

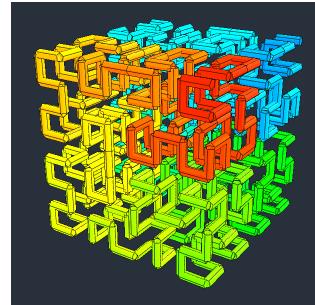
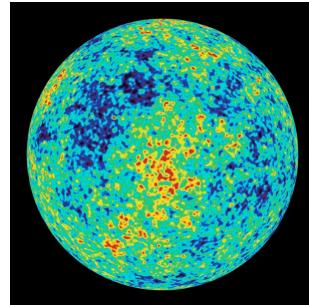
Precision cosmology: myth or reality?

A theoretical perspective

Romain Teyssier, Mischa Knabenhans, Laurent Legrand, Stefano Marelli,
Doug Potter, Aurel Schneider, Joachim Stadel, Bruno Sudret

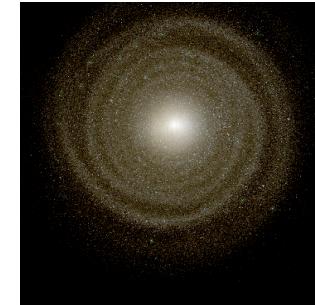
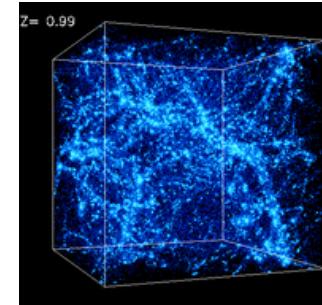


**University of
Zurich^{UZH}**



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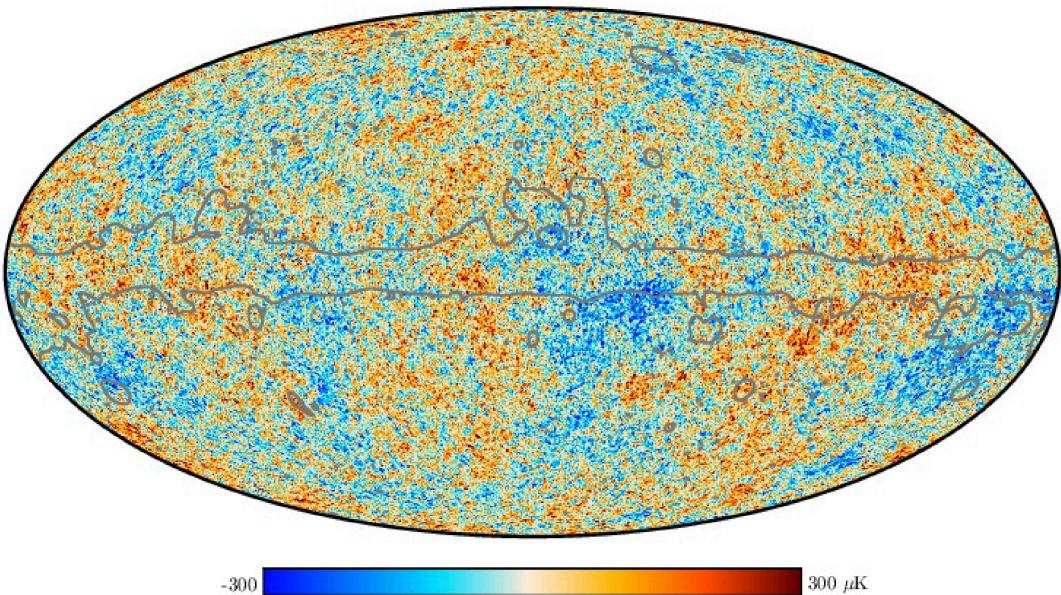
Eidgenössische Technische Hochschule Zürich
Swiss Federal Institute of Technology Zurich



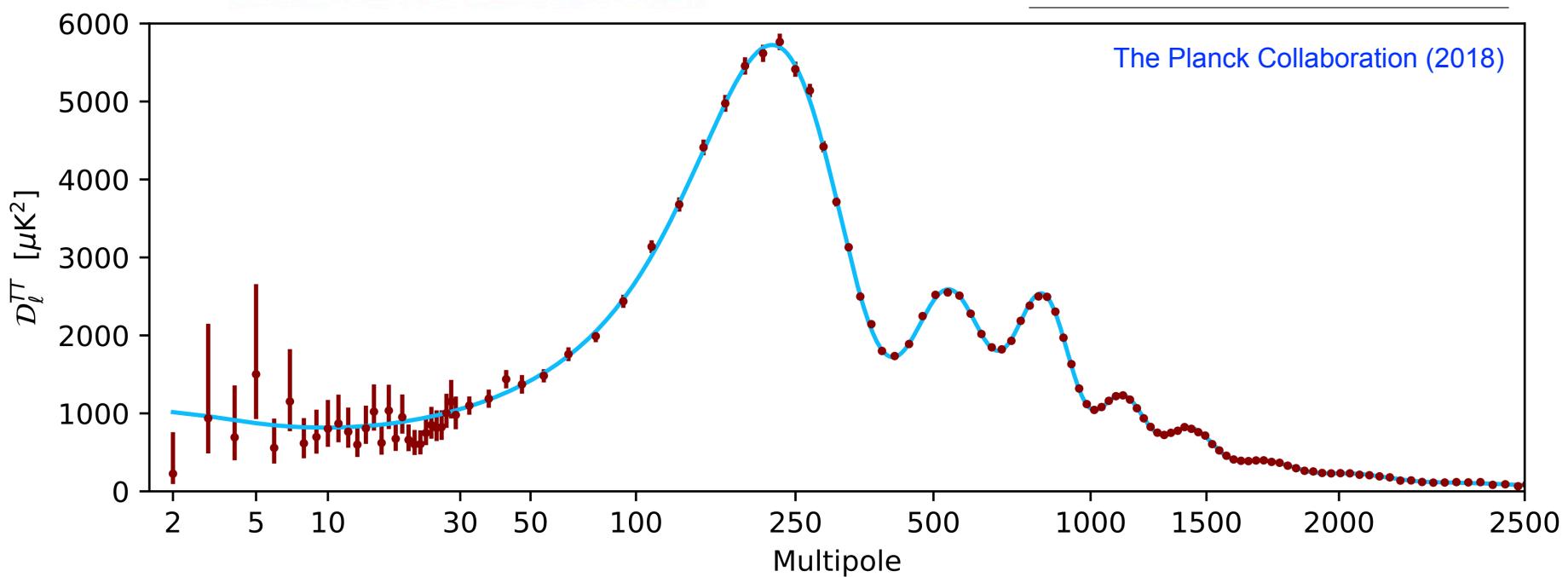
Context

Precision cosmology allows us to perform Bayesian inference on the standard model parameters, and may be discover exotic dark matter particles or modified gravity models.

Precision cosmology with the CMB sky



-300 300 μK



Parameter	<i>Planck alone</i>	<i>Planck + BAO</i>
$\Omega_b h^2$	0.022383	0.022447
$\Omega_c h^2$	0.12011	0.11923
$100\theta_{\text{MC}}$	1.040909	1.041010
τ	0.0543	0.0568
$\ln(10^{10} A_s)$	3.0448	3.0480
n_s	0.96605	0.96824
$H_0 [\text{km s}^{-1}\text{Mpc}^{-1}]$	67.32	67.70
Ω_Λ	0.6842	0.6894
Ω_m	0.3158	0.3106
$\Omega_m h^2$	0.1431	0.1424
$\Omega_m h^3$	0.0964	0.0964
σ_8	0.8120	0.8110
$\sigma_8(\Omega_m/0.3)^{0.5}$	0.8331	0.8253
z_{re}	7.68	7.90
Age [Gyr]	13.7971	13.7839

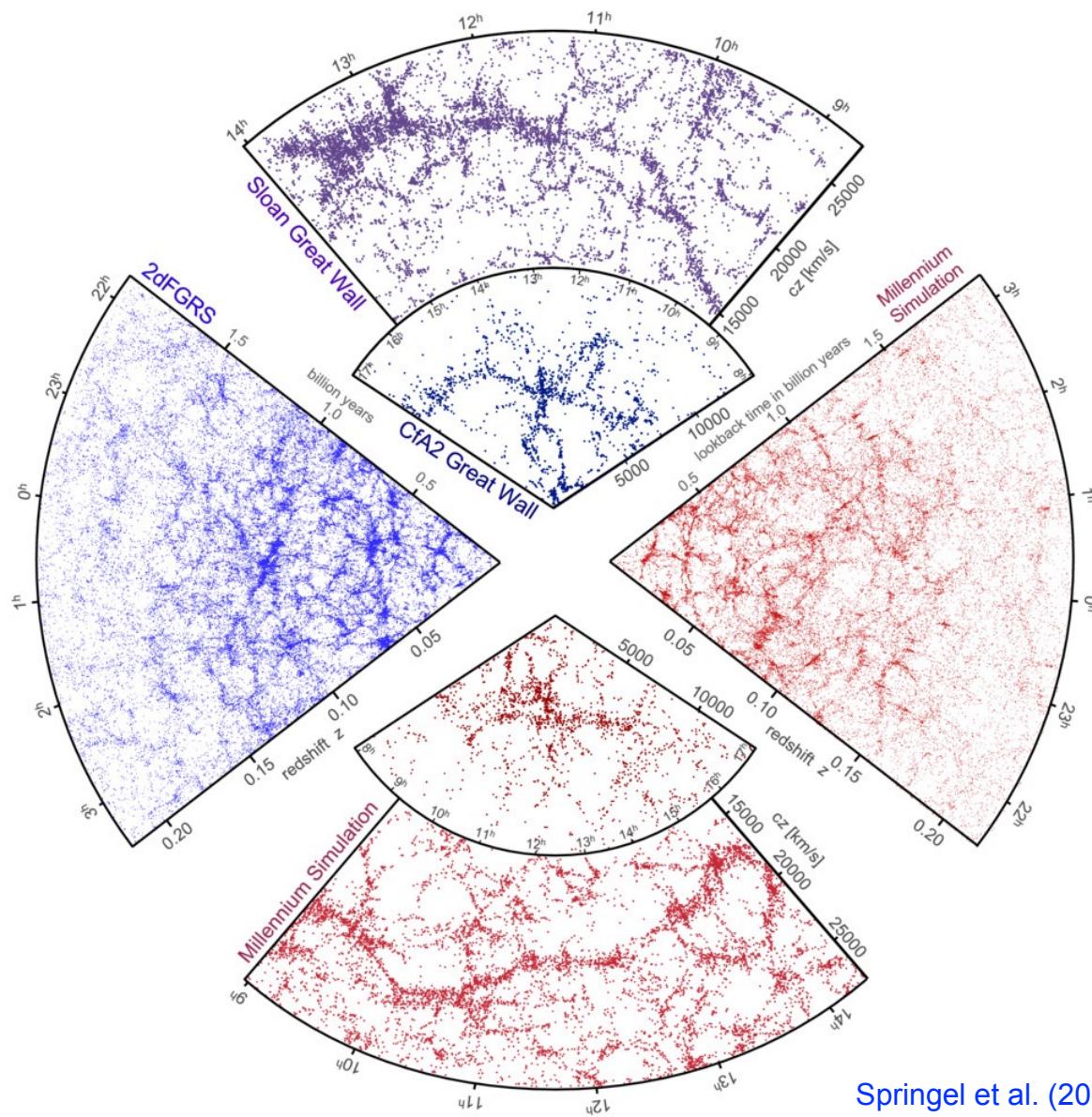
Context

Precision cosmology allows us to perform Bayesian inference on the standard model parameters, and may be discover exotic dark matter particles or modified gravity models.

Non-linear dark matter dynamics plays an important role in various cosmology probes of the dark sector.

- galaxy clustering: non-linear BAO shifts
- weak lensing: non-linear boost.

Large scale structures: towards precision cosmology?



Springel et al. (2005)

Context

Precision cosmology allows us to perform Bayesian inference on the standard model parameters, and may be discover exotic dark matter particles or modified gravity models.

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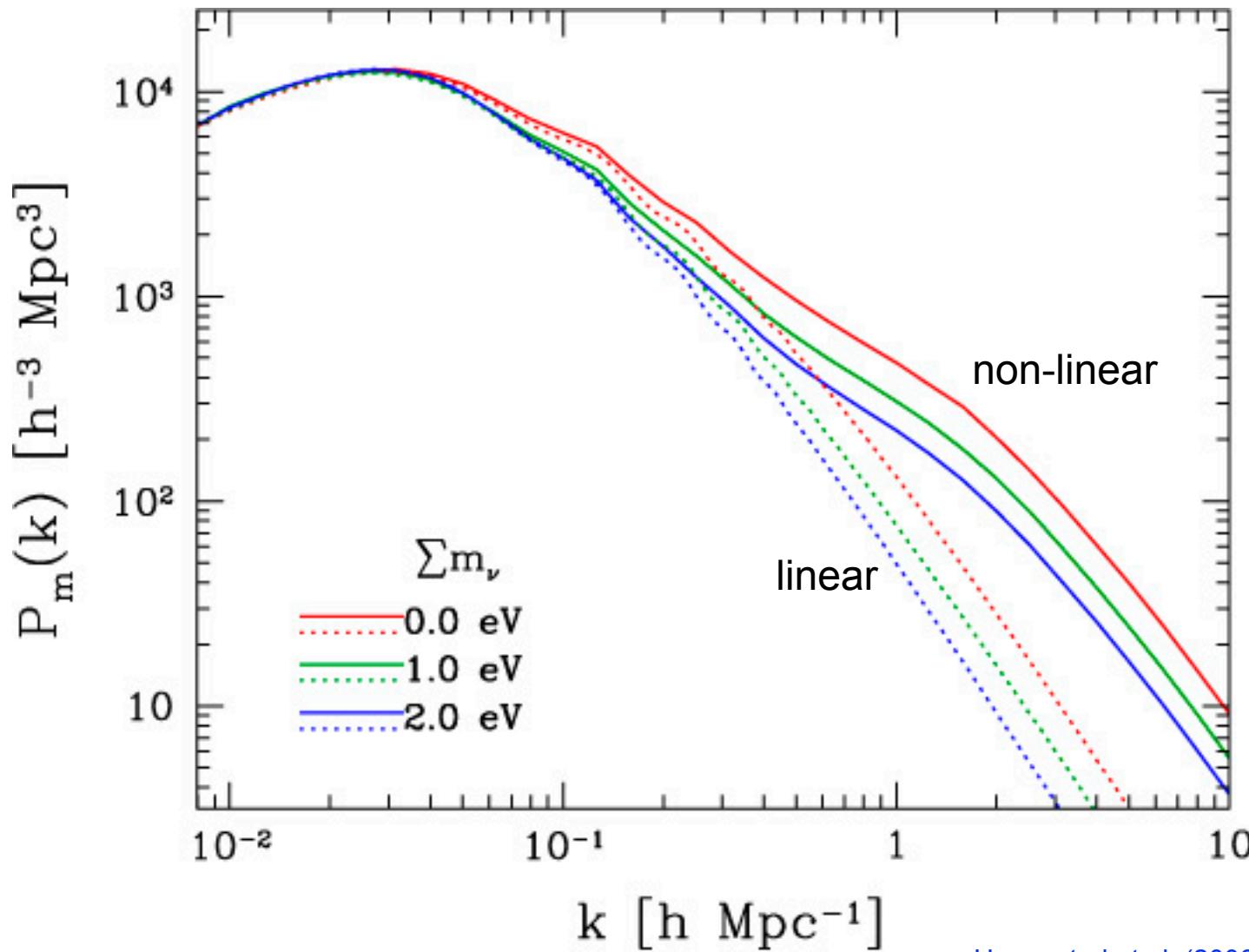
Non-linear dynamics can be described analytically using high-order perturbation theories on relatively large scales.

Small scales are affected by baryonic effects.

N-body models are the only viable alternative at intermediate scales.

How accurate are N-body simulations ? Can we use them to perform include non-linear corrections in expensive Monte Carlo Markov Chains and compute the Likelihood of the best fit parameters ?

Matter power spectrum: a tool to discriminate models



Hannestad et al. (2006)

Computing non-linear dark matter dynamics

N body codes solve the Vlasov-Poisson equations using a Lagrangian sampling of phase space. So far, the most competitive approach, but for how long ?

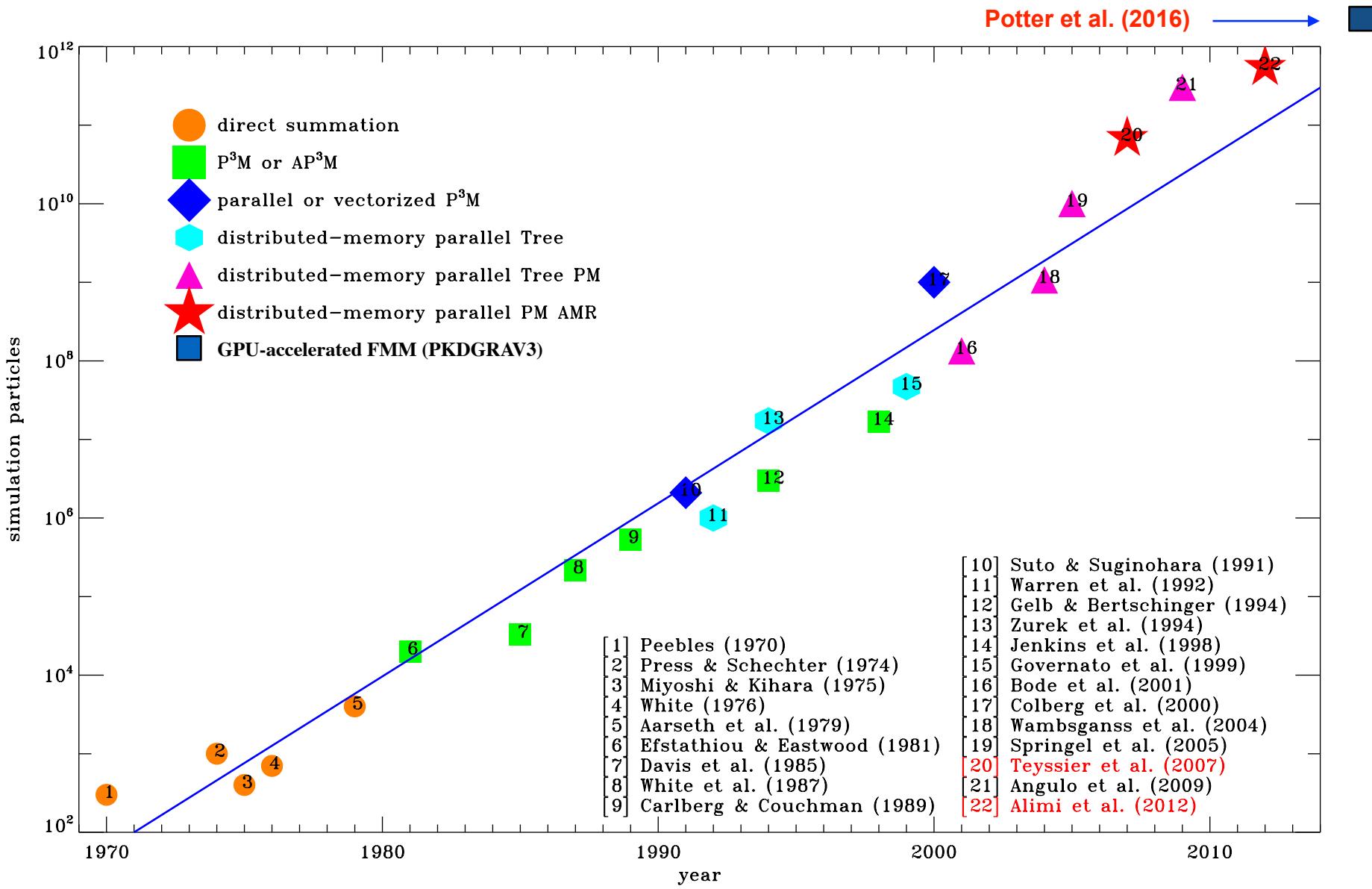
Main numerical limitations are:

1. Finite box size
2. Finite number of particles
3. Finite force resolution

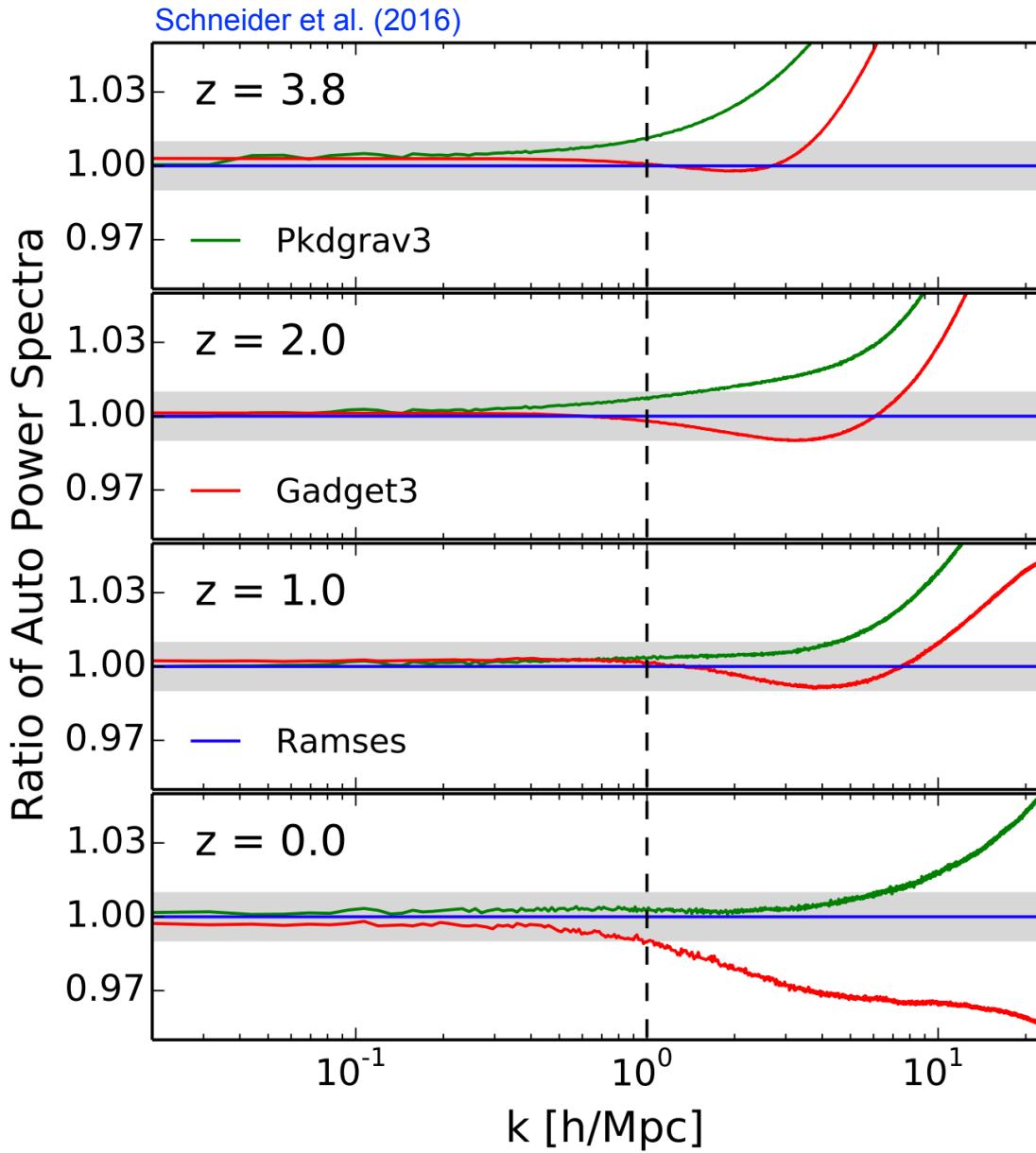
Popular techniques are, in chronological order:

1. Direct N body, scaling as N^2
2. PM: Fast Fourier Transform solvers, $N \log N$, low resolution
3. P3M (PP + PM): order $N \log N$ if large box, N^2 if small box, low resolution
4. Tree codes, $O(N \log N)$, high resolution, see also Tree-PM
5. Adaptive Mesh Refinement (AMR) and Multigrid solver, $O(N)$, high resolution
6. Fast Multipole Method (FMM), $O(N)$, high resolution

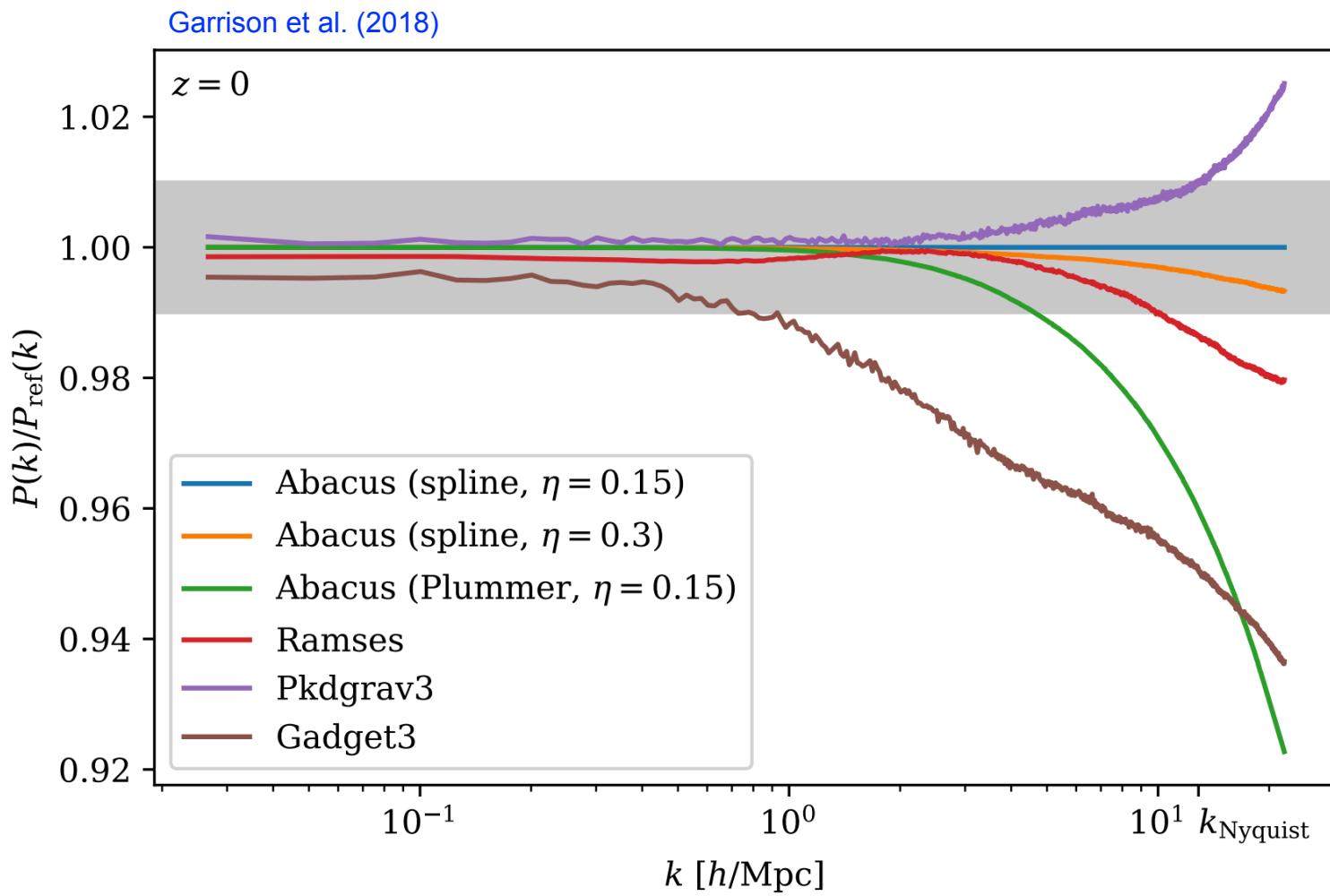
Performance of N body codes



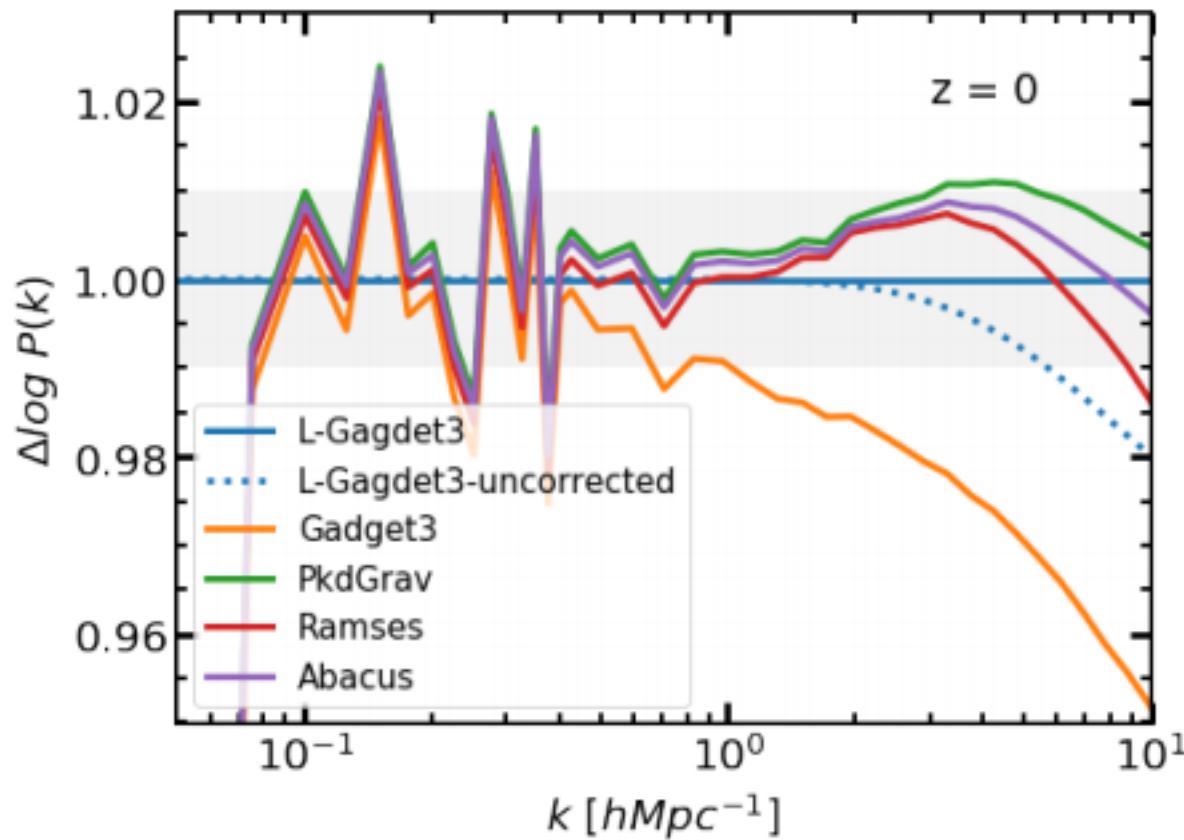
Systematic errors in N body codes ?



Systematic errors in N body codes ?

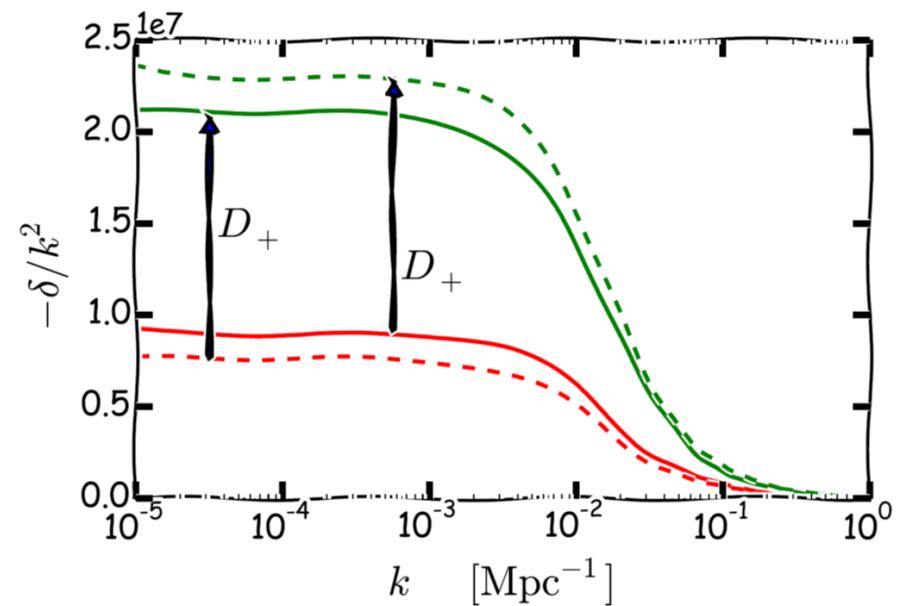
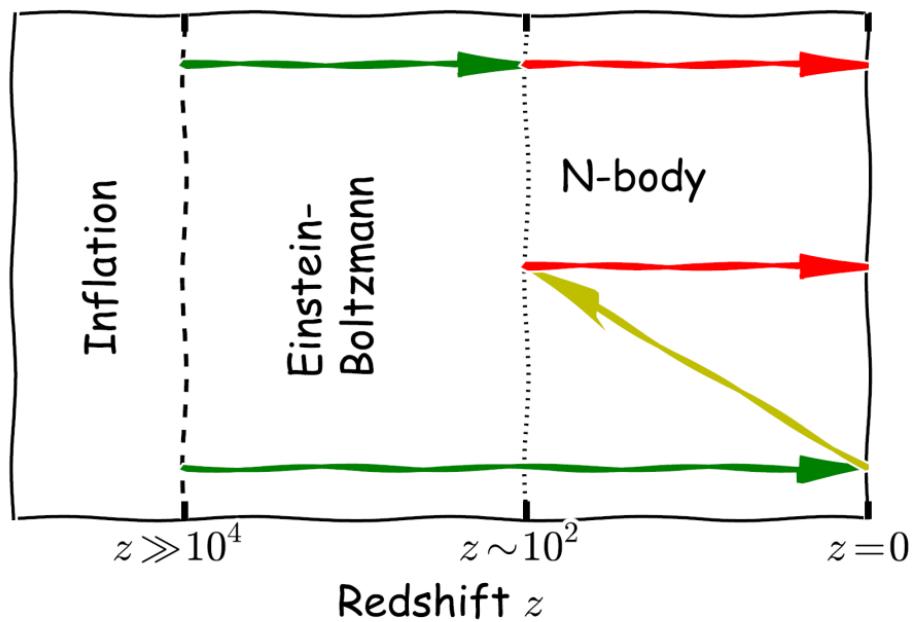


Systematic errors in N body codes ?

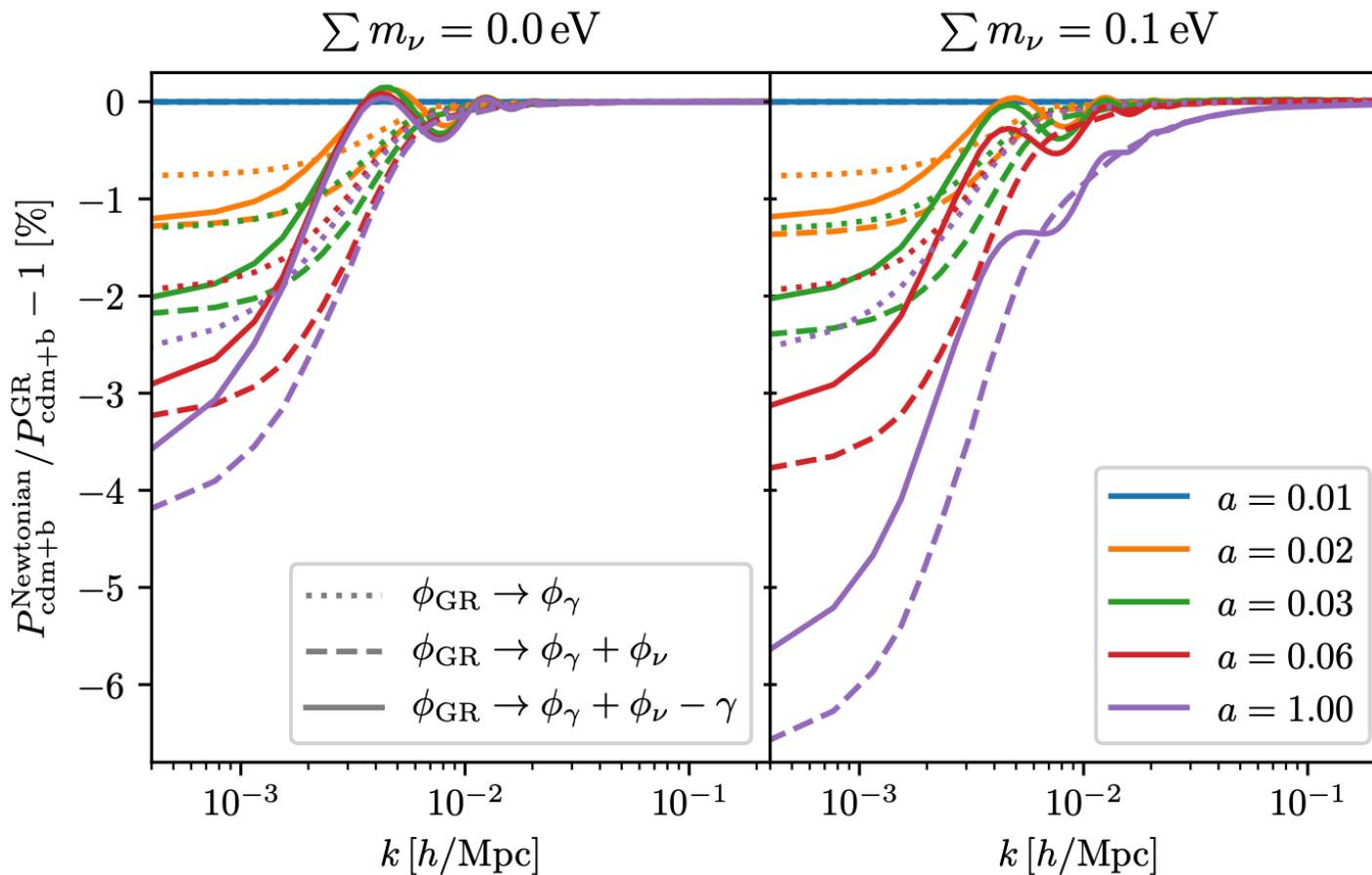


Back-scaled or forward initial conditions?

Fidler et al. 2017



Evolution of GR corrections over time

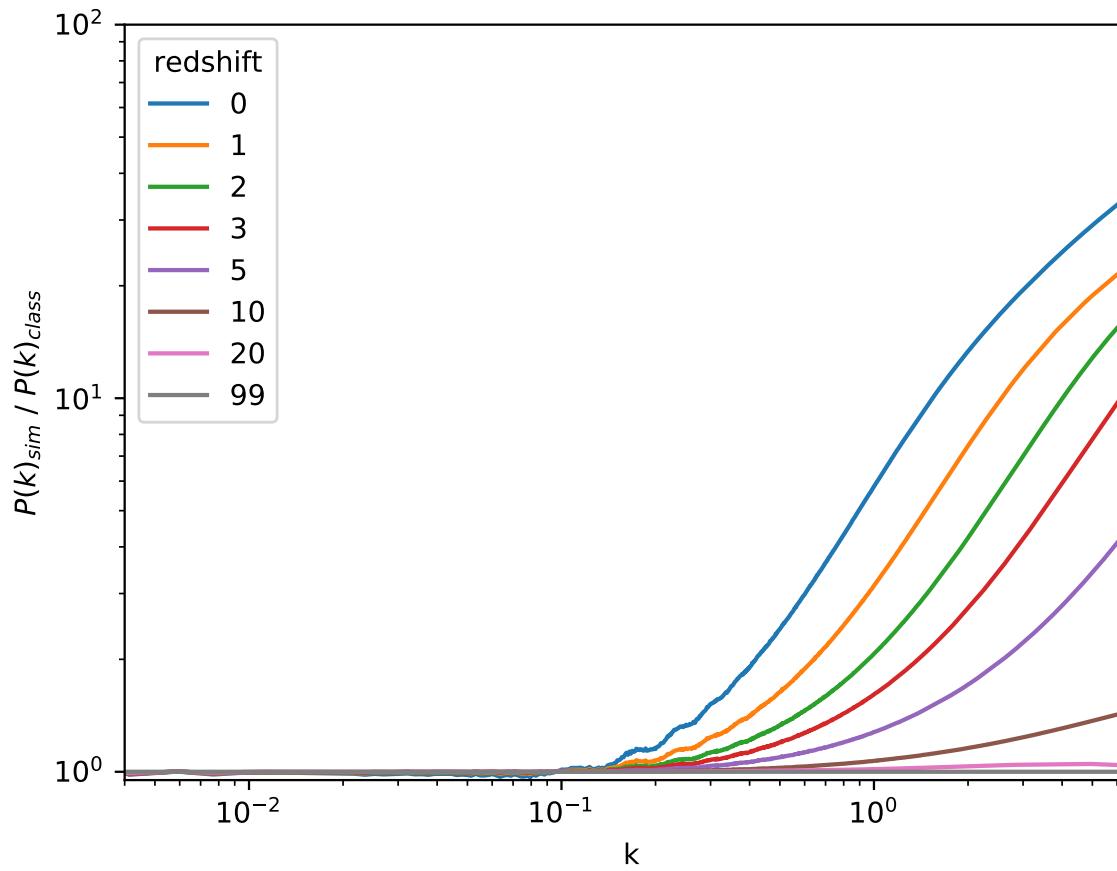


N-body gauge with photons, neutrinos and metric linear fluctuations generated with CLASS

Tram *et al.* 2018

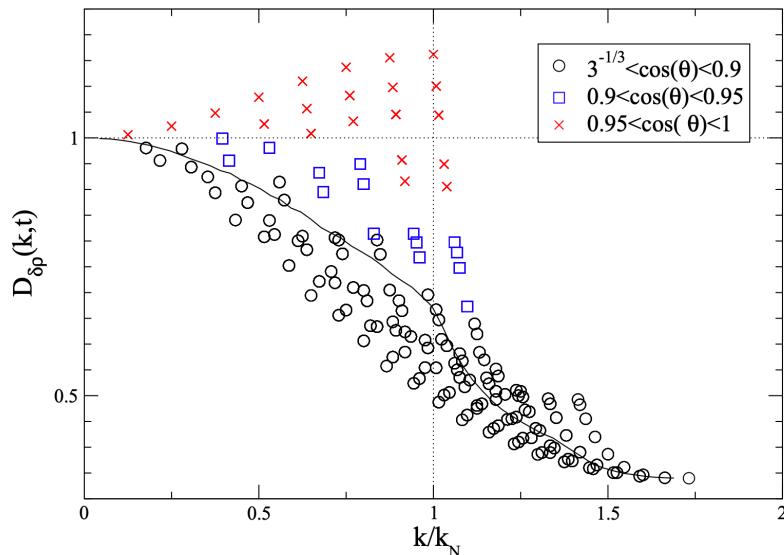
Bridging linear and non-linear scales

Flagship 2 power spectrum non-linear correction

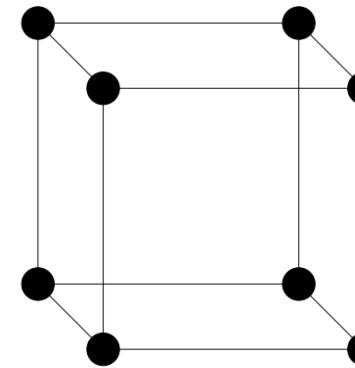


Particle discreteness effects

Joyce et al. 2008



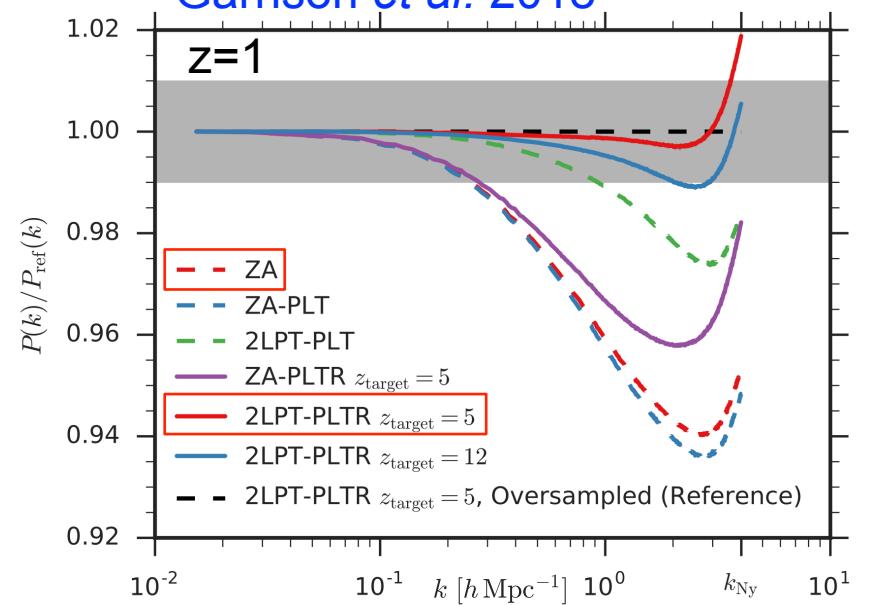
simple cubic lattice



Particle linear theory can be used to rescale the initial conditions and correct from discreteness effects.

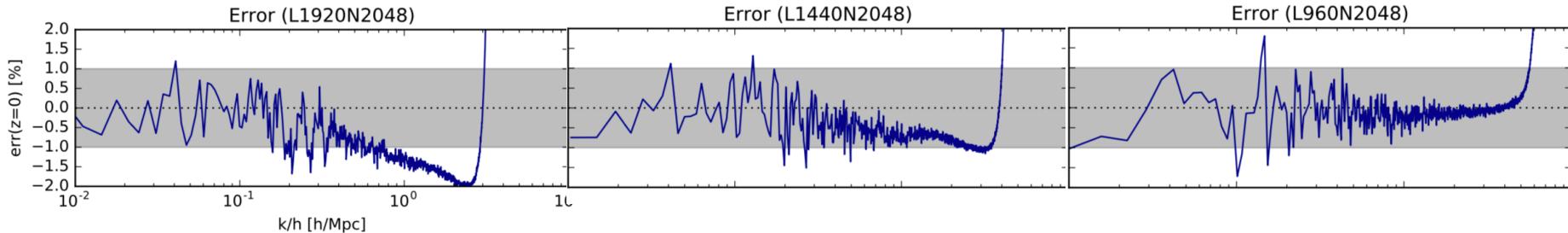
Sub-percent accuracy down to the Nyquist frequency

Garrison et al. 2018

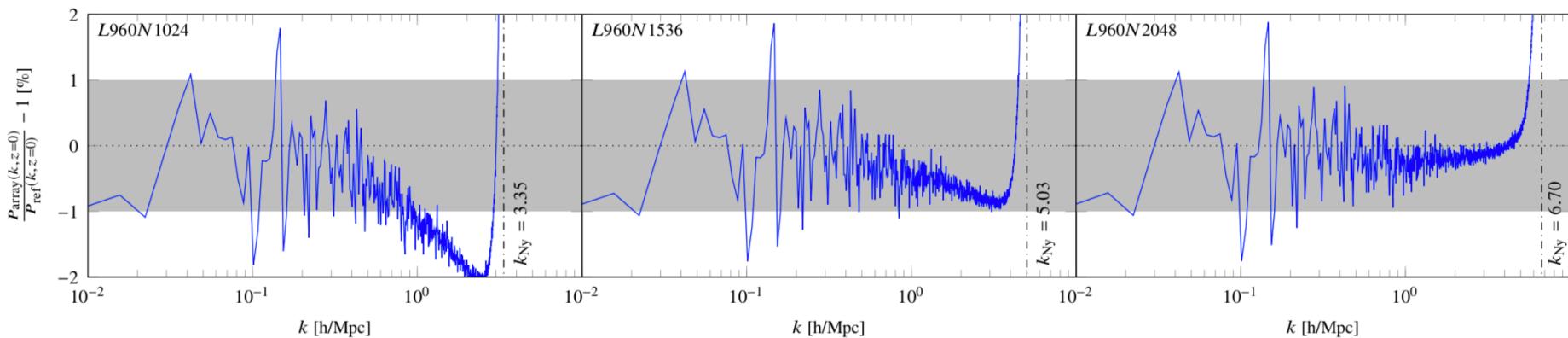


Designing the optimal N body simulation for P(k)

Effect of box size (cosmic variance and non-linear finite volume effects)



Effect of particle number

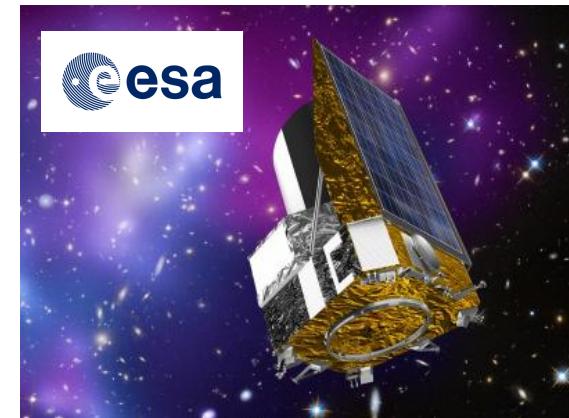


1. Box size larger than 2000 Mpc/h
2. N larger than 4096^3 particles
3. For each parameter set, use 2 simulations with pairing and fixing (Angulo and Pontzen 2016) to suppress variance on large scale

Emulating the power spectrum for Euclid

In order to fit galaxy survey data using non-linear corrections, one needs to explore the 6-dimensional parameter space to perform the likelihood analysis.

Each individual simulation is very expensive.



Goal: replace N body simulations by a surrogate model with the same sub-percent accuracy.

Uncertainty quantification techniques used routinely in engineering.

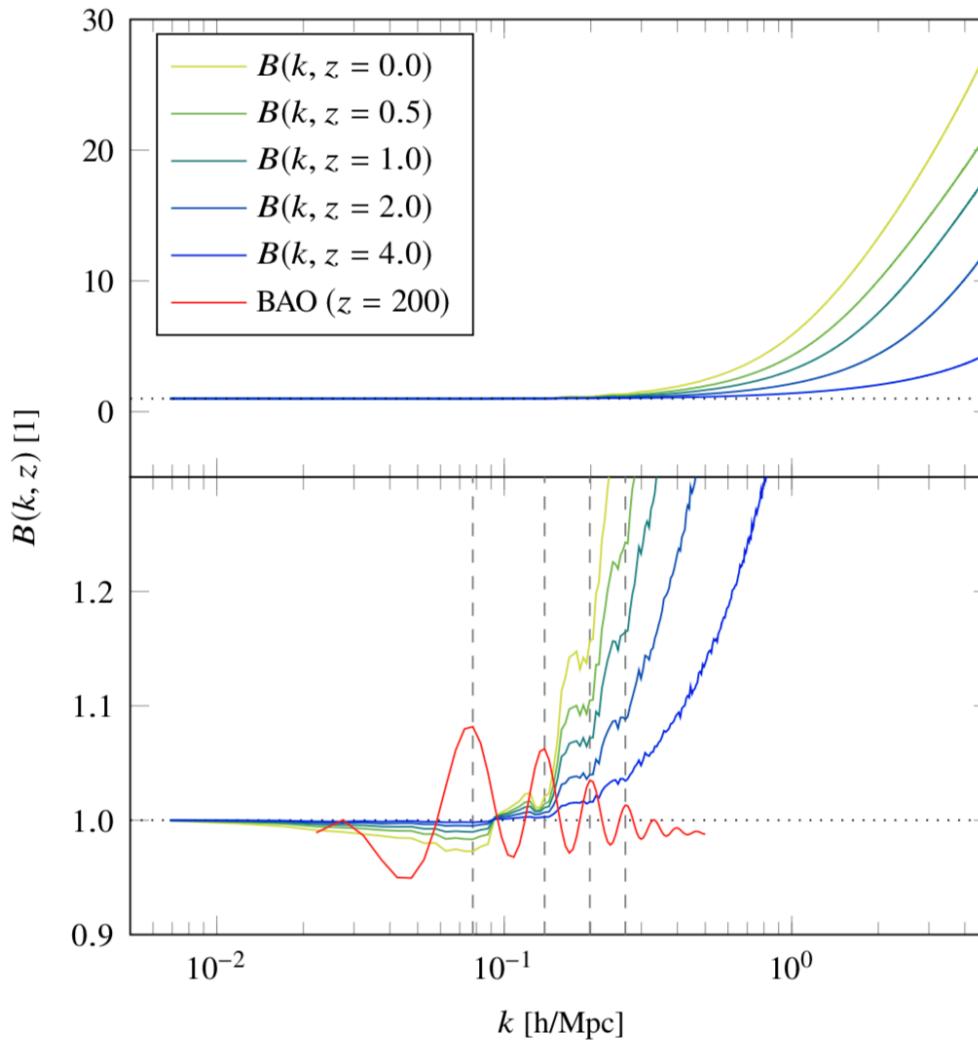
First implementation for cosmology by [Heitmann et al. \(2010\)](#): CosmicEmu.

This work: used the platform UQLab developed at ETH Zurich by Sudret and collaborators to design new, more accurate emulators for the Euclid mission.
[The Euclid collaboration, Knabenhans et al. \(2018\), arXiv:1809.04695](#): The EuclidEmulator

[Knabenhans et al. \(2020\)](#) EuclidEmulator2, under review

New approach: emulate the non-linear boost

$$B(k, z) = \frac{P_{\text{non-lin}}(k, z)}{P_{\text{lin}}(k, z)}$$



Optimising the experimental design

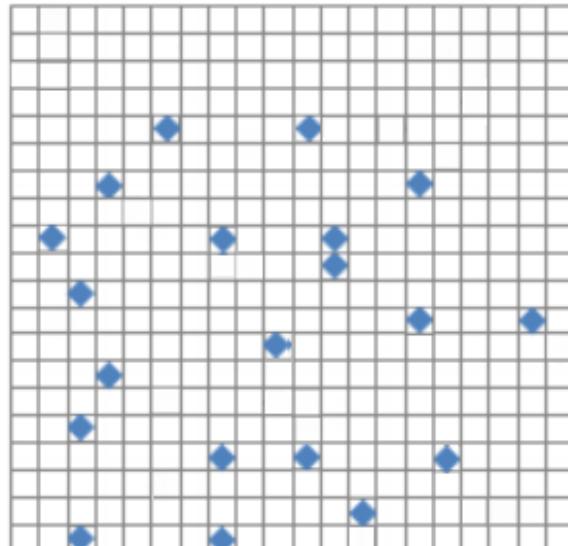
How many simulations do we need to perform ?

How do we sample the parameter space ?

How do we interpolate the power spectrum in between the sampling points ?

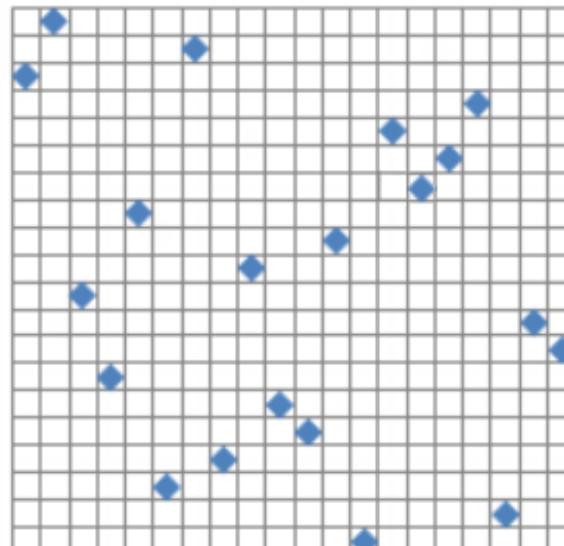
Monte Carlo Sampling

(Some shared rows or columns.)



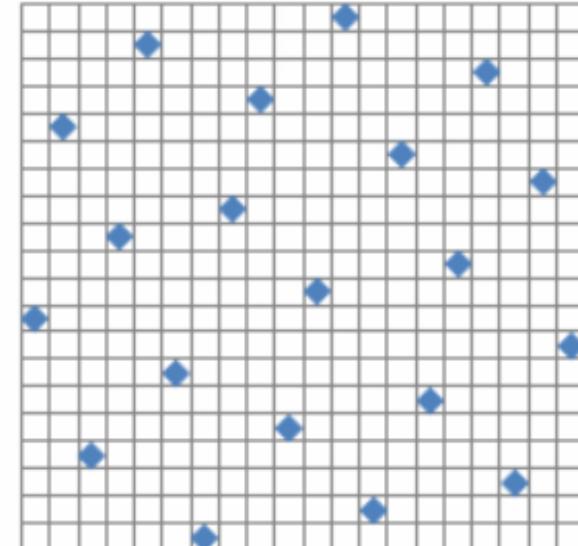
Latin Hypercube Sampling

(Monte Carlo Sampling with no shared rows or columns.)



Optimal Space-Filling Design Sampling

(Latin Hypercube sampling with even distribution of points)



Interpolation strategy

First, define the parameter space

Second, define n_{ED} , the number of points in the experimental design.

Third, perform a Principal Component Analysis

$$\mathbf{D} = \sum_{i=1}^{n_{ED}} \lambda_i(\omega_b, \omega_m, n_s, h, w_0, \sigma_8) \mathbf{PC}_i(k, z)$$

$$\omega_b \in [0.0215, 0.0235]$$

$$\omega_m \in [0.1306, 0.1546]$$

$$n_s \in [0.9283, 1.0027]$$

$$h \in [0.6155, 0.7307]$$

$$w_0 \in [-1.30, -0.70]$$

$$\sigma_8 \in [0.7591, 0.8707]$$

Fourth, interpolate the eigenvalues with Legendre polynomials of degree p .

$$\lambda_i(\omega_b, \omega_m, n_s, h, w_0, \sigma_8) \approx \sum_{\alpha \in \mathcal{A}} \eta_\alpha \Psi_\alpha(\mathbf{x})$$

Polynomial coefficients are computed by regression.

Key parameters:

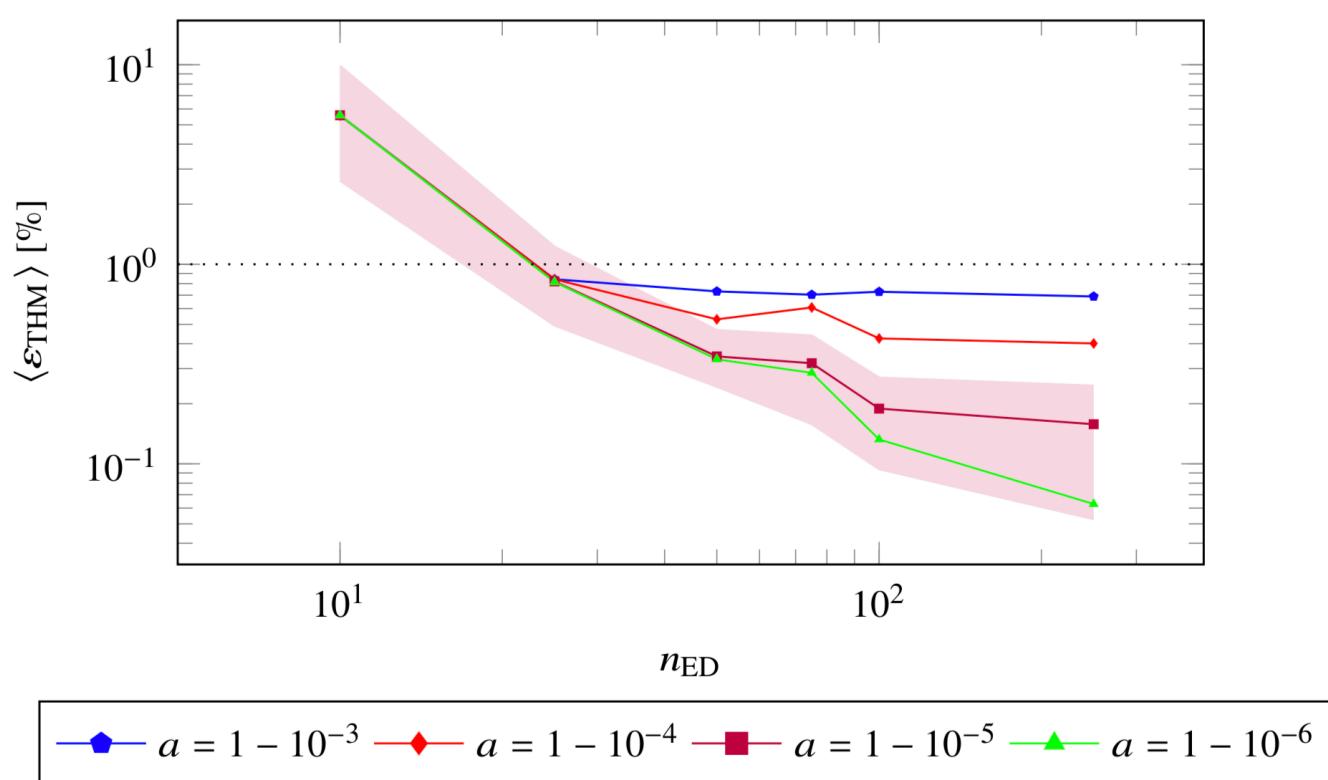
- Number of points in the experimental design n_{ED}
- Number of PCs N_{PC} or accuracy parameter a
- Polynomials degree p

Defining the emulation parameters

We use Takahashi's Halofit (THM) to mock N body simulations.

We estimate the Emulation-Only-Error (EOE).

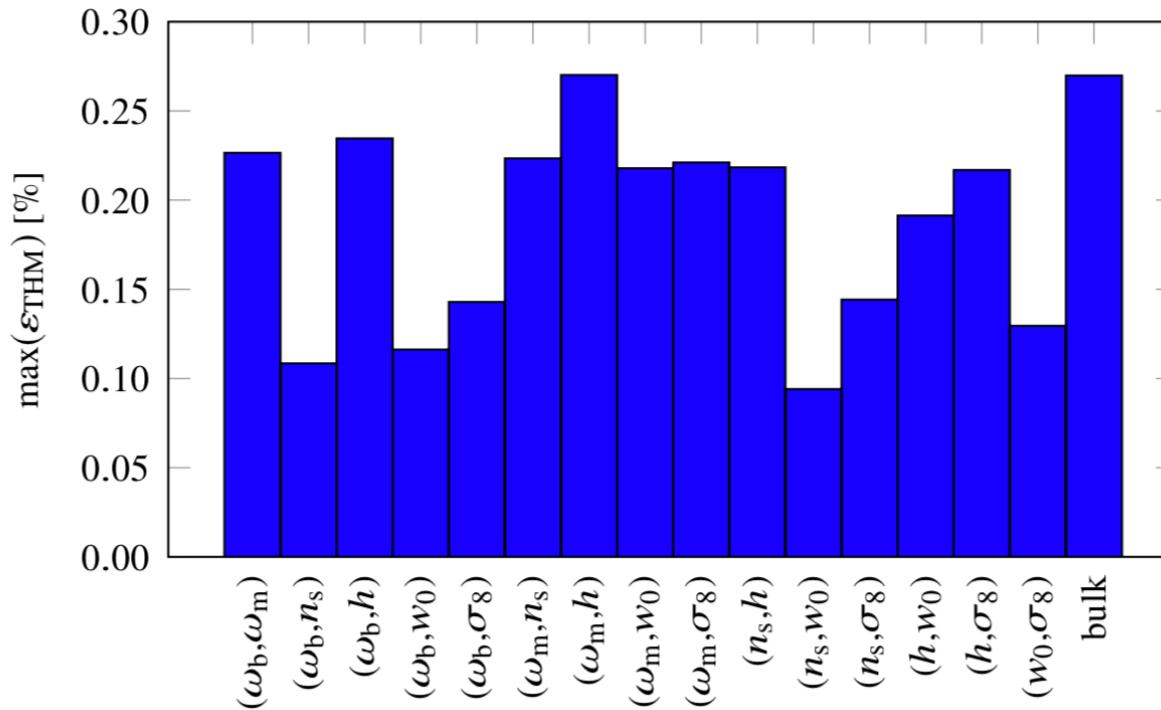
Required number of simulations depends on the PCA truncation parameter.



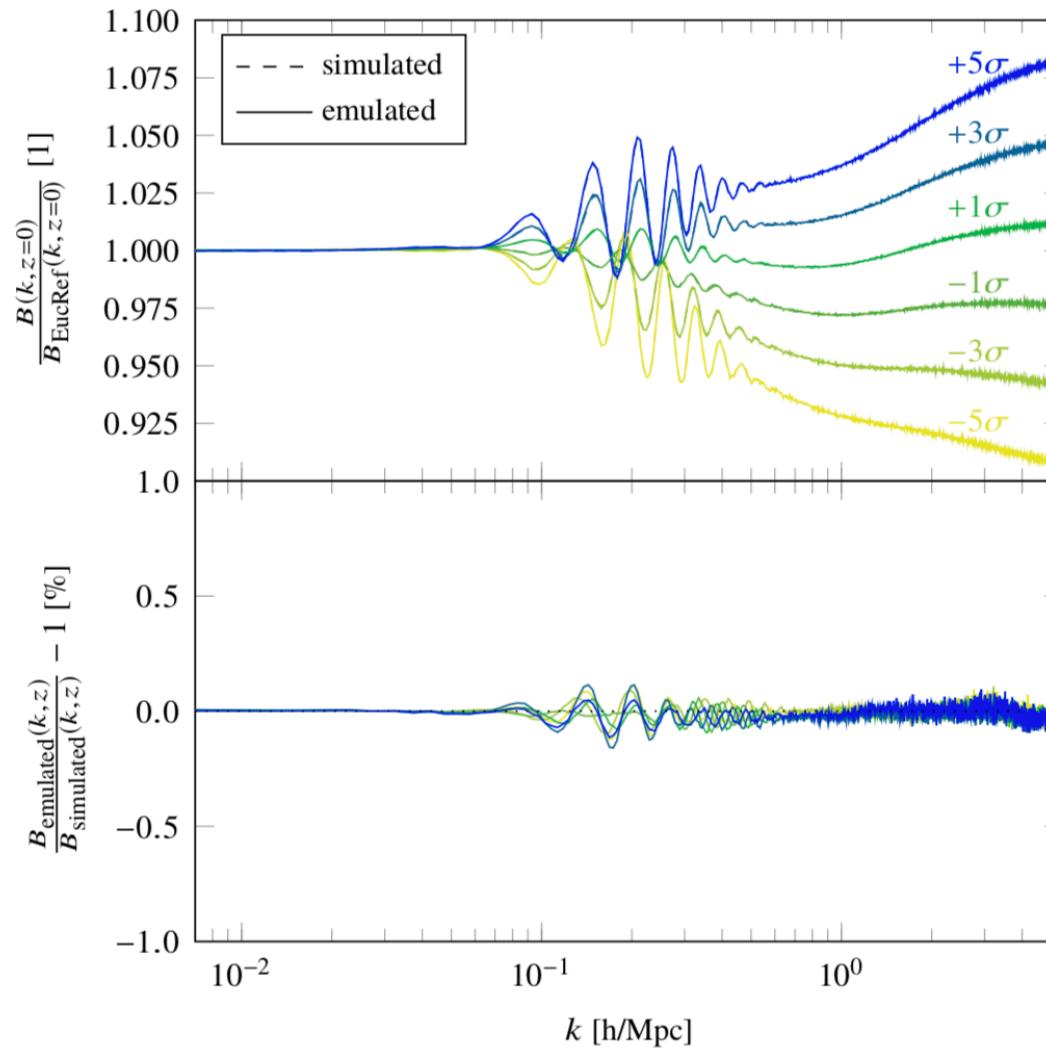
Characterising the final strategy

We set the polynomial degree to p=2, the size of the experimental design to 100 points and the number of PC to 11 ($a=1 \cdot 10^{-5}$).

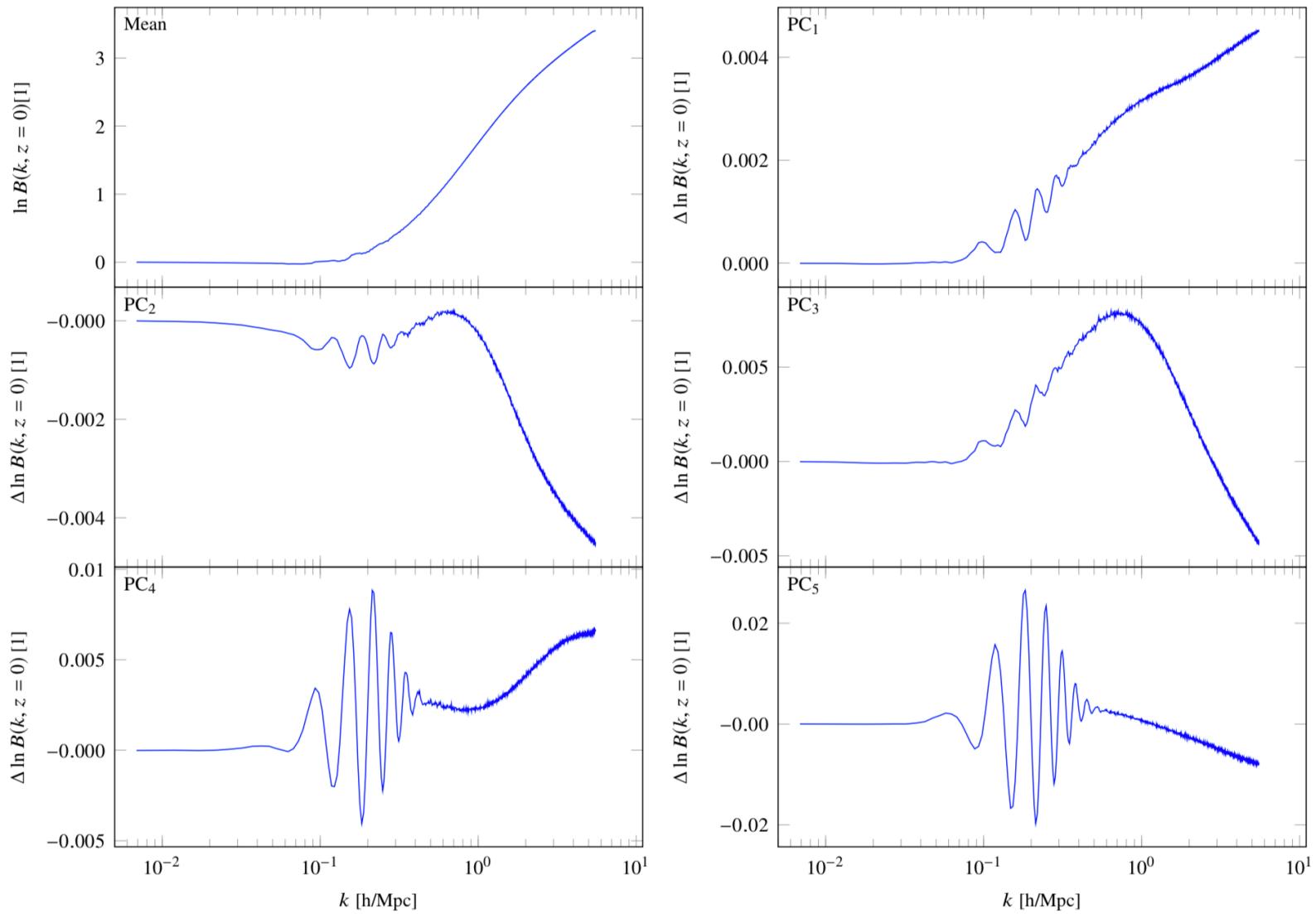
We test the corresponding emulator using Halofit in 2D cuts through the parameter space.



Testing the final Euclid Emulator

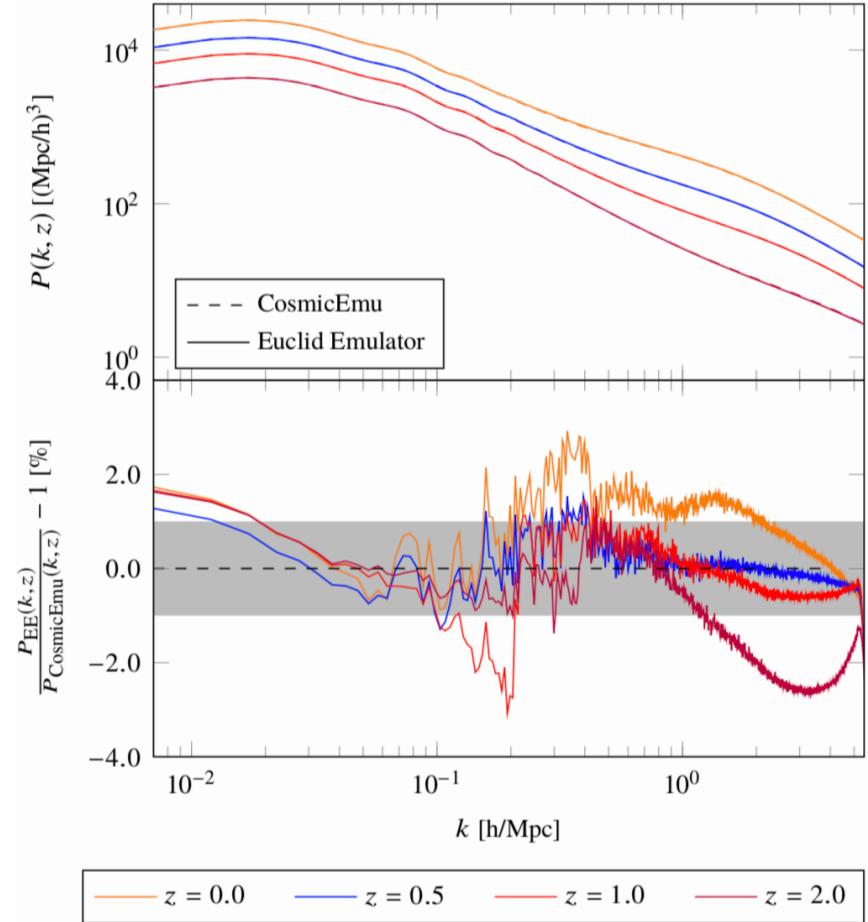
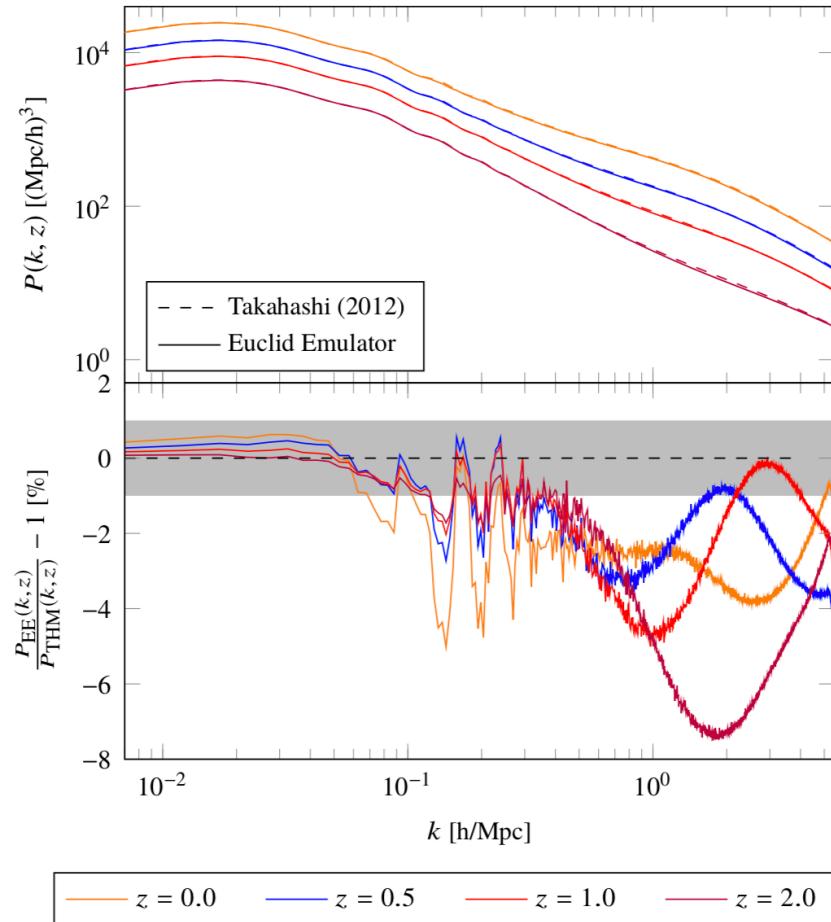


6 first PC in the final emulator



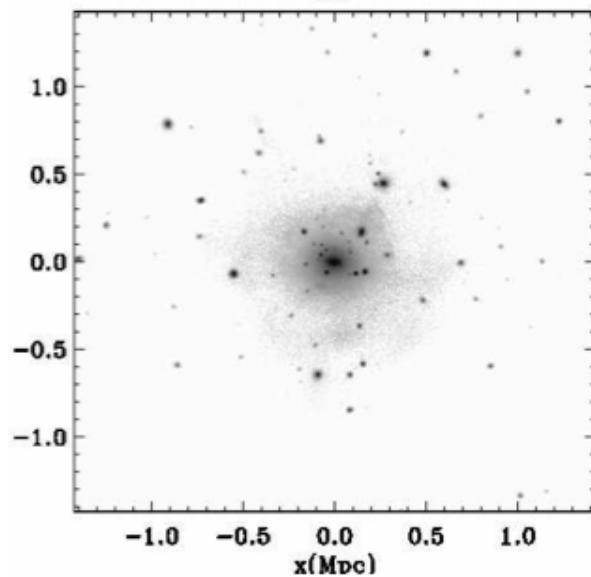
Comparison with previous work

Use the CLASS code to generate the linear power spectrum, then multiply by the non-linear correction (boost).

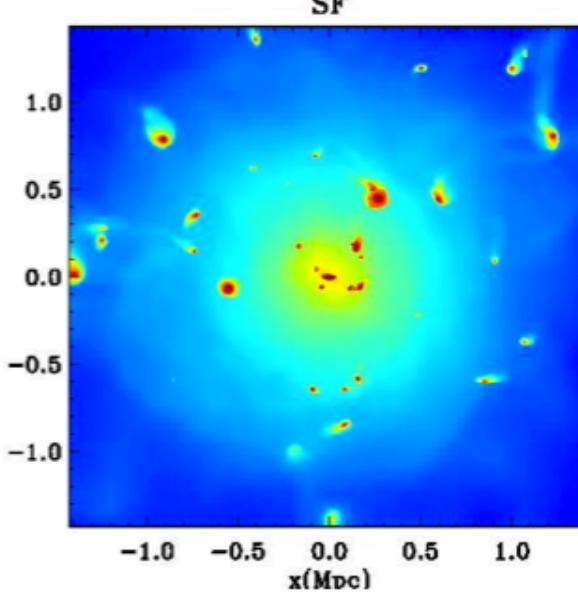


Galaxy formation in groups and clusters

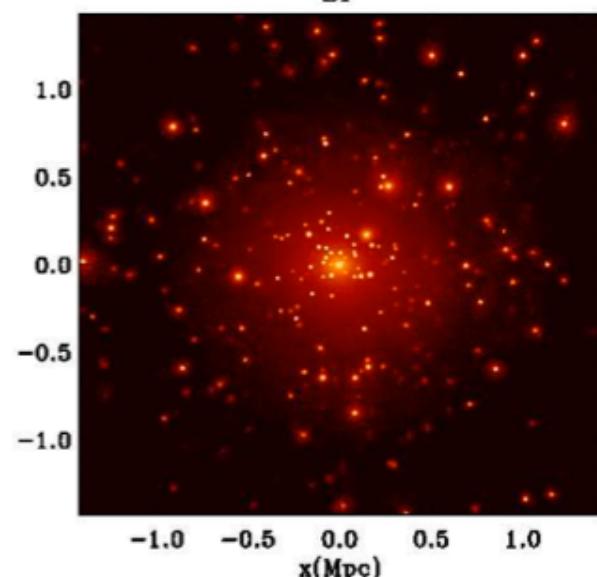
SF



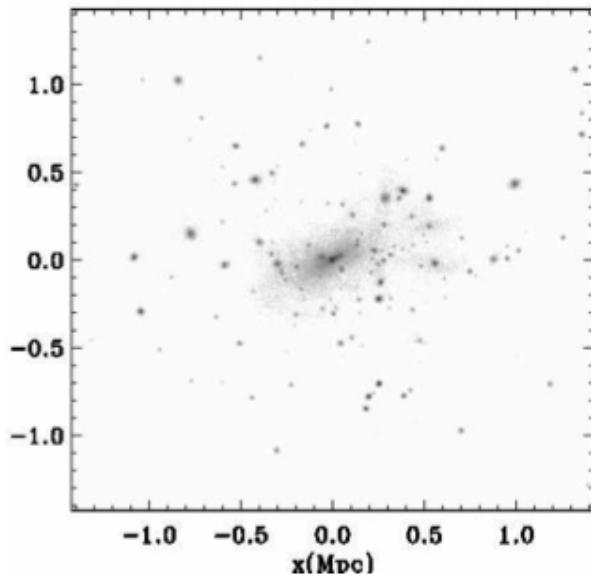
SF



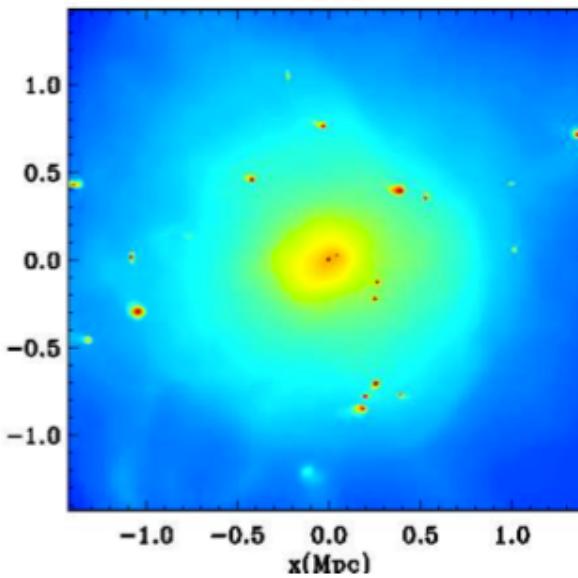
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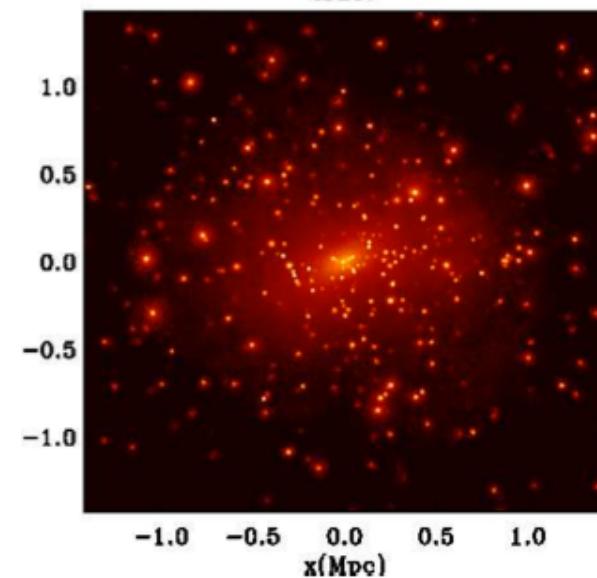
AGN



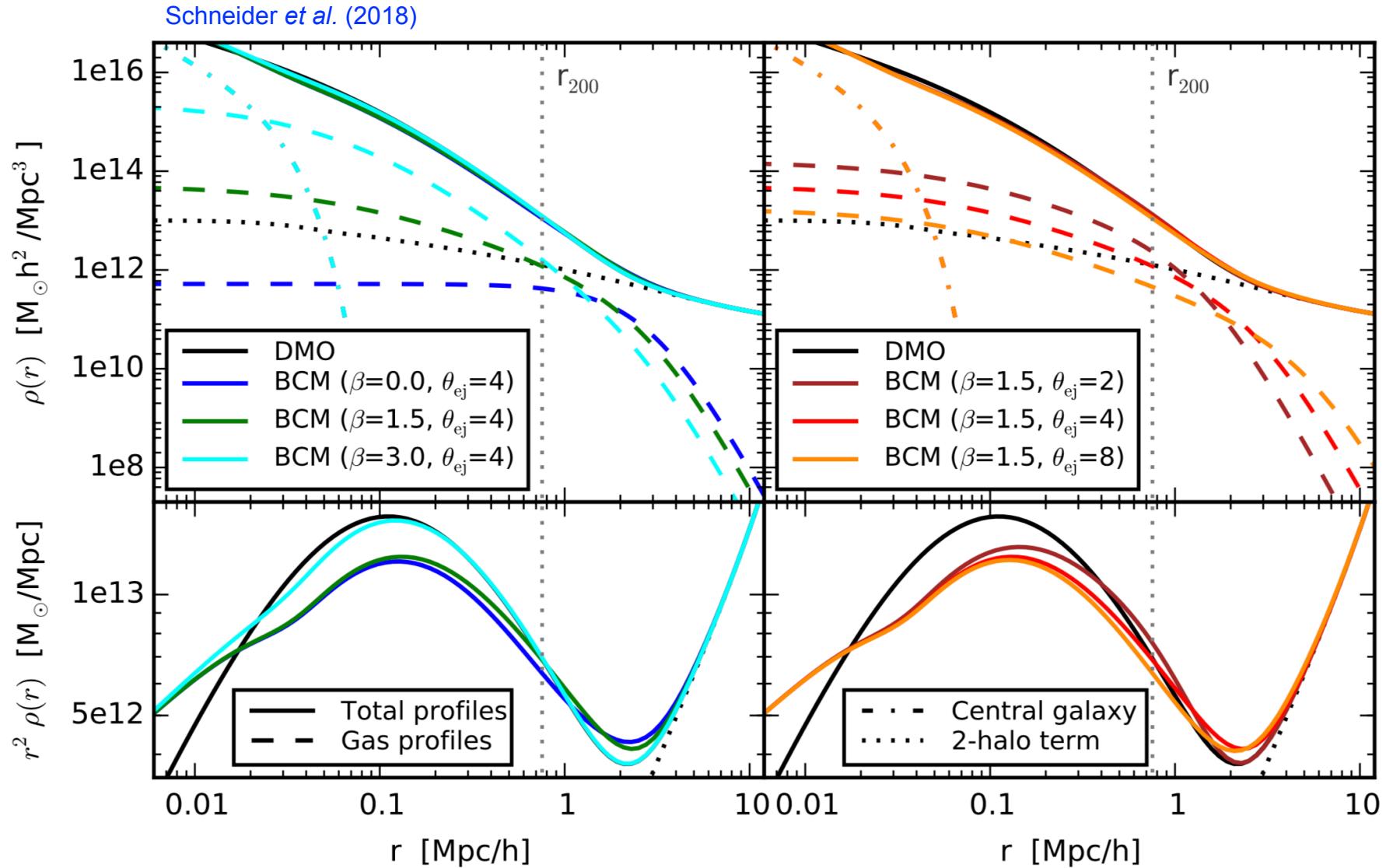
AGN



AGN



The Baryonic Correction Model (Schneider & Teyssier 2015)



Encoding baryonic effects on the power spectrum

$$\rho_{\text{gas}}(r) = \frac{\rho_{\text{gas},0}}{(1+u)^\beta(1+v^2)^{(7-\beta)/2}},$$

$u \equiv r/r_{\text{co}}$ and $v \equiv r/r_{\text{ej}}$.

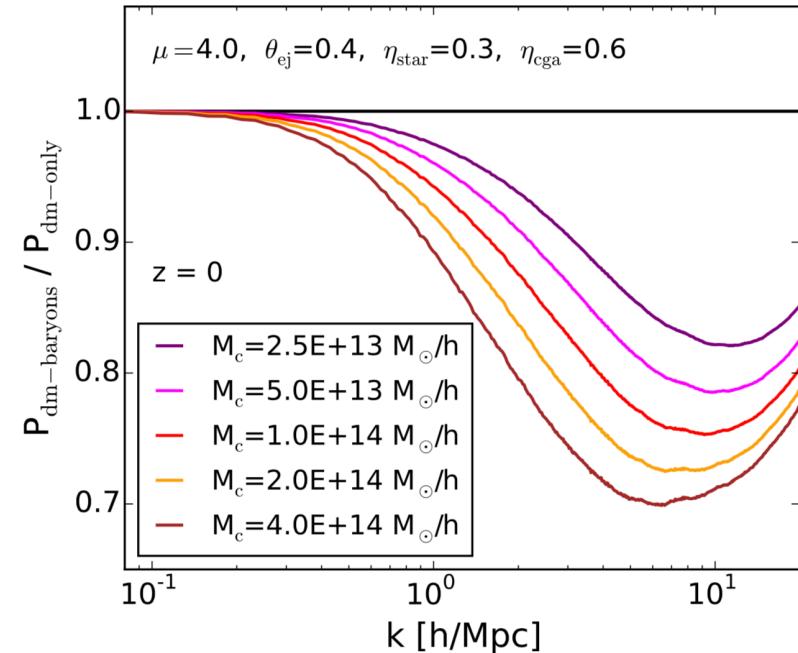
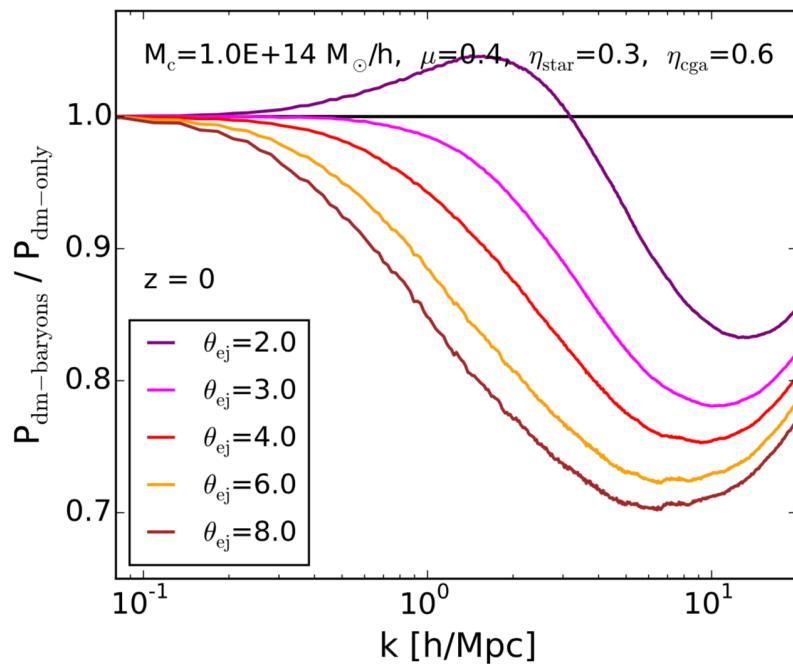
$$r_{\text{co}} \equiv \theta_{\text{co}} r_{\text{vir}}, \quad r_{\text{ej}} \equiv \theta_{\text{ej}} r_{\text{vir}},$$

$$\beta(M_{200}) = 3 - \left(\frac{M_c}{M_{200}}\right)^\mu$$

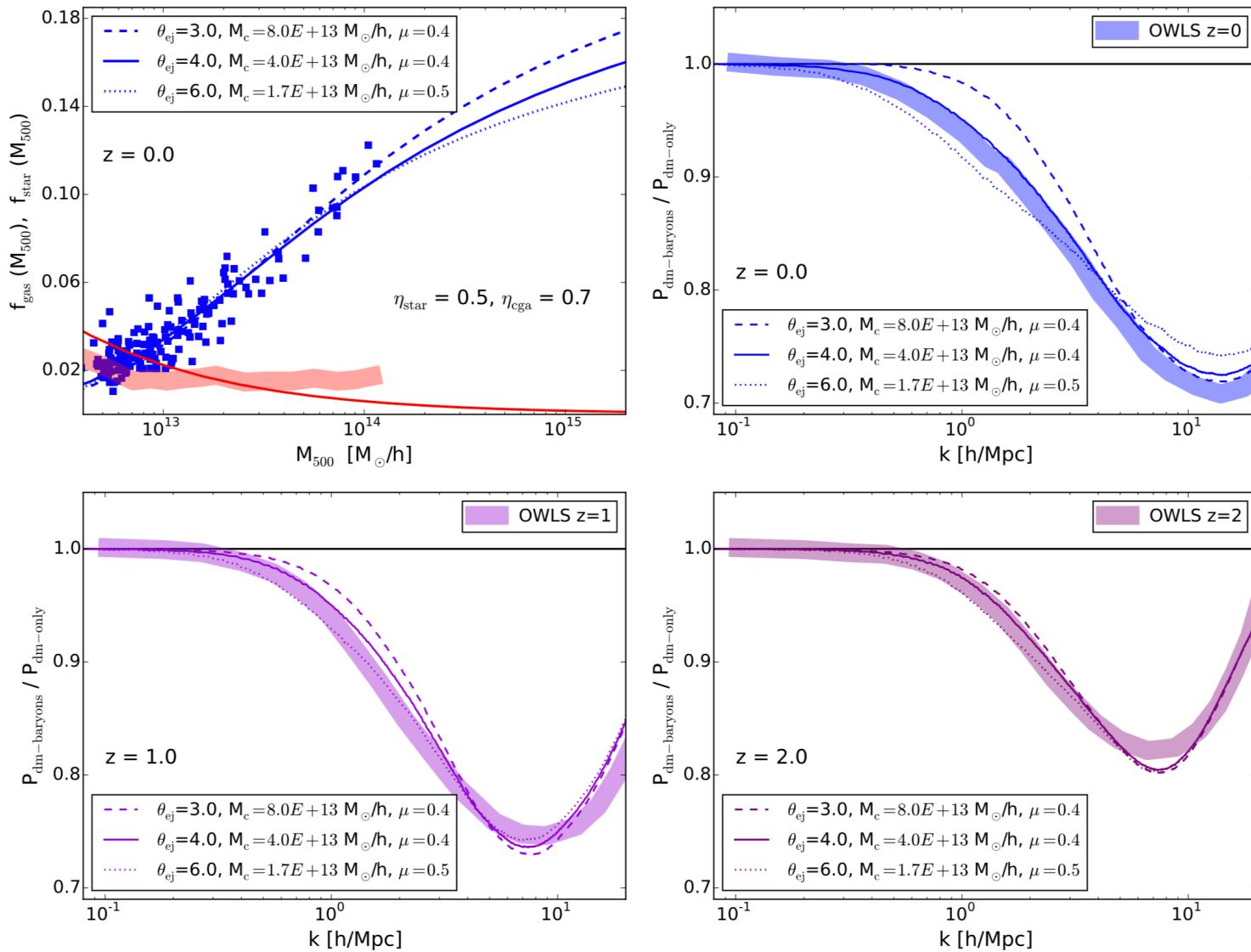
For each halo, we use an analytical spherically symmetric model for the total mass distribution.

Particle are displaced according to the mapping from pure NFW to BCM.

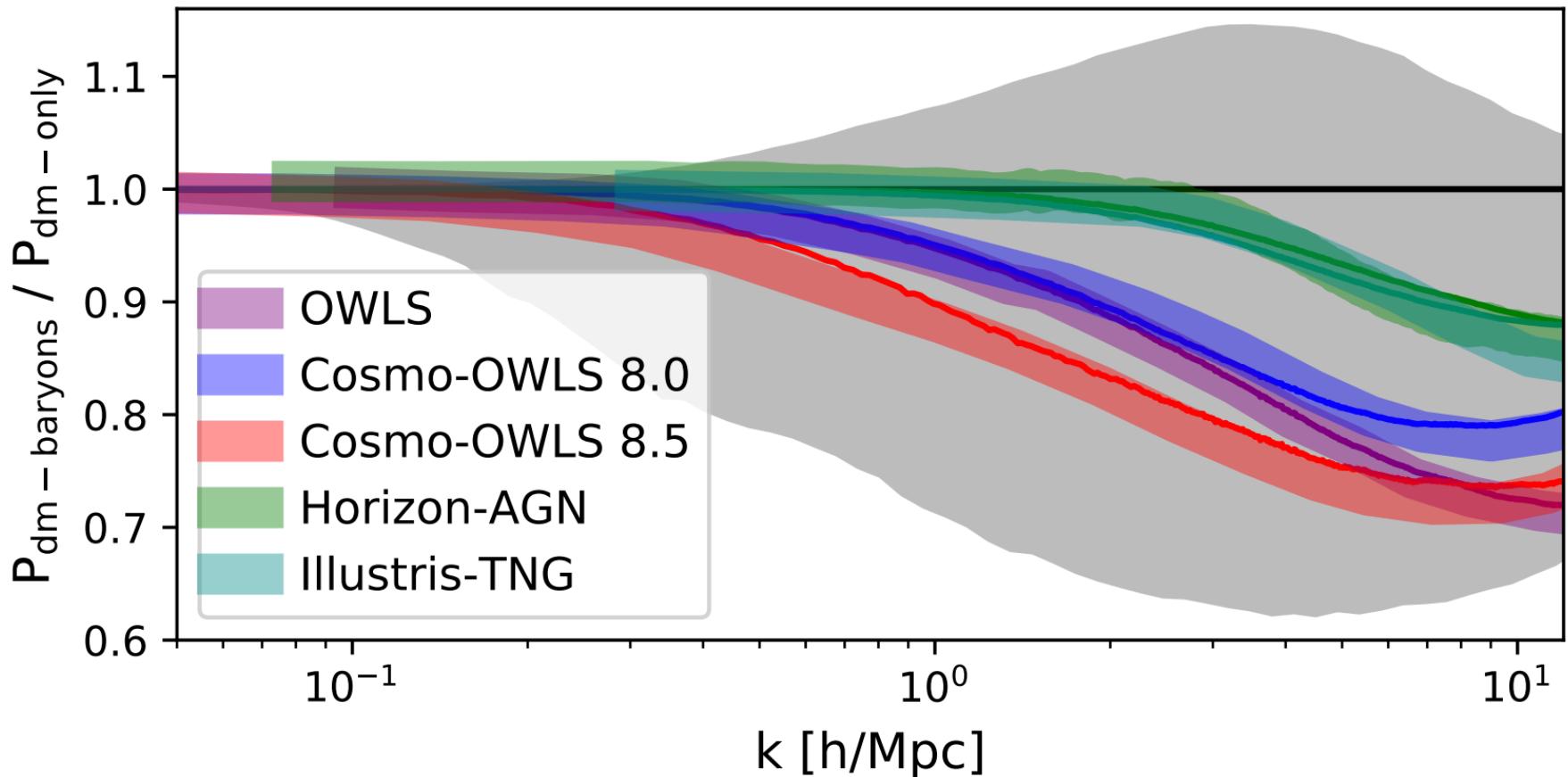
We then compute the baryonic correction on the power spectrum.



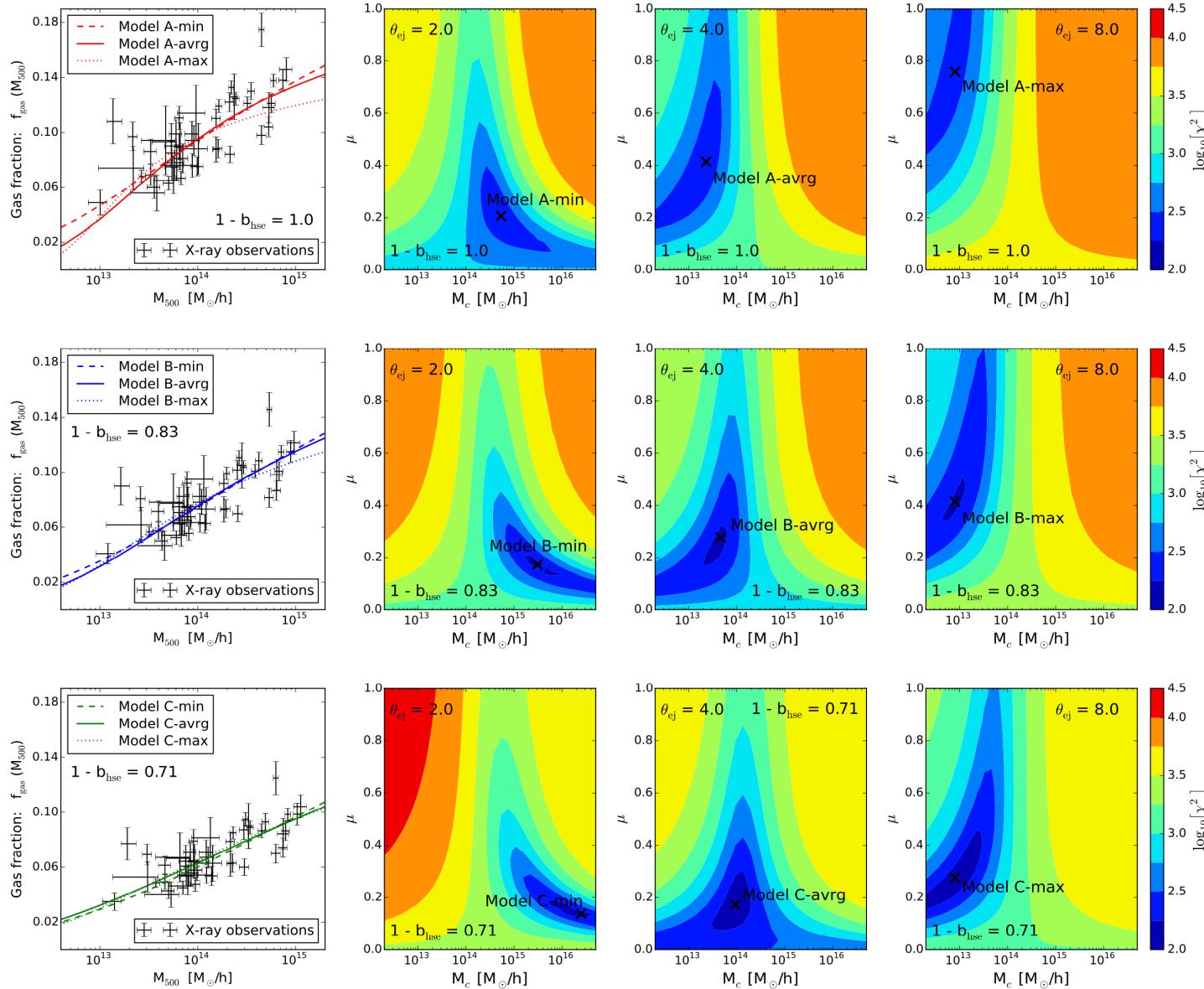
Matching the OWLS simulation (McCarthy et al. 2010)



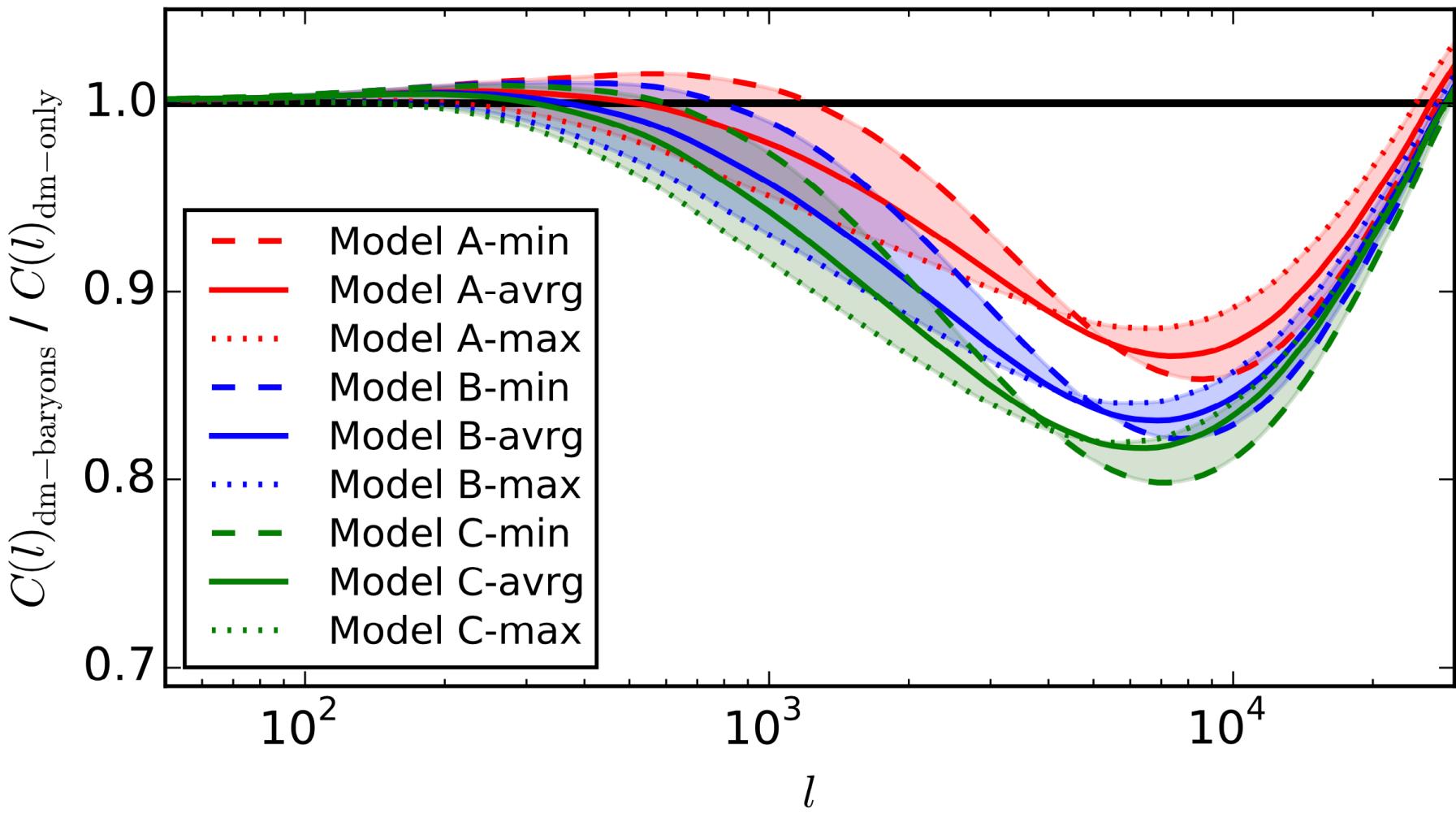
Matching other galaxy formation simulations



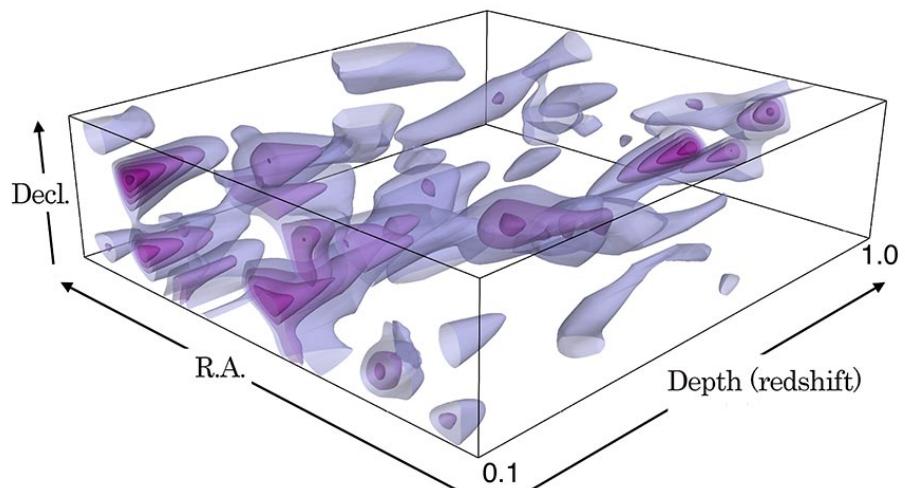
Matching real X-ray data with the BCM



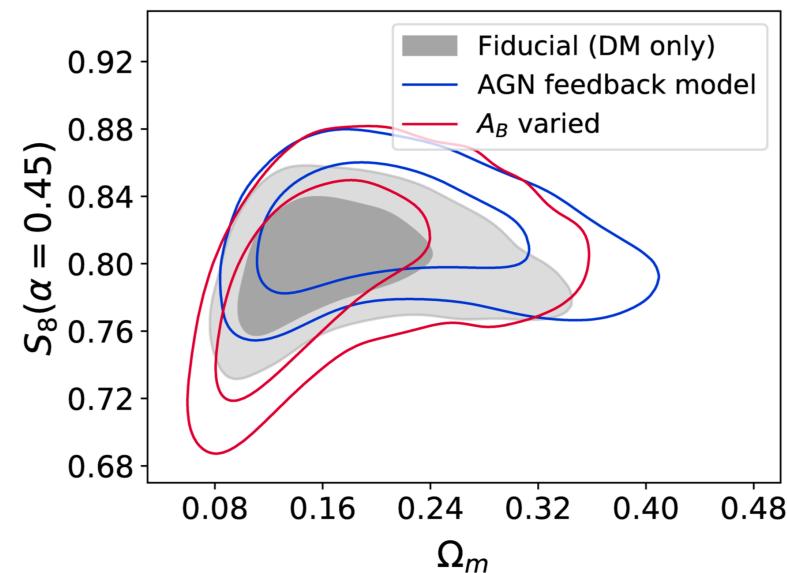
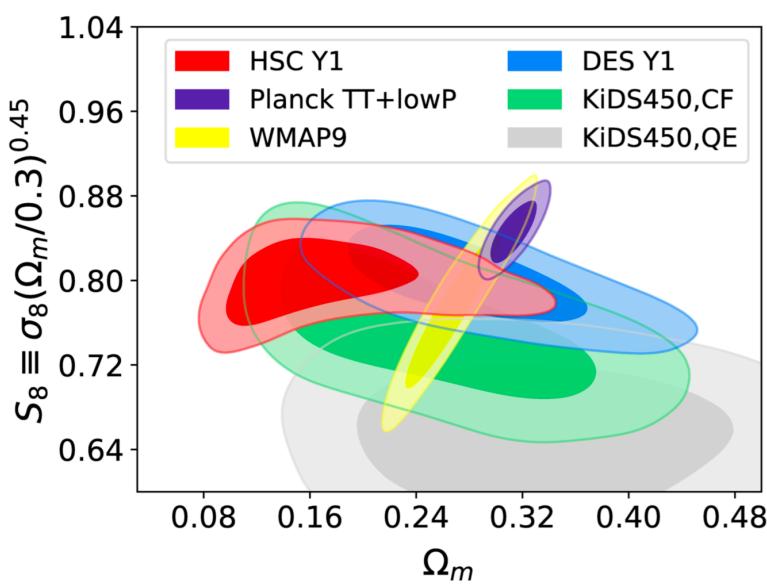
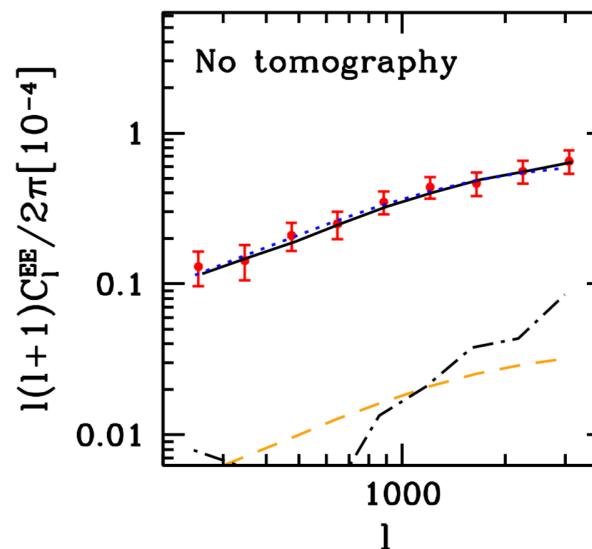
Baryonic corrections to weak shear power spectrum



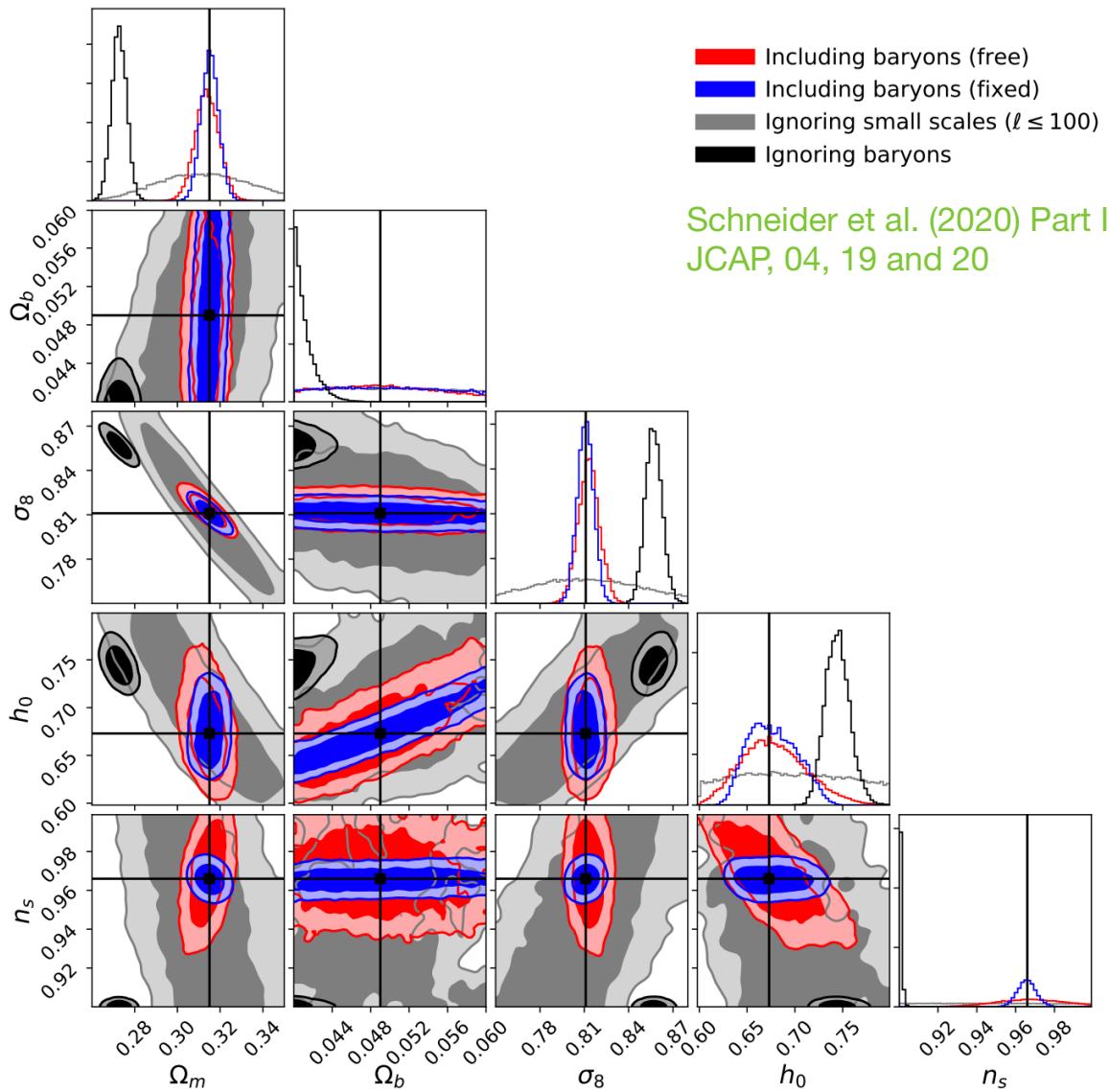
Interpreting current weak lensing surveys



Credit: University of Tokyo/NAOJ

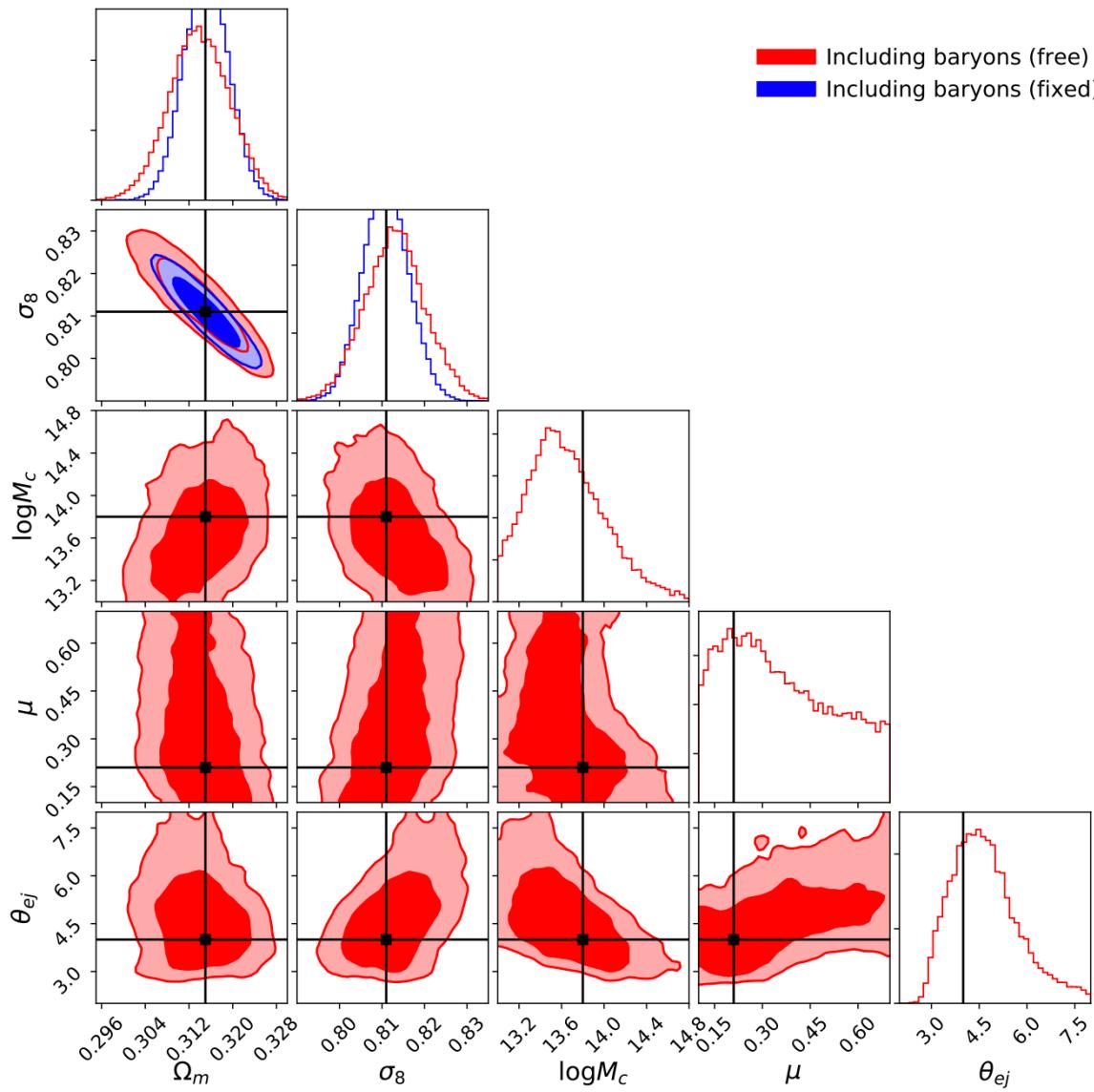


Forecasting future weak lensing surveys

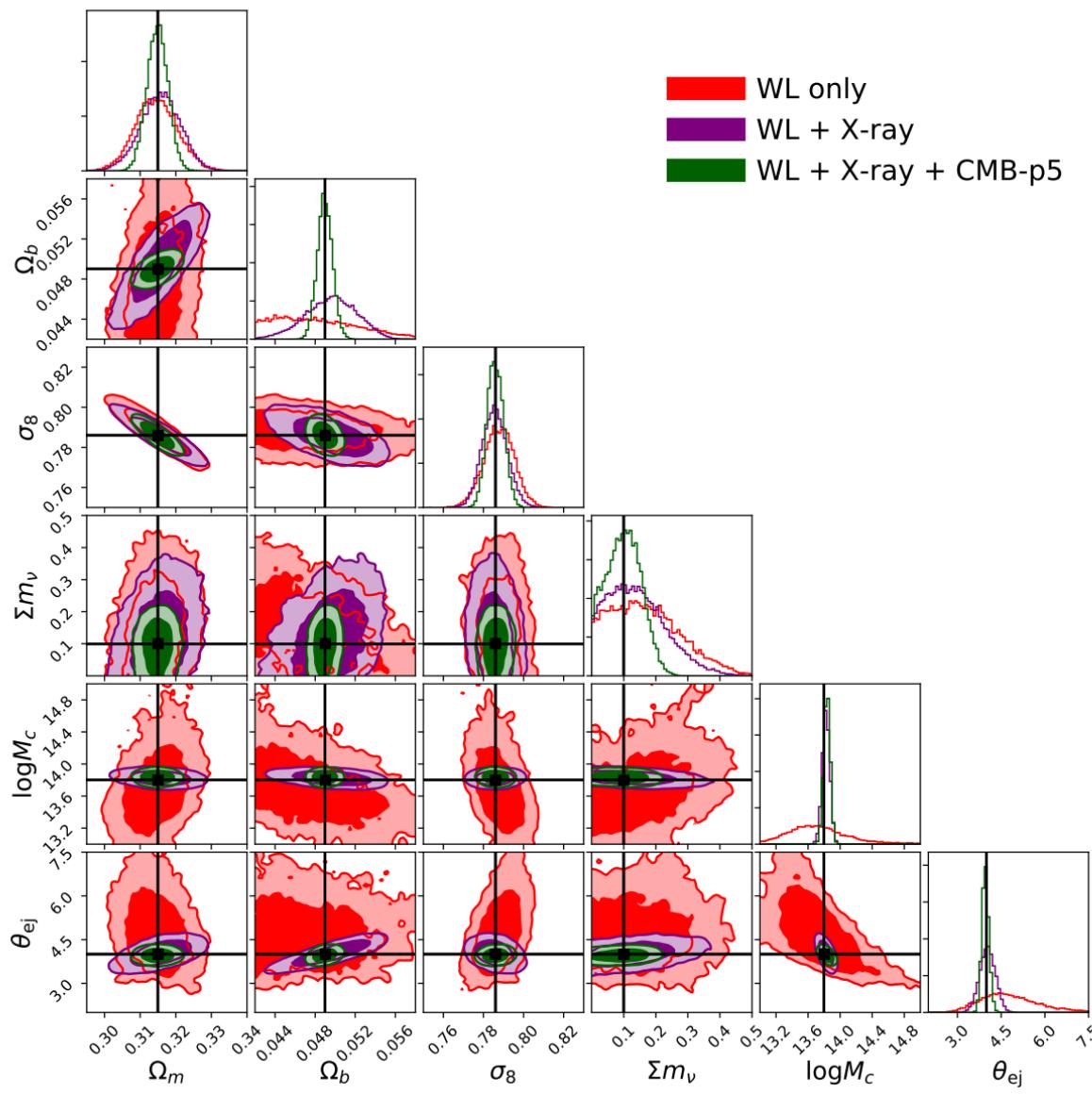


Schneider et al. (2020) Part I and Part II
JCAP, 04, 19 and 20

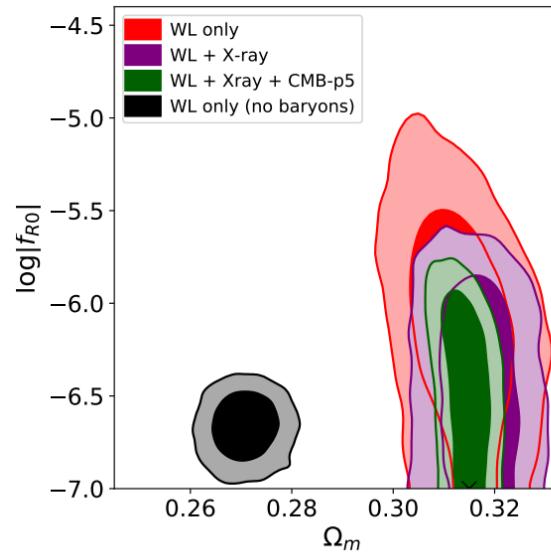
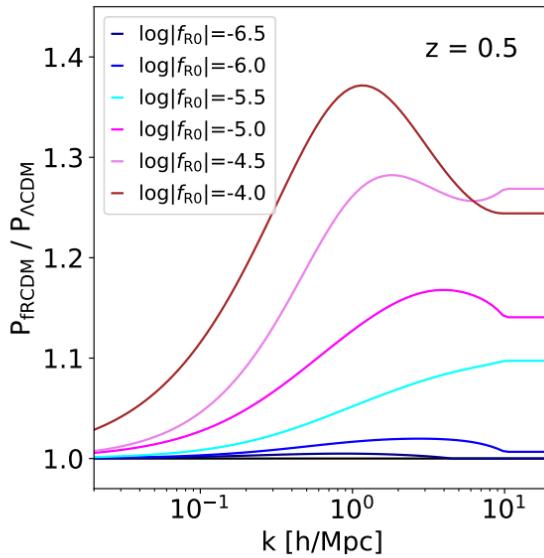
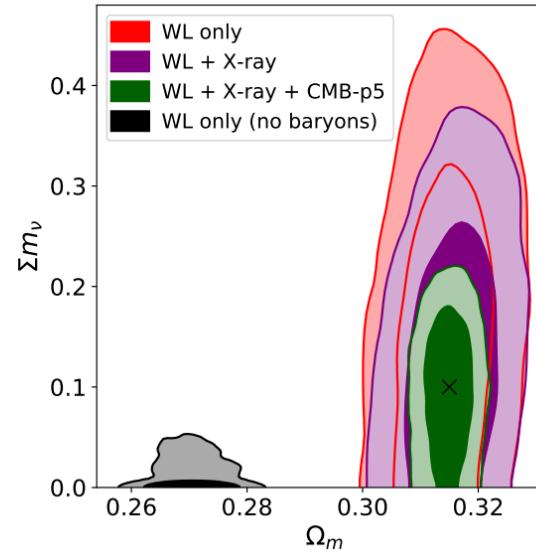
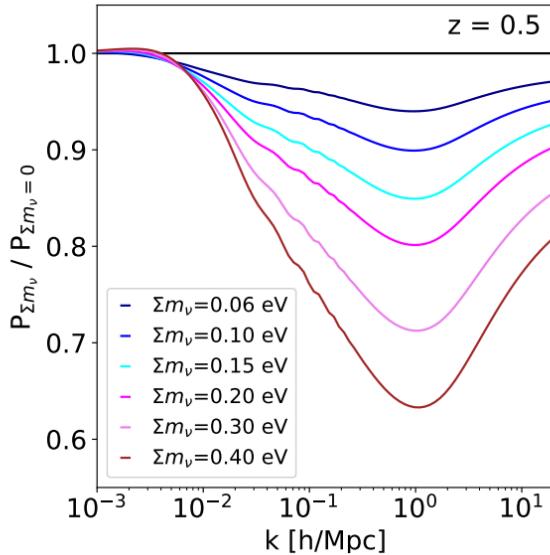
Weak-lensing as a probe of galaxy formation physics



Combining everything



Baryonic physics or new physics?



Conclusions

Systematic errors due to N body simulation below 1% at $k < 10 \text{ h/Mpc}$ ([Schneider et al. 2016](#), [Garrison et al. 2018](#)).

Given a reasonably sized parameter space, we can design an emulator for the non-linear boost of the matter power spectrum with « emulation-only-errors » as low as 0.2% ([The EuclidEmulator, Knabenhans et al. 2019](#)).

For weak lensing, baryonic effects kick in at $k = 1 \text{ h/Mpc}$ (pessimistic scenario). They can be encoded as a baryonic boost with additional parameters, and marginalised over ([Schneider et al. 2019, 2020](#)).

Forecast including baryonic effects show that combining CMB, X-ray and galaxy survey data would deliver maximum constraining power.

