Exotic hadrons in heavy ion collisions

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Aim

- A brief overview on the state-of-art of exotic states
- The exotics in a HIC: dependence on their interpretation
- Molecular, tetraquark and triangle singularities interpretations in HICs?

Case study of some of the most prominent exotic states

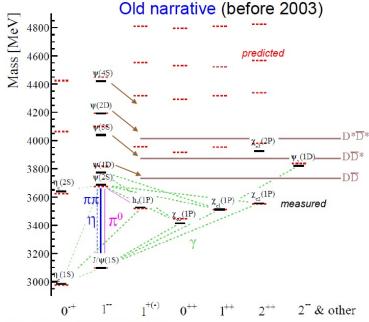


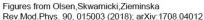
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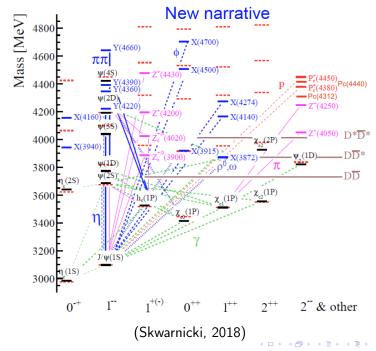
Case study of some of the most prominent exotic states











- ullet Proliferation of new states in the car c, bar b and ccq sectors
- Properties do not match standard quark-model predictions
- Decay properties require > 3 valence quarks

Ex.1: Z's are manifestly 4-quark states (J/ψ) not produced from $|0\rangle$

$$Z_c^+(3900) \to J/\psi \pi^+$$

Ex.2: isospin violation (if $X(3872) \sim c\bar{c}$: decay highly suppressed)

$$X(3872) \to \rho^0 J/\psi \to \pi^+ \pi^- J/\psi$$

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- 17 confirmed with complete assignment by PDG
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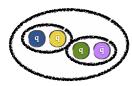
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Intepretations for composition and binding mechanisms?

Hadron Molecules



Hybrids



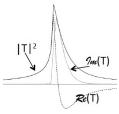
Glueballs



Tetraquarks



Kinematical effects (TS's)





(luciano.abreu@ufba.br)

The heavy exotics collection

Theoretical perspective

A compelling and unified understanding has not yet emerged

- No single theoretical framework explains all exotic candidates
- Candidates: different interpretations
- In some cases, properties (masses, decay widths) are well explained by different models or quantum-mechanical superposition of them
- Necessity of more observables to distinguish its internal structure
- Let us focus on some emblematic states



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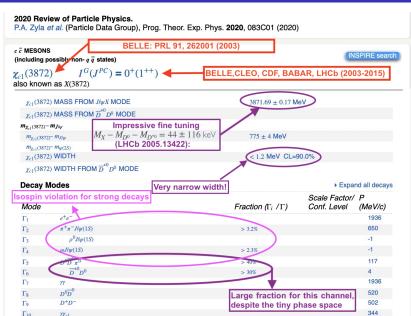
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X(3872): features



X(3872) story

74 years ago

```
K-mesons discovered -- associated production – strangeness – SU(3) -- model

Dec. 1947 ← 16 years → Jan 1964
```

18 years ago

(Adapted from S. Olsen, SCGP Workshop on Exotic Hadrons and Flavor Physics, May 2018)



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Recent dispute on Prompt production of X(3872)

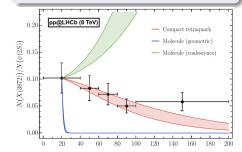
Esposito et al., 2006.15044

Comover interaction model:

$$\tau \frac{N_{\mathcal{Q}}}{d\tau} = -\langle v\sigma \rangle_{\mathcal{Q}} \rho_c N_{\mathcal{Q}};$$

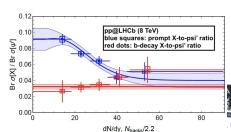
$$\begin{split} \langle v\sigma\rangle_{4q} \sim \pi r_{4q}^2 \simeq 11.6\,\mathrm{mb};\\ \langle v\sigma\rangle_{Mol} \sim \pi r_{Mol}^2 \simeq 1197\,\mathrm{mb} \end{split}$$

ullet Findings \Rightarrow tetraquark nature



Braaten et al., PRD (2021); 2012.13499

- $\langle v\sigma \rangle_{Mol}$: probability-weighted sum of $\langle v\sigma \rangle (\pi D^{(*)})$
- Insensitive to $E_b^{(X)}$
- $f_{out,Q}^{(prompt)}$: out of reach of comoving pions
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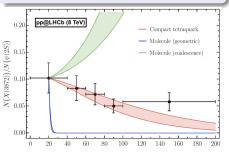
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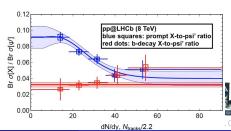
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Promising alternative: exotics in HICs

pp collisions

- Esposito et al., 2006.15044 \rightarrow Geometrical Cross sections: $\langle v\sigma\rangle_Q \sim \pi r_Q^2 (1-E_Q^{Thr}/E_\pi)^n$
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Relevance of hadronic medium interactions

Artoisenet and Braaten (PRD, 2011), Esposito et al. (JMP, 2013), Cho and Lee (PRD, 2013), LMA et al. (PRD, 2014), Guerrieri et al. (2014), ...

HICs

- End of the QGP phase: XYZ states interact with hadronic medium
- Absorption by comoving mesons or production from heavy mesons $(D^{(*)}\bar{D}^{(*)} \to X\pi, \rho, ...)$
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Breaking news: first evidence for X(3872) in HICs!

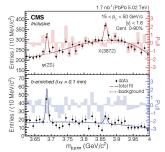
CMS-LHC, arXiv:2102.13048

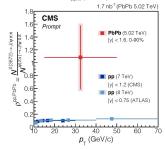
• Prompt X(3872)-production in PbPb collisions, $\sqrt{s} = 5.02$ TeV

$$X(3872)$$
 \rightarrow $J/\psi \pi^+ \pi^-$
 \rightarrow $\mu^+ \mu^- \pi^+ \pi^-$

$\rho^{(PbPb)} \simeq 10 \, \rho^{(pp)}$

Unique experimental input of the X(3872) production mechanism and its nature







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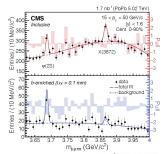
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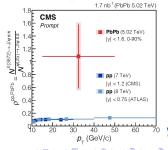
$$X$$
(3872) \rightarrow $J/\psi \pi^+ \pi^-$
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$$ho^{(PbPb)} = rac{N_{X(3872)}}{N_{\psi(25)}} = 1.08 \pm 0.9 \pm 0.52$$

$$\rho^{(PbPb)} \simeq 10 \, \rho^{(pp)}$$

Unique experimental input of the X(3872) production mechanism and its nature







Our contributions: LMA, Kamchandani, Nielsen, Navarra, Torres [PRD, PLB, ... (2014-2019)]

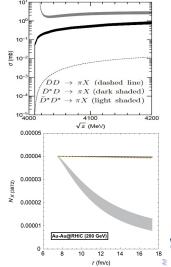
Hadronic effects on $N_{X(3872)}(\tau)$

- Use of Eff. Lagrangians based on relevant symmetries
- Inclusion of $g_{XD^{\pm}D^{*\mp}}$
- Anomalous: $\mathcal{L}_{\pi D^* \bar{D}^*}$, $\mathcal{L}_{XD^* \bar{D}^*}$
- Internal structure: coalescence

Bjorken picture - Rate equation:

$$\begin{array}{ll} \frac{dN_X(\tau)}{d\tau} & = & \sum_{\varphi} \left[\langle \sigma_{\bar{M}M' \to \varphi X} v_{\bar{M}M'} \rangle n_{\bar{M}}(\tau) N_{M'}(\tau) \\ \\ & - \langle \sigma_{\varphi X \to \bar{M}M'} v_{\varphi X} \rangle n_{\varphi}(\tau) N_X(\tau) \right]. \end{array}$$

 $N_{X(3872)}^{(4q)}$ reduced by a factor of 4



Exotic states with I=1

state	mass MeV	width MeV	QQ decay mode	phase space MeV	nearby threshold	ΔE MeV
X(3872) $Z_b(10610)$ $Z_b(10650)$ $Z_c(3900)$ $Z_c(4020)$ ×	3872 10608 10651 3900 4020	< 1.2 21 10 24 - 46 8 - 25	$J/\psi \pi^+\pi^ \Upsilon \pi$ $\Upsilon \pi$ $J/\psi \pi$ $J/\psi \pi$	495 1008 1051 663 783	ŪD* BB* B*B* ŪD* Ū*D* ŪD BB	< 1 2 ± 2 2 ± 2 24 6
1078	50					
1076 1066 N 1066 1056	50	X _b ² ∘ -X _b I=0	T	0650)	B*I BB	
410 400 New 400 390	50 - 00 X	(4014)	Z _c (402	•	D*I)*
390 386		(3872)	Z _c ⁺ (3	900)	DD	•
380	00 - 0	bserve	d opr	edict	ed	

$$Z_c^{\pm,0}(3900)$$
, $Z_c^{\pm,0}(4020)$

•
$$(D\bar{D}^*)^{\pm,0}, (D^*\bar{D}^*)^{\pm,0}$$

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- $I(J^{P(C)}) = 1(1^{+(-)})$
- CHARGED: no confusion with quarkonia $Q\bar{Q}$ states!
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$Z_c^{\pm,0}$ (3900), $Z_c^{\pm,0}$ (4020)

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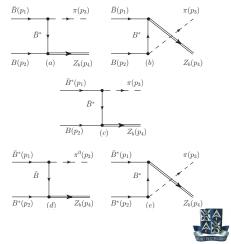
Bottomonium sector: $Z_b(10610)$ and $Z_b'(10650)$

LMA, Kamchandani, Torres, Nielsen, Navarra, Vasconcellos; (PRD, 2017; EPJ C, 2018, ...)

- $N_{Z_{b}^{(\prime)}}(au)$ in HIC environment
- Rate equation (Bjorken picture): absorption and regeneration effects in hadron gas phase
- Inclusion of the coupling:

$$Z_b^+(B^+\bar{B}^{*0}+\bar{B}^0B^{*+})$$

Inclusion of the term $-\Gamma_{Z_b}N_{Z_b}(au)$ $(\Gamma_{Z_b}\approx 20~{
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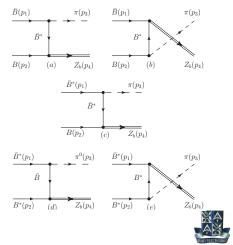
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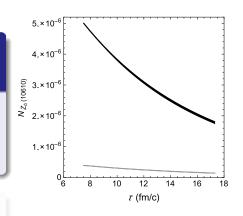
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Time evolution of the $Z_b(10610)^{\pm,0}$ abundance

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- Suppression by a factor two of the final yield
- Smaller suppression that than the one found for X(3872)
- Role played by the medium: "filtering" of the particles produced at the end of QGP



LHC (
$$\sqrt{s_{NN}} = 2.76 \text{ TeV}$$
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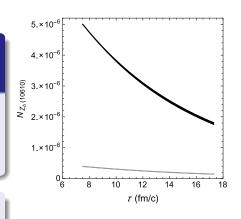
(Dark and light shaded bands: statistical and coalescence models (PRD, 2017; EPJ C, 2018,



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$X_J(2900)$: first fully open-charm exotic hadrons

LHCb, PRL 125, 242001 (2020)

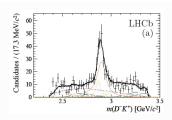
• Invariant mass spectrum of the $B^+ \to D^+ D^- K^+$ decay:

$$0^+: M = 2866 \text{MeV}, \Gamma = 57 \text{MeV}$$

 $1^-: M = 2904 \text{MeV}, \Gamma = 110 \text{MeV}$

Minimum valence quark contents: c̄sud

First observed unconventional hadrons which are fully open charm tetraquarks



- HIC: how they are affected by the medium during the expansion?
- Possible interpretations?
- Hadronic molecular states
 (J = 0 as a S-wave and J = 1 as a P-wave)?
- Tetraquark states?

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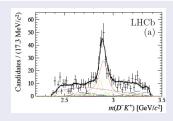
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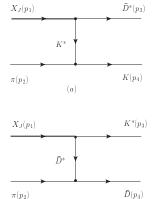


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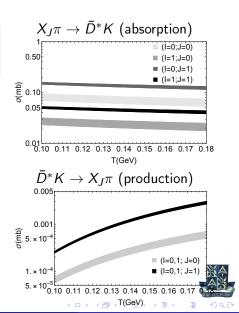
$X_J(2900)$: hadronic interactions in HICs

LMA, PRD103, 036013(2021)

• Analysis of the processes: $X_J\pi \leftrightarrow \bar{D}^*K, K^*\bar{D}$



(b)



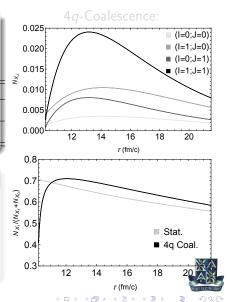
X_J(2900): Time evolution [LMA, PRD**103**, 036013(2021)]

- 4q coal.: initial yields $(\tau \sim \tau_C)$
- Hadron coal.: at $au \sim au_F$

State	$N_{X_J}^{0(Stat)}(au_H)$	$N_{X_J}^{0(4q)}(\tau_C)$	$N_{X_J}^{(Mol)}(au_F)$
J = 0, I = 0	6.9×10^{-3}	1.3×10^{-3}	4.5×10^{-4}
J = 0, I = 1	2.1×10^{-2}	3.9×10^{-3}	1.36×10^{-3}
J = 1, I = 0	1.7×10^{-2}	9.0×10^{-4}	7.2×10^{-3}
J = 1, I = 1	5.0×10^{-2}	2.7×10^{-3}	2.2×10^{-2}

At $au \sim au_{\it F}$:

- J = 0: $N_{X_0}^{(4q)} \approx 4N_{X_0}^{(Mol)}$
- J = 1: $N_{X_1}^{(Mol)} \approx 2.5 N_{X_1}^{(4q)}$
- HICs: might shed some light on the discrimination of their structure



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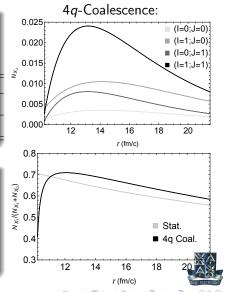
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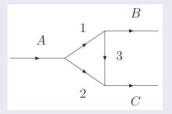
$Z_c(3900)$: triangle singularity or a new hadron?

Belle/BESIII PRL **110** (2013), 252002; 252001

$$Y(4260) \rightarrow Z_c \pi^{\mp} \rightarrow J/\psi \pi^{\pm} \pi^{\mp}$$

Guo, NPR 37 (2020), 406

 Possible kinematic accident of the (D₁*DD*) triangle satisfying the Coleman-Norton theorem



Triangle integral: singular when Coleman-Norton theorem applies

- All particles need to be on-shell
- All particles collinear
- Classically allowed: 3 can reach 2

$$t=t_0 \qquad \frac{1}{2}$$

$$t=t_1 \qquad \frac{B}{3} \qquad \frac{2}{2}$$

$$t=t_f \qquad \frac{B}{3} \qquad \frac{2}{3}$$



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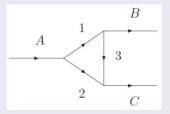
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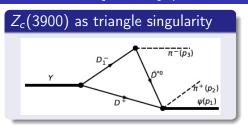
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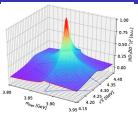
$$t=t_f \qquad \frac{B}{3} \qquad \frac{2}{C}$$



(luciano.abreu@ufba.br)

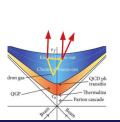
LMA and Llanes-Estrada, EPJ C **81**, 430 (2021); 2109.01015 [hep-ph] (to appear)

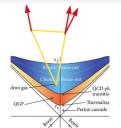




(Guo, NPR 37 (2020), 406)

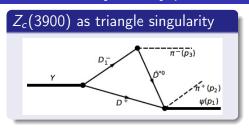
Can HICs help discern the correct interpretation? Two conditions

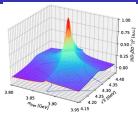




- Sufficient time to complete in a HIC
- Mass and/or width are modified from their vacuum values

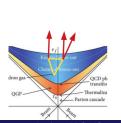
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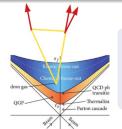




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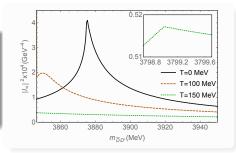




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$Z_c(3900)$ as a TS in HICs; LMA and Llanes-Estrada, EPJ C **81**, 430 (2021); 2109.01015 [hep-ph] (to appear)

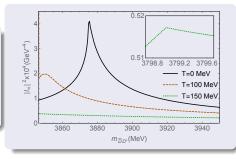
- Triangle loop: two requirements satisfied
- Calculations with thermal Matsubara formalism
- Use of Thermal M and Γ



- Singularity disappears at temperatures just below T_H
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- Medium: a spectroscopic filter and help distinguish actual hadrons from singularities

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- New particle zoo above $D^{(*)}\bar{D}^*, B^{(*)}\bar{B}^*$ thresholds: not $(\bar{q}q, qqq)$

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23 / 23

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Thank You!!!

Partial financial support:





