

# Precision measurements in the SM: why and what?



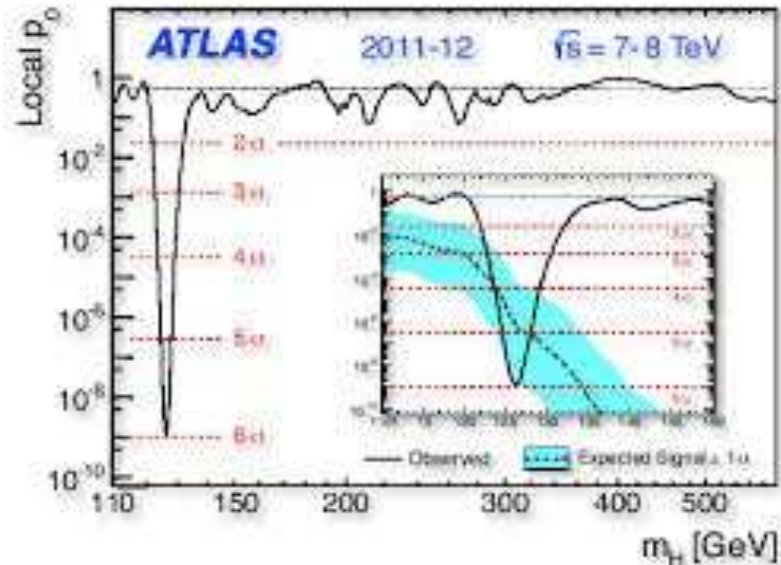
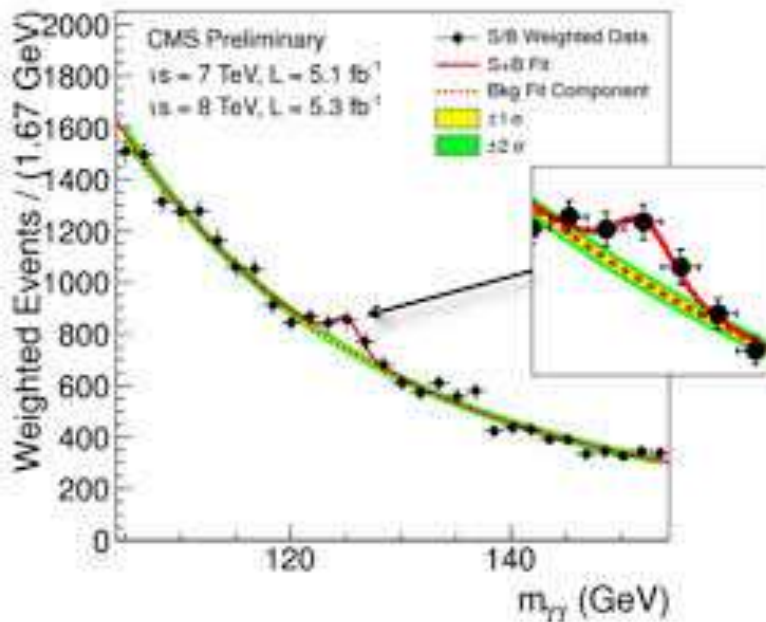
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**(CNRS & LAPTh)**



- 1. Standardissimo?!**
- 2. Trouble with minimal beyond the SM?**
- 3. Quo vadis? Further probe of Higgs**
- 4. Where and what else?**
- 5. Conclusion**

# 1. Standardissimo?!

The Higgs discovery in July 2012: a triumph for high-energy physics.



A very non-trivial check of the SM: test at the quantum/permille level:

– constraints from data:  $M_H = 92_{-26}^{+34} \text{ GeV} \lesssim 160 \text{ GeV}$  at 95% CL

– experimentally found to be:  $M_H \simeq 125 \pm 0.2 \text{ GeV}$  (ie within  $1\sigma$ ..)

In addition, it looks as it has the properties of the SM Higgs state:

The triumph of the SM model of particle physics or Standardissimo?!

# 1. Standardissimo?!

We have a theory for the strong+electroweak forces, the SM, that is:

- a relativistic quantum field theory based on a gauge symmetry,
- renormalisable as proved by 't Hooft and Veltman for SEWSB,
- unitary as we have now a Higgs and its mass is rather small,
- perturbative up to the Planck scale as again the Higgs is light,
- leads to a (meta)stable electroweak vacuum up to high scales,
- compatible with (almost) all precision data available to date...

Is the SM the “theory of everything” and should we be satisfied with it?

Maybe low energy manifestation of more fundamental theory that solve

- “Esthetical” problems with e.g. multiple and arbitrary parameters; gauge coupling unification:  $3 \neq g_i$  which do not meet a high scale.
- “Experimental” problems as it does not explain all seen phenomena:  $\nu$  masses/mixing, dark matter, baryon asymmetry in the universe ....  
(note:  $SO(10)$  at intermediate  $Q = 10^{11} \text{ GeV}$  and axions cure these pbs)
- “Theory” (or consistency) problem: the hierarchy/naturalness pbs.  
 $\Delta M_H^2 \propto \Lambda^2 \approx (10^{18} \text{ GeV})^2$ :  $M_H$  not stable against high scales.

All these may indicate that there is beyond the Standard Model.

# 1. Standardissimo?!

There were 3 avenues for solving the hierarchy/naturalness problem:

## I. Compositeness/substructure:

All particles are composite: Technicolor

⇒ **H bound state of two fermions**

(no more spin-0 fundamental state).

## II. Extra space-time dimensions

where at least  $s=2$  gravitons propagate.

⇒ **effective gravity scale  $\Lambda \approx 1\text{TeV}$ .**

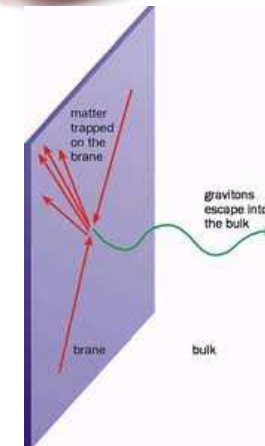
EWSB mechanism needed: H or not H!

## III. Supersymmetry: doubling the world.

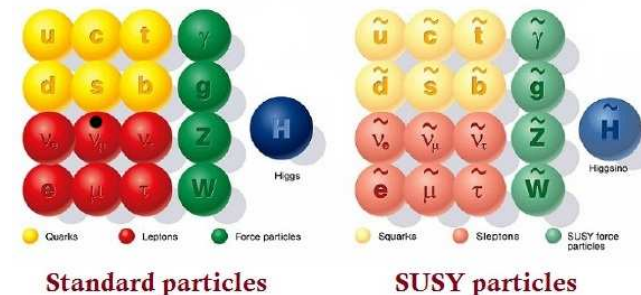
- links  $s=1/2$  fermions to  $s=1$  bosons,
- links internal/space-time symmetries,
- if made local, provides link to gravity,
- natural  $\mu^2 < 0$ : radiative EWSB,

⇒ **sparticle loops cancel  $\Lambda^2$  behavior**

**All predicted something at Terascale!**



### SUPERSYMMETRY



# 1. Standardissimo?

The problem is that:

A) we observe a Higgs with a mass of 125 GeV and no other Higgs:

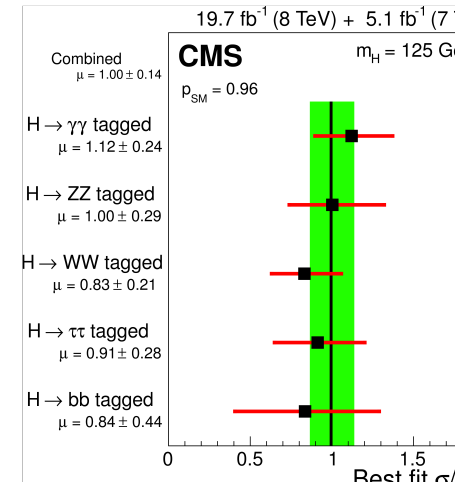
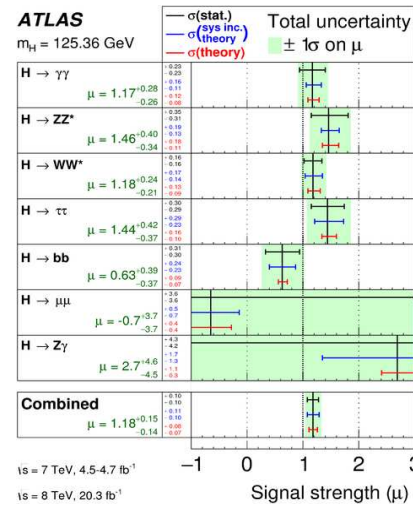
$\sigma \times \text{BR}$  rates compatible with those expected in the SM

Fit of all LHC Higgs data  $\Rightarrow$  agreement at 15–30% level

Results from the LHC run I campaign already give us:

$$\mu_{\text{tot}}^{\text{ATLAS}} = 1.18 \pm 0.15$$

$$\mu_{\text{tot}}^{\text{CMS}} = 1.00 \pm 0.14$$



we do not observe any new particle beyond those of SM with Higgs:

profound implications for most discussed BSM scenarios; they are in:

- “Mortuary”: Higgsless, 4th generation, fermio or gauge-phobic..
- “Hospital”: Technicolor, composite models (but some loopholes) ....
- “Trouble” and strongly constrained: extra-dimensions, SUSY, ...

As an example, let us see what it implies for SUSY and the MSSM.

## 2. Trouble with minimal BSM? The MSSM

In the MSSM we need two doublets of complex scalar fields:  $H_1, H_2$  to generate up/down-type fermion masses and no chiral anomalies. **after EWSB, three dof for  $W_L^\pm, Z_L \Rightarrow 5$  physical states:  $h, H, A, H^\pm$ .** Only 2 free parameters at tree-level to describe the system  $\tan\beta, M_A$ :

$$M_{h,H}^2 = \frac{1}{2} \left\{ M_A^2 + M_Z^2 \mp [(M_A^2 + M_Z^2)^2 - 4M_A^2 M_Z^2 \cos^2 2\beta]^{1/2} \right\}$$

$$M_{H^\pm}^2 = M_A^2 + M_W^2$$

$$\tan 2\alpha = \frac{-(M_A^2 + M_Z^2) \sin 2\beta}{(M_Z^2 - M_A^2) \cos 2\beta} = \tan 2\beta \frac{M_A^2 + M_Z^2}{M_A^2 - M_Z^2} \quad \left(-\frac{\pi}{2} \leq \alpha \leq 0\right)$$

$M_h \lesssim M_Z |\cos 2\beta| + RC \lesssim 130 \text{ GeV}$ ,  $M_H \approx M_A \approx M_{H^\pm} \lesssim M_{\text{EWSB}}$ .

- Couplings of  $h, H$  to  $VV$  are suppressed; no  $AVV$  couplings (CP).
- For  $\tan\beta \gg 1$ : couplings to  $b$  ( $t$ ) quarks enhanced (suppressed).

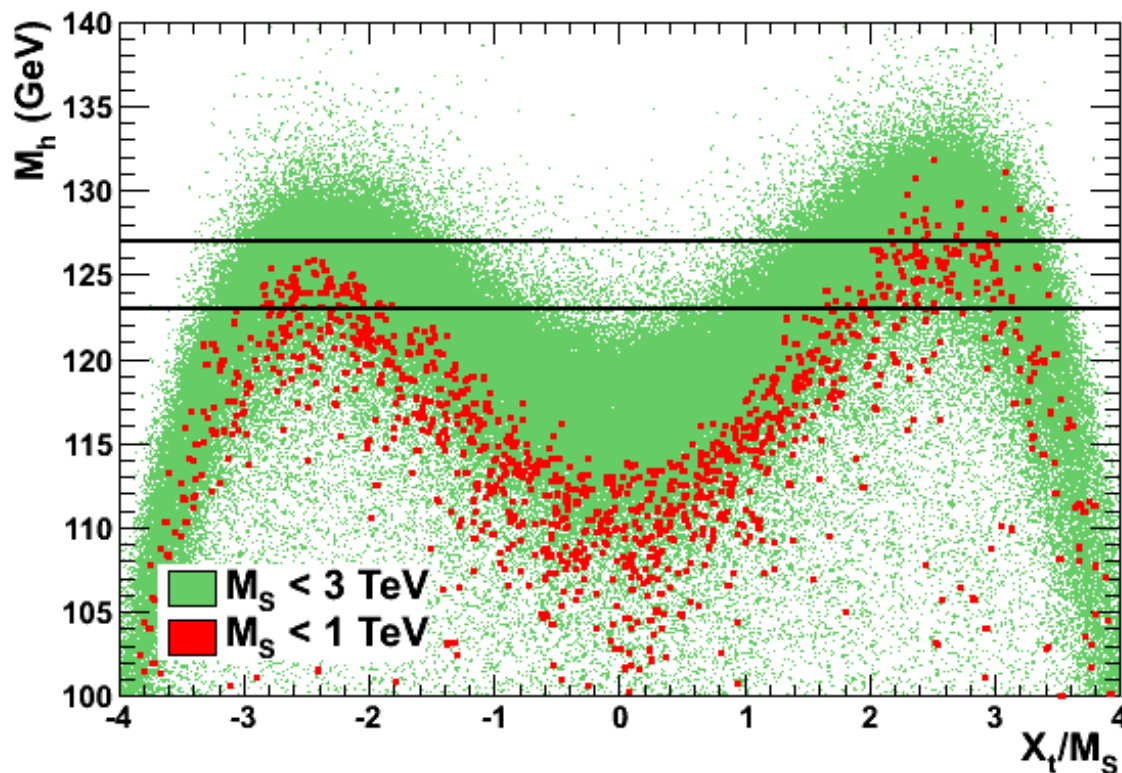
$\Phi$	$g_{\Phi\bar{u}u}$	$g_{\Phi\bar{d}d}$	$g_{\Phi VV}$
$h$	$\frac{\cos\alpha}{\sin\beta} \rightarrow 1$	$\frac{\sin\alpha}{\cos\beta} \rightarrow 1$	$\sin(\beta - \alpha) \rightarrow 1$
$H$	$\frac{\sin\alpha}{\sin\beta} \rightarrow 1/\tan\beta$	$\frac{\cos\alpha}{\cos\beta} \rightarrow \tan\beta$	$\cos(\beta - \alpha) \rightarrow 0$
$A$	$1/\tan\beta$	$\tan\beta$	$0$

In decoupling limit: MSSM Higgs sector reduces to SM with a light  $h$ .

# 1. Trouble with minimal BSM? The MSSM

There is first direct implication from the measurement  $M_h = 125\text{GeV}$ .

$$M_h^2 \xrightarrow{M_A \gg M_Z} M_Z^2 \cos^2 2\beta + \frac{3\bar{m}_t^4}{2\pi^2 v^2 \sin^2 \beta} \left[ \log \frac{M_S^2}{\bar{m}_t^2} + \frac{X_t^2}{M_S^2} \left( 1 - \frac{X_t^2}{12M_S^2} \right) \right] = (125)^2$$



Arbey, Battaglia, AD, Mahmoudi, Quevillon (2012)

$M_{\text{SUSY}} \gtrsim 1\text{TeV}$  in general MSSM and higher in constrained models.

# 2. Trouble with minimal BSM? The MSSM

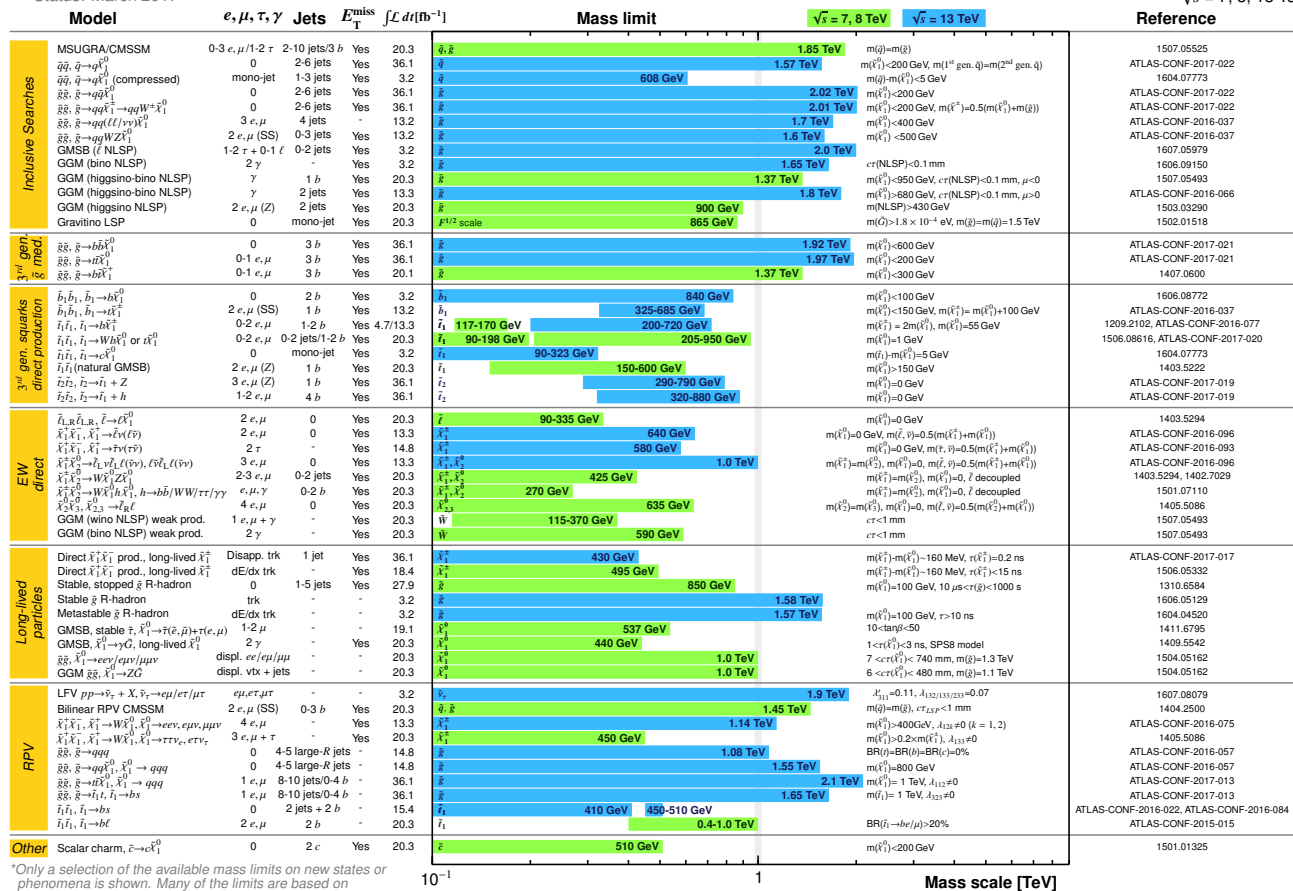
This is backed up by direct searches of SUSY particles at the LHC: the SUSY scale  $M_{\text{SUSY}} \gtrsim \mathcal{O}(1 \text{ TeV})$  in most experimental searches..

ATLAS SUSY Searches\* - 95% CL Lower Limits

Status: March 2017

ATLAS Preliminary

$\sqrt{s} = 7, 8, 13 \text{ TeV}$



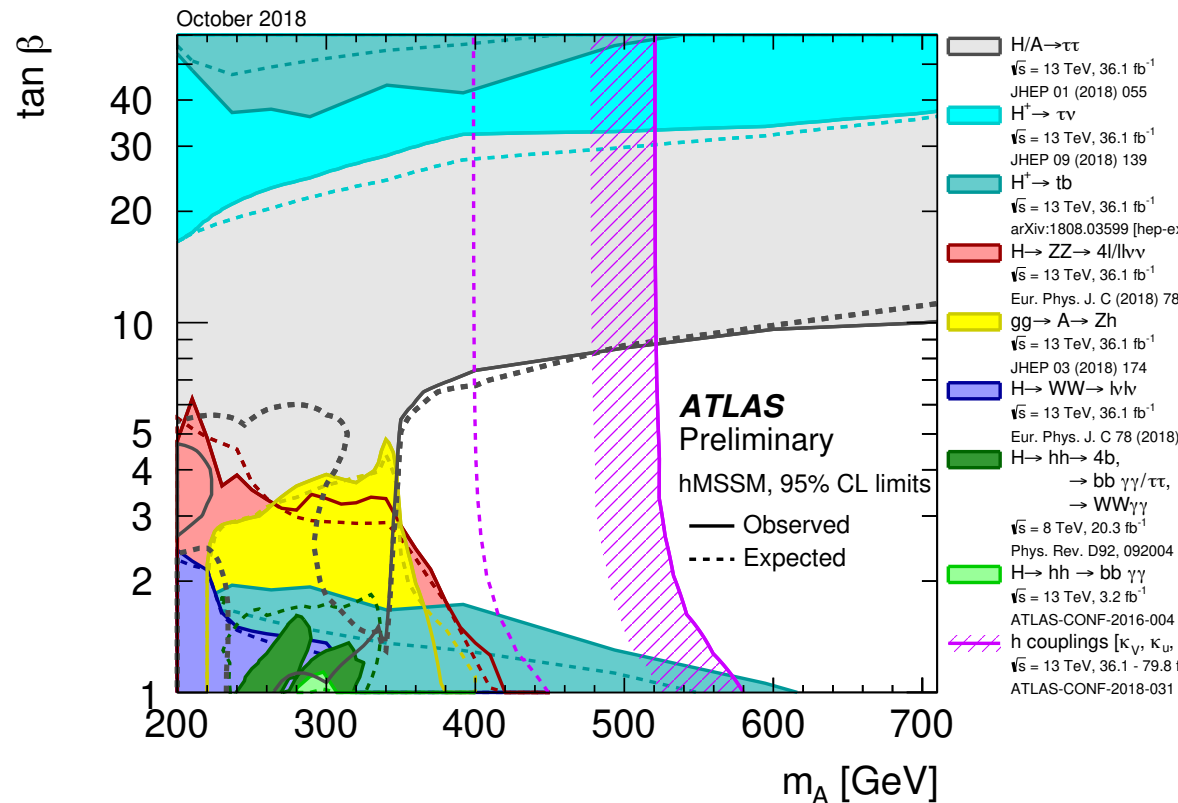
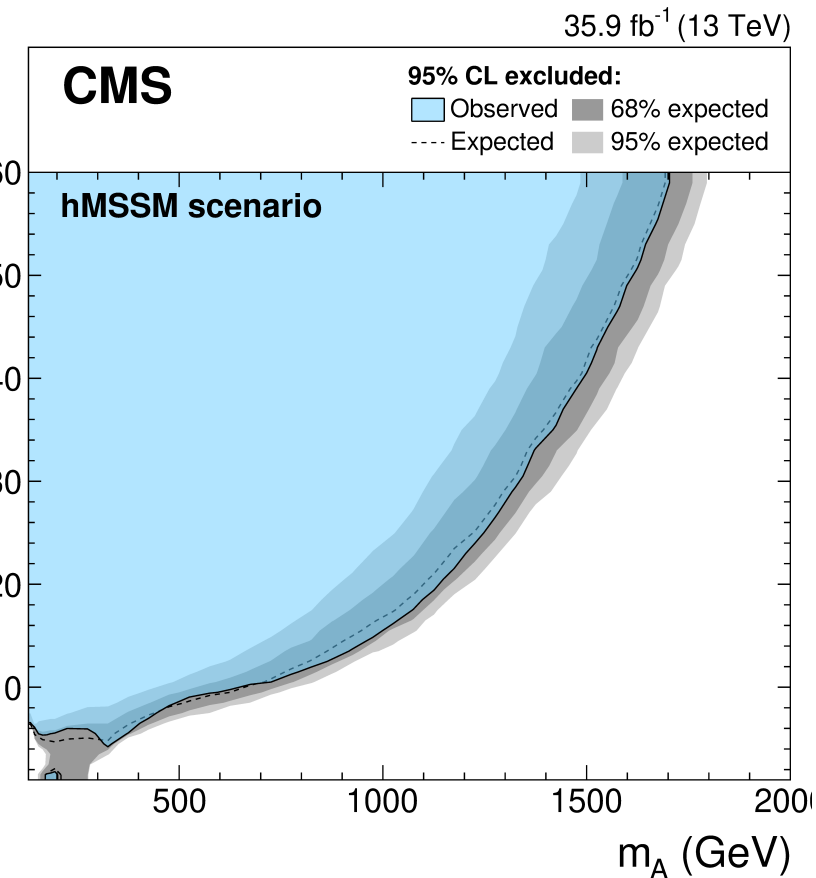
\*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

⇒ ATLAS/CMS depressing exclusion tables..



## 2. Trouble with minimal BSM? The MSSM

**Backed up by searches and measurements in the Higgs sector at LHC:**  
**fits of the  $h$  couplings  $\Rightarrow$  constraints on the MSSM  $[M_A, \tan\beta]$  plane:**  
**hMSSM:  $g_{h\bar{t}t} = \cos\alpha / \sin\beta$ ,  $g_{h\bar{b}b} = \cos\alpha / \sin\beta$ ,  $g_{hVV} = \sin(\beta - \alpha)$**



# 3. Quo vadis? Further probe of Higgs

So is Particle Physics “closed” and we should all go home? No!

Fully probe the TeV scale that is relevant for the hierarchy problem

⇒ continue searches for new particles in all possible channels.

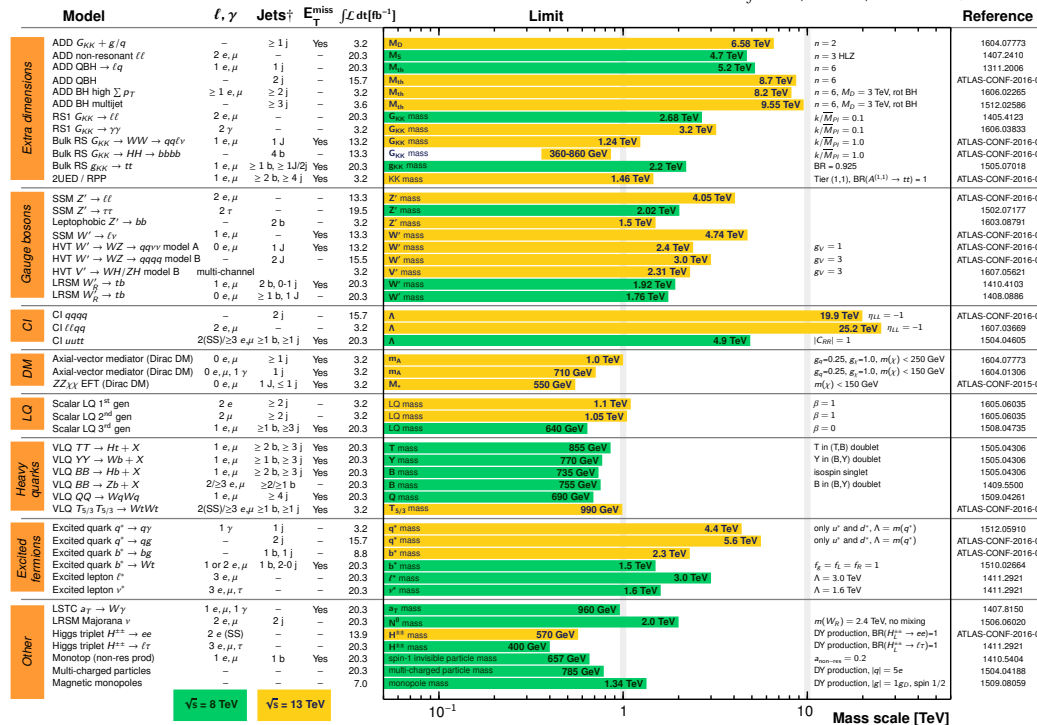
ATLAS Exotics Searches\* - 95% CL Exclusion

Status: August 2016

ATLAS Preliminary

$$\int \mathcal{L} dt = (3.2 - 20.3) \text{ fb}^{-1}$$

$$\sqrt{s} = 8, 13 \text{ TeV}$$



\*Only a selection of the available mass limits on new states or phenomena is shown. Lower bounds are specified only when explicitly not excluded.

$^\dagger$ Small-radius (large-radius) jets are denoted by the letter J (j).

Should be continued, extended, refined:

new states are simply around the corner and can be found tomorrow!

### 3. Quo vadis? Further probe of Higgs

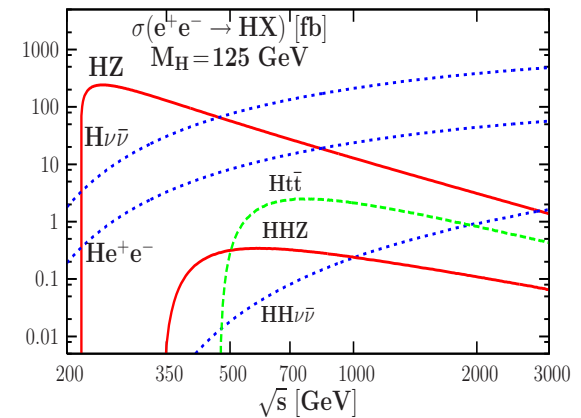
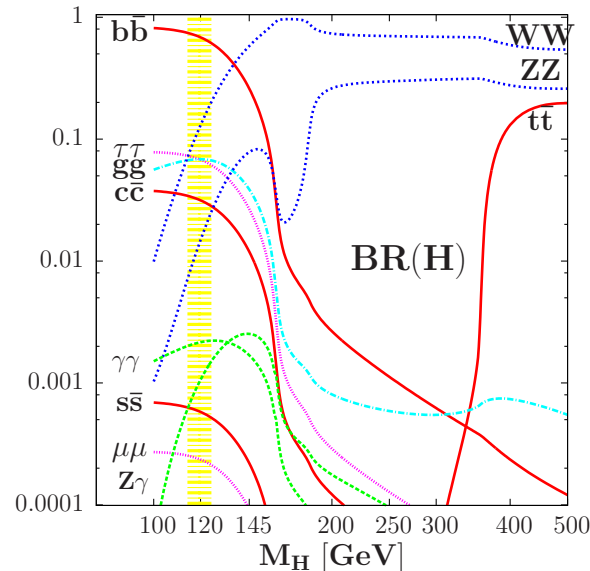
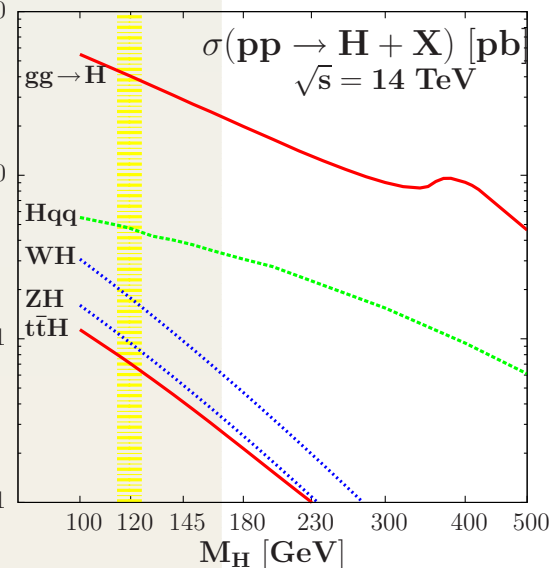
The next question is then: “is Particle Physics closed”? Answer is no!

2) Need to check that H is alone responsible of EWSB (SM-like?)

⇒ measure its fundamental properties in the most precise way:

- its mass and total decay width (invisible width from dark matter?),
- its spin–parity quantum numbers (CP violation for baryogenesis?),
- its couplings to fermions and gauge bosons and check if they are only proportional to particle masses (no new physics contributions?),
- its self-couplings to reconstruct  $V_S$  potential that makes EWSB.

Possible for  $M_H \approx 125$  GeV as all production/decay channels useful.



### 3. Quo vadis? Further probe of Higgs

A check of spin–parity quantum numbers and search for CP violation

**Spin:** clear situation (no suspense) as the new state decays into  $\gamma\gamma$   
 $\Rightarrow$  not  $s=1$  from Landau–Yang and  $s=2$  (KK graviton?) unlikely..

**CP numbers:** CP-even, CP-odd, or mixture?

(more important issue: CPV in Higgs sector.)

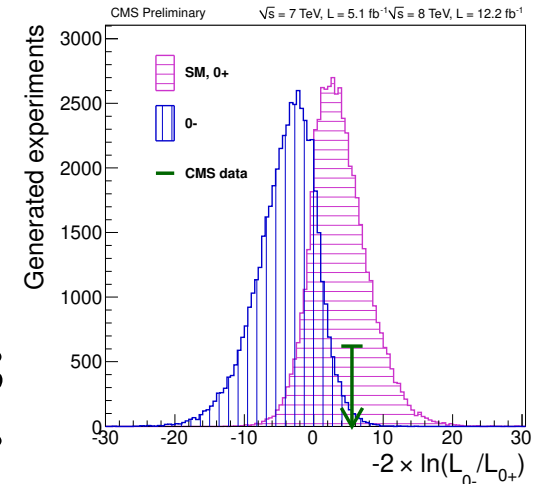
ATLAS and CMS MELA analyses for pure CP

$\Rightarrow$  pure CP-even favored at  $\gtrsim 3\sigma$  level.

But problems with this (too simple) picture:

pure CP-odd does not couple to  $VV$  @ tree-level;

in  $H \rightarrow ZZ^*$  only CP-even part is projected out.



• **Direct probe:** via production/decays in extensions like C2HDM:  
Ex: Undoubtable signs of CP-violation in Higgs decays at HL-LHC  
combined searches of  $h_i \rightarrow h_j Z$  and  $h_i \rightarrow ZZ$  with  $i, j = 1, 2, 3$ .

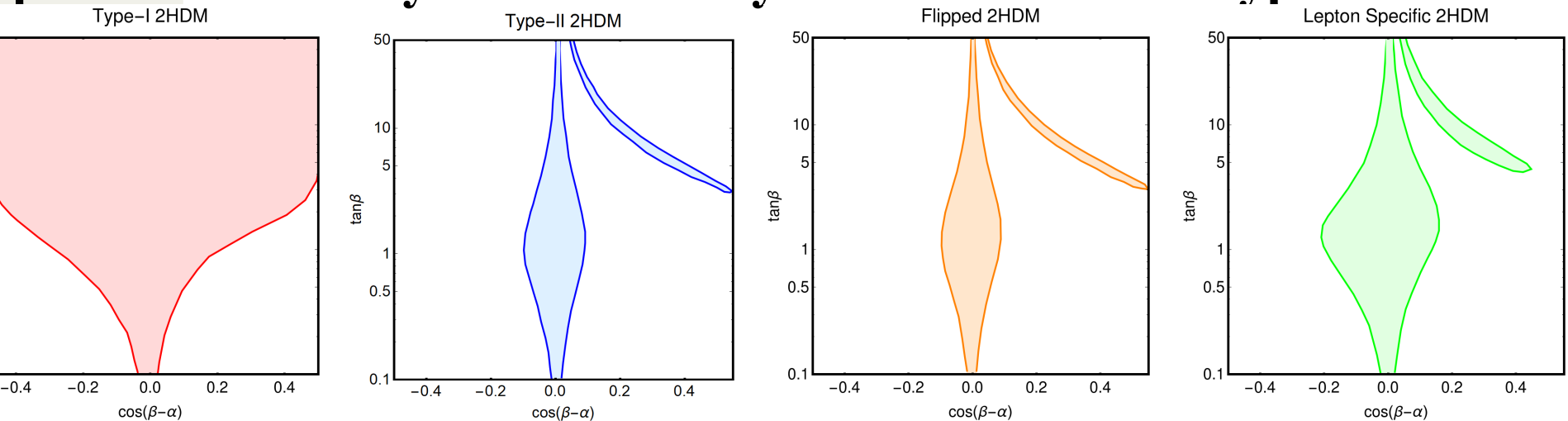
• **Indirect probe:**  $g_{Hff}$  more democratic  $\Rightarrow$  fermionic decays.  
ex: spin-correlations in  $q\bar{q} \rightarrow HZ \rightarrow b\bar{b}l\bar{l}$ ,  $q\bar{q}/gg \rightarrow Ht\bar{t} \rightarrow b\bar{b}t\bar{t}$ .

**Need to be lucky or is very challenging even at the HL-LHC...**

### 3. Quo vadis? Further probe of Higgs

Perform a much more precise measurement of the Higgs couplings  
⇒ would allow a better sensitivity to new physics virtual effects.

- In standard production+decay modes as  $gg \rightarrow H \rightarrow ZZ, WW, \gamma\gamma$  present sensitivity is low in many cases as 2HDM of type I to IV:



- In very rare decays that allow additional/unknown information:
  - $H \rightarrow \mu^+ \mu^-$  to probe second generation fermion couplings,
  - $H \rightarrow \Upsilon \gamma$  to probe the sign of some fermionic couplings (here b's),
  - $H \rightarrow Z \gamma$  with information that is complementary to  $H \rightarrow \gamma\gamma$ .

But will this be sufficient to probe BSM physics? (maybe ratios then?)

### 3. Quo vadis? Further probe of Higgs

example of search for New Physics via high precision measurements:  
the Higgs couplings in cleanest channels  $H \rightarrow \gamma\gamma$ ,  $H \rightarrow 4\ell^\pm$

channel	atlas/run1				cms/run1			
$\mu_{\gamma\gamma}$	1.17	+0.23	+0.16	(+0.12)	1.14	+0.21	+0.16	(+0.09)
		-0.23	-0.11	(-0.08)		-0.21	-0.10	(-0.05)
$\mu_{ZZ}$	1.46	+0.35	+0.19	(+0.18)	0.93	+0.26	+0.13	
		-0.31	-0.13	(-0.11)		-0.23	-0.09	

Is this enough to probe effects of new physics or BSM?

not in the case of weakly interacting theories like 2HDM, SUSY, etc...

expect effects at  $\approx \frac{C_{\text{new}} \alpha_W}{\pi} \approx \frac{M_h^2}{M_{\text{new}}^2} \approx 1\%$ ;

Is 1% accuracy possible at HL-LHC ( $3\text{ab}^{-1}$ )?

- statistical error:  $20\% / \sqrt{3 \times 100} \lesssim 1-2\%$
- systematical error: can be made  $\lesssim 1\%$  ?
- theoretical uncertainty is then crucial issue.

Less ambiguities in ratios of rates! Example:

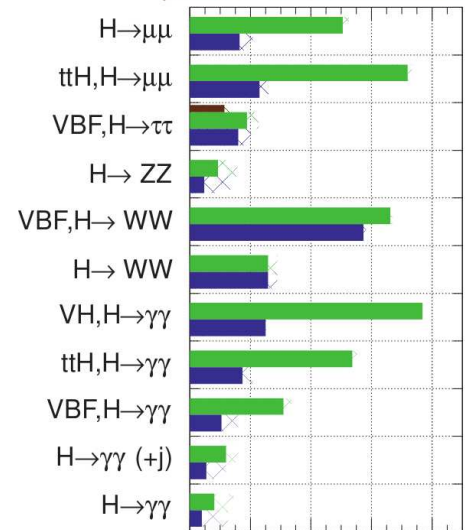
$$D_{\gamma\gamma} = \frac{\sigma(\text{pp} \rightarrow H \rightarrow \gamma\gamma)}{\sigma(\text{pp} \rightarrow H \rightarrow \text{VV})} = \frac{\Gamma(H \rightarrow \gamma\gamma)}{\Gamma(H \rightarrow \text{VV})} = d_{\gamma\gamma} c_\gamma^2 / c_V^2$$

Will  $D_{\gamma\gamma}$  be the g-2 of the LHC? Yes, at 1%!

ATLAS Simulation

$\sqrt{s} = 14 \text{ TeV}$ :  $\int \text{Ldt} = 300 \text{ fb}^{-1}$ ;  $\int \text{Ldt} = 3000 \text{ fb}^{-1}$

$\int \text{Ldt} = 300 \text{ fb}^{-1}$  extrapolated from 7+8 TeV



### 3. Quo vadis? Further probe of Higgs

- **Total width:**  $\Gamma_H = 4\text{MeV}$ , too small to be resolved experimentally.
  - very loose bound from interference  $gg \rightarrow ZZ$  (100% level).
  - difficult to access it indirectly (via production rates) very precisely...

• **Invisible width:** more accessible

#### Direct measurement of $H \rightarrow \text{inv}$

- $q\bar{q} \rightarrow HZ$  with  $Z \rightarrow ll$ ,  $H \rightarrow \text{inv}$
- similar  $E_T$  search in VBF mode
- and also in  $gg \rightarrow \text{Higgs} + \text{jet}$

Combined  $HZ + \text{VBF} + ggH$  now

$\text{BR}_{\text{inv}} \lesssim 20\% @ 95\% \text{ CL}$

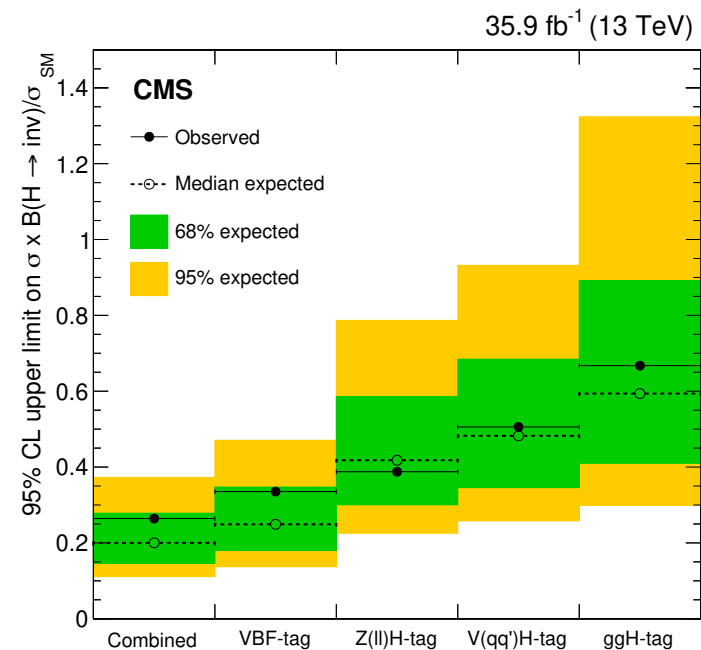
assuming a SM Higgs state

**few % @ HL-LHC possible?**

#### Indirect measurement of $H \rightarrow \text{inv}$

via Higgs BRs measurement: accuracy of  $O(\text{few}\%)$  at HL-LHC

but with TH assumptions: no other decays, SM-like Higgs, etc...



# 3. Quo vadis? Further probe of Higgs

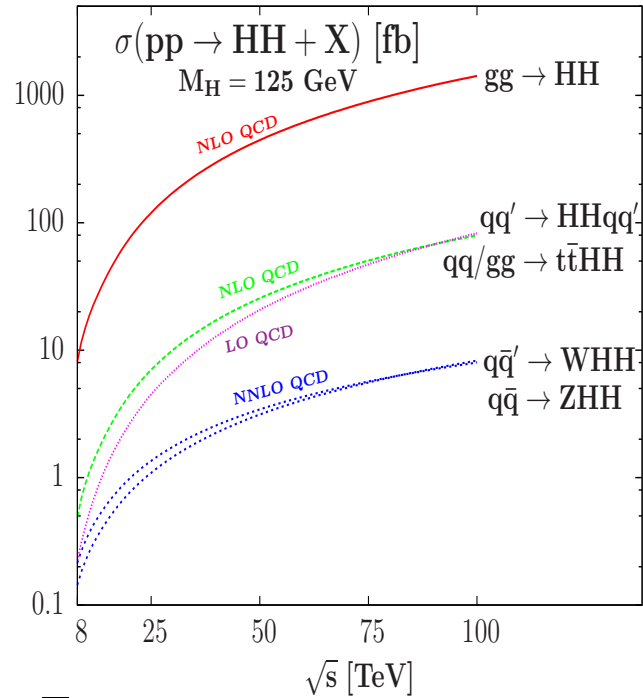
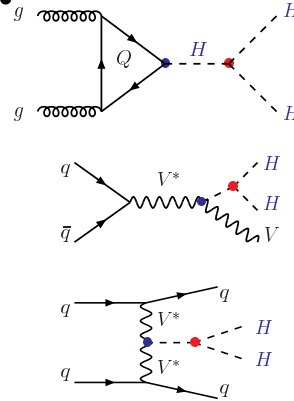
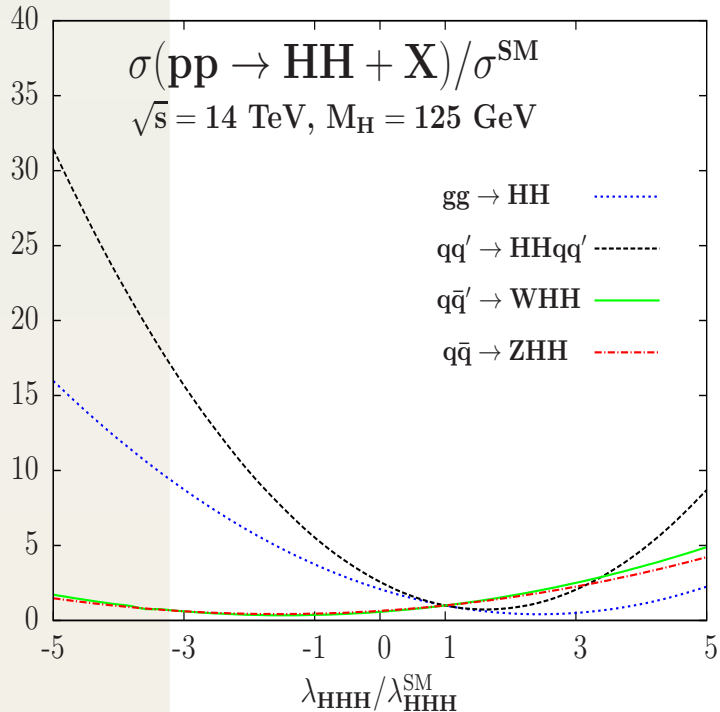
Important challenge: measure Higgs self-couplings and access to  $V_H$ .

$g_{H^3}$  from  $pp \rightarrow HH + X \Rightarrow$

$g_{H^4}$  from  $pp \rightarrow 3H + X$ , hopeless.

various processes for HH prod:

only  $gg \rightarrow HHX$  relevant...



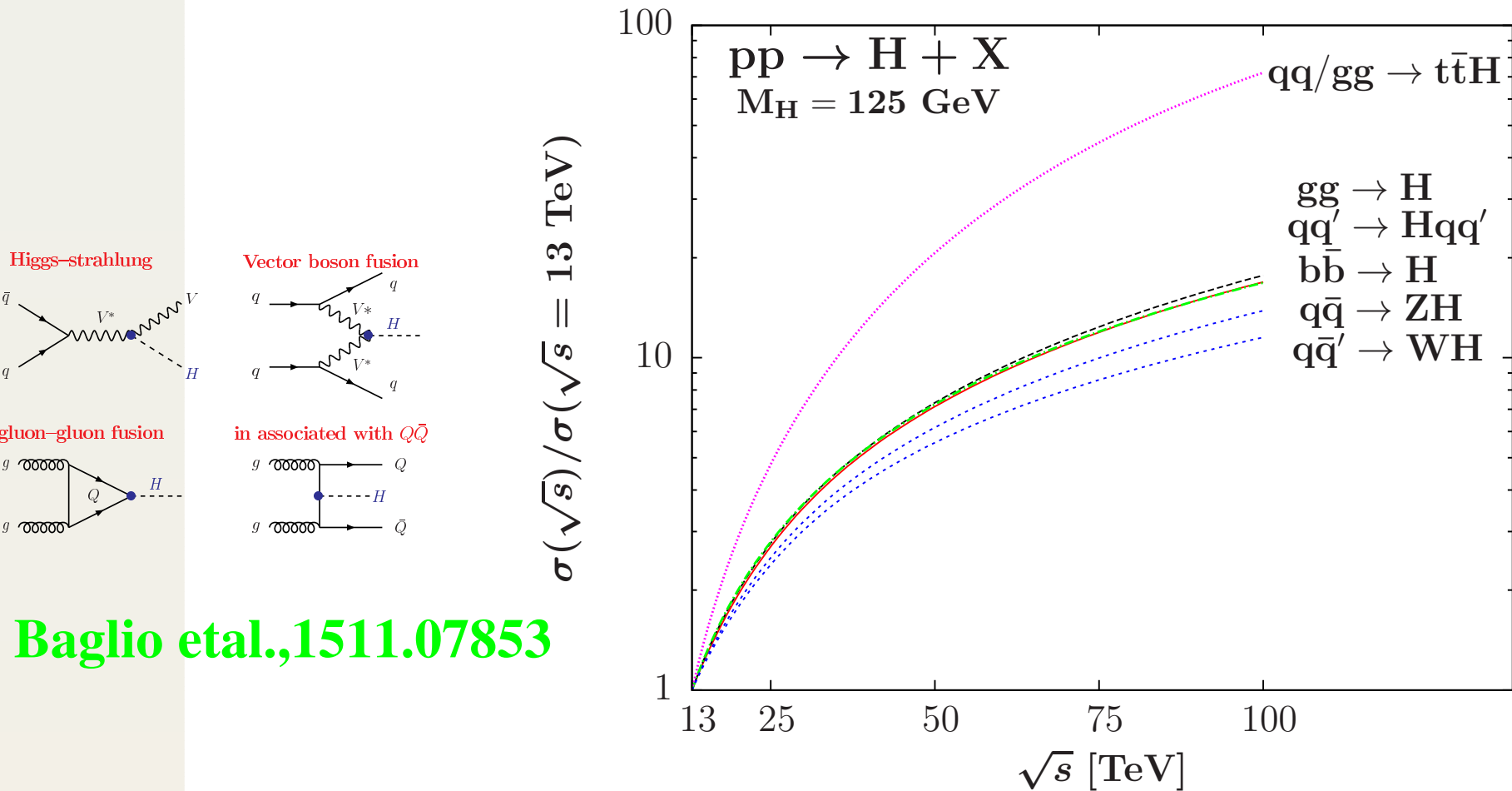
- $H \rightarrow b\bar{b}$  decay alone not clean
  - $H \rightarrow \gamma\gamma$  decay very rare,
  - $H \rightarrow \tau\tau$  would be possible?
  - $H \rightarrow WW$  not useful?
- $bb\tau\tau, bb\gamma\gamma$  viable? Maybe...  
 but needs very large luminosity.

Baglio et al., arXiv:1212.5581



## 4. Where and what else?

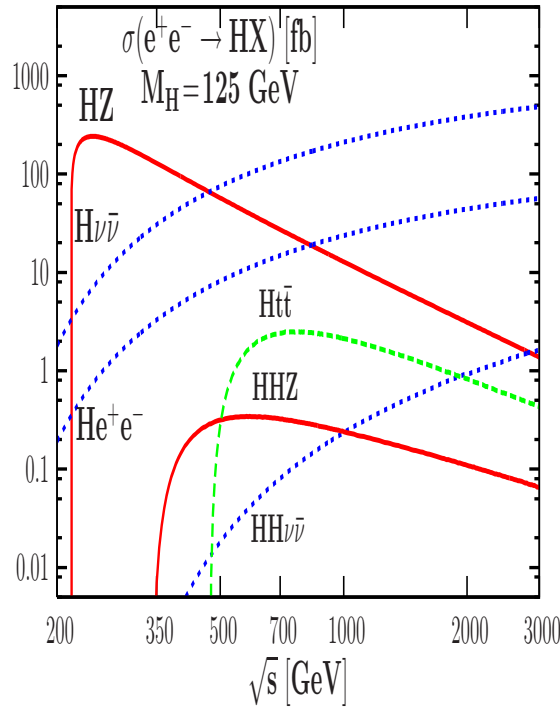
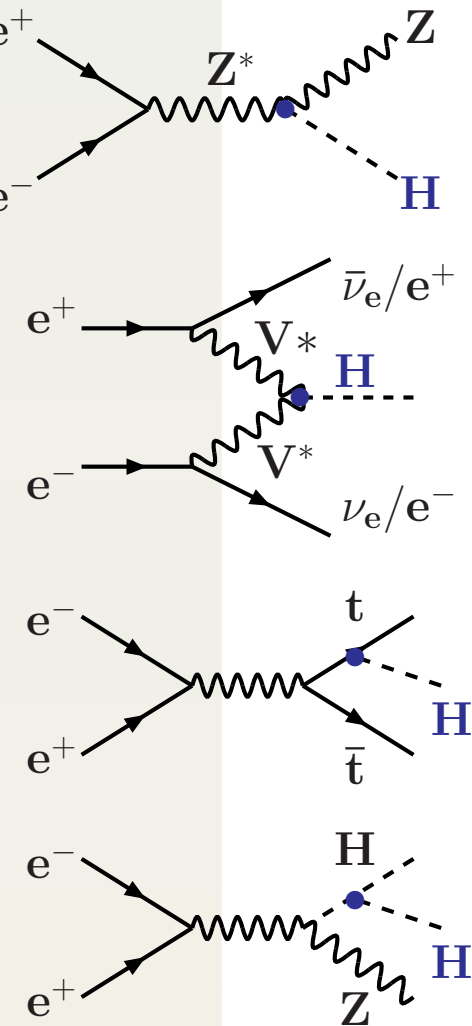
A large increase in sensitivity at high energy machines is possible as production cross section (especially in some cases) are larger:



Baglio et al., 1511.07853

Very interesting to move to 100 TeV (not only for this of course)!

## 4. Where and what else?



Very precise measurements  
 mostly at  $\sqrt{s} \lesssim 500$  GeV  
 and mainly in  $e^+e^- \rightarrow ZH$   
 (with  $\sigma \propto 1/s$ ) and  $ZHH, ttH$

$g_{HWW}$	$\pm 0.012$
$g_{HZZ}$	$\pm 0.012$
$g_{Hbb}$	$\pm 0.022$
$g_{Hcc}$	$\pm 0.037$
$g_{H\tau\tau}$	$\pm 0.033$
$g_{Htt}$	$\pm 0.030$
$\lambda_{HHH}$	$\pm 0.22$
$M_H$	$\pm 0.0004$
$\Gamma_H$	$\pm 0.061$
CP	$\pm 0.038$

$\Rightarrow$  best option for  $\approx 125$  GeV Higgs

## 4. Where and what else?

The SM is not only the Higgs, also other domains need to be tested.

A first one is the physics of the Z and W bosons!

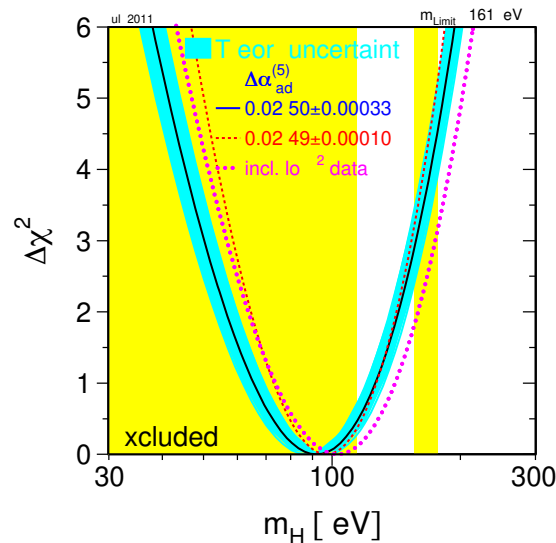
remember: pre-LHC constraints on the Higgs boson mass in the SM

- From various measurements of the Z observables at LEP1, SLC, ...

- From the very precise measurement of  $M_W$  at LEP2, Tevatron, ...

constraints on the Higgs mass  $M_H$  once we know precisely  $m_t, \alpha_s, \text{etc...}$

	Measurement	Fit	$\frac{O^{\text{meas}} - O^{\text{fit}}}{\sigma^{\text{meas}}}$
$\Delta\alpha_{\text{ad}}^{(5)}(m_Z)$	$0.0250 \pm 0.00033$	0.0259	
$m_Z$ [eV]	$91.185 \pm 0.0021$	91.184	
$\Gamma_Z$ [eV]	$2.4952 \pm 0.0023$	2.4959	
$\sigma_{\text{ad}}^0$ [nb]	$41.540 \pm 0.03$	41.48	
$\Gamma_{\text{fb}}^e$	$20.6 \pm 0.025$	20.42	
$\Gamma_{\text{fb}}^{\mu}$	$0.0114 \pm 0.00095$	0.01646	
$\Gamma_{\text{had}}^e$	$0.1465 \pm 0.0032$	0.1482	
$\Gamma_b$	$0.21629 \pm 0.00066$	0.2159	
$\Gamma_c$	$0.121 \pm 0.0030$	0.122	
$\Gamma_{\text{br}}^e$	$0.0992 \pm 0.0016$	0.1039	
$\Gamma_{\text{br}}^{\mu}$	$0.0 \pm 0.0035$	0.043	
$\Gamma_b$	$0.923 \pm 0.020$	0.935	
$\Gamma_c$	$0.60 \pm 0.02$	0.668	
$\sin^2\theta_{\text{eff}}^e(\text{SL})$	$0.1513 \pm 0.0021$	0.1482	
$\sin^2\theta_{\text{eff}}^e(\text{fb})$	$0.2324 \pm 0.0012$	0.2314	
$m_t$ [eV]	$80.399 \pm 0.023$	80.38	
$\Gamma_t$ [eV]	$2.085 \pm 0.042$	2.092	
$m_t$ [eV]	$13.20 \pm 0.90$	13.2	



Global fit:  $M_H = 92_{-26}^{+34} \text{ GeV} \lesssim 160(126) \text{ GeV}$  at 95% (68%) CL!

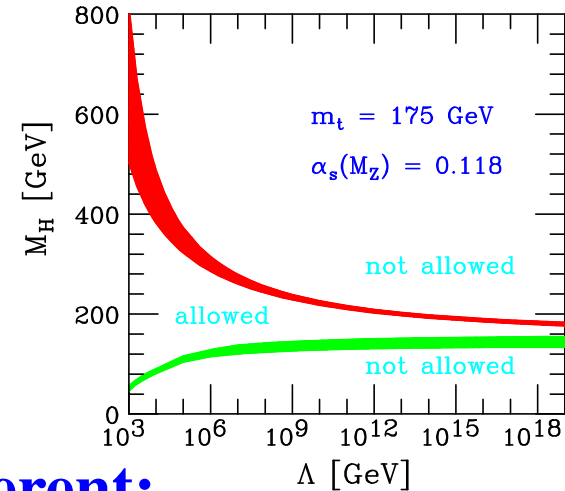
Now that the Higgs mass is known, we can really look for New Physics

The GigaZ for  $\sin^2\theta_W$  and the MegaW option for  $M_W$  are a must!

## 4. Where and what else?

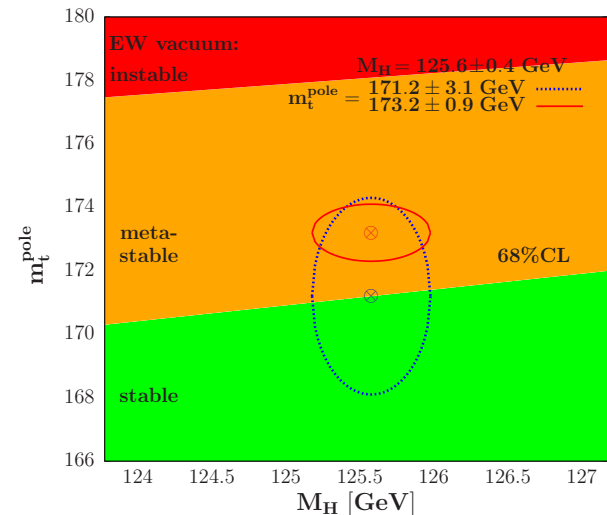
Another domain is the physics of the top quark! Just one example:  
remember again, pre-LHC constraints on the Higgs mass in the SM:

- included theory constraints on  $\lambda(Q^2)$  :
- $\lambda \gg 1$ : triviality (non-perturbativity)
- $\lambda \ll 1$ : (un)stability of EW vacuum
- strong constraint depending on cut-off
- $\Lambda \sim 1 \text{ TeV} : 70 \lesssim M_H \lesssim 700 \text{ GeV}$
- $\Lambda \sim M_{\text{Planck}} : 120 \lesssim M_H \lesssim 180 \text{ GeV}$



Now that we know  $M_H$ , the problem is different:

- Depending on  $m_t$  (meta)stability!**
- with usual top mass (pole,  $\overline{\text{MS}}$ , MC,?)
- vacuum metastable for  $\Lambda = M_P$
- but with  $m_t$  from top cross sections:
- vacuum stable up to  $\Lambda = M_P$
- There are also effects from  $\alpha_s, \dots$

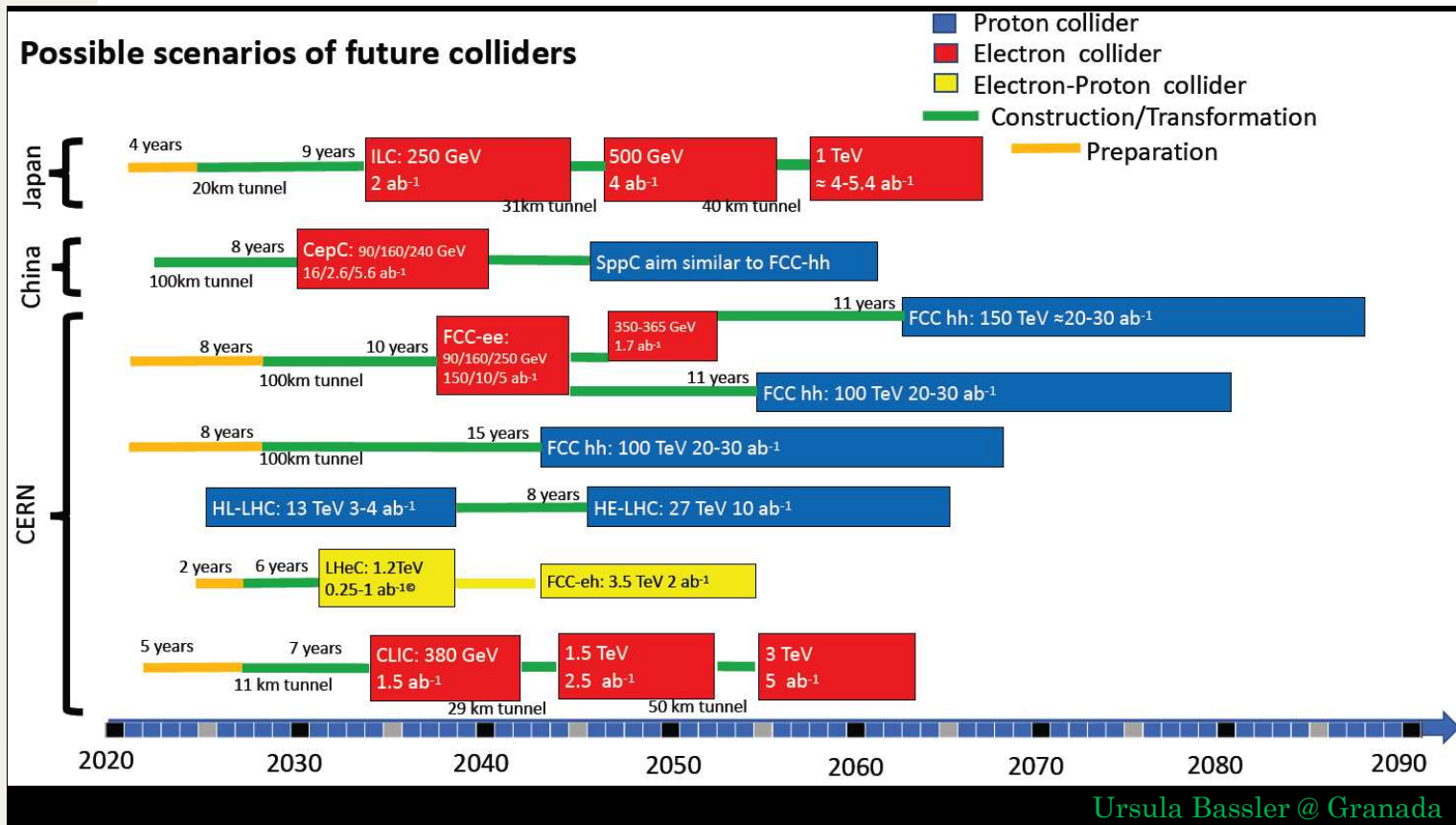


**We need a better and more "rigorous" determination of  $m_t, \alpha_s$**



# 5. Conclusion

- We need to continue to search for New Physics and falsify the SM:
- directly via new (heavy or light) particle searches with more data.
  - indirectly via high precision measurements in H/W/Z/top sectors,



**So let's move forward: it is still action time!**