

Well-posedness in modified theories of gravity

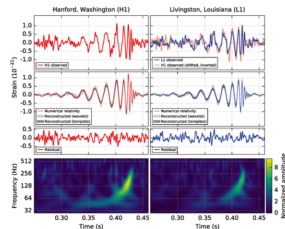
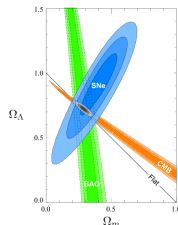
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18 November 2019

Modifying gravity: why?

- **Unexplained phenomena:** dark energy, dark matter, inflation related problems...
- **Theoretical shortcomings:** non-renormalizability of GR, unavoidable formation of singularities...
- **Brand-new observation channel with gravitational waves** → Model dependent tests



Well-posed or ill-posed?

System described by

- differential equation(s) determining its evolution
- initial data set

Well-posed Cauchy problem

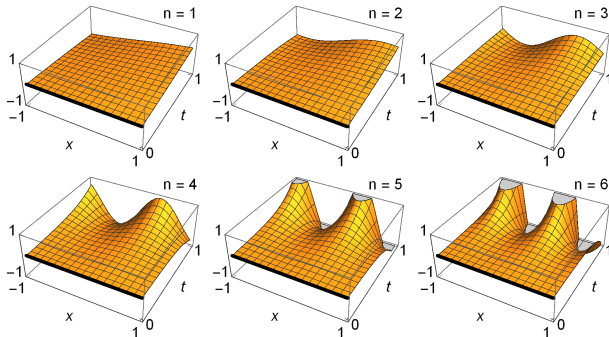
Hadamard 1902

The solution

- 1 exists (at least locally)
- 2 is unique
- 3 depends smoothly on the initial data

Laplace equation: $\partial_t^2 \phi + \partial_x^2 \phi = 0$

$$\phi_n(t, x) = \frac{e^{-\sqrt{n}}}{n} \sin(nx) \sinh(nt)$$



- $\phi_n(0, x) \xrightarrow{n \rightarrow \infty} 0$, $\partial_x \phi_n(0, x) \xrightarrow{n \rightarrow \infty} 0$
- **But** $\phi_n(t, x)$ blows up for $t \neq 0$ when $n \rightarrow \infty$

- 1 Hyperbolic PDEs
- 2 Horndeski theory: ill-posed?
- 3 Curing the ill-posedness

Second order PDEs

Linear PDE

$$P^{\mu\nu} \partial_\mu \partial_\nu u + Q^\mu \partial_\mu u + Ru = 0$$

- μ, ν : spacetime indices (0-3 here)
- u : N -dimensional vector
- $P^{\mu\nu} = P^{(\mu\nu)}$, Q^μ and R : $N \times N$ matrices

- Spatial Fourier transform $\tilde{u}(t, \xi^i)$
- $\tilde{w} \equiv \left(\sqrt{1 + \xi^2} \tilde{u}, -i\partial_t \tilde{u} \right)$

$$\rightarrow \tilde{w}(t, \xi^i) = e^{i\mathcal{M}(\xi^i)t} \tilde{w}(0, \xi^i)$$

Smooth dependence on initial data

$$\begin{aligned} \|\tilde{w}\|_{L^2}(t) < f(t) \|\tilde{w}\|_{L^2}(0) &\Leftrightarrow \|e^{i\mathcal{M}(\xi^i)t}\| < f(t) \\ &\Leftrightarrow \|e^{iM(\xi^i)t}\| < k \end{aligned}$$

where M is the high-frequency part of \mathcal{M}

$$M(\xi^i) = \begin{bmatrix} 0 & \mathbb{1} \\ -A^{-1}C(\xi^i) & -A^{-1}B(\xi^i) \end{bmatrix}$$

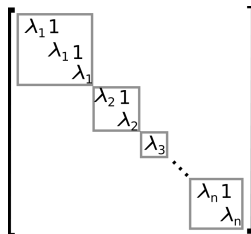
with $A \equiv P^{00}$, $B(\xi^i) \equiv P^{0i}\xi^i$, $C(\xi^i) \equiv P^{ij}\xi^i\xi^j$

Weak hyperbolicity

- M must have only **real eigenvalues**
- Otherwise, $\|e^{iM(\xi^i)t}\|$ grows like $e^{\#t}$

Jordan decomposition

$M = S^{-1}JS$ where J is like on the right



Strong hyperbolicity

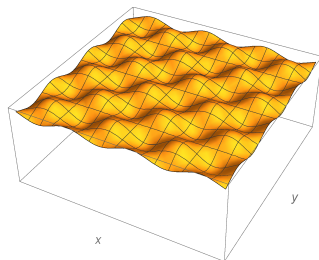
- M must have no non-trivial Jordan block
- $\Leftrightarrow M$ is diagonalizable
- Otherwise, $\|e^{iM(\xi^i)t}\|$ grows like t^p

Possible extensions

- Trivial extension to PDEs with non-constant coefficients
 $P^{\mu\nu}(t, x^i), \dots$
- Non-linear equations?

Non-linear PDEs

Small arbitrary deformation of the initial data set:



- Non-linearly well posed problem \Rightarrow All linearizations around small deviations well-posed
- **Converse result** holds in general: if all arbitrary small deviations yield well-posed linearized problems, then well-posed non linear problem

Kreiss & Lorenz '89

Gauge & constraints (1/2)



Example: electromagnetism in vacuum

- $\partial^\mu (\partial_\mu A_\nu - \partial_\nu A_\mu) = 0$: 4 equations and 4 unknowns
- **But** t component: $\nabla^2 A_0 - \vec{\nabla} \cdot (\partial_t \vec{A}) = 0$ ($\Leftrightarrow \vec{\nabla} \cdot \vec{E} = 0$)
→ Underdetermined evolution for A_μ
- Fine, because 1 function absorbed by gauge

Gauge & constraints (2/2)

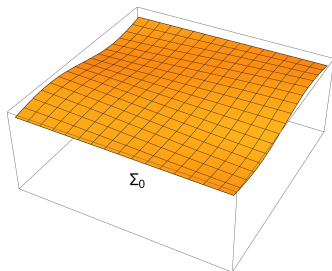
Good gauge?

Start with arbitrary initial data $(u, \partial_t u)$ on Σ_0 and gauge condition $G = 0$.

Then check:

- 1 Possible to impose $G = 0$ on Σ_0 ?
- 2 Constraints $\Rightarrow \partial_t G = 0$ on Σ_0 ?
- 3 $G = 0$ propagated by evolution equations?

+ Of course, are evolution equations well-posed in this gauge?



Well-posedness of GR (1/3)

$$g_{\mu\nu} \rightarrow g_{\mu\nu} + h_{\mu\nu}$$

(Linearized) harmonic gauge

$$H_\alpha \equiv g_{\alpha\beta}{}^{\mu\nu} \nabla^\beta h_{\mu\nu} = 0$$

$$g_{\alpha\beta}{}^{\mu\nu} = \frac{1}{2}(\delta_\alpha^\mu \delta_\beta^\nu + \delta_\alpha^\nu \delta_\beta^\mu - g_{\alpha\beta} g^{\mu\nu})$$

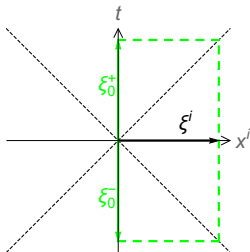
Good gauge in the above sense

Well-posedness of GR (2/3)

Einstein equations in harmonic gauge

$$P_{\alpha\beta}{}^{\rho\sigma\mu\nu} \partial_\mu \partial_\nu h^{\alpha\beta} + \dots = 0$$

- Principal symbol $P_{\alpha\beta}{}^{\rho\sigma\mu\nu} = g_{\alpha\beta}{}^{\rho\sigma} g^{\mu\nu}$
- $N = \frac{4(4+1)}{2} = 10$
- M has $2N$ eigenvectors: $v = (t^{(\alpha\beta)}, \xi_0^\pm t^{(\alpha\beta)})$; $Mv = \xi_0^\pm v$



→ **GR well-posed in harmonic gauge**

Well-posedness of GR (3/3)

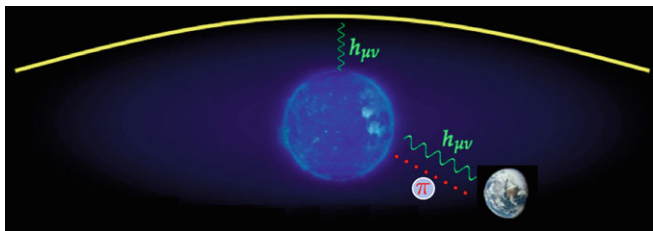
NB

- Well-posedness **depends on gauge**; GR not well-posed in ADM formalism when lapse and shift gauge-fixed
- Stronger statements about well-posedness in GR

- 1 Hyperbolic PDEs
- 2 Horndeski theory: ill-posed?
- 3 Curing the ill-posedness

Modifying gravity: why with a scalar field?

- Simplest additional degree of freedom
- Many theories related in specific regimes



Modifying gravity with a scalar field: how?

Horndeski action

$$S = \int \sqrt{-g} d^4x (\mathcal{L}_2 + \mathcal{L}_3 + \mathcal{L}_4 + \mathcal{L}_5)$$

$$\mathcal{L}_2 = G_2(\phi, X)$$

$$\mathcal{L}_3 = -G_3(\phi, X)\square\phi$$

$$\mathcal{L}_4 = G_4(\phi, X)R + G_{4X} [(\square\phi)^2 - (\nabla_\mu \nabla_\nu \phi)^2]$$

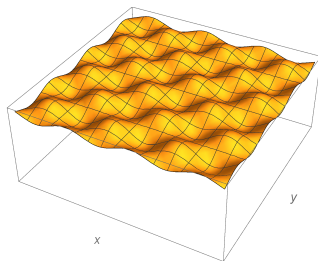
$$X = -\frac{1}{2}(\nabla\phi)^2$$

$$\mathcal{L}_5 = G_5(\phi, X)G_{\mu\nu}\nabla^\mu\nabla^\nu\phi$$

$$- \frac{1}{6}G_{5X} [(\square\phi)^3 - 3\square\phi(\nabla_\mu \nabla_\nu \phi)^2 + 2(\nabla_\mu \nabla_\nu \phi)^3]$$

→ **2nd** order field equations

Horndeski theory: ill-posed?



Claims

Papallo & Reall '17

- Horndeski models with $G_4(\phi, X)$ or $G_5(\phi, X)$ are **ill-posed** around a generic background in any **generalized harmonic gauge**
- Horndeski models with $G_4X = 0$ and $G_5 = 0$ are **well-posed** in the same context

Assumptions

Weak field background

- Technical conditions on the $G_{iX,\phi}$ (small non-Einstein terms)
- NB1: Weak field background **doesn't** imply well-posedness
- NB2: Highly symmetric backgrounds may still be well-posed

$$\begin{aligned}g_{\mu\nu} &\rightarrow g_{\mu\nu} + h_{\mu\nu} \\ \phi &\rightarrow \phi + \psi\end{aligned}$$

Generalized harmonic gauge

$$H_\alpha \equiv g_{\alpha\beta}^{\mu\nu} \nabla^\beta h_{\mu\nu} + f_{\alpha\beta} \nabla^\beta \psi = 0$$

Only other obvious “good gauge”

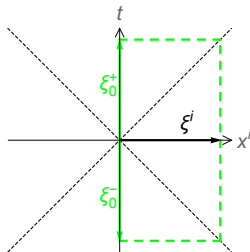
Results

Weak hyperbolicity

- Symmetries of the principal symbol \Rightarrow Weak hyperbolicity for **all Horndeski models**
- Independent of the gauge

Strong hyperbolicity

- Suitable $f_{\alpha\beta}$ for G_2 and G_3
- Impossible for G_4 and G_5 (2×2 Jordan block for ξ_0^+ and for ξ_0^-)





End of the story?

- 1 Hyperbolic PDEs
- 2 Horndeski theory: ill-posed?
- 3 Curing the ill-posedness

Disformal transformation

Einstein-scalar action

$$S = \frac{1}{2} \int d^4x \sqrt{-\tilde{g}} \left[M_{\text{Pl}}^2 \tilde{R} - (\nabla \tilde{\phi})^2 \right]$$

$$\downarrow \quad g_{\mu\nu} = \tilde{g}_{\mu\nu} - D \partial_\mu \phi \partial_\nu \phi$$

Horndeski action

$$G_2(X) = \frac{X}{\sqrt{1 - 2DX}}, \quad G_4(X) = \sqrt{1 - 2DX}$$

$$\phi \rightarrow \phi + \psi$$
$$g_{\mu\nu} \rightarrow g_{\mu\nu} + h_{\mu\nu} \Leftrightarrow \tilde{g}_{\mu\nu} \rightarrow \tilde{g}_{\mu\nu} + \tilde{h}_{\mu\nu}$$

Original harmonic gauge

$$\tilde{g}_{\alpha\beta}^{\mu\nu} \nabla^\beta \tilde{h}_{\mu\nu} = 0$$

↓ Disformal transformation

Non-generalized harmonic gauge!

$$\tilde{g}_{\alpha\beta}^{\mu\nu} \nabla^\beta h_{\mu\nu} + f_{\alpha\beta} \nabla^\beta \psi + \dots = 0 : (G)$$

$$\tilde{g}_{\alpha\beta}^{\mu\nu} \text{ instead of } g_{\alpha\beta}^{\mu\nu}$$

Good gauge?

- 1 Possible to impose $G = 0$ on Σ_0 ?
- 2 Constraints $\Rightarrow \partial_t G = 0$ on Σ_0 ?
- 3 $G = 0$ propagated by evolution equations?

Good gauge?

- 1 Possible to impose $G = 0$ on Σ_0 ? ✓
- 2 Constraints $\Rightarrow \partial_t G = 0$ on Σ_0 ?
- 3 $G = 0$ propagated by evolution equations? ✓



Conclusions

- Claims about generic ill-posedness of G_4 and G_5 Horndeski models
- Not generic enough gauge
- Ongoing proof of well-posedness for some G_4 model

Prospects

- How to generalize to more quartic and quintic models?
- Is well-posedness so essential?

Thank you!