

Numerical study of a dressed black hole solution in asymptotically AdS spacetime

Black Holes and Neutron Stars in Modified Gravity Workshop

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November 20, 2019



Context and objectives

Short introduction

- PhD student at LUTH since October 1st, supervised by Philippe Grandclément.
⇒ About the Evolution Problem in Numerical Relativity.
- **Today:** About my 3-month internship at the end of my Master (April-June 2019).

Context

- **Black hole no-hair theorem:** a stationary black hole is uniquely described by its mass, angular momentum and electric charge.
- Break asymptotic flatness hypothesis
⇒ **Asymptotically Anti-de Sitter (AAdS)** spacetime
 - Curved background (negative cosmological constant)
 - Stabilizes radiative systems (radiation reaches infinity in finite time)

Static dressed black hole

- Solve Einstein's equations with $\Lambda = -3/l^2$ and matter content obeying the Klein-Gordon equation.
- The matter consists in a minimally coupled scalar field with self-interacting potential:

$$V(\phi) = -\frac{3}{4\pi G l^2} \sinh^2 \left(\sqrt{\frac{4\pi G}{3}} \phi \right)$$

Static dressed black hole

- **Analytic solution** given by Martínez *et al.* (2004) [[hep-th/0406111](#)]:

$$ds^2 = \frac{r(r + 2G\mu)}{(r + G\mu)^2} \left[- \left(\frac{r^2}{l^2} - \left(1 + \frac{G\mu}{r} \right)^2 \right) dt^2 + \left(\frac{r^2}{l^2} - \left(1 + \frac{G\mu}{r} \right)^2 \right)^{-1} dr^2 + r^2 d\sigma^2 \right]$$

$$\phi = \sqrt{\frac{3}{4\pi G}} \operatorname{artanh} \left(\frac{G\mu}{r + G\mu} \right).$$

- $d\sigma^2$ is the line element of a surface of constant negative curvature.
- For $r \rightarrow \infty$ we recover AdS.
- One integration constant μ involved in the horizon radius and the intensity of the scalar field \Rightarrow mass parameter.

Static dressed black hole

- If $\mu > -l/4G$, event horizon located at

$$r_+ = \frac{l}{2} \left(1 + \sqrt{1 + \frac{4G\mu}{l}} \right)$$

- The horizon has negative curvature.
- There is a 1-to-1 correspondence between r_+ and $\phi(r_+)$.
- Coordinate singularity on the horizon (Schwarzschild-like): not well suited for numerical purposes.

Goals of the internship

- ① Obtain a numerical solution for the static black hole:
 - Parametrization of $d\sigma^2$,
 - Gauge and boundary conditions for regular coordinates.

- ② Use it as starting configuration to simulate a rotating black hole.

Building the system of equations

Bulk equations

- 3+1 formalism: well-suited for a numerical approach.
 - 11 unknowns: N (lapse), B^i (shift), γ_{ij} (3D metric), ϕ (scalar field).
 - 11 equations: $H = 0$, $M^i = 0$, $E_{ij} = 0$, Klein-Gordon. \Rightarrow 3+1 equations with Λ , projections of $T_{\mu\nu}$ and $\partial_t = 0$.

- Parametrization of the base manifold:

$$d\sigma^2 = \frac{1}{\cos^2 \theta} (d\theta^2 + \sin^2 \theta d\varphi^2), \quad \theta \in [0, \pi/2), \varphi \in [0, 2\pi).$$

\Rightarrow This is the *whole* hyperbolic plane \mathbb{H}^2 .

- **Advantage:** Easily implemented within the Kadath library.
 - **Drawback:** Not a compact base manifold.
- 3-vector basis: $(\partial_r, \frac{\cos \theta}{r} \partial_\theta, \frac{\cos \theta}{r \sin \theta} \partial_\varphi)$ orthonormal for our reference metric $\overline{ds}^2 = dr^2 + r^2 d\sigma^2$.

Gauge fixing conditions

- Diffeomorphism invariance: we need to fix the gauge for the numerical resolution.
- **Maximal slicing condition:** $K = 0$
(extrinsic curvature $K_{ij} = -\mathcal{L}_{\mathbf{m}}\gamma_{ij}/2N$, $K = \gamma^{ij}K_{ij}$)
- **Spatial harmonic gauge:** $V^k = 0$ in a modified system where we replace $\mathcal{R}_{ij} \leftarrow \mathcal{R}_{ij} - D_{(i}V_{j)}$ and we define $V^k = \gamma^{ij}(\Gamma_{ij}^k - {}^{(AdS)}\Gamma_{ij}^k)$
 \Rightarrow **Well-posed system** and we recover the solution of Einstein's equations in this gauge.
 \Rightarrow The ${}^{(AdS)}\Gamma_{ij}^k$ are fixed coordinate-dependent coefficients which act as gauge sources.
- Check a posteriori $K = 0$, $V = 0$ to assess numerical convergence.

Apparent horizon (Inner boundary conditions)

We follow the work done by Hugo Roussille at LUTH in 2018 for the simulation of a Kerr black hole.

- Apparent horizon vs. Event horizon (coincide for stationary systems)
- Non-expanding apparent horizon, shearless, at fixed $r = cst$:

$$\Theta = D_i \tilde{s}^i + K_{ij} \tilde{s}^i \tilde{s}^j = 0, \quad B^i = N \tilde{s}^i + \omega (\partial_\varphi)^i,$$

with \tilde{s}^i the unit normal to the horizon. In particular $B^r = \sqrt{\gamma^{rr}} N$.
 \Rightarrow **Everything about rotation is here:** ω parameter.

- 3 remaining degrees of freedom from the differential gauge:

$$N = 1/2, \quad \gamma_{r\theta} = \gamma_{r\varphi} = 0.$$

Apparent horizon (Inner boundary conditions)

- The system of equations is degenerate on the boundary.
- 1D example: $E(x) \equiv a(x)y''(x) + b(x)y'(x) + c(x)y(x) + d(x) = 0$ for $x \in [0, 1]$. If $a(0) = 0$, then $E(0)$ is 1st order and is its own boundary condition at $x = 0$.
- In our case, start from the 3+1 equations and write the matrix of the principal part normal to the horizon $\mathbf{A} \partial_{rr}^2 \mathbf{u} + \dots = 0$. \mathbf{A} has three vanishing eigenvalues $\propto (B^r)^2 - \gamma^{rr} N^2$ with eigenvectors close to $\gamma_{\theta\theta}$, $\gamma_{\theta\varphi}$ and $\gamma_{\varphi\varphi}$, so we use:

$$E_{\theta\theta} = E_{\theta\varphi} = E_{\varphi\varphi} = 0$$

- **Klein-Gordon also degenerate:** we don't need to impose $\phi = \phi_0$.

$$\frac{1}{N} D_i \left(N \left(\gamma^{ij} - \frac{B^i B^j}{N^2} \right) D_j \phi \right) - \frac{dV}{d\phi} = 0$$

AdS border (Outer boundary conditions)

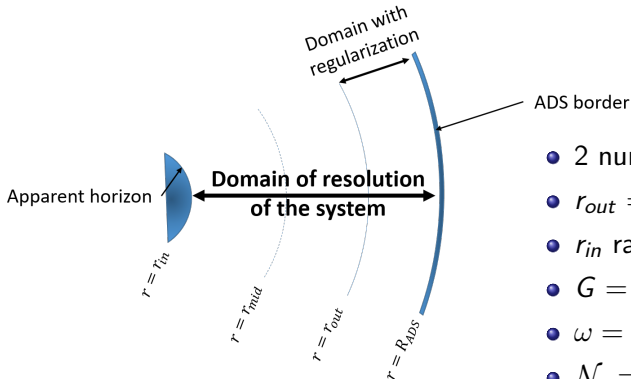
- Vanishing scalar field $\phi = 0$.
- AdS border: by definition, there exists $\Omega \rightarrow 0$ such that N^2 , $\gamma_{ij} \propto \Omega^{-2}$.
 \Rightarrow We need to **regularize** to deal with finite quantities.
- Define regularized quantities such as $\tilde{N} \equiv \Omega N$ and use the regularized equations in a domain reaching to the AdS border.
- Impose continuity of unknown fields and their normal derivatives.
- This process involves a change of variables to write the AdS metric as $ds_{AdS}^2 = -N^2 dt^2 + f(r')(dr'^2 + r'^2 d\sigma^2)$ where $f \propto \Omega^{-2}$.
 \Rightarrow This sets the AdS border at **finite radius**.

Numerical setting and resolution

Spectral methods and the Kadath library

- **Spectral methods:** discrete approximation of a function.
⇒ interpolation on a predefined basis of functions (e.g. Chebyshev polynomials, trigonometric functions).
- **Spectral convergence:** Very fast convergence to the true function (for smooth functions, faster than any polynomial).
⇒ Good precision with a limited number of points.
- **Kadath library:** numerical code developed at LUTH which implements spectral methods and a Newton-Raphson scheme to solve non-linear PDEs.
⇒ Very flexible in terms of geometry, equations to solve, designed with NR in mind, but no evolving systems yet (hence PhD).

Numerical setting

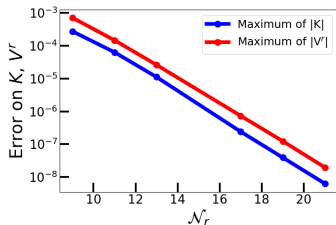


- 2 numerical domains,
- $r_{out} = 30$, $R_{AdS} = 40$,
- r_{in} ranging from 7 to 11,
- $G = 1$, $l = 20$,
- $\omega = 0$ (no rotation),
- $\mathcal{N}_r = 17$, $\mathcal{N}_\theta = 5$, $\mathcal{N}_\varphi = 4$.

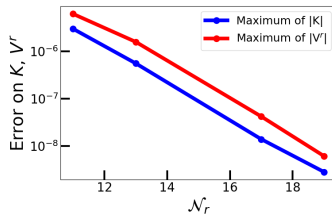
⇒ We vary slowly the horizon radius (0.2 step) and use the current solution as initial guess for the following one.

Tests and results

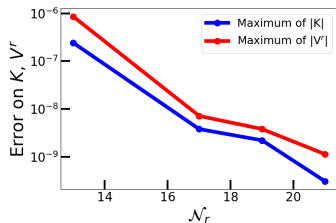
Convergence tests



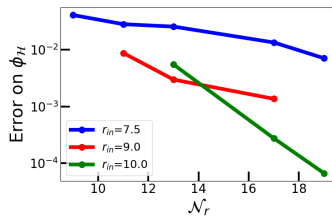
$r_{in} = 7.5$



$r_{in} = 9.0$

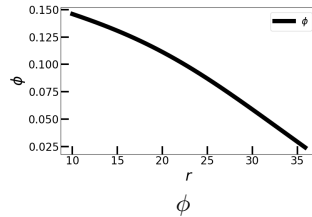
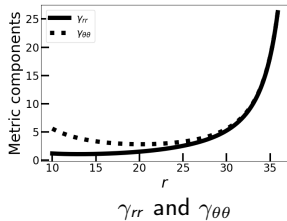
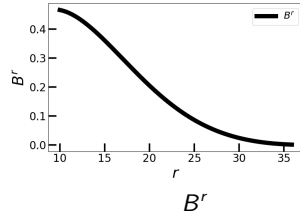
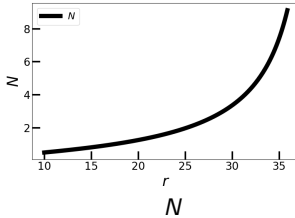


$r_{in} = 10.0$

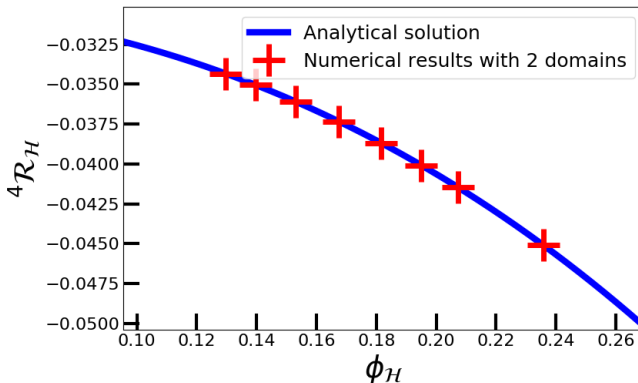


ϕ on horizon

Static Black Hole



Comparison with analytic solution



- Smaller radii correspond to more relativistic configurations.
- Analytic curve computed with SageManifolds.

Comparison with analytic solution

- Probably not a fully comprehensive test but still performs well.
- Global quantities?
 - ⇒ Can be tricky in asymptotically AdS
 - ⇒ Because the surface of \mathbb{H}^2 is infinite, the mass is infinite!

Rotating case and challenges

- Recall $B^\varphi = \omega \frac{r \sin \theta}{\cos \theta}$ on horizon: additional regularization required.
- More importantly, our base manifold is not compact!
⇒ Need a parametrization of a compact surface without boundary and with constant negative scalar curvature.
- We haven't found such a parametrization, and should we get it, it would need to be implemented into Kadath.
- In this case, the boundary conditions on the horizon may be changed.

⇒ **Suggestions are welcome!**

Conclusion

- We have obtained a numerical solution for the static black hole.
- To the extent of our tests, it is consistent with the analytic solution.
- Rotating case: need for a new parametrization of the base manifold, then probably redo the static solution and use further regularization.
- Seems quite involved and not the main purpose of the PhD: stand-by status.
- Suggestions and comments appreciated!

Analytic metric

$$S = \int d^4x \sqrt{-g} \left(\frac{\mathcal{R} + 6I^{-2}}{16\pi G} - \frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - V(\phi) \right), \quad (1)$$

$$V(\phi) = -\frac{3}{4\pi G I^2} \sinh^2 \sqrt{\frac{4\pi G}{3}} \phi. \quad (2)$$

$$G_{\mu\nu} - \frac{3}{I^2} g_{\mu\nu} = 8\pi G T_{\mu\nu}, \quad (3a)$$

$$g^{\mu\nu} \nabla_\mu \nabla_\nu \phi - \frac{dV}{d\phi} = 0, \quad (3b)$$

$$T_{\mu\nu} = \partial_\mu \phi \partial_\nu \phi - \frac{1}{2} g_{\mu\nu} g^{\alpha\beta} \partial_\alpha \phi \partial_\beta \phi - g_{\mu\nu} V(\phi). \quad (4)$$

Analytic metric

$$ds^2 = \frac{r(r + 2G\mu)}{(r + G\mu)^2} \left[- \left(\frac{r^2}{l^2} - \left(1 + \frac{G\mu}{r} \right)^2 \right) dt^2 + \left(\frac{r^2}{l^2} - \left(1 + \frac{G\mu}{r} \right)^2 \right)^{-1} dr^2 + r^2 d\sigma^2 \right] \quad (5a)$$

$$\phi = \sqrt{\frac{3}{4\pi G}} \operatorname{artanh} \left(\frac{G\mu}{r + G\mu} \right). \quad (5b)$$

$$d\sigma^2 = \frac{1}{\cos^2 \theta} (d\theta^2 + \sin^2 \theta \, d\varphi^2). \quad (6)$$

$$ds_{AdS}^2 = - \left(\frac{r^2}{l^2} - 1 \right) dt^2 + \left(\frac{r^2}{l^2} - 1 \right)^{-1} dr^2 + r^2 d\sigma^2. \quad (7)$$

3+1 equations

$$m^\alpha = N n^\alpha = (\partial_t)^\alpha - B^\alpha \quad (8)$$

$$\mathcal{L}_m = \partial_t - \mathcal{L}_B \quad (9)$$

$$\mathcal{R} + K^2 - K_{ij}K^{ij} - 2\Lambda = 16\pi G\rho, \quad (10)$$

$$D_k K_i^k - D_i K = 8\pi G p_i, \quad (11)$$

$$\mathcal{L}_m K_{ij} = -D_i D_j N + N \left[\mathcal{R}_{ij} + K K_{ij} - 2K_{ik} K_j^k - \Lambda \gamma_{ij} - 8\pi G \left(S_{ij} - \frac{S - \rho}{2} \gamma_{ij} \right) \right]. \quad (12)$$

$$\frac{1}{N} D_i \left(N \left(\gamma^{ij} - \frac{B^i B^j}{N^2} \right) D_j \phi \right) - \frac{dV}{d\phi} = 0. \quad (13)$$

Reference metric

$$\overline{ds}^2 = dr^2 + r^2 d\sigma^2 = dr^2 + \frac{r^2}{\cos^2 \theta} (d\theta^2 + \sin^2 \theta d\varphi^2) \quad (14)$$

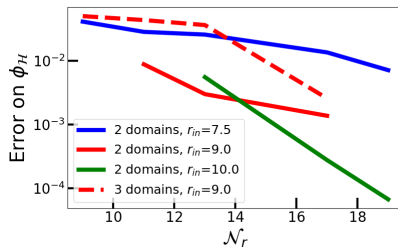
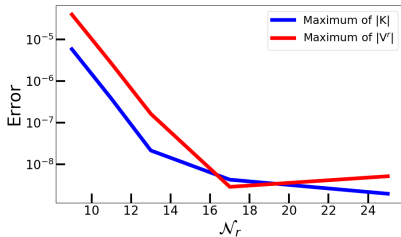
$$ds_{AdS}^2 = -N^2 dt^2 + f(r') (dr'^2 + r'^2 d\sigma^2). \quad (15)$$

$$\begin{cases} r' = R_{AdS} \exp \left[-\arctan \left(\frac{1}{\sqrt{\frac{r^2}{l^2} - 1}} \right) \right] \\ r = \frac{l}{\sin \ln \frac{R_{AdS}}{r}} \end{cases} \quad (16)$$

$$f(r') = \frac{l^2}{r'^2 \sin^2 \ln \frac{R_{AdS}}{r'}}, \quad N^2 = \frac{\cos^2 \ln R_{AdS}/r'}{\sin^2 \ln R_{AdS}/r'}. \quad (17)$$

$$\Omega \equiv \sin \ln \frac{R_{AdS}}{r'}. \quad (18)$$

Convergence tests with 3 domains



Static Black Hole with 3 domains

