

MIcro-tpc MAtrix of Chambers for Directional Dark Matter detection and Axion-Like-Particle Exploration

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MIMAC (MIcro-tpc MAtrix of Chambers)

LPSC (Grenoble) : D. Santos, F.Naraghi , N. Sauzet, C. Beaufort (Ph.D)

- -Technical Coordination, Gas circulation and detectors : **O. Guillaudin**
 - Electronics :

G. Bosson, J. Bouvier, J.L. Bouly, L.Gallin-Martel, F. Rarbi T. Descombes

- Data Acquisition:
- Mechanical Structure :
- COMIMAC (quenching) : J-F. Muraz
- J. Giraud J. F. Muraz

IRFU (Saclay): P. Colas, I. Giomataris

CCPM (Marseille): J. Busto, C. Tao

Tsinghua University (Beijing-China): Y. Tao (Ph.D)

Prototype hosted in IHEP (Beijing-China): ZhiminWang, Changgen Yang

WIMP Light Mass window MIMAC- NEWS-G complementarity



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WIMP Light Mass window MIMAC- NEWS-G complementarity



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Directional detection: principle



The signature able to correlate the rare events in a detector to the galactic halo !!

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Angular modulation of WIMP flux

Modulation is sidereal (tied to stars) not diurnal (tied to Sun)



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There are many "angles" for nuclear recoils... **3D** tracks are needed...



Map of recoils in galactic coordinates (HealPix)

10⁸ Events with $E_R = [5,50]$ keV

Robust with respect to Background events

100 WIMP evts + 100 Background evts



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Phenomenology: Discovery

J. Billard et al., PLB 2010
J. Billard et al., arXiv:1110.6079

Proof of discovery: Signal pointing toward the Cygnus constellation

Blind likelihood analysis in order to establish the galactic origin of the signal



Directional experiments around the world



MIMAC: Detection strategy



Scheme of a MIMAC µTPC

Evolution of the collected charges on the anode

Measurement of the ionization energy:

Charge integrator connected to the mesh coupled to a FADC sampled at 50 MHz

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MIMAC-bi-chamber module prototype



3D Tracks: Drift velocity

Magboltz Simulation



• New mixed gas MIMAC target : $CF_4 + x\% CHF_3$ (x=30)

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MIMAC (bi-chamber module)at Modane Underground Laboratory (France) since June 22nd 2012. Upgraded in June 2013, and in June 2014.

-working at 50 mbar (CF₄+28% CHF₃ + 2% C₄H₁₀)

-in a permanent circulating mode
 -Remote controlled

 and commanded

 -Calibration control twice per week

Many thanks to LSM staff

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In any case one needs to measure the ionization released by the particles in the active volume

Example of calibration (MIMAC)

X-ray generator producing fluorescence photons from Cd, Fe, Cu foils. Threshold ~ 1 keV

Circulation system:

Excelent Gain stability in time







MIMAC readout

Dedicated fast electronics (self-triggered) Based on the MIMAC chip (64 channels)

preamplifier signal + FADC: Energy



3D - track



Ionization Quenching Factors SRIM-Simulations (LPSC)



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Ionization Quenching Factor Measurements at LPSC-Grenoble





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Portable Quenching Facility (COMIMAC) (Electrons and Nuclei of known energies)



In a gas detector the IQF depends strongly on the quality of the gas. The IQF needs to be measured periodically (in-situ) in a long term run experiment.

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Electrons Performance







Electrons Performance









Ions Performance







Ionization Quenching Factor for Fluorine in pure CF4 at 50 mbar





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An alpha particle crossing the detector (as an illustration of the MIMAC observables)



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An Electron event (18 keV)



A "recoil event" (~34 keVee)



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Cathode Signal to place the 3D-track

- The cathode signal is produced by the primary electrons. It is produced before the anode signal produced by the avalanche.
- (C. Couturier, Q. Riffard, N. Sauzet et al. JINS 12 (11)P11020,2017)



Measurement in a MIMAC chamber of an alpha passing through the active volume parallel to the cathode at 10 cm distance.

MIMAC-Cathode Signal measurements (C. Couturier, Q. Riffard, N. Sauzet et al. JINS 12 (11)P11020,2017)



Figure 4. Measure of the time differences (TAC) between the grid signal and the delayed cathode signal in the "START Grid" configuration, as a function of the distance of the α source from the anode (green points); error bars correspond to the standard deviation of the mean. A linear fit of these points is superimposed in red and provides the values of the drift velocity and the additional delay.

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First Cathode Signals from the MIMAC bichamber background (O. Guillaudin, D.S. et al. October 2018)

Chamber 1

Chamber 2



Measuring the time between the "event production" and the avalanche signal !! Covering the 26 cm drift distance (13 us x 20 um/ns) !!

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Ionization Energy distribution of the events recorded with the Cathode Signal



Energy range: 1-60 keV

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3D event-localization in MIMAC



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Angular resolution measured with COMIMAC (¹⁹F ions at known kinetic energies) (Y. Tao, I. Moric, et al. (arXiv1903.02159)



Figure 8. MIMAC angular resolution as a function of 19 F ion kinetic energy. At lower energies, the ion tracks are shorter and have more straggling resulting in worse angular resolution and bigger error bars. The angular resolution is better than 20° down to a kinetic energy of 6.3 keV, and is below 10° for a kinetic energy of 9.3 keV. Error bars are derived from the pixel strips pitch and reconstructed track length as described in the text.

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New MIMAC low background detector





Gaz : MIMAC 50 mbar HT grille : -560 V Drift field : -150 V/cm

16,3 % FWHM (6 keV) Gain ~25 000 Energy threshold <1 keV D. Santos (LPSC Grenoble)

Kapton micromegas readout Piralux Pilar
The new 35 cm "new technology" MIMAC detector compared to the old one



$MIMAC - 2m^3 = 16 \text{ bi-chamber modules (}2x 35x35x52 \text{ cm}^3\text{)}$

New technology anode 35cmx35cm

Stretched thin (12 um) grid at 512um.

New electronic board (1792 channels)

Only one big chamber

To explore the New Physics...

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STRONG CP PROBLEM:

- \rightarrow QCD contains a CP-violating term, $\mathcal{L}_{QCD} \supset \bar{\theta} \frac{g_s^2}{32\pi^2} G_a^{\mu\nu} \tilde{G}_{a\mu\nu}$
- \rightarrow From theory $|\bar{\theta}| \in [0,\pi]$; but experimentally $|\bar{\theta}| < 10^{-10}$

Why is CP conserved in QCD?

PECCEI-QUINN SOLUTION:

- \rightarrow New global symmetry $U(1)_{PQ}$
- $\rightarrow U(1)_{PQ}$ breaking \Rightarrow axion
- \rightarrow The QCD vacuum structure is modified by the vev $\langle a \rangle \propto \bar{\theta}$

The axion dynamically compensates the CP-violating term



Standard axion (2/2)

 \rightarrow The axion can couple to photons through a coupling $g_{a\gamma\gamma}$ which depends on the scale of the symmetry breaking f_{PQ} :

$$\mathcal{L}_{QCD}^{eff} \supset \frac{g_{a\gamma\gamma}}{4} a F^{\mu\nu} \tilde{F}_{\mu\nu}$$

 $ightarrow f_{PQ}$ is constrained from astrophysical and cosmological concerns

AXION PROPERTIES:

- $10^{-5} \lesssim m_{PQ} \lesssim 10^{-2} \,\mathrm{eV}$
- $10^{30} \lesssim \tau_{a \to \gamma \gamma} \lesssim 10^{45} \,\mathrm{days}$

 \implies Dark matter candidate... ...but decay $a \rightarrow \gamma \gamma$ undetectable

Axion in extra dimensions (1/3)

FRAMEWORK:

- SM particles are confined into the observable $(1+3){\cal D}$ space
- Singlets under SM gauge group (axion, graviton, ...) can propagate into a $(1+3+\delta)D$ space
- + Extra dimensions δ are compactified with $R < 30\,\mu{\rm m}$



•
$$f_{PQ}^{4D} \gg f_{PQ}^{(4+\delta)D}$$

- f_{PQ}^{4D} is constrained experimentally - $f_{PQ}^{(4+\delta)D}$ sets the scale of the symmetry breaking

Axion in extra dimensions (2/3)

KALUZA-KLEIN DECOMPOSITION:

• In the observable space, the axion is seen as an infinite tower of KK excitations

•
$$\mathcal{L}_{4D}^{eff} \supset \frac{g_{a\gamma\gamma}}{4} \left(\sum_{\mathbf{n}=0}^{\infty} r_{\mathbf{n}} a_{\mathbf{n}}\right) F^{\mu\nu} \tilde{F}_{\mu\nu}$$



AXION-LIKE PARTICLE (ALP):

- The standard axion is identified to the fundamental mode a_0
- Non-diagonal mass matrix

 \Rightarrow the KK excitations pull m_{a_0} away from m_{PQ}

 $\Rightarrow m_{a_0}$ decouples from f_{PQ}

Decay $a_{\mathbf{n}} \to \gamma \gamma$:

$$\tau(a_{\mathbf{n}} \to \gamma \gamma) \sim \left(\frac{m_{PQ}}{m_{a_{\mathbf{n}}}}\right)^3 \tau(a_0 \to \gamma \gamma)$$

- Decay of every single KK excitation \Rightarrow continuum
- In the Sun, KK axions are produced in the range $m_{a_n} \sim (1-25) \text{ keV}$ $\implies 10^{11} \lesssim \tau_{a_n \rightarrow \gamma \gamma} \lesssim 10^{17} \text{ days}$

$a_{\mathbf{n}} \rightarrow \gamma \gamma$ is detectable with $a_{\mathbf{n}}$ produced in the Sun

Event rate (1/2)

THE SOLAR KK AXION MODEL:



- X-rays dark side of the Moon -



- Heating of the solar corona -

 \rightarrow Each parameter of the model is constrained according to observations (test of Newton's law, helioseismology, coronal heating, *etc.*) Di Lella *et al.*, 2003, Astroparticle Physics

Event rate (2/2)

 $a \rightarrow \gamma \gamma$ event rate



 $\sim 1~{\rm event}$ per day in a $2m^3$ terrestrial detector

AXION IN EXTRA DIMENSIONS:

- Increases its mass of ~ 7 orders of magnitude: $\mathcal{O}(10\,\mathrm{keV})$
- Decreases its lifetime of ~ 19 orders of magnitude : $\tau_a \sim \mathrm{age} \ \mathrm{of} \ \mathrm{the} \ \mathrm{universe}$
- Explains some observations (dark matter, coronal heating, etc.)

The event rate:

- Only depends on the detection volume
- Is a continuum in range $(1-25)\,\rm keV$
- Shows a peak around $9\,\rm keV$

SEARCH FOR KK AXION:

- For dark matter studies
- Can probe extra dimensions



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D. Santos (LPSC Grenoble) 44/14



Efficiency as a function of the Gas and pressure

Detectable events - MC simulations 2m³ (Solar KK axion model)



Mean free path of photons in Ar +5% isobutane (300 mbar) as a function of their energy (keV)



^{39/1}

Efficiency in the Argon mixture at 300 mbar vs. KK-axion energy



B: Santos (LPSC Grenoble)

Two sources of production in the Sun (Primakoff effect + photon Coalescence)



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KK-axion orbits



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Event rate expected from the direct flux and gravitationally trapped



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Event rate vs. KK-axion energy



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Exclusion curves

Sensitivity to trapped axions



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Two e⁻(4 keV) sent by COMIMAC (in less than 2 us)



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Conclusions

- A new directional detector of nuclear recoils at low energies has been developed giving a lot of flexibility on targets, pressure, energy range...
- Ionization quenching factor measurements have been determined experimentally and they can be checked in-situ.
- For the first time the 3D nuclear recoil tracks from Rn progeny have been observed.
- New degrees of freedom are available to discriminate electrons from nuclear recoils to improve the DM search for.
- Angular resolution and directional studies of 3D tracks have been performed with COMIMAC.
- MIMAC with its 3D tracks at high spatial resolution opens a new window in the exploration of rare events !
- The 1 m³ will be the validation of a new generation of a large DM high definition detector including directionality (a needed signature for DM discovery)
- An efficient ALP exploration can be performed with the same detector !

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Standard axion properties

$$\mathcal{L}^{eff} = \mathcal{L}_{QCD} + \frac{1}{2} (\partial_{\mu} a)^2 - \frac{1}{2} m_{PQ}^2 a^2 + \frac{a}{f_{PQ}} \xi \frac{g_s^2}{32\pi^2} G_a^{\mu\nu} \tilde{G}_{a\mu\nu} + \frac{g_{a\gamma\gamma}}{4} a F^{\mu\nu} \tilde{F}_{\mu\nu}$$

PARAMETERS:

• f_{PQ} : scale of the symmetry breaking

•
$$g_{a\gamma\gamma} = \frac{\alpha_{em}}{2\pi f_{PQ}} \left(\frac{E}{N} - 1.92(4) \right)$$

 $\frac{E}{N} = \begin{cases} 0 & \text{for KSVZ} \\ \frac{8}{3} & \text{for DFSZ} \end{cases}$

•
$$m_{PQ} = 5.70(7) \left(\frac{10^{12} \text{ GeV}}{f_{PQ}}\right) \mu \text{eV}$$

• $\tau_{a \to \gamma\gamma} = \frac{64\pi}{a^{-\frac{2}{2}}m_{PQ}^{-3}}$

CONSTRAINTS:

• $10^9 \lesssim f_{PQ} \lesssim 10^{12} \, {
m GeV}$

•
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 $\implies {\rm Dark \ matter \ candidate...} \\ ... {\rm but} \ a \rightarrow \gamma \gamma \ {\rm is \ undetectable}$

Axion in extra dimensions

EXTRA DIMENSIONS (ADD THEORY):



KALUZA-KLEIN DECOMPISITION:

• $a(x^{\mu}, \mathbf{y}) = \sum_{\mathbf{n}=0}^{\infty} a_{\mathbf{n}}(x^{\mu}) \cos\left(\frac{\mathbf{n}\mathbf{y}}{R}\right)$ • $\mathcal{L}_{4D}^{eff} \supset \frac{g_{a\gamma\gamma}}{4} \left(\sum_{\mathbf{n}=0}^{\infty} r_{\mathbf{n}}a_{\mathbf{n}}\right) F^{\mu\nu}\tilde{F}_{\mu\nu}$

- $M_P = (2\pi R M_S)^{n/2} M_S$ avec $M_S \sim \mathcal{O}(\text{TeV})$
- $\hat{f}_{PQ} \equiv (2\pi RM_S)^{\delta/2} f_{PQ}$
- $\hat{f}_{PQ} \gg f_{PQ}$
 - \hat{f}_{PQ} experimentally constrained
 - f_{PQ} scale of the $U(1)_{PQ}$ breakdown

The SM gauge fields couple to a linear combination of KK excitations

PECCEI-QUINN MECHANISM:

$$\begin{cases} \langle a_0 \rangle \ = \ \frac{\hat{f}_{PQ}}{\xi} \left(-\bar{\theta} + \ell \pi \right) \ , \quad \ell \in 2\mathbb{Z} \\ \langle a_k \rangle \ = \ 0 \ , \qquad \qquad \forall k > 0 \end{cases} \quad \text{and} \quad \begin{cases} a_0 \ \xrightarrow{U(1)_{PQ}} a_0 + \alpha f_{PQ} \\ a_k \ \xrightarrow{U(1)_{PQ}} a_k \ , \ \forall k > 0 \end{cases}$$

The standard axion is identified to the fundamental mode a_0

KK mass eigenstates

MASS MATRIX:

J

$$\mathcal{M}^{2} = m_{PQ}^{2} \begin{pmatrix} 1 & \sqrt{2} & \sqrt{2} & \sqrt{2} & \cdots \\ \sqrt{2} & 2 + y^{2} & 2 & 2 & \cdots \\ \sqrt{2} & 2 & 2 + 4y^{2} & 2 & \cdots \\ \sqrt{2} & 2 & 2 & 2 + 9y^{2} & \cdots \\ \vdots & \vdots & \vdots & \vdots & \ddots \end{pmatrix} \qquad \text{with:} \\ m_{PQ}^{2} = \xi^{2} \frac{g_{S}^{2}}{32\pi^{2}} \frac{A_{QCD}^{4}}{\hat{f}_{PQ}^{2}} \\ y \equiv \frac{1}{m_{PQ}R} \\ \end{bmatrix} \\ Mass EIGENSTATES: \\ \begin{cases} m_{a_{0}} \sim m_{PQ} & \text{if } m_{PQ}R \ll 1 \\ m_{a_{n}\neq0} \sim \frac{n}{R} & \text{if } m_{PQ}R \ll 1 \\ m_{a_{n}} \sim \frac{2n+1}{2R} & \text{if } m_{PQ}R \gg 1 \\ m_{a_{0}} \sim m_{PQ} & \text{if } m_{PQ}R \gg 1 \\ \end{cases}$$

KK AXION PRODUCTION IN THE SUN:



KK axions production with mass in range $\sim (1-25)\,{\rm keV}$

Trapped axions

AXIONS GRAVITATIONNALLY TRAPPED:

- Axions with $v_a < v_{esc} \sim 617\,{\rm km.s^{-1}}$ are trapped into the Sun's gravitational field \Rightarrow orbit
- Their number increases all along the Sun life



Trapped axions production in Sun

OBSERVATIONS AND PARAMETERS:

- Test of Newton's law: $R < 30\,\mu{\rm m}$
- Hierarchy problem solving: $M_S \sim \mathcal{O}(\text{TeV})$
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SOLAR KK AXION MODEL:

$$\begin{split} \delta &= n = 2 \quad ; \quad R = 10^3 \, \mathrm{keV^{-1}} \\ g_{a\gamma\gamma} &= 9.2 \times 10^{-14} \, \mathrm{GeV^{-1}} \quad ; \quad n_a = 4.07 \times 10^{13} \, \mathrm{m^{-3}} \end{split}$$

Di Lella et al., 2003, Astroparticle Physics
MIMAC-Exclusion limits



Directional detection: comparison of strategies

 Emulsion layers target = C (low masses), Ar, Br, Kr (high masses)



size 40±9 nm

 Anisotropic crystals target = O (low masses), Zn, W (high masses)



No tracks ; only statistical distributions (!) • Low pressure TPCs target = F





D'Ambrosio et al. 2014 Marseille- October 11th 2019

Directional detection: comparison of strategies

• Emulsion

- Anisotropic crystals
 - + 250 A - Depter vs. Z-Axis
- Low pressure TPCs



~100 nm

~10 nm

~1 mm (10⁵ times longuer !!)

(SRIM simulations)

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SRIM simulation of O (20 keV) in ZnO₄W showing the secondary recoils



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C (22 keV) in emulsion (SRIM simulation)



In emulsions and solids the transverse development is in general greater than the longitudinal !!

Directional detection: Directionality 'D'



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How big is a 1 tonne directional detector?

14 m x 14 m x 14 m directional dark matter detector



Mini-BooNE



MINOS

SNO



Super-Kamiokande

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TPC directional detectors

	DRIFT	MIMAC	NEWAGE	DMTPC
	Boulby	Modane	Kamioka	SNOLAB
Gas mix	73%CS2 +25%CF4 +2%O2	70%CF4 +28%CHF3 +2%C4H10	CF4	CF4
Current volume	800 L	6 L	37 L	1000 L
Drift	ion, 50 cm	e—, 25 cm	e—, 41 cm	e—, 27 cm
Threshold (keVee)	20	1	50	20
Readout	Multi-Wire Proportional Counters	Micromegas	micro-pixel chamber +GEM	CCD

Adapted from Mayet et al. [arXiv:1602.03781]

Radon Progeny





RPR: « In coincidence » events



First detection of 3D tracks of Rn progeny

Electron/recoil discrimination

Mesure: $\begin{cases} E_{ioni}(^{214}\text{Pb}) = 32.90 \pm 0.16 \text{ keVee} \\ E_{ioni}(^{210}\text{Pb}) = 45.60 \pm 0.29 \text{ keVee} \end{cases}$

First measurement of 3D nuclear-recoil tracks coming from radon progeny

MIMAC detection strategy validation





RPR events occur at different positions in the detector...



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