

Direct Photon Production In Association With A Heavy Quark - from the Tevatron to the LHC

Tzvetalina Stavreva

*Florida State University
Dept. of Physics*

September 24, 2009

Table of contents

Introduction

Theory Overview

Tevatron

LHC

Conclusions

Direct Photons and Heavy Quarks

What are Direct Photons and Why are they important?

- Any photon that is produced during the hard scattering process or via fragmentation
- Not to be confused with γ coming from the decay of hadrons, such as $\pi^0 \rightarrow \gamma\gamma$, $\eta \rightarrow \gamma\gamma$
- Carriers of electromagnetic force
- Escape confinement
- Photon acts as a probe of the hard scattering
- Charge coupling allows for a distinction between charm and bottom

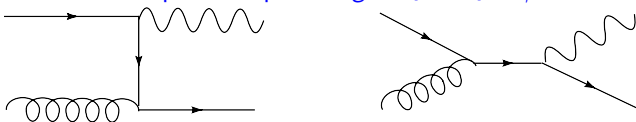
$\gamma + Q$ production

- Direct photons are produced in association with many different particles
- Look at one part of the cross section \rightarrow piece with heavy quarks
- Better understand the role of heavy quarks in high p_T collisions
- Possibility to better constrain Parton Distribution Functions of heavy quarks

Hardscattering Production

- Leading Order - $\mathcal{O}(\alpha\alpha_s)$ - Only one hard-scattering subprocess

Compton Subprocess $g + Q \rightarrow Q + \gamma$



- Next-to-Leading Order - $\mathcal{O}(\alpha\alpha_s^2)$

2 \rightarrow 3 hard-scattering subprocesses

$$g + g \rightarrow Q + \bar{Q} + \gamma$$

$$g + Q \rightarrow g + Q + \gamma$$

$$Q + q \rightarrow q + Q + \gamma$$

$$Q + \bar{q} \rightarrow Q + \bar{q} + \gamma$$

$$Q + Q \rightarrow Q + Q + \gamma$$

$$Q + \bar{Q} \rightarrow Q + \bar{Q} + \gamma$$

$$q + \bar{q} \rightarrow Q + \bar{Q} + \gamma$$

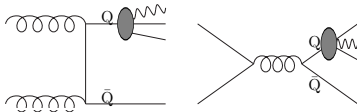
- Also need to include Direct Photons which are produced via fragmentation

Photon Fragmentation

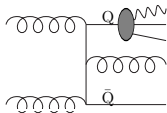
- If photon is emitted collinearly to a quark \rightarrow singularity
- Absorb singularity in $D_{\gamma/q,g}(z, \mu_F)$; resum large logs in $D_{\gamma/q,g}(z, \mu_F)$ FF via DGLAP
- Photon couples to quark, responsible for behavior of $D_{\gamma/q,g}(z, \mu_F) \sim \mathcal{O}(\alpha/\alpha_s)$

Fragmentation Effects

- LO: include all $2 \rightarrow 2$ subprocesses $\sim \mathcal{O}(\alpha_s^2)$,
 $\mathcal{O}(\alpha_s^2) \otimes D_{\gamma/q,g} \sim \alpha_s^2 \alpha / \alpha_s = \alpha \alpha_s$



- NLO: same idea as in LO case, convolute all $2 \rightarrow 3 \sim \mathcal{O}(\alpha_s^3)$ with γ FF

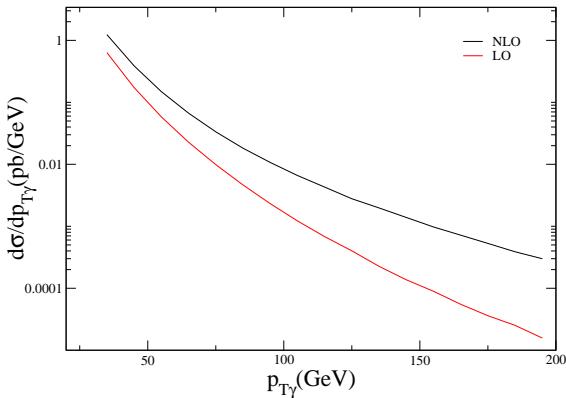


Tevatron Predictions

- $D\bar{D}$ cuts: $p_{T\gamma} > 30 \text{ GeV}$, $p_{Tb} > 15 \text{ GeV}$, $|y_\gamma| < 1$, $|y_Q| < 0.8$

$$p + \bar{p} \rightarrow \gamma + b + X$$

$$\sqrt{S} = 1.96 \text{ TeV}$$

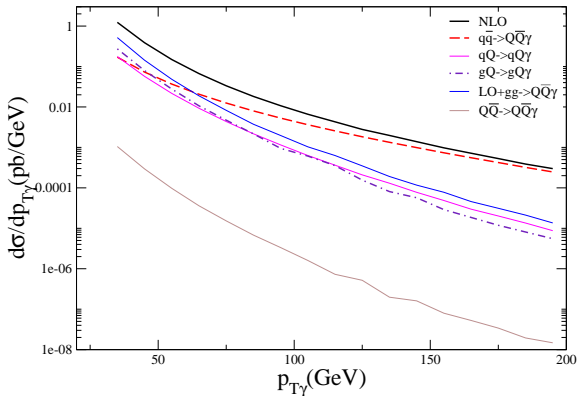


- As $p_{T\gamma}$ increases the difference between LO and NLO grows
- What drives this difference?

Subprocess Contributions

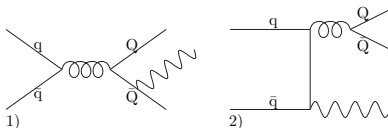
$$p + \bar{p} \rightarrow \gamma + b + X$$

$$\sqrt{S} = 1.96 \text{ TeV}$$

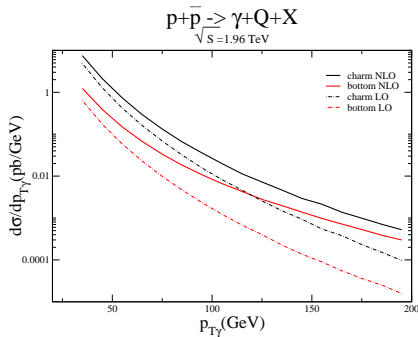


- At $p_{T\gamma} \sim 70$ GeV $q\bar{q} \rightarrow Q\bar{Q}\gamma$ starts to dominate
- PDFs of light valance quarks are prevalent
- Take a closer look at the annihilation subprocess

Annihilation Subprocess

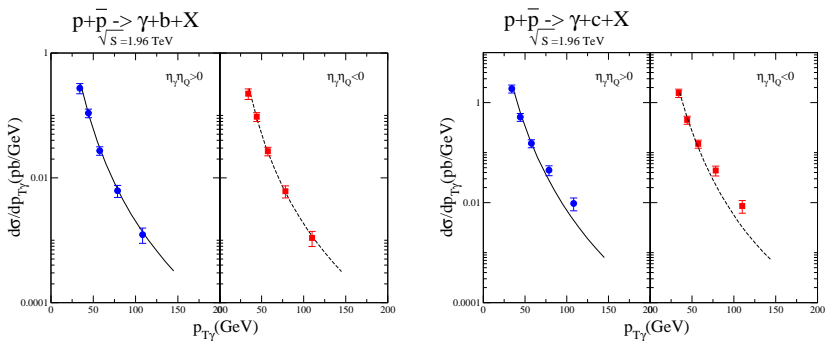


- Diagram 1) $\sim e_Q^2$ - photon couples to heavy quark, Diagram 2) $\sim e_q^2$ - photon couples to initial quarks
- Diagram 2) - dominant \rightarrow difference between c and b decreases with p_T



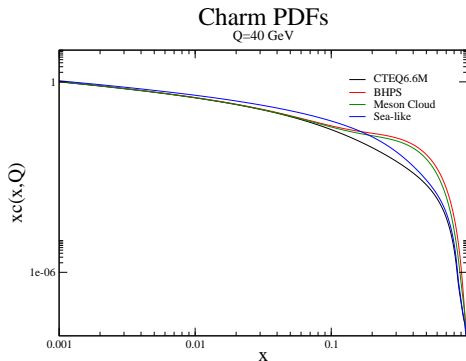
Comparison between theory and data

Measurements by DØ Collaboration



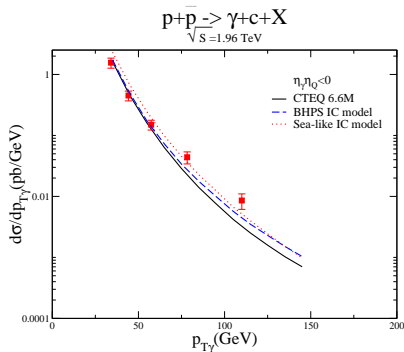
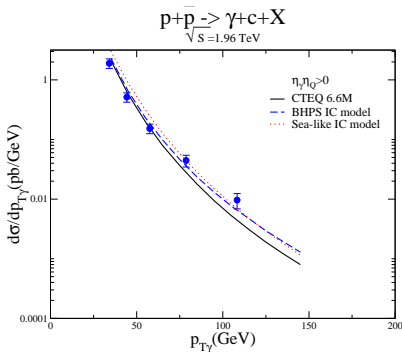
- There is really good agreement between data and theory for the bottom cross section
- For charm the data points at large $p_{T\gamma}$ lie above the theory curve \rightarrow possible explanation - existence of intrinsic charm

Intrinsic Charm



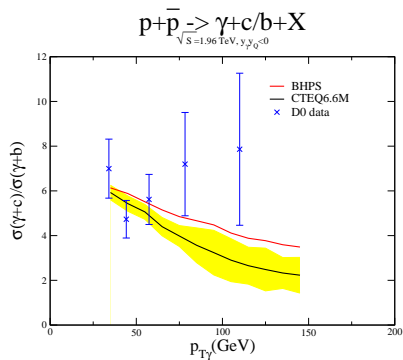
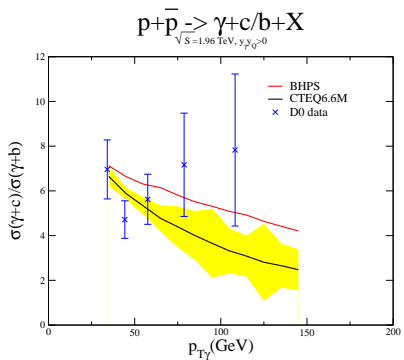
- CTEQ6.6M - radiatively generated charm - $c(x, \mu = m_c) = 0$
- Three intrinsic charm models - Non-perturbative charm component of the nucleon
 - Sea-like model - $c(x) \sim \bar{u}(x) + \bar{d}(x)$
 - BHPS - IC appears mainly at large x
 - Meson Cloud model - IC appears mainly at large x

Comparison between theory and data - IC $c + \gamma$



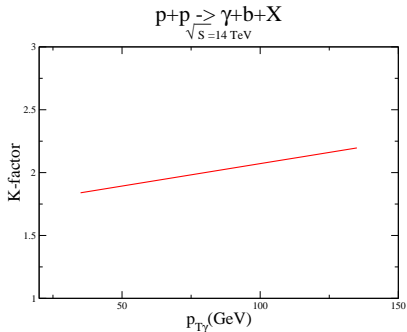
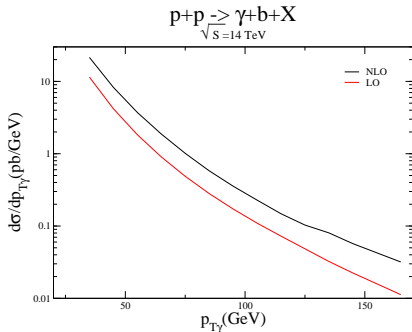
- With the use of the BHPS PDFs the cross section grows at large $p_{T\gamma}$, but is still below the data
- However if we are to look at the ratio of the c to b cross section ...

Ratio of c and b



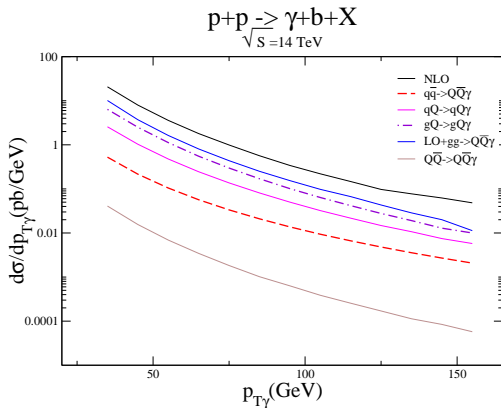
Things look better

LHC Predictions



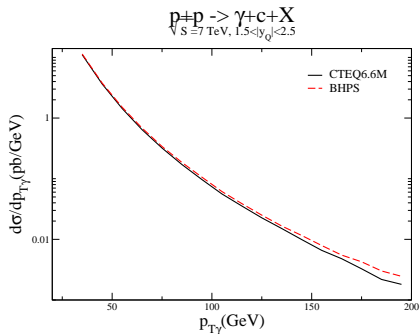
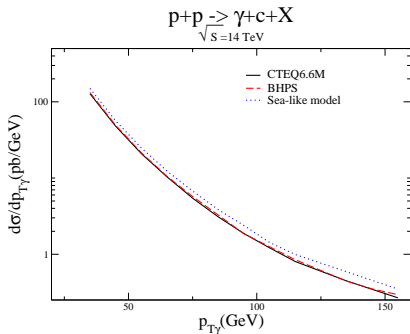
The difference between NLO and LO no longer grows at large $p_{T\gamma}$

Subprocess Contributions



- $q\bar{q}$ no longer dominates
 - No valence \bar{q}
- LO and subprocesses with initial g dominate \rightarrow
 - Larger center of mass energy probes lower values of x ($\sim \frac{p_T}{\sqrt{S}}$), region where gluon PDF dominates

Intrinsic Charm at the LHC



- Due to smaller x probed at the LHC can still test IC, but mainly the Sea-like model
- At 7 TeV and forward rapidity can slightly differentiate between BHPS and radiatively generated charm

Conclusions

Tevatron

- At Tevatron energies $q\bar{q}$ dominates the cross section at large $p_{T\gamma}$
- Good distinction between different IC models, can test for BHPS, Sea-like
- Good comparison between data and experiment for $\gamma + b$, for $\gamma + c$ some discrepancy at high p_T

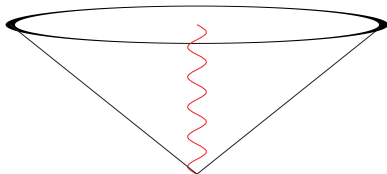
LHC

- At the LHC (pp 14 TeV) subprocesses with initial gluons and heavy quarks dominate
- Sensitivity to initial state heavy quarks
- Can test the low x behavior of heavy quark PDFs once LHC data is available

Isolation

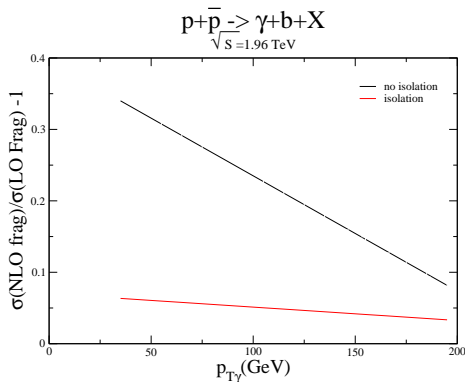
To be detected a direct photon needs to be isolated

- Helps acquire a reliable measurement of the photon's energy
- Helps reduce background from photons coming from the decay of hadrons



- Hadronic energy less than $E_h = \epsilon * E_\gamma$ in $R = \sqrt{\Delta\eta^2 + \Delta\phi^2}$
- Isolation requirements modeling the DØ requirements: $R_1 < 0.2$, $\epsilon_1 < 0.04$ and $R_2 < 0.4$, $\epsilon_2 < 0.07$

NLO Fragmentation



- No isolation - NLO fragmentation increases the cross section up to $\sim 30\%$
- Isolation requirement decreases the NLO fragmentation contribution to a few %

Photon Fragmentation Function

$$\frac{dD_{\gamma/Q}(z)}{dt} = \frac{\alpha}{2\pi} P_{\gamma \leftarrow Q}(z) + \frac{\alpha_s}{2\pi} \int \frac{dy}{y} [D_{\gamma/Q}(z/y) P_{Q \leftarrow Q}(y) + D_{\gamma/g}(z/y) P_{g \leftarrow Q}(y)]$$

$$\frac{dD_{\gamma/g}(z)}{dt} = \frac{\alpha_s}{2\pi} \int \frac{dy}{y} [D_{\gamma/Q}(z/y) P_{Q \leftarrow g}(y) + D_{\gamma/g}(z/y) P_{g \leftarrow g}(y)]$$