

Tree Level Gauge Mediation

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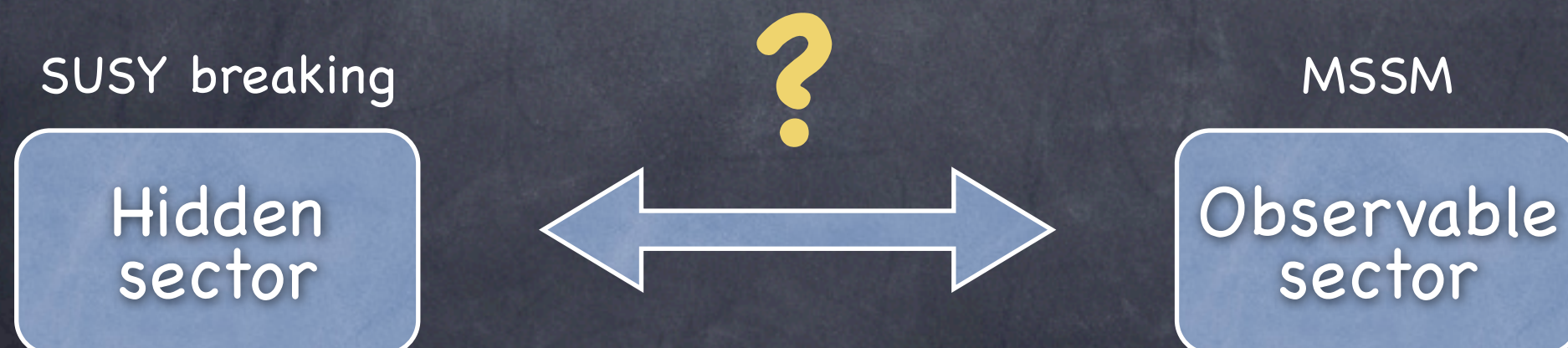
with A. Romanino and R. Ziegler
arXiv:0912.5482 (JHEP)
arXiv:0909.3058 (JHEP)

MSSM and SUSY Breaking

- The MSSM, despite its name, is not supersymmetric: $\mathcal{L}^{\text{MSSM}} = \mathcal{L}^{\text{MSSM}}_{\text{susy}} + \mathcal{L}^{\text{MSSM}}_{\text{soft}}$

$$\begin{aligned} \mathcal{L}^{\text{MSSM}}_{\text{soft}} = & -\frac{1}{2} \left(M_3 \tilde{g} \tilde{g} + M_2 \tilde{W} \tilde{W} + M_1 \tilde{B} \tilde{B} + \text{c.c.} \right) \\ & - \left(\tilde{u} \mathbf{a}_u \tilde{Q} H_u - \tilde{d} \mathbf{a}_d \tilde{Q} H_d - \tilde{e} \mathbf{a}_e \tilde{L} H_d + \text{c.c.} \right) \\ & - \tilde{Q}^\dagger \mathbf{m}_Q^2 \tilde{Q} - \tilde{L}^\dagger \mathbf{m}_L^2 \tilde{L} - \tilde{u} \mathbf{m}_u^2 \tilde{u}^\dagger - \tilde{d} \mathbf{m}_d^2 \tilde{d}^\dagger - \tilde{e} \mathbf{m}_e^2 \tilde{e}^\dagger \\ & - m_{H_u}^2 H_u^* H_u - m_{H_d}^2 H_d^* H_d - (b H_u H_d + \text{c.c.}) \end{aligned}$$

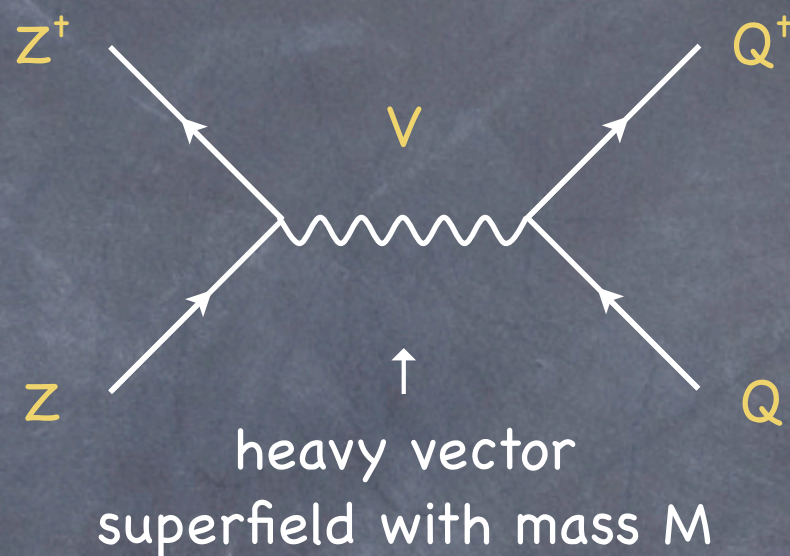
- Around 100 new physical parameters, which can give rise to FCNC and CP violation
- The soft sector keeps into account the unknown dynamics of the SUSY breaking and its mediation to the SM superfields



Tree Level Gauge Mediation

- In Tree Level Gauge Mediation the supersymmetry breaking is communicated to the scalars of the MSSM through TREE level gauge interactions:

Z chiral superfield
 $\langle Z \rangle = F\theta^2$
 $F \gg (M_Z)^2$
 SM singlet



$$\int d^4\theta \frac{Z^\dagger Z Q^\dagger Q}{M^2}$$

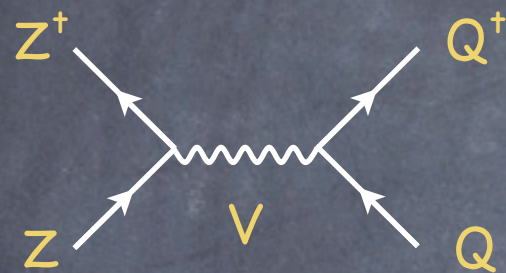
$$\rightarrow m^2 \tilde{Q}^\dagger \tilde{Q}, \quad m^2 = \frac{F^2}{M^2}$$

- Two arguments seem to prevent this possibility: **supertrace formula** and **gaugino masses**
- To satisfy the constraints imposed by the supertrace formula we need two ingredients: an extra U(1) gauge interactions and heavy d.o.f. which receive negative supersymmetry breaking contributions
- Gaugino Masses arise at 1-loop through the usual loop gauge mediation, the expected hierarchy between sfermion and gaugino masses can be reduced through model dependent factors

A concrete example

- $G = SO(10)$ "minimal" GUT (V heavy SM singlet means rank ≥ 5)

- V associated to the $SU(5)$ -invariant generator "X"



gives

$$\tilde{m}_Q^2 \propto X_Q X_Z$$

- The (usual) embedding of a MSSM family in a single 16 does not work (whatever the sign of X_Z)
- The three MSSM families are embedded in $16_i + 10_i$, $i=1,2,3$ (needs $X_Z > 0$)

SO(10)	SU(5)		
16_i	$= 5_i +$	$\boxed{10_i} +$	1_i
X	-3	1	5

SO(10)	SU(5)	
10_i	$= \boxed{5_i} +$	5_i
X	2	-2

- Prediction: $\tilde{m}_q^2 = \tilde{m}_{u^c}^2 = \tilde{m}_{e^c}^2 = \tilde{m}_{10}^2 = \frac{1}{10} m^2$, $\tilde{m}_l^2 = \tilde{m}_{d^c}^2 = \tilde{m}_{\bar{5}}^2 = \frac{1}{5} m^2$, $m = \frac{F}{M}$

An example of spectrum

Higgs:	m_{h^0}	114
	m_{H^0}	1543
	m_A	1543
	m_{H^\pm}	1545
Gluinos:	$M_{\tilde{g}}$	448
Neutralinos:	$m_{\chi_1^0}$	62
	$m_{\chi_2^0}$	124
	$m_{\chi_3^0}$	1414
	$m_{\chi_4^0}$	1415
Charginos:	$m_{\chi_1^\pm}$	124
	$m_{\chi_2^\pm}$	1416
Squarks:	$m_{\tilde{u}_L}$	1092
	$m_{\tilde{u}_R}$	1027
	$m_{\tilde{d}_L}$	1095
	$m_{\tilde{d}_R}$	1494
	$m_{\tilde{t}_1}$	1007
	$m_{\tilde{t}_2}$	1038
	$m_{\tilde{b}_1}$	1069
	$m_{\tilde{b}_2}$	1435
Sleptons:	$m_{\tilde{e}_L}$	1420
	$m_{\tilde{e}_R}$	1091
	$m_{\tilde{\tau}_1}$	992
	$m_{\tilde{\tau}_2}$	1387
	$m_{\tilde{\nu}_e}$	1418
	$m_{\tilde{\nu}_\tau}$	1382

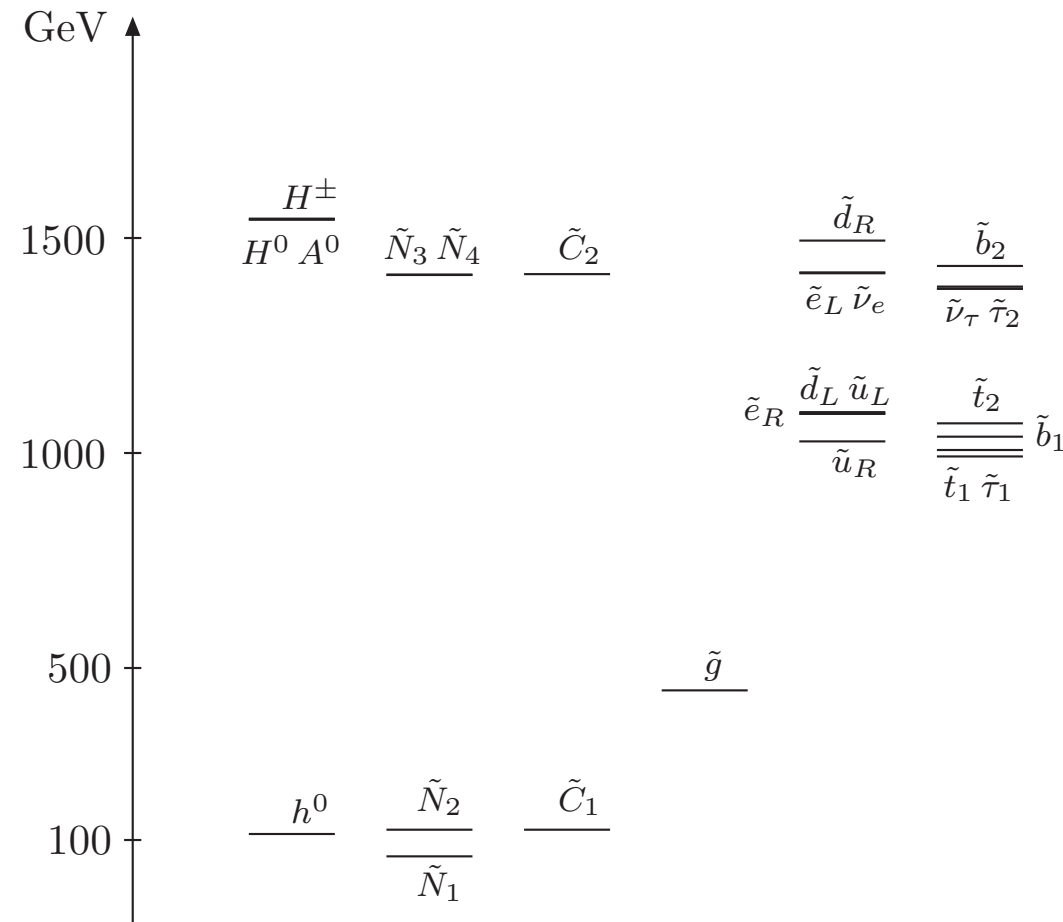


Figure 2: An example of spectrum, corresponding to $m = 3.2 \text{ TeV}$, $M_{1/2} = 150 \text{ GeV}$, $\theta_d = \pi/6$, $\tan \beta = 30$ and $\text{sign}(\mu) = +$, $A = 0$, $\eta = 1$. All the masses are in GeV, the first two families have an approximately equal mass.

Conclusions

- Tree level gauge mediation is one simple way to communicate supersymmetry breaking: through the tree level renormalizable exchange of a heavy gauge messenger
- This possibility is viable, despite the well known arguments associated to the supertrace formula
- It offers new model building avenues (including μ -problem)
- It solves the susy flavour problem
- It leads to a description of supersymmetry breaking in terms of few parameters, with peculiar relations among sfermion masses that can be tested at the LHC