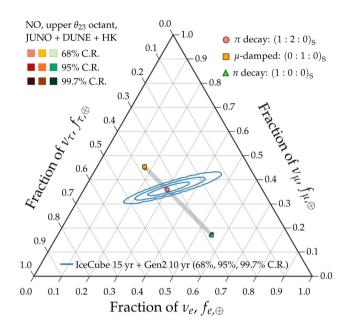
2020





2030

2040



Allowed regions: overlapping Measurement: imprecise

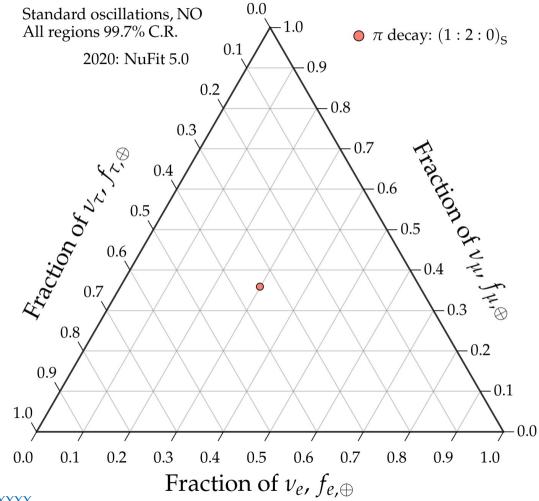
Not ideal

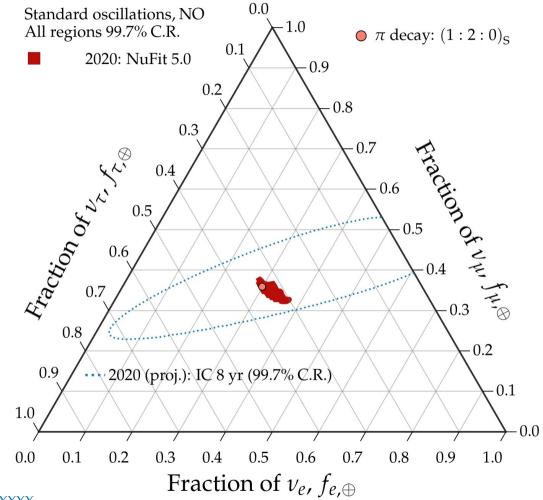
Allowed regions: well separated Measurement: improving

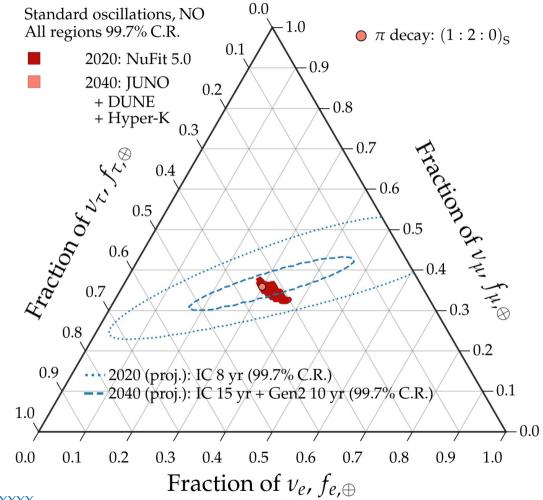
Nice

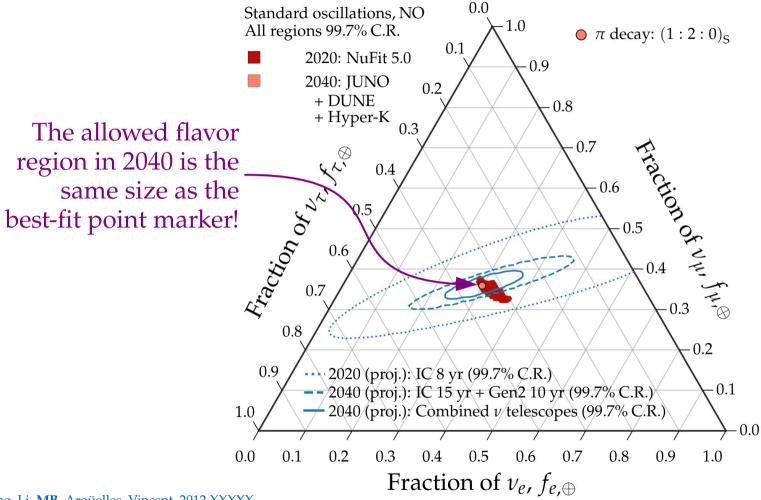
Allowed regions: well separated Measurement: precise

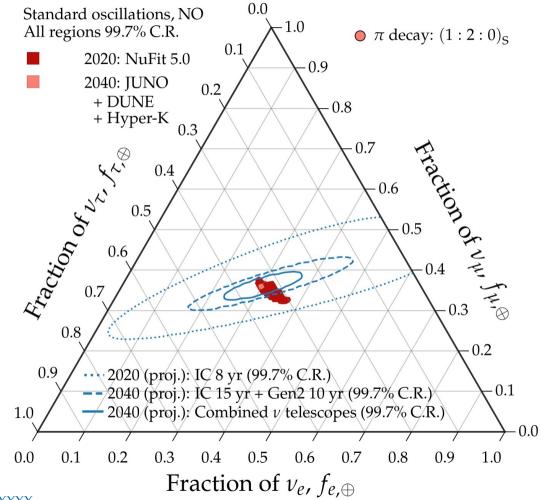
Success

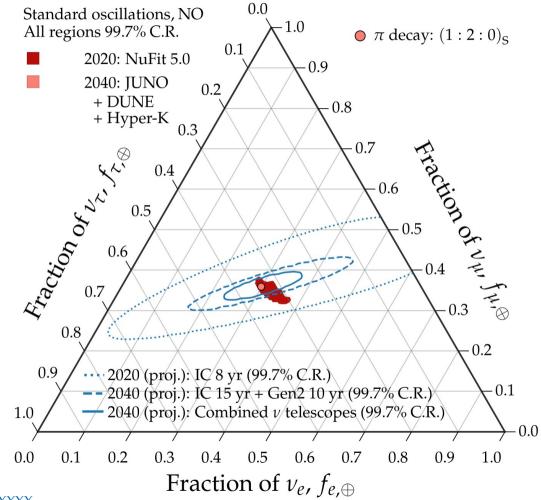


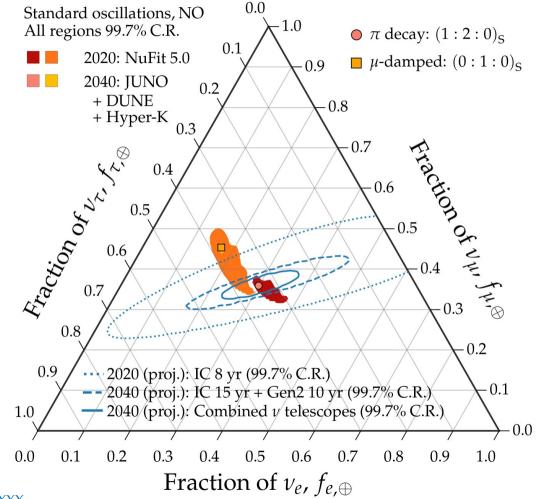


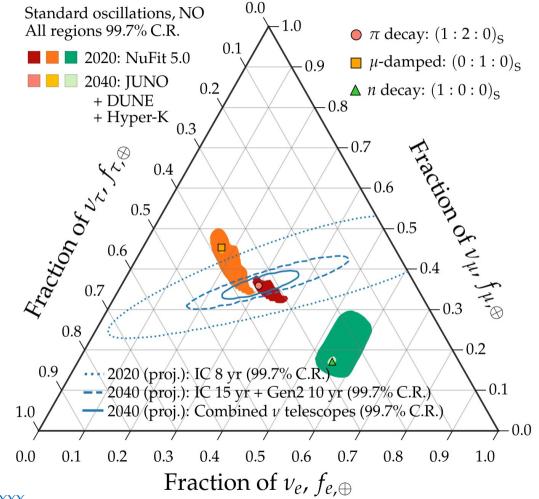


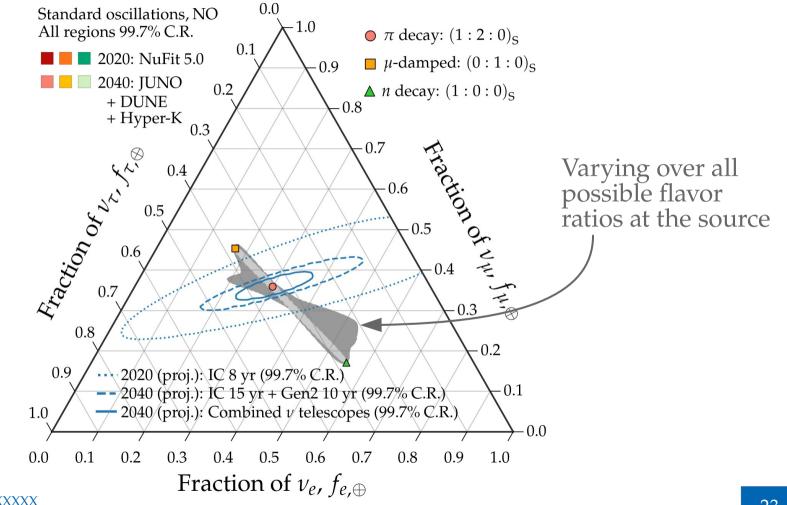


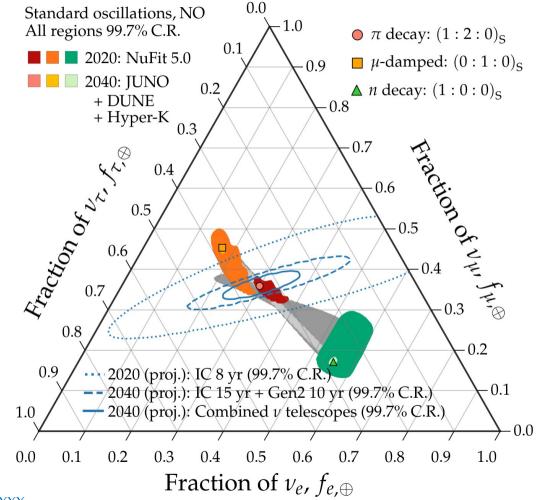


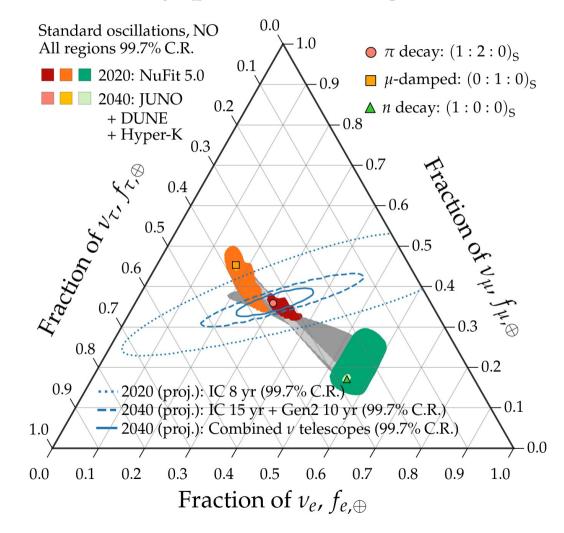












#### By 2040:

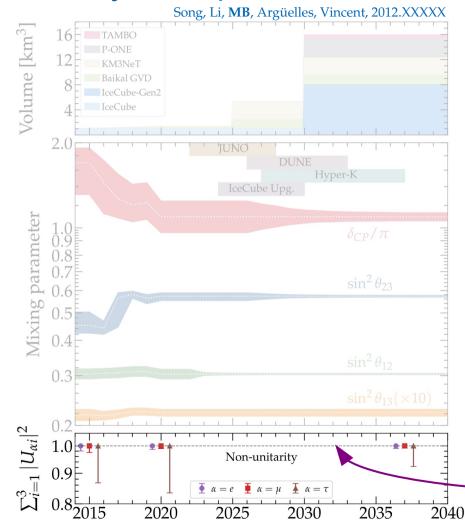
#### Theory –

Mixing parameters known precisely: allowed flavor regions are *almost* points (already by 2030)

*Measurement of flavor ratios* – Can distinguish between similar predictions at 99.7% C.R. (3σ)

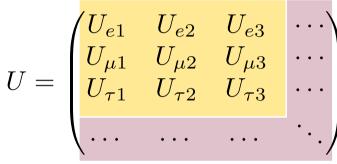
Can finally use the full power of flavor composition for astrophysics and neutrino physics

### No unitarity? *No problem*



The  $3 \times 3$  active mixing matrix is a non-unitary sub-matrix of a bigger one:

Active flavors



#### Additional sterile flavors

The elements  $|U_{\alpha i}|^2$  for active flavors can be measured *without* assuming unitarity

Because the sub-matrix is not-unitary  $(U_{3\nu}^{\dagger}U_{3\nu} \neq 1)$ , the "row sum" may be <1

Ellis, Kelly, Li, 2008.01088 Parke & Ross-Lonergan, *PRD* 2016

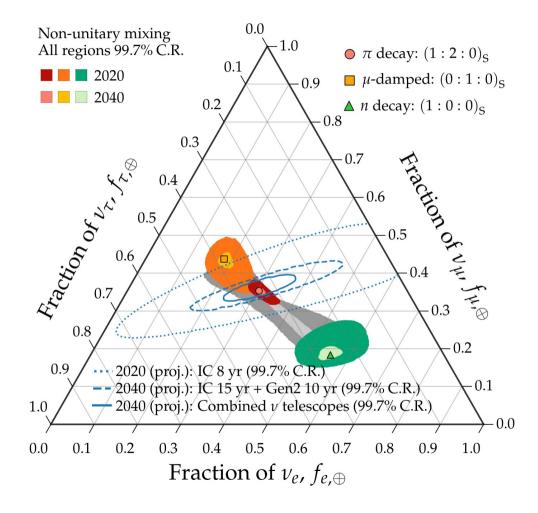
## No unitarity? No problem

Flavor ratios at Earth:

$$f_{\alpha,\oplus} = \sum_{\beta=e,\mu,\tau} P_{\beta\alpha} f_{\beta,\mathrm{S}}$$

Same as for standard oscillations...

... but the probability is computed directly using the values of the  $|U_{\alpha i}|^2$ (instead of the mixing angles)

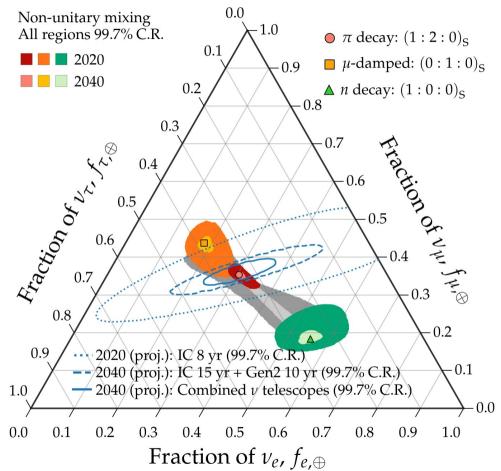


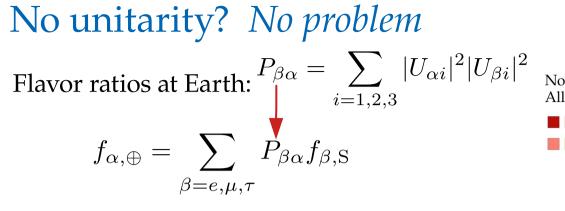
No unitarity? No problem  
Flavor ratios at Earth: 
$$P_{\beta\alpha} = \sum_{i=1,2,3} |U_{\alpha i}|^2 |U_{\beta i}|^2$$

$$f_{\alpha,\oplus} = \sum_{\beta=e,\mu,\tau} P_{\beta\alpha} f_{\beta,S}$$

Same as for standard oscillations...

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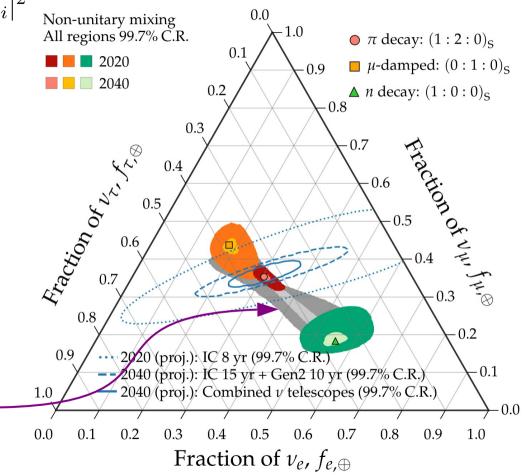




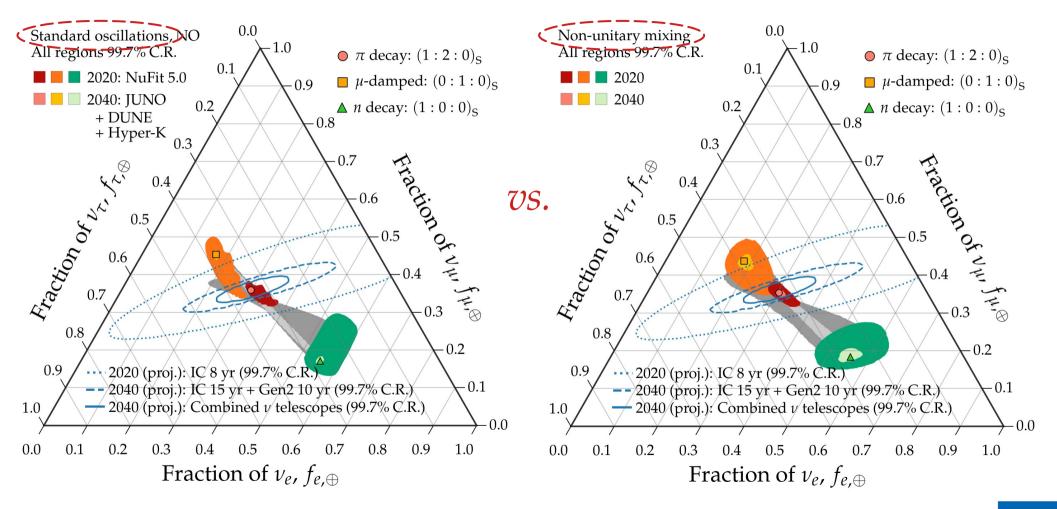
Same as for standard oscillations...

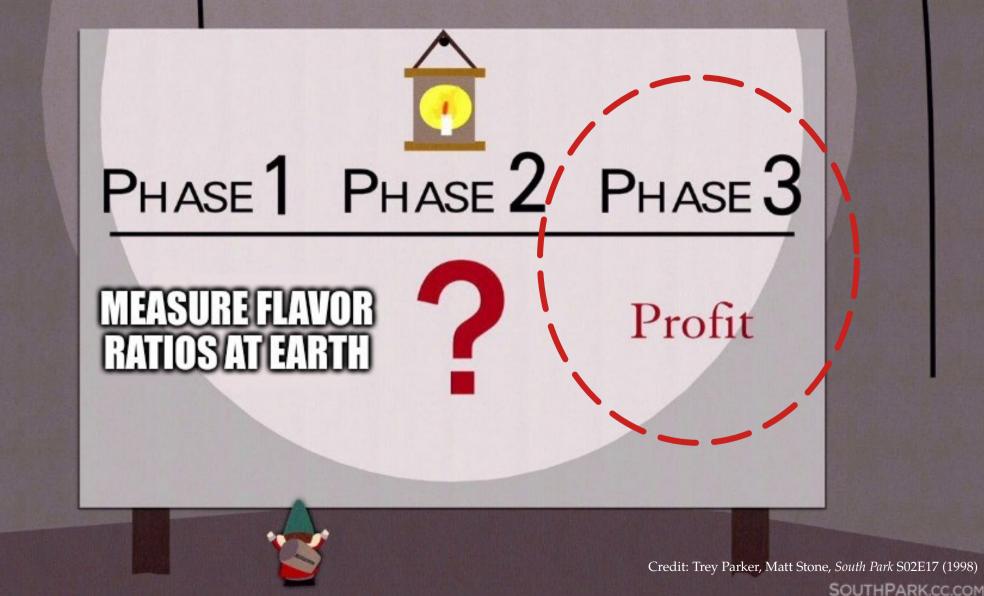
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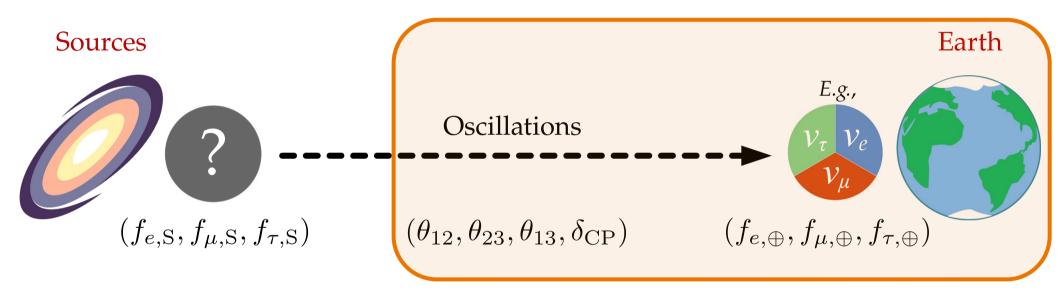
The allowed flavor regions \_\_\_\_\_\_are bigger, but *not much bigger*!



# No unitarity? No problem

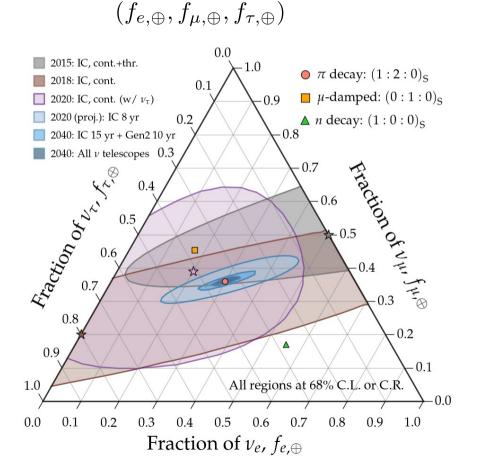




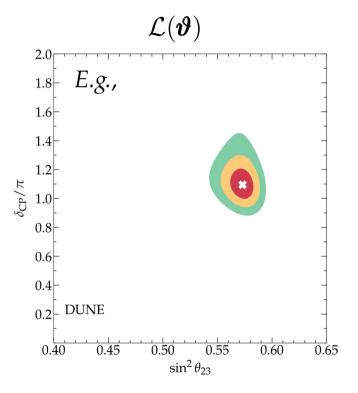


*From Earth to sources:* we let the data teach us about  $f_{\alpha,S}$ 

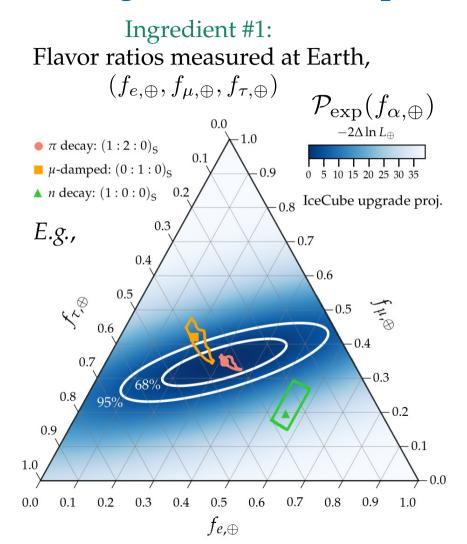
#### Ingredient #1: Flavor ratios measured at Earth,



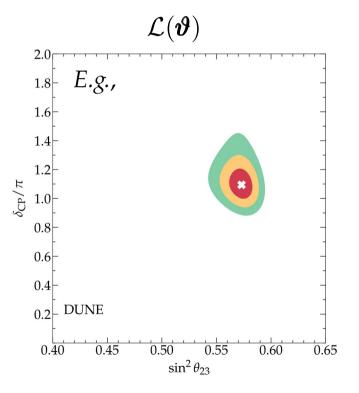
Ingredient #2: Probability density of mixing parameters ( $\theta_{12}, \theta_{23}, \theta_{13}, \delta_{CP}$ )



Song, Li, **MB**, Argüelles, Vincent, 2012.XXXXX **MB** & Ahlers, *PRL* 2019



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Song, Li, **MB**, Argüelles, Vincent, 2012.XXXXX **MB** & Ahlers, *PRL* 2019

Ingredient #1: Flavor ratios measured at Earth,  $(f_{e,\oplus}, f_{\mu,\oplus}, f_{\tau,\oplus})$  Ingredient #2: Probability density of mixing parameters ( $\theta_{12}, \theta_{23}, \theta_{13}, \delta_{CP}$ )

Posterior probability of  $f_{\alpha,S}$  [MB & Ahlers, *PRL* 2019]:

$$\mathcal{P}(\boldsymbol{f}_s) = \int d\boldsymbol{\vartheta} \mathcal{L}(\boldsymbol{\vartheta}) \mathcal{P}_{\mathrm{exp}}(\boldsymbol{f}_{\oplus}(\boldsymbol{f}_{\mathrm{S}}, \boldsymbol{\vartheta}))$$

30

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Oscillation experiments Neutrino telescopes

30

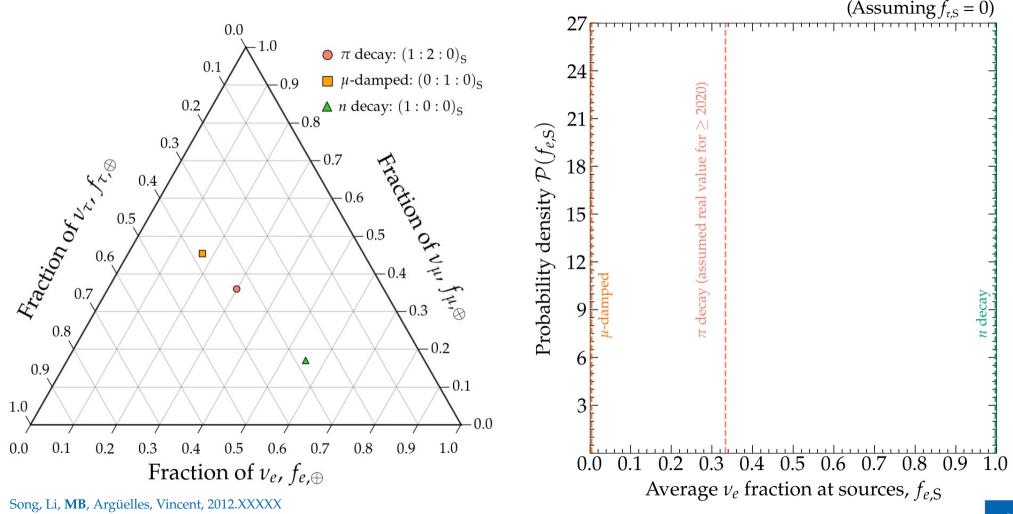
Ingredient #1: Flavor ratios measured at Earth,  $(f_{e,\oplus}, f_{\mu,\oplus}, f_{\tau,\oplus})$  Ingredient #2: Probability density of mixing parameters ( $\theta_{12}, \theta_{23}, \theta_{13}, \delta_{CP}$ )

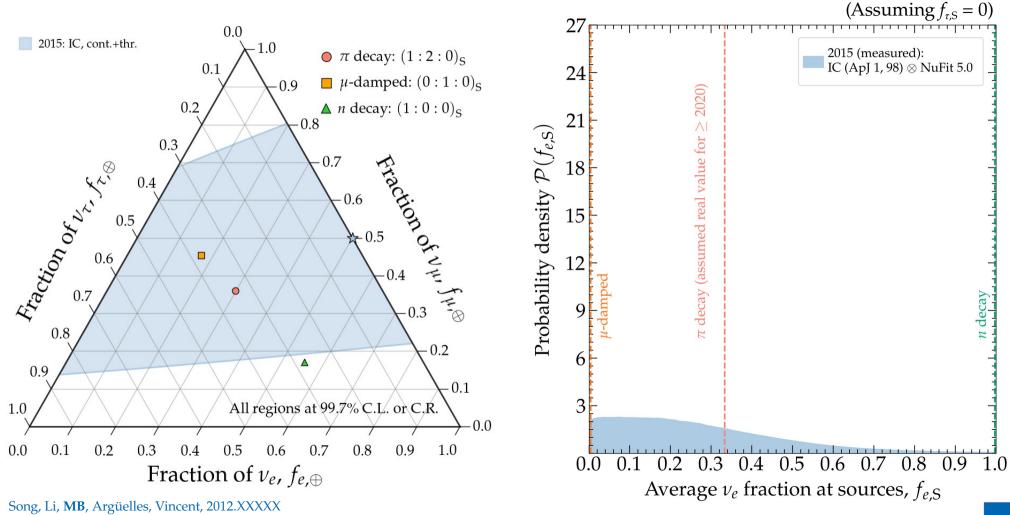
Posterior probability of  $f_{\alpha,S}$  [MB & Ahlers, *PRL* 2019]:

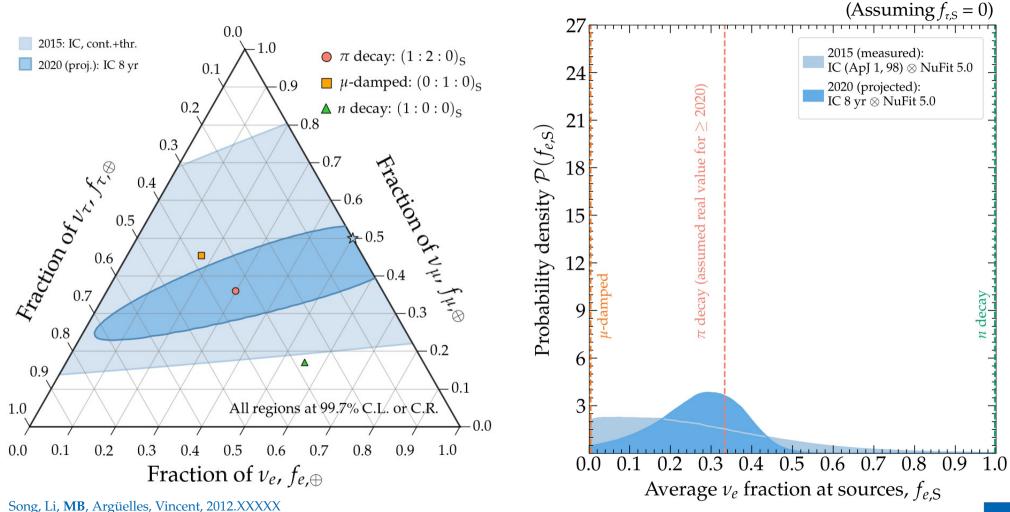
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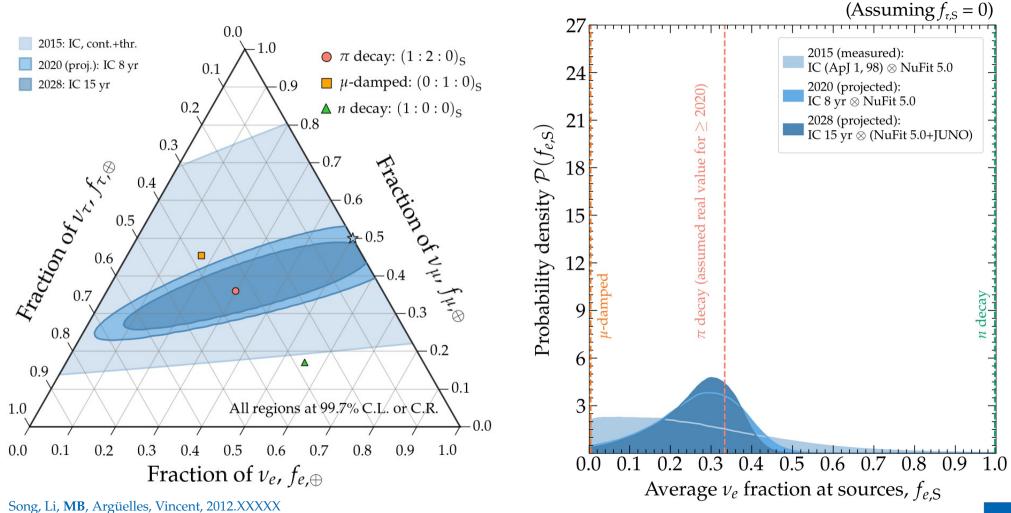
Oscillation experiments Neutrino telescopes

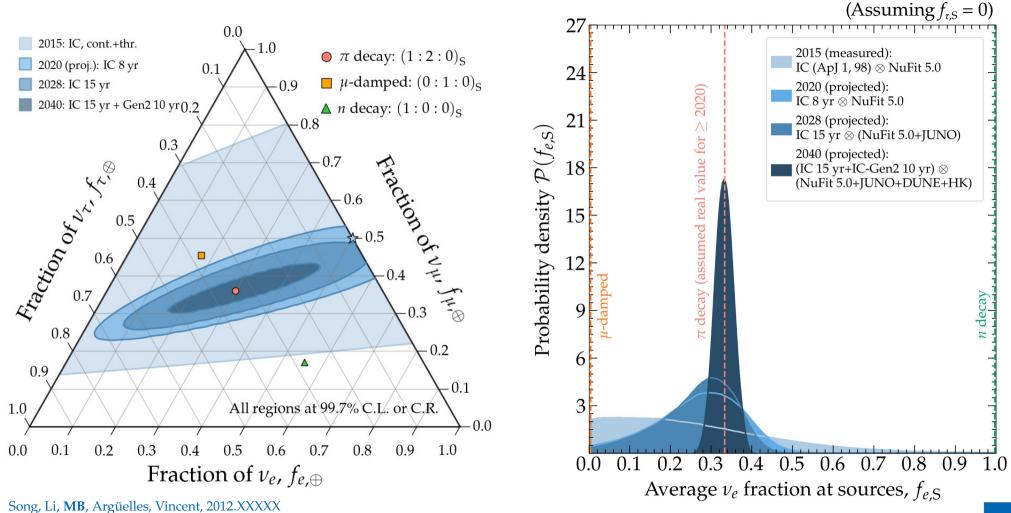
30

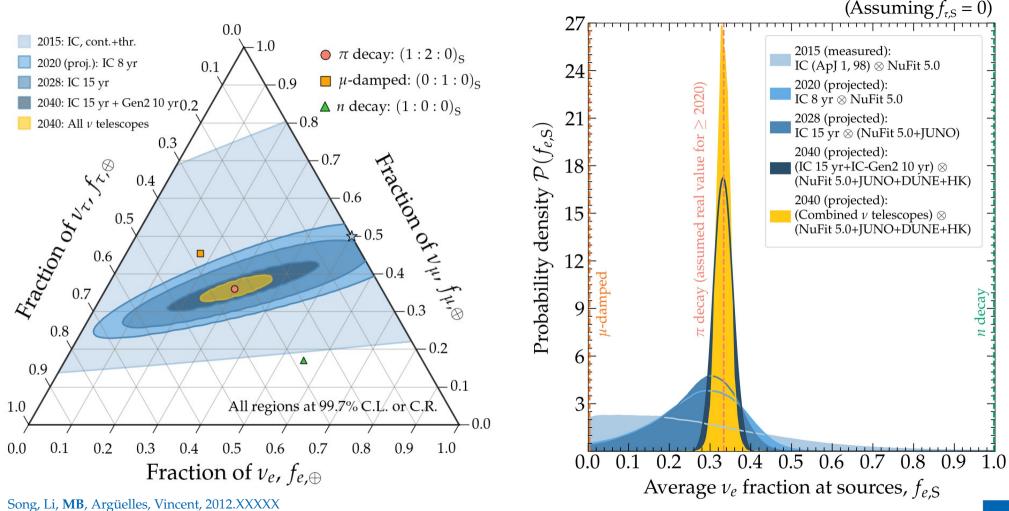




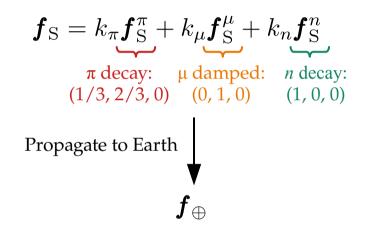






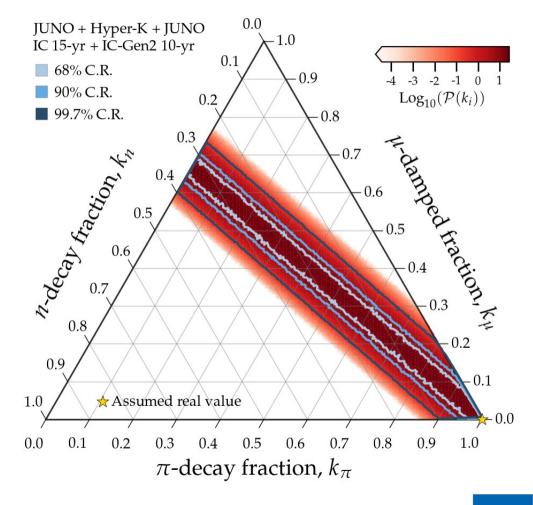


Can we detect the contribution of multiple v production mechanisms?

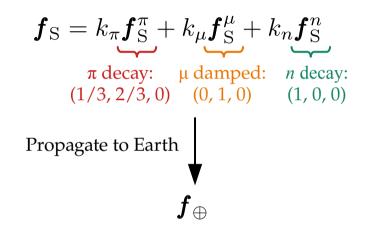


Assume real value  $k_{\pi} = 1$  ( $k_{\mu} = k_n = 0$ )

*By 2040, how well will we recover the real value?* [Adding spectrum information (not shown) will likely help]

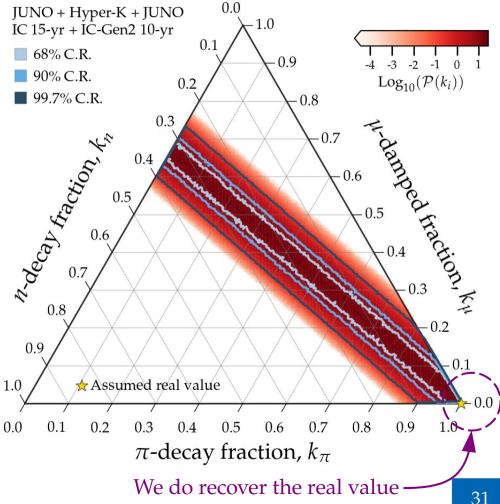


Can we detect the contribution of multiple v production mechanisms?

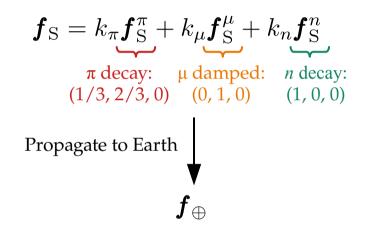


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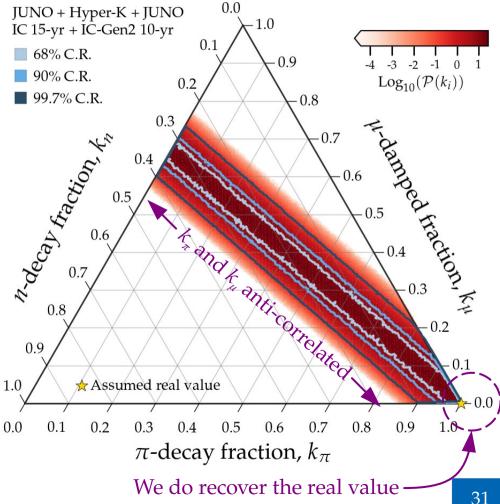


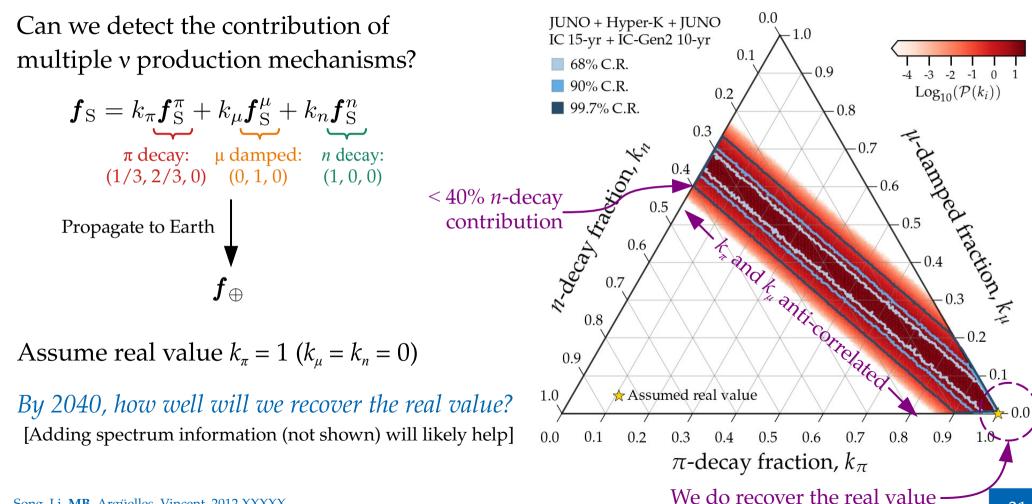
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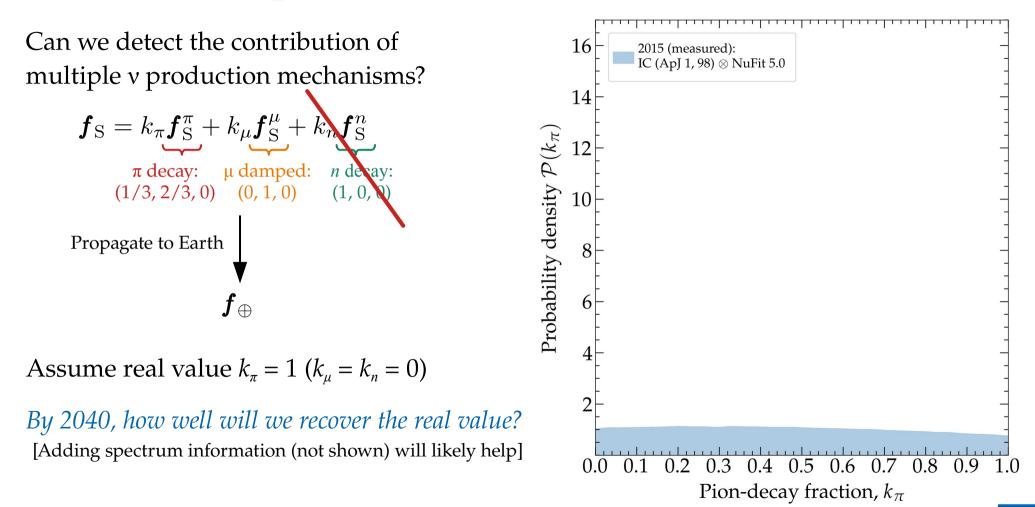
Can we detect the contribution of multiple v production mechanisms?

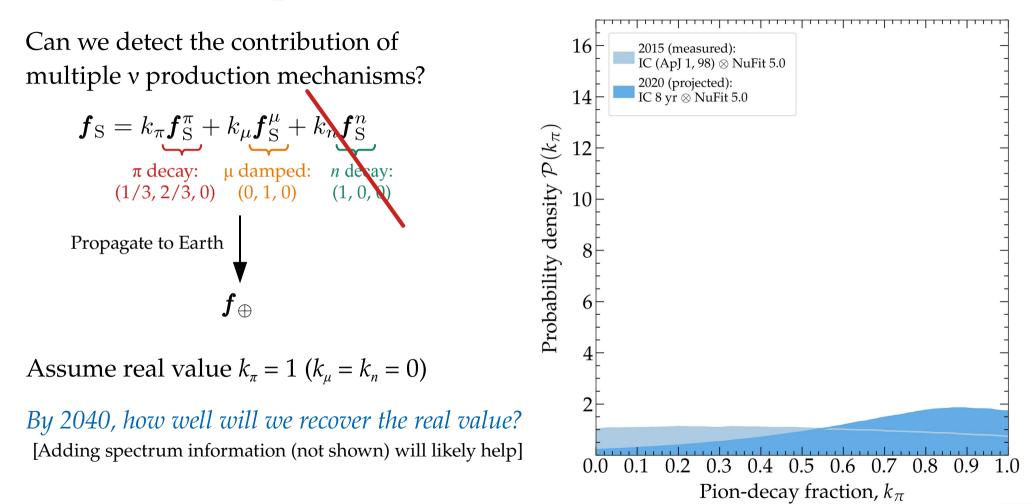
$$f_{\rm S} = k_{\pi} f_{\rm S}^{\pi} + k_{\mu} f_{\rm S}^{\mu} + k_{n} f_{\rm S}^{n}$$

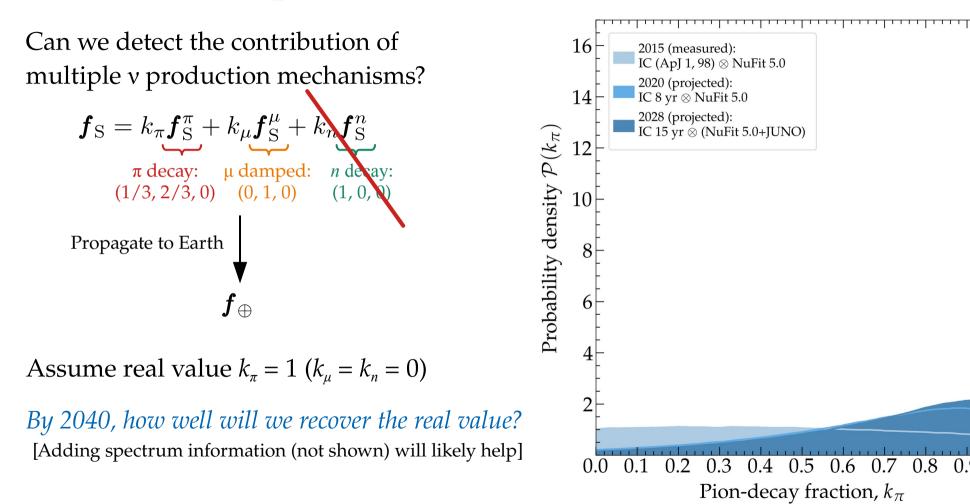
$$\frac{\pi \text{ decay: } \mu \text{ damped: } n \text{ decay: } (1/3, 2/3, 0) \quad (0, 1, 0) \quad (1, 0, 0)$$
Propagate to Earth
$$f_{\oplus}$$

Assume real value  $k_{\pi} = 1$  ( $k_{\mu} = k_n = 0$ )

*By 2040, how well will we recover the real value?* [Adding spectrum information (not shown) will likely help]



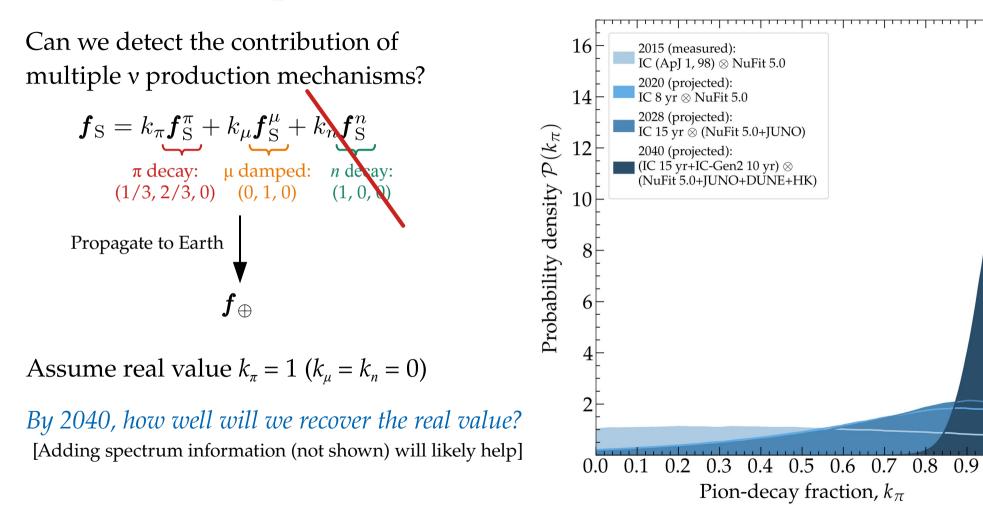




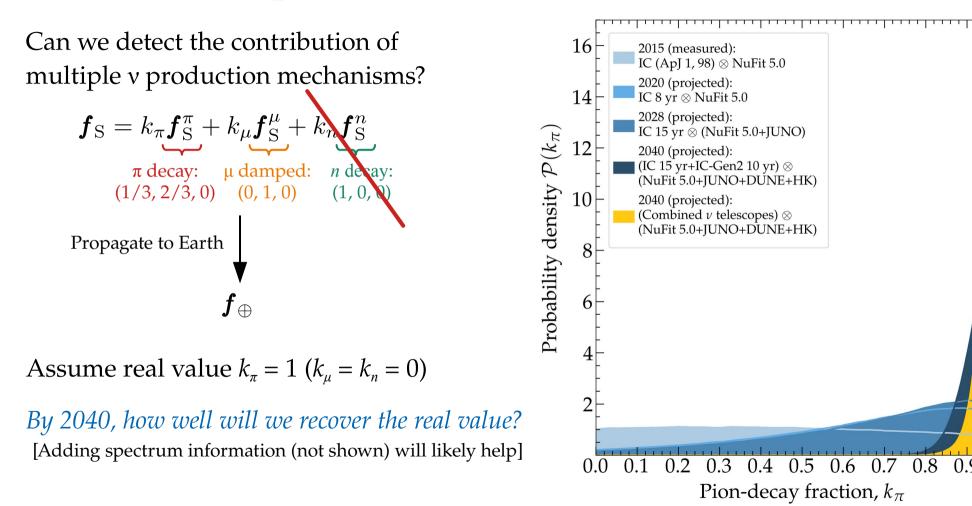
1.0

0.8

0.9



1.0



Song, Li, MB, Argüelles, Vincent, 2012.XXXXX

1.0

0.8 0.9

#### Using high-energy neutrinos as magnetometers

If sources have strong magnetic fields, charged particles cool via synchrotron:

**MB**, Tamborra, *PRD* 2020 See also: Winter, *PRD* 2013

# Using high-energy neutrinos as magnetometers

If sources have strong magnetic fields, charged particles cool via synchrotron:

## Proton cooling

Induce a high-energy cut-off in the emitted v spectrum:

$$\begin{split} E_{\nu}^{\prime 2} \frac{dN_{\nu}}{dE_{\nu}^{\prime}} &\propto E_{\nu}^{\prime 2-\alpha_{\nu}} e^{-E_{\nu}^{\prime}/E_{\nu}^{\prime \max}} \\ E_{\nu}^{\max} &\approx \frac{10^{10} \Gamma \text{ GeV}}{\sqrt{B^{\prime}/\text{G}}} \qquad \overbrace{p+\gamma(p) \rightarrow \pi^{+} \rightarrow \mu^{+} + \nu_{\mu}} \\ & \downarrow \bar{\nu}_{\mu} + e^{+} + \nu_{e} \end{split}$$

# Using high-energy neutrinos as magnetometers

( *p* -

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$$E_{\nu}^{\max} \approx \frac{10^{10} \Gamma \text{ GeV}}{\sqrt{B^{\prime}/\text{G}}}$$

Muon cooling

Change flavor composition:

$$(f_{e,\mathrm{S}}, f_{\mu,\mathrm{S}}, f_{\tau,\mathrm{S}}) = \begin{cases} (\frac{1}{3}, \frac{2}{3}, 0), & \text{if } E_{\nu} < E_{\nu,\mu}^{\mathrm{sync}} \\ (0, 1, 0), & \text{if } E_{\nu} \ge E_{\nu,\mu}^{\mathrm{sync}} \end{cases}$$
$$E_{\nu,\mu}^{\mathrm{sync}} \approx 10^{9} \Gamma \frac{\mathrm{G}}{B'} \mathrm{GeV}$$
$$\rightarrow \pi^{+} \rightarrow \mu^{+} + \nu_{\mu}$$
$$\downarrow \overline{\nu}_{\mu} + e^{+} + \nu_{e}$$

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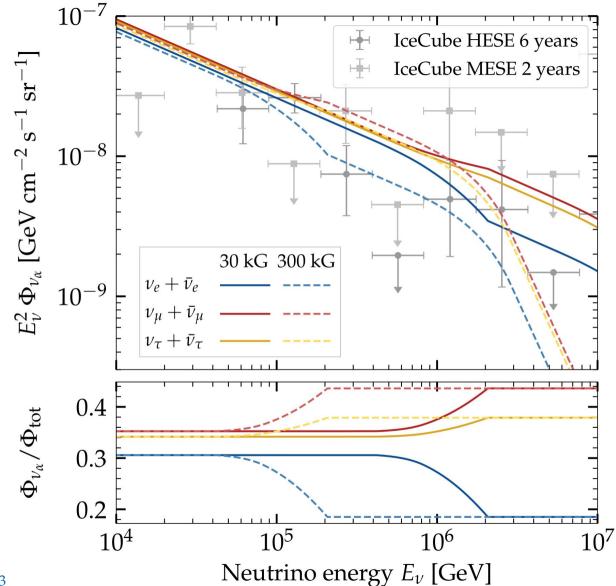
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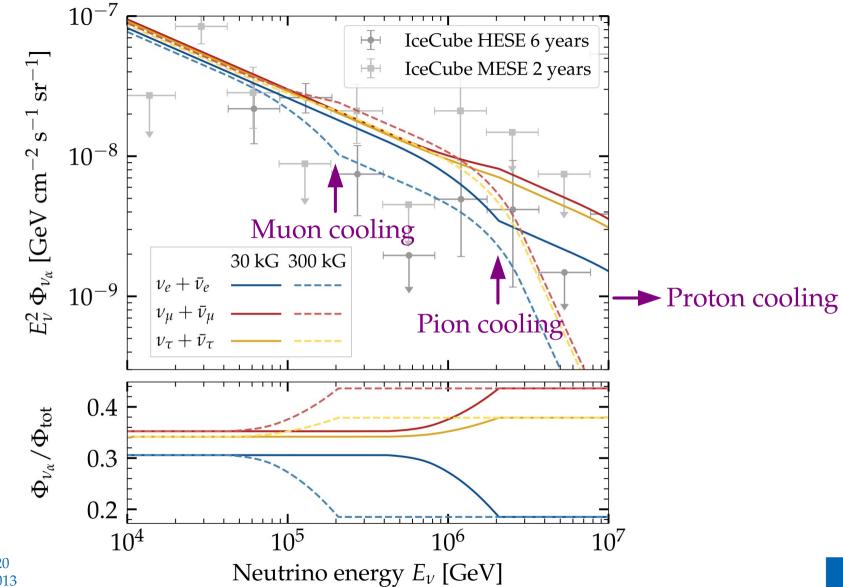
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$$\pi^{+} \rightarrow \mu^{+} + \nu_{\mu}$$
$$\rightarrow \bar{\nu}_{\mu} + e^{+} + \nu_{e}$$

# Pion cooling Steepen the v spectrum: $\alpha_{\nu} = \begin{cases} \gamma, & \text{if } E_{\nu} < E_{\nu,\pi}^{\text{sync}} \\ \gamma+2, & \text{if } E_{\nu} \ge E_{\nu,\pi}^{\text{sync}} \end{cases}$ $E_{\nu,\pi}^{\text{sync}} \approx 10^{10} \Gamma \frac{\text{G}}{B'} \text{ GeV}$

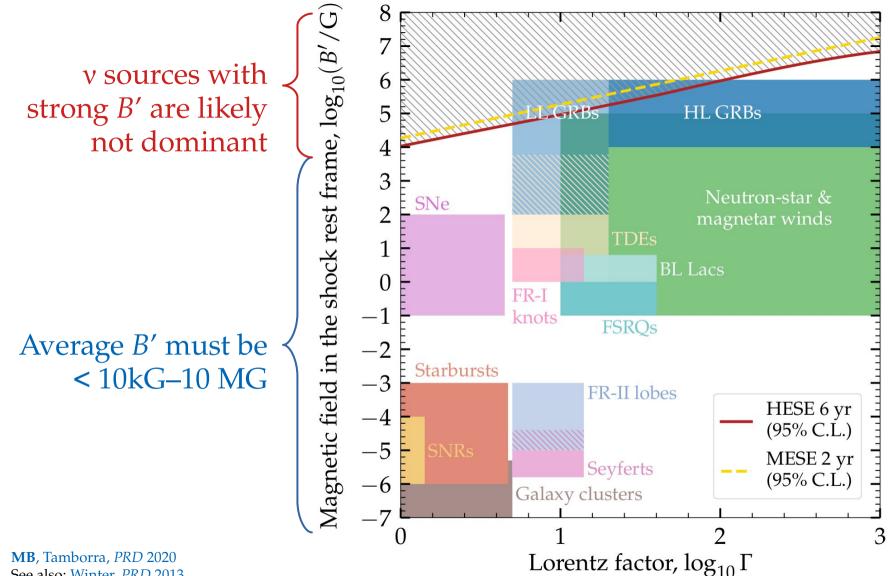
**MB**, Tamborra, *PRD* 2020 See also: Winter, *PRD* 2013



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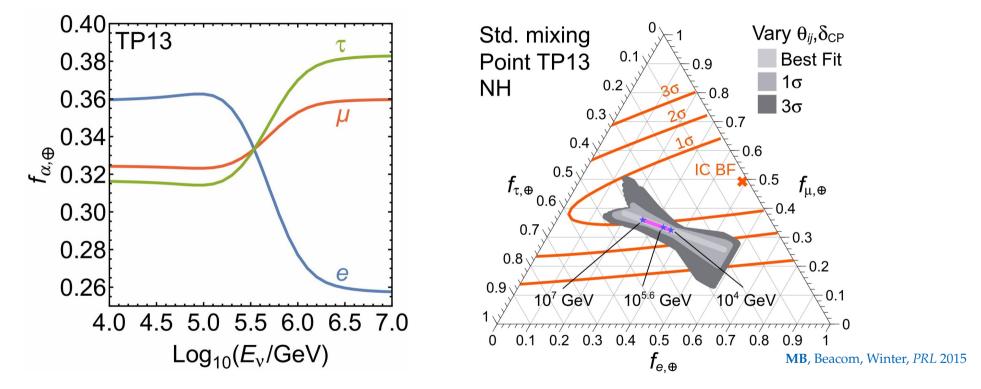
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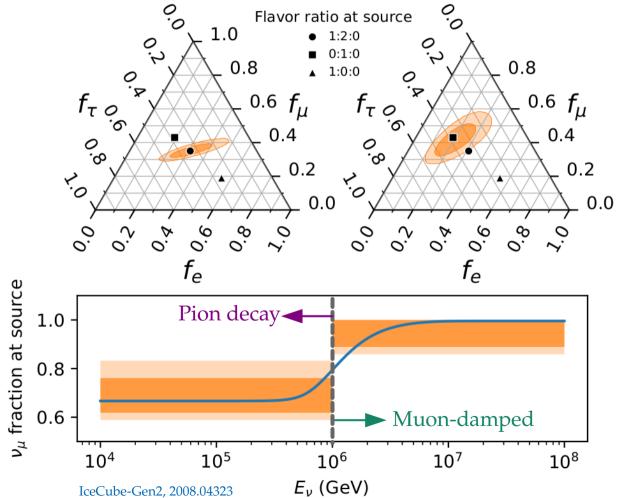
# Energy dependence of the flavor composition?

Different neutrino production channels accessible at different energies -

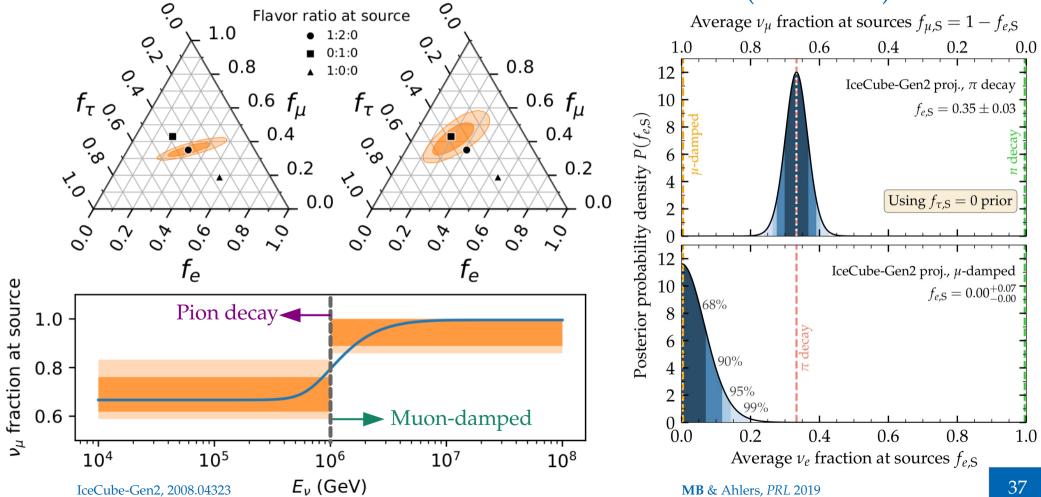


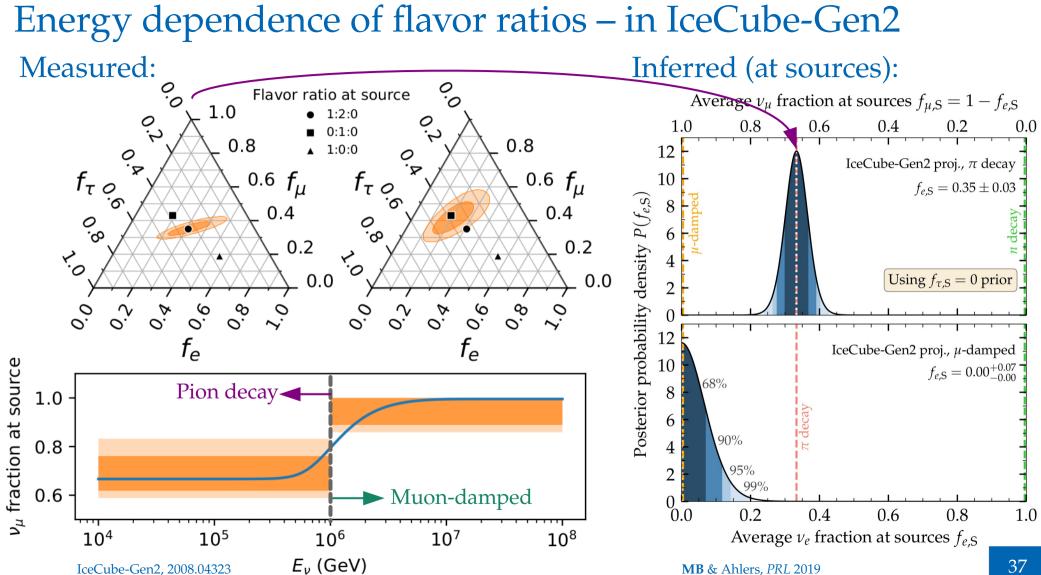
TP13: py model, target photons from e<sup>-</sup>e<sup>+</sup> annihilation [Hümmer+, Astropart. Phys. 2010]
 Will be difficult to resolve [Kashti, Waxman, PRL 2005; Lipari, Lusignoli, Meloni, PRD 2007]

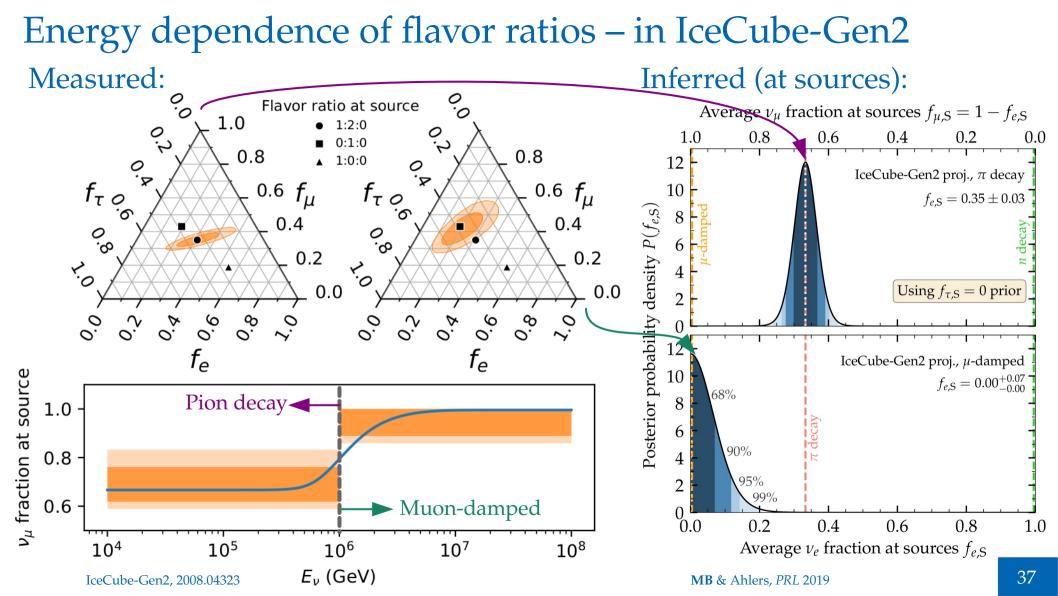
# Energy dependence of flavor ratios – in IceCube-Gen2 Measured:



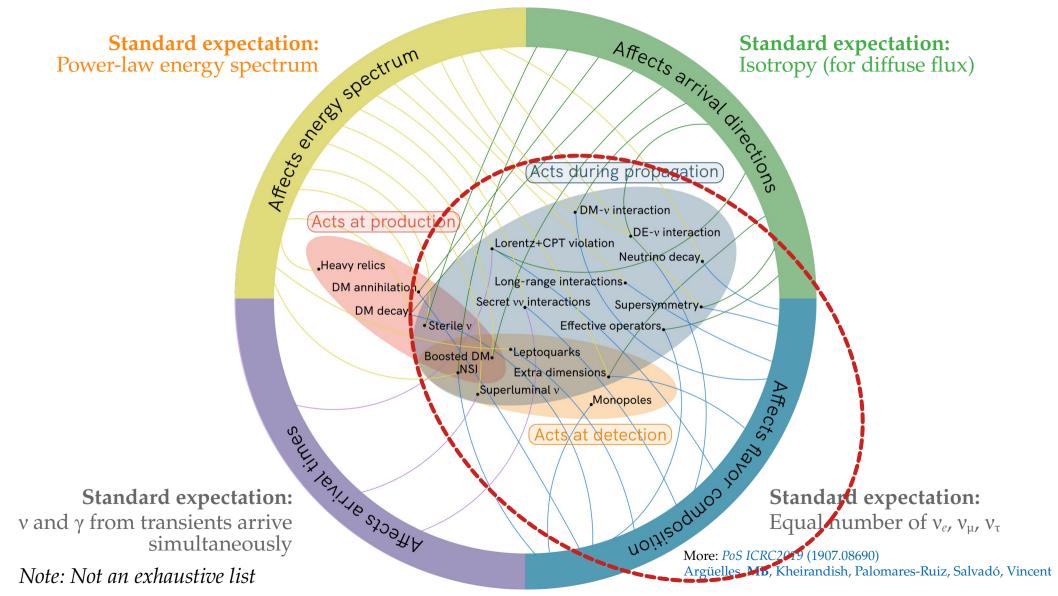
# Energy dependence of flavor ratios – in IceCube-Gen2 Measured: Inferred (at sources):







Repurpose the flavor sensitivity to test new physics:

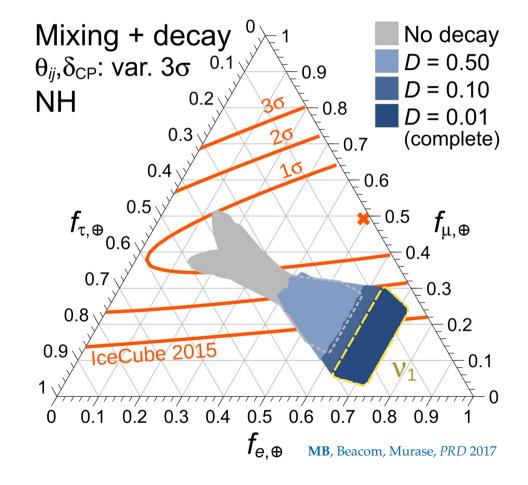


Repurpose the flavor sensitivity to test new physics:

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#### Neutrino decay

[Beacom *et al.*, *PRL* 2003; Baerwald, **MB**, Winter, JCAP 2010; **MB**, Beacom, Winter, *PRL* 2015; **MB**, Beacom, Murase, *PRD* 2017]



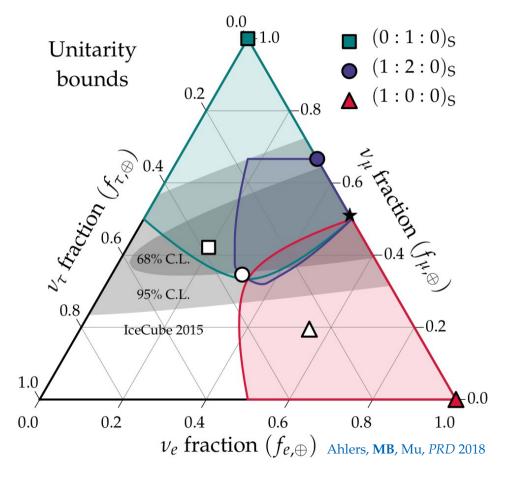
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[Xu, He, Rodejohann, *JCAP* 2014; Ahlers, **MB**, Mu, *PRD* 2018; Ahlers, **MB**, Nortvig, 2009.01253]



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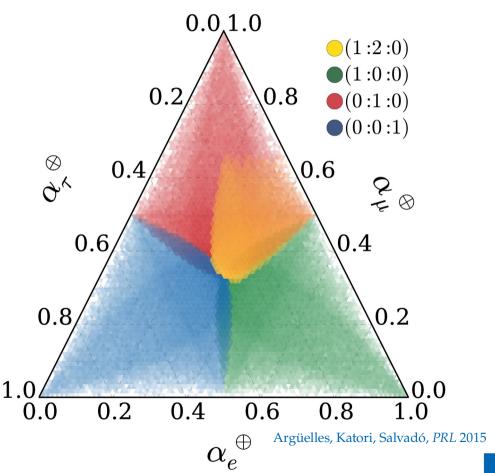
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Lorentz- and CPT-invariance violation

[Barenboim & Quigg, *PRD* 2003; **MB**, Gago, Peña-Garay, *JHEP* 2010; Kostelecky & Mewes 2004; Argüelles, Katori, Salvadó, *PRL* 2015]



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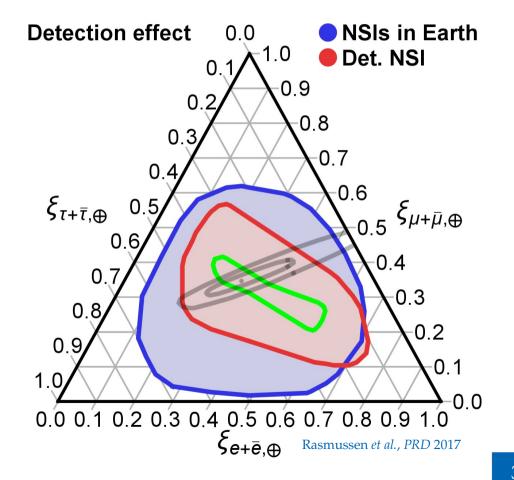
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[González-García *et al., Astropart. Phys.* 2016; Rasmussen *et al., PRD* 2017]



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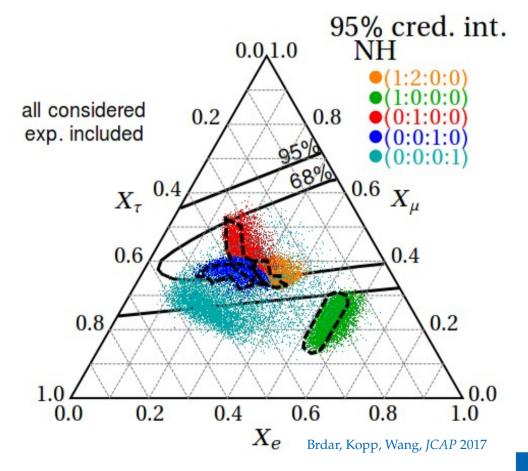
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#### Active-sterile v mixing

[Aeikens *et al., JCAP* 2015; Brdar, Kopp, Wang, *JCAP* 2017; Argüelles *et al., JCAP* 2020; Ahlers, **MB**, Nortvig, 2009.01253]



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[Xu, He, Rodejohann, *JCAP* 2014; Ahlers, **MB**, Mu, *PRD* 2018; Ahlers, **MB**, Nortvig, 2009.01253]

#### Lorentz- and CPT-invariance violation

[Barenboim & Quigg, *PRD* 2003; **MB**, Gago, Peña-Garay, *JHEP* 2010; Kostelecky & Mewes 2004; Argüelles, Katori, Salvadó, *PRL* 2015]

#### Non-standard interactions

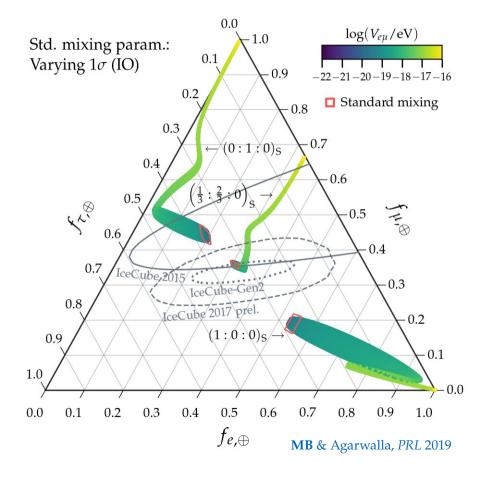
[González-García *et al., Astropart. Phys.* 2016; Rasmussen *et al., PRD* 2017]

#### Active-sterile v mixing

[Aeikens *et al.*, *JCAP* 2015; Brdar, Kopp, Wang, *JCAP* 2017; Argüelles *et al.*, *JCAP* 2020; Ahlers, **MB**, Nortvig, 2009.01253]

Long-range ev interactions [MB & Agarwalla, PRL 2019]

```
Reviews:
Mehta & Winter, JCAP 2011; Rasmussen et al., PRD 2017
```



# Assorted answerable questions

## *Measure flavor ratios better in water-based Cherenkov v telescopes?*

Neutron and muon echoes separate better showers from  $v_e$  and  $v_{\tau}$  [Li, MB, Beacom, PRL 2019] Might also help distinguish e.m. *vs.* hadronic showers

## Flavor ratios from point sources?

Need enough events, probably only doable with IceCube-Gen2 or a global v observatory (PLEvUM) Trigger with tracks, then look for cascades in the same direction

## *Flavor ratios of v vs.* $\bar{v}$ ?

Hard: cross sections and inelasticity distributions are very similar Glashow resonance helps – from 1 candidate:  $\bar{v}_e/v_e = 1$  [Lu, UHECR 2018],  $\bar{v}/v = 0.64 \pm 0.23$  [MB, 2004.06844]

## Can we measure the flavor composition of ultra-high-energy neutrinos?

Using in-ice radio (RNO-G, IceCube-Gen2): promising! [García-Fernández, Nelles, Glaser, PRD 2020] Other techniques (atmospheric radio, fluorescence, *etc.*): remains to be seen



imgflip.com

# Backup slides

## Flavor-transition probability: the quick and dirty of it

• In matrix form: 
$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1}^* & U_{e2}^* & U_{e3}^* \\ U_{\mu1}^* & U_{\mu2}^* & U_{\mu3}^* \\ U_{\tau1}^* & U_{\tau2}^* & U_{\tau3}^* \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

▶ Pontecorvo-Maki-Nakagawa-Sakata matrix ( $c_{ij} = \cos \theta_{ij}, s_{ij} = \sin \theta_{ij}$ ):

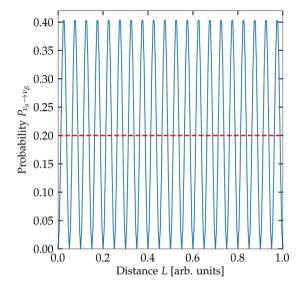
$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} e^{i\alpha_1/2} & 0 & 0 \\ 0 & e^{i\alpha_2/2} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$
  
Atmospheri  
Cross  
mixing  
Solar Majorana CP phases  
Probability for  $\mathbf{v}_{a} \rightarrow \mathbf{v}_{\beta}$ :  $P_{\nu_{\alpha} \rightarrow \nu_{\beta}} = \delta_{\alpha\beta} - 4 \sum_{i>j} \operatorname{Re}(U_{\alpha i}^{*}U_{\beta i}U_{\alpha j}U_{\beta j}^{*}) \sin^{2}\left(\Delta m_{ij}^{2}\frac{L}{4E}\right) + 2 \sum_{i>j} \operatorname{Im}(U_{\alpha i}^{*}U_{\beta i}U_{\alpha j}U_{\beta j}^{*}) \sin\left(\Delta m_{ij}^{2}\frac{L}{2E}\right)$   
Mauricio Bustamante (NBI)

## Flavor-transition probability: the quick and dirty of it

# ... But high-energy neutrinos oscillate *fast*

$$P_{\nu_{\alpha} \to \nu_{\beta}} = \delta_{\alpha\beta} - 4 \sum_{i>j} \operatorname{Re}(U_{\alpha i}^{*} U_{\beta i} U_{\alpha j} U_{\beta j}^{*}) \sin^{2} \left(\Delta m_{ij}^{2} \frac{L}{4E}\right) + 2 \sum_{i>j} \operatorname{Im}(U_{\alpha i}^{*} U_{\beta i} U_{\alpha j} U_{\beta j}^{*}) \sin\left(\Delta m_{ij}^{2} \frac{L}{2E}\right)$$

Oscillation length for 1-TeV v:  $2\pi \times 2E/\Delta m^2 \sim 0.1$  pc



~ 8% of the way to Proxima Centauri≪ Distance to Galactic Center (8 kpc)

- ≪ Distance to Andromeda (1 Mpc)
- ≪ Cosmological distances (few Gpc)

We cannot resolve oscillations, so we use instead the average probability:

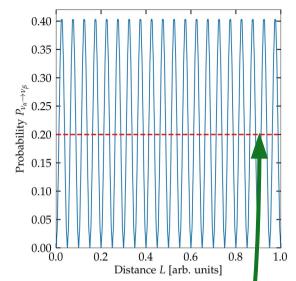
$$\langle P_{\nu_{\alpha} \to \nu_{\beta}} \rangle = \sum_{i=1}^{3} |U_{\alpha i}|^2 |U_{\beta i}|^2$$

Mauricio Bustamante (NBI)

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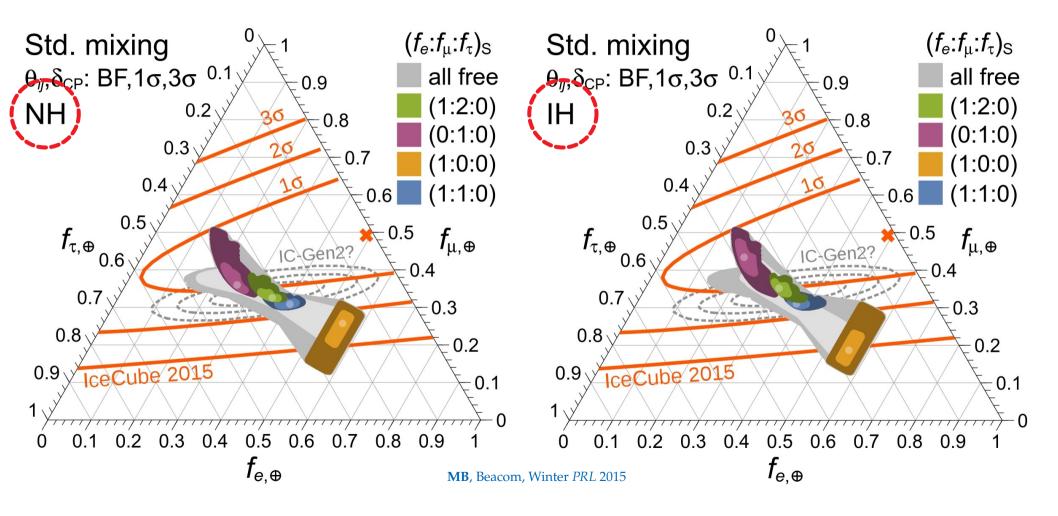
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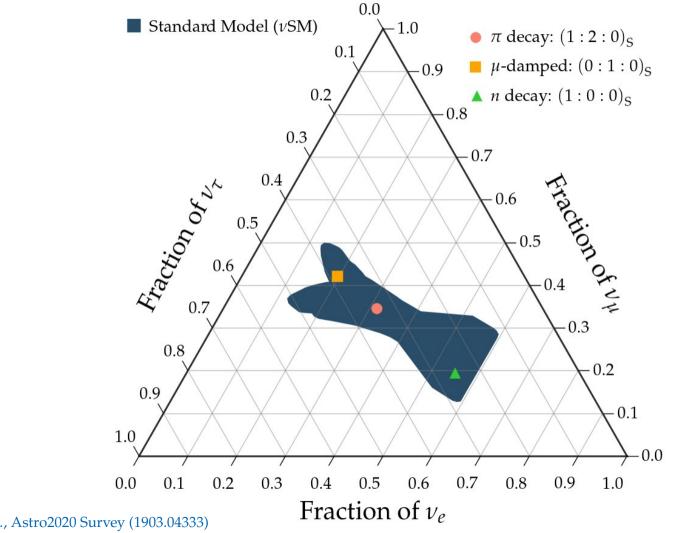
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Mauricio Bustamante (NBI)

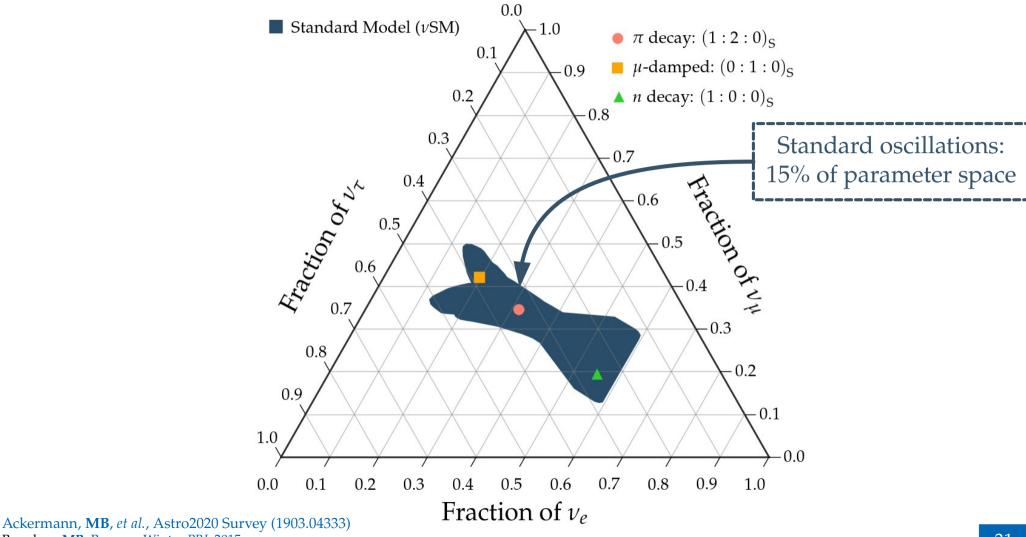
# Flavor composition – a few source choices

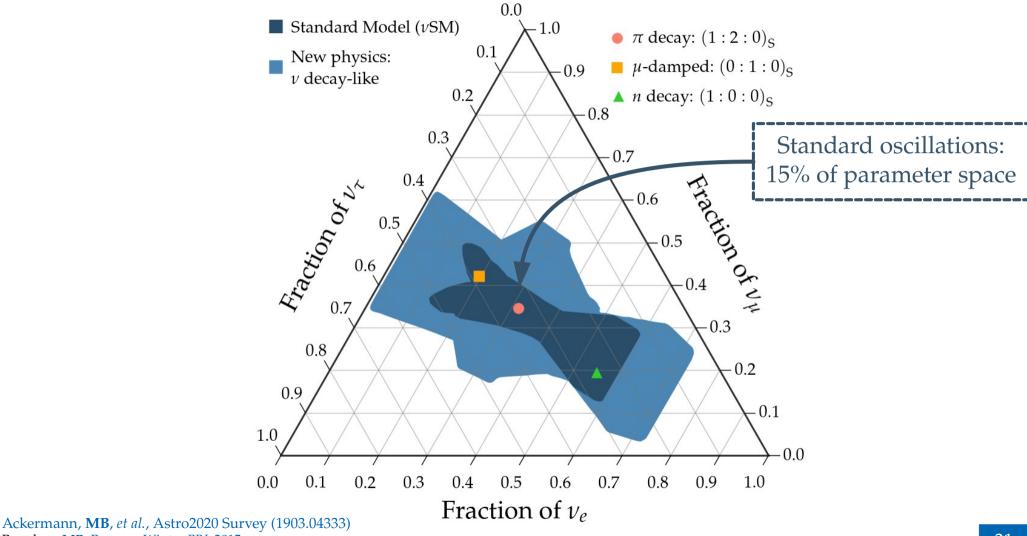
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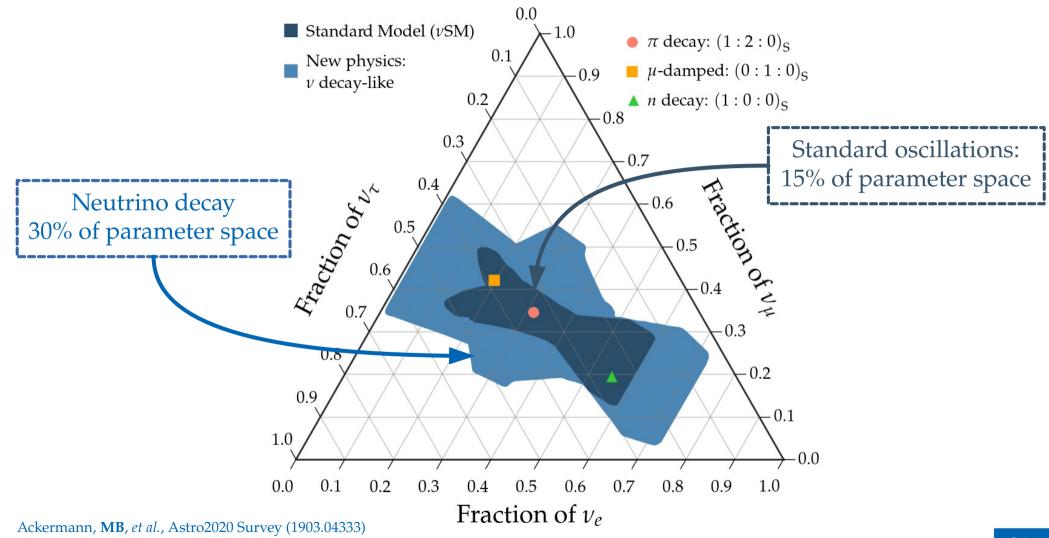


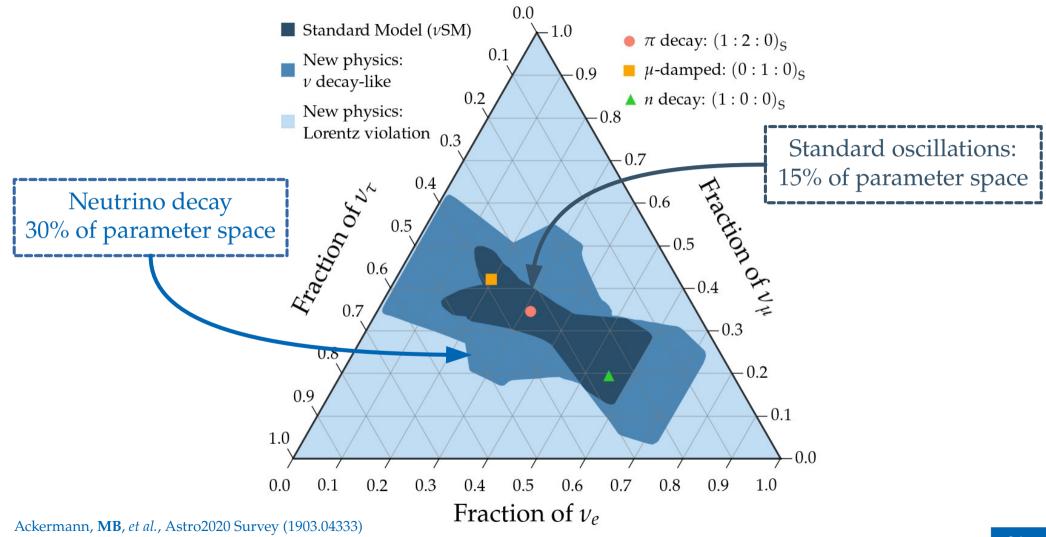


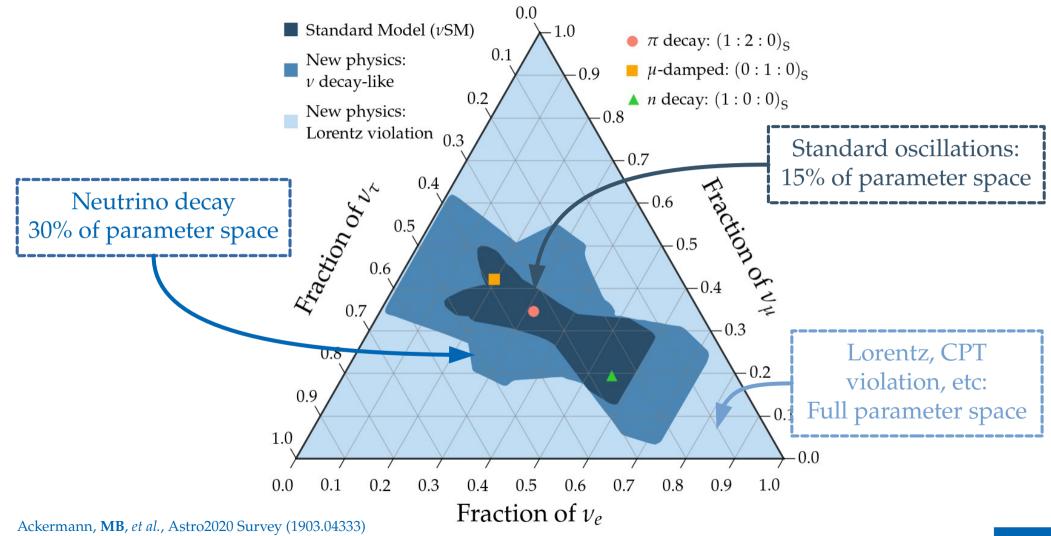
Ackermann, **MB**, *et al.*, Astro2020 Survey (1903.04333) Based on: **MB**, Beacom, Winter *PRL* 2015

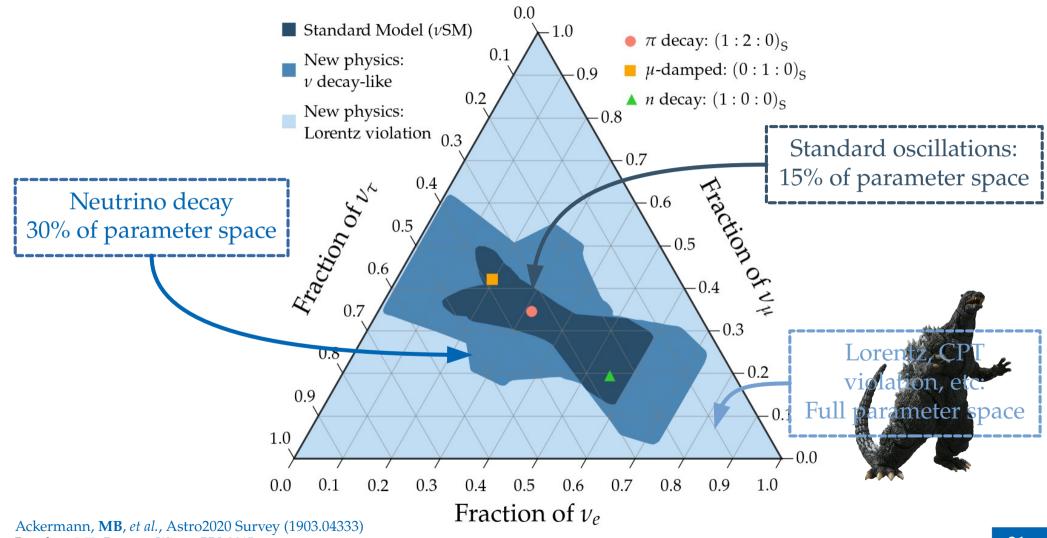






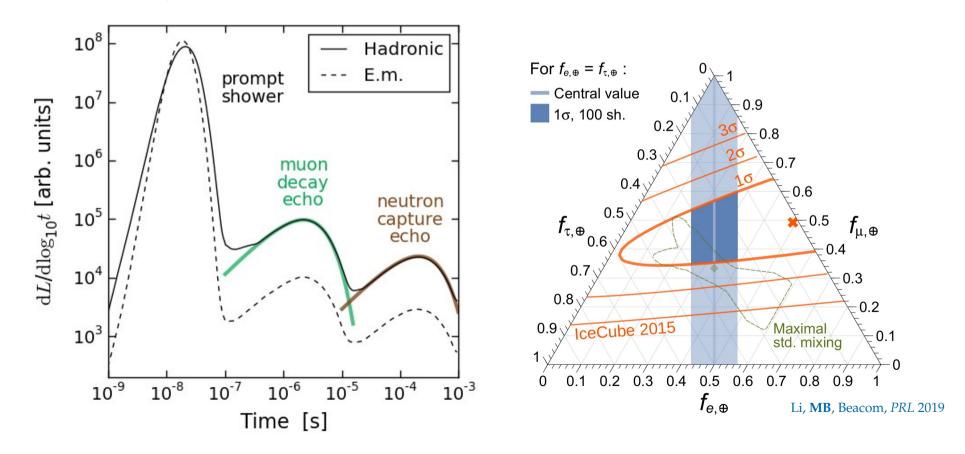






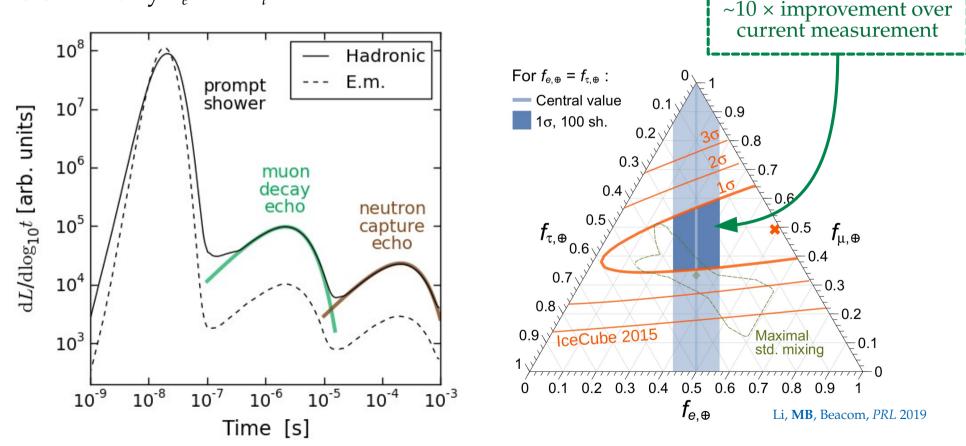
## Side note: Improving flavor-tagging using *echoes*

Late-time light (*echoes*) from muon decays and neutron captures can separate showers made by  $v_e$  and  $v_{\tau}$  –



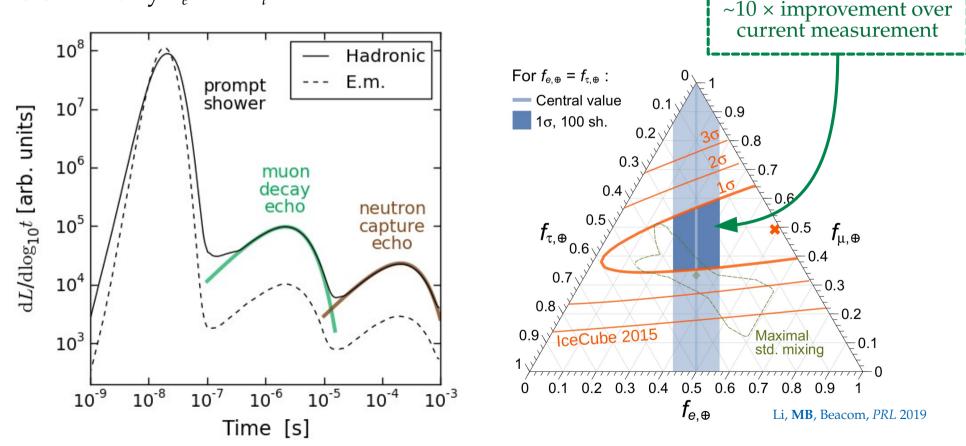
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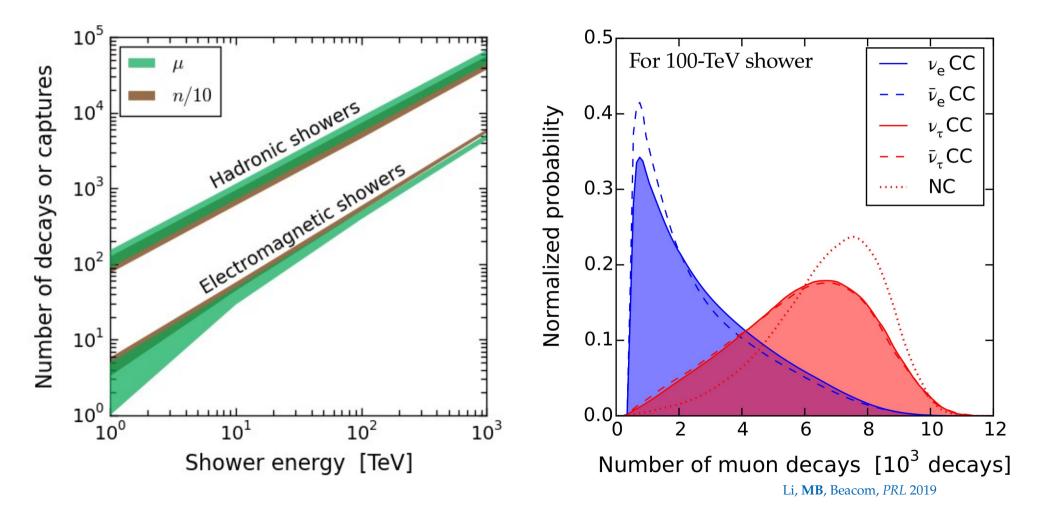


## Side note: Improving flavor-tagging using echoes

Late-time light (*echoes*) from muon decays and neutron captures can separate showers made by  $v_e$  and  $v_{\tau}$  –



### Hadronic vs. electromagnetic showers



### Are neutrinos forever?

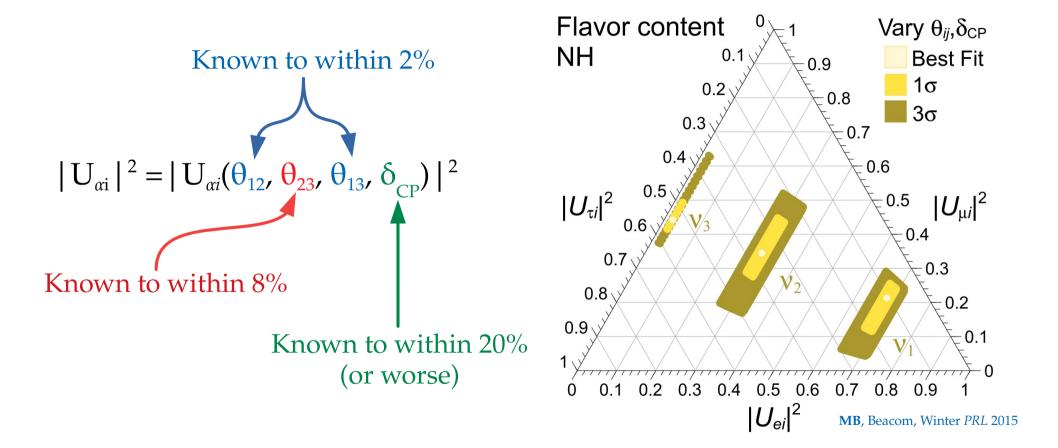
In the Standard Model (vSM), neutrinos are essentially stable (τ > 10<sup>36</sup> yr):
 One-photon decay (v<sub>i</sub> → v<sub>j</sub> + γ): τ > 10<sup>36</sup> (m<sub>i</sub>/eV)<sup>-5</sup> yr
 Two-photon decay (v<sub>i</sub> → v<sub>j</sub> + γ + γ): τ > 10<sup>57</sup> (m<sub>i</sub>/eV)<sup>-9</sup> yr
 Three-neutrino decay (v<sub>i</sub> → v<sub>i</sub> + v<sub>k</sub> + v<sub>k</sub>): τ > 10<sup>55</sup> (m<sub>i</sub>/eV)<sup>-5</sup> yr

► BSM decays may have significantly higher rates:  $v_i \rightarrow v_j + \phi$ 

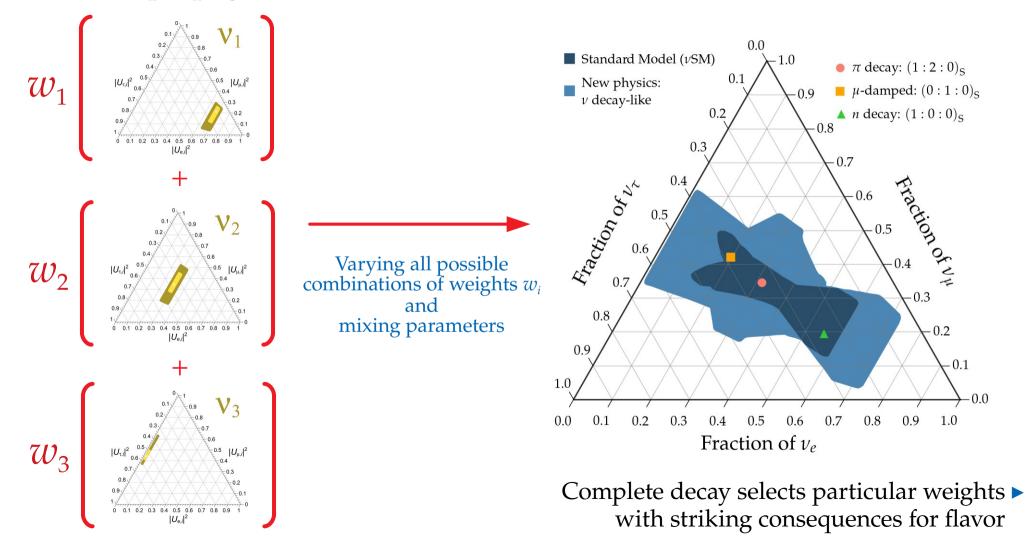
**φ**: Nambu-Goldstone boson of a broken symmetry (*e.g.*, Majoron)

We work in a model-independent way: the nature of φ is unimportant if it is invisible to neutrino detectors

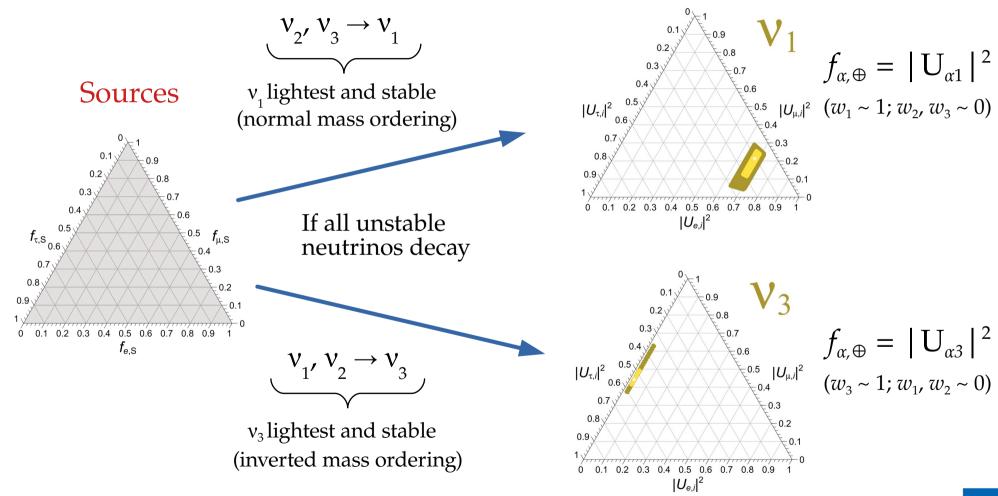
### Flavor content of neutrino mass eigenstates



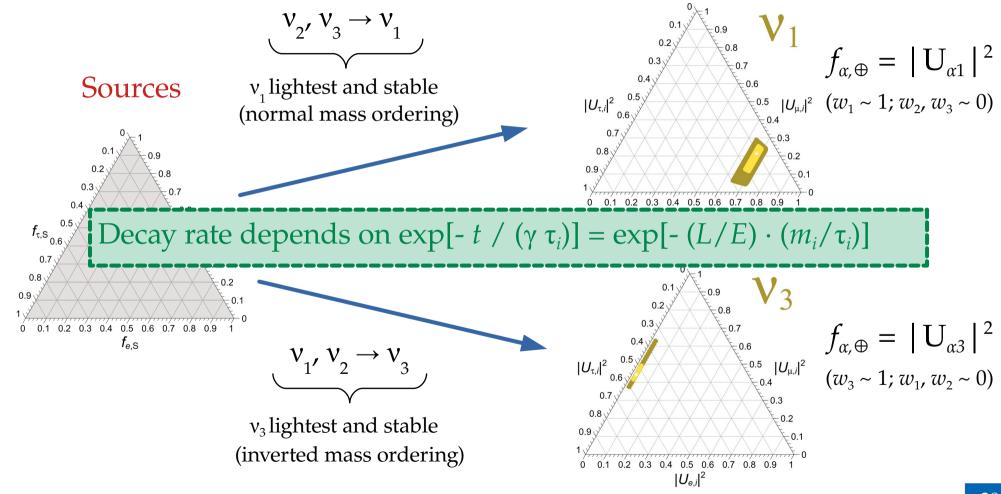
Neutrinos propagate as an incoherent mix of  $v_1$ ,  $v_2$ ,  $v_3$  —

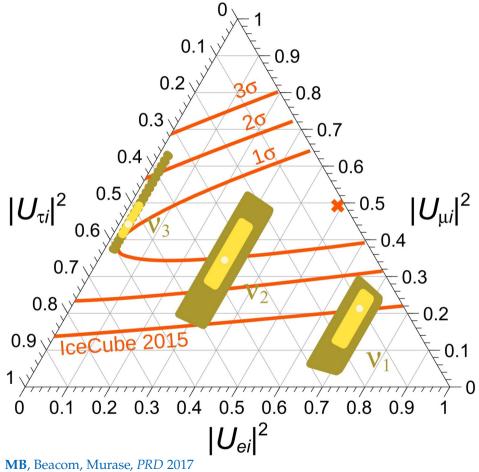


### Measuring the neutrino lifetime

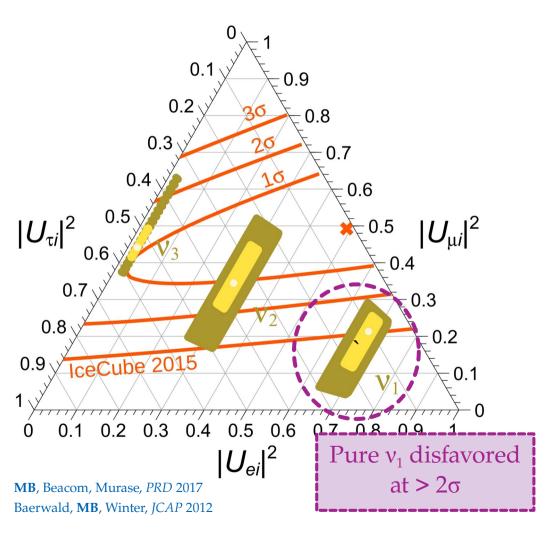


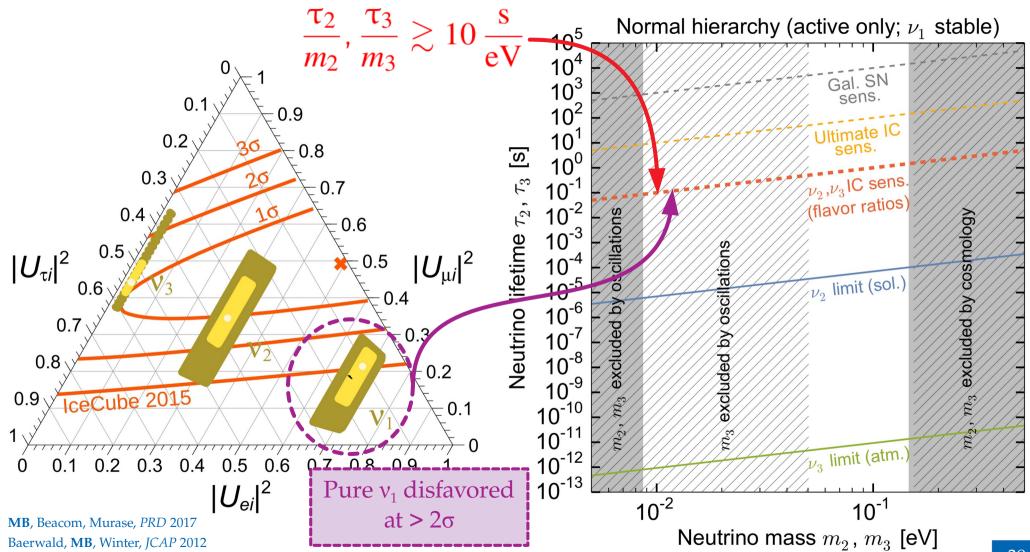
### Measuring the neutrino lifetime





Baerwald, **MB**, Winter, *JCAP* 2012





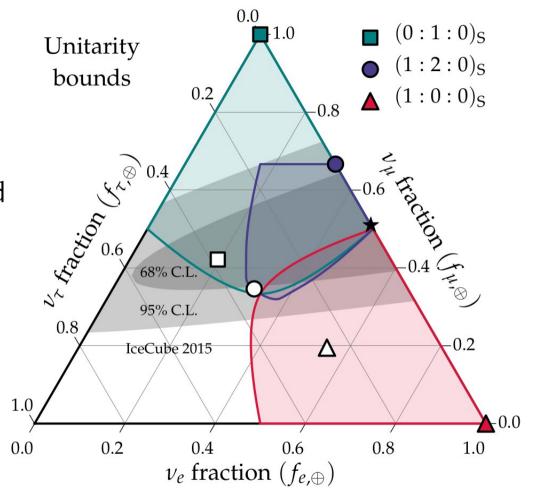
## Using unitarity to constrain new physics

 $H_{tot} = H_{std} + H_{NP}$ 

New mixing angles unconstrained

- Use unitarity  $(U_{NP}U_{NP}^{\dagger} = 1)$  to bound all possible flavor ratios at Earth
- Can be used as prior in new-physics searches in IceCube

Ahlers, **MB**, Mu, *PRD* 2018 See also: Xu, He, Rodejohann, *JCAP* 2014



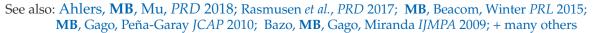
## How to fill out the flavor triangle?

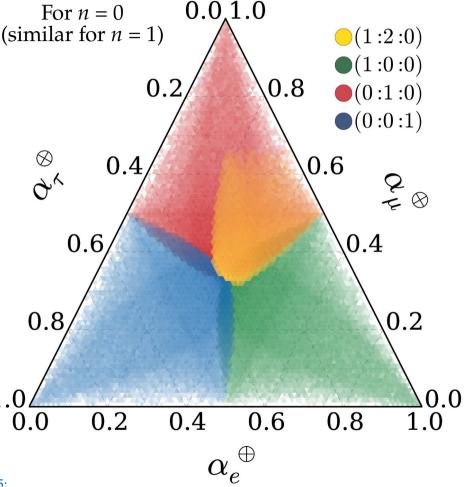
 $H_{\text{tot}} = H_{\text{std}} + H_{\text{NP}} \qquad (I$   $H_{\text{std}} = \frac{1}{2E} U_{\text{PMNS}}^{\dagger} \operatorname{diag} \left(0, \Delta m_{21}^{2}, \Delta m_{31}^{2}\right) U_{\text{PMNS}} \qquad (I$   $H_{\text{NP}} = \sum_{n} \left(\frac{E}{\Lambda_{n}}\right)^{n} U_{n}^{\dagger} \operatorname{diag} \left(O_{n,1}, O_{n,2}, O_{n,3}\right) U_{n}$ 

This can populate *all* of the triangle –

► Use current atmospheric bounds on  $O_{n,i}$ :  $O_0 < 10^{-23}$  GeV,  $O_1/\Lambda_1 < 10^{-27}$  GeV

Sample the unknown new mixing angles





Argüelles, Katori, Salvadó, PRL 2015

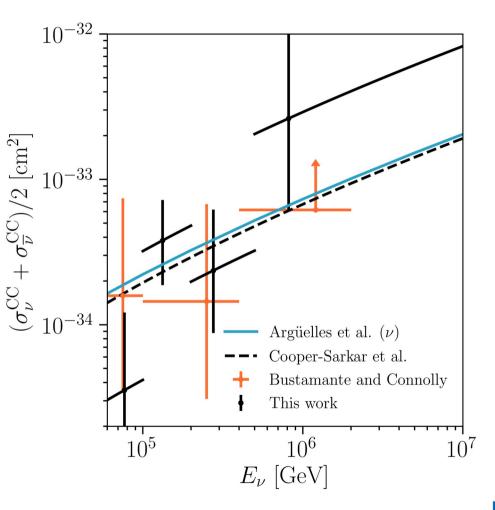
## How to fill out the flavor triangle?

0.0.1.0 For n = 0 $H_{\text{tot}} = H_{\text{std}} + H_{\text{NP}}$ (similar for n = 1) (1:2:0)(1:0:0) $H_{\text{std}} = \frac{1}{2F} U_{\text{PMNS}}^{\dagger} \operatorname{diag} \left(0, \Delta m_{21}^2, \Delta m_{31}^2\right) U_{\text{PMNS}}$ 8.0 ().2(0:1:0)(0:0:1) $H_{\rm NP} = \sum \left(\frac{E}{\Lambda_n}\right)^n U_n^{\dagger} \operatorname{diag}\left(O_{n,1}, O_{n,2}, O_{n,3}\right) U_n$ 0.4 0.6 Q E Ø This can populate *all* of the triangle – 0.6 0.4• Use current atmospheric bounds on  $O_{n,i}$ :  $O_0 < 10^{-23}$  GeV,  $O_1/\Lambda_1 < 10^{-27}$  GeV 0.8 0.2Sample the unknown new mixing angles 0.00.2 0.40.60.8 0.0 1.0 $lpha_{e}^{\,\oplus}$ See also: Ahlers, MB, Mu, PRD 2018; Rasmusen et al., PRD 2017; MB, Beacom, Winter PRL 2015;

**MB**, Gago, Peña-Garay *JCAP* 2010; Bazo, **MB**, Gago, Miranda *IJMPA* 2009; + many others

### *New (IC 7.5 yr):* Updated cross section measurement

- ► Uses 7.5 years of IceCube data
- Uses starting showers + tracks
  - Vs. starting showers only in Bustamante & Connolly 2017
  - ► *Vs.* throughoing muons in IceCube 2017
- Extends measurement to 10 PeV
- Still compatible with Standard Model predictions
- Higher energies? Work in progress by Valera & Bustamante



#### Earth

Galactic (kpc) or extragalactic (Mpc – Gpc) distance

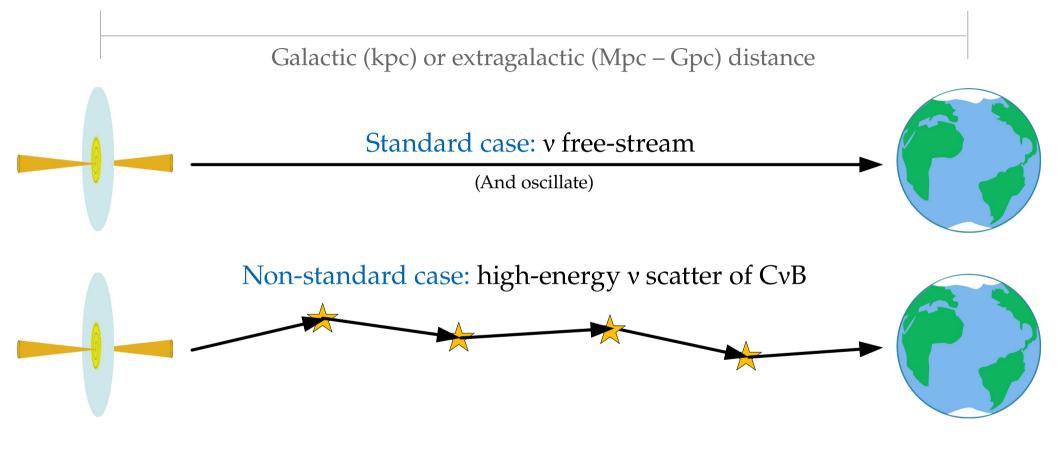
#### Earth

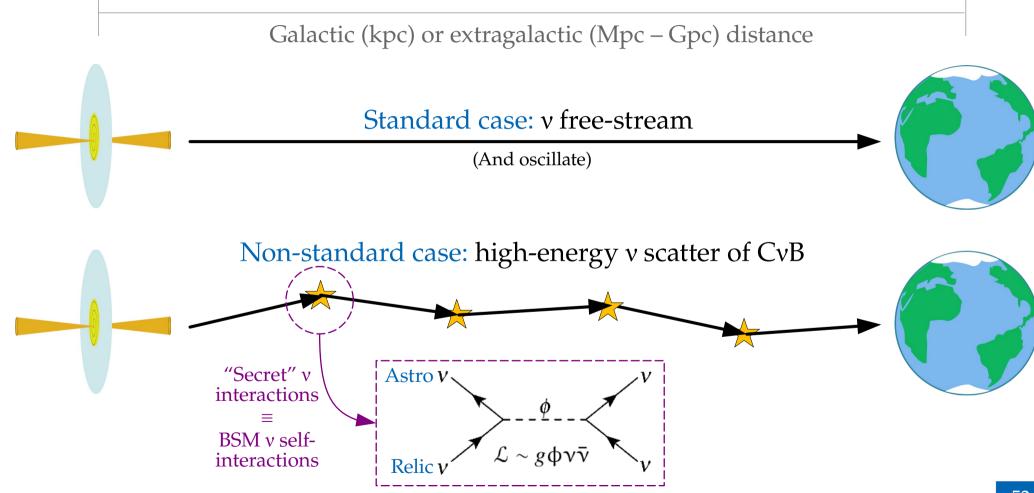
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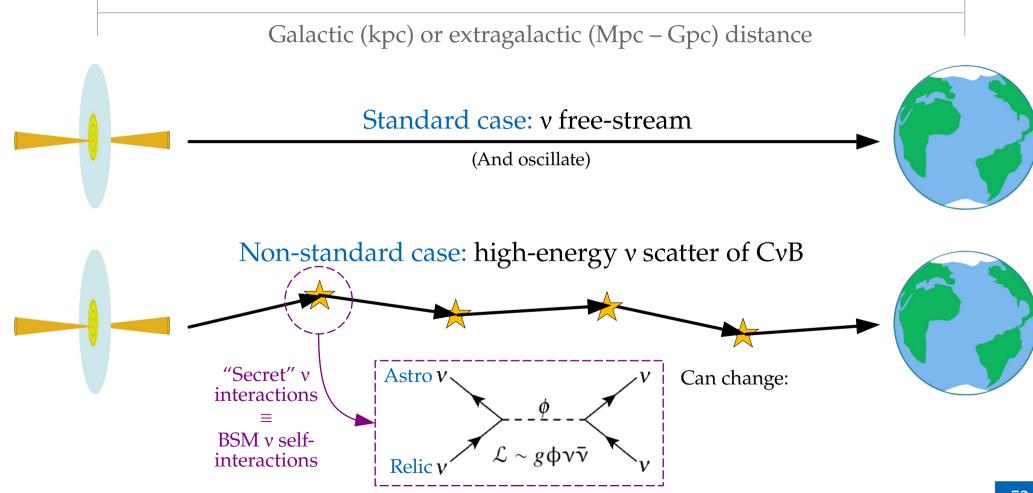
Standard case: v free-stream

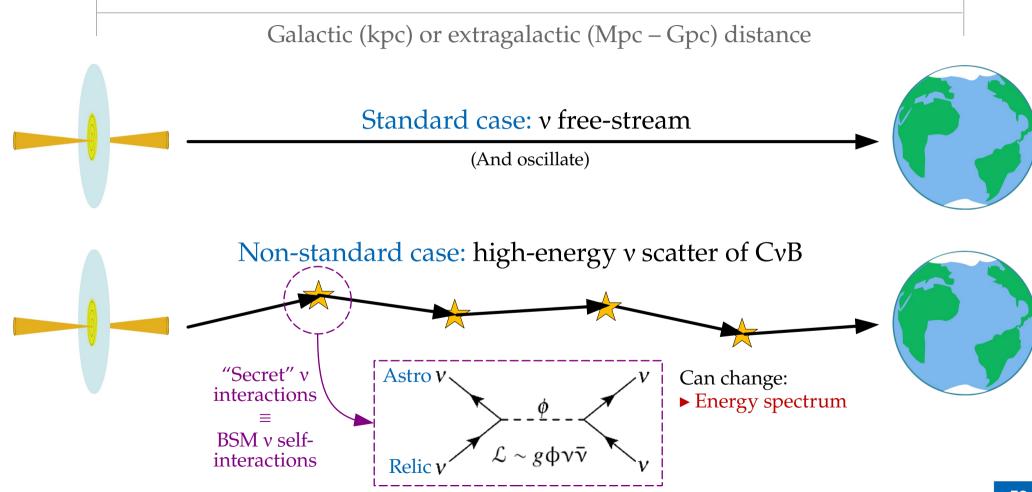
(And oscillate)

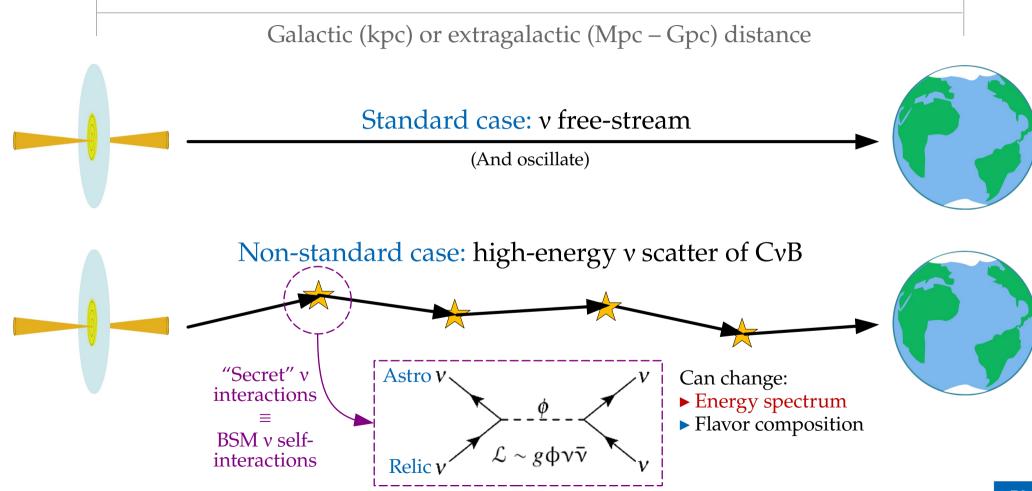


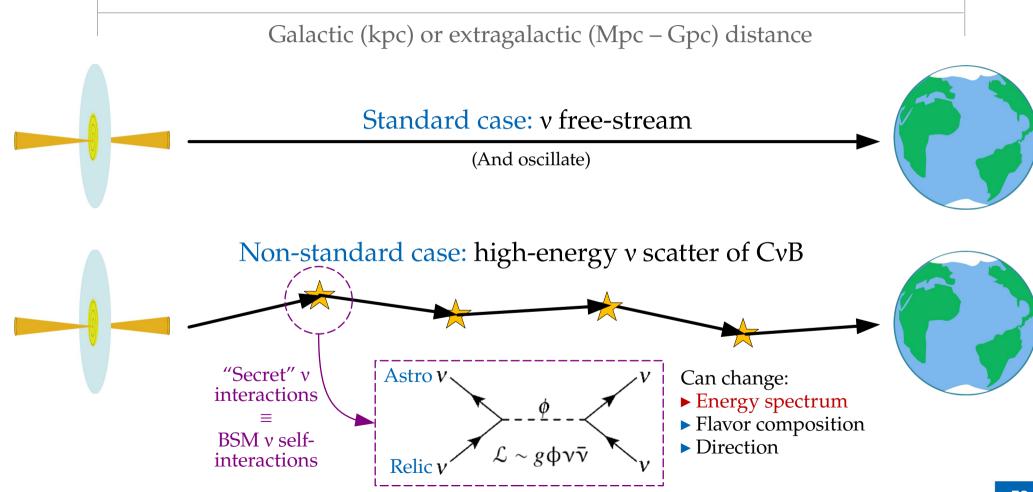


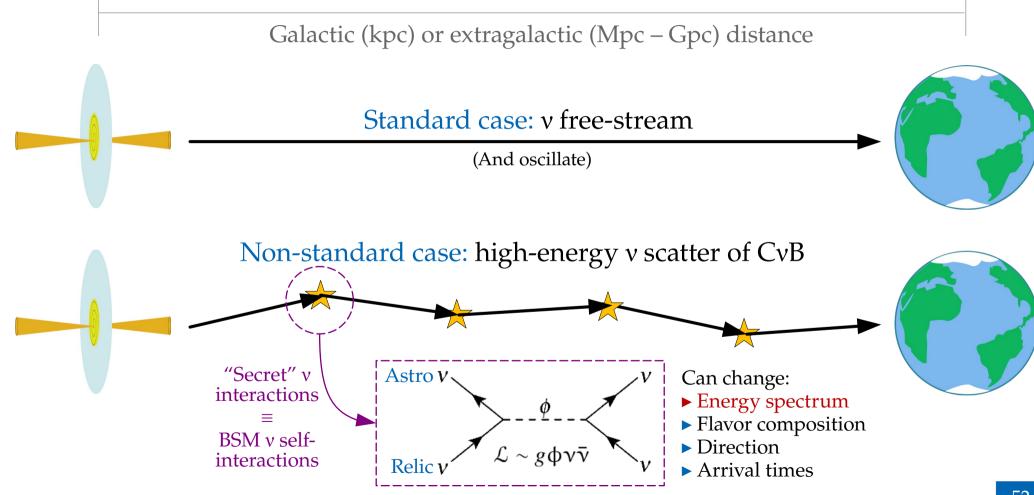












## Fundamental physics with HE cosmic neutrinos

Numerous new-physics effects grow as ~  $\kappa_n \cdot E^n \cdot L$ 

So we can probe  $\kappa_n \sim 4 \cdot 10^{-47} \, (E/PeV)^{-n} \, (L/Gpc)^{-1} \, PeV^{1-n}$ 

► Improvement over limits using atmospheric v:  $\kappa_0 < 10^{-29}$  PeV,  $\kappa_1 < 10^{-33}$ 

Fundamental physics can be extracted from four neutrino observables:

- Spectral shape
- Angular distribution
- Flavor composition
- Timing

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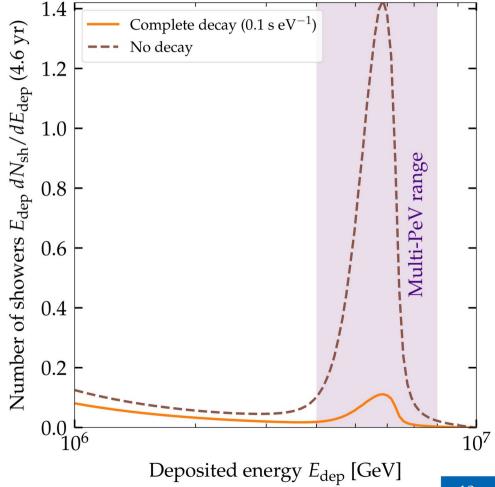
► Spectral shape

► Timing

 Angular distribution
 Flavor composition
 Timing & astrophysical unknowns

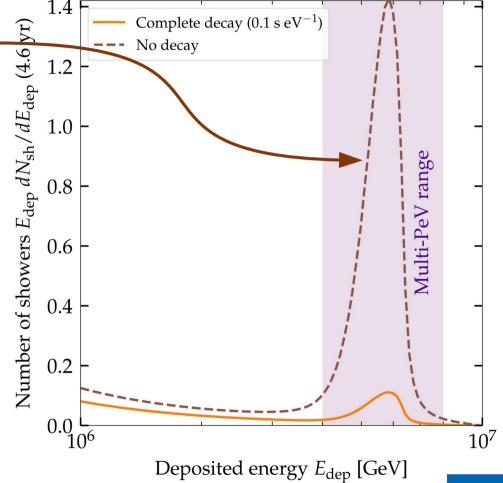
At 6.3 PeV, the Glashow resonance  $(v_e + e \rightarrow W)$  should trigger showers in IceCube

- ... unless v<sub>1</sub>, v<sub>2</sub> decay to v<sub>3</sub> en route to Earth (the surviving v<sub>3</sub> have little electron content)
- IceCube has seen 1 shower in the 4–8 PeV range, so v<sub>1</sub>, v<sub>2</sub> must make it to Earth
- So we set *lower* limits on their lifetimes (in the inverted mass ordering)
- Translated into *upper* limits on coupling



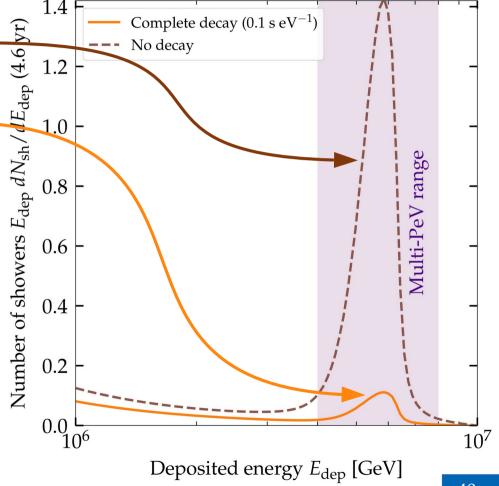
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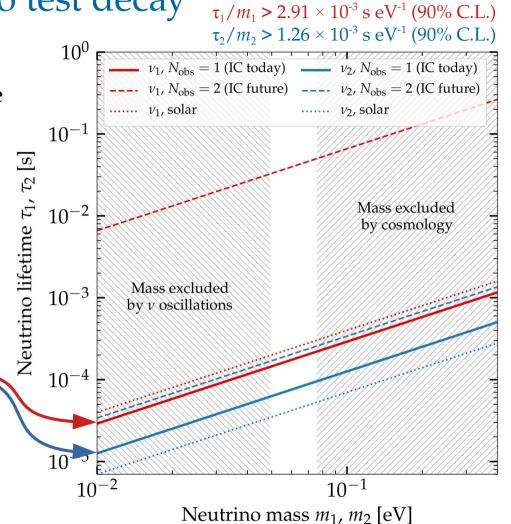


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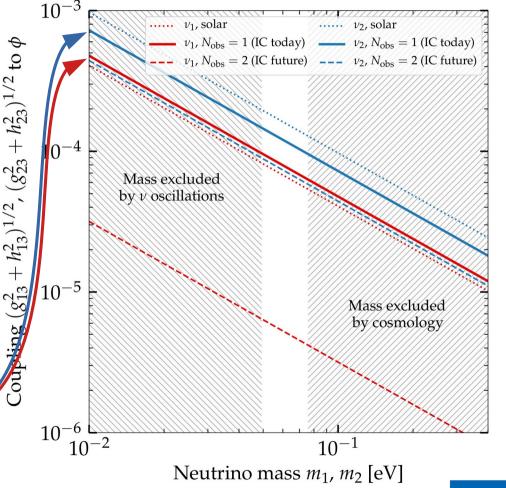
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► Translated into *upper* limits on coupling •  $\mathcal{L} = g_{ij}\bar{\nu}_i\nu_j\phi + h_{ij}\bar{\nu}_j\gamma_5\nu_j\phi + h.c.$ 



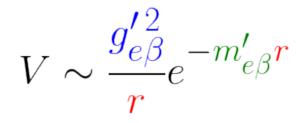
## Ultra-long-range flavorful interactions

- ► Simple extension of the SM: Promote the global lepton-number symmetries  $L_e$ - $L_\mu$ ,  $L_e$ - $L_\tau$  to local symmetries
- They introduce new interaction between electrons and v<sub>e</sub> and v<sub>μ</sub> or v<sub>τ</sub> mediated by a new neutral vector boson (Z'):
  - Affects oscillations
  - ► If the *Z*′ is *very* light, *many* electrons can contribute

X.-G. He, G.C. Joshi, H. Lew, R. R. Volkas, *PRD* 1991 / R. Foot, X.-G. He, H. Lew, R. R. Volkas, *PRD*A. Joshipura, S. Mohanty, *PLB* 2004 / J. Grifols & E. Massó, *PLB* 2004 / A. Bandyopadhyay, A. Dighe, A. Joshipura, *PRD*M.C. González-García, P..C. de Holanda, E. Massó, R. Zukanovich Funchal, *JCAP* 2007 / A. Samanta, *JCAP*S.-S. Chatterjee, A. Dasgupta, S. Agarwalla, *JHEP*

### The new potential sourced by an electron

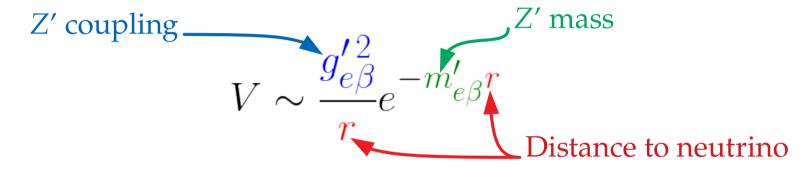
Under the  $L_e$ - $L_\mu$  or  $L_e$ - $L_\tau$  symmetry, an electron sources a Yukawa potential —



A neutrino "feels" all the electrons within the interaction range  $\sim (1/m')$ 

### The new potential sourced by an electron

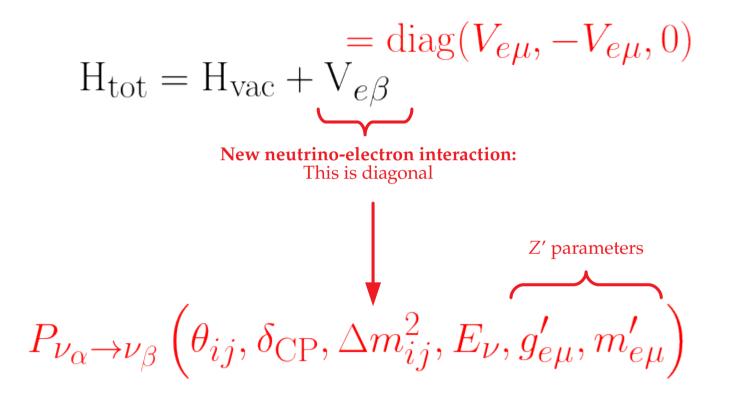
Under the  $L_e$ - $L_\mu$  or  $L_e$ - $L_\tau$  symmetry, an electron sources a Yukawa potential —

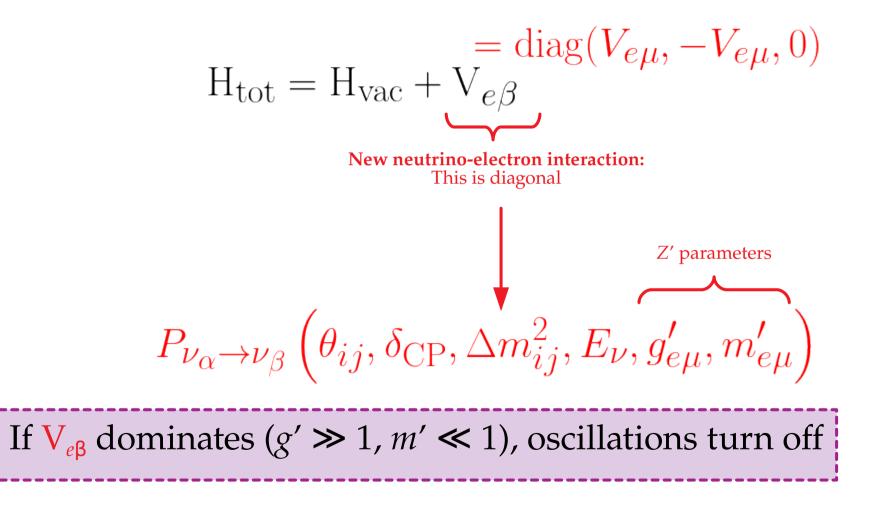


A neutrino "feels" all the electrons within the interaction range  $\sim (1/m')$ 

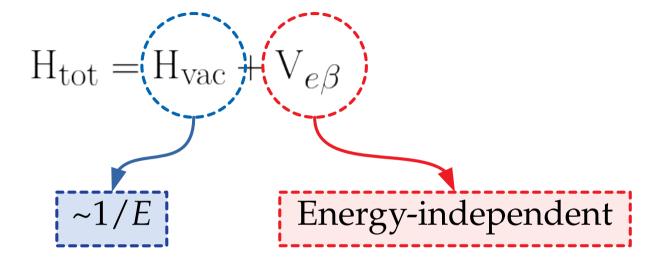
 $H_{tot} = H_{vac}$ Standard oscillations: Neutrinos change flavor because this is non-diagonal

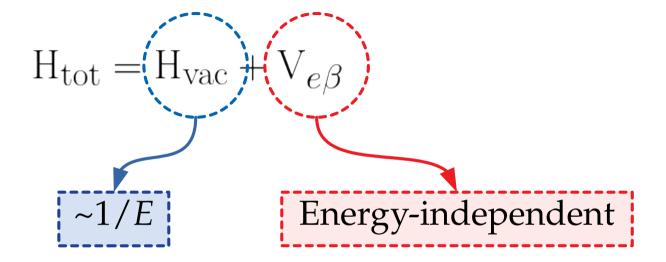
 $H_{tot} = H_{vac}$ Standard oscillations: Neutrinos change flavor because this is non-diagonal  $P_{\nu_{\alpha} \to \nu_{\beta}} \left( \theta_{ij}, \delta_{\rm CP} \right)$ 



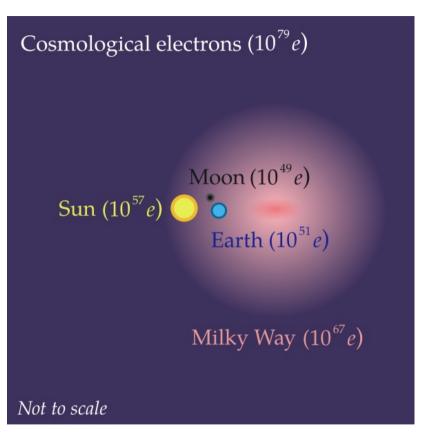


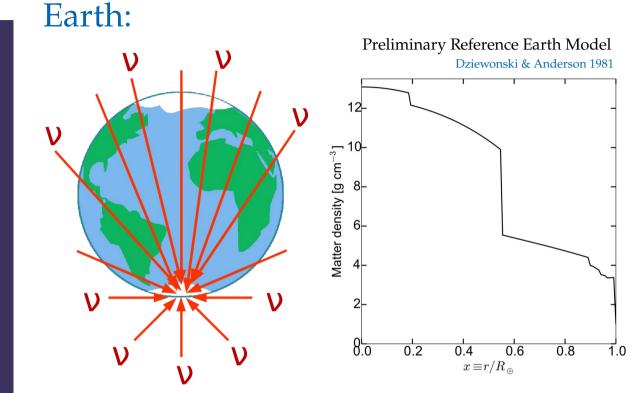
$$H_{tot} = H_{vac} + V_{e\beta}$$



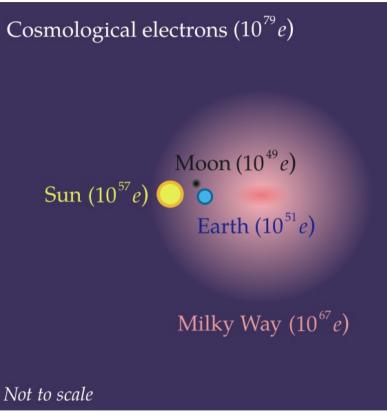


: We can use high-energy astrophysical neutrinos

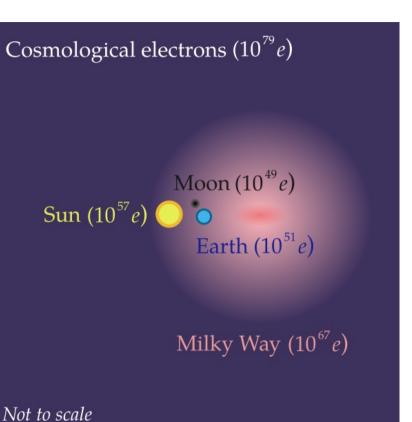




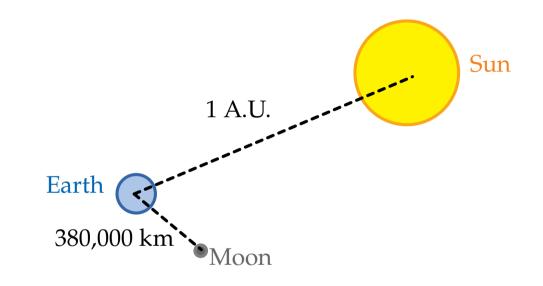
Neutrinos traverse different electron column depths



$$V_{e\beta} = V_{e\beta}^{\oplus}$$

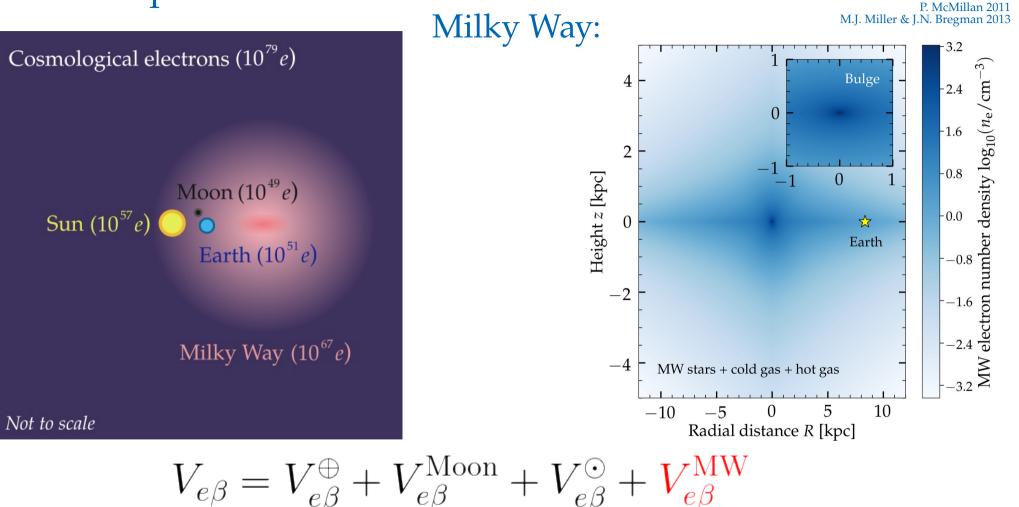


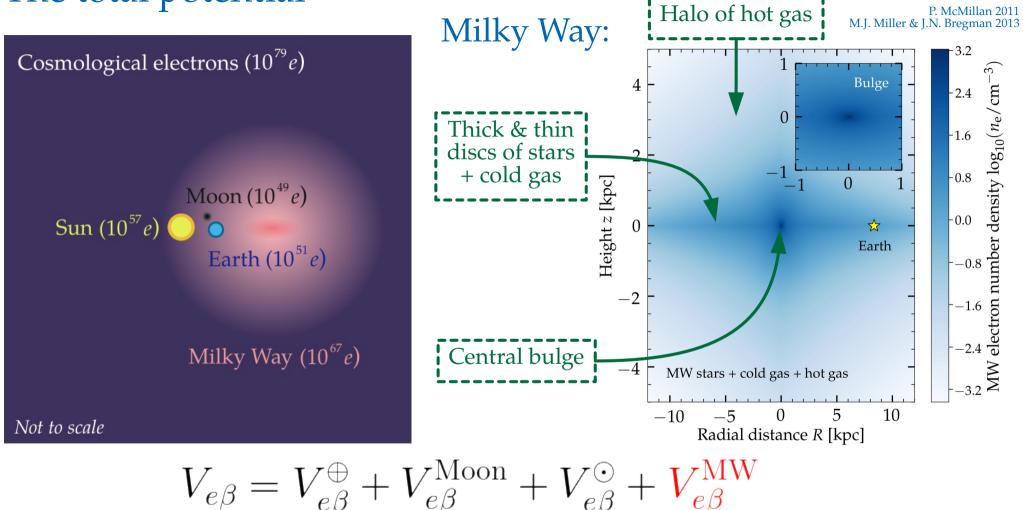
## Moon and Sun:

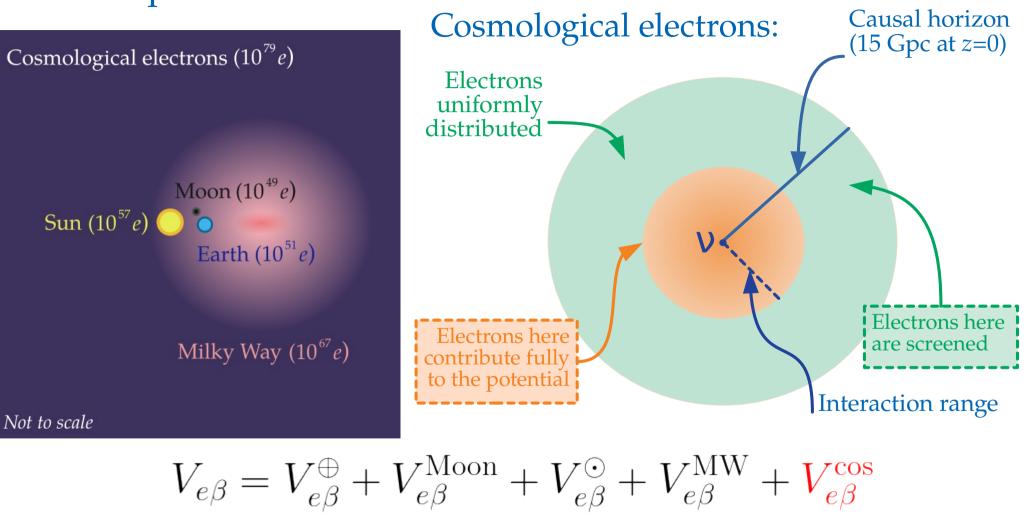


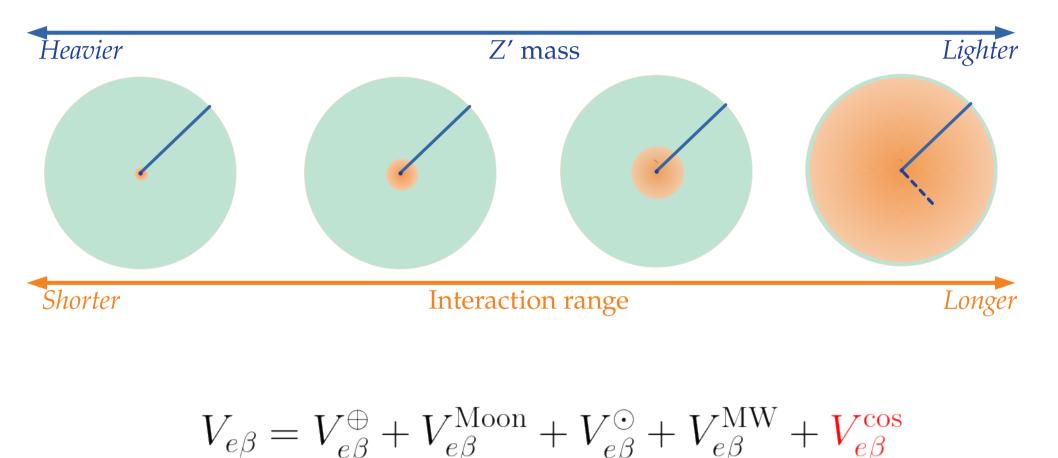
#### Treated as point sources of electrons

 $V_{e\beta} = V_{e\beta}^{\oplus} + V_{e\beta}^{\text{Moon}} + V_{e\beta}^{\odot}$ 









Potential:

$$V_{e\beta} \propto \frac{1}{r} e^{-m'_{e\beta}r}$$

Potential:

