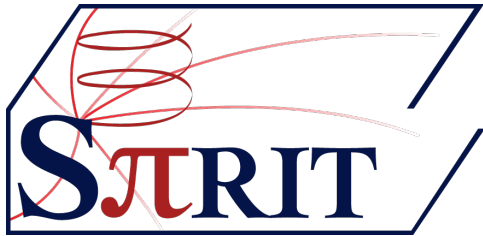


Isoscaling effect with $Z=1$ and 2 particles in Sn+Sn at 270 MeV/u

Jung Woo Lee
for S π RIT-collaboration
Korea University



Introduction to Isoscaling

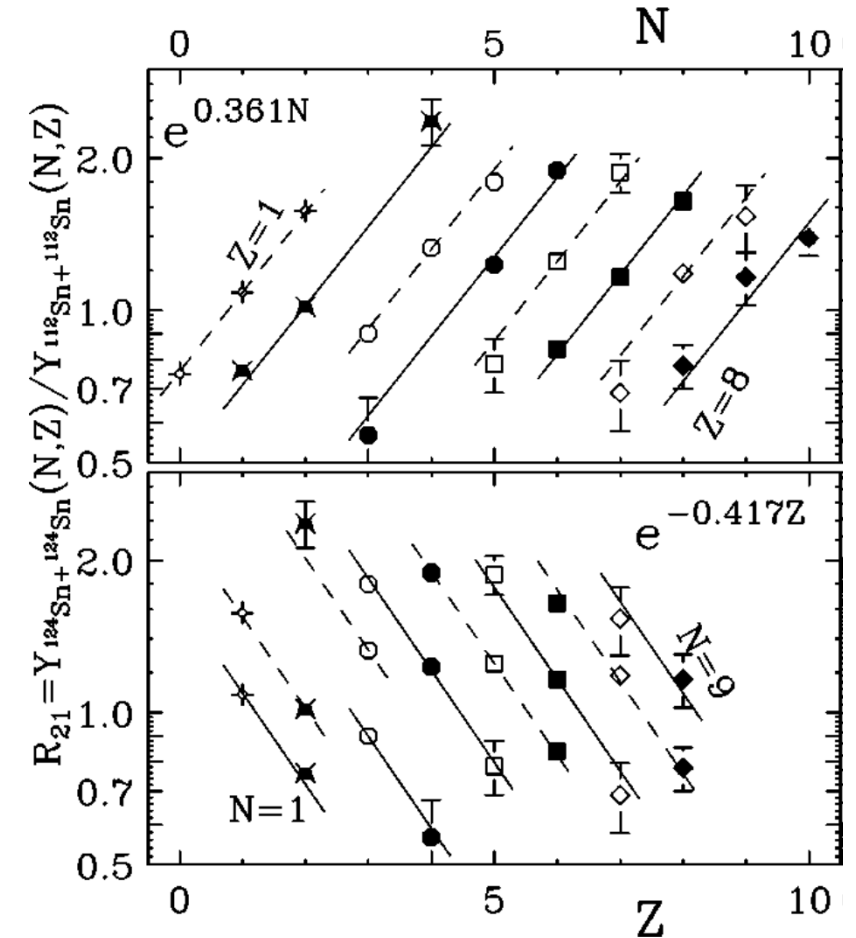
- Grand canonical ensemble, thermal equilibrium assumption in heavy-ion collision the fragment yield can be written down as:

- $$Y^{prim.}(N, Z, T) = \frac{VA^{3/2}}{\lambda_T^3} \omega(N, Z, T) \exp \left[\frac{E_B(N, Z) + N\mu_n + Z\mu_p}{T} \right].$$
- $$Y^{final}(N, Z, T) \simeq f(N, Z, T) Y^{prim.}(N, Z, T).$$

* S. Albergo et al. *Nuovo Cimento Soc. Ital. Fis., A* 89 (1985) 1

* H. S. Xu et al. *Phys. Rev. Lett.* 85 (2000) 716

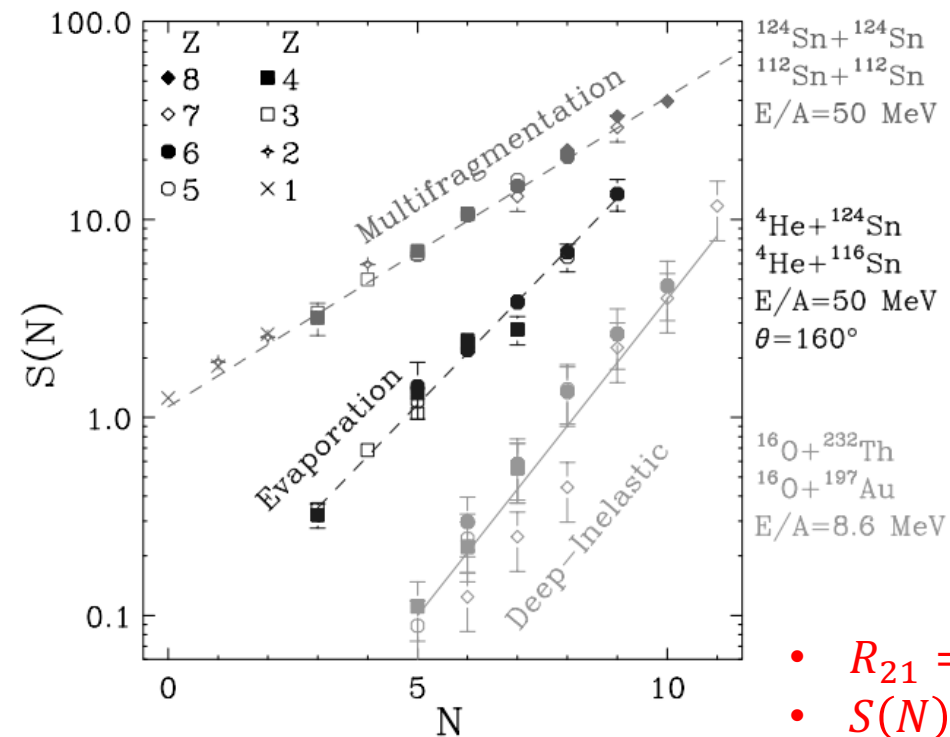
- We compare two collision system 2 and 1.
 - Two system have similar size and temperature.
 - System 2 is neutron rich system compared to system 1.
- The fragment yield ratio between two systems exhibit exponential behavior:
 - $$R_{21} = Y_{system-2}(N, Z) / Y_{system-1}(N, Z).$$
 - $$R_{21} = C \exp(\alpha N + \beta Z), \quad \alpha = \frac{\Delta\mu_n}{T}, \quad \beta = \frac{\Delta\mu_p}{T}.$$



* M. B. Tsang et al. (MSU) *Phys. Rev. C* 64 (2001) 041603

Introduction to Isoscaling

- Isoscaling was reported in many other experiments not only from the multi-fragmentation in heavy-ion collisions but in non symmetric collisions such as evaporation, deeply inelastic scattering and light ion induced reactions.



* M. B. Tsang et al. Phys. Rev. C 64 (2001) 041603

* M. B. Tsang et al. Phys. Rev. Lett. 86 (2001) 5023

* A. S. Botvina et al. Phys. Rev. C 65 (2002) 044610

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* Akira Ono et al. Phys. Rev. C 68 (2003) 051601

* S. R. Souza et al. Phys. Rev. C 80 (2009) 044606

- $R_{21} = C \exp(\alpha N + \beta Z)$
- $S(N) = R_{21} \exp(-\beta Z)$

Introduction to Isoscaling

- Isoscaling was reported in many other experiments not only from the multi-fragmentation in heavy-ion collisions but in non symmetric collisions such as evaporation, deeply inelastic scattering and light ion induced reactions.
- Isoscaling was also reported in many model calculations
 - SMM (Statistical multi-fragmentation model),
 - EES (expanding emission source model),
 - AMD (antisymmetrized molecular dynamics).
- Scaling parameter α and symmetry energy C_{sym} :
 - $$\alpha = \frac{4C_{sym}}{T} \left[\left(\frac{Z_1}{A_1} \right)^2 - \left(\frac{Z_2}{A_2} \right)^2 \right], (T < 5 \text{ MeV}).$$

* M. B. Tsang et al. *Phys. Rev. C* 64 (2001) 041603

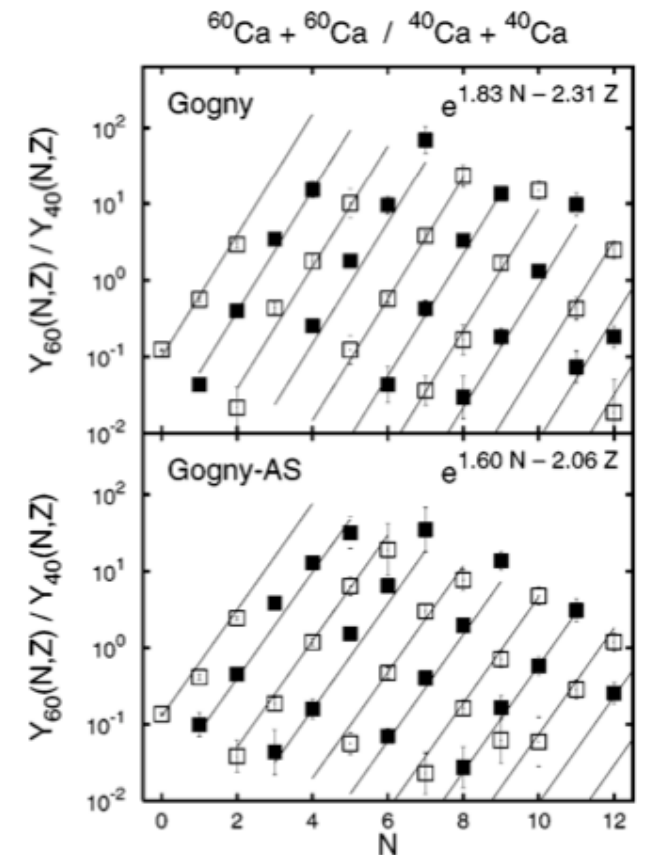
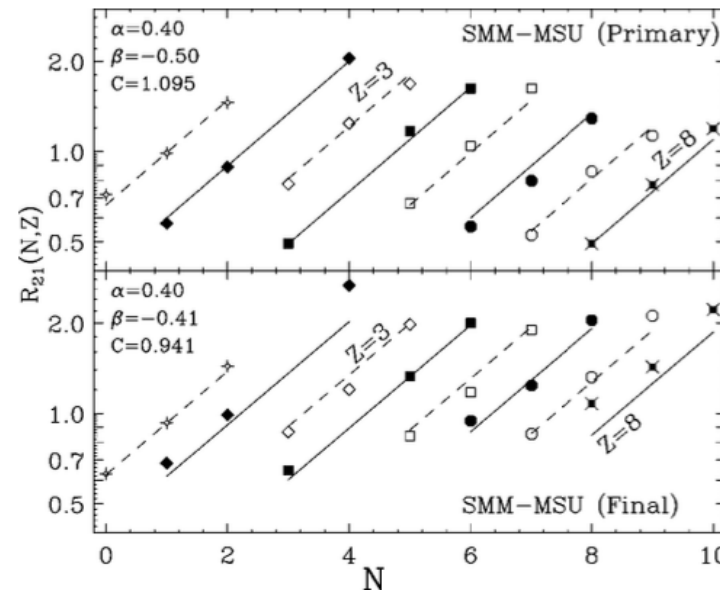
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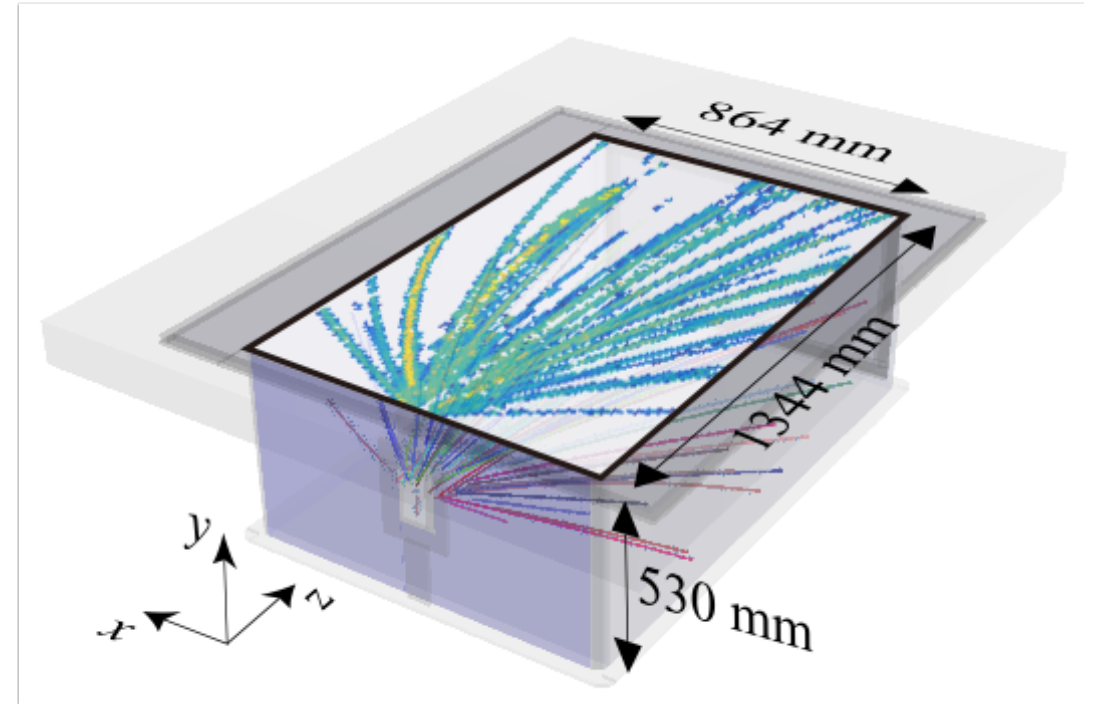
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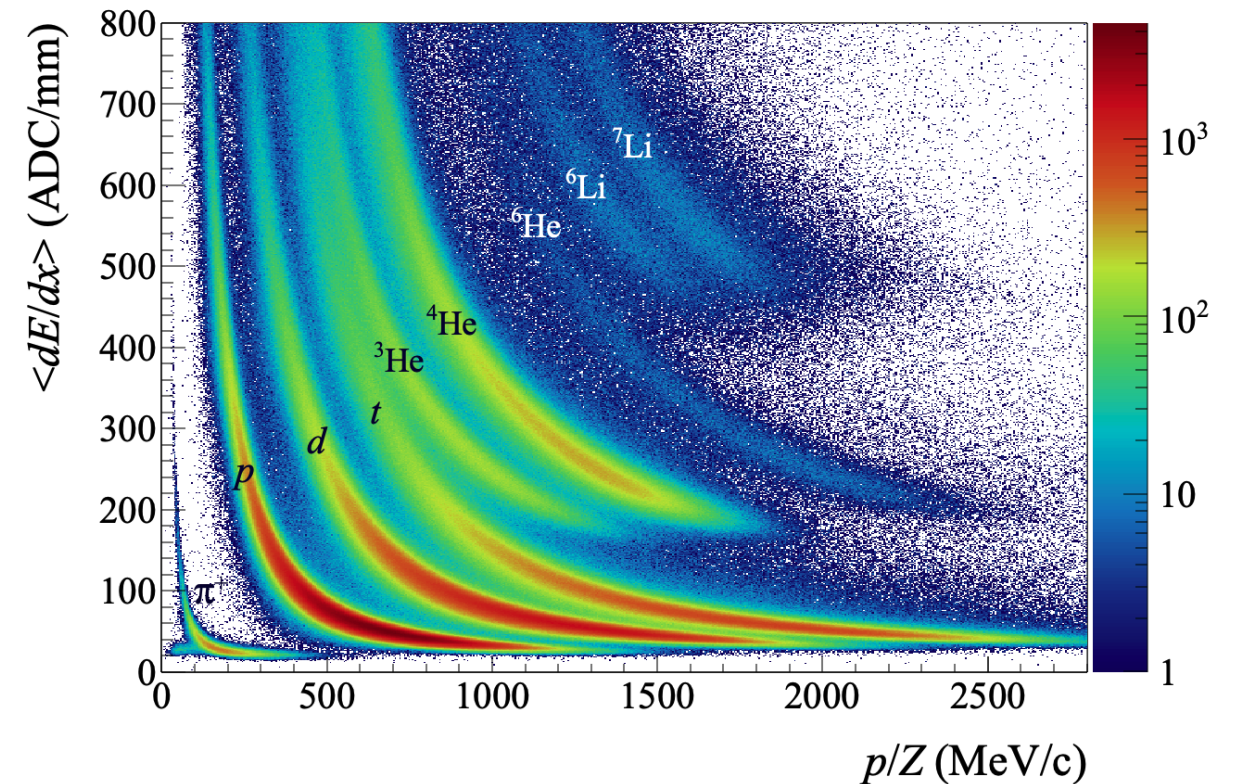
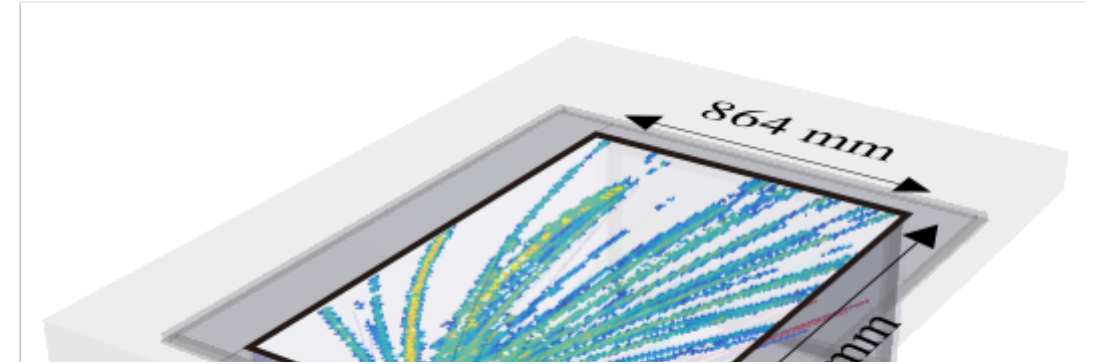
S π RIT Experiment

- Sn + Sn collisions at 270 MeV/u measured with S π RIT-TPC inside the SAMURAI magnet producing 0.5 T at RIBF RIKEN (2015-2016).
 - $^{132}\text{Sn} + ^{124}\text{Sn}$ (N/Z = 1.56)
 - $^{108}\text{Sn} + ^{112}\text{Sn}$ (N/Z = 1.20)
 - $^{124}\text{Sn} + ^{112}\text{Sn}$ (N/Z = 1.36)
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S π RIT Experiment

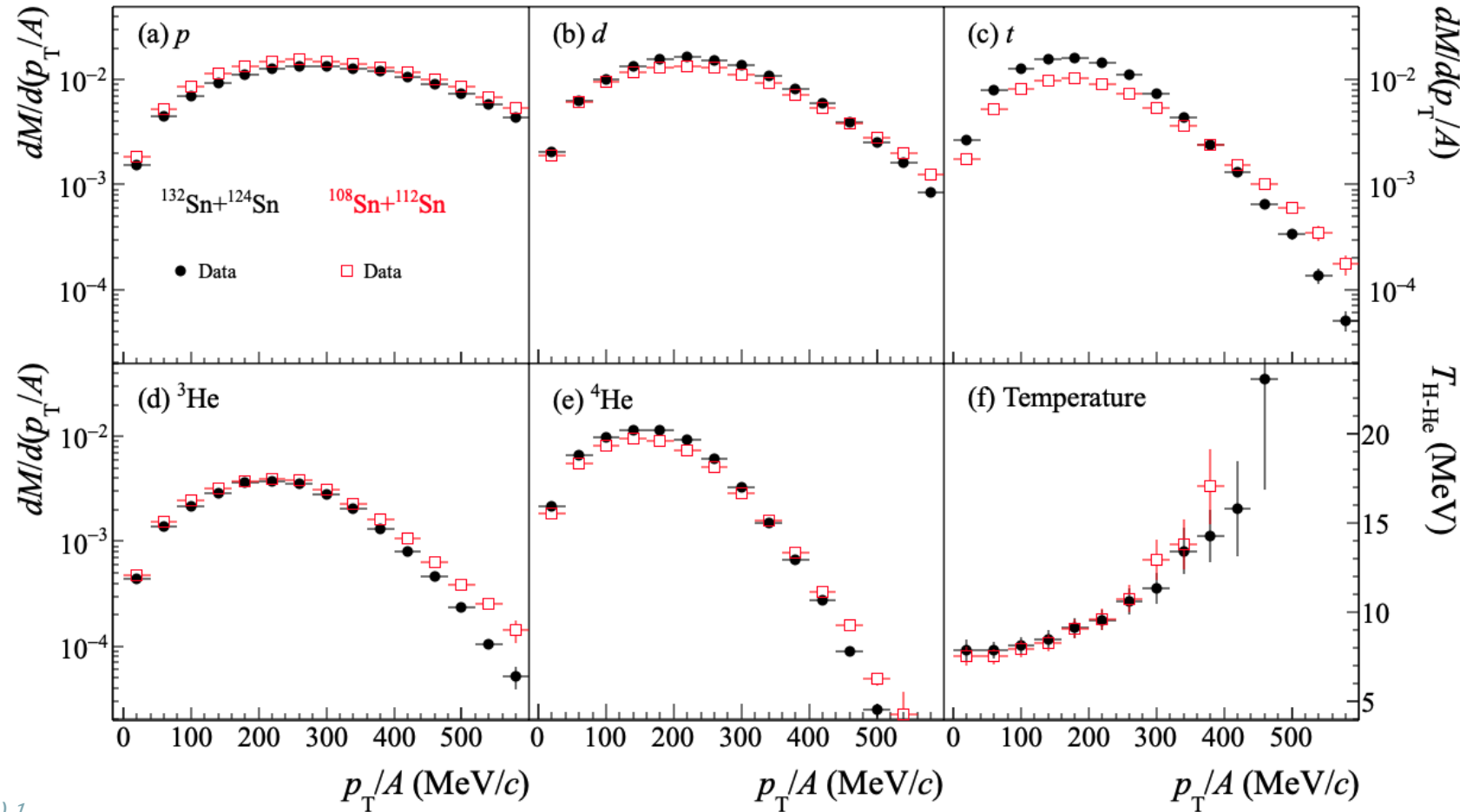
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- Pions, Hydrogen and Helium isotopes are found.
- Cut
 - Rapidity $y_0 = 0 - 0.4$
 - Impact factor $b < 1.5$ fm



p_T/A Spectra

- Different trend of p_T/A is shown for (N-Z) groups.
- N-Z=1 (p, ^3He) particles show more yield in neutron deficient system.
- N-Z=0 (d, ^4He) particles show little more yield is shown for neutron rich system.
- N-Z=1 (t) particle show larger yield is shown for neutron rich system.
- H-He double ratio temperature was calculated for reference.

$$T = 1.43 \ln^{-1} \left[1.6 \frac{Y(d) Y(^4\text{He})}{Y(t) Y(^3\text{He})} \right]$$



* S. Albergo et al. Nuovo Cimento Soc. Ital. Fis., A 89 (1985) 1

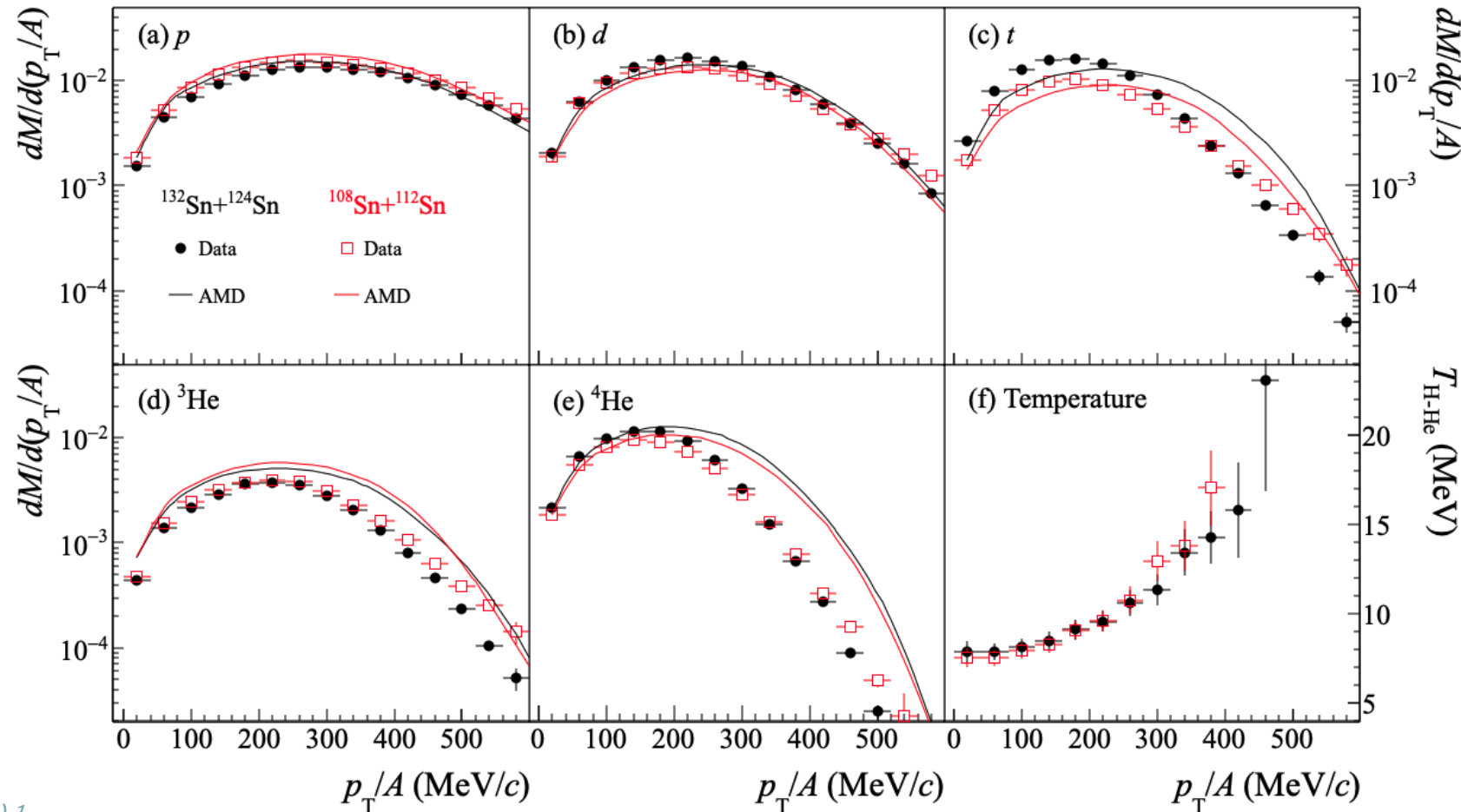
p_T/A Spectra + AMD

AMD : Soft symmetry energy

* M. Kaneko et al. Phys. Lett. B 822 (2021) 136681

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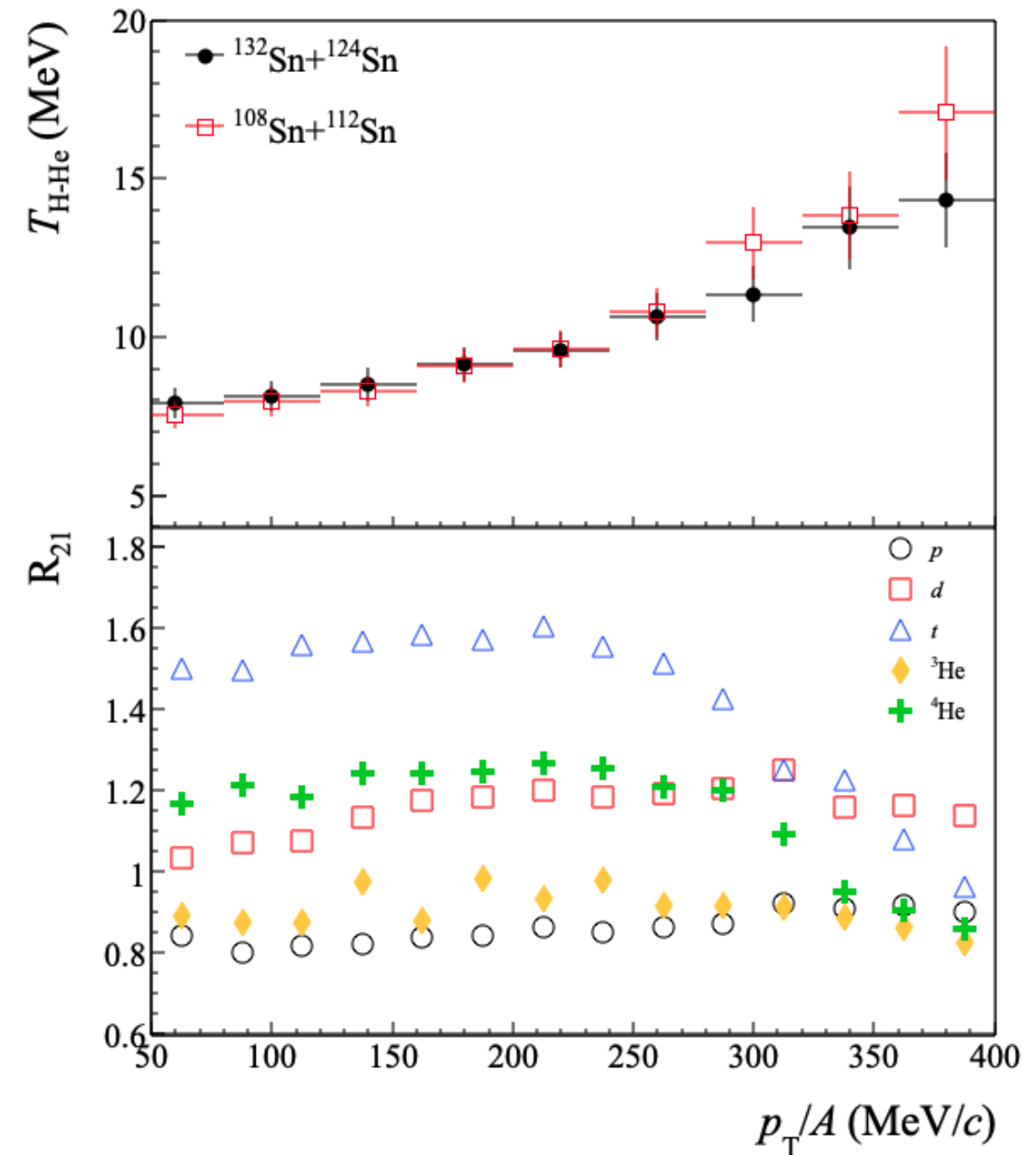
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Spectral Yield Ratio

- N-Z group is clearly shown in the spectra.
 - $R_{21} = C \exp(\alpha N + \beta Z)$.
 - $\alpha \simeq -\beta$
- Valid at p_T/A below 250 MeV/c and temperature exceed over 10 MeV at this point.

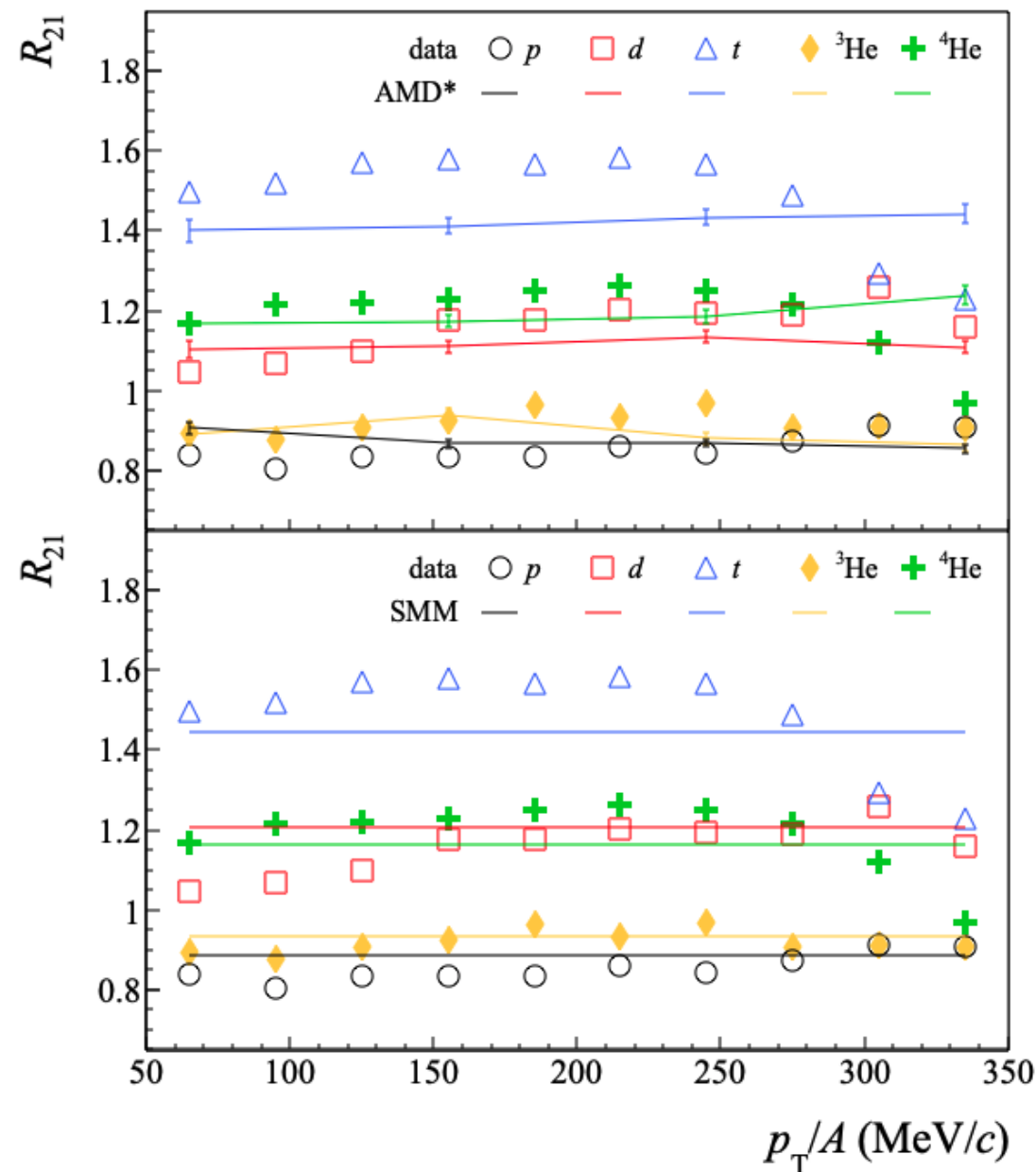


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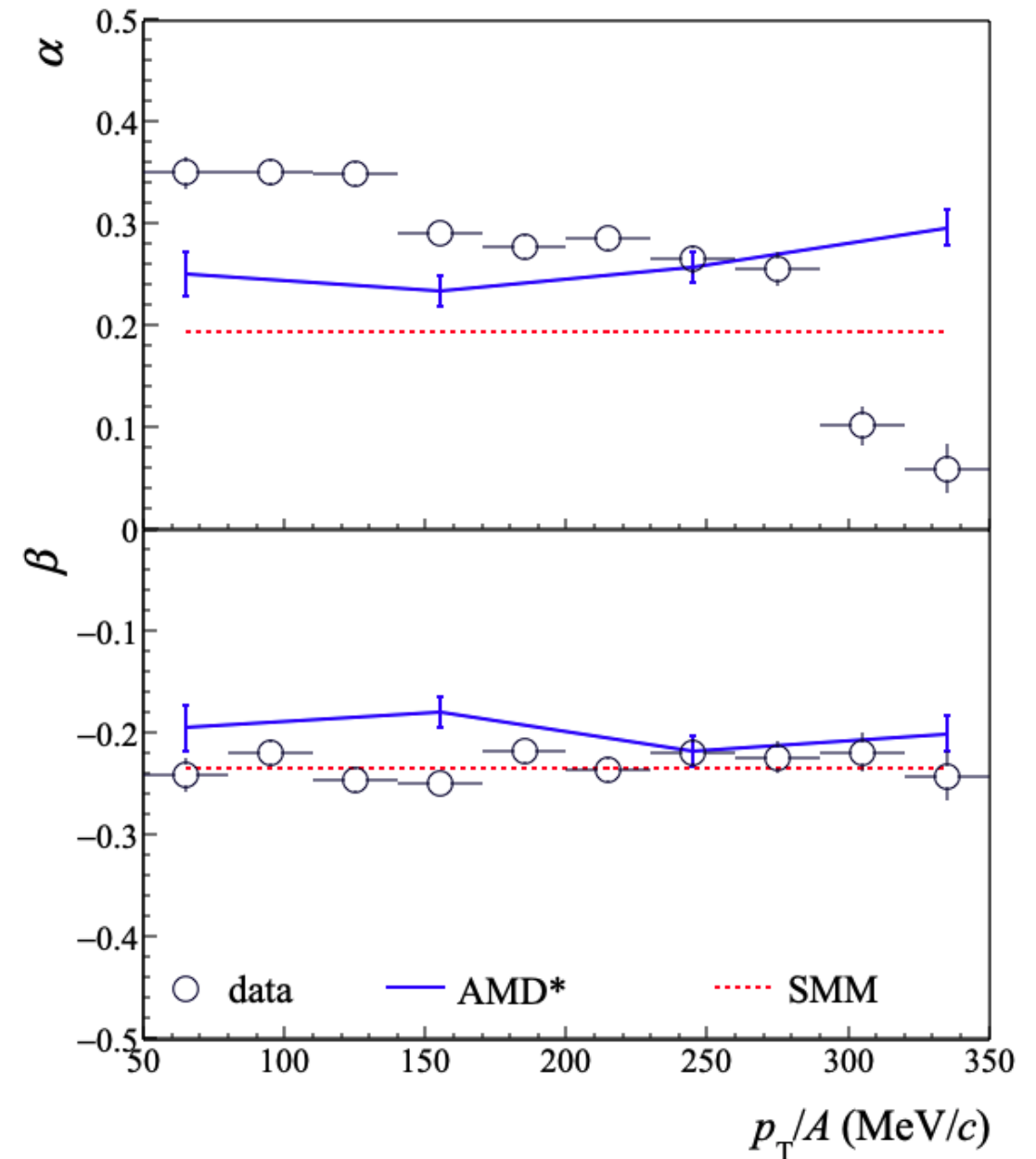
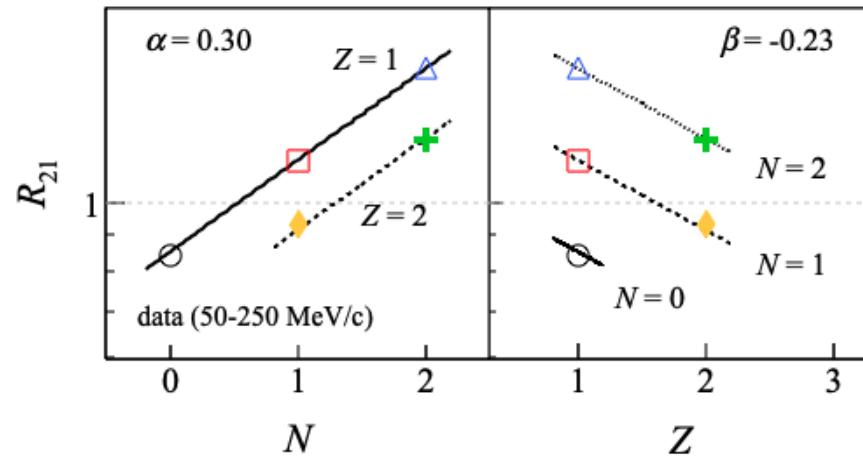
**S. R. Souza et al. Nucl. Phys. A 989 (2019) 69*

- SMM calculation where provided with A and Z of the source:
 - Neutron-rich system : A=172, Z=79
 - Neutron-deficient system : A=142, Z=72
- Compared AMD and SMM calculation explain the data up to 250 MeV/c.
- The overall trend breaks above 250 MeV/c.



Isoscaling with data

- Isoscaling fit to data, AMD and SMM where performed to get the α and β values through p_T/A .
- Nice isoscaling is given below 250 MeV/c.
- Isoscaling above 250 MeV/c is shown to break down.
- α from data decrease even below 250 MeV/c while AMD stays constant. β values stay constant for all cases.



Conclusion

- Isoscaling properties are studied in $^{132}\text{Sn}+^{124}\text{Sn}$ and $^{108}\text{Sn}+^{112}\text{Sn}$ systems at 270 MeV/u. compared to the previous studies on isoscaling, we explored collision systems with comparably large energy (270 MeV/u) and large isospin difference ($N/Z_2 = 1.56$ and $N/Z_1 = 1.20$).
- Since temperature increase as a function of p_T/A , we performed isoscaling fit for different p_T/A bins.
- The spectral yield ratio R21 show constant trend up to $p_T/A = 250 \text{ MeV}/c$ but triton and ^4He starts to decrease above this region. Consequently, isoscaling breaks above this boundary ($p_T/A = 250 \text{ MeV}/c$).
- In the reliable p_T/A region (50– 250 MeV/c), we obtained $\alpha = 0.30$ and $\beta = -0.23$.
- Thermal equilibrium approximation might not be appropriate for high energy particles in heavy-ion collisions at energy range around 270 MeV/u.
- The temperature of the systems for AMD (H-He temperature) show constant trend through p_T/A and temperature for SMM is fixed to 8 MeV. This means that calculations is assuming thermal equilibrium through out the system and this might be the reason for different scaling trend compared to data.
- For the feature studies, more cases are needed be explored in other similar systems, and include other produced particles (neutrons, ^6He , Li isotopes) if possible.