

# Universality in many-body systems

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# On the point

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## What are we talking about:

- + Structure of “many”-particle non relativistic systems (e.g., nucleons).
- + Build a interaction that describe such systems.
- + Make ab initio calculations (numerical problem)

## The goals

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- + Understanding of the **mechanisms** of nuclear properties;
- + Support to **experiments**;
- + **Precision description** of nuclear observables;

# What is universality? – unitarity (2-body only)

\*nonrelativistic & quantum

Size of the **two-body\*** system

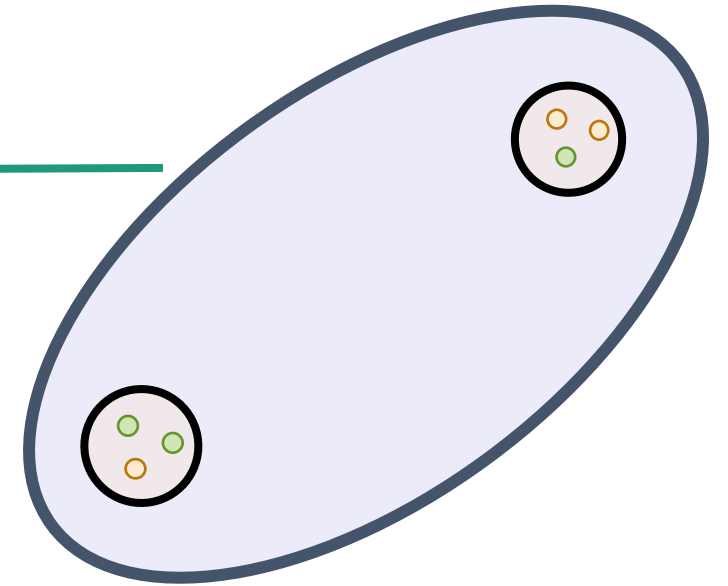
>>

range of the **interaction**  
and **particle size**.

Ideally:

$a_0 \rightarrow \infty$  (Scattering length)

$r_0 \rightarrow 0$  (Effective range)



# What is universality? – unitarity (2-body only)

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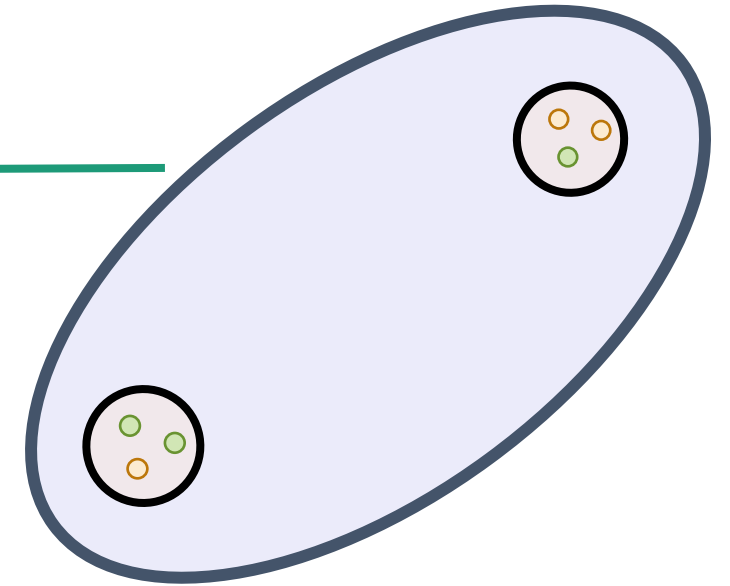
Ideally:

$a_0 \rightarrow \infty$  (Scattering length)

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Systems close to the **Unitary limit** can be found in

- **Atomic physics** (Feshbach resonances,  ${}^6\text{Li} - {}^6\text{Li}$ ,  ${}^{40}\text{K} - {}^{40}\text{K}$  atoms)
- **Nuclear physics** ( $n - p$  interaction )
- **Lattice nuclei** (Unphysically large  $m_\pi$ )
- **Hypernuclei** ( $\Lambda - n$  interaction)
- **Hadronic physics** ( $X(3872)$  Particles)
- **Cold Atom Physics** ( Unitary two-specie fermions )



**Atoms (experiments):**

C.A. Regal (2003)

M.W. Zwierlein (2003)

M. E. Gehm (2003)

J. T. Stewart (2007)

**Nuclei (theory):**

U. van Kolck (1999)

S. König (2017)

**Hypernuclei (theory):**

H.-W. Hammer (2001)

L.C. (2018)

**Hadrons (theory):**

E. Braaten et al (2003)

**Lattice Nuclei (theory):**

N. Barnea et al (2015)

L.C. et al (2017)

# Universality

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We observe very different systems to have **similar few-body** properties



2B scattering parameters;  
Few-body states;



We can use a **similar theory**

# Universality

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Same microscopic  
**symmetries;**  
**Separation of scales**



We observe very  
different systems to  
have **similar few-**  
**body** properties



Same **many-body**  
phenomena

## The Idea

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If you know that  
**two systems** with the  
same fundamental symmetries  
or the same few-body properties,  
you can design a **common theory**.

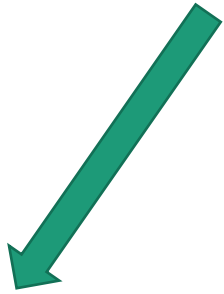
**2B scattering**  
**parameters;**  
**Few-body states;**



We can use a  
**similar theory**

# Simple and intuitive: Contact theory

- Treat **particles as degrees of freedom** (elementary particles)
- They can interact only **short-range**  
(Short range structure is irrelevant: no quark structure)  
(Long range interactions are negligible: no pion exchange)



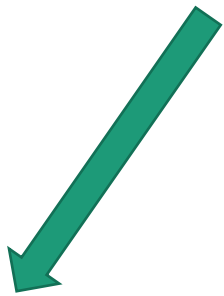
- Works for a limited set of energies



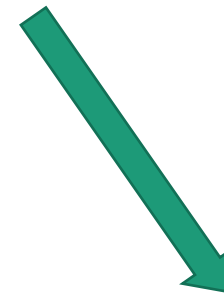
- Easy to understand
- Clear limitations
- Expandible
- Minimal inputs required
- Universally transposable

# Simple and intuitive: Contact theory

- Treat **particles as degrees of freedom** (elementary particles)
- They can interact only **short-range**  
(Short range structure is irrelevant: no quark structure)  
(Long range interactions are negligible: no pion exchange)



- Works for a limited set of energies
- Tricky to be properly implemented
- Clear limitations only in the known cases
- Not trivial to be practically expanded beyond 1st order
- Minimal inputs required at the first orders



- Easy to understand
- Clear limitations
- Expandible
- Minimal inputs required
- Universally transposable



# A complete theory

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Contact theory formally:

$$L = N^\dagger \left( i\partial_0 + \frac{\hbar^2}{2m} \nabla^2 \right) N - C_0 N^\dagger N^\dagger N N$$


$$V(r_{ij}) = \delta(r_{ij})$$

$r_{ij} = r_i - r_j$

# A complete theory

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$$L^{N>0LO} = C_2 (N^\dagger \nabla^2 N N^\dagger N + h.c.) + C_{11} (N^\dagger \vec{\nabla} N N^\dagger \vec{\nabla} N) +$$
$$C_4 (N^\dagger \nabla^4 N N^\dagger N + h.c.) + \dots$$
$$D_0 (N^\dagger N^\dagger N^\dagger N N N) + E_0 (N^\dagger N^\dagger N^\dagger N^\dagger N N N N) + \dots$$

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Including all the derivative/many-body operators one can **express any interaction**

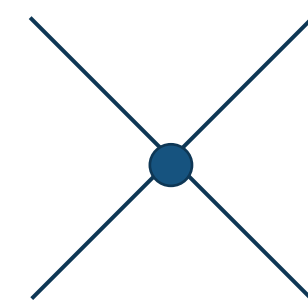
# A complete theory

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Contact theory formally:

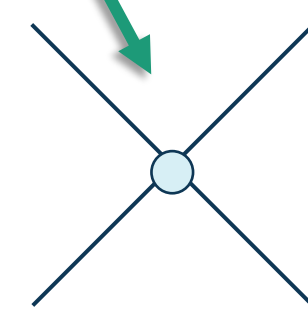
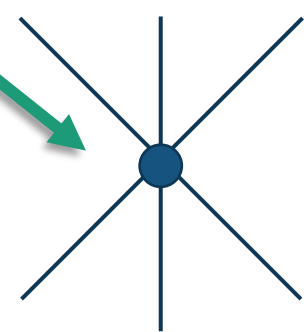
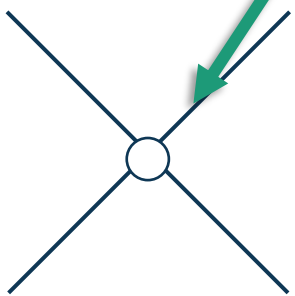
$$L = N^\dagger \left( i\partial_0 + \frac{\hbar^2}{2m} \nabla^2 \right) N - C_0 N^\dagger N^\dagger N N$$

$$V(r_{ij}) = \delta(r_{ij})$$



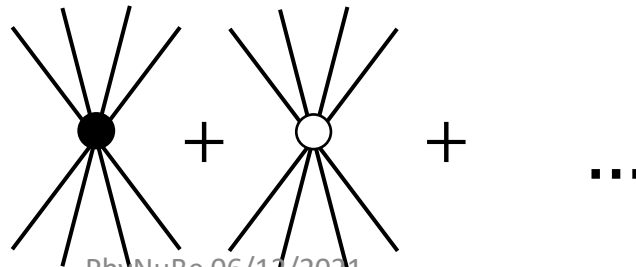
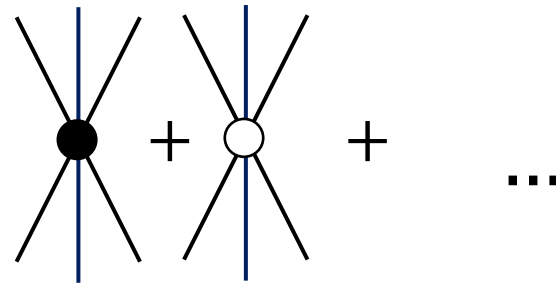
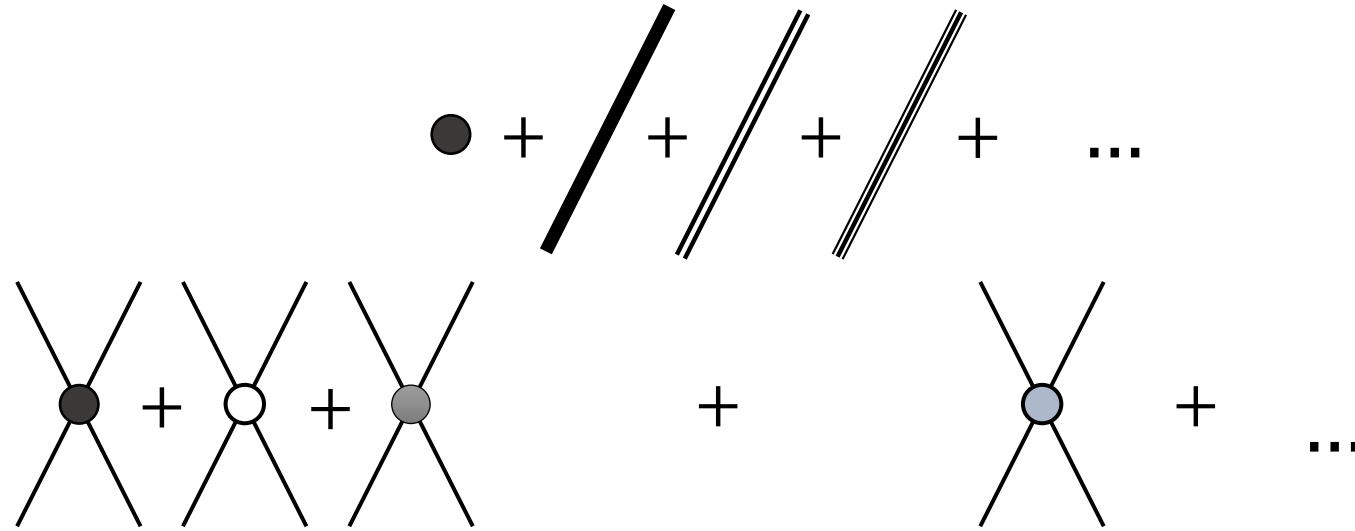
$$L^{N>0LO} = C_2 (N^\dagger \nabla^2 N N^\dagger N + h.c.) + C_{11} (N^\dagger \vec{\nabla} N N^\dagger \vec{\nabla} N) + C_4 (N^\dagger \nabla^4 N N^\dagger N + h.c.) + \dots$$

$$D_0 (N^\dagger N^\dagger N^\dagger N N N) + E_0 (N^\dagger N^\dagger N^\dagger N^\dagger N N N N) + \dots$$



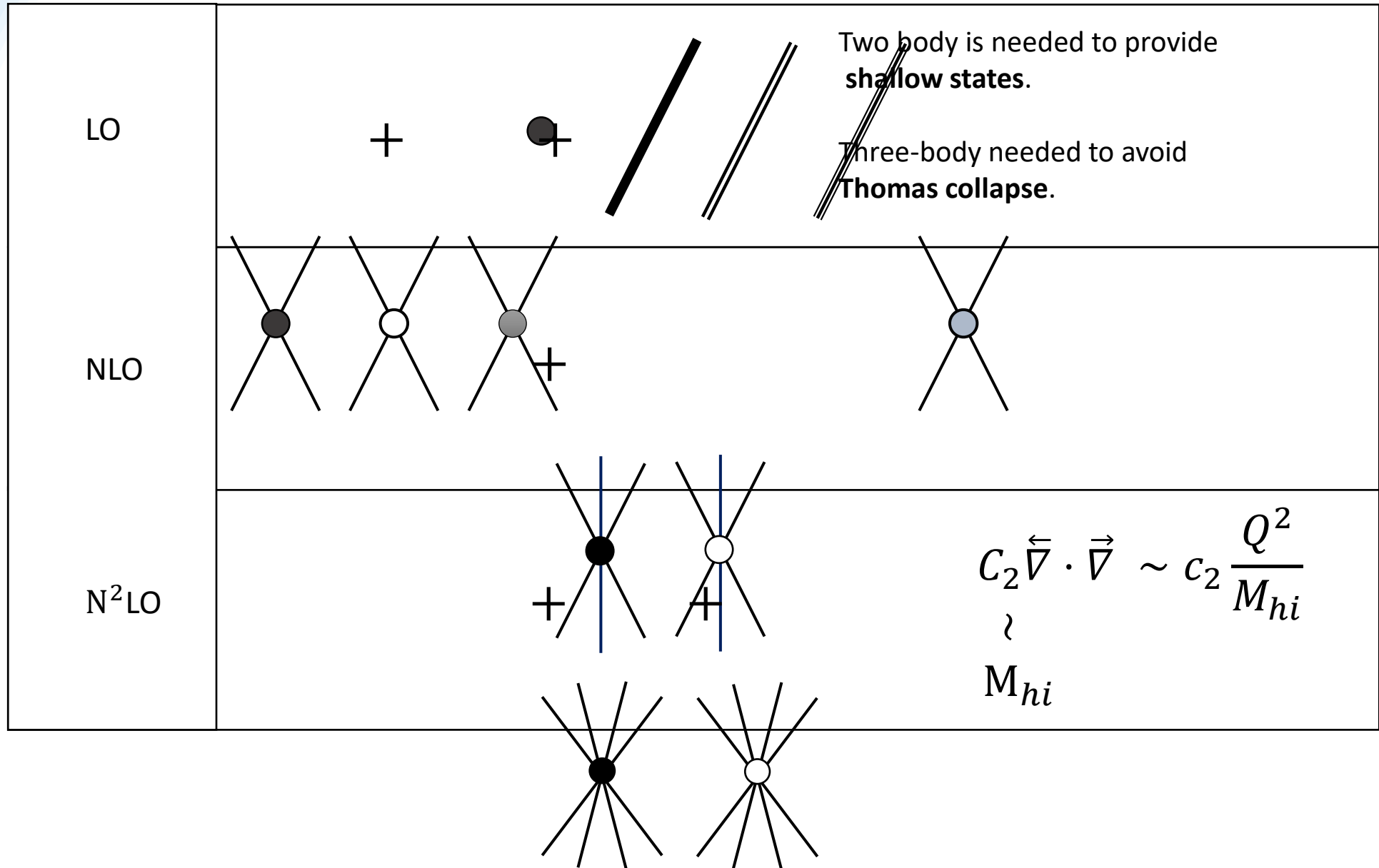
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# Pionless EFT powercounting



PhyNuBe 06/12/2021

# Pionless EFT powercounting

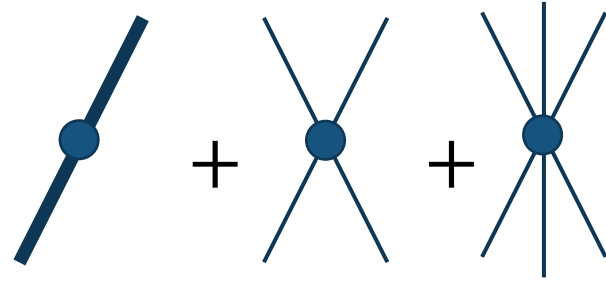


# Pionless EFT powercounting

In the nuclear case:  $\Gamma_{NN} = \frac{Q}{m_\pi} = 0.5 \sim 0.8$

Momentumless 2-3 body

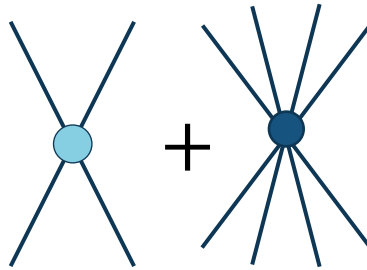
LO



1

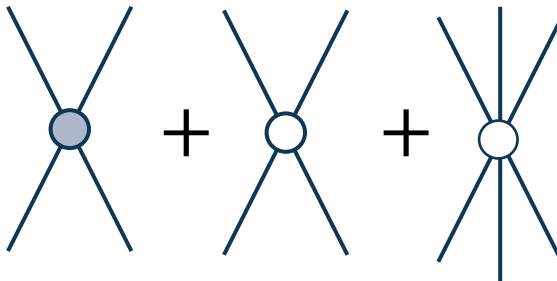
Momentum dependent / 4-body

NLO



$O(\Gamma)$

N<sup>2</sup>LO



$O(\Gamma^2)$

G.P. Lepage, How to renormalize the Schrödinger equation (1997)

U. van Kolck, Nucl.Phys. A645 273-302 (1999)

J.-W. Chen, et al. Nucl.Phys. A653 (1999)

S. König, H. W. Griesshammer, H. W. Hammer, and U. van Kolck, J. Phys. G43, 055106 (2016)

B. Bazak, PRL 122.143001 (2019)

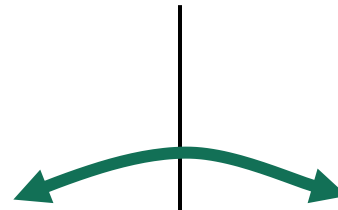
$O(\Gamma^{\geq 3})$

# Duality

universality

Unitary limit:  $a_0 = \infty$   
 $r_0 = 0$

Finite three-body scale:  $0 > E_3 > -\infty$



(contact) EFT  
(nonrelativistic)

$$\mathcal{L} = N^\dagger \left( \partial_0 + \frac{\nabla^2}{2m} \right) N + \quad \underline{\text{LO}}$$
$$+ C_0 N^\dagger N^\dagger N N + D_0 N^\dagger N^\dagger N^\dagger N N N$$



# Duality

## universality

Unitary limit:  $a_0 = \infty$   
 $r_0 = 0$   
Finite three-body scale:  $0 > E_3 > -\infty$

## (contact) EFT (nonrelativistic)

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$$+ C_0 N^\dagger N^\dagger N N + D_0 N^\dagger N^\dagger N^\dagger N N N$$

However, no physical system is perfectly in the unitary limit

S. König (2016)

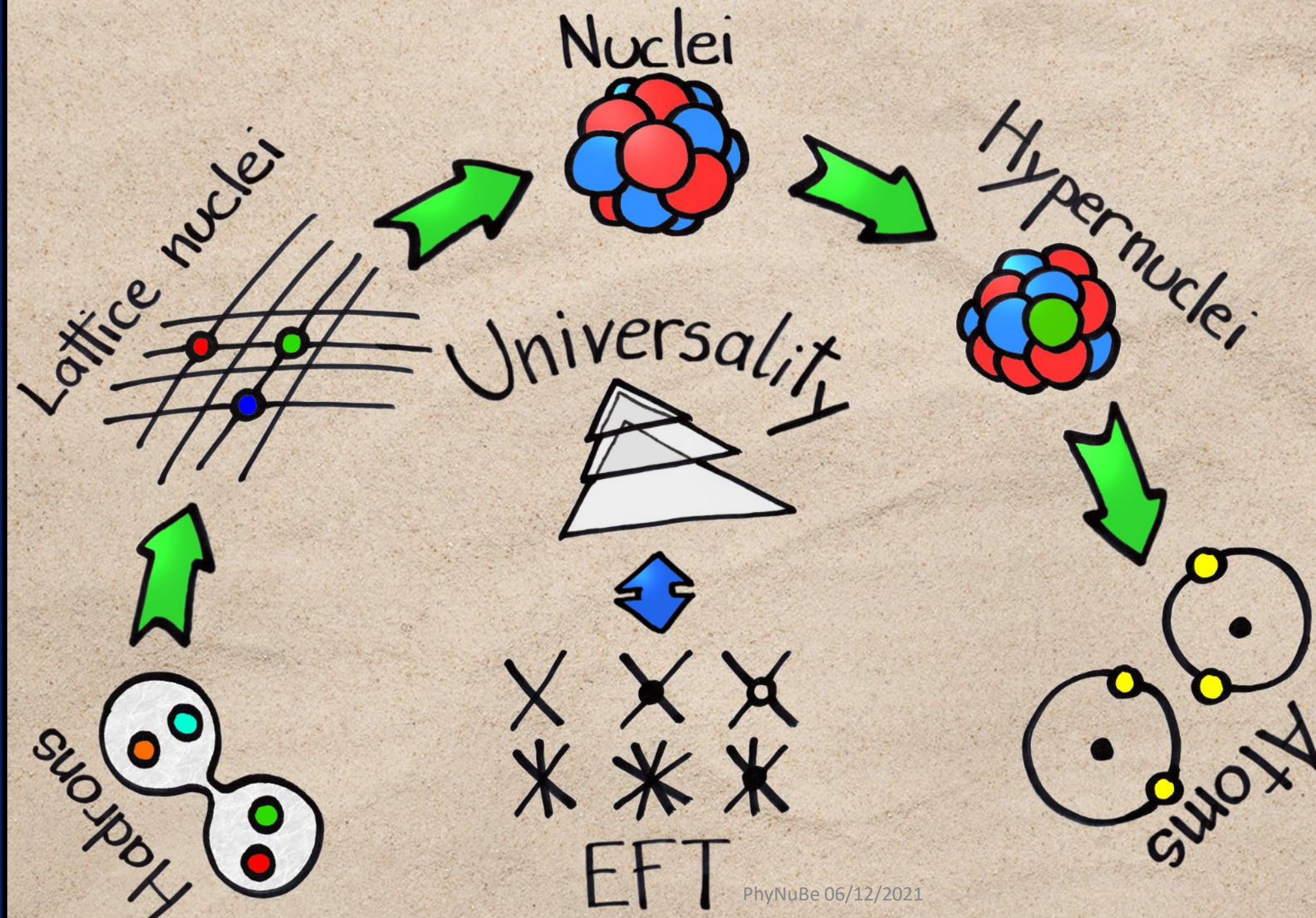
Physical systems can be close to the limit:  
e.g.  $|a_{n-n}| = (|-23. | \text{ fm}) \gg (r_0 \sim 2.7 \text{ fm})$

**Deviation from the universal limit**  
are needed to predict physical phenomena.

$N^n \text{LO}$

Effective field theory **powercounting**

**i.e. subleading perturbative corrections**  
define the specific physical system.



Interested in any of these systems?

Come for a chat!

Other examples are:  
 + Alpha clusters  
 + Condensed matter  
 + Neutron dorps  
 & neutron matter

# Hadrons

With a simple theory and only knowing  
That  $\mathcal{D} - \bar{\mathcal{D}}$  interaction is **unitary**  
The **range** of such interaction

A LO theory can be fitted

**Predict bound 3X, 4X  
(qualitative prediction)**

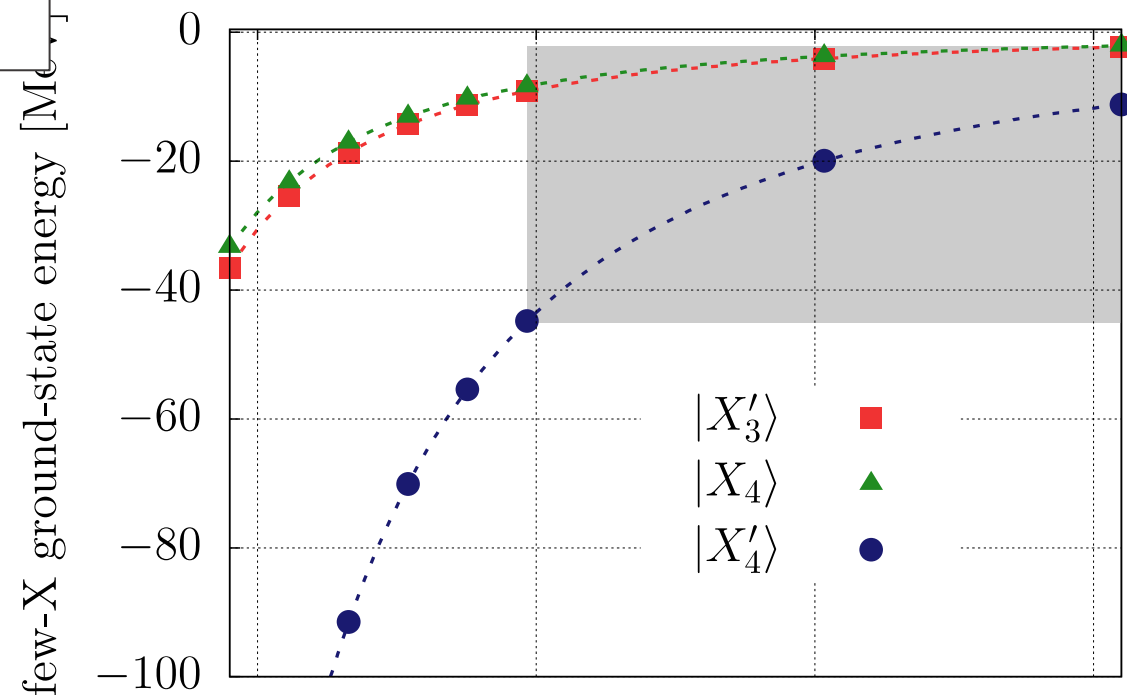
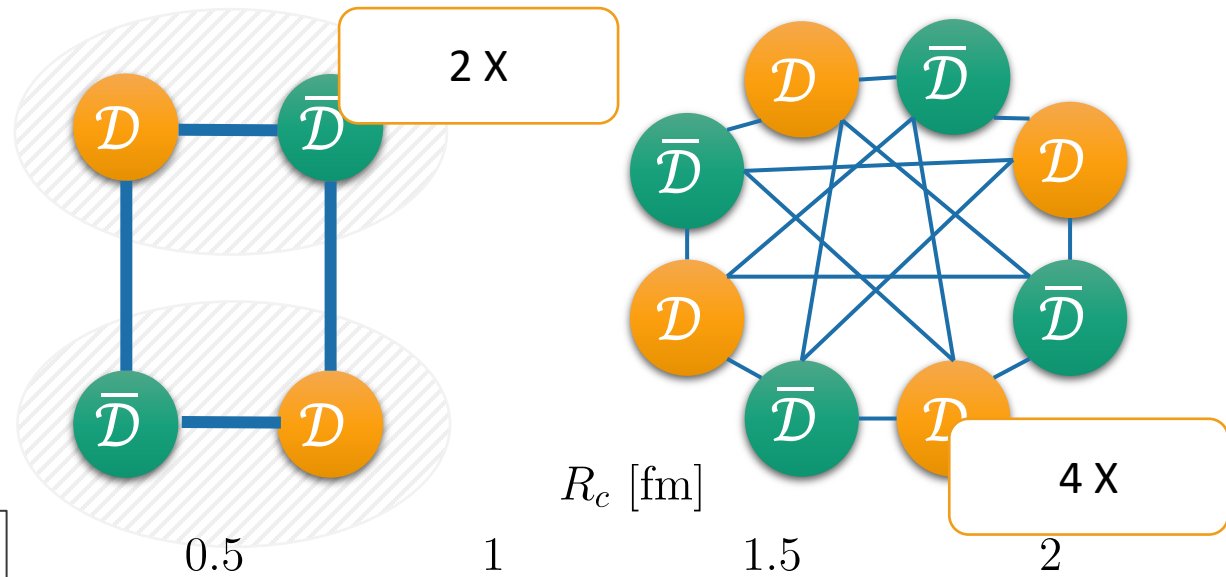
**2X is uncertain  
(bound only for certain cut-offs)**

This might describe a  
Brunnian system!

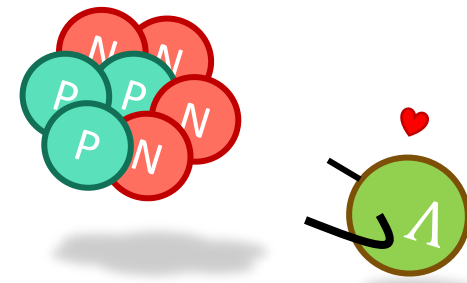
Minimal input parameters + underlying symmetries



Allow to predict effects unknown experimentally!



# Hyperons ( $\Lambda$ – Hypernuclei)



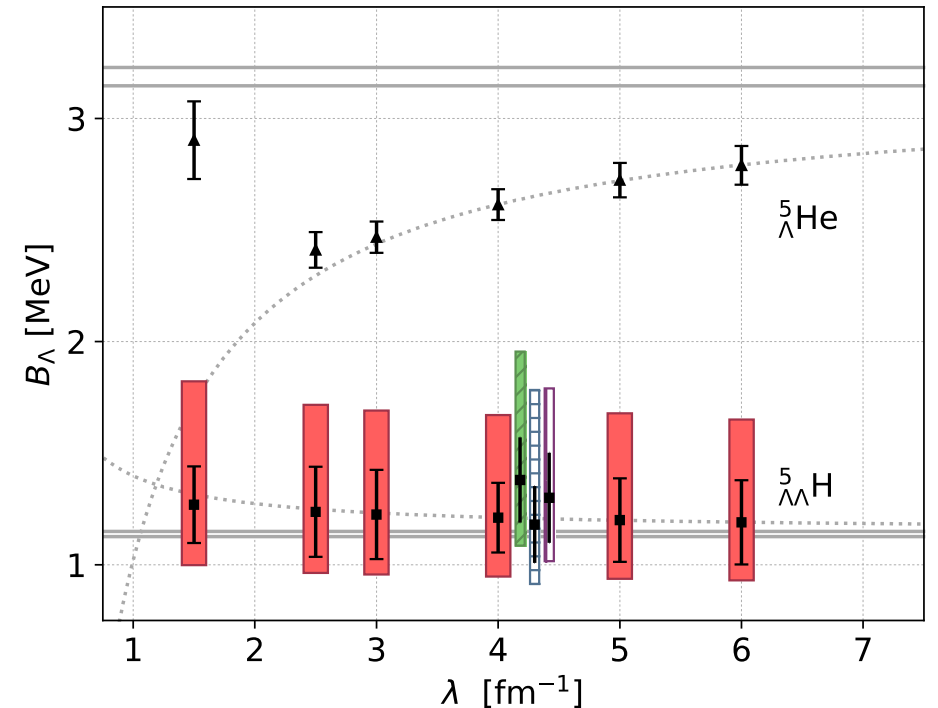
$\Lambda - N$  interaction is short range  
 $\Lambda - \#\#$  is a shallowly bound system
 } Perfect example of interaction close to unitary limit

**Single  $\Lambda$ :** very few experimental data  
**Double  $\Lambda$ :** hard to be created experimentally  
**Few-body hypernuclei:** theoretically hard to be described all together

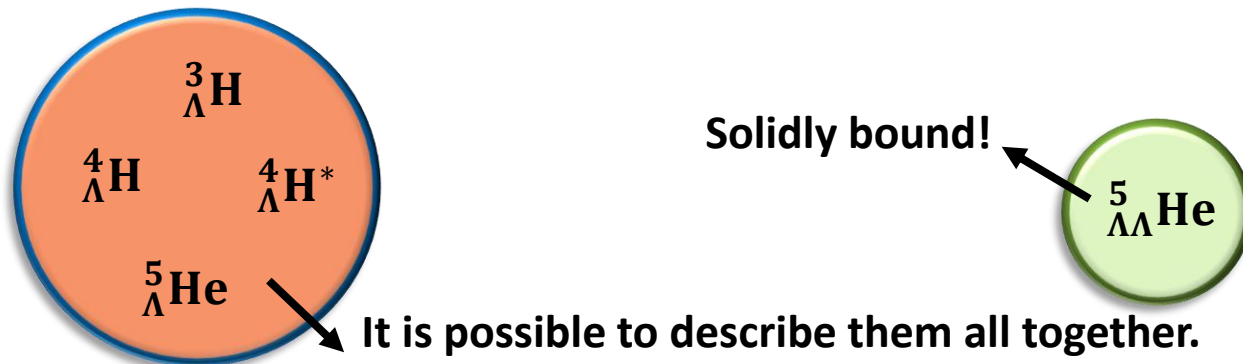
**LO EFT has few input:** can be fitted from existing experiments;  
**Theoretical errors** are relatively small;  
**The powercounting guides** in the choice of the relevant operators:  
**solves the overbinding problem;**

Is proved only for shallowly bound (**few-body**) hypernuclei

$$B_{\Lambda}({}^5_{\Lambda\Lambda}\text{He}) = 1.14(1)^{+(44)}_{-(26)}$$



L. C., M. Schäfer, N. Barnea, A. Gal, J. Mareš  
 Phys. Lett. B 797 (2019) 134893



**Solidly bound!**  
**It is possible to describe them all together.**  
**( No overbinding problem! )**

# Nuclei

Nuclear physics is “almost” at the unitary limit

$$a_{np} \sim 5.5 \text{ fm} \quad r_{np} \sim 1.5 \text{ fm}$$

$$a_{nn} \sim 20 \text{ fm} \quad r_{nn} \sim 2.5 \text{ fm}$$

Deviations may play an **important role!**

$$E^{\text{LO}}(^4\text{He}) = 29.2 \pm 0.5 \text{ MeV}$$

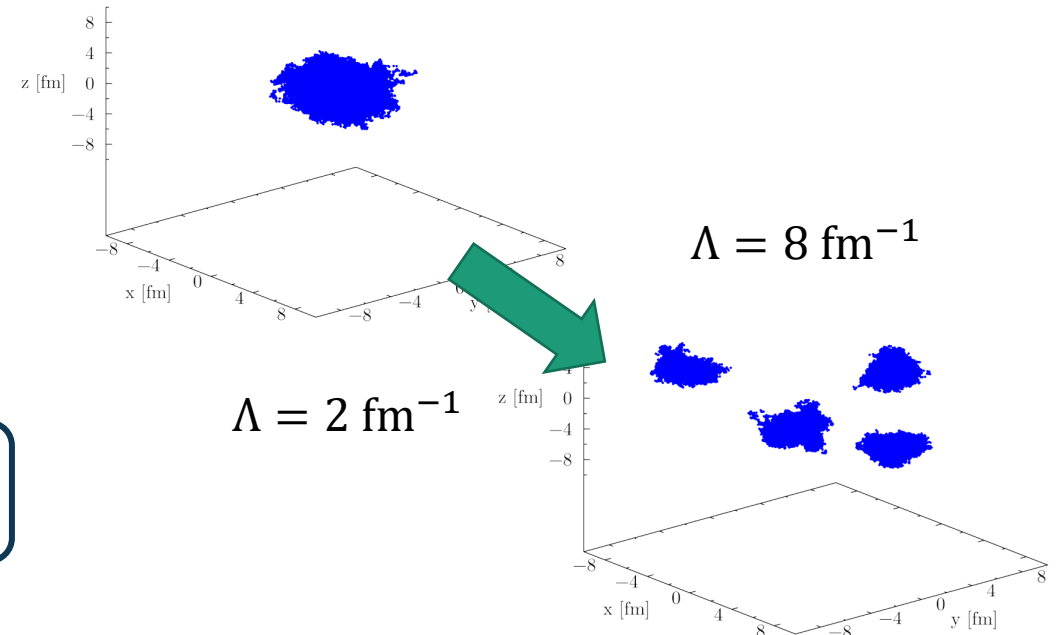
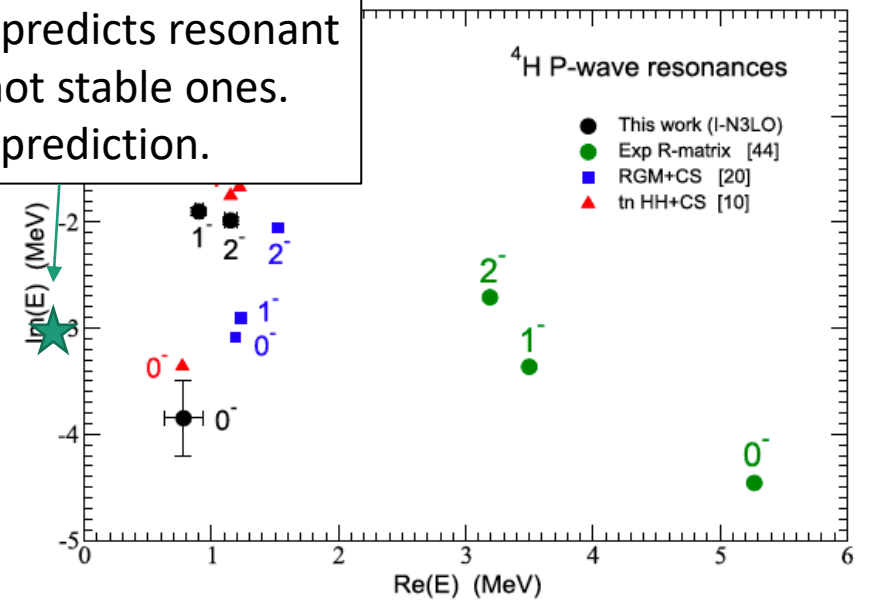
Is **close** to the experimental data

$$E^{\text{Exp}}(^4\text{He}) 28.3 \text{ MeV.}$$

But larger nuclei (systems with mixed symmetry)  
Appear to be **never stable in the universal regime!**

What kind of **deviations** should we include to stabilize them?

The theory predicts resonant states but not stable ones.  
Here is our prediction.



$\Lambda$  is the cut-off of the interaction  
( $\Lambda \rightarrow \infty$  is the universal limity)

# Lattice nuclei

L. C., A. Lovato, F. Pederiva, A. Roggero, J. Kirscher,  
U. van Kolck Phys. Lett. B **772**, 839-848 (2017)

With Lattice QCD you can access  
**high pion masses** physics.

Completely **ab initio** calculations.

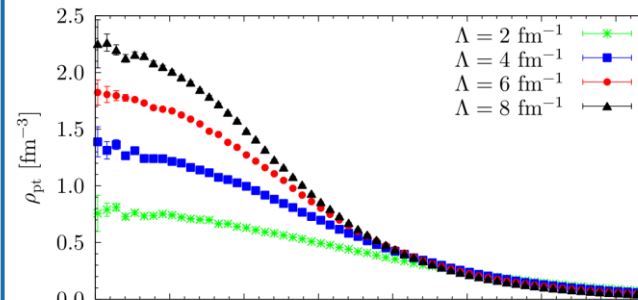
Would be the physics as we know it?

The interaction range is **shorter**  
The binding momentum is **the same**

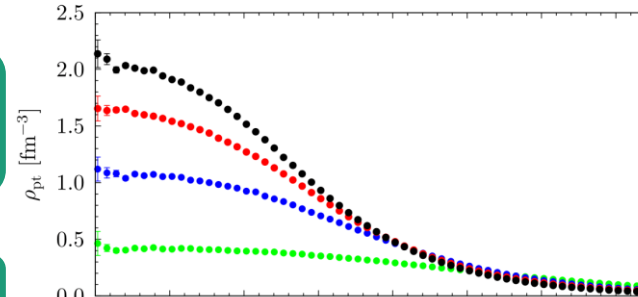
**Contact interaction**

**No stable P-shell nuclei**  
(as universality predicts)

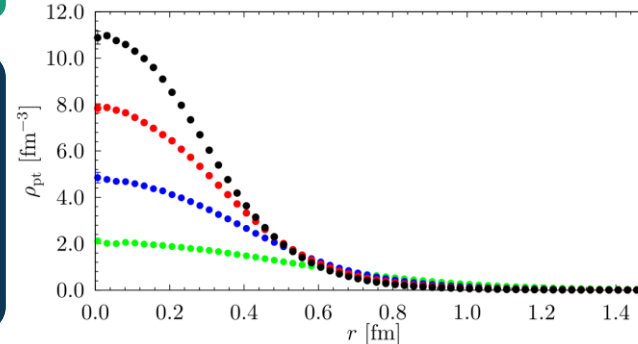
**The world would be very different!**



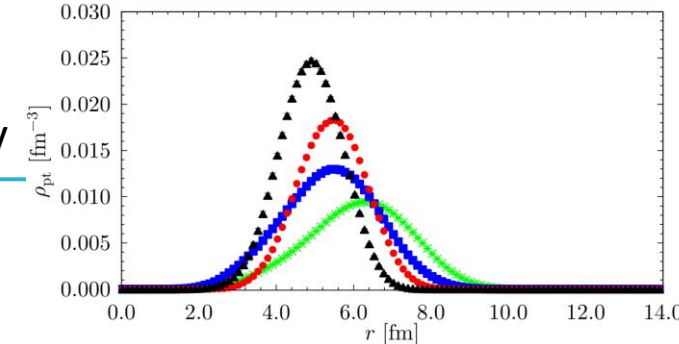
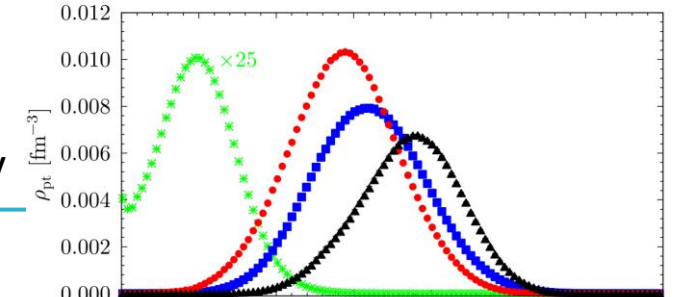
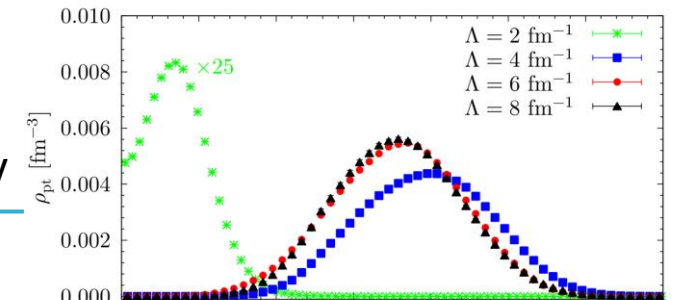
$m_\pi = 140$  MeV



$m_\pi = 510$  MeV



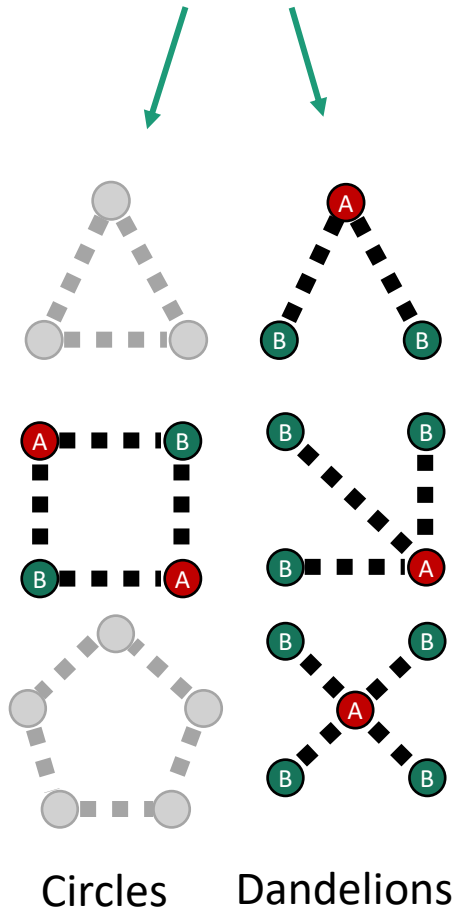
$m_\pi = 805$  MeV



# Atoms (Bosons)

Systems with only a **subset of unitary interactions**

We see two different kind of behaviours

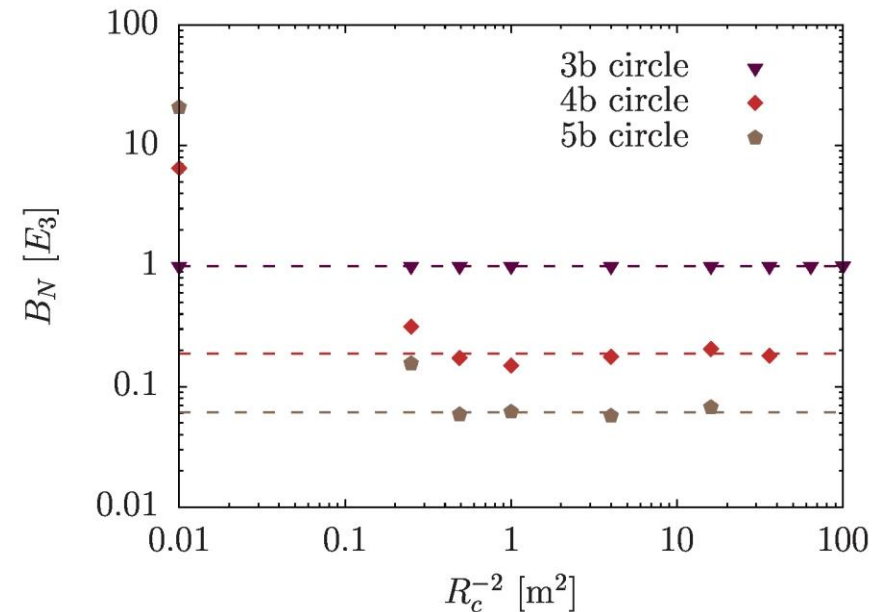


Circles

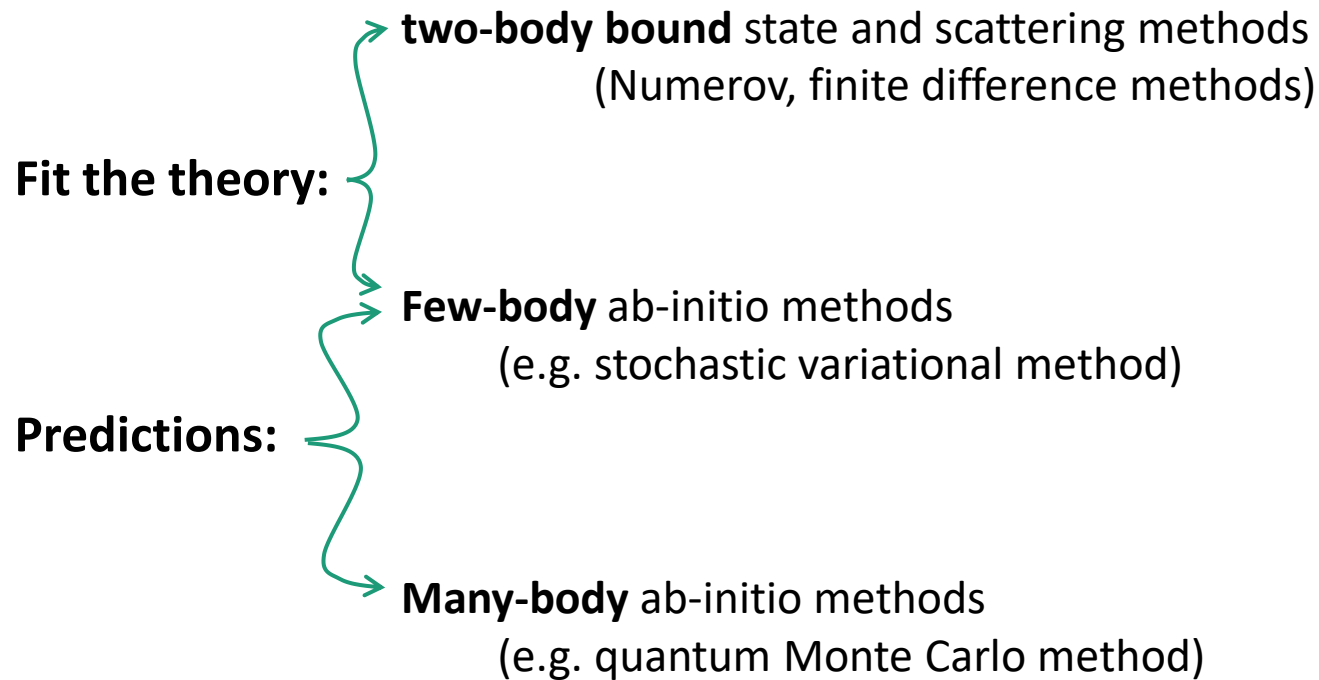
The square system **behave differently** compared with the other systems composed by two-atom species.

**Can this be seen in experiments?**

The  $\square$  system energy depends from the  $\triangle$  but the  $\triangle$  is formed by different kind of particles



# Numerics



## SVM:

Variational diagonalization method based on random gaussian basis expansion

## Monte Carlo:

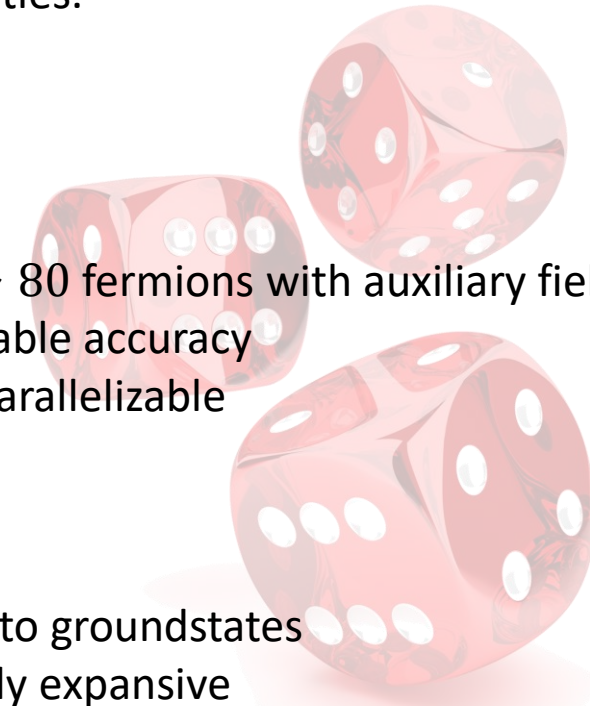
**Quantum Monte Carlo** is a class of ab initio, numerical, stochastic many-body methods able to solve the Schrödinger equation with improvable uncertainties.

### Pros:

- + Up to  $\sim 80$  fermions with auxiliary field technique
- + Improvable accuracy
- + Easily parallelizable

### Cons:

- Limited to groundstates
- Generally expensive
- Sign problem makes things hard

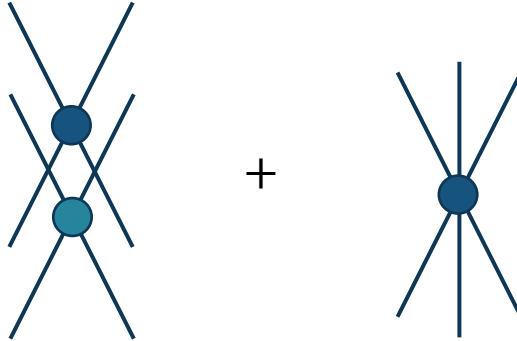




# Nuclei: a personal story

## Leading order theory:

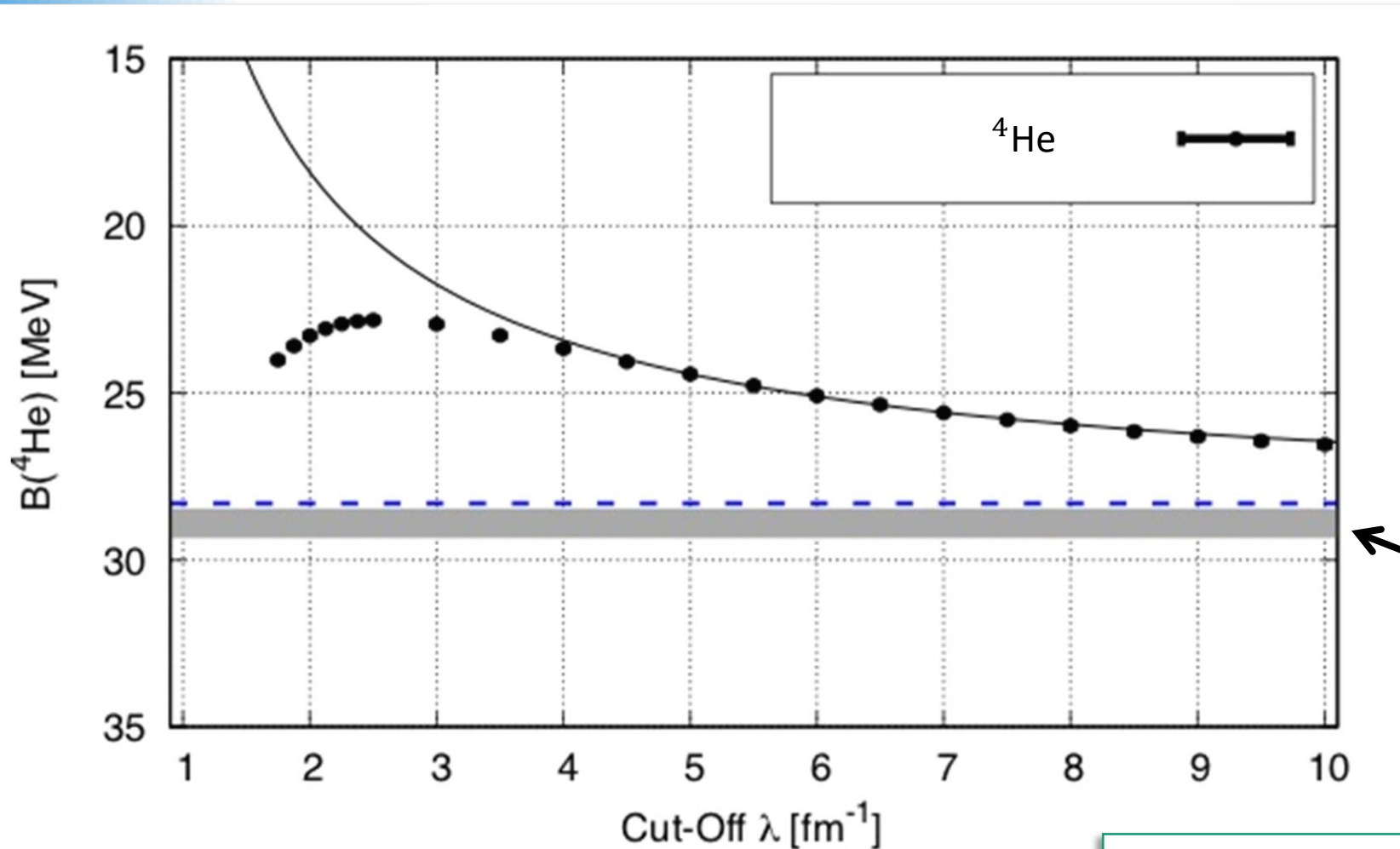
2-body + 3-body fitted to reproduce (n-n) scattering, deuterium, and triton binding energy

$$H^{LO} = -\frac{\hbar^2}{2m} \sum_i \nabla^2 +$$


The diagram shows two Feynman diagrams representing interactions. The first diagram on the left shows two particles (represented by blue dots) interacting via a two-body force, with two lines connecting them. The second diagram on the right shows a single particle (represented by a blue dot) interacting via a three-body force, with three lines connecting it to a central point.

No coulomb, no tensor force, no spin orbit at LO

# A practical example: few nucleons



LO pionless EFT theory fitted on **two- and three-body** observables predicts well  $^4\text{He}$  energy!

Fitted on:  $a_{n-n} = -18.63$  fm

$B_d = -2.22$  MeV

$B_t = -8.48$  MeV

--- Experimental data: 28.3 MeV

— Extrapolation:  $29.2 \pm 0.5$  MeV

Calculations done with **few-body stochastic variational diagonalization method**: Y. Suzuki, K. Varga (2003)

# $^{16}\text{O}$ - Monte Carlo calculation

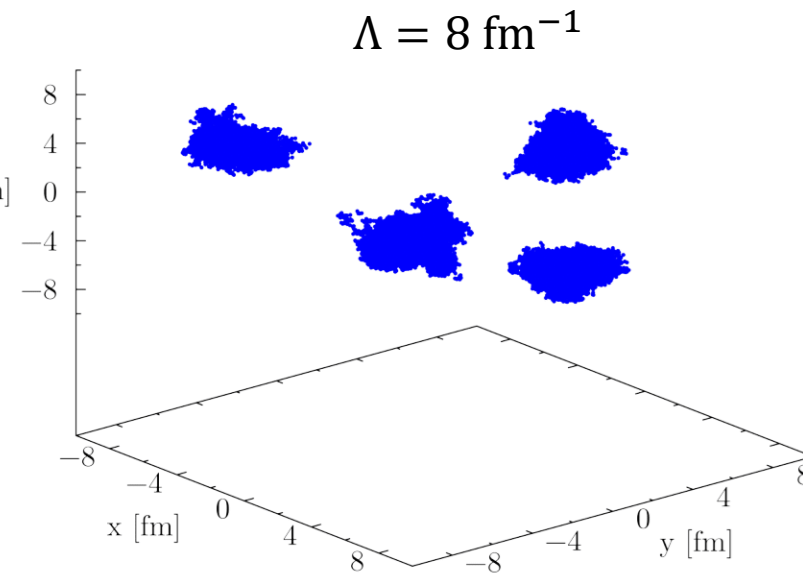
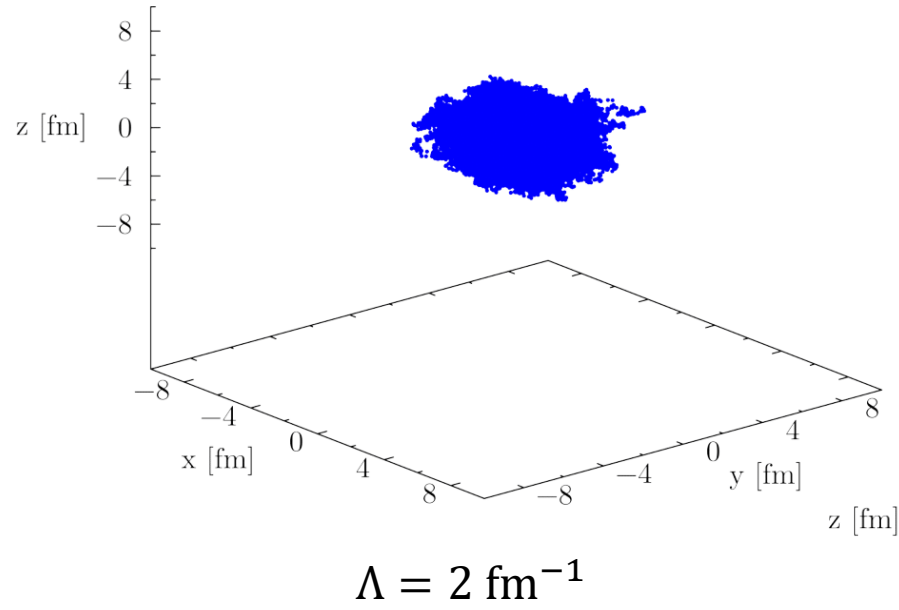
*Phys.Lett.B 772 (2017) 839-848*

S-wave system		P-wave system		
$\Lambda$ [fm <sup>-1</sup> ]	<sup>4</sup> He Energy [MeV]	$\Lambda$ [fm <sup>-1</sup> ]	<sup>16</sup> O Energy [MeV]	4 $\alpha$ treshold [MeV]
2	-23.17(2)	2	-97.19(6)	-92.68(8)
4	-23.63(3)	4	-92.23(14)	-94.52(9)
6	-24.06(2)	6	-97.51(14)	-100.24(8)
8	-26.04(5)	8	-100.97(20)	-104.2(2)
$\infty$	$-30^{0.3(sys)}_{2.0(stat)}$	$\infty$	$-115^{1(sys)}_{8(stat)}$	$-120^{1(sys)}_{8(stat)}$
Exp	-28.296			

- All the errors shown are statistical errors from Monte Carlo method.

Be( $^{16}\text{O}$ )  $\sim$  127 MeV  
 Be(4 $\alpha$ )  $\sim$  113 MeV  
 It is only 10% difference!

# Oxygen density ( $m_\pi = 140$ MeV)



5He

6He

J. Kirscher, H. W. Griesshammer, D. Shukla,  
H. M. Hofmann: arXiv:0909.5606

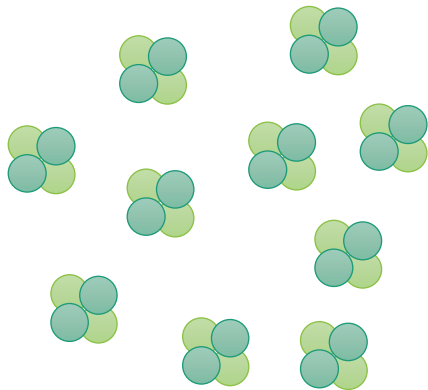
Breaks in  $\alpha + n$  and  $\alpha + n + n$



40Ca

QMC calculation suggests the breaking in:

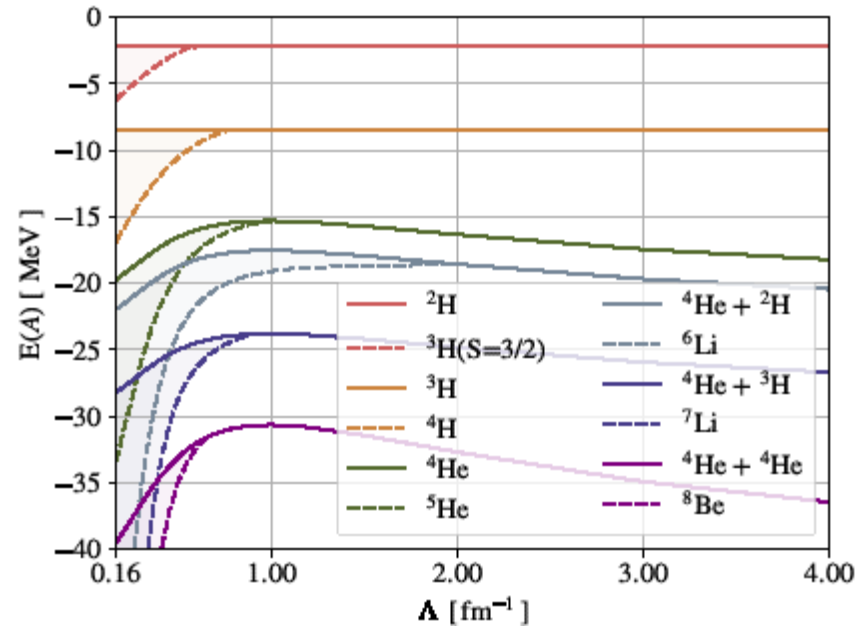
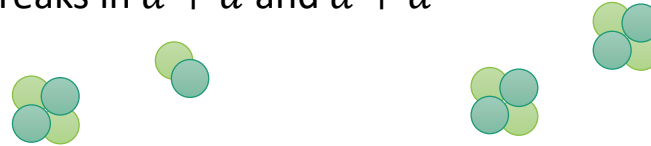
Breaks in  $\alpha + \alpha + \alpha + \dots$



7Li 8Be

Our calculations in SU(4) symmetry

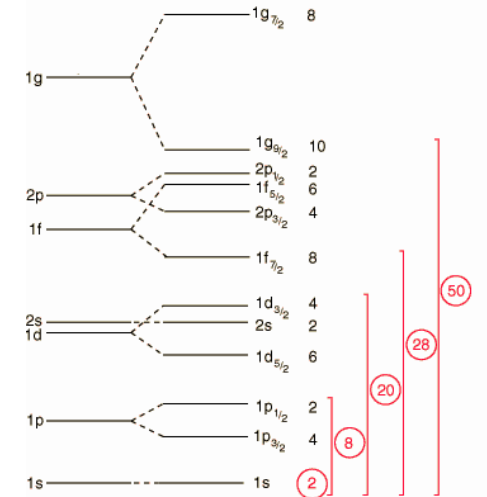
Breaks in  $\alpha + d$  and  $\alpha + \alpha$



P-wave systems

In a shell model representation

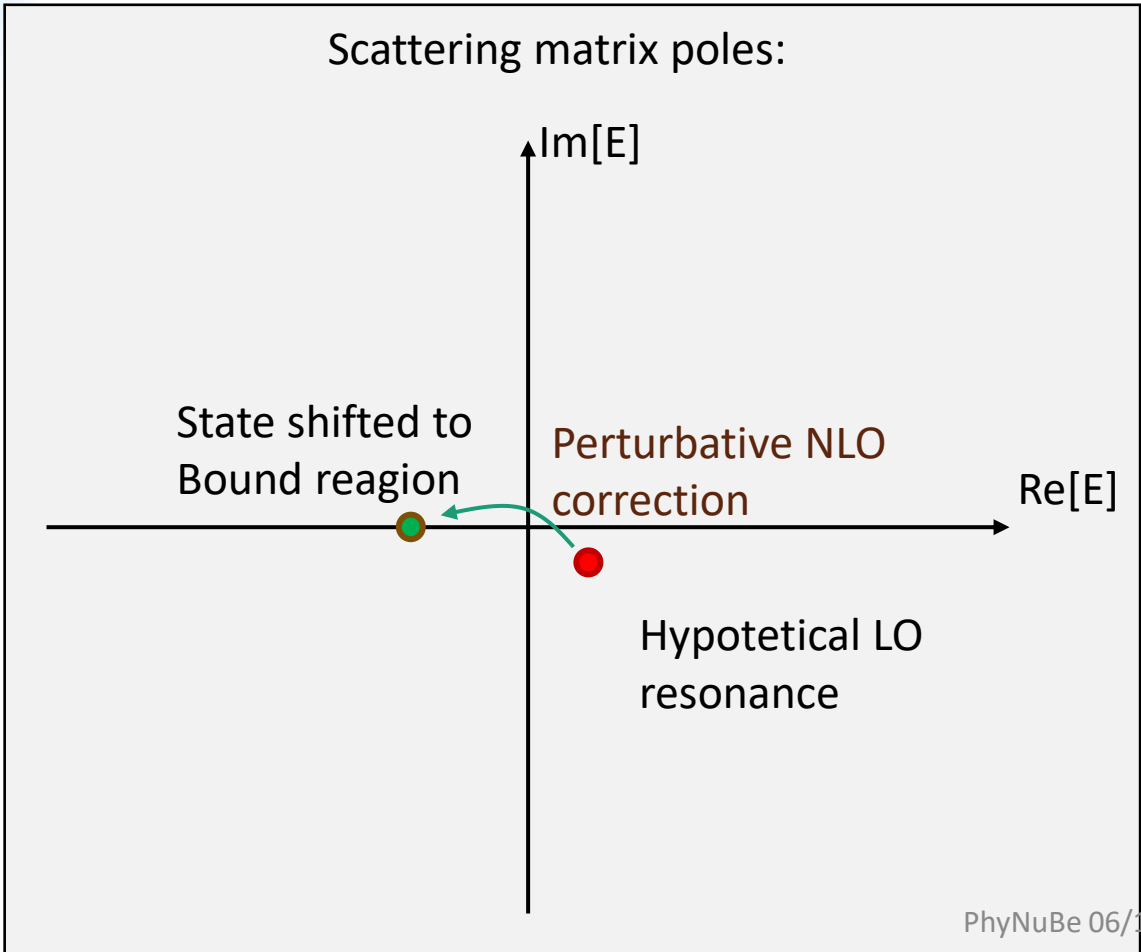
Relation between shell model and magic numbers



Will they ever bind?

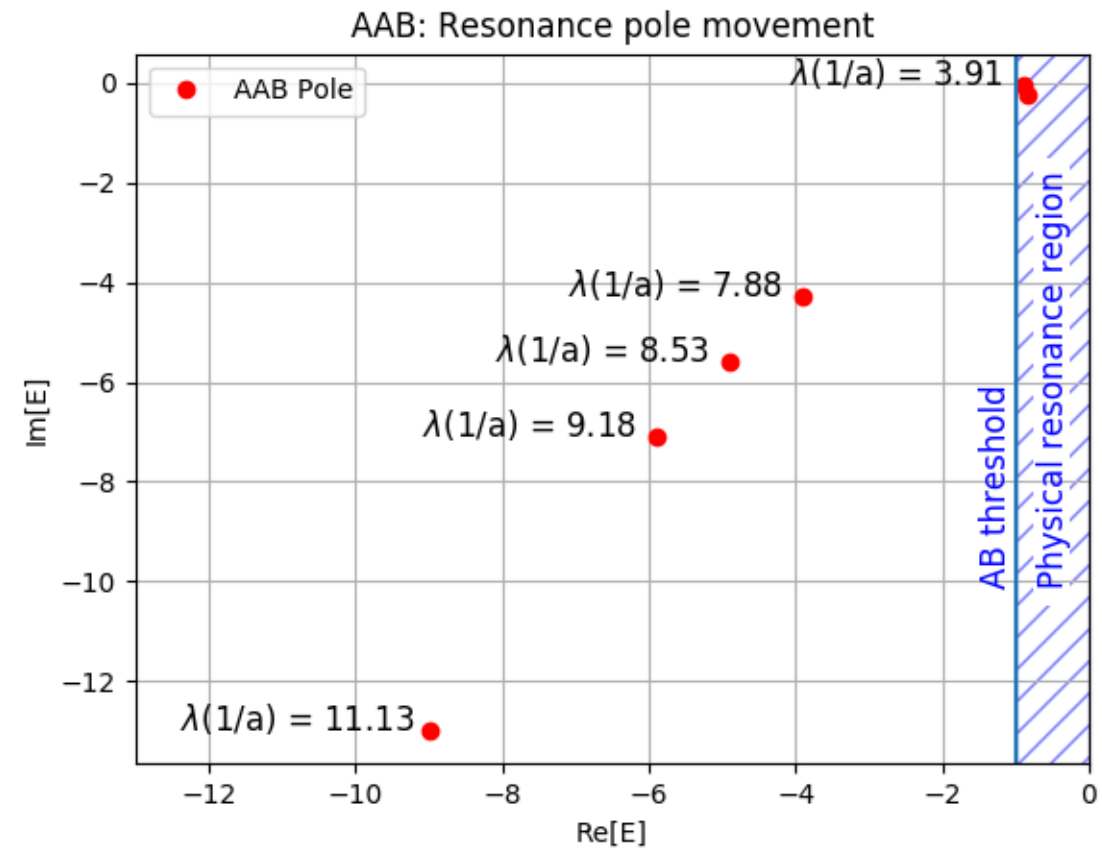
One little step further is necessary:

If a resonance is close to the **threshold**, it might be possible to move it with a **subleading correction** (there is no proof this is possible, nor proof this is not possible)

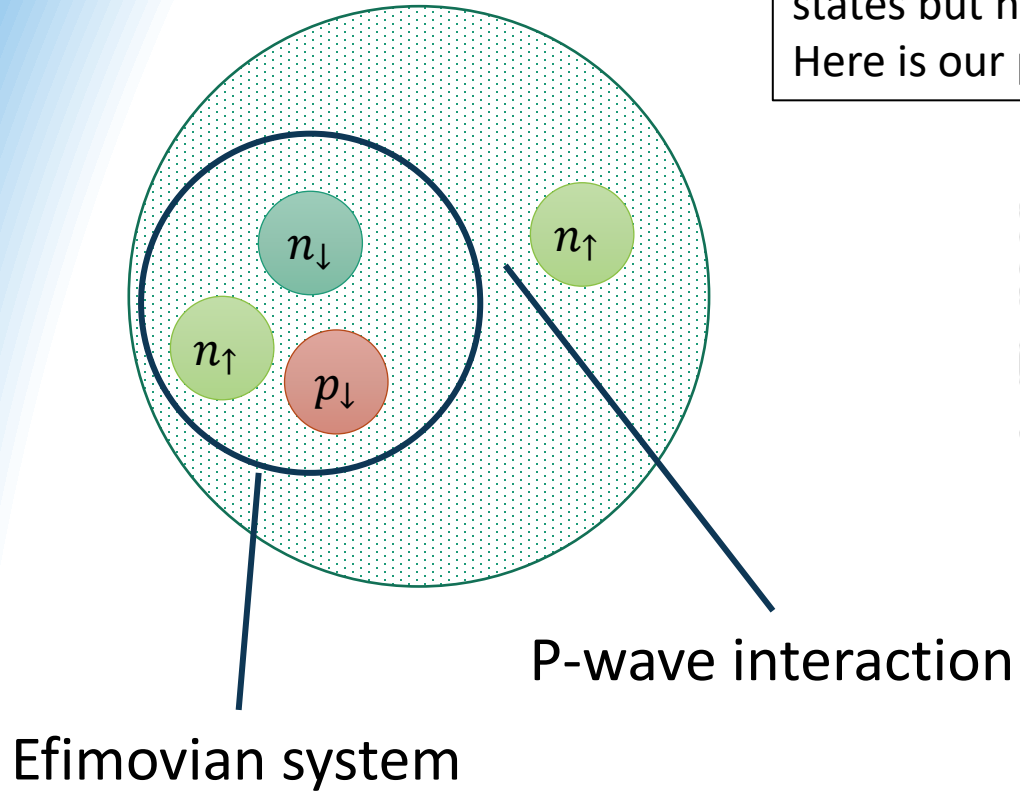


Known **three-fermion** case:  
No physical resonance is found.

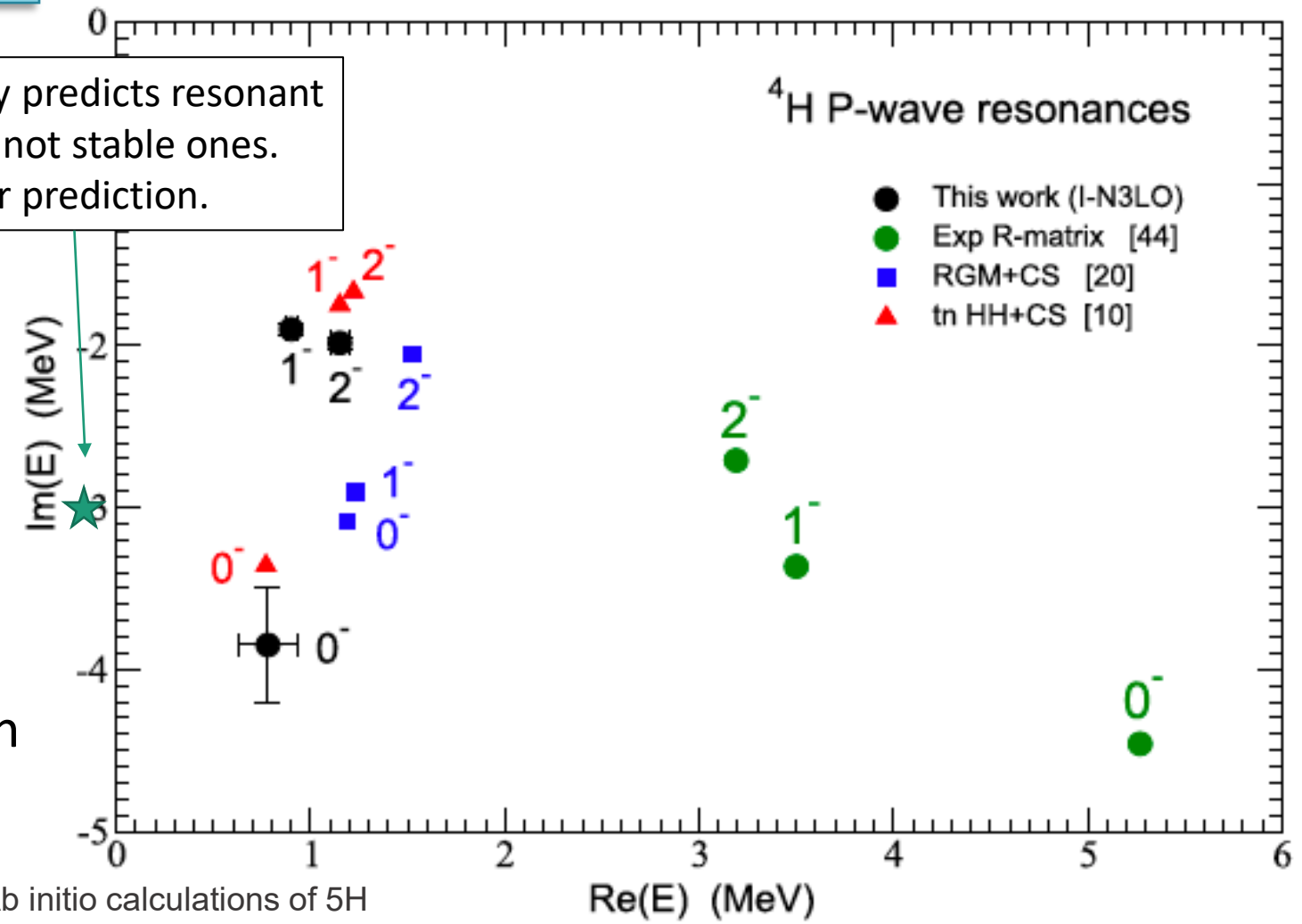
No scale invariance breaking,  
Three-body force might change picture.



# 4H resonance: the minimal nuclear system with an Efimovian component



The theory predicts resonant states but not stable ones. Here is our prediction.



R. Lazauskas, E. Hiyama, and J. Carbonell, "Ab initio calculations of 5H resonant states," Phys. Lett. B, vol. 791, pp. 335–341, 2019.

# Contact EFT: a sub-threshold resonance is present

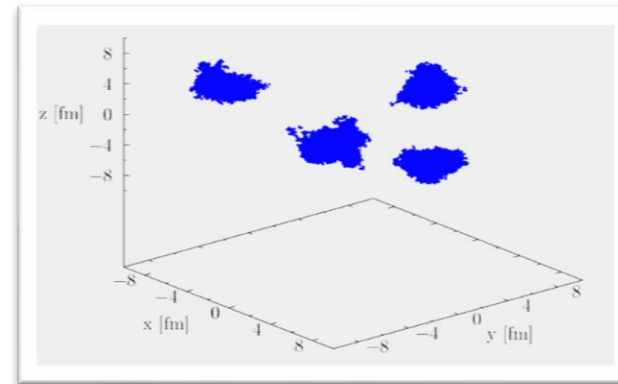
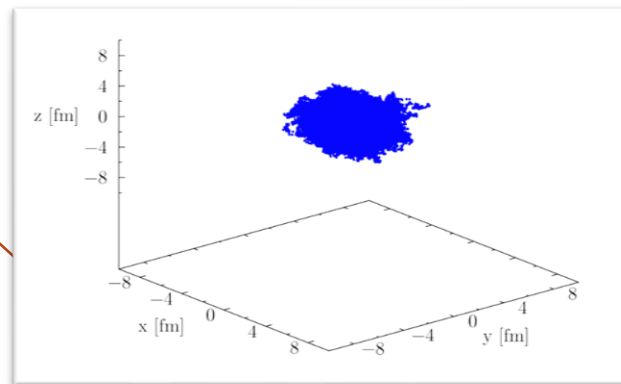
## Contact theory

- everything fine in **S-wave**
- **no P-wave boundstates**

A resonance is found in  ${}^4\text{H}$   
(not in the correct physical position)

- many-body **P-shell poles can be created**

Can the resonant pole be moved to the bounded region with a **perturbative** NLO insertion?  
Preliminary calculation showed that either  $\mathbf{P}^2$  or  $\mathbf{P} \cdot \mathbf{P}'$  can do the job.





# What can be learned

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## About universality:

- + Many systems show a **universal behaviour**.
- + **LO** of EFT takes care of **universality**;  
**Sub-leading** orders account for **deviations**.
- + Similar systems share **common EFT** description and you can transfer the knowledge to **unknown physics**.

- 
- + Many other **universality classes** are out-there.  
**What** are? **Where** to find them? Can they be **useful**?



The guarding "Chatula"  
Sept 2017 Jerusalem

# A weird concept: scale invariance (this is quite theoretical)

Define the **typical momentum** you are interested  
(probe momentum)

Two-body system is **infinitely large** with respect to this probe

The range of the interaction is **infinitesimal**

**No scale in the system: everything should be 0 or  $\infty$**   
[no measure units]

## Problem:

In real physics the three body energy is not **0** nor  $\infty$  ...  
Moreover, if you make the calculation in universality

$$E_3 \rightarrow -\infty$$

# Discrete scale invariance: 3+ body

Unitary system with 3 particles in S-wave show:

## Thomas Collapse / Efimov Effect

In the **unitary limit** a system of **3 bosons/distinguishable particles** collapses

$$\text{Unitarity} \Rightarrow E_3 \rightarrow -\infty$$

A **repulsion** is needed to stabilize the system to a finite energy  $E_3$ .

$E_3$  breaks the scale invariance of the system!

### Conditions to (Efimov) unitarity

Range of the interaction	$(r_0 \rightarrow 0)$
Size of the two body system	$(a_0 \rightarrow \infty)$
Size of the three body system	$(R_3 < \infty)$

L. H. Thomas (1935)  
G. Skorniakov and  
K. Ter-Martirosian (1957)

