

# B-decay discrepancies after Moriond 2019

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# Outline

## 1 B-decay discrepancies

## 2 New physics interpretations

- $b \rightarrow s\ell\ell$  in the weak effective theory
- The global picture in the Standard Model effective field theory (SMEFT)
- Simplified leptoquark model

## 3 Conclusions

### Based on:

Jason Aebischer, Wolfgang Altmannshofer, Diego Guadagnoli, Méril Reboud, PS, David M. Straub [arXiv:1903.10434]

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# $b \rightarrow s \mu^+ \mu^-$ anomaly

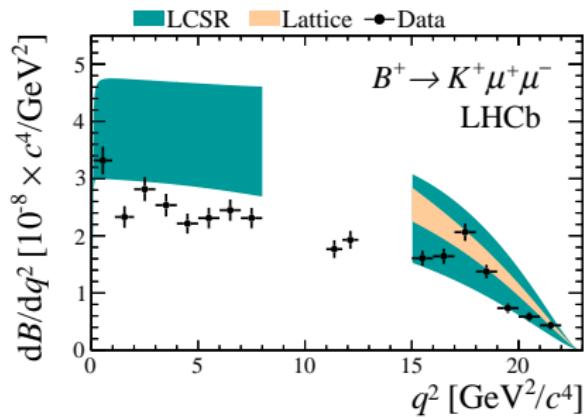
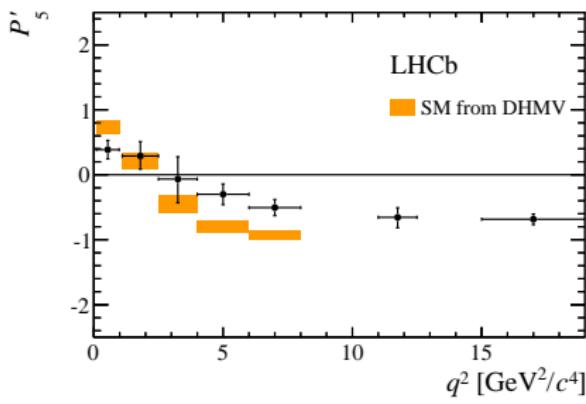
Several LHCb measurements deviate from Standard model (SM) predictions by 2-3 $\sigma$ :

- ▶ Angular observable  $P'_5$  in  $B \rightarrow K^* \mu^+ \mu^-$ .

LHCb, arXiv:1512.04442

- ▶ Branching ratios of  $B \rightarrow K \mu^+ \mu^-$ ,  $B \rightarrow K^* \mu^+ \mu^-$ , and  $B_s \rightarrow \phi \mu^+ \mu^-$ .

LHCb, arXiv:1403.8044, arXiv:1506.08777, arXiv:1606.04731

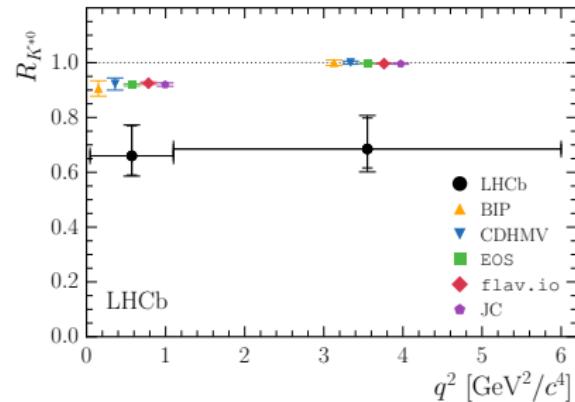
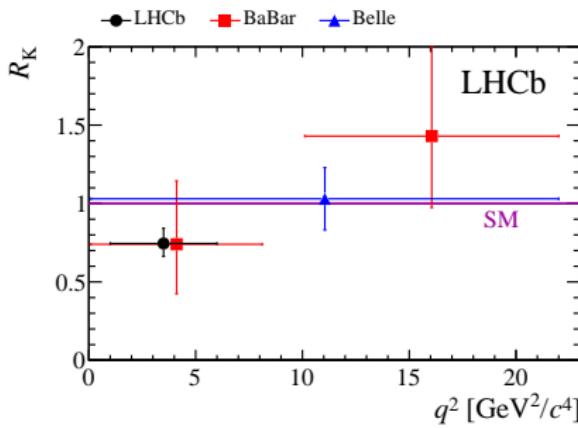


# Hints for LFU violation in $b \rightarrow s \ell^+ \ell^-$ decays

Measurements of lepton flavour universality (LFU) ratios  $R_K^{[1,6]}$ ,  $R_{K^*}^{[0.045, 1.1]}$ ,  $R_{K^*}^{[1.1, 6]}$  showed deviations from SM by about  $2.5\sigma$  each.

LHCb, arXiv:1406.6482, arXiv:1705.05802

$$R_{K^{(*)}} = \frac{BR(B \rightarrow K^{(*)}\mu^+\mu^-)}{BR(B \rightarrow K^{(*)}e^+e^-)}$$



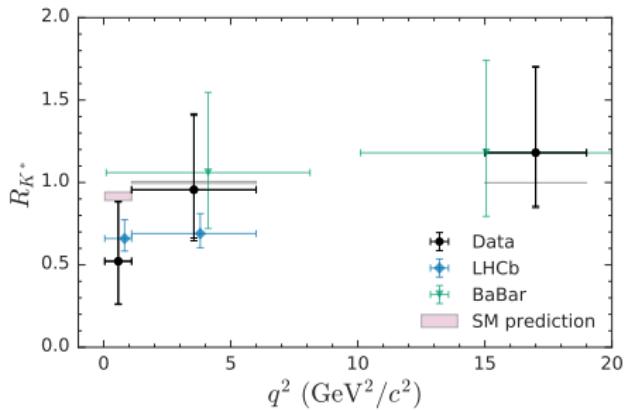
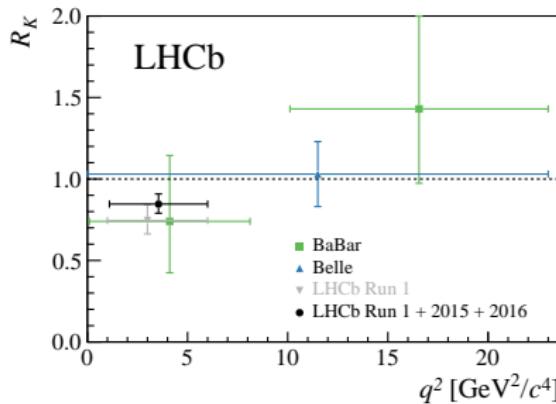
# Hints for LFU violation in $b \rightarrow s \ell^+ \ell^-$ decays

## New results at Moriond 2019

- ▶ Updated measurement of  $R_K$  by LHCb
- ▶ New measurement of  $R_{K^*}$  by Belle

LHCb, arXiv:1903.09252

Belle, arXiv:1904.02440



# Hints for LFU violation in $b \rightarrow c \ell \nu$ decays

Measurements of LFU ratios  $R_D$  and  $R_{D^*}$  by BaBar, Belle, and LHCb showed combined deviation from SM by  $3.8\sigma$ .

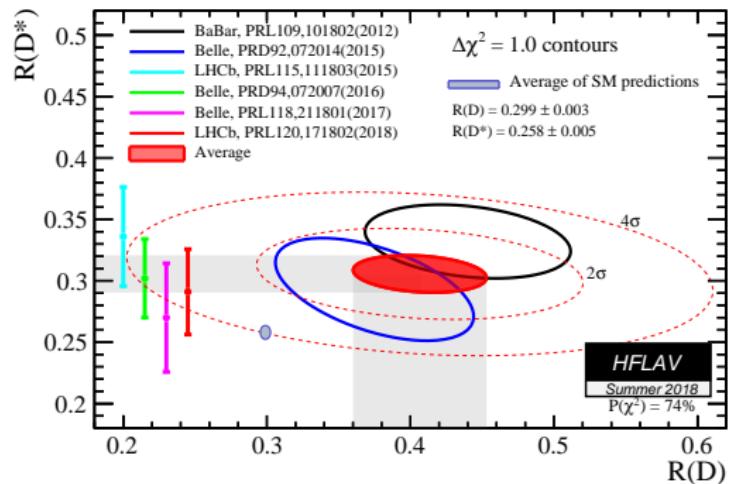
BaBar, arXiv:1205.5442, arXiv:1303.0571

LHCb, arXiv:1506.08614, arXiv:1708.08856

Belle, arXiv:1507.03233, arXiv:1607.07923, arXiv:1612.00529

$$R_{D^{(*)}} = \frac{BR(B \rightarrow D^{(*)}\tau\nu)}{BR(B \rightarrow D^{(*)}\ell\nu)}$$

$$\ell \in \{e, \mu\}$$



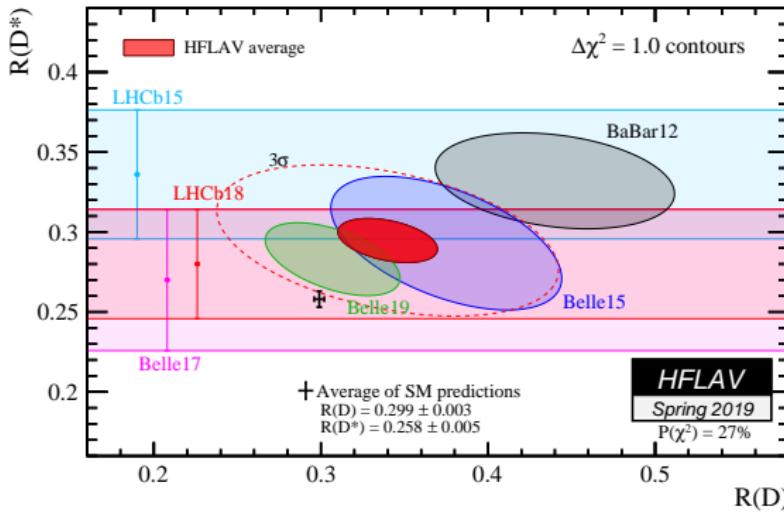
HFLAV, arXiv:1612.07233

# Hints for LFU violation in $b \rightarrow c \ell \nu$ decays

## New results at Moriond 2019

- Updated measurements of  $R_D$  and  $R_{D^*}$  by Belle

Belle, arXiv:1904.08794



HFLAV, hflav.web.cern.ch

see talks by Adam Morris and Dawid Gerstel and discussion session on  $R_D$  and  $R_{D^*}$

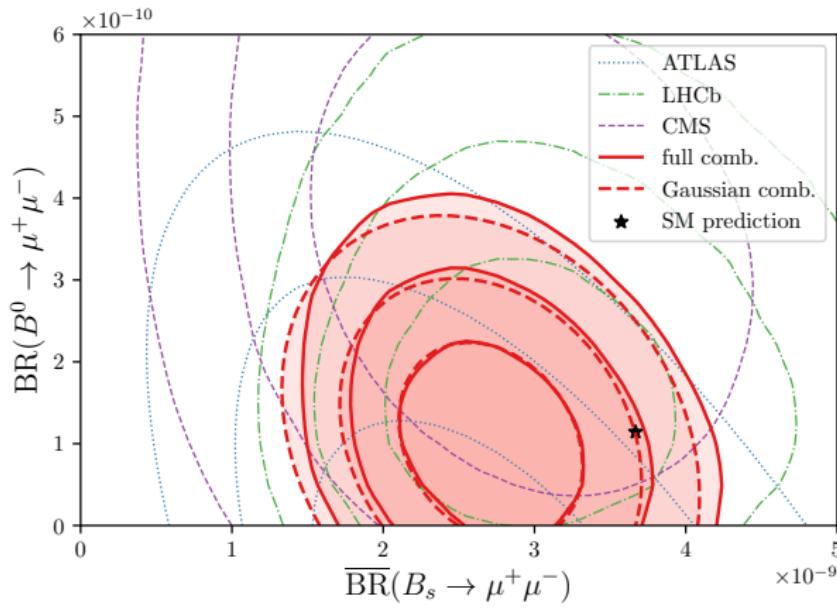
# Combination of $B_{s,d} \rightarrow \mu^+ \mu^-$ measurements

Measurements of  $\text{BR}(B_{s,d} \rightarrow \mu^+ \mu^-)$  by LHCb, CMS, and ATLAS show combined deviation from SM by about  $2\sigma$ .

LHCb, arXiv:1703.05747

CMS, arXiv:1307.5025

ATLAS, arXiv:1812.03017



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# Setup

- ▶ Global likelihood from **smelli** python package for comparing theory predictions to experimental data
- Aebischer, Kumar, PS, Straub, arXiv:1810.07698
- ▶ Quantify agreement between theory and experiment by likelihood  $L$ ,  $\Delta\chi^2$ , and pull

$$\text{pull}_{1D} = 1\sigma \cdot \sqrt{\Delta\chi^2}, \quad \text{where } -\frac{1}{2}\Delta\chi^2 = \ln L(\vec{0}) - \ln L(\vec{C}_{\text{best fit}}).$$

$$\text{pull}_{2D} = 1\sigma, 2\sigma, 3\sigma, \dots \quad \text{for } \Delta\chi^2 \approx 2.3, 6.2, 11.8, \dots$$

- ▶ New physics scenarios in effective field theories:
  - ▶ **Weak Effective Theory (WET)** at scale  $m_b$
  - ▶ **Standard Model Effective Field Theory (SMEFT)** at scale 2 TeV
- ▶ **Simplified leptoquark model** matched to SMEFT at 2 TeV

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# **$b \rightarrow s\ell\ell$ in the weak effective theory**

- Effective Hamiltonian at scale  $m_b$ :  $\mathcal{H}_{\text{eff}}^{bs\ell\ell} = \mathcal{H}_{\text{eff, SM}}^{bs\ell\ell} + \mathcal{H}_{\text{eff, NP}}^{bs\ell\ell}$

$$\mathcal{H}_{\text{eff, NP}}^{bs\ell\ell} = -\mathcal{N} \sum_{\ell=e,\mu} \sum_{i=9,10,S,P} (C_i^{bs\ell\ell} O_i^{bs\ell\ell} + C_i'^{bs\ell\ell} O_i'^{bs\ell\ell}) + \text{h.c.}$$

- Operators considered here ( $\ell = e, \mu$ )

$$\begin{aligned} O_9^{bs\ell\ell} &= (\bar{s}\gamma_\mu P_L b)(\bar{\ell}\gamma^\mu \ell), & O_9'^{bs\ell\ell} &= (\bar{s}\gamma_\mu P_R b)(\bar{\ell}\gamma^\mu \ell), \\ O_{10}^{bs\ell\ell} &= (\bar{s}\gamma_\mu P_L b)(\bar{\ell}\gamma^\mu \gamma_5 \ell), & O_{10}'^{bs\ell\ell} &= (\bar{s}\gamma_\mu P_R b)(\bar{\ell}\gamma^\mu \gamma_5 \ell), \\ O_S^{bs\ell\ell} &= m_b(\bar{s}P_R b)(\bar{\ell}\ell), & O_S'^{bs\ell\ell} &= m_b(\bar{s}P_L b)(\bar{\ell}\ell), \\ O_P^{bs\ell\ell} &= m_b(\bar{s}P_R b)(\bar{\ell}\gamma_5 \ell), & O_P'^{bs\ell\ell} &= m_b(\bar{s}P_L b)(\bar{\ell}\gamma_5 \ell). \end{aligned}$$

- Not considered here

- Dipole operators: strongly constrained by radiative decays. e.g. [arXiv:1608.02556]
- Four quark operators: dominant effect from RG running above  $m_B$ . Jäger, Leslie, Kirk, Lenz [arXiv:1701.09183]

# Scenarios with a single Wilson coefficients

Coefficient	type	best fit	$1\sigma$	$\text{pull}_{1D} = \sqrt{\Delta\chi^2}$
$C_9^{bs\mu\mu}$	$L \otimes V$	-0.97	[-1.12, -0.81]	<b>5.9<math>\sigma</math></b>
$C_9'^{bs\mu\mu}$	$R \otimes V$	+0.14	[-0.03, +0.32]	0.8 $\sigma$
$C_{10}^{bs\mu\mu}$	$L \otimes A$	+0.75	[+0.62, +0.89]	<b>5.7<math>\sigma</math></b>
$C_{10}'^{bs\mu\mu}$	$R \otimes A$	-0.24	[-0.36, -0.12]	2.0 $\sigma$
$C_9^{bs\mu\mu} = C_{10}^{bs\mu\mu}$	$L \otimes R$	+0.20	[+0.06, +0.36]	1.4 $\sigma$
$C_9^{bs\mu\mu} = -C_{10}^{bs\mu\mu}$	$L \otimes L$	-0.53	[-0.61, -0.45]	<b>6.6<math>\sigma</math></b>

Only small pull for

- ▶ Coefficients with  $\ell = e$  (cannot explain  $b \rightarrow s\mu\mu$  anomaly)
- ▶ Scalar coefficients (can only reduce tension in  $B_s \rightarrow \mu\mu$ )

see also similar fits by other groups:

Algueró et al., arXiv:1903.09578

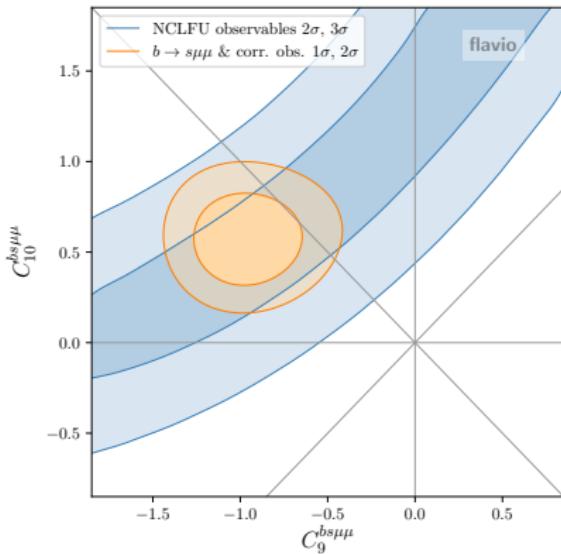
Ciuchini et al., arXiv:1903.09632

Datta et al., arXiv:1903.10086

Kowalska et al., arXiv:1903.10932

Arbey et al., arXiv:1904.08399

# Scenarios with two Wilson coefficients

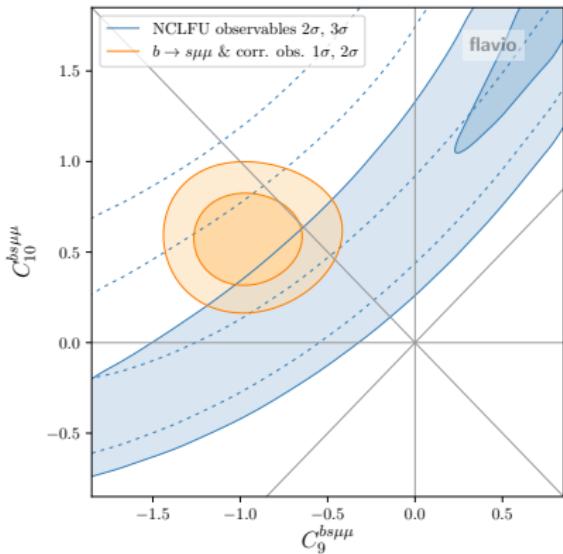


## ► Before Moriond 2019:

Very good agreement between fits to  
 $b \rightarrow s \mu \mu$  observables and  $R_K$  &  $R_{K^*}$

WET at 4.8 GeV

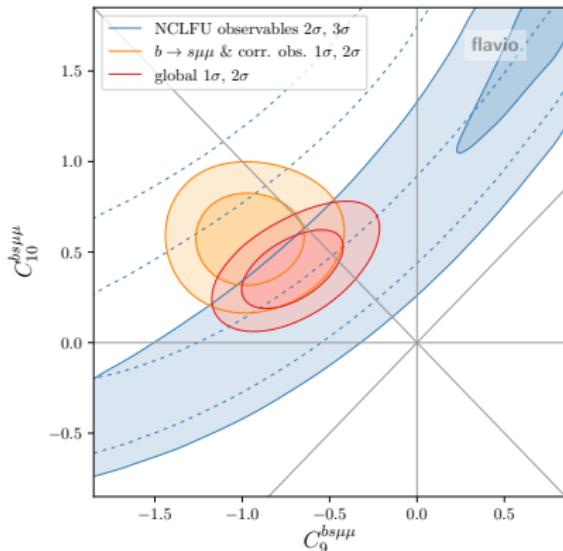
## Scenarios with two Wilson coefficients






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# Scenarios with two Wilson coefficients



- ▶ **Before Moriond 2019:**  
Very good agreement between fits to  $b \rightarrow s\mu\mu$  observables and  $R_K$  &  $R_{K^*}$
- ▶ **After Moriond 2019:**  
Updated  $R_K$  measurement by LHCb and new  $R_{K^*}$  measurement by Belle closer to SM value    [LHCb, arXiv:1903.09252](#)  
[Belle, arXiv:1904.02440](#)
- Tension between fits to  $R_K$  &  $R_{K^*}$  and  $b \rightarrow s\mu\mu$  observables in  $C_9$  direction
- ▶ **Global likelihood:**  
Contribution to purely left-handed  
 $C_9^{bs\mu\mu} = -C_{10}^{bs\mu\mu}$  yields very good fit to experimental data

# Scenarios with two Wilson coefficients

- ▶ **LFU contribution** only affects  $b \rightarrow s\mu\mu$  observables
- ▶ Tension between fits to  $b \rightarrow s\mu\mu$  observables and  $R_K$  &  $R_{K^*}$  could be reduced by **LFU** contribution to  $C_9$
- ▶ Perform two-parameter fit in space of  $C_9^{\text{univ.}}$  and  $\Delta C_9^{bs\mu\mu} = -C_{10}^{bs\mu\mu}$ :

$$C_9^{bsee} = C_9^{\text{univ.}}$$

$$C_{10}^{bsee} = 0$$

$$C_9^{bs\mu\mu} = C_9^{\text{univ.}} + \Delta C_9^{bs\mu\mu}$$

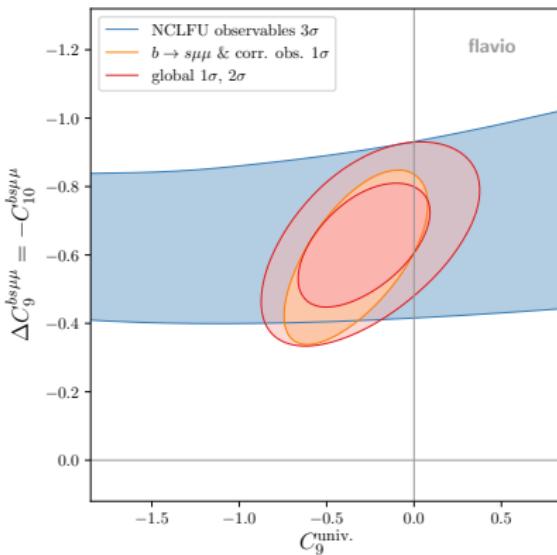
$$C_{10}^{bs\mu\mu} = -\Delta C_9^{bs\mu\mu}$$

$$C_9^{bst\tau\tau} = C_9^{\text{univ.}}$$

$$C_{10}^{bst\tau\tau} = 0$$

scenario first considered in  
Algúeró et al., arXiv:1809.08447

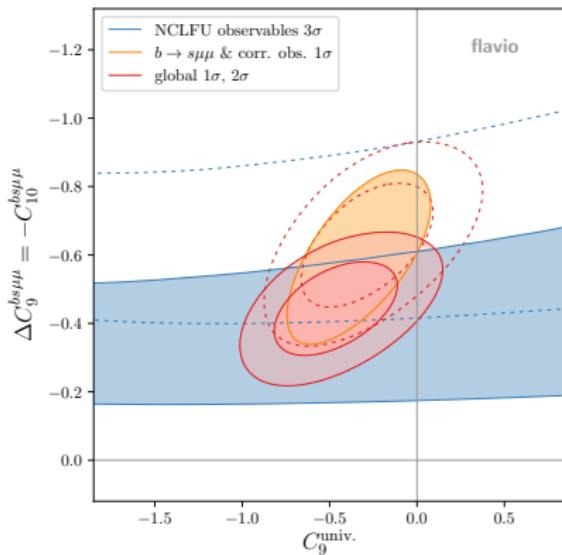
# Scenarios with two Wilson coefficients



- **Before Moriond 2019:**  
Fit compatible with  $C_9^{\text{univ.}} = 0$  and only contribution to  $C_9^{\text{bs}\mu\mu} = -C_{10}^{\text{bs}\mu\mu}$

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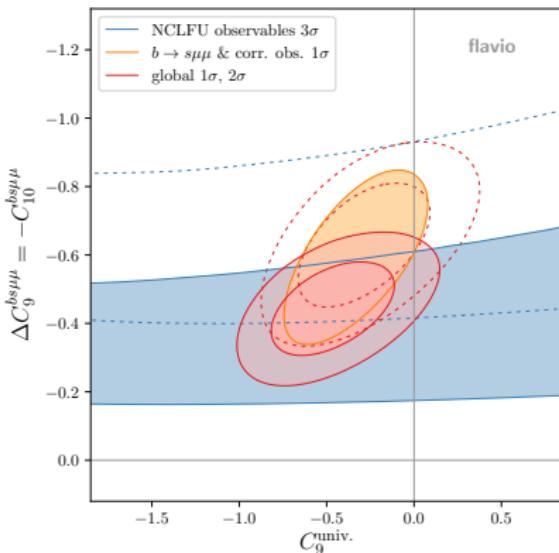
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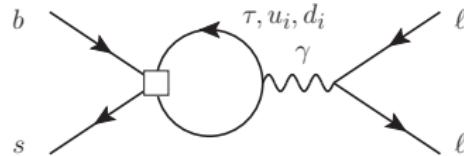
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- ▶ **After Moriond 2019:**  
Preference for **non-zero  $C_9^{\text{univ.}}$**

# Scenarios with two Wilson coefficients



WET at 4.8 GeV

- ▶ **Before Moriond 2019:**  
Fit compatible with  $C_9^{\text{univ.}} = 0$  and only contribution to  $C_9^{bs\mu\mu} = -C_{10}^{bs\mu\mu}$
- ▶ **After Moriond 2019:**  
Preference for **non-zero  $C_9^{\text{univ.}}$**
- ▶  $C_9^{\text{univ.}}$  can arise from RG effects:



Bobeth, Haisch, arXiv:1109.1826  
 Crivellin, Greub, Müller, Saturnino, arXiv:1807.02068

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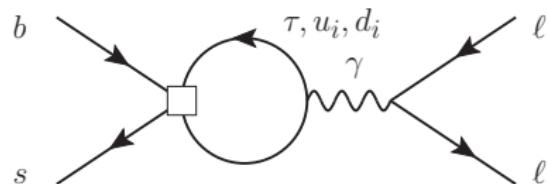
- $b \rightarrow s\ell\ell$  in the weak effective theory
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## 3 Conclusions

# The global picture in the SMEFT

RG effects require scale separation

- ▶ Consider **SMEFT at 2 TeV**



Possible operators:

- ▶  $[O_{lq}^{(3)}]_{3323} = (\bar{l}_3 \gamma_\mu \tau^a l_3)(\bar{q}_2 \gamma^\mu \tau^a q_3)$ :

Can also **explain  $R_{D^{(*)}}$  anomalies!**

- ▶  $[O_{lq}^{(1)}]_{3323} = (\bar{l}_3 \gamma_\mu l_3)(\bar{q}_2 \gamma^\mu q_3)$ :

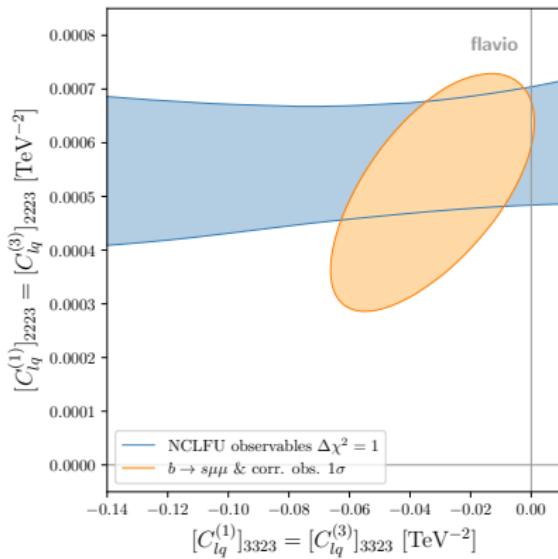
Strong constraints from  $B \rightarrow K \nu \nu$  require  $[C_{lq}^{(1)}]_{3323} \approx [C_{lq}^{(3)}]_{3323}$

Buras et al., arXiv:1409.4557

- ▶  $[O_{qe}]_{2333} = (\bar{q}_2 \gamma_\mu q_3)(\bar{e}_3 \gamma^\mu e_3)$  cannot explain  $R_{D^{(*)}}$

- ▶ Four-quark operators cannot explain  $R_{D^{(*)}}$ , models yielding large enough contributions already in tension with data

# The global picture in the SMEFT

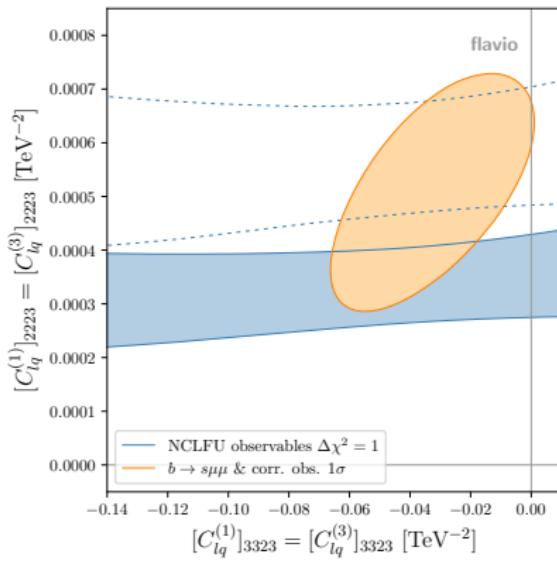


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$$[C_{lq}^{(1)}]_{3323} = [C_{lq}^{(3)}]_{3323} \Rightarrow C_9^{\text{univ.}} \quad (\text{RG effect})$$

$$[C_{lq}^{(1)}]_{2223} = [C_{lq}^{(3)}]_{2223} \Rightarrow \Delta C_9^{bs\mu\mu} = -C_{10}^{bs\mu\mu}$$

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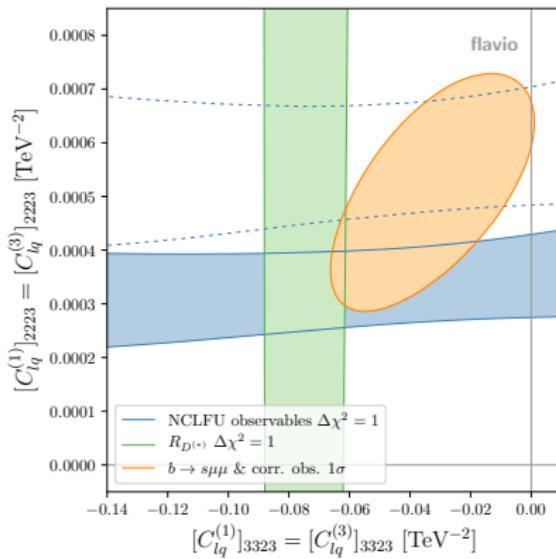


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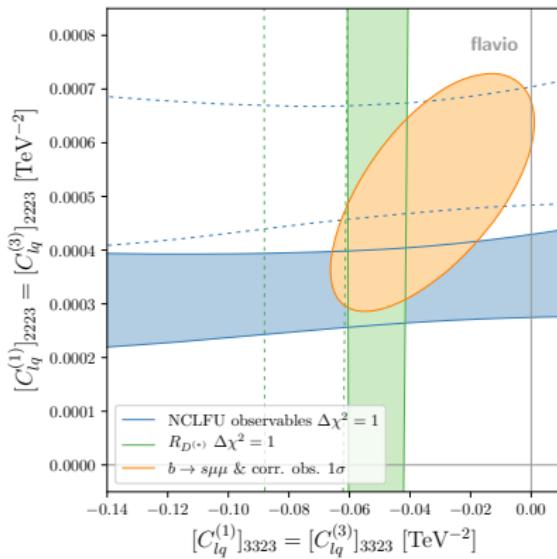


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Agreement with combined  $R_{K^{(*)}}$  and  
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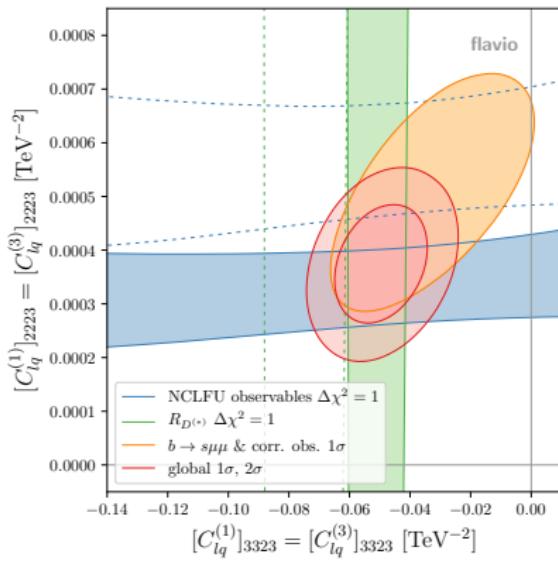


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# Predictions from global likelihood in SMEFT scenario

Observable	$1\sigma$	SM
$R_{K^*}^{[0.045, 1.1]}$	$0.88^{+0.01}_{-0.01}$	$0.926 \pm 0.004$
$R_{K^*}^{[1.1, 6.0]}$	$0.81^{+0.04}_{-0.04}$	$0.9964 \pm 0.0006$
$R_{K^*}^{[0.1, 8.0]}$	$0.83^{+0.04}_{-0.03}$	$0.995 \pm 0.002$
$R_{K^*}^{[15, 19]}$	$0.79^{+0.04}_{-0.04}$	$0.99807 \pm 0.00004$
$R_K^{[1.0, 6.0]}$	$0.80^{+0.04}_{-0.04}$	$1.0008 \pm 0.0003$
$R_\phi^{[1.0, 6.0]}$	$0.81^{+0.04}_{-0.04}$	$0.9970 \pm 0.0003$
$\langle P'_5 \rangle^{[4.0, 6.0]}$	$-0.58^{+0.13}_{-0.12}$	$-0.763 \pm 0.072$
$R_D$	$0.34^{+0.01}_{-0.01}$	$0.303 \pm 0.006$
$R_{D^*}$	$0.29^{+0.01}_{-0.01}$	$0.255 \pm 0.004$
$\overline{\text{BR}}(B_s \rightarrow \mu^+ \mu^-)$	$2.98^{+0.20}_{-0.19} \times 10^{-9}$	$(3.67 \pm 0.16) \times 10^{-9}$
$\text{BR}(B^\pm \rightarrow K^\pm \tau^+ \tau^-)$	$3.05^{+1.78}_{-1.06} \times 10^{-5}$	$(1.66 \pm 0.19) \times 10^{-7}$
$\overline{\text{BR}}(B_s \rightarrow \tau^+ \tau^-)$	$1.41^{+0.80}_{-0.47} \times 10^{-4}$	$(7.78 \pm 0.33) \times 10^{-7}$

see parallel session on  $b \rightarrow s \tau^+ \tau^-$ , talk by Olcyr Sumensari

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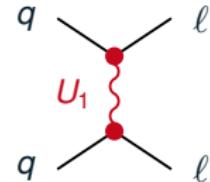
## 3 Conclusions

# Simplified $U_1$ -leptoquark model

see talk by Jonathan Kriewald

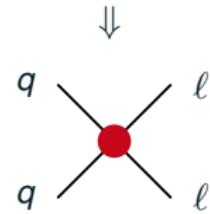
- $U_1$  vector leptoquark  $(3, 1)_{2/3}$  couples quarks and leptons

$$\mathcal{L}_{U_1} \supset g_{lq}^{ij} (\bar{q}^i \gamma^\mu l^j) U_\mu + \text{h.c.}$$



- Generates **semi-leptonic operators at tree-level**

$$[C_{lq}^{(1)}]_{ijkl} = [C_{lq}^{(3)}]_{ijkl} = -\frac{g_{lq}^{jk} g_{lq}^{il*}}{2M_U^2}$$

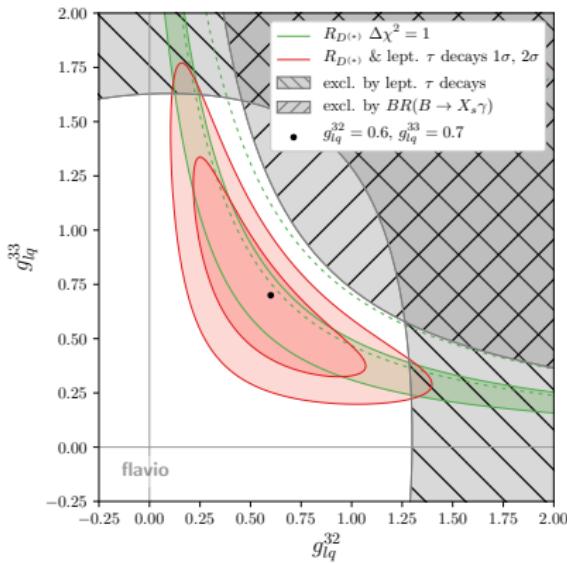


- And **dipole operators at one-loop**, e.g.

$$[O_{dV}]_{ij} = (\bar{q}_i \sigma^{\mu\nu} V_{\mu\nu} q_j) \varphi, \quad V \in \{W, B, G\}:$$

$$[C_{dV}]_{23} = \kappa_V \frac{Y_b}{16\pi^2} \sum_i \frac{g_{lq}^{i2} g_{lq}^{i3*}}{M_U^2}, \quad \kappa_W = \frac{g}{6}, \quad \kappa_B = \frac{-4g'}{9}, \quad \kappa_V = \frac{-5g_s}{12}$$

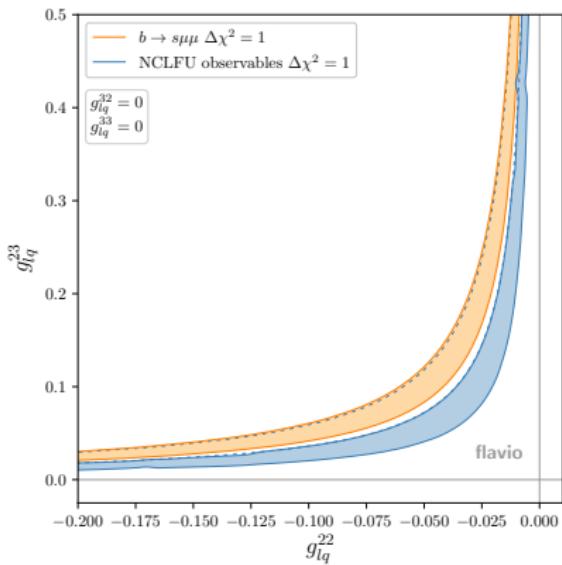
# Simplified $U_1$ -leptoquark model



- ▶  $R_{D(*)}$  mostly depends on **tauonic couplings  $g_{lq}^{32}, g_{lq}^{33}$**
- ▶ Dipole operators contribute to  $\text{BR}(B \rightarrow X_s \gamma)$
- ▶ RG running contributes to **leptonic  $\tau$  decays**
- ▶ Well defined allowed region for explaining  $R_{D(*)}$ , select **benchmark point**

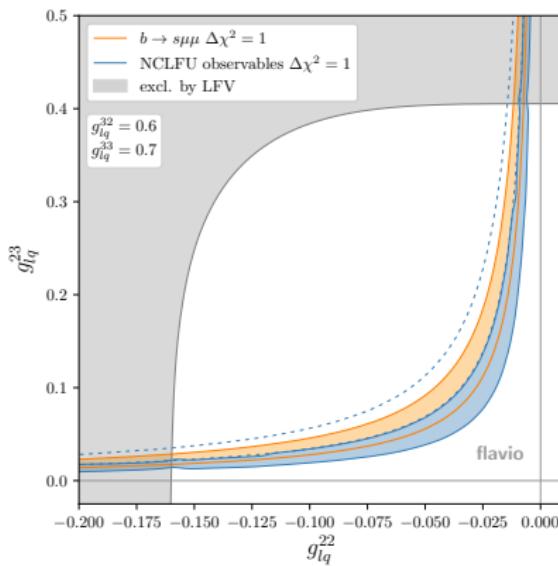
$$g_{lq}^{32} = 0.6, \quad g_{lq}^{33} = 0.7$$

# Simplified $U_1$ -leptoquark model



- ▶  $R_{K(*)}$  can be explained by **muonic couplings**  $g_{lq}^{22}, g_{lq}^{23}$
- ▶ **Vanishing tauonic couplings:** Tension between fits to  $R_{K(*)}$  and  $b \rightarrow s \mu \mu$  observables after Moriond 2019

# Simplified $U_1$ -leptoquark model



- ▶  $R_{K(*)}$  can be explained by **muonic couplings**  $g_{lq}^{22}, g_{lq}^{33}$
  - ▶ **Vanishing tauonic couplings:** Tension between fits to  $R_{K(*)}$  and  $b \rightarrow s\mu\mu$  observables after Moriond 2019
  - ▶ Benchmark point explaining  $R_{D(*)}$ ,
- $$g_{lq}^{32} = 0.6, \quad g_{lq}^{33} = 0.7,$$
- implies non-zero  $C_9^{\text{univ.}}$ ,  $R_{K(*)}$  and  $b \rightarrow s\mu\mu$  in good agreement after Moriond 2019
- ▶ Constraint from **LFV observables**

# Outline

## 1 B-decay discrepancies

## 2 New physics interpretations

- $b \rightarrow s\ell\ell$  in the weak effective theory
- The global picture in the Standard Model effective field theory (SMEFT)
- Simplified leptoquark model

## 3 Conclusions

# Conclusions

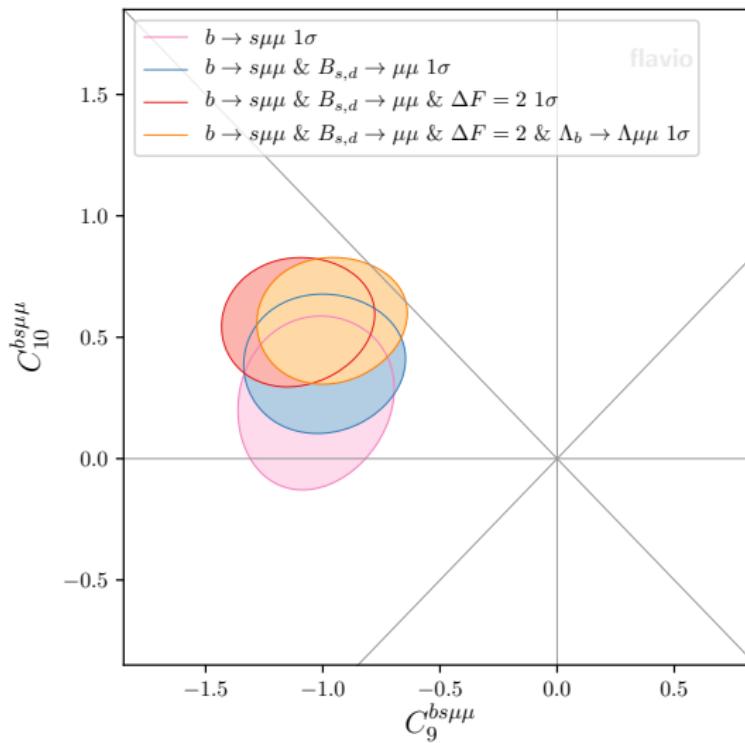
- ▶ New and updated measurements of  $R_{K^{(*)}}$  and  $R_{D^{(*)}}$  as well as  $B_s \rightarrow \mu\mu$ .
- ▶ New physics in the single muonic Wilson coefficients  $C_9^{bs\mu\mu}$ ,  $C_{10}^{bs\mu\mu}$ , and  $C_9^{bs\mu\mu} = -C_{10}^{bs\mu\mu}$  gives a clearly better fit to data than the SM (pull  $\approx 6\sigma$ ).
- ▶ Slight tension between  $R_{K^{(*)}}$  fit and  $b \rightarrow s\mu\mu$  fit with only muonic Wilson coefficients can be reduced by lepton flavor universal  $C_9^{\text{univ}}$ .
- ▶ Lepton flavor universal  $C_9^{\text{univ}}$  can be generated through RG effects from semi-tauonic Wilson coefficients that can explain  $R_{D^{(*)}}$ .
- ▶  $U_1$ -leptoquark can generate these semi-tauonic Wilson coefficients in addition to semi-muonic ones that explain  $R_{K^{(*)}}$ .

# Backup slides

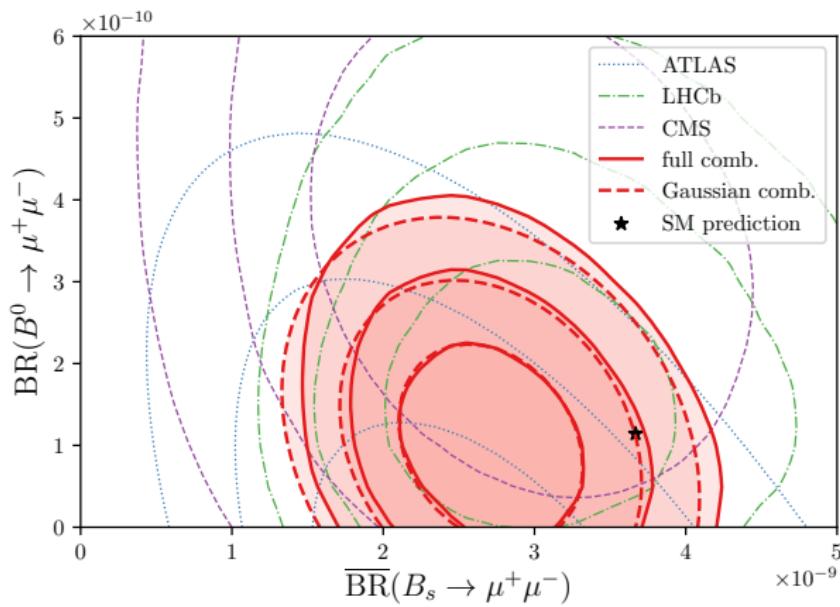
# New physics in individual Wilson coefficients

Coefficient	type	best fit	$1\sigma$	pull
$C_9^{bs\mu\mu}$	$L \otimes V$	-0.97	[-1.12, -0.81]	<b>5.9<math>\sigma</math></b>
$C_9'^{bs\mu\mu}$	$R \otimes V$	+0.14	[-0.03, +0.32]	0.8 $\sigma$
$C_{10}^{bs\mu\mu}$	$L \otimes A$	+0.75	[+0.62, +0.89]	<b>5.7<math>\sigma</math></b>
$C_{10}'^{bs\mu\mu}$	$R \otimes A$	-0.24	[-0.36, -0.12]	2.0 $\sigma$
$C_9^{bs\mu\mu} = C_{10}^{bs\mu\mu}$	$L \otimes R$	+0.20	[+0.06, +0.36]	1.4 $\sigma$
$C_9^{bs\mu\mu} = -C_{10}^{bs\mu\mu}$	$L \otimes L$	-0.53	[-0.61, -0.45]	<b>6.6<math>\sigma</math></b>
$C_9^{bsee}$	$L \otimes V$	+0.93	[+0.66, +1.17]	3.5 $\sigma$
$C_9'^{bsee}$	$R \otimes V$	+0.39	[+0.05, +0.65]	1.2 $\sigma$
$C_{10}^{bsee}$	$L \otimes A$	-0.83	[-1.05, -0.60]	3.6 $\sigma$
$C_{10}'^{bsee}$	$R \otimes A$	-0.27	[-0.57, -0.02]	1.1 $\sigma$
$C_9^{bsee} = C_{10}^{bsee}$	$L \otimes R$	-1.49	[-1.79, -1.18]	3.2 $\sigma$
$C_9^{bsee} = -C_{10}^{bsee}$	$L \otimes L$	+0.47	[+0.33, +0.59]	3.5 $\sigma$
$(C_S^{bs\mu\mu} = -C_P^{bs\mu\mu}) \times \text{GeV}$	$\bar{L}R \otimes \bar{R}L$	-0.006	[-0.009, -0.003]	2.8 $\sigma$
$(C_S'^{bs\mu\mu} = C_P'^{bs\mu\mu}) \times \text{GeV}$	$\bar{R}L \otimes \bar{L}R$	-0.006	[-0.009, -0.003]	2.8 $\sigma$

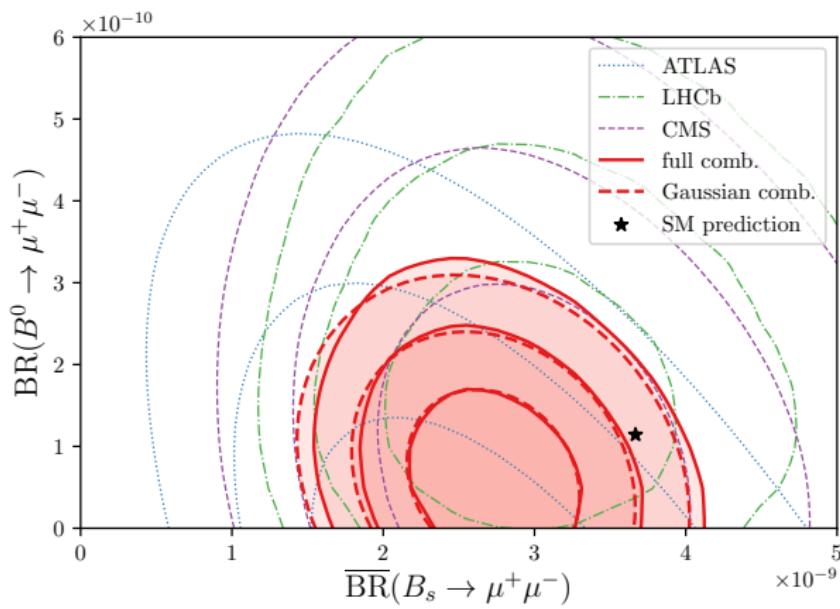
# $C_9$ vs. $C_9 = -C_{10}$



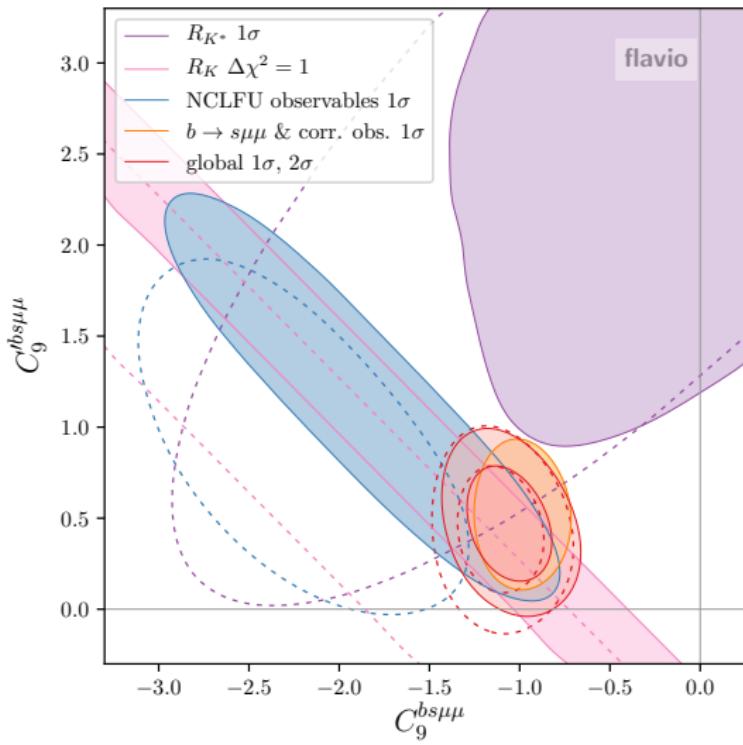
# Combination of $B_{s,d} \rightarrow \mu^+ \mu^-$ measurements



# Combination of $B_{s,d} \rightarrow \mu^+ \mu^-$ measurements



# New physics in right-handed quark current



# Direct constraints on $U_1$ leptoquark

