Complex multi-loop results via finite-field techniques

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Based on:

T. P., JHEP 1612 (2016) 030, arXiv:1608.01902

T. P., JHEP 1907 (2019) 031, arXiv:1905.08019

Introduction & motivation

Experiments at LHC

- high-accuracy (% level)
- large SM background
- high c.o.m. energy ⇒ multi-particle states

We need scattering amplitudes

- high accuracy \Rightarrow loops (% level \sim 2 loops)
- multi-particle ⇒ high multiplicity

Theoretical studies of amplitudes

• structures of QFT/gauge theories



State of the art of scattering amplitudes

- Tree-level and one loop
 - today, mostly numeric
 - essentially solved
 - automated
- Two and higher loops
 - many calculations in recent years . . .
 - ... but still some open issues
 - ullet until recently, restricted to $2 \to 2$ processes
 - beyond MPLs not well understood

Two and higher loops

- Algebraic calculations for multi-loop amplitudes
 - preferred strategy $0 \ \ell \ge 2$ loops
 - faster/more stable evaluation
 - better suited for many multi-loop techniques
 - allows more tests, studies, etc. . . and better control
 - often characterized by high complexity

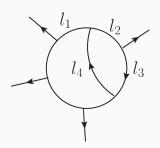
- Complexity can be a combination of
 - number of loops for high accuracy
 - number of legs for high multiplicity
 - numbers of scales (invariants, external/internal masses)

Loop amplitudes

ullet An integrand contribution to ℓ -loop amplitude

$$\mathcal{A} = \int_{-\infty}^{\infty} \left(\prod_{i=1}^{\ell} d^d k_i \right) \frac{\mathcal{N}}{D_1 D_2 D_3 \cdots}$$

- \bullet rational function in the components of loop momenta k_j
- ullet polynomial numerator ${\cal N}$
- quadratic denominators corresp. to loop propagators



$$D_j = l_j^2 - m_j^2$$

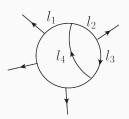
Computing amplitudes: Step 1/3

Write amplitudes as I.c. of Feynman integrals

$$\mathcal{A} = \sum_{j} a_{j} I_{j}$$

- Dependence on particle-content in rational coeff.s a_j
- The integrals should have a "nice" / "standard" form

$$I = \int_{-\infty}^{\infty} \left(\prod_{i=1}^{\ell} d^d k_i \right) \frac{1}{D_1^{\alpha_1} D_2^{\alpha_2} D_3^{\alpha_3} \cdots}, \qquad \alpha_j \leq 0$$



$$D_j = \begin{cases} l_j^2 - m_j^2 \\ l_j \cdot v_j - m_j^2 \end{cases}$$

Hard to do at high multiplicity

Computing amplitudes: Step 2/3

Chetyrkin, Tkachov (1981), Laporta (2000)

• Feynman integrals obey linear relations, e.g. IBPs

$$\int \left(\prod_j d^d k_j\right) \frac{\partial}{\partial k_j^{\mu}} v^{\mu} \frac{1}{D_1^{\alpha_1} D_2^{\alpha_2} \cdots} = 0, \qquad v^{\mu} = \begin{cases} p_i^{\mu} & \text{external} \\ k_i^{\mu} & \text{loop} \end{cases}$$

- Very large and sparse linear systems
- Reduce to linearly independent Master Integrals (MIs) $\{G_1, G_2, \ldots\} \subset \{I_i\}$

$$I_j = \sum_k c_{jk} G_k$$

Computing amplitudes: Step 3/3

- The MIs can often be computed analytically
 - in terms of special functions (MPLs, elliptic, ...)
 - most effective method is differential equations (DEs)
 Kotikov (1991), Gehrmann, Remiddi (2000)
 - can be simplified by the choice of MIs, e.g. UT integrals Henn (2013)
- Numerical methods may work depending on the process
 - the most successful is sector decomposition Binoth, Heinrich (2000)
 - can be improved via IBP reduction to a "better" basis of MIs

Computing amplitudes

Computing amplitudes (summary)

- 1. Integral representation $\mathcal{A} = \sum_j a_j I_j$
- 2. IBP reduction $I_j = \sum_k c_{jk} G_k$
- 3. Compute MIs G_k

A major bottleneck

- Large intermediate expressions
- Intermediate stages much more complicated than final result

Main idea of the talk

- Reconstruct analytic results from "numerical" evaluations
- Can be used for steps 1, 2 and help with step 3 (e.g. using DEs)

Finite fields and functional reconstruction

Functional reconstruction

- reconstruct analytic results from numerical evaluations
 - evaluation over finite fields \mathcal{Z}_p (i.e. modulo prime integers p)
 - ullet use machine-size integers, $p < 2^{64} \Rightarrow$ fast and exact
 - collect numerical evaluations and infer analytic result
- sidesteps large intermediate expressions & highly parallelizable
- applicable to any rational algorithm
- first applications
 - IBPs and univ. reconstruction von Manteuffel, Schabinger (2014)
 - helicity amplitudes and multivariate reconstruction T.P. (2016)

Some notable examples

- FINRED (private) [von Manteuffel]
 - several results for 4-loop form factors [von Manteuffel, Schabinger]
- FINITEFLOW [T.P.]
 - Several two-loop five-point amplitudes
 [Badger, Brønnum-Hansen, Hartanto, T.P.;
 Badger, Chicherin, Gehrmann, Heinrich, Henn, T.P., Wasser, Zhang, Zoia]
 - Matter dependence of the four-loop cusp anomalous dimension [Henn, T.P., Stahlhofen, Wasser]
- CARAVEL (private)
 - [Abreu, Dormans, Febres Cordero, Ita, Page, Sotnikov, Zeng]
 - analytic five-parton amplitudes
- FIRE 6 [A.V. Smirnov, F.S. Chuharev]
 - Four-loop quark form factor with quartic fundamental colour factor [Lee, Smirnov, Smirnov, Steinhauser]

The black-box interpolation problem

Given a rational function f in the variables $z = (z_1, \ldots, z_n)$ over Q

ullet Reconstruct analytic form of f, given a numerical procedure

$$(\boldsymbol{z},p) \longrightarrow \boxed{f} \longrightarrow f(\boldsymbol{z}) \bmod p.$$

- ullet evaluate f numerically for several $oldsymbol{z}$ and p
- efficient multivariate reconstruction algorithms exist
 e.g. T.P. (2016,2019), Klappert, Lange (2019)
- ullet upgrade analytic f over $\mathcal Q$ using rational reconstruction algorithm [Wang (1981)] and Chinese remainder theorem

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Question in this talk

How to build the black box?

Example: Scattering amplitudes over finite fields

T.P. (2016)

- External states (momenta and polarizations)
 - rational parametrization with momentum twistors variables Hodges (2009), Badger, Frellesvig, Zhang (2013), Badger (2016)
- Tree-level
 - diagrams or recursion relations (e.g. Berends-Giele)
- Loop integrands
 - Feynman diagrams and t'Hooft algebra
 - Unitarity cuts sewing tree-level currents
 - higher finite-dim. representation of internal states in dim. reg.
- Integrand reduction
 - linear fit to a "nice" integrand basis

How to build the black box?

How to build a code for fast numerical evaluations of finite fields? We can consider a few options:

- 1. Low-level coding (e.g. in C/C++/FORTRAN)?
 - ✓ very good runtime efficiency
 - X harder to program
 - X limits usability
- 2. Low-level coding + high-level interfaces?
 - common algorithms in C++ (e.g. linear solvers, fits, etc...)
 - high-level wrapper (e.g. for MATHEMATICA/PYTHON)
 - ✓ good efficiency and usability
 - X not flexible
 - X these algorithms are often intermediate steps

How to build the black box?

Observations:

- A typical multi-loop algorithm involves several steps
 - solving linear systems
 - substitutions / changes of variables
 - etc. . .
- Large simplifications often occur at the very last stages
 - it's best to do everything numerically
 - only the final expression reconstructed analytically
- Many algorithms share common "building blocks"

FiniteFlow: using data flow graphs

FINITEFLOW [T.P. (2019)] has three main components

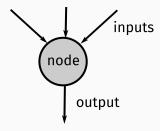
- 1. "basic" algorithms in C++ over finite fields
 - dense/sparse linear solvers, linear fits, evaluating rat. functions, list manipulations, etc...
- 2. higher-level framework to combine them into complex ones
 - output of a basic algorithm is input of others
 - graphical representation of your calculation (dataflow graphs)
- 3. multivariate reconstruction algorithms

FiniteFlow

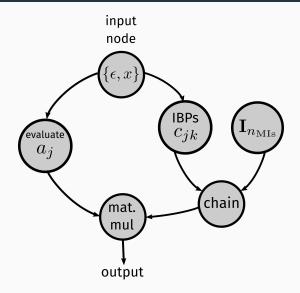
- build complex algorithms without any low-level programming (e.g. from MATHEMATICA interface)
- many methods for amplitudes can be cast in this framework

FiniteFlow: using data flow graphs

- FINITEFLOW uses (simplified) data flow graphs
 - Nodes represent numerical algorithms
 - Arrows represent lists of numerical values
- In my implementation, a node has
 - 0 or more lists (arrows) of input values
 - 1 list (arrow) of output values



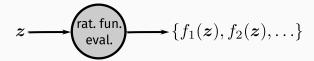
Example of a graph



Example: Evaluation of rational functions

- input: a list of values $z = (z_1, \dots, z_n)$
- ullet output: a list of rational functions $\{f_1,f_2,\ldots\}$ at z

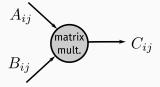
$$f_i(\boldsymbol{z}) = rac{p_i(\boldsymbol{z})}{q_i(\boldsymbol{z})} = rac{\sum_{lpha} n_{i,lpha} \, \boldsymbol{z}^{lpha}}{\sum_{eta} d_{i,eta} \, \boldsymbol{z}^{eta}},$$



Example: Matrix multiplication

- Two lists as input
 - 1. entries of a matrix A
 - 2. entries of a matrix B
- use row-major order to store them as a list
- ullet ouput: entries of matrix C such that

$$C_{ij} = \sum_{k} A_{ik} B_{kj}$$



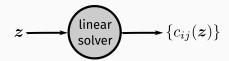
Example: Linear solver

• A $n \times m$ linear system with parametric rational entries

$$\sum_{i=1}^{m} A_{ij} x_j = b_i, \quad (i = 1, \dots, n), \qquad A_{ij} = A_{ij}(z), \quad b_i = b_i(z)$$

- ullet input: list of values for paramers $oldsymbol{z}=(z_1,\ldots,z_n)$
- ullet output: solution $c_{ij}=c_{ij}(oldsymbol{z})$ such that

$$x_i = \sum_{j \in \mathsf{indep}} c_{ij} \, x_j + c_{i0} \qquad (i \not\in \mathsf{indep})$$



Learning algorithms

- Some algorithms have a learning phase
 - used to learn information for defining its output
 - must be completed before using them
- Example: linear solver
 - learn: its rank, dep. and indep. unknowns, indep. eq.s
 - learning phase: solve the system numerically a few times
 - optional: mark & sweep equations (sparse solver)
- ⇒ It can be used to simplify the algorithm see also e.g. Kira: Maierhöfer, Usovitsch, Uwer (2017)

IBP reduction

- IBPs are large and sparse linear systems
- ullet they reduce Feynman integrals I_j to a lin. indep. set of MIs G_j

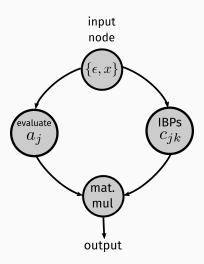
$$I_i = \sum_j c_{ij} G_j$$

amplitudes and other multi-loop objects can be reduced mod IBPs

$$A = \sum_{j} a_{j} I_{j} = \sum_{jk} a_{j} c_{jk} G_{k} = \sum_{j} A_{j} G_{j}$$
, with $A_{j} = \sum_{k} a_{k} c_{kj}$

- ullet final results for A_k often much simpler than c_{ij}
- \Rightarrow solve IBPs numerically and compute A_j via a matrix multiplication

IBP reduction



Differential equations for MIs

ullet The MIs G_k satisfy differential equations Kotikov (1991), Gehrmann, Remiddi (2000)

$$\partial_x G_i = \sum_j A_{ij}^{(x)} G_j$$

- Identify MIs G_i (e.g. by solving IBPs numerically)
- Compute their derivatives in terms of (non-master) loop integrals

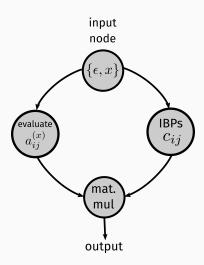
$$\partial_x G_i = \sum_j a_{ij}^{(x)} I_j$$

- Reduce the needed integrals modulo IBPs: $I_i = \sum_j c_{ij} G_j$
- The final result is given by a matrix multiplication

$$A_{ij}^{(x)} = \sum_{k} a_{ik}^{(x)} c_{kj}$$

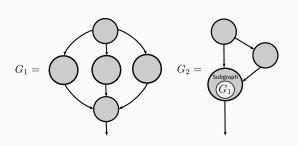
• Reconstruct $A_{ij}^{(x)}$ analytically from its numerical evaluations

Differential equations for MIs



Subgraphs

- Any graph G_1 can be used as a subgraph by an algorithm (a node) A belonging to another graph G_2
 - A will evaluate G_1 several times to compute its output
 - input of G_1 = auxiliary variables chained with inputs of A



Examples:

- Laurent expansion
- maps: evaluate G₁
 for several inputs
- partial reconstructions
- (total or partial) fits w.r.t. an ansatz

Coefficients of the ϵ -expansion

 $\bullet\,$ If MIs are known analytically in terms of special functions f_k

$$G_j = \sum_k g_{jk}(\epsilon, x) f_k + \mathcal{O}(\epsilon),$$

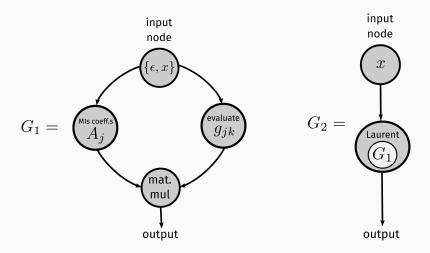
we can compute

$$A = \sum_{k} u_k(\epsilon, x) f_k + O(\epsilon), \text{ where } u_k(\epsilon, x) = \sum_{j} A_j(\epsilon, x) g_{jk}(\epsilon, x)$$

• what we want is the ϵ -expansion of the $u_k(\epsilon,x)$

$$u_k(\epsilon, x) = \sum_{j=-p}^{0} u_k^{(j)}(x) \, \epsilon^j + \mathcal{O}(\epsilon),$$

Coefficients of the ϵ -expansion



Reconstruction of amplitudes

Observations:

- we can detect linear relations btw. the coefficients of the amplitude and reconstruct a simpler subset of linearly independent ones
- we can subtract IR divergencies predicted from lower orders by

$$\mathcal{A} = \mathcal{Z} \, \mathcal{A}^f$$

- this can also be used to subtract a finite contribution
- significantly simplifies result and reconstruction

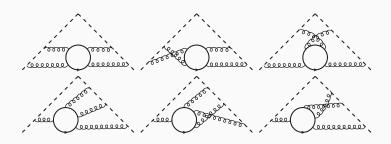
Open question:

Can we improve this subtraction of finite pieces?

Cutting-edge applications of FiniteFlow

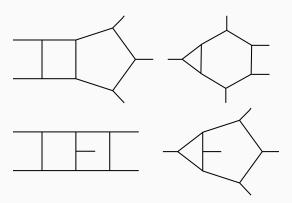
 Matter dependence of the 4-loop cusp anomalous dimension Henn, T.P., Stahlhofen, Wasser (2019)

(see also: Lee, Smirnov, Smirnov, Steinhauser (2019))



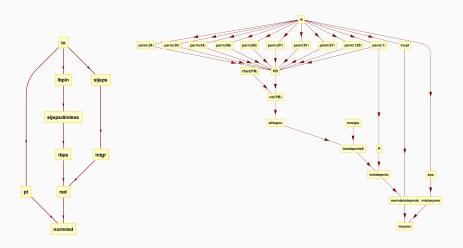
Cutting-edge applications of FiniteFlow

- Five-point two-loop amplitudes
 - \bullet Several planar results for five partons and W+4 partons [Badger, Brønnum-Hansen, Hartanto, T.P. (2017-2019)]
 - all-plus five gluon non-planar [Badger, Chicherin, Gehrmann, Heinrich, Henn, T.P., Wasser, Zhang, Zoia (2019)]



Example of graphs in FiniteFlow

Piecing together the all-plus five gluon amplitude (only planar contributions are shown)



Other notable applications of FiniteFlow

 Simplifying analytic expressions for NNLO QCD corrections to three-photon production at the LHC

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[Chawdhry, Czakon, Mitov, Poncelet (2019)]
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Analytic simplification of IBP systems
 [Xin Guan, Xiao Liu, Yan-Qing Ma (2019)]

 Deriving canonical differential equations for Feynman integrals from a single uniform weight integral (INITIAL public code)

[Christoph Dlapa, Johannes Henn, Kai Yan (2020)]

Public codes

• FINITEFLOW

https://github.com/peraro/finiteflow

- C++ code
- MATHEMATICA interface (strongly recommended)
- FINITEFLOW MATHTOOLS

https://github.com/peraro/finiteflow-mathtools

- packages FFUTILS, LITEMOMENTUM, LITEIBP, SYMBOLS
- examples (amplitudes, IBPs, diff. equations and many more)

Summary & Outlook

Summary

- Finite fields and functional reconstruction
 - enhance the possibilities of our theoretical predictions
 - new results unattainable with traditional computer algebra
 - public code FINITEFLOW
- Progress on 2-loop 5-point and other complex processes

Outlook

- More applications
 - massive processes, phase-space integrals, . . .
- High level of automation for higher-loop predictions