

Experimental Methods and Physics at the LHC - II

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26th Vietnam School of Physics: Particles and Dark Matter 29 Nov - 11 Dec 2020, Quy Nhon & virtual



Background estimation Matrix – or ABCD - method

When there exist two variables x and y for which the BG is uncorrelated, i.e. factorizable:

 $f^{BG}(x,y) = f^{BG}(x) \cdot f^{BG}(y)$

- Apply all cuts except those on x and y on data
- Divide the x-y plane into 4-regions:
- When there is no signal, we have

 $\frac{N_A^{BG}}{N_B^{BG}} = \frac{N_C^{BG}}{N_D^{BG}}, \quad \frac{N_A^{BG}}{N_C^{BG}} = \frac{N_B^{BG}}{N_D^{BG}}$

 In the presence of signal, A will be contaminated by the signal. But we can estimate the number of BG events in A from

$$N_A^{BG} = \frac{N_C^{BG} N_B^{BG}}{N_D^{BG}}$$



Note: Always beware the signal
 contamination in the control regions. Add it as a systematic.



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CMS tt+jets cross section measurement in the muon+jets channel.



Note: Always beware the signal contamination in the control regions. Add it as a systematic.



The ratios of objects found by a tight identification over objects found by a loose identification is widely used as a BG estimation tool.



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$$\begin{array}{rcl} \mbox{Get these counts} & & & & \\ \mbox{from data} & & & & \\ \mbox{SR:} & & & & \\ \mbox{SR:} & & & & \\ \mbox{k} \equiv N_{tight}^k & & & \\ \mbox{k} \equiv N_{tight}^k / N_{loose}^k \rightarrow & = & \\ \mbox{$\epsilon^{k} \equiv N_{tight}^k / N_{loose}^k \rightarrow $ \\ \mbox{k} = & \\ \mbox{$\epsilon^{real} N_{loose}^{real} + \\ \mbox{$\epsilon^{fake} N_{loose}^{fake} \rightarrow $ \\ \mbox{k} = \\ \mbox{$\epsilon^{real} N_{loose}^{real} + \\ \mbox{$\epsilon^{fake} N_{loose}^{fake} \rightarrow $ \\ \mbox{k} = \\ \mbox{$\epsilon^{real} N_{loose}^{real} + \\ \mbox{$\epsilon^{fake} N_{loose}^{fake} \rightarrow $ \\ \mbox{k} = \\ \mbox{$\epsilon^{real} N_{loose}^{real} + \\ \mbox{$\epsilon^{fake} N_{loose}^{fake} \rightarrow $ \\ \mbox{$\epsilon^{real} N_{loose}^{real} + \\ \mbox{$\epsilon^{fake} N_{loose}^{fake} \rightarrow $ \\ \mbox{$\epsilon^{real} N_{loose}^{real} + \\ \mbox{$\epsilon^{fake} N_{loose}^{fake} \rightarrow $ \\ \mbox{$\epsilon^{real} N_{loose}^{real} + \\ \mbox{$\epsilon^{fake} N_{loose}^{fake} \rightarrow $ \\ \mbox{$\epsilon^{real} N_{loose}^{real} + \\ \mbox{$\epsilon^{fake} N_{loose}^{fake} \rightarrow $ \\ \mbox{$\epsilon^{real} N_{loose}^{real} + \\ \mbox{$\epsilon^{fake} N_{loose}^{fake} = \\ \mbox{$\epsilon^{real} N_{loose}^{real} + \\ \mbox{$\epsilon^{real} N_{loose}$$



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Suppose we would like to estimate QCD BG in a signal region that has leptons with tight ID. Signal contribution comes from real leptons and QCD contribution comes from fake leptons (jets faking leptons). We define a control region with looser lepton ID. CR and SR can be decomposed as:



Finally obtain the number of BG events from

$$\epsilon^{fake} N_{loose}^{fake} = N_{tight}^{fake} = N_{BG}$$



Background estimation Tag-and probe method

- Tag and probe (TP) is a data-driven method used for measuring particle efficiencies. It is used for obtaining trigger, reconstruction, identification efficiencies. Mainly used for leptons.
- For TP, we need a mass resonance decaying to the object whose efficiency we want to measure (e.g. J/psi, upsilon, Z)
- We select two objects, a tag object and a probe object.
 - Tag object : Tight selection/ID criteria- we assume this is a real object.
 - Probe object: Very loose selection/ID criteria.
- We compute the diobject invariant mass of the tag object + probe object.
 - If the invariant mass is close to the resonance mass value, we assume that the probe object was a real object. Otherwise it should be a fake object.
- We take the real leptons inside the resonance mass window and apply on them the criteria of the selection, whose efficiency we want to measure
- Selection efficiency is computed as

$$\epsilon_{\rm selection} = N_{\rm selection}^{\rm in\,mass\,window} / N_{\rm total}^{\rm in\,mass\,window}$$

Background estimation Validating the estimates

BG estimation methods must always be validated with closure tests or independent validation regions, or alternative methods.

Closure tests : Validate the internal consistency of the method, e.g. validate the method using purely MC events.

Validation regions : Validate the method in independent dedicated regions. These can have a composition similar to the signal regions but be dominated by BG. Estimate should be equal to data.

Alternative methods : Estimate the BG with multiple methods and compare the results.





Lecture 3: Systematic uncertainties, results, interpretation A few highlights from LHC searches

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(6) Systematic uncertainties

Systematic uncertainties Statistical vs. systematic

- Statistical uncertainties are the result of stochastic fluctuations arising from the fact that a measurement is based on a finite set of observations.
 - Repeated measurements of the same phenomenon will therefore result in a set of observations that will differ. Statistical uncertainty is a measure of the range of this variation.
- Systematic uncertainties or "systematics" arise from uncertainties associated with the nature of the experiment, assumptions made by the experimenter, or the theoretical model used to make inferences based on the observed data.
 - They cause a shift in the mean of a measurement from the true value due to effects from the experimental setup / theoretical calculations. They are different from statistical uncertainties.
 - They can effect either the normalization or the shape of a given distribution.

Effects of both uncertainties on the measured quantity is included in an experimental result.

Systematic uncertainties Types of systematics

- Experimental uncertainties (coming from the accelerator/detector setup) on:
 - Luminosity measurements
 - Trigger efficiency measurements
 - Jet energy scale, jet energy resolution measurements
 - Lepton, photon, b-jet, W-jet, top-jet, etc. efficiency measurements
 - Uncertainties derived in background estimations, e.g. from closure tests, validation regions or comparison of alternative estimation methods
 - ..
- Theoretical uncertainties (due to insufficient accuracy of calculations) on:
 - Cross section and branching ratio calculations
 - Parton distribution functions
 - Initial state / final state radiation, renormalization / factorization scale effects
 - ...

When using MC in an analysis, systematics must be applied on the MC events to reflect the uncertainties in the predictions.

Systematic uncertainties Types of systematics

Source	Signal	Lost lepton	1ℓ (not from t)	$Z ightarrow \nu \bar{ u}$
Data statistical uncertainty		5-50%	4-30%	
Simulation statistical uncertainty	6–36%	3-68%	5-70%	4–41%
tī $p_{\rm T}^{\rm miss}$ modeling		3–50%		
Signal $p_{\rm T}^{\rm miss}$ modeling	1–25%			
QCD scales	1-5%	0–3%	2–5%	1-40%
Parton distribution		0–4%	1-8%	1–12%
Pileup	1-5%	1-8%	0–5%	0–7%
Luminosity	2.3-2.5%			2.3–2.5%
$W + b(\overline{b})$ cross section			20–40%	
ttZ cross section				5-10%
System recoil (ISR)	1–13%	0–3%		
Jet energy scale	2-24%	1-16%	1-34%	1–28%
$p_{\rm T}^{\rm miss}$ resolution		1-10%	1–5%	
Trigger	2–3%	1–3%		2–3%
Lepton efficiency	3–4%	2–12%		1–2%
Merged t tagging efficiency	3–6%			5-10%
Resolved t tagging efficiency	5-6%			3–5%
b tagging efficiency	0–2%	0–1%	1–7%	1-10%
Soft b tagging efficiency	2–3%	0–1%	0–1%	0–5%

Example set of systematic uncertainties and their effects on expected signal yields and predicted background yields from a CMS SUSY analysis.



(8) Results and (9) Interpretation (including statistical model (7))

Experimental results What do we measure?

An experimental result is the empirical outcome of the experiment, the measurement of some physical quantity, such as:

- Event counts < mainly in new physics search analyses.
- Masses
- Cross sections/branching ratios
- Signal strengths
- Couplings
- Kinematic shapes, peaks, edges, endpoints (usually to derive masses or mass differences)
- Decay widths
- Charge asymmetry
- Spin correlations
- etc.

Experimental results Event counts



The results of the SUSY razor search consist of event counts in 2-dimensional bins of razor variables M_R and R^2 .

- Observed data counts
- Estimated yields from 4 different backgrounds.
- Systematic uncertainties on the estimated BG yields (gray shaded band).

In a search analysis, we compare observed data and estimated BG yields to see if there is agreement or discrepancy.

We quantify this with a statistical analysis.

It is customary to show expected / MC distributions from a few signal points, but these are not a part of the results.



Experimental results Higgs mass measurements



Experimental results Cross section measurements

For a signal process, the number of expected signal events is given by

$$N_{signal} = \sigma_{signal} \cdot \int \mathcal{L}dt \cdot \epsilon_{signal} = N^{observed} - N_{BG}$$

where ε_{signal} is the product of all branching ratios, geometrical and kinematical acceptance, efficiencies for trigger, object reconstruction and identification. The cross section can be obtained by reverting the equation:

$$\sigma_{signal} = \frac{N^{observed} - N_{BG}}{\int \mathcal{L}dt \cdot \epsilon_{signal}}$$

Cross sections for SM processes are continuously being refined at the LHC.

Experimental results SM cross section measurements

	m				$[fb^{-1}]$	Reference
	$\sigma = 96.07 \pm 0.18 \pm 0.91$ mb (data) COMPETE HPR1R2 (theory)				50×10 ⁻⁸	PLB 761 (2016) 158
	$\sigma = 95.35 \pm 0.38 \pm 1.3$ mb (data) COMPETE HPR1R2 (theory)	ATLAS Preliminary	o		8×10 ⁻⁸	NPB 889, 486 (2014)
	$\sigma = 190.1 \pm 0.2 \pm 6.4 \text{ nb (data)}$ DYNNLO + CT14NNLO (theory)	,	þ	6	0.081	PLB 759 (2016) 601
	$\sigma = 112.69 \pm 3.1$ nb (data) DYNNLO + CT14NNLO (theory)	Bup 1 2 $\sqrt{c} = 7.8.12^{\circ}$			20.2	EPJC 79, 760 (2019)
	σ = 98.71 ± 0.028 ± 2.191 nb (data) DYNNLO + CT14NNLO (theory)	$\pi u = 1, 2, \forall S = 1, 0, 13$			4.6	EPJC 77, 367 (2017)
	$\sigma = 58.43 \pm 0.03 \pm 1.66$ nb (data) DYNNLO+CT14 NNLO (theory)		b '		3.2	JHEP 02 (2017) 117
	$\sigma = 34.24 \pm 0.03 \pm 0.92$ nb (data) DYNNLO+CT14 NNLO (theory)		Å		20.2	JHEP 02 (2017) 117
	$\sigma = 29.53 \pm 0.03 \pm 0.77$ nb (data) DYNNLO+CT14 NNLO (theory)		ó		4.6	JHEP 02 (2017) 117
	$\sigma = 826.4 \pm 3.6 \pm 19.6 \text{ pb} (\text{data})$ top++ NNLO+NNLL (theory)	Ċ		i i i i i i i i i i i i i i i i i i i	36.1	arXiv: 1910.08819
	$\sigma = 242.9 \pm 1.7 \pm 8.6 \text{ pb (data)}$	Å			20.2	EPJC 74, 3109 (201-
	$\sigma = 182.9 \pm 3.1 \pm 6.4 \text{ pb (data)}$	ō'			4.6	EPJC 74, 3109 (2014
	$\sigma = 247 \pm 6 \pm 46 \text{ pb (data)}$	6			3.2	JHEP 04 (2017) 086
chan	$\sigma = 89.6 \pm 1.7 + 7.2 - 6.4 \text{ pb (data)}$	▲ -			20.3	EPJC 77, 531 (2017
	$\sigma = 68 \pm 2 \pm 8 \text{ pb (data)}$	ò			4.6	PRD 90, 112006 (20
	$\sigma = 130.04 \pm 1.7 \pm 10.6 \text{ pb (data)}$				36.1	EPJC 79, 884 (2019
wl	$\sigma = 68.2 \pm 1.2 \pm 4.6 \text{ pb (data)}$	<u>م</u>			20.3	PLB 763, 114 (2016
	$\sigma = 51.9 \pm 2 \pm 4.4 \text{ pb} (\text{data})$	6			4.6	PRD 87, 112001 (20
	$\sigma = 61.7 \pm 2.8 + 4.3 - 3.6 \text{ pb} (data)$	6			79.8	PRD 101 (2020) 012
	$\sigma = 27.7 \pm 3 + 2.3 - 1.9 \text{ pb (data)}$	<u>ل</u>			20.3	EPJC 76, 6 (2016)
	$\sigma = 22.1 + 6.7 - 5.3 + 3.3 - 2.7$ pb (data)	6	Theory		4.5	EPJC 76, 6 (2016)
	$\sigma = 94 \pm 10 + 28 - 23 \text{ pb (data)}$				3.2	JHEP 01 (2018) 63
t	$\sigma = 23 \pm 1.3 + 3.4 - 3.7 \text{ pb (data)}$	Δ.	$1 \text{ HC } \text{ pp } \sqrt{6} = 12 \text{ TeV}$		20.3	JHEP 01, 064 (2016
- 	$\sigma = 16.8 \pm 2.9 \pm 3.9 \text{ pb} \text{ (data)}$	D	$Effor pp \forall s = 15 fev$		2.0	PLB 716, 142-159 (2
	$\sigma = 51 \pm 0.8 \pm 2.3 \text{ pb (data)}$		Data –		36.1	EPJC 79, 535 (2019 PLB 761 (2016) 179
7	$\sigma = 24.3 \pm 0.6 \pm 0.9 \text{ pb} \text{ (data)}$	\mathbf{A}^{+}	stat		20.3	PRD 93, 092004 (20 PLB 761 (2016) 179
-	$\sigma = 19 + 1.4 - 1.3 \pm 1 \text{ pb (data)}$	o di	Stat & Syst		4.6	EPJC 72, 2173 (201 PLB 761 (2016) 179
	$\sigma = 17.3 \pm 0.6 \pm 0.8 \text{ pb} (\text{data})$		LHC pp $\sqrt{s} = 8$ TeV	1 1	36.1	PRD 97 (2018) 0320
,	$\sigma = 7.3 \pm 0.4 + 0.4 - 0.3 \text{ pb (data)}$	A ⁺	Data		20.3	JHEP 01, 099 (2017
-	$\sigma = 6.7 \pm 0.7 + 0.5 - 0.4 \text{ pb} \text{ (data)}$	0	stat		4.6	JHEP 03, 128 (2013)
chan	$\sigma = 4.8 \pm 0.8 \pm 1.6 \pm 1.3 \text{ pb} \text{ (data)}$	A	stat ⊕ syst		20.3	PLB 756, 228-246 (2
Chan	$\sigma = 870 \pm 130 \pm 140 \text{ fb} \text{ (data)}$		$1 \text{ HC nn } \sqrt{s} = 7 \text{ TeV}$		36.1	PRD 99, 072009 (20
N	$\sigma = 369 + 86 - 79 \pm 44 \text{ fb} \text{ (data)}$	-			20.3	JHEP 11, 172 (2015
	$\sigma = 950 \pm 80 \pm 100 \text{ fb} \text{ (data)}$	6	• Data —		36.1	PRD 99, 072009 (20
Z	$\sigma = 176 + 52 - 48 \pm 24 \text{ fb} (\text{data})$	۲	stat e syst		20.3	JHEP 11, 172 (2015
w/w/	$\sigma = 0.65 + 0.16 - 0.15 + 0.16 - 0.14 \text{ pb} (data)$	0	51al - 5y51		79.8	PLB 798 (2019) 134
W/7	$\sigma = 0.55 \pm 0.14 + 0.15 - 0.13 \text{ pb (data)}$	6			79.8	PLB 798 (2019) 134
	Snerpa 2.2.2 (tneory)		الماسم ۸۸۸ جاستین استین استین		10.0	. 12 100 (2010) 104
1	0^{-4} 10^{-3} 10^{-2} 10^{-1}	$1 10^1 10^2 10^3$	$10^4 \ 10^5 \ 10^6 \ 10^1$	1 0.5 1.0 15 20		
T	0 10 10 10	- 10 10 10				

Experimental results Higgs cross section measurements



Experimental results Differential ttbar cross section measurement



Experimental results Signal strength

Signal strength in a search region:

 $\mu = \frac{\text{Observed event yield for a process}}{\text{Expected SM event yield for the process}}$

- μ = 0: no signal
- $\mu = 1$: signal consistent with the SM
- $\mu > 1$: signal different from the SM -> BSM



Higgs cross sections

Higgs branching ratios



Experimental results Signal strength - let's get fancier

Fitting the signal strength and the Higgs mass together to find their common best values.

Comparing production signal strengths for different decay channels.



All these help to check the consistency of the discovered Higgs with the SM and and see if there are any hints for BSM

Experimental results Couplings



Interpretation The idea

- Interpretation is the comparison of experimental results with the expectations of a given a theoretical model.
- CAUTION! Interpretation is NOT the experimental result.
- We use a statistical model and likelihood to interpret the experimental result.
 - The statistical model of an analysis provides the complete mathematical description of that analysis.
 - It relates the observed quantities x to the theoretical model's parameters θ_i through the probability density $p(x|\theta_i)$. Parameters θ_i include signal model parameters θ_s and background model parameters θ_b .
 - The likelihood $L(\theta_i) = p(X_0|\theta_i)$ is the probability density $p(x|\theta_i)$ evaluated at the observed values X_0 of the observables x.
 - The likelihood is the starting point of any interpretation.
- We estimate parameters θ using statistical procedures, and test the validity of the model.
- An experimental result can be interpreted with multiple theoretical models.

Limits: Suppose we are testing a theory with a free parameter θ by an experimental analysis. For example, this parameter could be the mass m of a particle.

Suppose our analysis did not observe a sign of this particle. Data look consistent with the background hypothesis. What can we say about the theory and m?

We can say up to which value of m our analysis could have observed m, if the theory were true. This is called a limit.

For a theory with free parameters θ_i , a limit establishes the boundaries defining the range where the experiment can make a statement about the theory.

Limits are obtained using elaborate statistical methods based on likelihoods.

- Observed limit: Obtained by comparing observed data with 1) signal MC + estimated BG and 2) with only estimated BG.
 - Checks the consistency of the observation with the signal + BG hypothesis and compares it to the BG-only hypothesis.
- Expected limit: Obtained by comparing estimated BG with signal + estimated BG.
 - Useful for predicting the analysis sensitivity.

Interpretation

Limits

Interpretation Limits for 1 free parameter



CMS search for high mass dijet resonances.

Interpretation Limits for 2 free parameters

Interpreting SUSY razor analysis results in terms of a simplified SUSY model, with 2 particles with free masses and cross sections.





Exactly like the 1 free parameter case on the previous page, but for 2 free parameters!

Interpretation Interpreting with multiple models

SUSY razor analysis results interpreted for multiple theoretical models:



Interpretation SUSY mass limits

Mass limits on SUSY particles gluinos, squarks in various decay channels to neutralinos. Interpretations were based on several analyses.

Interpreting the same model with multiple analyses provides complementarity.



Interpretation Limits on exotic models

Overview of CMS EXO results



Interpretation Limits on long-lived particles



Now the free parameter is the long-lived particle's proper lifetime. Particle masses are fixed.

Different searches exclude different lifetimes, providing complementary information.

Interpretation Limits on long-lived particles

ATLAS Long-lived Particle Searches* - 95% CL Exclusion

Status: May 2020

ATLAS Preliminary





"Data are coming! Data are coming!"

Actually, it is the other way around...



What we learn from the absence of discovery and excluding many models takes us closer to discovering the true nature of the Universe.



Hands-on exercises 1 and 2: Looking into signal and background events, plotting variable distributions, applying simple event selections.