

Can we improve energy density functionals? —A perturbative method—

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Collaborators in This Talk

- Giacomo Accorto (U. Zagreb)
- Haozhao Liang (U. Tokyo)
- Tamara Nikšić (U. Zagreb)
- Daisuke Ohashi (U. Tokyo → Private Company)
- Dario Vretenar (U. Zagreb)

Density Functional Theory (DFT) and Energy Density Functional (EDF)

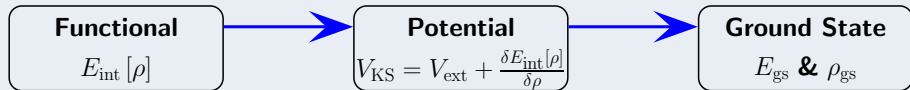
- Hohenberg-Kohn thm. guarantees existence of the unique EDF E_{int}
- Kohn-Sham scheme provides a practical method
- Kohn-Sham orbital φ_j obeys

$$\left[-\frac{\hbar^2}{2M} \Delta + V_{\text{KS}}(\mathbf{r}) \right] \varphi_j(\mathbf{r}) = \varepsilon_j \varphi_j(\mathbf{r})$$

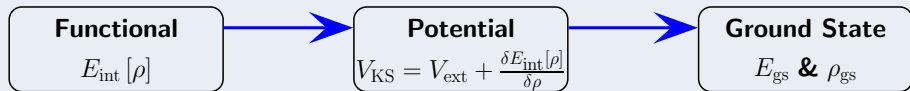
$$V_{\text{KS}}(\mathbf{r}) = V_{\text{ext}}(\mathbf{r}) + \left. \frac{\delta E_{\text{int}}[\rho]}{\delta \rho} \right|_{\rho=\rho_{\text{gs}}(\mathbf{r})} \quad \rho_{\text{gs}}(\mathbf{r}) = \sum |\varphi_j(\mathbf{r})|^2$$

- If the exact E_{int} were known, DFT would provide the exact ρ_{gs} and E_{gs}
→ Practically, approximation is needed to derive E_{int}
- Accuracy of E_{int} governs accuracy of DFT calc.

Usual DFT Calculation



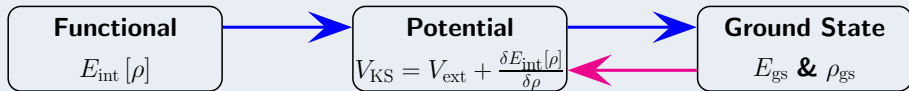
Usual DFT Calculation



Q1. If we know ρ_{gs} , can we derive V_{KS} ?

Introduction

Inverse Kohn-Sham Method



Q1. If we know ρ_{gs} , can we derive V_{KS} ?

A1. Yes! Inverse Kohn-Sham Method

Inverse Kohn-Sham Method

- Assume that ρ_{gs} of a system (e.g., ^{208}Pb) is known
- V_{KS} for the system (e.g., ^{208}Pb) reads

$$V_{\text{KS}}(\mathbf{r}) = \frac{\sum \varphi_j^*(\mathbf{r}) \left(\varepsilon_j + \frac{\hbar^2}{2M} \Delta \right) \varphi_j(\mathbf{r})}{\rho_{\text{gs}}(\mathbf{r})}$$

Practically, V_{KS} is derived iteratively

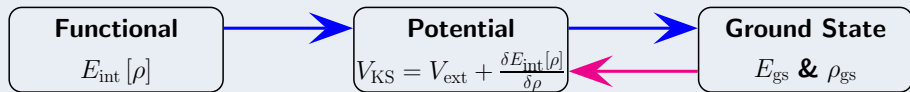
$$V_{\text{KS}}^n(\mathbf{r}) = \frac{\sum \varphi_j^{n*}(\mathbf{r}) \left(\varepsilon_j^n + \frac{\hbar^2}{2M} \Delta \right) \varphi_j^n(\mathbf{r})}{\rho_{\text{gs}}(\mathbf{r})} = \frac{\rho^n(\mathbf{r})}{\rho_{\text{gs}}(\mathbf{r})} V_{\text{KS}}^{n-1}(\mathbf{r})$$

$$\left[-\frac{\hbar^2}{2M} \Delta + V_{\text{KS}}^{n-1}(\mathbf{r}) \right] \varphi_j^n(\mathbf{r}) = \varepsilon_j^n \varphi_j^n(\mathbf{r}) \quad \rho^n(\mathbf{r}) = \sum |\varphi_j^n(\mathbf{r})|^2$$

- This potential V_{KS} can only be used for the system (e.g., ^{208}Pb)
→ This V_{KS} cannot be used directly for the other systems (e.g., ^{48}Ca)

Wang and Parr. *Phys. Rev. A* **47**, R1591 (1993)

Inverse Kohn-Sham Method

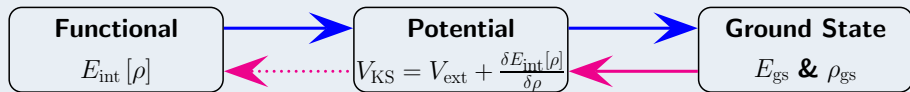


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Introduction

Inverse Kohn-Sham Method

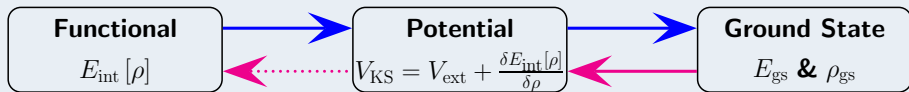


Q1. If we know ρ_{gs} , can we derive V_{KS} ?

A1. Yes! Inverse Kohn-Sham Method

Q2. Can we derive E_{int} from V_{KS} ?

Inverse Kohn-Sham Method



Q1. If we know ρ_{gs} , can we derive V_{KS} ?

A1. Yes! Inverse Kohn-Sham Method

Q2. Can we derive E_{int} from V_{KS} ?

A2. Today's Topic!! "IKS-DFPT"

Brief Idea of IKS-DFPT

- DFT calculation is succeeded
→ We assume that a “known” EDF \tilde{E}_{int} is good enough
- Improve \tilde{E}_{int} using some “exact” G.S. densities
- Difference between \tilde{E}_{int} & $E_{\text{int}}^{\text{exact}}$ is treated as 1st order perturbation
- Idea of **Density Functional Perturbation Theory (DFPT)** is used
Original DFPT is usually used for phonon calculation for solid with considering perturbation for V_{ext}
- E_{gs} is formally derived by two methods
 - Inverse Kohn-Sham method
 - Density functional perturbation theory
- Using the fact that two E_{gs} should be identical, $E_{\text{int}}^{\text{exact}} - \tilde{E}_{\text{int}}$ is derived

Naito, Ohashi, and Liang. *J. Phys. B* **52**, 245003 (2019)

Perturbation Expansion for $E_{\text{int}}, \rho, \varphi_j$

$$E_{\text{int}}^{\text{exact}}[\rho] = \tilde{E}_{\text{int}}[\rho] + \lambda E_{\text{int}}^{(1)}[\rho] + O(\lambda^2)$$

$$V_{\text{ext}}^{\text{exact}}(\mathbf{r}) = \tilde{V}_{\text{ext}}(\mathbf{r})$$

$$\rho_{\text{gs}}^{\text{exact}}(\mathbf{r}) = \tilde{\rho}_{\text{gs}}(\mathbf{r}) + \lambda \rho_{\text{gs}}^{(1)}(\mathbf{r}) + O(\lambda^2)$$

$$\varphi_j^{\text{exact}}(\mathbf{r}) = \tilde{\varphi}_j(\mathbf{r}) + \lambda \varphi_j^{(1)}(\mathbf{r}) + O(\lambda^2)$$

Note: $\rho_{\text{gs}}^{\text{exact}}$ is known

Assumption: Orthogonal Condition

$$\int \tilde{\varphi}_j^*(\mathbf{r}) \varphi_j^{(1)}(\mathbf{r}) d\mathbf{r} = 0$$

Density

$$\tilde{\rho}_{\text{gs}}(\mathbf{r}) = \sum |\tilde{\varphi}_j(\mathbf{r})|^2$$

$$\rho_{\text{gs}}^{(1)}(\mathbf{r}) = \sum \left[\varphi_j^{(1)*}(\mathbf{r}) \tilde{\varphi}_j(\mathbf{r}) + \tilde{\varphi}_j^*(\mathbf{r}) \varphi_j^{(1)}(\mathbf{r}) \right]$$

Ground-State Energy from IKS

$$E_{\text{gs}}^{\text{exact}} = \sum \varepsilon_j^{\text{exact}} + E_{\text{int}}^{\text{exact}}[\rho_{\text{gs}}^{\text{exact}}] - \int \frac{\delta E_{\text{int}}^{\text{exact}}[\rho_{\text{gs}}^{\text{exact}}]}{\delta \rho(\mathbf{r})} \rho_{\text{gs}}^{\text{exact}}(\mathbf{r}) d\mathbf{r}$$

=

$\varepsilon_j^{\text{exact}}$ can be derived by inverse Kohn-Sham method

Ground-State Energy from IKS

$$\begin{aligned}
 E_{\text{gs}}^{\text{exact}} &= \sum \varepsilon_j^{\text{exact}} + E_{\text{int}}^{\text{exact}}[\rho_{\text{gs}}^{\text{exact}}] - \int \frac{\delta E_{\text{int}}^{\text{exact}}[\rho_{\text{gs}}^{\text{exact}}]}{\delta \rho(\mathbf{r})} \rho_{\text{gs}}^{\text{exact}}(\mathbf{r}) d\mathbf{r} \\
 &= \sum \varepsilon_j^{\text{exact}} + \tilde{E}_{\text{int}}[\rho_{\text{gs}}^{\text{exact}}] + \lambda E_{\text{int}}^{(1)}[\rho_{\text{gs}}^{\text{exact}}] \\
 &\quad - \int \frac{\delta \tilde{E}_{\text{int}}[\rho_{\text{gs}}^{\text{exact}}]}{\delta \rho(\mathbf{r})} \rho_{\text{gs}}^{\text{exact}}(\mathbf{r}) d\mathbf{r} - \lambda \int \frac{\delta E_{\text{int}}^{(1)}[\rho_{\text{gs}}^{\text{exact}}]}{\delta \rho(\mathbf{r})} \rho_{\text{gs}}^{\text{exact}}(\mathbf{r}) d\mathbf{r} + O(\lambda^2)
 \end{aligned}$$

$\varepsilon_j^{\text{exact}}$ can be derived by inverse Kohn-Sham method

Ground-State Energy from DFPT

$$E_{\text{gs}}^{\text{exact}} = \tilde{E}_{\text{gs}} + \lambda \left. \frac{dE_{\text{int}}^{\text{exact}}[\tilde{\rho}_{\text{gs}}]}{d\lambda} \right|_{\lambda=0} + O(\lambda^2)$$

=

=

Please see our paper for the deviation of this equation, since it is complicated

Ground-State Energy from DFPT

$$\begin{aligned} E_{\text{gs}}^{\text{exact}} &= \tilde{E}_{\text{gs}} + \lambda \left. \frac{dE_{\text{int}}^{\text{exact}}[\tilde{\rho}_{\text{gs}}]}{d\lambda} \right|_{\lambda=0} + O(\lambda^2) \\ &= \tilde{E}_{\text{gs}} + \lambda E_{\text{int}}^{(1)}[\tilde{\rho}_{\text{gs}}] + O(\lambda^2) \\ &= \end{aligned}$$

Please see our paper for the deviation of this equation, since it is complicated

Ground-State Energy from DFPT

$$\begin{aligned}
 E_{\text{gs}}^{\text{exact}} &= \tilde{E}_{\text{gs}} + \lambda \left. \frac{dE_{\text{int}}^{\text{exact}}[\tilde{\rho}_{\text{gs}}]}{d\lambda} \right|_{\lambda=0} + O(\lambda^2) \\
 &= \tilde{E}_{\text{gs}} + \lambda E_{\text{int}}^{(1)}[\tilde{\rho}_{\text{gs}}] + O(\lambda^2) \\
 &= \sum \tilde{\epsilon}_j + \tilde{E}_{\text{int}}[\tilde{\rho}_{\text{gs}}] - \int \frac{\delta \tilde{E}_{\text{int}}[\tilde{\rho}_{\text{gs}}]}{\delta \rho(\mathbf{r})} \tilde{\rho}_{\text{gs}}(\mathbf{r}) d\mathbf{r} + \lambda E_{\text{int}}^{(1)}[\tilde{\rho}_{\text{gs}}] + O(\lambda^2)
 \end{aligned}$$

Please see our paper for the deviation of this equation, since it is complicated

Equation for IKS-DFPT

Compare two expression and neglect $O(\lambda^2)$ term

$$\begin{aligned} \tilde{E}_{\text{gs}} + \lambda E_{\text{int}}^{(1)}[\tilde{\rho}_{\text{gs}}] &\simeq \sum \varepsilon_j^{\text{exact}} + \tilde{E}_{\text{int}}[\rho_{\text{gs}}^{\text{exact}}] + \lambda E_{\text{int}}^{(1)}[\rho_{\text{gs}}^{\text{exact}}] \\ &\quad - \int \frac{\delta \tilde{E}_{\text{int}}[\rho_{\text{gs}}^{\text{exact}}]}{\delta \rho(\mathbf{r})} \rho_{\text{gs}}^{\text{exact}}(\mathbf{r}) d\mathbf{r} - \lambda \int \frac{\delta E_{\text{int}}^{(1)}[\rho_{\text{gs}}^{\text{exact}}]}{\delta \rho(\mathbf{r})} \rho_{\text{gs}}^{\text{exact}}(\mathbf{r}) d\mathbf{r} \end{aligned}$$

Therefore, We Should Solve

$$\begin{aligned} \lambda E_{\text{int}}^{(1)}[\tilde{\rho}_{\text{gs}}] - \lambda E_{\text{int}}^{(1)}[\rho_{\text{gs}}^{\text{exact}}] + \lambda \int \frac{\delta E_{\text{int}}^{(1)}[\rho_{\text{gs}}^{\text{exact}}]}{\delta \rho(\mathbf{r})} \rho_{\text{gs}}^{\text{exact}}(\mathbf{r}) d\mathbf{r} \\ \simeq \sum \varepsilon_j^{\text{exact}} + \tilde{E}_{\text{int}}[\rho_{\text{gs}}^{\text{exact}}] - \int \frac{\delta \tilde{E}_{\text{int}}[\rho_{\text{gs}}^{\text{exact}}]}{\delta \rho(\mathbf{r})} \rho_{\text{gs}}^{\text{exact}}(\mathbf{r}) d\mathbf{r} - \tilde{E}_{\text{gs}} \end{aligned}$$

Equation for IKS-DFPT

Compare two expression and neglect $O(\lambda^2)$ term

$$\begin{aligned} \tilde{E}_{\text{gs}} + \lambda E_{\text{int}}^{(1)}[\tilde{\rho}_{\text{gs}}] &\simeq \sum \varepsilon_j^{\text{exact}} + \tilde{E}_{\text{int}}[\rho_{\text{gs}}^{\text{exact}}] + \lambda E_{\text{int}}^{(1)}[\rho_{\text{gs}}^{\text{exact}}] \\ &\quad - \int \frac{\delta \tilde{E}_{\text{int}}[\rho_{\text{gs}}^{\text{exact}}]}{\delta \rho(\mathbf{r})} \rho_{\text{gs}}^{\text{exact}}(\mathbf{r}) d\mathbf{r} - \lambda \int \frac{\delta E_{\text{int}}^{(1)}[\rho_{\text{gs}}^{\text{exact}}]}{\delta \rho(\mathbf{r})} \rho_{\text{gs}}^{\text{exact}}(\mathbf{r}) d\mathbf{r} \end{aligned}$$

Therefore, We Should Solve

$$\begin{aligned} \lambda E_{\text{int}}^{(1)}[\tilde{\rho}_{\text{gs}}] - \lambda E_{\text{int}}^{(1)}[\rho_{\text{gs}}^{\text{exact}}] + \lambda \int \frac{\delta E_{\text{int}}^{(1)}[\rho_{\text{gs}}^{\text{exact}}]}{\delta \rho(\mathbf{r})} \rho_{\text{gs}}^{\text{exact}}(\mathbf{r}) d\mathbf{r} \\ \simeq \sum \varepsilon_j^{\text{exact}} + \tilde{E}_{\text{int}}[\rho_{\text{gs}}^{\text{exact}}] - \int \frac{\delta \tilde{E}_{\text{int}}[\rho_{\text{gs}}^{\text{exact}}]}{\delta \rho(\mathbf{r})} \rho_{\text{gs}}^{\text{exact}}(\mathbf{r}) d\mathbf{r} - \tilde{E}_{\text{gs}} = C[\rho_{\text{gs}}^{\text{exact}}] \end{aligned}$$

Right-hand side can be calculated exactly as $C[\rho_{\text{gs}}^{\text{exact}}]$ (constant)

Our Task

Derive $\lambda E_{\text{int}}^{(1)}$ from

$$\lambda E_{\text{int}}^{(1)}[\tilde{\rho}_{\text{gs}}^{\text{exact}}] - \lambda E_{\text{int}}^{(1)}[\rho_{\text{gs}}^{\text{exact}}] + \lambda \int \frac{\delta E_{\text{int}}^{(1)}[\rho_{\text{gs}}^{\text{exact}}]}{\delta \rho(\mathbf{r})} \rho_{\text{gs}}^{\text{exact}}(\mathbf{r}) d\mathbf{r} = C[\rho_{\text{gs}}^{\text{exact}}]$$

This equation is a functional equation: Difficult to solve!!

How to Solve It?

Our Task

Derive $\lambda E_{\text{int}}^{(1)}$ from

$$\lambda E_{\text{int}}^{(1)}[\tilde{\rho}_{\text{gs}}^{\text{exact}}] - \lambda E_{\text{int}}^{(1)}[\rho_{\text{gs}}^{\text{exact}}] + \lambda \int \frac{\delta E_{\text{int}}^{(1)}[\rho_{\text{gs}}^{\text{exact}}]}{\delta \rho(\mathbf{r})} \rho_{\text{gs}}^{\text{exact}}(\mathbf{r}) d\mathbf{r} = C[\rho_{\text{gs}}^{\text{exact}}]$$

This equation is a functional equation: Difficult to solve!!

How to Solve It?

- 1 Assume the form of $\lambda E_{\text{int}}^{(1)}$
- 2 If $\lambda E_{\text{int}}^{(1)}$ has n parameters, substitute ρ_{gs} and ε_j of n kind of systems to the equation
- 3 Solve simultaneously

Benchmark Calculation (1): Atoms

Assumption for $\lambda E_{\text{int}}^{(1)}$

$$\lambda E_{\text{int}}^{(1)}[\rho] = \lambda A \int [\rho(\mathbf{r})]^\alpha d\mathbf{r}$$

Unknown parameters: λA and α

Equation for IKS-DFPT1

$$\lambda A \int \left\{ [\tilde{\rho}_{\text{gs}}(\mathbf{r})]^\alpha + (\alpha - 1) [\rho_{\text{gs}}^{\text{exact}}(\mathbf{r})]^\alpha \right\} d\mathbf{r} = C [\rho_{\text{gs}}]$$

To derive λA and α , two densities are required

Iterative IKS-DFPT: Calculated Functional at n -th Step

$$E_{\text{int}}^{n\text{-th}}[\rho] = \tilde{E}_{\text{int}}[\rho] + \sum_{j=1}^n \lambda A_j \int [\rho(\mathbf{r})]^{\alpha_j} d\mathbf{r}$$

$$E_{\text{int}}[\rho] = E_{\text{H}}[\rho] + E_{\text{x}}[\rho] + E_{\text{c}}[\rho]$$

E_{H} Coulomb Hartree EDF

E_{x} Coulomb exchange EDF (LDA is used in this work)

E_{c} Coulomb correlation EDF (LDA is used in this work)

Setup

- Noble gas atoms are used for ρ_{gs}
- Two patterns are tested (Today I only show pattern 1)

Pattern 1

Known \tilde{E}_{int} E_{H}

“Exact” E_{int} $E_{\text{H}} + E_{\text{x}}$

Can we reproduce E_{x} ?

Pattern 2

Known \tilde{E}_{int} $E_{\text{H}} + E_{\text{x}}$

“Exact” E_{int} $E_{\text{H}} + E_{\text{x}} + E_{\text{c}}$

Can we reproduce E_{c} ?

Benchmark Calculation (1): Atoms

Energy: $E_{\text{int}}^{\text{exact}}$ is Hartree & LDA exchange (Hartree atomic unit)

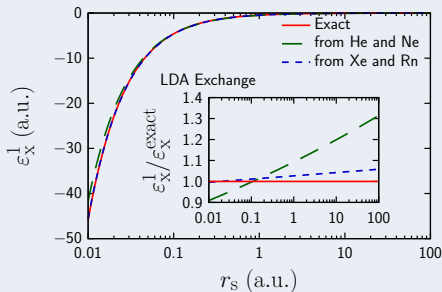
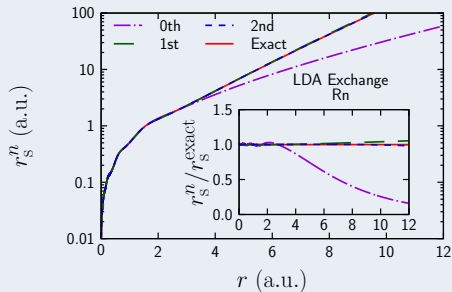
n	α_n	λA_n	$E_{\text{gs}}^{n\text{-th}}$ of Xe	$E_{\text{gs}}^{n\text{-th}}$ of Rn
0			-7054.6485	-21479.344
1	1.3311445	-0.7558229	-7224.9365	-21852.010
2	1.0436323	0.0306234	-7223.0601	-21848.894
Target	1.3333333	-0.7385588	-7223.1853	-21848.954

Hartree Atomic Unit

$$m_e = 1, \quad \hbar = 1, \quad e^2 = 1, \quad 4\pi\epsilon_0 = 1, \quad c = \frac{1}{\alpha}$$

Benchmark Calculation (1): Atoms

Density and energy density: $E_{\text{int}}^{\text{exact}}$ is Hartree & LDA exchange



$$r_s = \left(\frac{3}{4\pi\rho} \right)^{1/3}, \quad E_x[\rho] = \int \epsilon_x^1(r_s(\mathbf{r})) \rho(\mathbf{r}) d\mathbf{r}$$

Relativistic Extension

- There are scalar KS potential S_{KS} and vector KS potential V_{KS}

- There are scalar density ρ_{S} and vector density ρ_{V}

$$\rho_{\text{S}}(\mathbf{r}) = \sum \varphi_j^\dagger(\mathbf{r}) \beta \varphi_j(\mathbf{r}) = \sum \bar{\varphi}_j(\mathbf{r}) \varphi_j(\mathbf{r}) \quad \rho_{\text{V}}(\mathbf{r}) = \sum \varphi_j^\dagger(\mathbf{r}) \varphi_j(\mathbf{r})$$

- S.P. orbital obeys $[\boldsymbol{\alpha} \cdot \mathbf{p} + \beta(m + S_{\text{KS}}(\mathbf{r})) + V_{\text{KS}}(\mathbf{r})] \varphi_j(\mathbf{r}) = \varepsilon_j(\mathbf{r})$

Non-Relativistic Inverse Kohn-Sham Method

$$V_{\text{KS}}(\mathbf{r}) = -\frac{1}{\rho(\mathbf{r})} \sum \varphi_j^*(\mathbf{r}) \left(-\frac{\hbar^2}{2M} \Delta - \varepsilon_j \right) \varphi_j(\mathbf{r})$$

Wang and Parr. *Phys. Rev. A* **47**, R1591 (1993)

Relativistic Inverse Kohn-Sham Method

$$V_{\text{KS}}(\mathbf{r}) + S_{\text{KS}}(\mathbf{r}) = -M - \frac{1}{\rho_{\text{V}}(\mathbf{r}) + \rho_{\text{S}}(\mathbf{r})} \sum [\varphi_j^\dagger(\mathbf{r}) + \bar{\varphi}_j(\mathbf{r})] (\boldsymbol{\alpha} \cdot \mathbf{p} - \varepsilon_j) \varphi_j(\mathbf{r})$$

$$V_{\text{KS}}(\mathbf{r}) - S_{\text{KS}}(\mathbf{r}) = +M - \frac{1}{\rho_{\text{V}}(\mathbf{r}) - \rho_{\text{S}}(\mathbf{r})} \sum [\varphi_j^\dagger(\mathbf{r}) - \bar{\varphi}_j(\mathbf{r})] (\boldsymbol{\alpha} \cdot \mathbf{p} - \varepsilon_j) \varphi_j(\mathbf{r})$$

Accorto, Naito, Liang, Nikšić, and Vretenar. *Phys. Rev. C* **103**, 044304 (2021)

Non-Relativistic IKS-DFPT

$$\begin{aligned} & \lambda E_{\text{int}}^{(1)}[\tilde{\rho}_{\text{gs}}] - \lambda E_{\text{int}}^{(1)}[\rho_{\text{gs}}^{\text{exact}}] + \lambda \int \frac{\delta E_{\text{int}}^{(1)}[\rho_{\text{gs}}^{\text{exact}}]}{\delta \rho(\mathbf{r})} \rho_{\text{gs}}^{\text{exact}}(\mathbf{r}) d\mathbf{r} \\ & \simeq \sum \varepsilon_j^{\text{exact}} - \tilde{E}_{\text{gs}} + \tilde{E}_{\text{int}}[\rho_{\text{gs}}^{\text{exact}}] - \int \frac{\delta \tilde{E}_{\text{int}}[\rho_{\text{gs}}^{\text{exact}}]}{\delta \rho(\mathbf{r})} \rho_{\text{gs}}^{\text{exact}}(\mathbf{r}) d\mathbf{r} \end{aligned}$$

Naito, Ohashi, and Liang. *J. Phys. B* **52**, 245003 (2019)

Relativistic IKS-DFPT

$$\begin{aligned} & \lambda E_{\text{int}}^{(1)}[\tilde{\rho}_{\text{V, gs}}, \tilde{\rho}_{\text{S, gs}}] - \lambda E_{\text{int}}^{(1)}[\rho_{\text{V, gs}}^{\text{exact}}, \rho_{\text{S, gs}}^{\text{exact}}] \\ & + \lambda \int \frac{\delta E_{\text{int}}^{(1)}[\rho_{\text{V, gs}}^{\text{exact}}, \rho_{\text{S, gs}}^{\text{exact}}]}{\delta \rho_{\text{V}}(\mathbf{r})} \rho_{\text{V, gs}}^{\text{exact}}(\mathbf{r}) d\mathbf{r} + \lambda \int \frac{\delta E_{\text{int}}^{(1)}[\rho_{\text{V, gs}}^{\text{exact}}, \rho_{\text{S, gs}}^{\text{exact}}]}{\delta \rho_{\text{S}}(\mathbf{r})} \rho_{\text{S, gs}}^{\text{exact}}(\mathbf{r}) d\mathbf{r} \\ & \simeq \sum \varepsilon_j^{\text{exact}} - \tilde{E}_{\text{gs}} + \tilde{E}_{\text{int}}[\rho_{\text{V, gs}}^{\text{exact}}, \rho_{\text{S, gs}}^{\text{exact}}] \\ & - \int \frac{\delta \tilde{E}_{\text{int}}[\rho_{\text{V, gs}}^{\text{exact}}, \rho_{\text{S, gs}}^{\text{exact}}]}{\delta \rho_{\text{V}}(\mathbf{r})} \rho_{\text{V, gs}}^{\text{exact}}(\mathbf{r}) d\mathbf{r} - \int \frac{\delta \tilde{E}_{\text{int}}[\rho_{\text{V, gs}}^{\text{exact}}, \rho_{\text{S, gs}}^{\text{exact}}]}{\delta \rho_{\text{S}}(\mathbf{r})} \rho_{\text{S, gs}}^{\text{exact}}(\mathbf{r}) d\mathbf{r} \end{aligned}$$

Accorto, Naito, Liang, Nikšić, and Vretnar. *Phys. Rev. C* **103**, 044304 (2021)

DD-PC1 (without Coulomb)

$$S_{\text{DD-PC1}} [\rho_V, \rho_S, \vec{\rho}_{\text{TV}}] = \alpha_S(\rho_V) \rho_S + \delta_S \Delta \rho_S$$

$$V_{\text{DD-PC1}} [\rho_V, \rho_S, \vec{\rho}_{\text{TV}}] = \alpha_V(\rho_V) \rho_V + \alpha_{\text{TV}}(\rho_V) \vec{\rho}_{\text{TV}}$$

$$\Sigma_{\text{R}} [\rho_V, \rho_S, \vec{\rho}_{\text{TV}}] = \frac{1}{2} \frac{\partial \alpha_S}{\partial \rho} \rho_S^2 + \frac{1}{2} \frac{\partial \alpha_V}{\partial \rho} \rho_V^2 + \frac{1}{2} \frac{\partial \alpha_{\text{TV}}}{\partial \rho} \rho_{\text{TV}}^2$$

$$\alpha_S(\rho_V) = a_S + \left(b_S + c_S \frac{\rho_V}{\rho_{\text{sat}}} \right) e^{-d_S \rho_V / \rho_{\text{sat}}}$$

$$\alpha_V(\rho_V) = a_V + b_V e^{-d_V \rho_V / \rho_{\text{sat}}}$$

$$\alpha_{\text{TV}}(\rho_V) = b_{\text{TV}} e^{-d_{\text{TV}} \rho_V / \rho_{\text{sat}}}$$

Nikšić, Vretenar, and Ring. *Phys. Rev. C* **78**, 034318 (2008)

Benchmark Calculation

Start “LDA terms (only a_S and a_V)” of DD-PC1

Goal All terms of DD-PC1 Can we reproduce b and c ? (d are fixed)

Systems: $N = Z$ nuclei w/o Coulomb (${}^8_8\text{16}$, ${}^{28}_{28}\text{56}$, ${}^{50}_{50}\text{100}$)

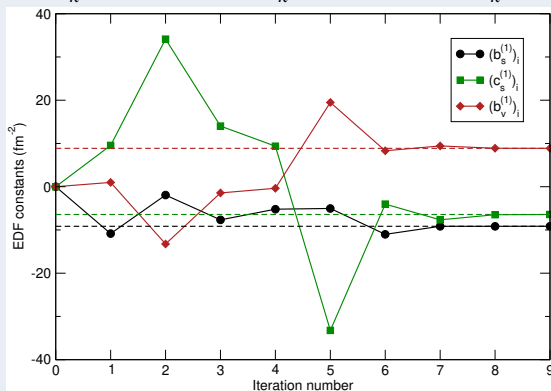
Benchmark Results

Iterative method is used

$$b_S = \sum_k b_S^{k\text{-th}}$$

$$c_S = \sum_k c_S^{k\text{-th}}$$

$$b_V = \sum_k b_V^{k\text{-th}}$$



Accorto, Naito, Liang, Nikšić, and Vretenar. *Phys. Rev. C* **103**, 044304 (2021)

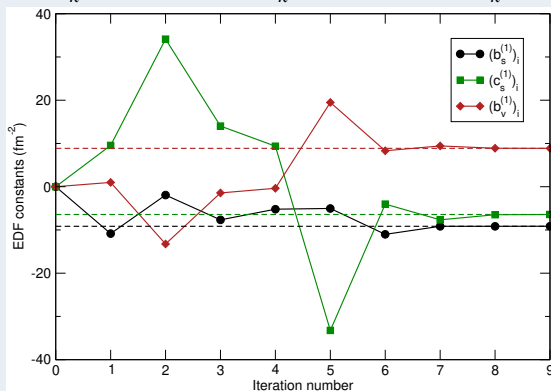
Benchmark Results

Iterative method is used

$$b_S = \sum_k b_S^{k\text{-th}}$$

$$c_S = \sum_k c_S^{k\text{-th}}$$

$$b_V = \sum_k b_V^{k\text{-th}}$$



Well reproduced!!

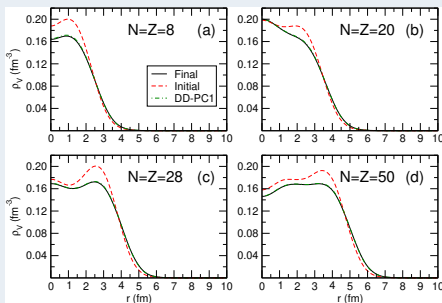
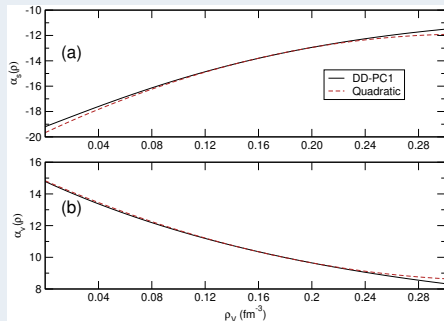
Accorto, Naito, Liang, Nikšić, and Vretenar. *Phys. Rev. C* **103**, 044304 (2021)

Benchmark Results (2)

Iterative method with **different ansatz** is used

$$\alpha_S(\rho_V) = a_S + \sum_k b_S^{k\text{-th}} \left(\frac{\rho_V}{\rho_{\text{sat}}} - 1 \right) + \sum_k c_S^{k\text{-th}} \left(\frac{\rho_V}{\rho_{\text{sat}}} - 1 \right)^2$$

$$\alpha_V(\rho_V) = a_V + \sum_k b_V^{k\text{-th}} \left(\frac{\rho_V}{\rho_{\text{sat}}} - 1 \right) + \sum_k c_V^{k\text{-th}} \left(\frac{\rho_V}{\rho_{\text{sat}}} - 1 \right)^2$$



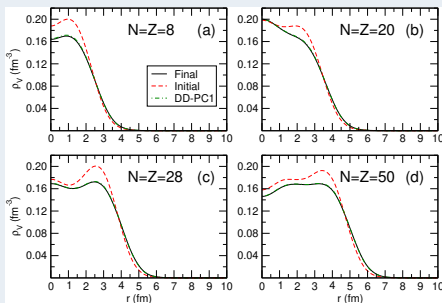
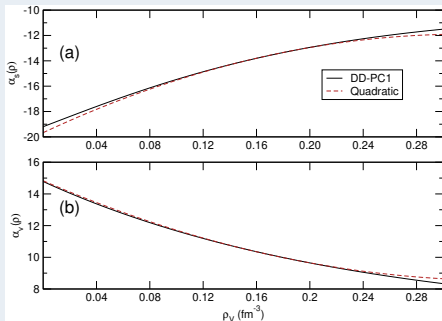
Accorto, Naito, Liang, Nikšić, and Vretenar. *Phys. Rev. C* **103**, 044304 (2021)

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Even with different ansatz, well reproduced!!

Accorto, Naito, Liang, Nikšić, and Vretenar. *Phys. Rev. C* **103**, 044304 (2021)

Conclusion and Perspectives

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- Inverse KS method and IKS-DFPT are extended to relativistic systems
- IKS-DFPT works well both for atomic systems and for nuclear systems in benchmark calculations

Perspectives

- How can we constrain spin-orbit potential?
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Thank you for attention!!