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## Marco Rocchini

## INFN - Istituto Nazionale di Fisica Nucleare

 FIRENZE DIVISIONAxtal Shape Asymmetiry and Configuration Coexistence in Neutron-Rich 74Zn: Suggestion for a Northern Extension of the $N=40$ Island of Inversion


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## Islands of Inversion

Iols, Shape Coexistence and Triaxiality

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Our Experiment on ${ }^{74} \mathrm{Zn}$
$\gamma-\gamma$ Angular Correlations

Experimental Results

LSSM
Calculations

Shapes and Occupations

Conclusions

- Iols: Regions of the nuclide chart in which the energy gained through correlations (e.g., quadrupole) can offset the spherical meanfield gaps, leading to the appearance of unexpected deformed ground states
- Their study permits investigating correlation energies and phenomena such as deformation and shape coexistence
- 4 Iols identified: $N=8,20,28,40$
- 1 new lol theorised at $N=50$ F. Nowacki et al., PRL 117, 272501 (2016)


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$$
g \text { factor of the exotic } N=21 \text { isotope }{ }^{34} \mathrm{Al} \text { : probing the } N=20
$$ and $N=28$ shell gaps at the border of the "island of inversion"

P. Himpe ${ }^{\text {a }}$, G. Neyens ${ }^{\text {a,* }, ~ D . L . ~ B a l a b a n s k i ~}{ }^{\text {b }}$, G. Bélier ${ }^{\text {c }}$, J.M. Daugas ${ }^{\text {c }}$, F. de Oliveira Santos ${ }^{\text {d }}$, M. De Rydt ${ }^{a}$, K.T. Flanagan ${ }^{a}$, I. Matea ${ }^{e}$, P. Morel ${ }^{c}$, Yu.E. Penionzhkevich ${ }^{f}$, L. Perrot ${ }^{d}$ N.A. Smirnova ${ }^{\text {g.1 }}$, C. Stodel ${ }^{\text {d }}$, J.C. Thomas ${ }^{\text {d }}$, N. Vermeulen ${ }^{\text {a }}$, D.T. Yordanov ${ }^{\text {a }}$, The search for the boundaries of the lols
 Y. Utsuno ${ }^{\text {h }}$, T. Otsuka ${ }^{\text {i.j }}$


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"K






Summit of the $N=40$ island of inversion: Precision mass measurements and ab initio calculations of neutron-rich chromium isotopes







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$$
28,40
$$





## Shape Coexistence between Z = 28 and Z = 50

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Shape Coexistence is predicted in 78 Ni but still no firm proof $\quad N=50$
nature ${ }^{78} \mathrm{Ni}$ revealed as a doubly magic stronghold against nuclear deformation
R. Taniuchil ${ }^{1,2}$, C. Santamaria ${ }^{2,3}$, P. Doornenbal ${ }^{2 *}$, A. Obertelli ${ }^{2,3,4}$, K. Yoneda ${ }^{2}$, G. Authelet ${ }^{3}$, H. Baba ${ }^{2}$, D. Calvet ${ }^{3}$, F. Château ${ }^{3}$, R. Taniuchi ${ }^{1,2}$, C. Santamaria ${ }^{2,3}$, P. Doornenbal ${ }^{2 *}$, A. Obertelli ${ }^{2,3,4}$, K. Yoneda ${ }^{2}$, G. Authelet ${ }^{3}$, H. Baba ${ }^{2}$, D. Calvet ${ }^{3}$, F. C
A. Corsi ${ }^{3}$, A. Delbart ${ }^{3}$, J.-M. Gheller ${ }^{3}$, A. Gillibert ${ }^{3}$, J. D. Holt ${ }^{5}$, T. Isobe ${ }^{2}$, V. Lapoux ${ }^{3}$, M. Matsushita ${ }^{6}$, J. Menéndez ${ }^{6}$, A. Corsi ${ }^{3}$, A. Delbart ${ }^{3}$, J.-M. Gheller ${ }^{3}$, A. Gillibert ${ }^{3}$, J. D. Holt ${ }^{5}$, T. Isobe ${ }^{2}$, V. Lapoux ${ }^{3}$, M. Matsushita ${ }^{6}$, J. Menéndez ${ }^{6}$, ${ }^{10}$,
S. Momiyama ${ }^{1,2}$, T. Motobayashi ${ }^{2}$, M. Niikura ${ }^{1}$, F. Nowacki ${ }^{7}$ K. Ogata ${ }^{8,9}$, H. Otsu ${ }^{2}$, T. Otsuka ${ }^{1,2,6}$, C. Péron ${ }^{3}$, S. Péru ${ }^{10}$, S. Momiyama ${ }^{1,2}$, T. Motobayashi ${ }^{2}$, M. Niikura ${ }^{1}$, F. Nowacki ${ }^{7}$, K. Ogata ${ }^{8,9}$, H. Otsu ${ }^{2}$, T. Otsuka ${ }^{1,2,6}$, C. Péron ${ }^{3}$, S. Péru ${ }^{10}$
A. Peyaud ${ }^{3}$, E. C. Pollacco ${ }^{3}$ A. Poves ${ }^{11}$ J.-Y. Rousse ${ }^{3}$ H. S. Sakurai ${ }^{1,2}$ A. Schwenk ${ }^{4,12,13}$ Y. Shiga ${ }^{2,14}$, J. Simonis ${ }^{4}, 12,15$,
 S. R. Stroberg ${ }^{5}, 16$, S. Takeuchi ${ }^{2}$, Y. Tsunoda ${ }^{6}$, T. Uesaka ${ }^{2}$, H. Wang ${ }^{2}$, F. Browne ${ }^{17}$, L. X. Chung , ${ }^{18}$, Z. Dombradi ${ }^{19}$, S. Franchoo ${ }^{20}$, F. Giacoppo ${ }^{21}$, A. Gottardo ${ }^{20}$, K. Hadyńska-Klęk ${ }^{21}$, Z. Korkulu ${ }^{19}$, S. Koyama ${ }^{1,2}$, Y. Kubota²,6, J. Lee ${ }^{22}$, M. Lettmann ${ }^{4}$, C. Louchart ${ }^{4}$ R. Lozeva², ${ }^{23}$, K. Matsui ${ }^{1,2}$, T. Miyazaki ${ }^{1,2}$, S. Nishimura ${ }^{2}$, L. Olivier ${ }^{20}$, S. Ota ${ }^{6}$, Z. Patel ${ }^{24}$, E. Şahin ${ }^{21}$, C. Shand ${ }^{24}$, P.-A. Söderström I. Stefan ${ }^{20}$, D. Steppenbeck ${ }^{6}$, T. Sumikama ${ }^{25}$, D. Suzuki ${ }^{20}$, Z. Vajta ${ }^{19}$, V. Werner ${ }^{4}$, J. Wu ${ }^{2}, 26$ \& Z. Y. Xu ${ }^{22}$



## Shape Coexistence between $\mathrm{Z}=28$ and $\mathrm{Z}=50$



## Triaxiality in the $\mathrm{Se}, \mathrm{Ge}$ and Zn Isotopes



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PHYSICAL REVIEW C 103, 014311 (2021)
Onset of triaxial deformation in ${ }^{66} \mathrm{Zn}$ and properties of its first excited $0^{+}$state studied by means of Coulomb excitation
M. Rocchini ${ }^{1,2,3, *}$ K. Hadyńska-Kleç, ${ }^{4,5}$ A. Nannini ${ }^{1,2}$ A. Goasduff ${ }^{6,6,7}$ M. Zielińska, ${ }^{8}$ D. Testov $\odot^{6,7}$ T. R. Rodríguez, ${ }^{9}$ A. Gargano, ${ }^{10}$ F. Nowackio, ${ }^{11}$ G. De Gregorio ${ }^{10,12}{ }^{12}$ H. Naïdja, ${ }^{13}$ P. Sona, ${ }^{1,2}$ J. J. Valiente-Dobón, ${ }^{4}$ D. Mengoni, ${ }^{6}$
P. R. John, ${ }^{6,7,14}$ D. Bazzacco ${ }^{6,7}$ G. Benzoni ${ }^{15}{ }^{15}$ A. Boso, ${ }^{6,7}$ P. Cocconi, ${ }^{4}$ M. Chiari ${ }^{1,1,2}$ D. T. Doherty, ${ }^{16}$ F. Galtarossa, ${ }^{4}$ G. Jaworski, ${ }^{4}$ M. Komorowska, ${ }^{5}$ N. Marchini, ${ }^{1 / 17}$ M. Matejska-Minda ${ }^{5,18}$, ${ }^{5}$ B. Melon, ${ }^{1}$ R. Menegazzo, ${ }^{6,}{ }^{6}$ P. J. Napiorkowski, ${ }^{5}$ D. Napoli,${ }^{4}$ M. Ottanelli, ${ }^{1}$ A. Perego, ${ }^{1,2}$ L. Ramina, ${ }^{7}$ M. Rampazzo, ${ }^{7}$ F. Recchia ${ }^{6,7}{ }^{, 7}$ S. Riccetto, ${ }^{19,20}$




| $Z=50$ |  | $N=50$ |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | ${ }^{2005} \mathrm{~S}$ | ${ }^{102} \mathrm{Sn}$ | ${ }^{1 a} \mathrm{Sn}$ | ${ }^{100} \mathrm{Sn}$ | ${ }^{200} \mathrm{Sn}$ | ${ }^{112} \mathrm{Sn}$ | ${ }^{12} \mathrm{~S}$ n | MSn | ${ }^{1: C S n}$ | ${ }^{110} \mathrm{Sn}$ |
| ${ }^{\circ} \mathrm{Cd}$ | ${ }^{3} \mathrm{Cd}$ | ${ }^{\circ} \mathrm{CA}$ Cd | ${ }^{10 \mathrm{Cd}}$ | ${ }^{100} \mathrm{Cd}$ | ${ }^{10 \times C d}$ | ${ }^{100} \mathrm{Cd}$ | ${ }^{100} \mathrm{Cd}$ | ${ }^{19} \mathrm{Cd}$ | PCd | ${ }^{124} \mathrm{Cd}$ | ${ }^{115} \mathrm{Cd}$ |
| ${ }^{2 \times P d}$ | ${ }^{4.4 P d}$ | ${ }^{\text {rippd }}$ | ${ }^{39 P d}$ | $1{ }^{10 \mathrm{~Pa}}$ | ${ }^{1020} \mathrm{Pd}$ | $\mathrm{rapd}^{\text {coper }}$ | ILSPd | $\mathrm{mpapd}$ | ${ }_{112} \mathrm{Pd}$ | ${ }^{127 P d}$ | \|14.Pd |
| ${ }^{20} \mathrm{RU}$ | ${ }^{22 R} \mathrm{R}$ | 2RRu | ${ }^{\text {s*RU }}$ R | 9RU | ${ }^{100 \mathrm{RU}}$ | ExRU | ${ }^{29 R R u}$ | ${ }^{106 R U}$ | ${ }^{\text {1asRu }}$ | ${ }^{\text {H.ORU }}$ | 112RU |
| ${ }^{\text {asmo }}$ | 9 Mo | ${ }^{22} \mathrm{MO}$ | MMO | 8 Mo | *Mo | 1100 Mo | ${ }^{102 \mathrm{Mo}}$ | 13 Mo | ${ }^{100 \mathrm{MO}}$ | 1209 Mo | ${ }^{120} \mathrm{MO}$ |
| ewr | a ${ }^{\text {ar }}$ r | ${ }^{30} \mathrm{Zr}$ | $\stackrel{\text { P2r }}{ }$ | *Zr | ${ }^{\text {s/2, }}$ r | sezr | ${ }_{120 \mathrm{Zr}}$ | ${ }^{102} \mathrm{Zr}$ | ${ }^{10+2} \mathrm{Zr}$ | ${ }^{100 \mathrm{Z}} \mathrm{Z}$ | ${ }^{108} \mathrm{Zr}$ |
| ${ }^{\text {eas }}$ S | Sr | ${ }^{3} \mathrm{Sr}$ | ${ }^{30} \mathrm{Sr}$ | ${ }^{2} \mathrm{Sr}$ | ${ }^{3}+\mathrm{Sr}$ | sosr | wsr | ${ }^{\text {nowSr }}$ | ${ }^{120} 5 \mathrm{Sr}$ | ${ }^{12045}$ | ${ }^{100} \mathrm{Sr}$ |
| ${ }_{8 \times \mathrm{Kr}}$ | ${ }^{3} \mathrm{Kr} \mathrm{r}^{\text {r }}$ | ${ }^{35 \mathrm{Kr}}$ r | ${ }_{80 \mathrm{Kr} \mathrm{r}}$ | ${ }^{*} \mathrm{Kr}$ | ${ }^{2} 2 \mathrm{Kr}$ | 84 Kr | ${ }^{6} \mathrm{Kr}$ r | ${ }^{99 \mathrm{Kr}}$ | ${ }^{100 \mathrm{Kr}}$ |  |  |
| ${ }^{205} \mathrm{Se}$ | ${ }^{23} \mathrm{Se}$ | ${ }^{84} 5 \mathrm{e}$ | ${ }^{365}$ S | ${ }^{3} \mathrm{~S}$ e | ${ }^{\text {cose }}$ | \%Se | ${ }^{45} 5$ |  |  |  |  |
| ${ }^{7 \times G}$ Ge | me | ${ }^{12 \mathrm{Ge}}$ | ${ }^{4} \mathrm{Ge}$ | ${ }^{\text {moGe }}$ | ${ }^{34 G e}$ | ${ }^{\text {senge }}$ |  |  |  |  |  |
| 702n | ${ }^{7} \mathrm{Zn}$ | $\mathrm{svz}_{\mathrm{Zn}}$ | 8 zn | ${ }^{4} \mathrm{Zn}$ | eezn |  |  |  |  |  |  |
| ${ }^{74} \mathrm{Ni}$ | ${ }^{\text {s/ }} \mathrm{Ni}$ | ${ }^{88} \mathrm{Ni}$ | ${ }^{\infty} \mathrm{Ni}$ | ${ }^{32} \mathrm{Ni}$ | $=28$ |  |  |  |  |  |  |

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Axial Shape Asymmetry and Configuration Coexistence in Neutron-Rich 74 Zn

## Triaxiality in the Se, Ge and Zn Isotopes

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- ISAC facility: Radioactive Ion Beam production using the ISOL technique
- ISAC-I $\Rightarrow$ Non-reaccelerated beams $(20-40 \mathrm{keV}) \Rightarrow$ GRIFFIN
- ISAC-II $\Rightarrow$ Post-accelerated beams (up to $\sim 10 \mathrm{MeV} / \mathrm{A}) \Rightarrow$ TIGRESS
- GRIFFIN (Gamma-Ray Infrastructure For Fundamental Investigations of Nuclei): High-efficiency $\gamma$-ray spectrometer equipped with many ancillary devices


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ISAC-I and ISAC-II Facility
EMM


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## Our Experiment on ${ }^{74} \mathrm{Zn}$ with GRIFFIN

- ${ }^{74} \mathrm{Zn}$ via ${ }^{74} \mathrm{Cu} \beta$-decay $\left[\mathrm{T}_{1 / 2}=1.63(5) \mathrm{s}\right]$, Beam intensity $\approx 1.5 \cdot 10^{3} \mathrm{pps}$
- GRIFFIN: 12 of 16 available clovers at 14.5 cm from the target
- $\varepsilon_{Y}(1332.5 \mathrm{keV})=7.8 \%, \varepsilon_{\mathrm{Y}}(300 \mathrm{keV})=16.6 \%$
- P/T (addback + BGO suppressors) $=45.5 \%$
- Tape cycle: $5 \mathrm{~T}_{1 / 2}$ on, 1 s off, 0.5 s background, 1 s tape movement





## Y- $\mathbf{\gamma}$ Angular Correlations

- $\quad \mathrm{Y}-\mathrm{Y}$ Angular Correlations with GRIFFIN: J.K. Smith et al., NIMA 922, 47 (2019)
- Rhombicuboctahedron geometry $\Rightarrow$ Up to 52 opening angles

- Event mixing technique $\Rightarrow$ No need to know \# of pairs for each opening angle and relative efficiencies of the detectors
- Finite sizes of the detectors $\Rightarrow$ Detailed GEANT4 simulations
- Definitive spin assignments at the 99\% CL


## Previous Knowledge

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- ${ }^{74} \mathrm{Cu} \beta$-decay @ORNL, 3 clovers J.L. Tracy Jr. et al., PRC 98, 034309 (2018)
- 170 y-rays
- 50 levels (29 new)
- log(ft) values
- Tentative spin assignments based on decay patterns, intensities, $\log (\mathrm{ft})$ values, and in a few cases model predictions (shell model and vibrational model)



## $\mathrm{y}-\mathrm{Y}$ Angular Correlations: the $\left(\mathrm{O}_{2}{ }^{+}\right)$




## Experimental Results in a Nutshell

Iols, Shape Coexistence and Triaxiality

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- New, definitive spin assignment for:
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\text { - } 2_{3^{+}} \longrightarrow 4_{1^{+}} \text {and } 2_{3^{+}} \longrightarrow 0_{2^{+}}
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- From measured branching ratios and $\delta(E 2 / M 1)$ mixing ratios $\Rightarrow$ Relative $B(E 2)$ values



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Strong transitions observed, indicative of band structures at low-spin in 74 Zn


## New Large-Scale Shell-Model Calculations

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- By Silvia Lenzi, Frédéric Nowacki, Duc D. Dao
- LNPS interaction, pf shell for protons, $1 p_{3 / 2} \mathrm{Of}_{5 / 2} 1 \mathrm{p}_{1 / 2} \mathrm{Og}_{9 / 2} 1 \mathrm{~d}_{5 / 2}$ orbitals for neutrons
- DNO-SM: Constrained Hartree-Fock shell-model calculations D.D. Dao and F. Nowacki, PRC 105, 054314 (2022)
- Excellent agreement both for the spectra and the relative $B(E 2)$ values


In-band transitions in the calculated spectrum are labelled with the respective $B(E 2)$ values in Weisskopf units

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TABLE I. Energies $E_{\gamma}$, branching ratios $I_{\gamma}$, mixing ratios $\delta(E 2 / M 1)$, and relative $B(E 2)$ values $B_{\text {rel }}(E 2)$ measured in the present Letter, together with branching ratios from Ref. [40]. Relative and absolute $B(E 2)$ values obtained from the present LSSM calculations (full diagonalization) are also given. Relative $B(E 2)$ values of 100 are assumed for normalizing transitions.

| $J_{i}^{\pi} \rightarrow J_{f}^{\pi}$ | $E_{\gamma}(\mathrm{keV})$ | $I_{\gamma}$ | $I_{\gamma}^{\text {prev }}$ [40] | $\delta(E 2 / M 1)$ | $B_{\text {rel }}^{\text {exp }}(E 2)$ | $B_{\text {rel }}^{\text {SM }}(E 2)$ | $B_{\text {abs }}^{\mathrm{SM}}(E 2)$ (W.u.) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $2_{2}^{+} \rightarrow 2_{1}^{+}$ | 1064.32(10) | 100.0(12) | 100.0(6) | -1.13(6) | 100(5) | 100 | 9.7 |
| $2_{2}^{+} \rightarrow 0_{1}^{+}$ | 1670.07(20) | 49.3(10) | 49.4(4) |  | 9.24(19) | 22 | 2.1 |
| $3_{1}^{+} \rightarrow 2_{2}^{+}$ | 428.73(18) | 6.5(4) | 9.3(4) | $-0.8_{-1.5}^{+0.2}$ | $100_{-30}^{+120}$ | 100 | 40 |
| $3_{1}^{+} \rightarrow 4_{1}^{+}$ | 680.75(15) | 7.10 (19) | 10.5(4) | $-1.0_{-0.8}^{+0.3}$ | $14_{-5}^{+7}$ | 7.8 | 3.1 |
| $3_{1}^{+} \rightarrow 2_{1}^{+}$ | 1493.2(3) | 100.0(18) | 100.0(11) | $\begin{gathered} -0.57_{-0.07}^{+0.06} \\ -2.7(5)^{a} \\ -2 . \end{gathered}$ | $\begin{gathered} 1.9_{-0.3}^{+0.4} \\ 6.8(4) \end{gathered}$ | 8.8 | 3.5 |
| $2_{3}^{+} \rightarrow 0_{2}^{+}$ | 359.2(6) | 2.0(4) |  |  | 100(20) | 100 | 17 |
| $2_{3}^{+} \rightarrow 2_{2}^{+}$ | 478.13(15) | 6.8(7) | 6.5(10) | $+0.9{ }_{-0.3}^{+0.8}$ | $37_{-15}^{+24}$ | 15 | 2.6 |
| $2_{3}^{+} \rightarrow 4_{1}^{+}$ | 729.94(19) | 3.1(7) |  |  | 4.5(10) | 2.4 | 0.4 |
| $2_{3}^{+} \rightarrow 2_{1}^{+}$ | 1542.5(3) | 37(3) | 29.4(14) | $+2.4{ }_{-1.0}^{+1.8}$ | $1.09_{-0.26}^{+0.15}$ | 0.18 | 0.03 |
| $2_{3}^{+} \rightarrow 0_{1}^{+}$ | 2148.73(16) | 100(8) | 100.0(27) |  | 0.66 (5) | 0.18 | 0.03 |

[^0]
## Calculated Shapes from Shell Model

Iols, Shape
Coexistence and
Triaxiality
GRIFFIN
@TRIUMF
Our Experiment
on 74 Zn
Y-y Angular
Correlations
Experimental
Results
LSSM
Calculations
Shapes and
Occupations
Conclusions

- Shapes of the $0_{1}{ }^{+}, 0_{2}{ }^{+}$states calculated from quadrupole sum rules A. Poves, F. Nowacki, and Y. Alhassid, PRC 101, 054307 (2020)



## Calculated Shapes from Shell Model

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A. Poves, F. Nowacki, and Y. Alhassid, PRC 101, 054307 (2020)
- Triaxial ground state (similar dispersions to those in ${ }^{76} \mathrm{Ge}$ !!)

PHYSICAL REVIEW LETTERS 123, 102501 (2019)


Evidence for Rigid Triaxial Deformation in ${ }^{76} \mathbf{G e}$ from a Model-Independent Analysis
A. D. Ayangeakaa@, ${ }^{1,{ }^{*}}$ R. V.F. Janssens, ${ }^{2,3, \dagger}$ S. Zhu, ${ }^{4,+\hbar}$ D. Little ${ }^{2,3}$ J. Henderson, ${ }^{5}$ C. Y. Wu, ${ }^{5}$ D. J. Hartley, ${ }^{1}$ M. Albers, ${ }^{4}$ K. Auranen, ${ }^{4}$ B. Bucher,,${ }^{5,8}$ M. P. Carpenter, ${ }^{4}$ P. Chowdhury, ${ }^{6}$ D. Cline, ${ }^{7}$ H. L. Crawford, ${ }^{8}$ P. Fallon, ${ }^{8}$ A. M. Forney, ${ }^{,}$ A. Gade, ${ }^{10,11}$ A. B. Hayes, ${ }^{7}$ F. G. Kondev, ${ }^{4}{ }^{4}$ Krishichayan, ${ }^{3,12}$ T. Lauritsen, ${ }^{4}$ J. Li, ${ }^{4}$ A. O. Macchiavelli, ${ }^{4}$ D. Rhodes, ${ }^{10}$ D. Seweryniak, ${ }^{4}$ S. M. Stolze, ${ }^{4}$ W. B. Walters, ${ }^{9}$ and J. Wu ${ }^{4}$


## Calculated Shapes from Shell Model

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A. Poves, F. Nowacki, and Y. Alhassid, PRC 101, 054307 (2020)
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- Different shape for the $\mathrm{O}_{2}{ }^{+}$: Less deformed and more prolate



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- "T-Plot"-like analysis: Shape for each individual state
- Normalized probability to find a deformation ( $\beta, \gamma$ ) superimposed to the potential energy surface (GCM)
- Results from quadrupole sum rules confirmed



## Calculated Shapes from Shell Model

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- Results from quadrupole sum rules confirmed



## Occupation Numbers from Shell Model

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- Occupation difference with respect to the Fermi configuration for each orbital in the chosen model space:



SM

Axial Shape Asymmetry and Configuration Coexistence in Neutron-Rich ${ }^{74}$ Zn

## Occupation Numbers from Shell Model



Axial Shape Asymmetry and Configuration Coexistence in Neutron-Rich 74 Zn

## Occupation Numbers from Shell Model

Iols, Shape Coexistence and Triaxiality

## GRIFFIN

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Our Experiment on ${ }^{74} \mathrm{Zn}$
$y-y$ Angular
Correlations
Experimental Results

LSSM
Calculations
Shapes and
Occupations
Conclusions

- Occupation difference with respect to the Fermi configuration for each orbital in the chosen model space:

- Which is the intruder state? Neutrons from the pf shell and percentage of 0pOh configuration:



The ground state in
74 Zn seems to be the intruder state 1529

Typical behaviour of a nucleus in an Island of Inversion

## Axial Shape Asymmetry and Configuration Coexistence in Neutron-Rich ${ }^{74} \mathrm{Zn}$

## Conclusions

## Iols, Shape

Coexistence and Triaxiality

## GRIFFIN

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Our Experiment on ${ }^{74} \mathrm{Zn}$

## $\gamma-\gamma$ Angular

 CorrelationsExperimental Results

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Conclusions

- Strong $2_{3}{ }^{+} \longrightarrow \mathrm{O}_{2}{ }^{+} \Longrightarrow$ Hint of Configuration Coexistence
- Strong $31^{+} \longrightarrow 2_{2}{ }^{+} \Longrightarrow$ Hint of a quasi $\gamma$-band at low excitation energy and Triaxiality
- New Large-Scale Shell-Model calculations support this interpretation
- Inversion of "normal" and intruder configurations $\Rightarrow{ }^{74} \mathrm{Zn}$ seems to be in the $\mathrm{N}=40$


Island of Inversion, which extends further north in the chart of the nuclides

## PHYSICAL REVIEW LETTERS 130, 122502 (2023)

First Evidence of Axial Shape Asymmetry and Configuration Coexistence in ${ }^{74} \mathrm{Zn}$ : Suggestion for a Northern Extension of the $N=40$ Island of Inversion
M. Rocchini®,,${ }^{1, *}$ P. E. Garrett®, ${ }^{1}$ M. Zielińska®, ${ }^{2}$ S. M. Lenzi® ${ }^{3,4}$ D. D. Dao®, ${ }^{5}$ F. Nowacki, ${ }^{5}$ V. Bildstein, ${ }^{1}$ A. D. MacLean, ${ }^{1}$ B. Olaizola $\odot,^{6, \dagger}$ Z. T. Ahmed, ${ }^{1}$ C. Andreoiu ${ }^{1}{ }^{7}$ A. Babu, ${ }^{6}$ G. C. Ball, ${ }^{6}$ S. S. Bhattacharjee, ${ }^{6,+}$ H. Bidaman, ${ }^{1}$ C. Cheng, ${ }^{6}$ R. Coleman, ${ }^{1}$ I. Dillmann $\odot{ }^{6,8}$ A. B. Garnsworthy, ${ }^{6}$ S. Gillespie, ${ }^{6}$ C. J. Griffin®, ${ }^{6}$ G. F. Grinyer®, ${ }^{9}$ G. Hackman, ${ }^{6}$ M. Hanley©,${ }^{10}$ A. Illana $\odot{ }^{11}$ S. Jones, ${ }^{12}$ A. T. Laffoley, ${ }^{1}$ K. G. Leach $\odot,{ }^{10}$ R. S. Lubna, ${ }^{6,8}$ J. McAfee,,${ }^{6,13}$ C. Natzke, ${ }^{6,10}$ S. Pannu, ${ }^{1}$ C. Paxman $\odot,{ }^{6,13}$ C. Porzio $\odot,{ }^{6,14,15, \|}$ A. J. Radich, ${ }^{1}$ M. M. Rajabali, ${ }^{16}$ F. Sarazin $\odot,{ }^{10}$ K. Schwarz, ${ }^{6}$



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## THANK YOU FOR THE ATTENTION

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11 University of Jyväskylä, Jyväskylä, Finland 12 University of Tennessee, Knoxville, USA 13 Univesity of Surrey, Guildford, UK 14 INFN Sezione di Milano, Milano, Italy

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Marco Rocchini

## INFN - Istituto Nazionale di Fisica Nucleare

 FIRENZE DIVISION
## BACKUP

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## $\mathrm{O}_{2}{ }^{+}$States in the $\mathrm{Ge}, \mathrm{Zn}$ and Ni Isotopes

- Stable $\mathrm{Ge}(\mathrm{N}=38-44) \Longrightarrow \mathrm{v}$ excitations across $\mathrm{N}=40$, small-deformation M. Honma, T. Otsuka, T. Mizusaki, and M. Hjorth-Jensen, PRC 80, 064323 (2009)
(Radioactive ${ }^{82} \mathrm{Ge} \Rightarrow \mathrm{v}$ excitations across $\mathrm{N}=50$, large-deformation

Multitude of p-h excitations across shell and sub shell gaps, with
different features and deformations
J. K. Hwang et al., PRC 84, 024305 (2011)


- $\quad \mathrm{Ni}(\mathrm{N}=36-42)$, Multiple Shape Coexistence:
S. Leoni et al., PRL 118, 162502 (2017) - S. Leoni et al., Acta Physica Polonica B 50, 605 (2019) - N. Marginean et al., PRL 125, 102502 (2020)
- $\mathrm{O}_{2}{ }^{+}$in ${ }^{64,66,68} \mathrm{Ni} \Longrightarrow \mathrm{v}$ excitations across $\mathrm{N}=40$, small-deformation
- $\mathrm{O}_{4}{ }^{+}$in ${ }^{64,66} \mathrm{Ni}, \mathrm{O}_{3}{ }^{+}$in ${ }^{68} \mathrm{Ni}, \mathrm{O}_{2}{ }^{+}$in ${ }^{70} \mathrm{Ni} \Rightarrow$ excitations across $\mathrm{Z}=28$, large-deformation


## Energy Systematics



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$$
\mathbf{3 1}^{+} \rightarrow \mathbf{2 1}^{+} \rightarrow \mathbf{0}_{1}{ }^{+}
$$



## Extension of the $\mathrm{N}=40 \mathrm{lol}$

TABLE I: Occupation of the neutron intruder orbitals and percentage of particle-hole excitations across the $N=40$ gap in the ground states and first excited $0^{+}$states of the Zinc isotopes. The red bold number indicates places where the extra intruder population is larger than 1 particule in average in the ground state and delimits the contour of the Island of Inversion.

| Nucleus state | IPM |  |  |  |  |  |  | SM |  |  | 0 p 0 h |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\nu g_{9 / 2}$ |  |  |  |  |  |  | $\nu d_{5 / 2}$ | $\nu g_{9 / 2}$ | $\nu d_{5 / 2}$ |  |
| ${ }^{68} \mathrm{Ni}$ | $0_{1}^{+}$ | 0 | 0 | 0.80 | 0.10 | 62 |  |  |  |  |  |
|  | $0_{2}^{+}$ | 0 | 0 | 2.10 | 0.10 | 17 |  |  |  |  |  |
| ${ }^{70} \mathrm{Ni}$ | $0_{1}^{+}$ | 2 | 0 | 2.86 | 0.17 | 55 |  |  |  |  |  |
|  | $0_{2}^{+}$ | 2 | 0 | 4.47 | 0.41 | 21 |  |  |  |  |  |
| ${ }^{72} \mathrm{Ni}$ | $0_{1}^{+}$ | 4 | 0 | 4.52 | 0.34 | 61 |  |  |  |  |  |
|  | $0_{2}^{+}$ | 4 | 0 | 5.18 | 0.18 | 12 |  |  |  |  |  |
| ${ }^{74} \mathrm{Ni}$ | $0_{1}^{+}$ | 6 | 0 | 5.96 | 0.43 | 58 |  |  |  |  |  |
|  | $0_{2}^{+}$ | 6 | 0 | 5.65 | 0.73 | 44 |  |  |  |  |  |
| ${ }^{76} \mathrm{Ni}$ | $0_{1}^{+}$ | 8 | 0 | 7.84 | 0.36 | 90 |  |  |  |  |  |
|  | $0_{2}^{+}$ | 8 | 0 | 7.39 | 0.77 | 92 |  |  |  |  |  |
| ${ }^{78} \mathrm{Ni}$ | $0_{1}^{+}$ | 10 | 0 | 9.83 | 0.19 | 99 |  |  |  |  |  |
|  | $0_{2}^{+}$ | 10 | 0 | 7.79 | 2.33 | 94 |  |  |  |  |  |
| ${ }^{70} \mathrm{Zn}$ | $0_{1}^{+}$ | 0 | 0 | 2.60 | 0.17 | 4 |  |  |  |  |  |
|  | $0_{2}^{+}$ | 0 | 0 | 2.66 | 0.25 | 12 |  |  |  |  |  |
| ${ }^{72} \mathrm{Zn}$ | $0_{1}^{+}$ | 2 | 0 | 3.99 | 0.35 | 4 |  |  |  |  |  |
|  | $0_{2}^{+}$ | 2 | 0 | 3.97 | 0.29 | 12 |  |  |  |  |  |
| ${ }^{74} \mathrm{Zn}$ | $0_{1}^{+}$ | 4 | 0 | 5.13 | 0.47 | 13 |  |  |  |  |  |
|  | $0_{2}^{+}$ | 4 | 0 | 4.35 | 0.50 | 34 |  |  |  |  |  |
| ${ }^{76} \mathrm{Zn}$ | $0_{1}^{+}$ | 6 | 0 | 5.79 | 0.52 | 54 |  |  |  |  |  |
|  | $0_{2}^{+}$ | 6 | 0 | 6.12 | 0.47 | 47 |  |  |  |  |  |
| ${ }^{78} \mathrm{Zn}$ | $0_{1}^{+}$ | 8 | 0 | 7.67 | 0.48 | 93 |  |  |  |  |  |
|  | $0_{2}^{+}$ | 8 | 0 | 7.73 | 0.44 | 92 |  |  |  |  |  |
| ${ }^{80} \mathrm{Zn}$ | $0_{1}^{+}$ | 10 | 0 | 9.72 | 0.30 | 100 |  |  |  |  |  |
|  | $0_{2}^{+}$ | 10 | 0 | 9.74 | 0.28 | 100 |  |  |  |  |  |

TABLE II. Occupation number $n^{\nu}$ of neutron intruder orbitals from the shell-model calculations (SM) of this work compared to the independent particle model (IPM). The percentage of particlehole excitations across the $N=40$ gap in the ground state of Cr isotopes is also given. The last column features the pairing correlations energy differences $\Delta E_{\text {Pairing }}^{*}$ (in MeV ) evaluated between the ground state and the $2_{1}^{+}$state.

| Nucleus | $n^{\nu}\left(g_{9 / 2}+d_{5 / 2}\right)$ | $0 p 0 h$ | $2 p 2 h$ | $4 p 4 h$ | $6 p 6 h$ | $\Delta E_{\text {Pairing }}^{*}$ |  |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | IPM | SM |  |  |  |  |  |
| ${ }^{60} \mathrm{Cr}$ | 0 | 1.8 | 14 | 75 | 7 | 0 | 1.84 |
| ${ }^{62} \mathrm{Cr}$ | 0 | 3.5 | 1 | 25 | 71 | 3 | 1.49 |
| ${ }^{64} \mathrm{Cr}$ | 0 | 4.3 | 0 | 8 | 71 | 20 | 1.25 |
| ${ }^{66} \mathrm{Cr}$ | 2 | 5.2 | 0 | 40 | 56 | 3 | 1.13 |
| ${ }^{68} \mathrm{Cr}$ | 4 | 6.0 | 6 | 79 | 11 | 0 | 1.24 |

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## Staggering

$$
\begin{aligned}
& S(J) \\
& =\frac{\left\{E\left(J_{\gamma}^{+}\right)-E\left[(J-1)_{\gamma}^{+}\right]\right\}-\left\{E\left[(J-1)_{\gamma}^{+}\right]-E\left[(J-2)_{\gamma}^{+}\right]\right\}}{E\left(2_{1}^{+}\right)},
\end{aligned}
$$



## $\mathrm{N}=40$

PHYSICAL REVIEW C 80, 064323 (2009)

## New effective interaction for $f_{5} p g_{9}$-shell nuclei

M. Honma, ${ }^{1}$ T. Otsuka, ${ }^{2,3,4}$ T. Mizusaki, ${ }^{5}$ and M. Hjorth-Jensen ${ }^{6}$


FIG. 24. (Color online) Occupation numbers of the neutron $g_{9 / 2}$ orbit in the shell-model wave functions. The filling configuration corresponds to the dot-dashed line.

## ${ }_{28}^{68} \mathrm{Ni}_{40}$ : Magicity versus Superfluidity

O. Sorlin, ${ }^{1}$ S. Leenhardt, ${ }^{1}$ C. Donzaud, ${ }^{1}$ J. Duprat, ${ }^{1}$ F. Azaiez, ${ }^{1}$ F. Nowacki, ${ }^{2}$ H. Grawe, ${ }^{3}$ Zs. Dombrádi, ${ }^{4}$
F. Amorini, ${ }^{5}$ A. Astier, ${ }^{6}$ D. Baiborodin, ${ }^{7}$ M. Belleguic, ${ }^{1}$ C. Borcea, ${ }^{8}$ C. Bourgeois, ${ }^{1}$ D. M. Cullen, ${ }^{9},{ }^{*}$ Z. Dlouhy, ${ }^{7}$ E. Dragulescu, ${ }^{8}$ M. Górska, ${ }^{3}$ S. Grévy, ${ }^{10}$ D. Guillemaud-Mueller, ${ }^{1}$ G. Hagemann, ${ }^{11}$ B. Herskind, ${ }^{11}$ J. Kiener, ${ }^{12}$ R. Lemmon, ${ }^{13}$ M. Lewitowicz, ${ }^{14}$ S. M. Lukyanov, ${ }^{15}$ P. Mayet, ${ }^{3}$ F. de Oliveira Santos, ${ }^{14}$ D. Pantalica, ${ }^{7}$ Yu.-E. Penionzhkevich, ${ }^{15}$ F. Pougheon, ${ }^{1}$ A. Poves, ${ }^{16}$ N. Redon, ${ }^{6}$ M. G. Saint-Laurent, ${ }^{14}$ J. A. Scarpaci, ${ }^{1}$
G. Sletten, ${ }^{11}$ M. Stanoiu, ${ }^{14}$ O. Tarasov, ${ }^{15, \dagger}$ and Ch. Theisen ${ }^{17}$

TABLE II. Shell-model and experimental energies of the first excited state $E\left(2^{+}\right)[\mathrm{MeV}]$, calculated $B\left(E 2 \uparrow ; 0_{1}^{+} \rightarrow 2^{+}\right)$ $\left[e^{2} \mathrm{fm}^{4}\right]$, and the number $\left\langle n_{9 / 2}\right\rangle_{\text {extra }}$ of excessive neutrons occupying the $g_{9 / 2}$ orbit are compiled.

|  | ${ }^{62} \mathrm{Ni}$ | ${ }^{64} \mathrm{Ni}$ | ${ }^{66} \mathrm{Ni}$ | ${ }^{68} \mathrm{Ni}$ | ${ }^{70} \mathrm{Ni}$ | ${ }^{72} \mathrm{Ni}$ | ${ }^{74} \mathrm{Ni}$ |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $E\left(2^{+}\right)_{\text {calc }}$ | 1.11 | 1.24 | 1.49 | 1.73 | 1.50 | 1.42 | 1.33 |
| $E\left(2^{+}\right)_{\text {exp }}$ | 1.173 | 1.346 | 1.425 | 2.033 | 1.259 |  |  |
| $B(E 2 \uparrow)_{\text {calc }}$ | 775 | 755 | 520 | 265 | 410 | 505 | 690 |
| $\left\langle n_{9 / 2}\right\rangle_{\text {ex tra }}$ | 0.24 | 0.43 | 0.67 | 1.19 | 0.73 | 0.45 | 0.27 |

## 66Zn Coulex

PHYSICAL REVIEW C 103, 014311 (2021)

## Onset of triaxial deformation in ${ }^{66} \mathrm{Zn}$ and properties of its first excited $0^{+}$state studied by means of Coulomb excitation

M. Rocchini $\odot,{ }^{1,2,3, *}$ K. Hadyńska-Klęk, ${ }^{4,5}$ A. Nannini $\odot,{ }^{1,2}$ A. Goasduff $\odot,{ }^{6,7}$ M. Zielińska, ${ }^{8}$ D. Testov $\odot,{ }^{6,7}$ T. R. Rodríguez, ${ }^{9}$ A. Gargano, ${ }^{10}$ F. Nowacki $\odot,{ }^{11}$ G. De Gregorio $\odot,{ }^{10,12}$ H. Naïdja, ${ }^{13}$ P. Sona, ${ }^{1,2}$ J. J. Valiente-Dobón, ${ }^{4}$ D. Mengoni, ${ }^{6}$ P. R. John,,${ }^{6,7,14}$ D. Bazzacco, ${ }^{6,7}$ G. Benzoni ${ }^{,}{ }^{15}$ A. Boso, ${ }^{6,7}$ P. Cocconi, ${ }^{4}$ M. Chiari $\odot,{ }^{1,2}$ D. T. Doherty, ${ }^{16}$ F. Galtarossa, ${ }^{4}$ G. Jaworski, ${ }^{4}$ M. Komorowska, ${ }^{5}$ N. Marchini, ${ }^{1,17}$ M. Matejska-Minda@ ${ }^{5,18}$ B. Melon, ${ }^{1}$ R. Menegazzo, ${ }^{6,7}$ P. J. Napiorkowski, ${ }^{5}$
D. Napoli $\odot{ }^{4}$ M. Ottanelli, ${ }^{1}$ A. Perego, ${ }^{1,2}$ L. Ramina, ${ }^{7}$ M. Rampazzo, ${ }^{7}$ F. Recchia, ${ }^{6,7}$ S. Riccetto, ${ }^{19,20}$

## D. Rosso ${ }^{4}$ and M. Siciliano ${ }^{4,6}$



FIG. B. Potential energy surfaces for stable Zn isotopes resulting from detormation-constrained Hatree-Fock calculations with
hie particle number projection method (PN-VAP) and Gogny DIS interaction.


FIG. 9. Same as Fig. 8 for ${ }^{66} \mathrm{Zn}$ after projecting onto angular momentum $J=0$. A triaxial shape with a finite dispersion charac lerize the isotope and a second prolate minimum results from the calculation.


FIG. I1. Collective wave functions for selected states in ${ }^{2} \mathrm{Zn}$,
 keV is given on topleff of each C CW. The ecolored frumes are used
to present the suggested tand assignments.

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## Zn: Relativistic Mean-Field

Eur. Phys. J. A 25, 29-39 (2005) DOI 10.1140/ep.ja/i2004-10235-

The European PHYSICAL JOURNAL A

Relativistic mean-field study for Zn isotopes
W.Z. Jiang ${ }^{1,3, \mathrm{a}}, ~ Z . Z$. Ren $^{2}$, T.T. Wang ${ }^{1}$, Y.L. Zhao ${ }^{1}$, and Z.Y. Zhu ${ }^{1,3}$

ig. 8. The Routhian for ${ }^{60} \mathrm{Zn}$ with respect to the quadrupole deformation with the NL-SH set. The coefficient labeled in the gure is the reduction factor of the pairing gap constant. The coefficient 1.0 means no reduction.

$\beta_{2}$

Fig. 9. The same as shown in fig. 8 , but for ${ }^{66} \mathrm{Zn}$.

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## Zn: Covariant Density Functional theory

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Ground state properties of $\mathrm{Zn}, \mathrm{Ge}$, and Se isotopic chains in covariant density functional theory
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Ni Isotopes \& Z = 28
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Novel shape evolution in exotic Ni isotopes and configuration-dependent shell structure
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FGG. 3. (Color online) Potential energy surfuces (PESS) of Ni isotopes, coordinated by the usual $\underline{Q}_{\mathrm{a}}$ and $Q_{\text {, (or }} y$ ). The energy relaive to



FIG. 4: ${ }^{64} \mathrm{Ni}$ potential energy surfaces with (a) full, original interaction used in MCSM calculations [11], and (b) monopole-frozen interaction (i.e., the monopole component is subtracted from the proton-neutron interaction, and singleparticle energies are adjusted to original effective values of the spherical minimum [6]).

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## Possible excited deformed rotational bands in ${ }^{82} \mathrm{Ge}$

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## $0^{+}$States in Ni Isotopes




[^0]:    ${ }^{\mathrm{a}}$ Second solution.

