

Probing New Physics in $\Lambda_b \rightarrow \Lambda_c(\rightarrow \Lambda\pi)\ell\bar{\nu}$

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- LFU in the SM: $\mathcal{L}_{\text{Yukawa}} = -\bar{q}_L \mathcal{Y}_u u_R \phi - \bar{q}_L \mathcal{Y}_d d_R \phi - \bar{\ell}_L \mathcal{Y}_\ell \ell_R \phi + \text{h.c.}$
 $\rightarrow$ After SSB: **lepton mass** and (negligible) Higgs interaction.
- τ reconstruction difficult due to hadronic resonances ~ 1 GeV.
- Experimental measurements hint at violation of LFU.

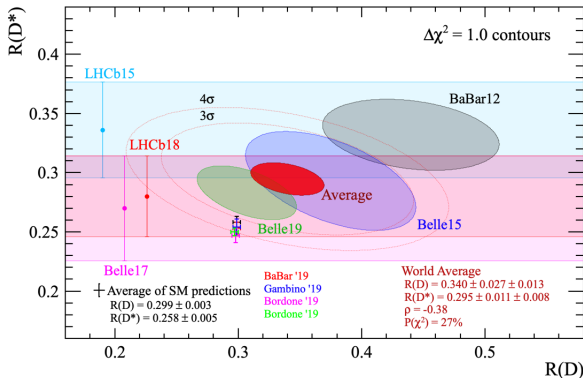
Semileptonic decay:

$$\bar{B} \rightarrow D^{(*)} \ell \bar{\nu}$$

[Rotondo, priv. com.]

$$\mathcal{R}(D^{(*)}) = \frac{\mathcal{B}(B \rightarrow D^{(*)} \tau \bar{\nu})}{\mathcal{B}(B \rightarrow D^{(*)} \ell \bar{\nu})}$$

$\sim 3.9\sigma$



- Model-independent analysis of New Physics is desirable.
- We consider the electroweak Hamiltonian at the scale $\mu \sim m_b$

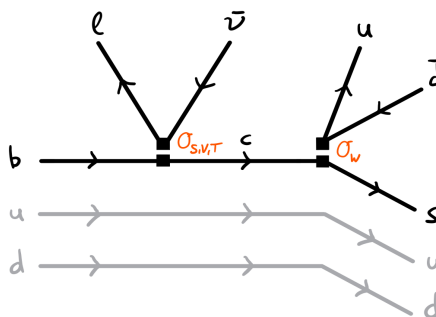
$$\mathcal{H}_{\text{eff}} = \frac{4\tilde{G}_F\tilde{V}_{cb}}{\sqrt{2}} \sum_i \mathcal{C}_i \mathcal{O}_i .$$

- Most general set of operators for $b \rightarrow c\ell\bar{\nu}_L$ up to dim-6 is

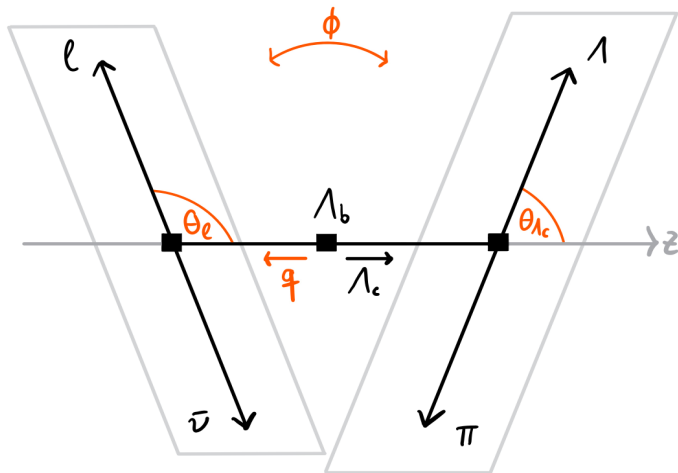
$$\begin{aligned} \mathcal{O}_{S,L} &= [\bar{c}P_L b][\bar{\ell}P_L \nu_\ell] , & \mathcal{O}_{S,R} &= [\bar{c}P_R b][\bar{\ell}P_L \nu_\ell] \\ \mathcal{O}_{V,L(R)} &= [\bar{c}\gamma^\mu P_{L(R)} b][\bar{\ell}\gamma_\mu P_L \nu_\ell] , & \mathcal{O}_{V,R} &= [\bar{c}\gamma^\mu P_R b][\bar{\ell}\gamma_\mu P_L \nu_\ell] \\ \mathcal{O}_T &= [\bar{c}\sigma^{\mu\nu} b][\bar{\ell}\sigma_{\mu\nu} P_L \nu_\ell] . \end{aligned}$$

- Motivation: SMEFT is able to generate all operators above.
[Catà, Jung, 1505.05804].

$$\Lambda_b^0 \rightarrow \Lambda_c^+ (\rightarrow \Lambda^0 \pi^+) \ell^- \bar{\nu}_\ell$$



- Λ_b assumed to be unpolarized.
- Narrow width approximation of Λ_c since $c\tau_{\Lambda_c} \approx 60\mu m$.



$$m_\ell^2 \leq q^2 \leq (m_{\Lambda_b} - m_{\Lambda_c})^2, \quad \cos \theta_\ell, \theta_{\Lambda_c} \in [-1, 1], \quad \phi \in [0, 2\pi]$$

Angular Distribution

$$\frac{d^4\Gamma}{dq^2 d\cos\theta_\ell d\cos\theta_{\Lambda_c} d\phi} \equiv \frac{3}{8\pi} \frac{d\Gamma}{dq^2} K(q^2, \cos\theta_\ell, \cos\theta_{\Lambda_c}, \phi)$$

$$\begin{aligned} K = & (K_{1ss} \sin^2 \theta_\ell + K_{1cc} \cos^2 \theta_\ell + K_{1c} \cos \theta_\ell) \\ & + (K_{2ss} \sin^2 \theta_\ell + K_{2cc} \cos^2 \theta_\ell + K_{2c} \cos \theta_\ell) \cos \theta_{\Lambda_c} \\ & + (K_{3sc} \sin \theta_\ell \cos \theta_\ell + K_{3s} \sin \theta_\ell) \sin \theta_{\Lambda_c} \sin \phi \\ & + (K_{4sc} \sin \theta_\ell \cos \theta_\ell + K_{4s} \sin \theta_\ell) \sin \theta_{\Lambda_c} \cos \phi \end{aligned}$$

- 10 angular coefficients (observables).
- Number of K_i 's fixed by **angular momentum conservation**.

Results can be cast into compact expressions, e.g.

$$\begin{aligned} \frac{d\Gamma}{dq^2} K_{2ss} = & \frac{\alpha}{2} \operatorname{Re} \left\{ \left(1 + \frac{m_\ell^2}{q^2} \right) A_{\perp 1} A_{\parallel 1}^* + 2 A_{\perp 0} A_{\parallel 0}^* + 2 \frac{m_\ell^2}{q^2} A_{\perp t} A_{\parallel t}^* \right\} \\ & - 2 \alpha \frac{m_\ell}{\sqrt{q^2}} \operatorname{Re} \left\{ A_{\perp 0}^T A_{\parallel 0}^* + A_{\perp 1}^T A_{\parallel 1}^* + (\perp \leftrightarrow \parallel) \right\} \\ & + 2 \alpha \operatorname{Re} \left\{ \left(1 + \frac{m_\ell^2}{q^2} \right) A_{\perp 1}^T A_{\parallel 1}^{T*} + 2 \frac{m_\ell^2}{q^2} A_{\perp 0}^T A_{\parallel 0}^{T*} \right\}. \end{aligned}$$

- Transversity amplitudes:

$$A_{\perp t} \equiv \sqrt{2} \mathcal{N} \sqrt{s_+} \frac{m_{\Lambda_b} - m_{\Lambda_c}}{m_\ell} f_t^V(q^2) \left[\frac{m_\ell}{\sqrt{q^2}} \mathcal{C}_{V,+} + \frac{\sqrt{q^2}}{m_b - m_c} \mathcal{C}_{S,+} \right].$$

- In total 10, all of factorized form

$$A_{\perp i / \parallel i} \sim f * \mathcal{C}.$$

- ① $\Lambda_b \rightarrow \Lambda_c$ form factors ($3 + 3 + 4 = 10$),

$$\begin{aligned} \langle \Lambda_c | \bar{c} \gamma^\mu b | \Lambda_b \rangle \equiv & \bar{u}_{\Lambda_c} \left[f_t^V(q^2) (m_{\Lambda_b} - m_{\Lambda_c}) \frac{q^\mu}{q^2} \right. \\ & + f_0^V(q^2) \frac{m_{\Lambda_b} + m_{\Lambda_c}}{s_+} \left(p^\mu + k^\mu - \frac{q^\mu}{q^2} (m_{\Lambda_b}^2 - m_{\Lambda_c}^2) \right) \\ & \left. + f_\perp^V(q^2) \left(\gamma^\mu - \frac{2m_{\Lambda_c}}{s_+} p^\mu - \frac{2m_{\Lambda_b}}{s_+} k^\mu \right) \right] u_{\Lambda_b} . \end{aligned}$$

Lattice data available: [Detmold, Lehner, Meinel, 1503.01421].

- ② $\Lambda_c \rightarrow \Lambda \pi$ angular asymmetry parameter from [PDG, BESIII]:

$$\frac{d\Gamma^{\Lambda_c \rightarrow \Lambda \pi}}{d \cos \Theta_{\Lambda_c}} \sim 1 + \alpha \cos \Theta_{\Lambda_c} , \quad \alpha = -0.82 \pm 0.09 .$$

What would be the benefit of measuring this decay?

- Measurement of all angular observables by LHCb is possible:
[LHCb, 1709.01920]

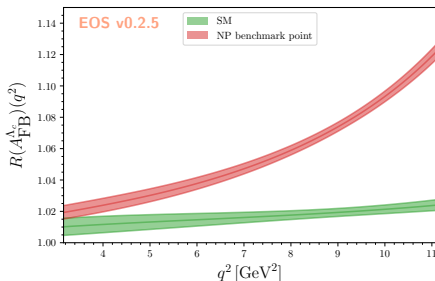
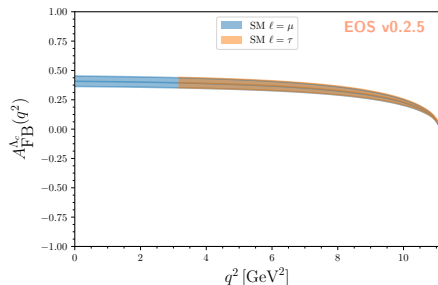
$$\begin{aligned}
 & N(\Lambda_b^0 \rightarrow \Lambda_c^+ (\rightarrow \Lambda^0 \pi^+) \mu^- \bar{\nu}) \\
 & \simeq N(\Lambda_b^0 \rightarrow \Lambda_c^+ (\rightarrow p K^- \pi^+) \mu^- \bar{\nu}) \times \frac{\mathcal{B}(\Lambda_c^+ \rightarrow \Lambda^0 \pi^+)}{\mathcal{B}(\Lambda_c^+ \rightarrow p K^- \pi^+)} \\
 & \simeq 0.5 \cdot 10^6 \simeq 10 \times N(\Lambda_b \rightarrow \Lambda_c^*(2625) \mu \bar{\nu}_\mu) .
 \end{aligned}$$

[Böer, Bordone, Graverini, Owen, Rotondo, van Dyk, 1801.08367].

- **25 real parameters** from non-redundant combinations of the $\mathcal{C}_i$'s.
 - For $m_\ell = 0$, angular observables are sensitive to 12 out of 25, for $m_\ell \neq 0$ to 22 out of 25 combinations. **High sensitivity!**
 - $\bar{B} \rightarrow D^{(*)} \tau \bar{\nu}$ distribution not sensitive to $\mathcal{C}_{V,-}$ and $\mathcal{C}_{S,-(+)}$,
 $\mathcal{C}_{i,\pm} = \mathcal{C}_{i,L} \pm \mathcal{C}_{i,R}$.

$$A_{\text{FB}}^{\Lambda_c} = K_{2ss} + \frac{1}{2}K_{2cc}$$

$$R(A_{\text{FB}}^{\Lambda_c}) \equiv \frac{[A_{\text{FB}}^{\Lambda_c}]_{\ell=\tau}}{[A_{\text{FB}}^{\Lambda_c}]_{\ell=\mu}}$$



- New Physics benchmark point [Murgui et. al., 1904.09311]:

$$\mathcal{C}_{V,L}^{(\tau)} = 1.15, \quad \mathcal{C}_{V,R}^{(\tau)} = 0, \quad \mathcal{C}_{S,L}^{(\tau)} = -0.3, \quad \mathcal{C}_{S,R}^{(\tau)} = +0.3, \quad \mathcal{C}_T^{(\tau)} = 0.$$

- $\delta_R \sim 1\%$, compared to $\delta_{\mathcal{R}(D)} \sim 2\%$.

- Complete four-fold distribution of $\Lambda_b \rightarrow \Lambda_c(\rightarrow \Lambda\pi)\ell\bar{\nu}$
→ For the **first time** with $m_\ell \neq 0$ in presence of all weak effective operators.
- **High sensitivity** to BSM parameters in the cascade decay.
→ Accessible set of combinations from Wilson coefficients large, 22 out of 25 combinations for $m_\ell \neq 0$.
- **Hadronic forward-backward asymmetry**
→ Attractive observable for LHCb for a first measurement.
→ Small theory uncertainty.

Thanks for your attention!

obs.	SM		BMP #1	BMP #2	BMP #3
	$\ell = \mu$	$\ell = \tau$	$\ell = \tau$	$\ell = \tau$	$\ell = \tau$
K_{1cc}	$+0.206 \pm 0.004$	$+0.310 \pm 0.001$	$+0.311 \pm 0.000$	$+0.307 \pm 0.001$	$+0.343 \pm 0.001$
K_{1c}	-0.134 ± 0.004	$+0.016 \pm 0.003$	$+0.037 \pm 0.002$	-0.073 ± 0.002	$+0.080 \pm 0.003$
K_{2ss}	$+0.288 \pm 0.032$	$+0.221 \pm 0.024$	$+0.228 \pm 0.025$	$+0.167 \pm 0.018$	$+0.252 \pm 0.027$
K_{2cc}	$+0.115 \pm 0.013$	$+0.183 \pm 0.020$	$+0.193 \pm 0.021$	$+0.122 \pm 0.013$	$+0.130 \pm 0.013$
K_{2c}	-0.164 ± 0.018	-0.031 ± 0.004	-0.017 ± 0.003	-0.123 ± 0.013	-0.080 ± 0.009
K_{4sc}	$+0.063 \pm 0.008$	$+0.023 \pm 0.003$	$+0.022 \pm 0.002$	$+0.026 \pm 0.003$	$+0.001 \pm 0.000$
K_{4s}	$+0.125 \pm 0.015$	$+0.063 \pm 0.007$	$+0.065 \pm 0.007$	$+0.161 \pm 0.017$	$+0.039 \pm 0.004$

$$\begin{aligned}
 \sigma_{V,1}^{\pm} &= \frac{1}{2} |\mathcal{C}_{V,\pm}|^2, & \sigma_{V,2} &= -\mathcal{C}_{V,-} \mathcal{C}_{V,+}^*, & \sigma_{T,1} &= \frac{1}{2} |\mathcal{C}_T|^2, \\
 \sigma_{S,1}^{\pm} &= \frac{1}{2} |\mathcal{C}_{S,\pm}|^2, & \sigma_{S,2} &= -\mathcal{C}_{S,-} \mathcal{C}_{S,+}^*, & \sigma_{VT,1}^{\pm} &= \frac{1}{2} \mathcal{C}_{V,\pm} \mathcal{C}_T^*, \\
 \sigma_{VS,1}^{\pm} &= \frac{1}{2} \mathcal{C}_{V,\pm} \mathcal{C}_{S,\pm}^*, & \sigma_{VS,2} &= -\mathcal{C}_{V,-} \mathcal{C}_{S,+}^*, & \sigma_{ST,1}^{\pm} &= \frac{1}{2} \mathcal{C}_{S,\pm} \mathcal{C}_T^*. \\
 & & \sigma_{SV,2} &= -\mathcal{C}_{S,-} \mathcal{C}_{V,+}^*, & &
 \end{aligned}$$

$$\Sigma(m_\ell = 0) \equiv \{\sigma_{V,1}^{\pm}, \sigma_{S,1}^{\pm}, \sigma_{T,1}, \text{Re } \sigma_{V,2}, \text{Im } \sigma_{V,2}, \text{Re } \sigma_{S,2}, \text{Re } \sigma_{ST,1}^{\pm}, \text{Im } \sigma_{ST,1}^{\pm}\}$$

$$\Sigma \equiv \Sigma(m_\ell = 0) \cup \{\text{Re } \sigma_{VS,1}^{\pm}, \text{Im } \sigma_{VS,1}^{\pm}, \text{Re } \sigma_{VS,2}, \text{Re } \sigma_{SV,2}, \text{Re } \sigma_{VT,1}^{\pm}, \text{Im } \sigma_{VT,1}^{\pm}\}$$

We crosschecked our results with

- [Datta, Kamali, Meinel, Rashed, 1702.02243],
- [Shivashankara, Wu, Datta, 1502.07230]*,
- [Gutsche, Ivanov, Körner, Lyubovitskij, 1502.04864].