

Measurement of an excess in the yield of J/ψ at very low p_T in Pb-Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV

- ❖ J/ψ production in heavy-ion collisions
- ❖ Vector meson photoproduction in ultra-peripheral collisions
- ❖ Analysis in Pb-Pb at 5.02 TeV
- ❖ Systematics on the tracking efficiency
- ❖ Conclusion and outlooks

Quarkonia production in heavy-ion collisions

- ▶ **Quarkonia** are bound states of (anti-)charm and (anti-)beauty quarks
 - Charmonia ($c\bar{c}$) : $J/\psi(1S)$, $\chi_{c1}(1P)$ and $\psi(2S)$
 - Bottomonia ($b\bar{b}$) : $\Upsilon(1S)$, $\Upsilon(2S)$ and $\Upsilon(3S)$

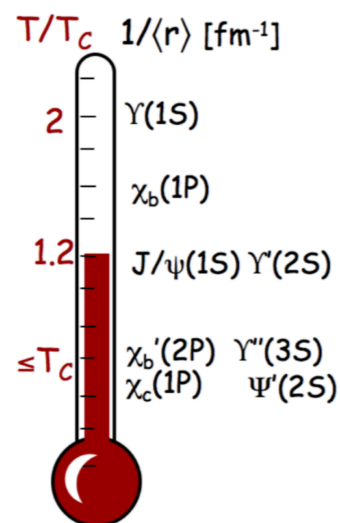
Quarkonia production in heavy-ion collisions

- ▶ **Quarkonia** are bound states of (anti-)charm and (anti-)beauty quarks
 - Charmonia ($c\bar{c}$) : $J/\psi(1S)$, $\chi_{c1}(1P)$ and $\psi(2S)$
 - Bottomonia ($b\bar{b}$) : $\Upsilon(1S)$, $\Upsilon(2S)$ and $\Upsilon(3S)$
- ▶ Heavy quarks produced at the **early stages of AA collisions** → experience the hot QCD medium evolution :

Quarkonia production in heavy-ion collisions

- ▶ **Quarkonia** are bound states of (anti-)charm and (anti-)beauty quarks
 - Charmonia ($c\bar{c}$) : $J/\psi(1S)$, $\chi_{c1}(1P)$ and $\psi(2S)$
 - Bottomonia ($b\bar{b}$) : $\Upsilon(1S)$, $\Upsilon(2S)$ and $\Upsilon(3S)$

- ▶ Heavy quarks produced at the **early stages of AA collisions** → experience the hot QCD medium evolution :
 - **Sequential suppression** via color-charge screening inside the QGP
 Phys. Lett. B178 (1986) 416

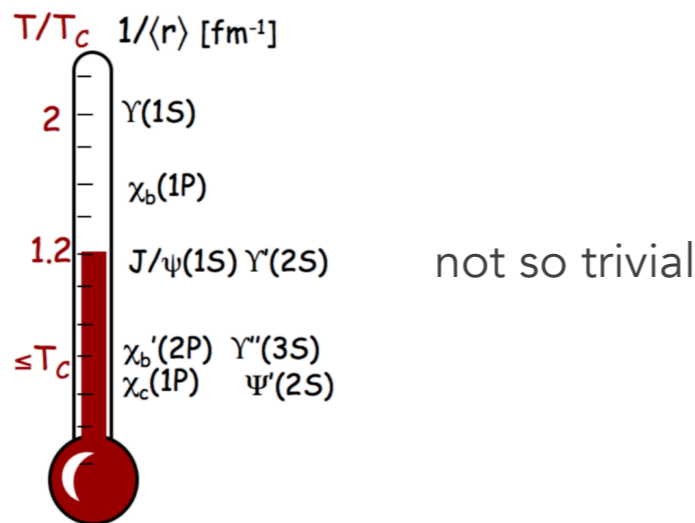


Eur. Phys. J. C61 (2009) 705

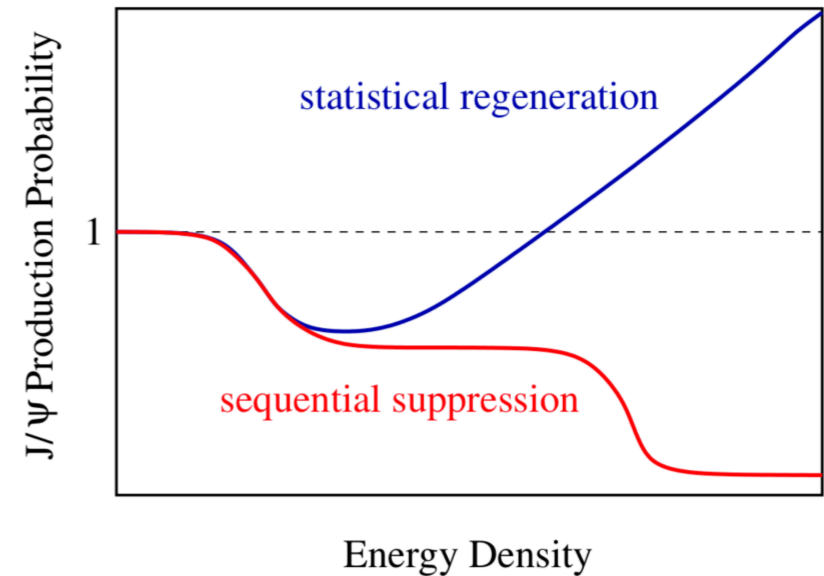
Quarkonia production in heavy-ion collisions

- ▶ **Quarkonia** are bound states of (anti-)charm and (anti-)beauty quarks
 - Charmonia ($c\bar{c}$) : $J/\psi(1S)$, $\chi_{c1}(1P)$ and $\psi(2S)$
 - Bottomonia ($b\bar{b}$) : $\Upsilon(1S)$, $\Upsilon(2S)$ and $\Upsilon(3S)$

- ▶ Heavy quarks produced at the **early stages of AA collisions** → experience the hot QCD medium evolution :
 - **Sequential suppression** via color-charge screening inside the QGP
Phys. Lett. B178 (1986) 416
 - **Regeneration** inside the QGP and/or during hadronisation phase
Phys. Lett. B490 (2000) 196, Phys. Rev. C63 (2001) 054904



Eur. Phys. J. C61 (2009) 705

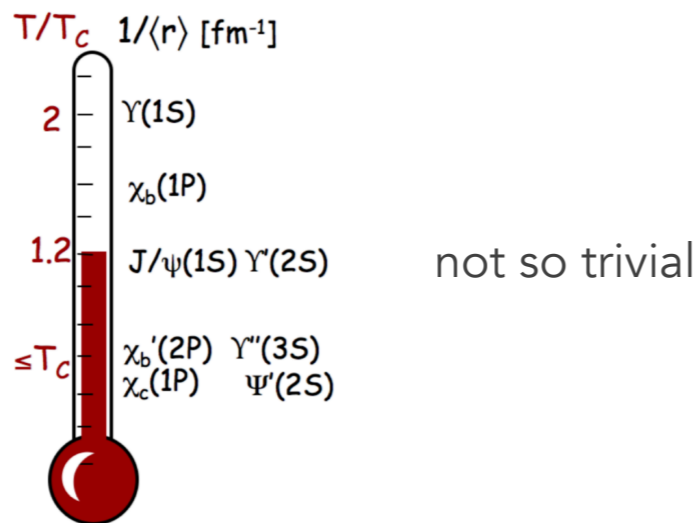


Nucl. Phys. Proc. Suppl. 214 (2011) 3

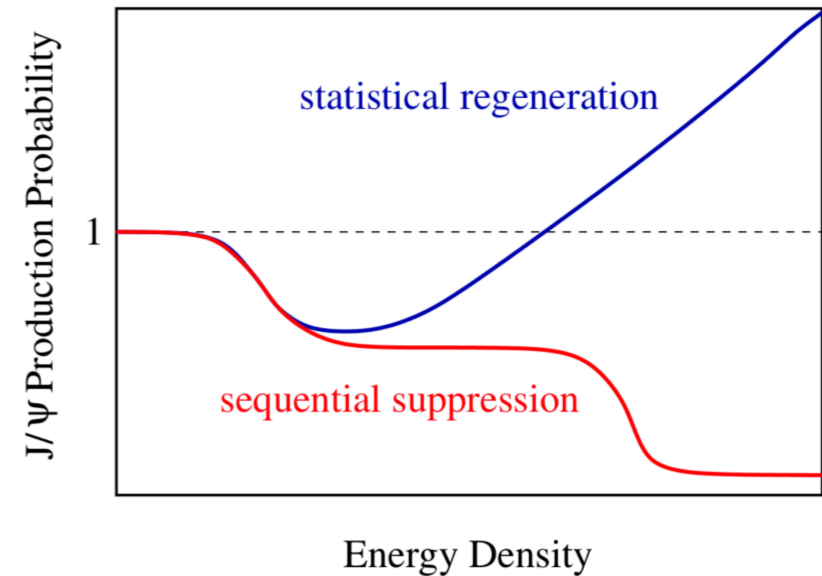
Quarkonia production in heavy-ion collisions

- ▶ **Quarkonia** are bound states of (anti-)charm and (anti-)beauty quarks
 - Charmonia ($c\bar{c}$) : $J/\psi(1S)$, $\chi_{c1}(1P)$ and $\psi(2S)$
 - Bottomonia ($b\bar{b}$) : $\Upsilon(1S)$, $\Upsilon(2S)$ and $\Upsilon(3S)$

- ▶ Heavy quarks produced at the **early stages of AA collisions** → experience the hot QCD medium evolution :
 - **Sequential suppression** via color-charge screening inside the QGP
Phys. Lett. B178 (1986) 416
 - **Regeneration** inside the QGP and/or during hadronisation phase
Phys. Lett. B490 (2000) 196, Phys. Rev. C63 (2001) 054904



Eur. Phys. J. C61 (2009) 705



Nucl. Phys. Proc. Suppl. 214 (2011) 3

❖ Need a well understanding of the production mechanisms

J/ψ production in different collision systems

- ▶ **pp collisions**
Reference for charmonium production mechanisms

J/ψ production in different collision systems

- ▶ **pp collisions**
Reference for charmonium production mechanisms
- ▶ **pA collisions**
study the cold nuclear matter effects

J/ψ production in different collision systems

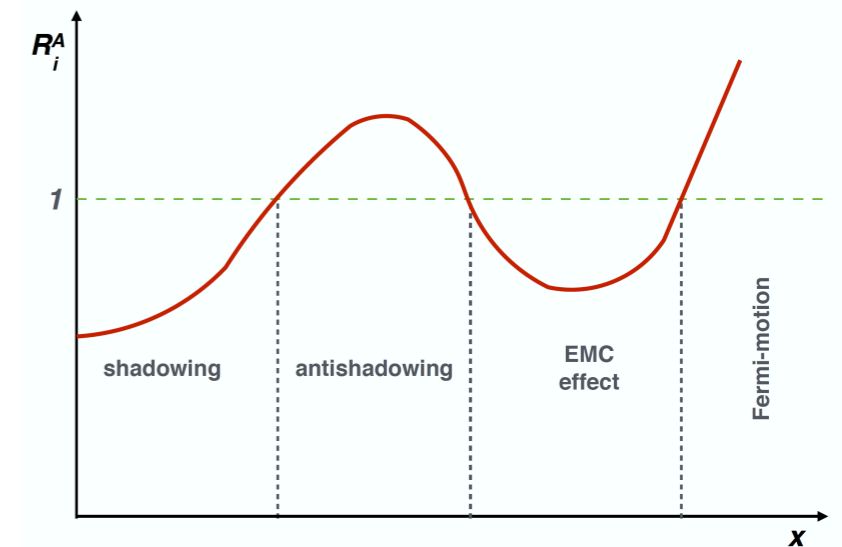
- ▶ **pp collisions**
Reference for charmonium production mechanisms
- ▶ **pA collisions**
study the cold nuclear matter effects
- ▶ **AA collisions**
access to hot nuclear matter effects and QGP properties

J/ ψ production in different collision systems

- ▶ **pp collisions**
Reference for charmonium production mechanisms
- ▶ **pA collisions**
study the cold nuclear matter effects
- ▶ **AA collisions**
access to hot nuclear matter effects and QGP properties

Measure the probes in the three systems and compare

← pp vs pA →



J. Phys. G32 (2006) R367

J/ ψ production in different collision systems

- ▶ **pp collisions**
Reference for charmonium production mechanisms

- ▶ **pA collisions**
study the cold nuclear matter effects

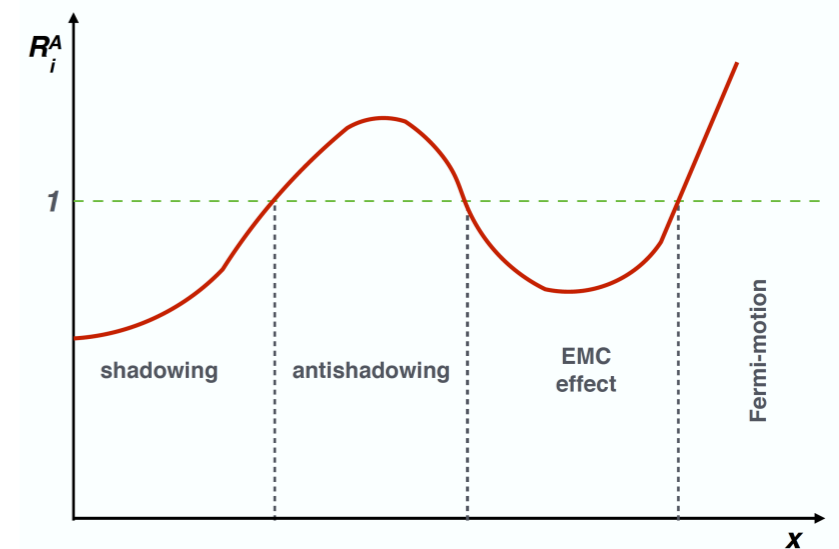
- ▶ **AA collisions**
access to hot nuclear matter effects and QGP properties

Measure the probes in the three systems and compare

← pp vs pA →

← pp vs AA

⚠ cold nuclear effects also in PbPb



J. Phys. G32 (2006) R367

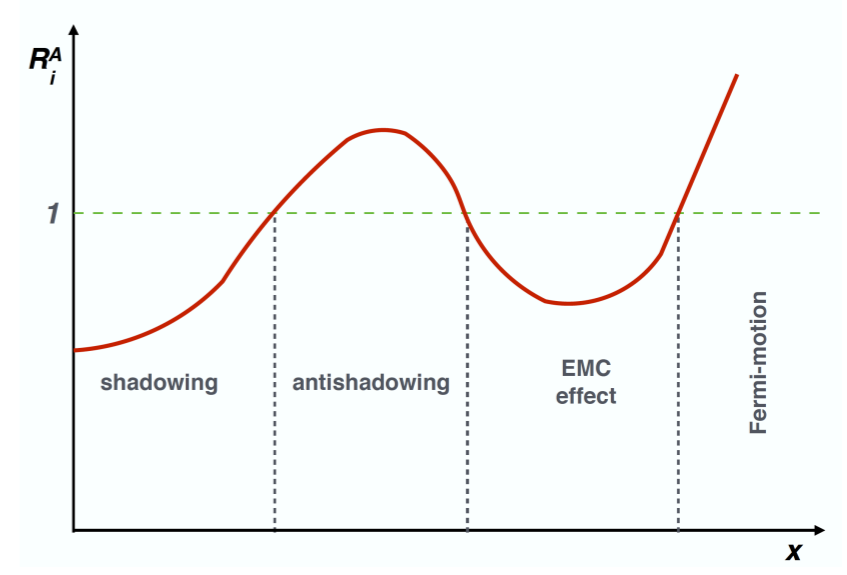
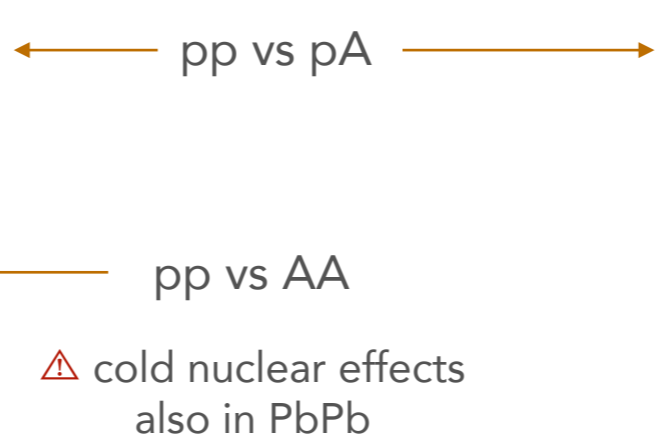
J/ψ production in different collision systems

▶ **pp collisions**
Reference for charmonium production mechanisms

▶ **pA collisions**
study the cold nuclear matter effects

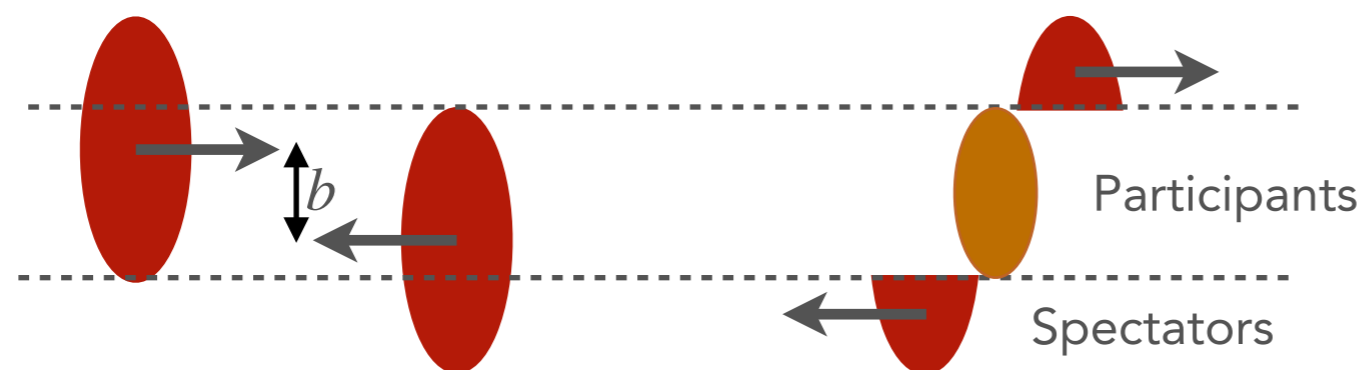
▶ **AA collisions**
access to hot nuclear matter effects and QGP properties

Measure the probes in the three systems and compare



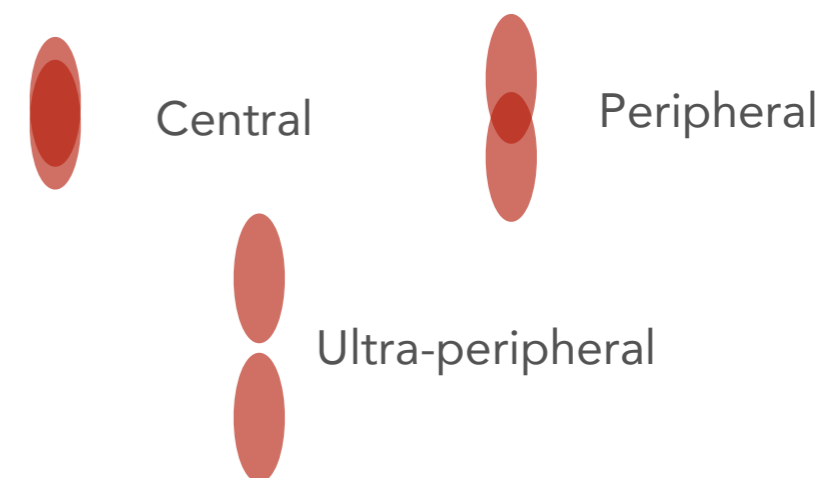
J. Phys. G32 (2006) R367

▶ Glauber Model to describe heavy ion collisions



Centrality class ↑ with *b*

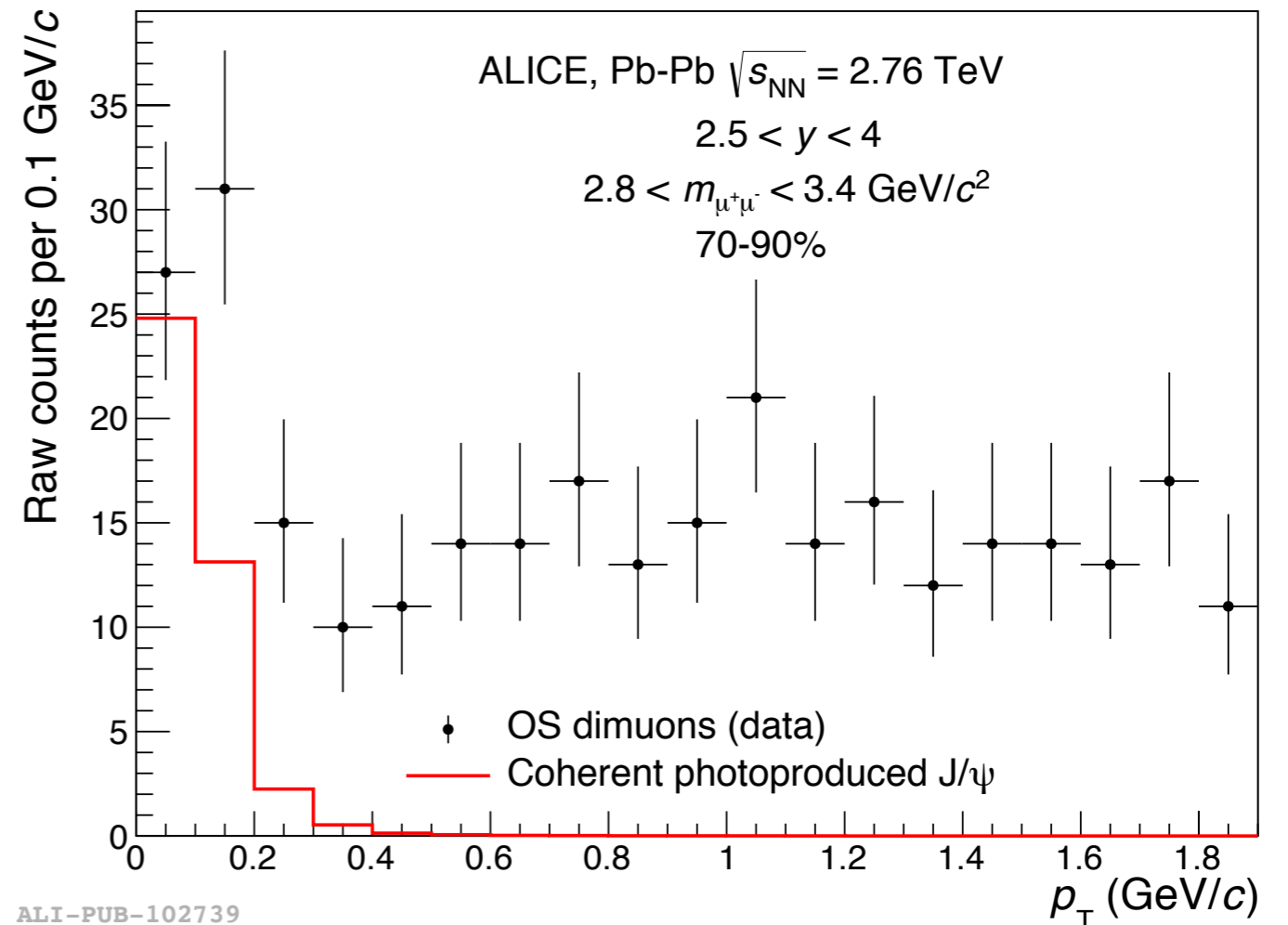
0% to 90%



J/ψ excess at very low p_T in Pb-Pb collisions at 2.76 TeV

❖ First measurement of an excess at very low p_T in peripheral Pb-Pb collisions

- Around the J/ψ mass
 $M_{J/\psi} = 3.096 \text{ GeV}/c$
- At very low p_T
 $\langle p_T \rangle \simeq 0.055 \text{ GeV}/c$
- For peripheral events only
- Observed only in dimuon unlike sign events
- Never reported in pp collisions
↳ will affect the R_{AA}

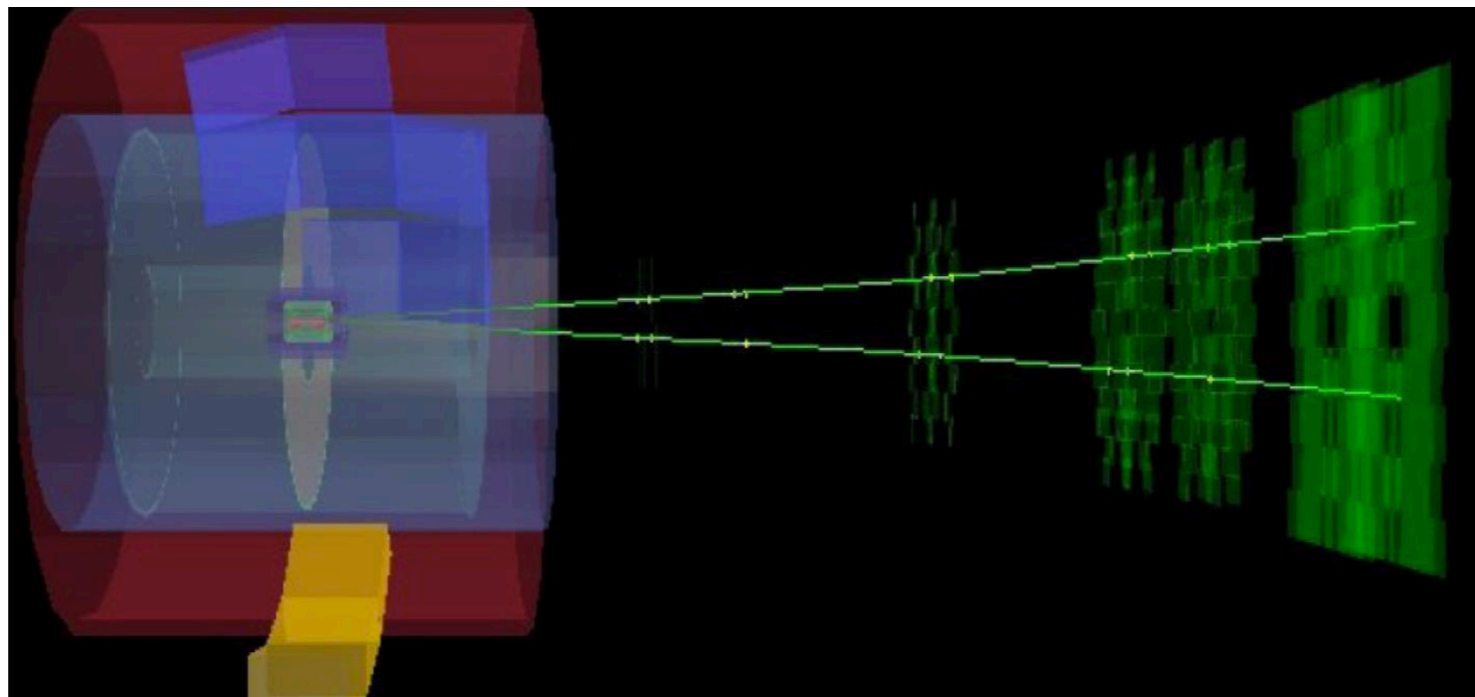
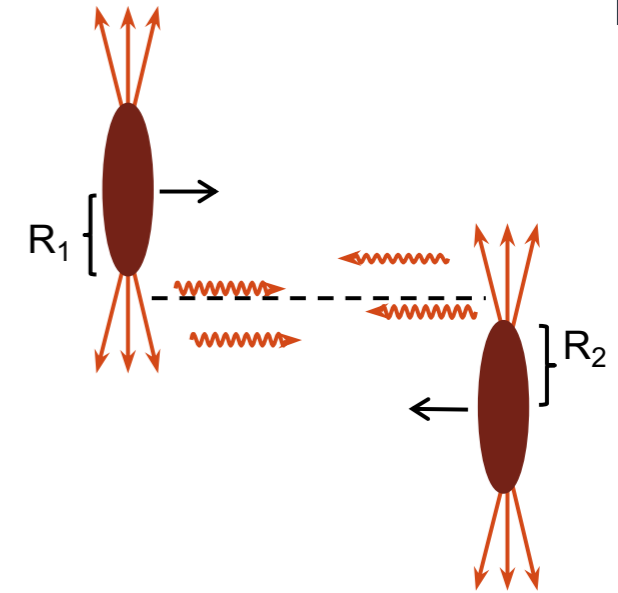


Phys. Rev. Lett. 116 (2016) 222301

Ultra-peripheral collisions

UPC = impact parameter $b > R_1 + R_2$

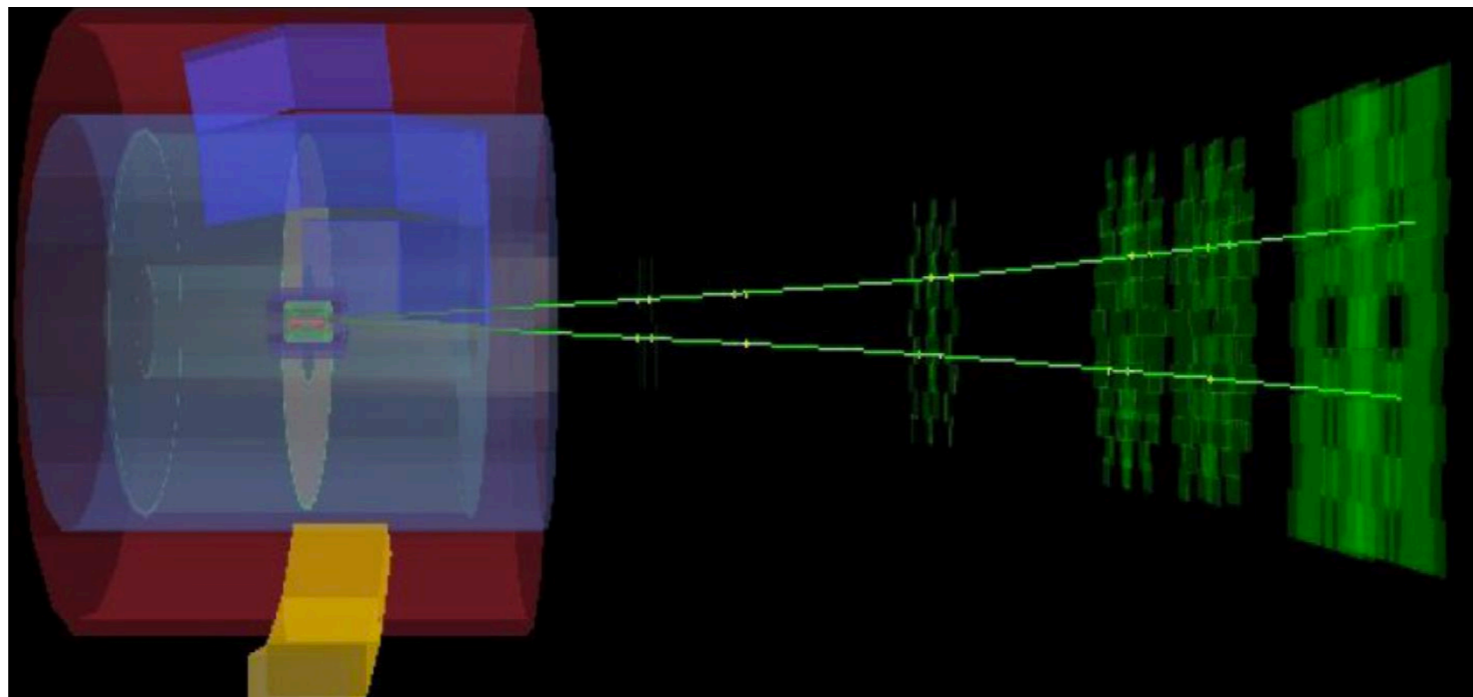
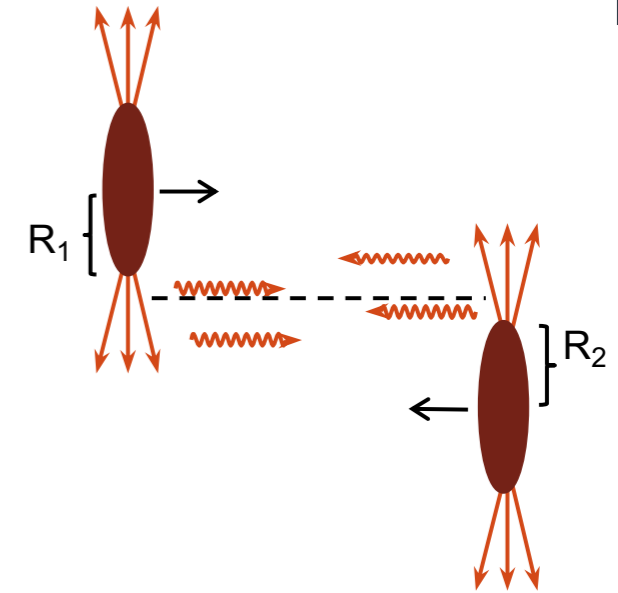
- ▶ Hadronic interactions are strongly suppressed
- ▶ High electromagnetic field from ultra-relativistic Pb nuclei
 - Treated as quasi-real photons flux by the Weizsäcker-Williams approximation
 - Photon-flux $\propto Z^2$



Ultra-peripheral collisions

UPC = impact parameter $b > R_1 + R_2$

- ▶ Hadronic interactions are strongly suppressed
- ▶ High electromagnetic field from ultra-relativistic Pb nuclei
 - Treated as quasi-real photons flux by the Weizsäcker-Williams approximation
 - Photon-flux $\propto Z^2$
- ❖ High cross-section for **photon induced processes** in exclusive events



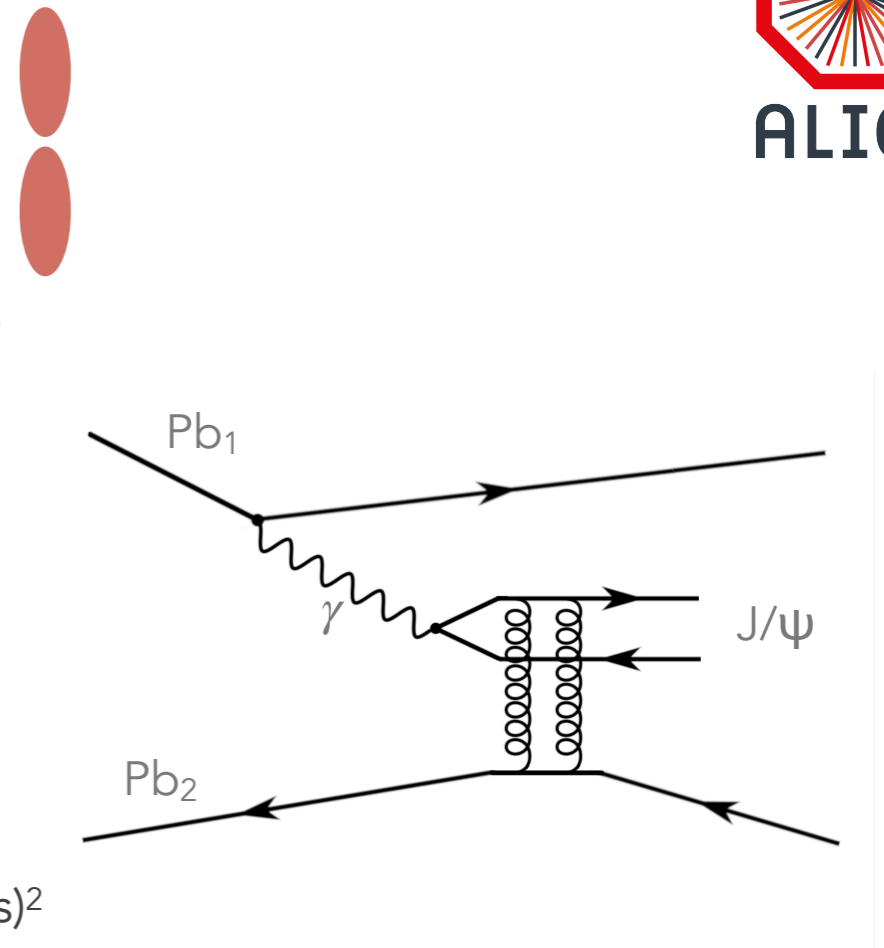
Coherent J/ψ photoproduction in UPC

► QED part

- Intense electromagnetic field of Pb_1 acts as quasi-real photon source
- Flux of photons with a low transfer momentum Q^2

► QCD part

- Photon fluctuates in a $q\bar{q}$ pair
- $q\bar{q}$ pair scatters off $Pb_2 \rightarrow$ emerges as vector meson
- Treated at Leading Order (LO) perturbative QCD (pQCD)
- Photo-nuclear cross section $\sigma_{\gamma A} \propto (\text{gluon density of the target nucleus})^2$



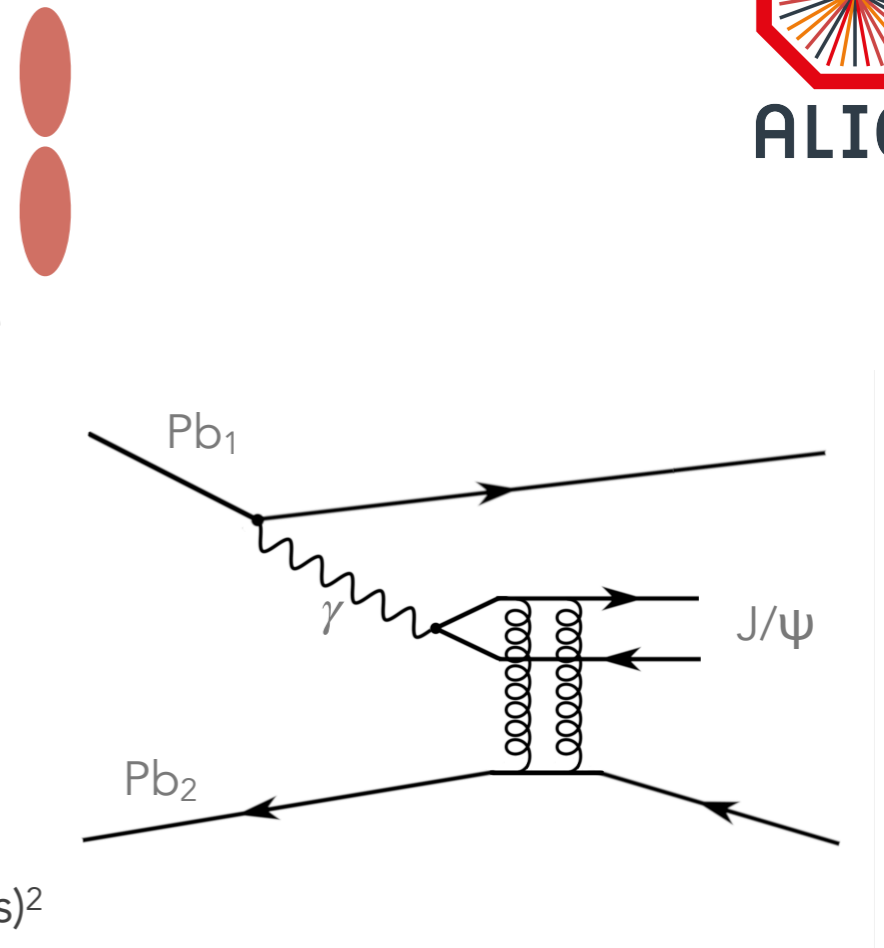
Coherent J/ψ photoproduction in UPC

► QED part

- Intense electromagnetic field of Pb_1 acts as quasi-real photon source
- Flux of photons with a low transfer momentum Q^2

► QCD part

- Photon fluctuates in a $q\bar{q}$ pair
- $q\bar{q}$ pair scatters off $Pb_2 \rightarrow$ emerges as vector meson
- Treated at Leading Order (LO) perturbative QCD (pQCD)
- Photo-nuclear cross section $\sigma_{\gamma A} \propto (\text{gluon density of the target nucleus})^2$



❖ Used to probe the gluon distribution of the target nucleus at low Bjorken- x

$$x \approx \frac{M_{J/\psi}}{\sqrt{s_{NN}}} e^{\pm y}$$

Covering from $x = 10^{-5}$ to $x = 10^{-2}$ at LHC energies

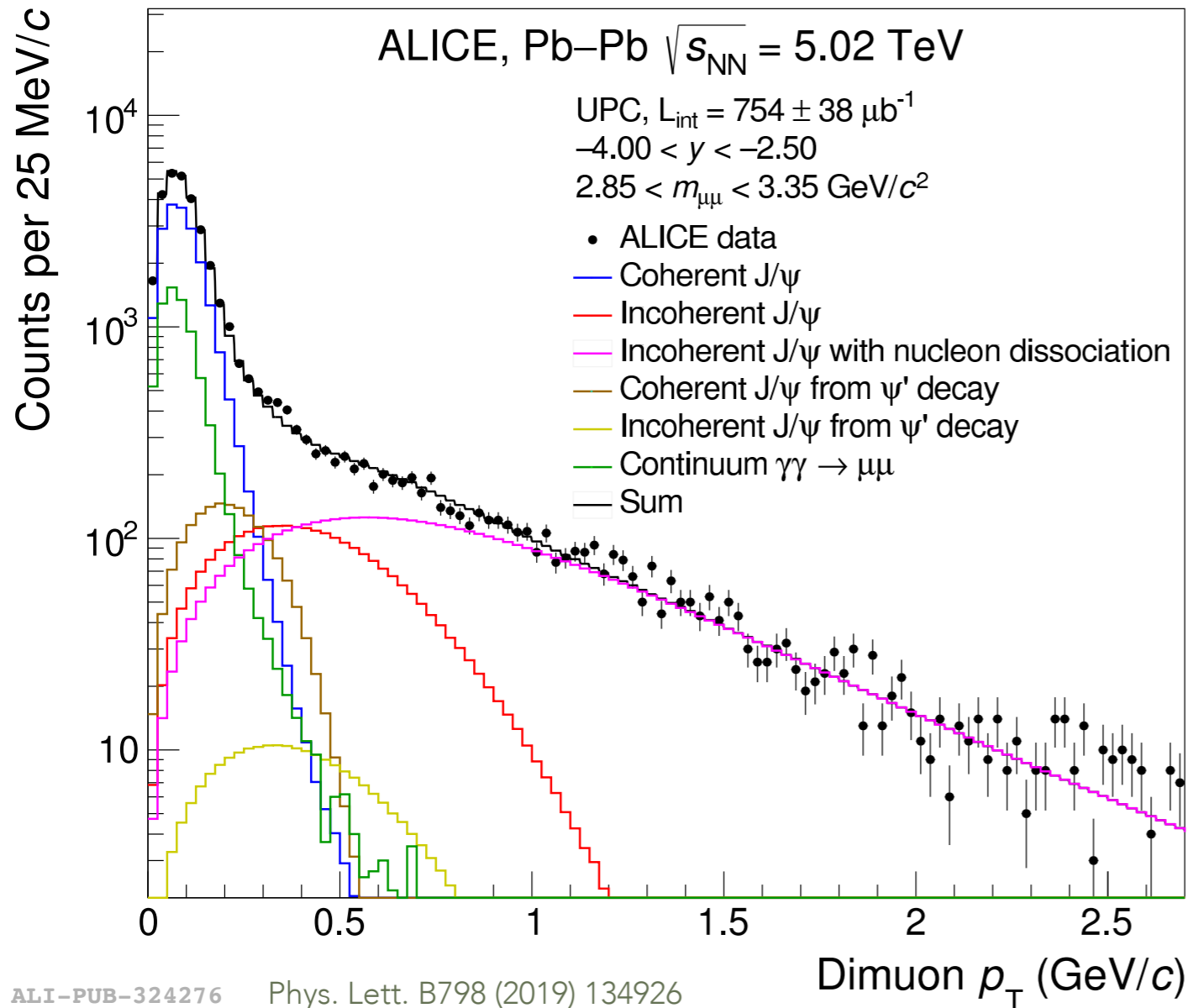
- With the dipole model : study the gluon shadowing (nPDF)
- If not maybe more generalized parton distribution (GPD)

Coherent J/ψ photoproduction in UPC

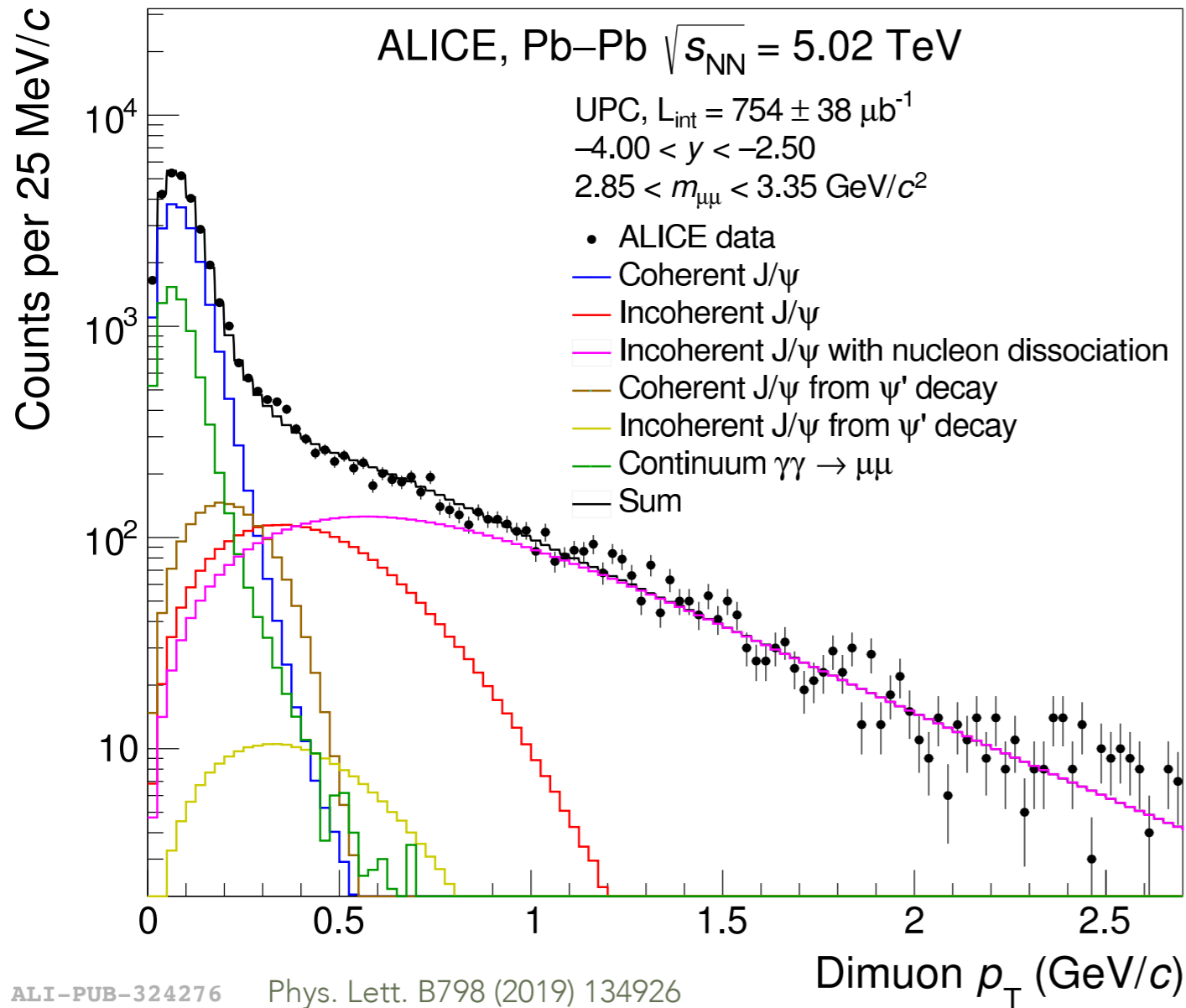


Clear experimental signature :

- Exclusive vector meson with very **low** p_T
- $\langle p_T \rangle \approx 1/R_T$



Coherent J/ψ photoproduction in UPC

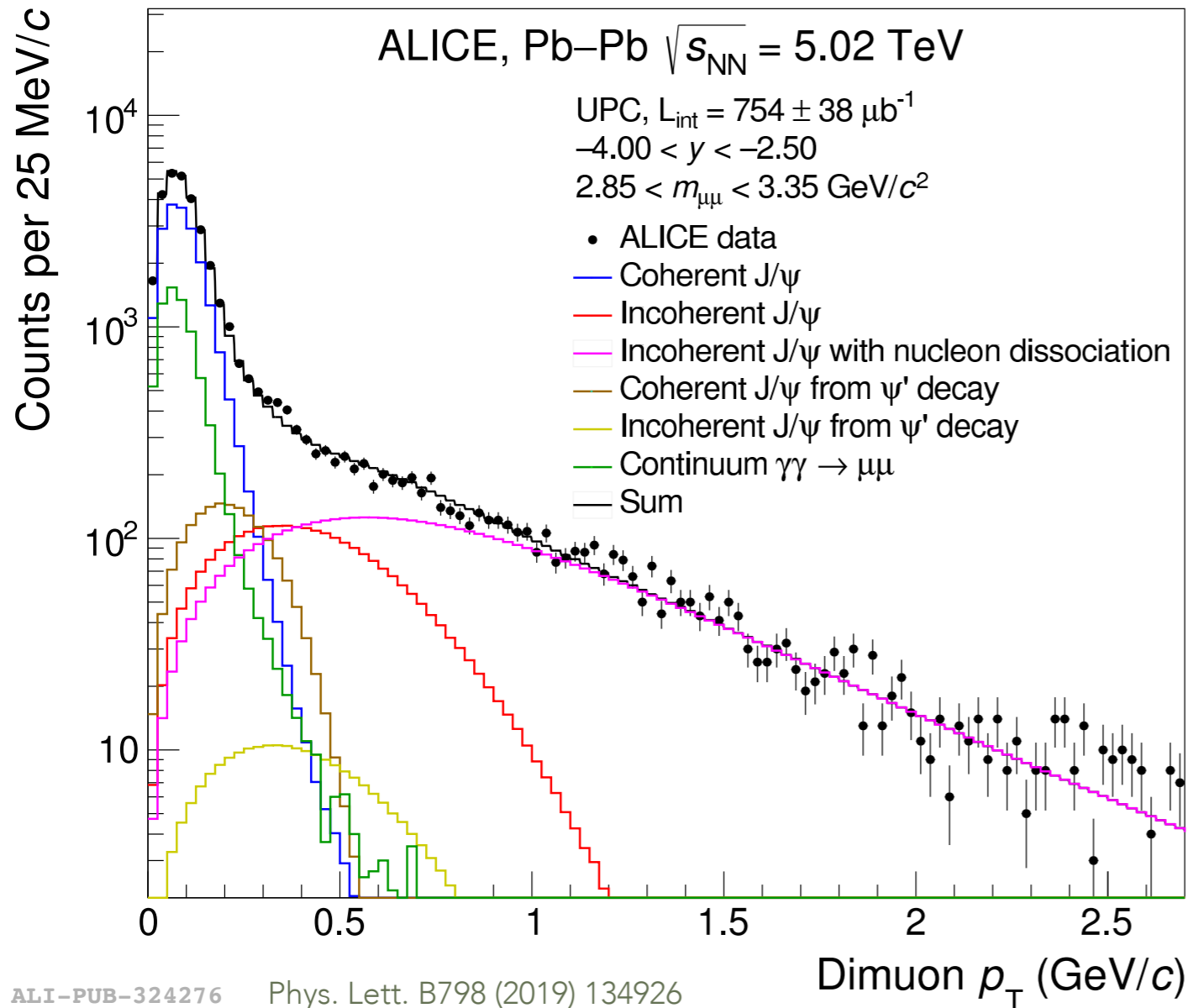


Clear experimental signature :

- Exclusive vector meson with very **low** p_T
- $\langle p_T \rangle \approx 1/R_T$



Coherent J/ψ photoproduction in UPC



Clear experimental signature :

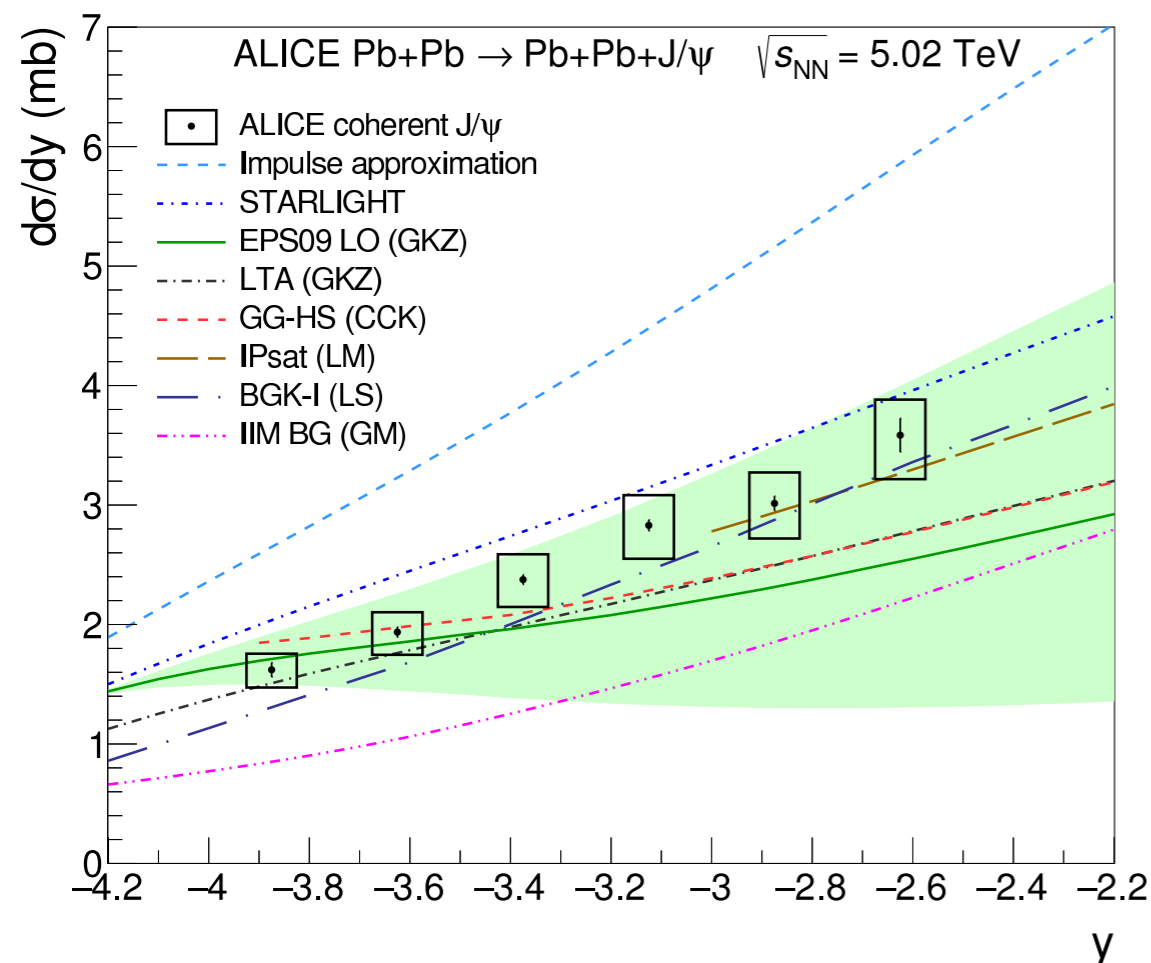
- Exclusive vector meson with very **low** p_T
- $\langle p_T \rangle \approx 1/R_T$



Distinction between :

- ▶ **Coherent** photoproduction
 - Photon interacts with **all nucleons** in the nucleus
 - $\langle p_T \rangle \simeq 60$ MeV/c
 - Target nucleus don't breaks up
- ▶ **Incoherent** photoproduction
 - Photons interacts with one **single nucleon**
 - $\langle p_T \rangle \simeq 500$ MeV/c
 - Target nucleus usually breaks up

Coherent J/ψ photo production in UPC



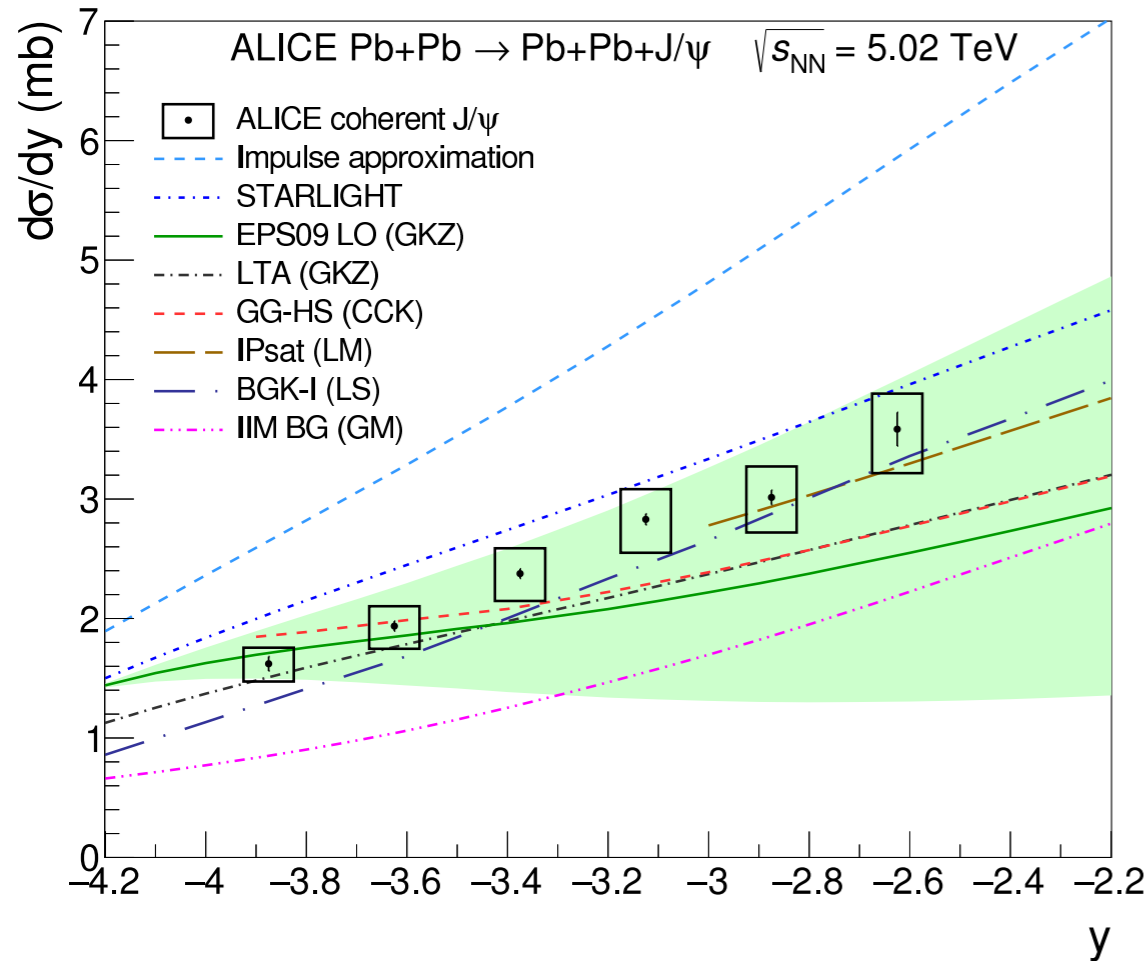
❖ Last results from UPC at 5.02 TeV

- Cross section increase with energy $\sqrt{s_{NN}}$
- Cross section increase with rapidity y

❖ Models

- Without accounting for gluon shadowing overpredict data
- With gluon shadowing underpredict data but remains in theoretical errors

Coherent J/ψ photo production in UPC



❖ Last results from UPC at 5.02 TeV

- Cross section increase with energy $\sqrt{s_{NN}}$
- Cross section increase with rapidity y

❖ Models

- Without accounting for gluon shadowing overpredict data
- With gluon shadowing underpredict data but remains in theoretical errors



ALI-PUB-324284 Phys. Lett. B798 (2019) 134926

For a measure at a given y : cross section is a combination of **two contributions** according to which nucleus emits the γ

$$\frac{d\sigma_{PbPb}}{dy} = n_\gamma(+y, \{b\}) \sigma_{\gamma Pb} (+y) + n_\gamma(-y, \{b\}) \sigma_{\gamma Pb} (-y)$$

high energy γ

$x = 10^{-5}$

low energy γ

$x = 10^{-2}$

Phys. Rev. C 96 (2017) 015203



Conclusion on motivations

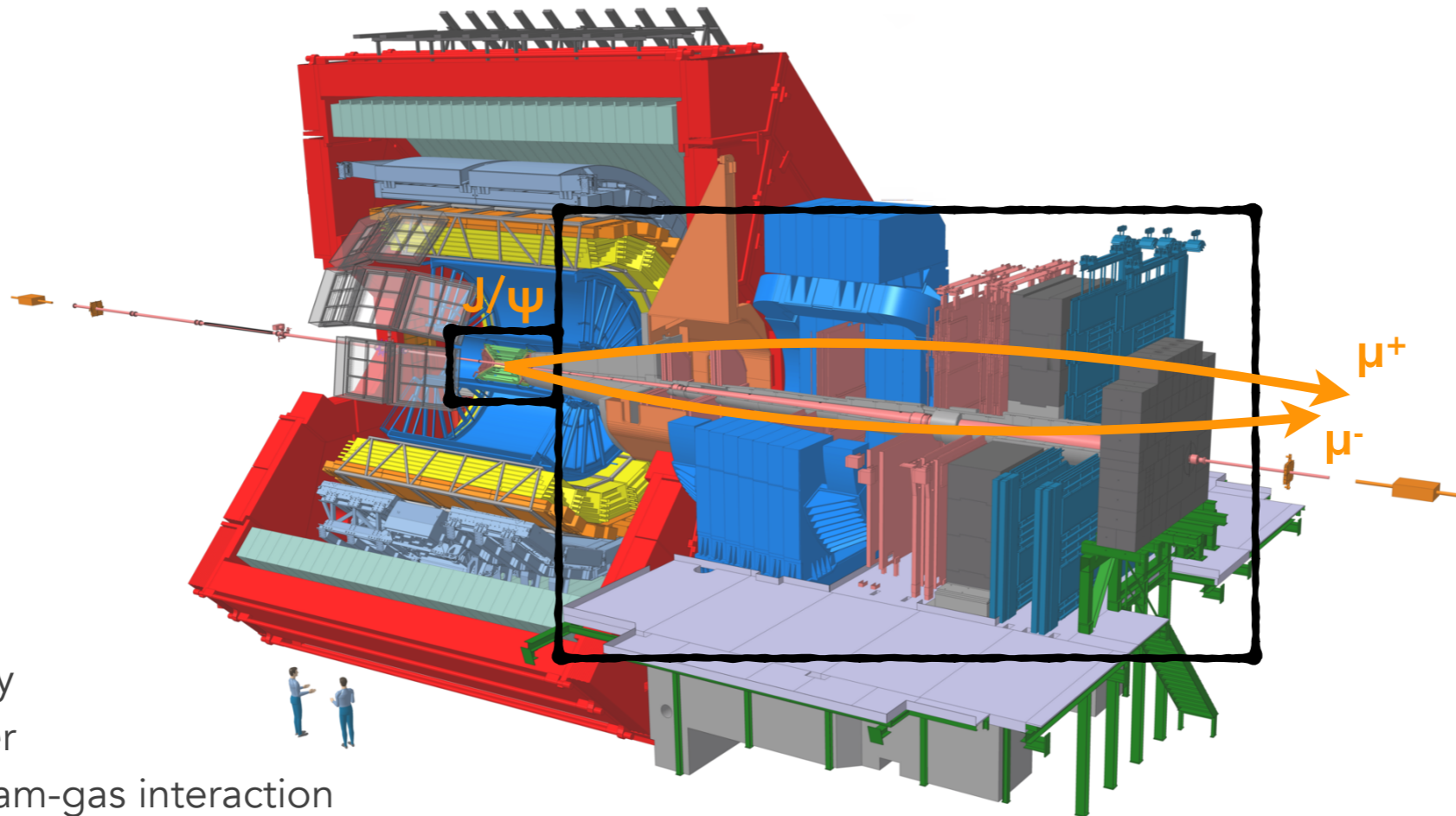
- ❖ Excess of J/ψ at very low p_T unexpected in Pb-Pb peripheral collisions
Coherent photoproduction suggested as underlying mechanism

- ▶ Affect the R_{AA} and hot medium properties study

- ▶ How the excess can have a similar shape as in UPC despite the **nuclear overlap** ?
How the coherence can exist in hadronic collisions ?
 - Coherence with the entire nucleus ? $\rightarrow \langle p_T \rangle \approx 1/R_{\text{nucleus}}$
 - Coherence with the spectators ? $\rightarrow \langle p_T \rangle \approx 1/R_{\text{spectator}}$

- ▶ Improve the statistics to study low Bjorken- x physics

Pb-Pb collisions at 5.02 TeV in ALICE



▶ V0

- Centrality
- Luminosity
- MB trigger
- Reject beam-gas interaction

▶ SPD

- Vertex reconstruction

▶ ZDC

- Electromagnetic background rejection
- Spectator nucleons detection

▶ Muon spectrometer

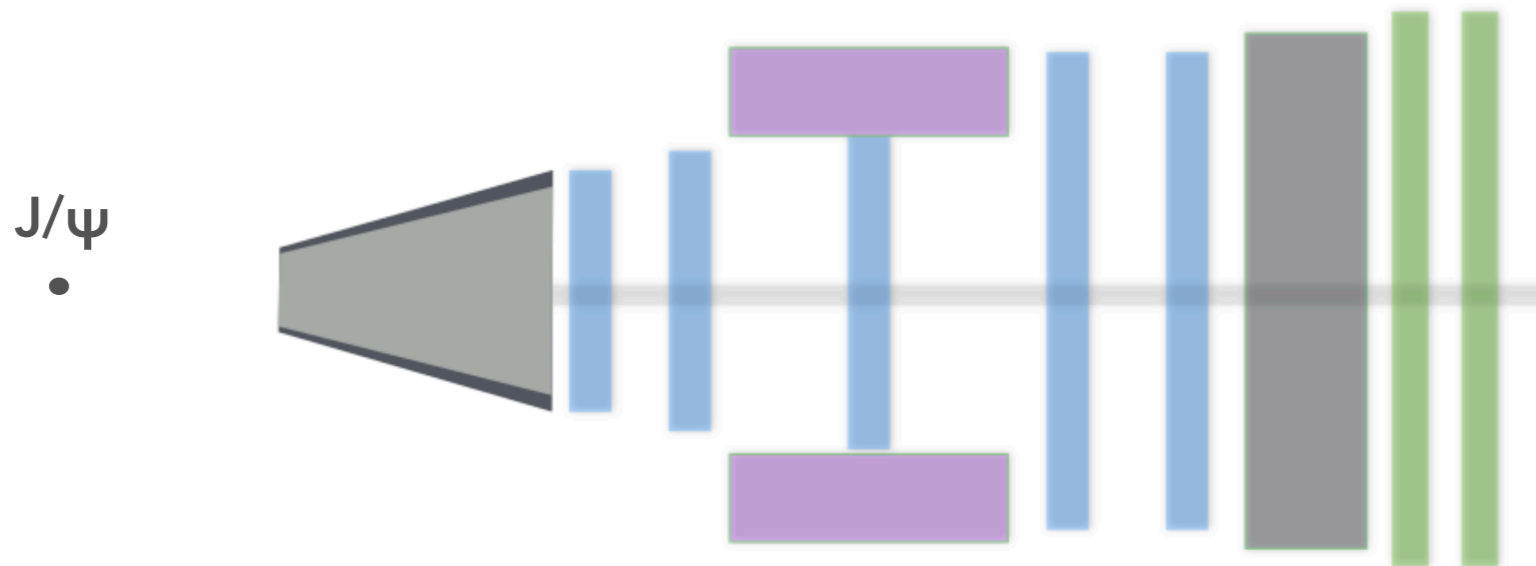
- Front absorber → reduce hadron contamination
- Muon track reconstruction
- Dimuon trigger
→ two opposite-sign muons with $p_T > 1 \text{ GeV}/c$

▶ Acceptance

- $-4 < \eta_\mu < -2.5$
- $0^\circ < \phi_\mu < 360^\circ$

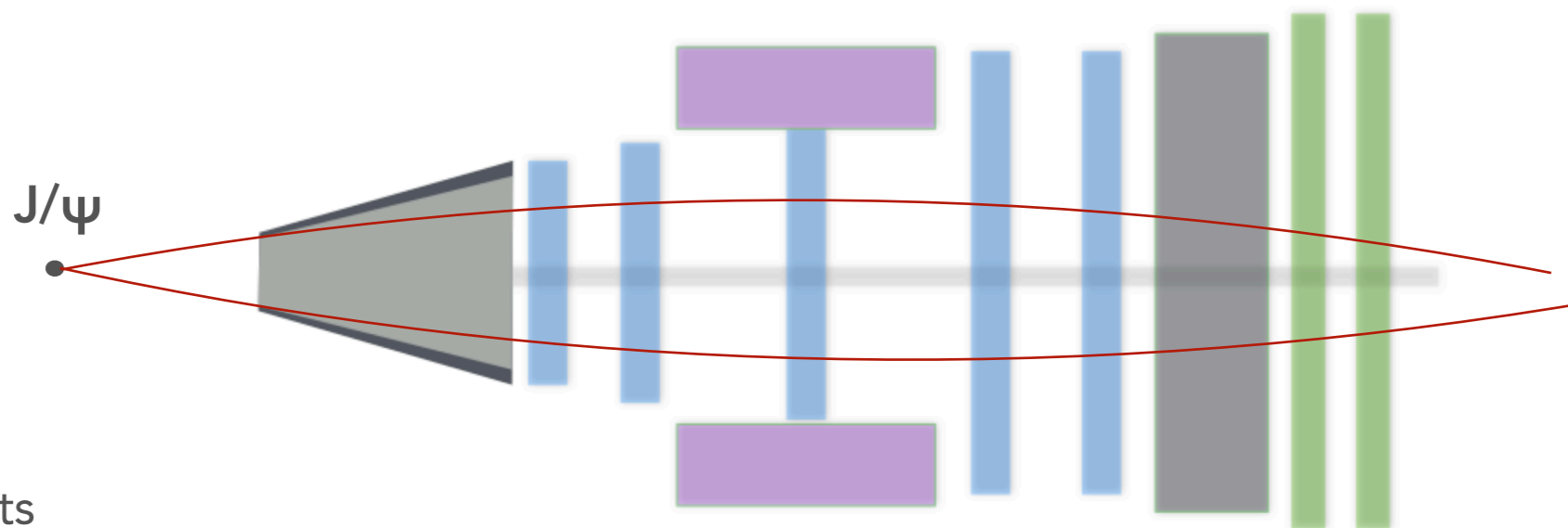
Data sample and event selection

- ▶ Data Sample : Pb-Pb collisions at 5 TeV
 - 2015 : $L_{\text{int}} = 225 \mu\text{b}^{-1}$
 - 2018 : $L_{\text{int}} = 525 \mu\text{b}^{-1}$
- ▶ J/ψ are obtained by combining opposite-sign muons tracks



Data sample and event selection

- ▶ Data Sample : Pb-Pb collisions at 5 TeV
 - 2015 : $L_{\text{int}} = 225 \mu\text{b}^{-1}$
 - 2018 : $L_{\text{int}} = 525 \mu\text{b}^{-1}$
- ▶ J/ψ are obtained by combining opposite-sign muons tracks



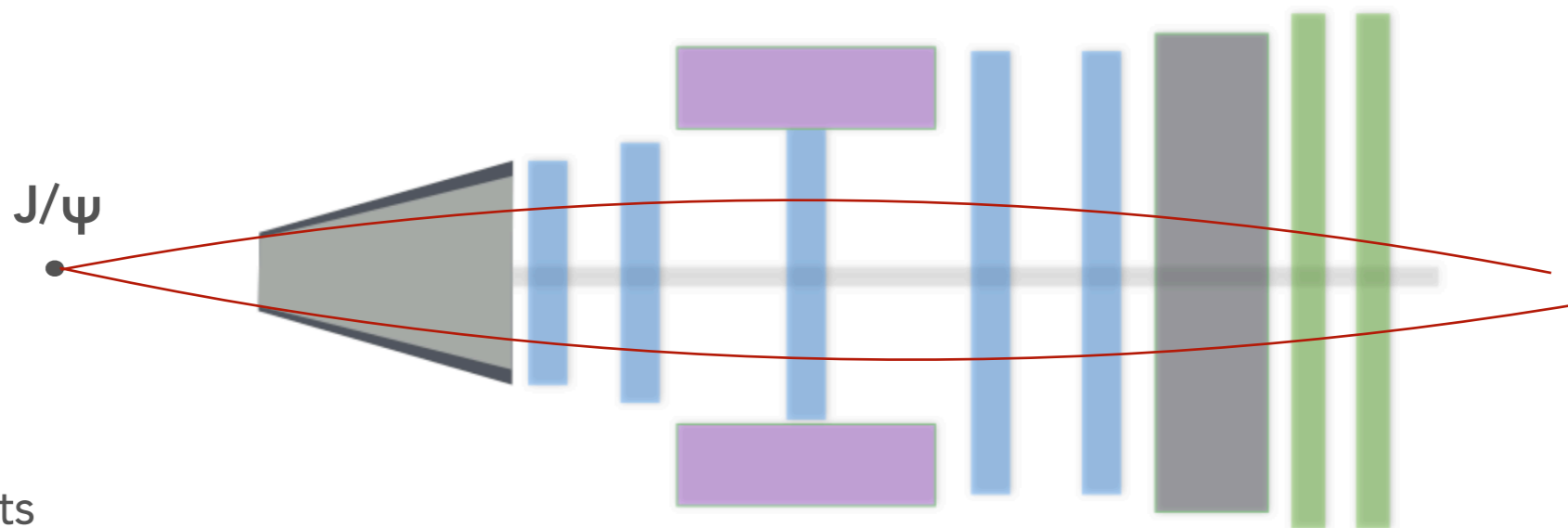
- ▶ Single muon cuts
 - Matched tracks

Data sample and event selection

► Data Sample : Pb-Pb collisions at 5 TeV

- 2015 : $L_{\text{int}} = 225 \mu\text{b}^{-1}$
- 2018 : $L_{\text{int}} = 525 \mu\text{b}^{-1}$

► J/ψ are obtained by combining opposite-sign muons tracks



► Single muon cuts

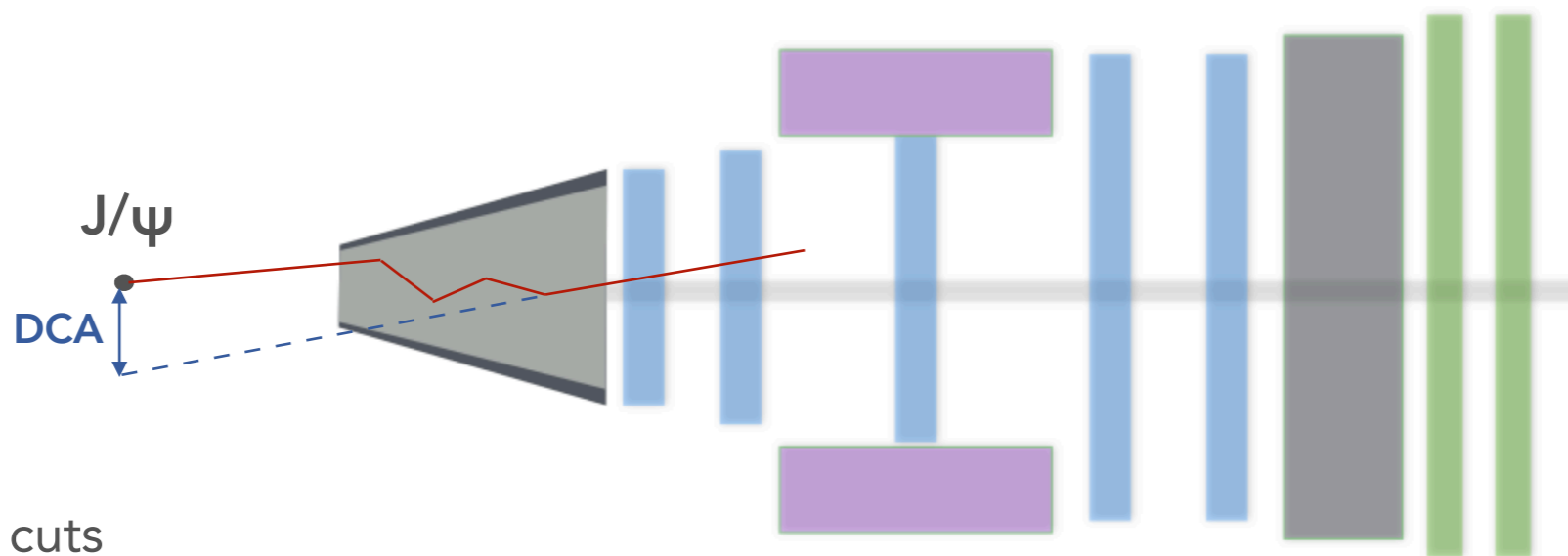
- Matched tracks
- Pseudo-rapidities : $-4 < \eta_{\mu} < -2.5$

Data sample and event selection

► Data Sample : Pb-Pb collisions at 5 TeV

- 2015 : $L_{\text{int}} = 225 \mu\text{b}^{-1}$
- 2018 : $L_{\text{int}} = 525 \mu\text{b}^{-1}$

► J/ ψ are obtained by combining opposite-sign muons tracks



► Single muon cuts

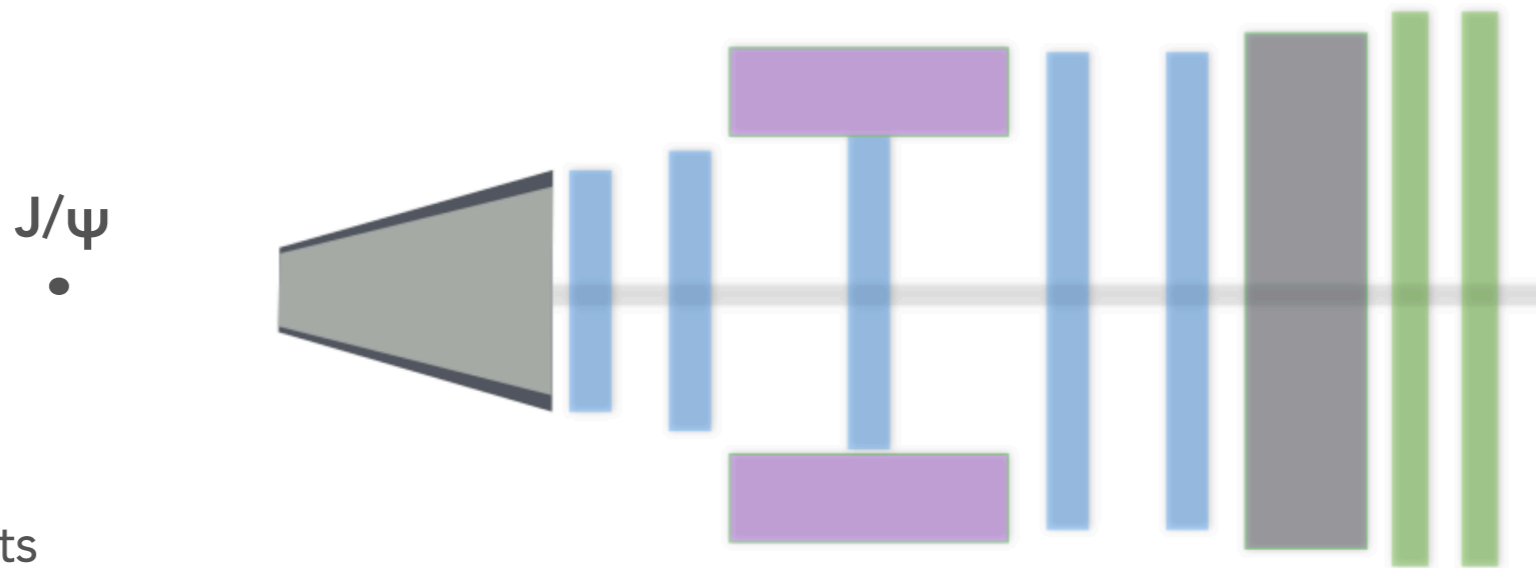
- Matched tracks
- Pseudo-rapidities : $-4 < \eta_{\mu} < -2.5$
- Angle at the end of the front absorber : $2^{\circ} < \theta_{\text{abs}} < 10^{\circ}$
- pDCA

Data sample and event selection

► Data Sample : Pb-Pb collisions at 5 TeV

- 2015 : $L_{\text{int}} = 225 \mu\text{b}^{-1}$
- 2018 : $L_{\text{int}} = 525 \mu\text{b}^{-1}$

► J/ψ are obtained by combining opposite-sign muons tracks



► Single muon cuts

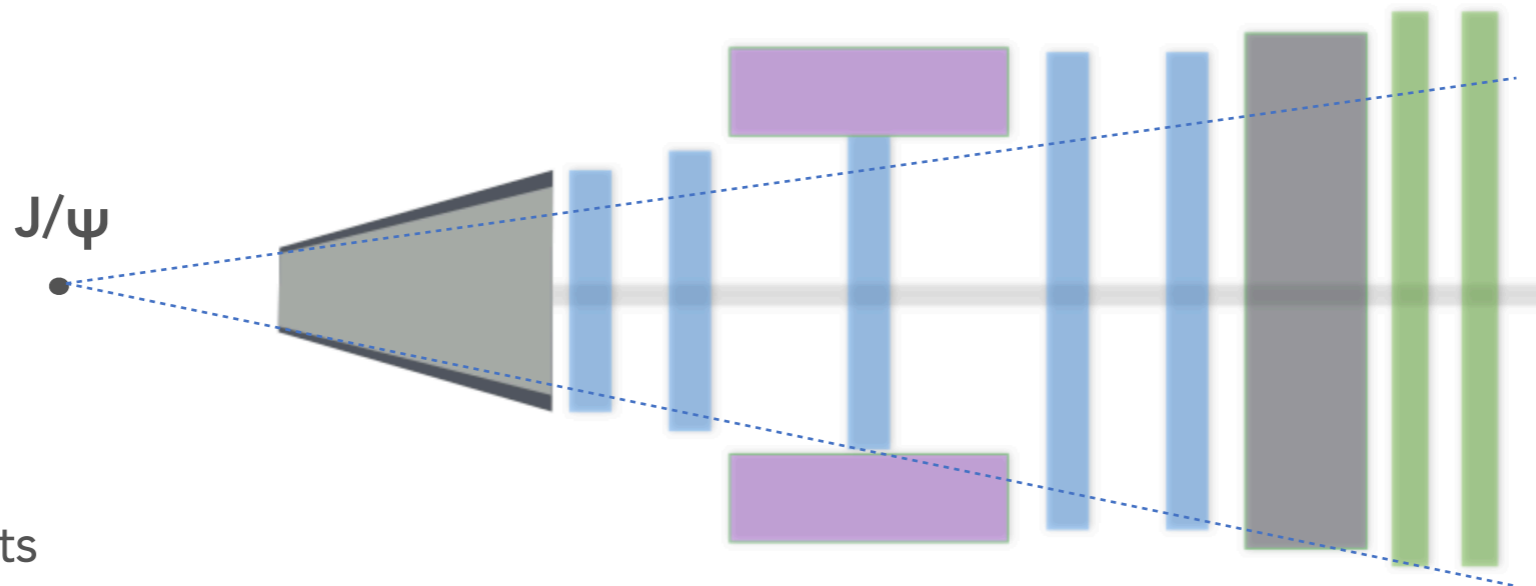
- Matched tracks
- Pseudo-rapidities : $-4 < \eta_{\mu} < -2.5$
- Angle at the end of the front absorber : $2^{\circ} < \theta_{\text{abs}} < 10^{\circ}$
- pDCA

Data sample and event selection

▶ Data Sample : Pb-Pb collisions at 5 TeV

- 2015 : $L_{\text{int}} = 225 \mu\text{b}^{-1}$
- 2018 : $L_{\text{int}} = 525 \mu\text{b}^{-1}$

▶ J/ψ are obtained by combining opposite-sign muons tracks



▶ Single muon cuts

- Matched tracks
- Pseudo-rapidities : $-4 < \eta_{\mu} < -2.5$
- Angle at the end of the front absorber : $2^{\circ} < \theta_{\text{abs}} < 10^{\circ}$
- pDCA

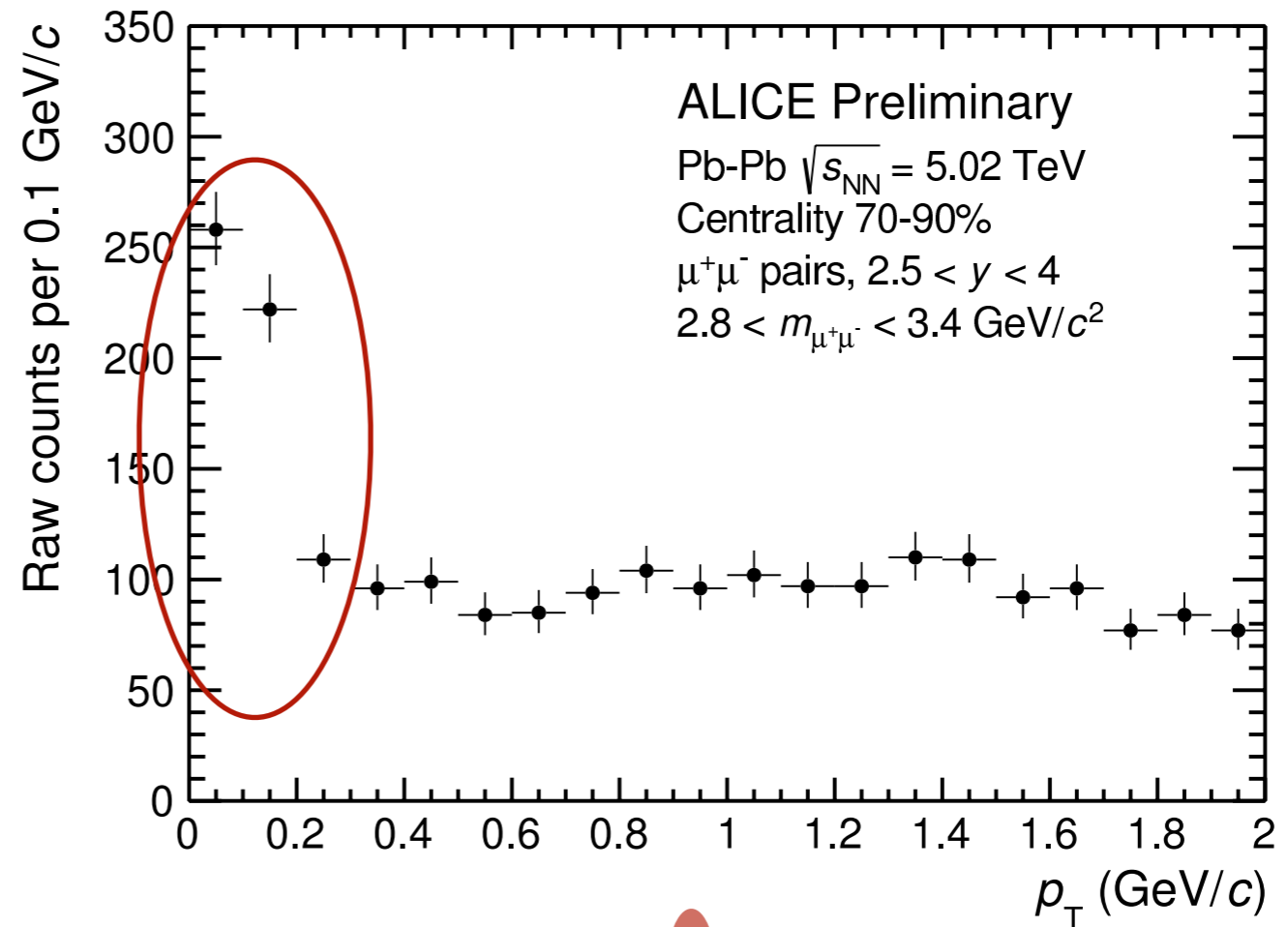
▶ Reconstructed dimuon cut

- Rapidity : $-4 < y_{\text{lab}} < -2.5$

Transverse momentum distribution

▶ Looking at the p_T distribution, the excess is also observed in peripheral collisions at 5.02 TeV

- In J/ψ mass range
 $2.85 < m_{\mu^+\mu^-} < 3.35 \text{ GeV}/c^2$
- For $p_T < 0.3 \text{ GeV}/c$
- Observed only for centrality class from 50% to 90%

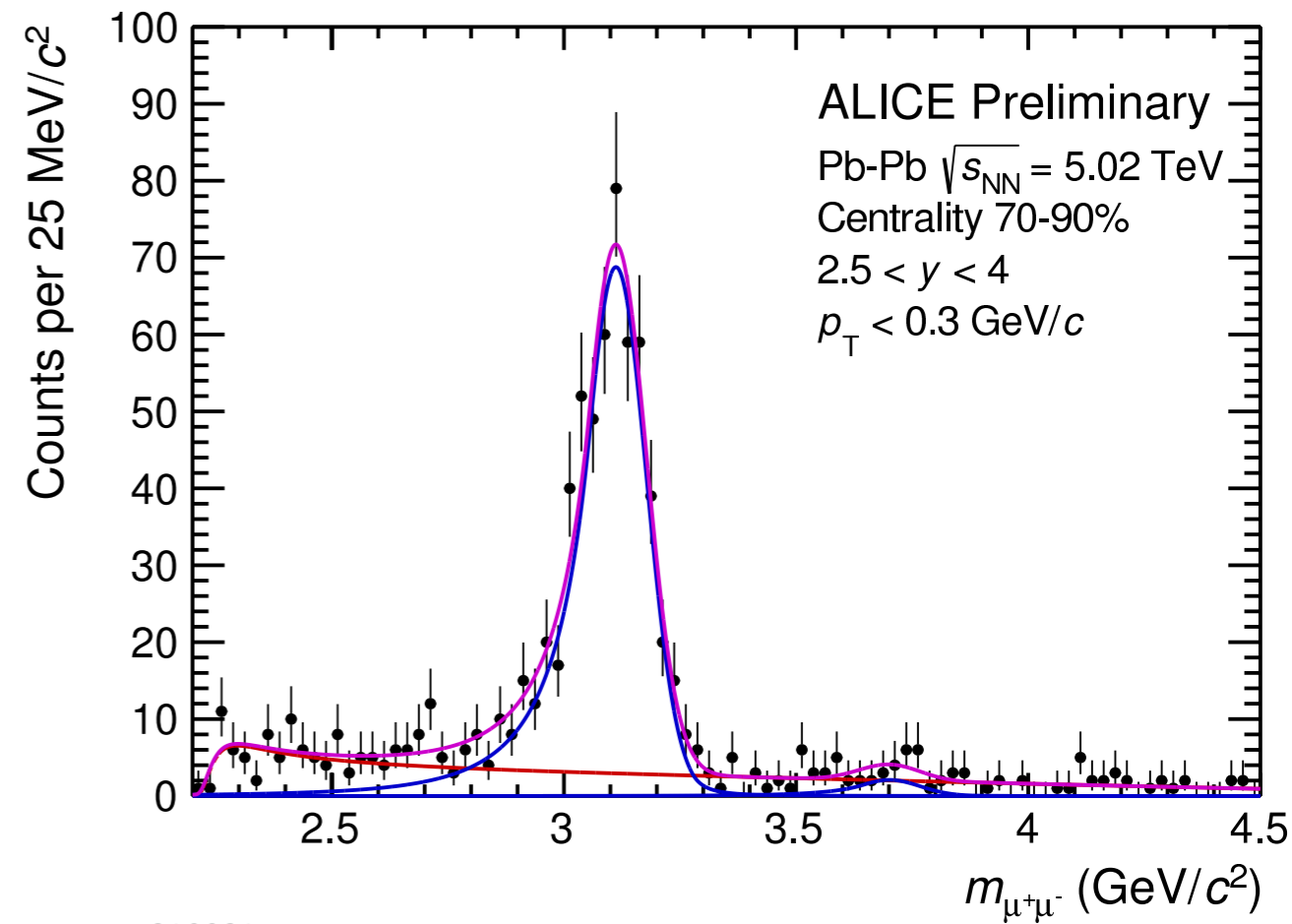


ALI-PREL-309896

Signal extraction

Raw J/ψ number is extracting by fitting the opposite-sign invariant mass distribution

- 2 functions for signal and background + 3 sets of parameters



ALI-PREL-309920

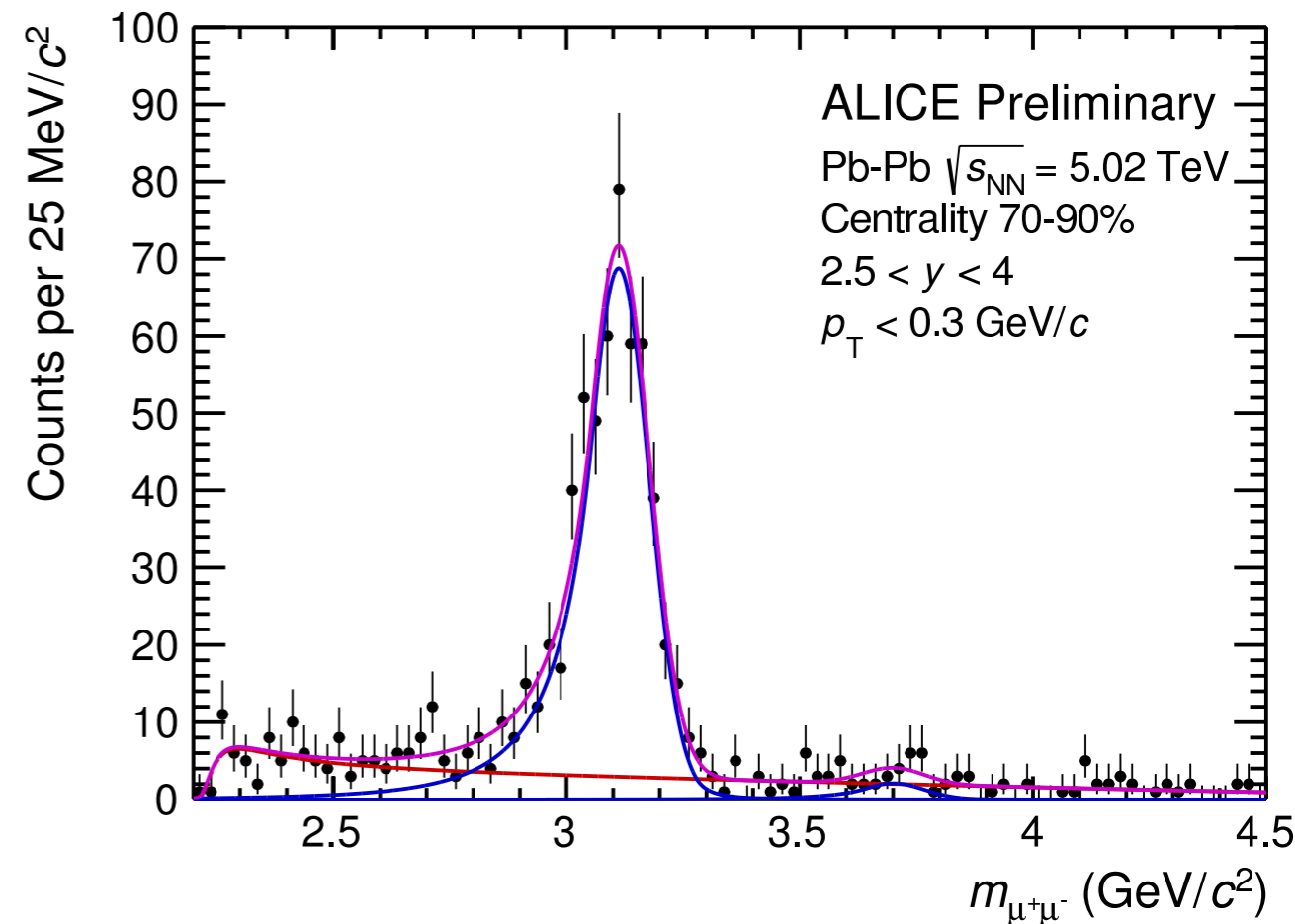
Signal extraction

Raw J/ψ number is extracting by fitting the opposite-sign invariant mass distribution

- 2 functions for signal and background + 3 sets of parameters

In 3 p_T ranges to study different J/ψ production processes :

- Coherent production dominant : $p_T < 0.3$ GeV/c
- Incoherent production dominant : $p_T < 1$ GeV/c
- Hadronic production dominant : $p_T < 8$ GeV/c



ALI-PREL-309920

Estimation of the hadronic component

Number of J/ψ from **hadroproduction** is given by integrating the following parametrization for $p_T < 0.3 \text{ GeV}/c$

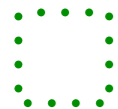
$$\frac{dN_{AA}^{h J/\psi}}{dp_T} = \boxed{\mathcal{N}} \times \boxed{\frac{d\sigma_{pp}^{h J/\psi}}{dp_T}} \times \boxed{R_{AA}^{h J/\psi}} \times \boxed{(\mathcal{A} \times \epsilon)_{AA}^{h J/\psi}}$$



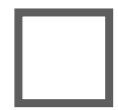
J/ψ cross-section measured by ALICE in pp collisions a 5 TeV



J/ψ nuclear modification factor in Pb-Pb collisions at 5 TeV



Hadronic J/ψ acceptance times efficiency obtained by simulations



Normalization factor as the number of J/ψ in $1 < p_T < 8 \text{ GeV}/c$ where hadronic component is expected to be dominant

$$\mathcal{N} = N_{J/\psi} (1 - 8 \text{ GeV}/c) / \int_{1 \text{ GeV}/c}^{8 \text{ GeV}/c} \frac{d\sigma_{pp}^{J/\psi}}{dp_T} \times R_{AA}^{J/\psi} \times (\mathcal{A} \times \epsilon)_{AA}^{J/\psi} (p_T) dp_T$$



Coherent J/ψ photoproduction cross-section

- ▶ Extraction of the J/ψ number from coherent photoproduction

$$N_{\text{coh}}^{J/\psi} = \frac{N_{\text{excess}}^{J/\psi}}{1 + f_{\text{incoh}} + f_{\text{coh}}^{\psi'}} = \frac{N_{\text{raw}}^{J/\psi} - N_{\text{had}}^{J/\psi}}{1 + f_{\text{incoh}} + f_{\text{coh}}^{\psi'}}$$

f_{incoh} is the fraction of J/ψ from incoherent photoproduction

$f_{\text{coh}}^{\psi'}$ is the fraction of J/ψ from the decay of coherently photoproduced ψ'

- ▶ Cross section evaluation

$$\frac{d\sigma_{\text{coh}}^{J/\psi}}{dp_T} = \frac{N_{\text{coh}}^{J/\psi}}{BR_{J/\psi \rightarrow \mu^+\mu^-} \times (\mathcal{A} \times \epsilon)_{\text{coh}}^{J/\psi} \times \mathcal{L}_{\text{int}} \times \Delta y}$$

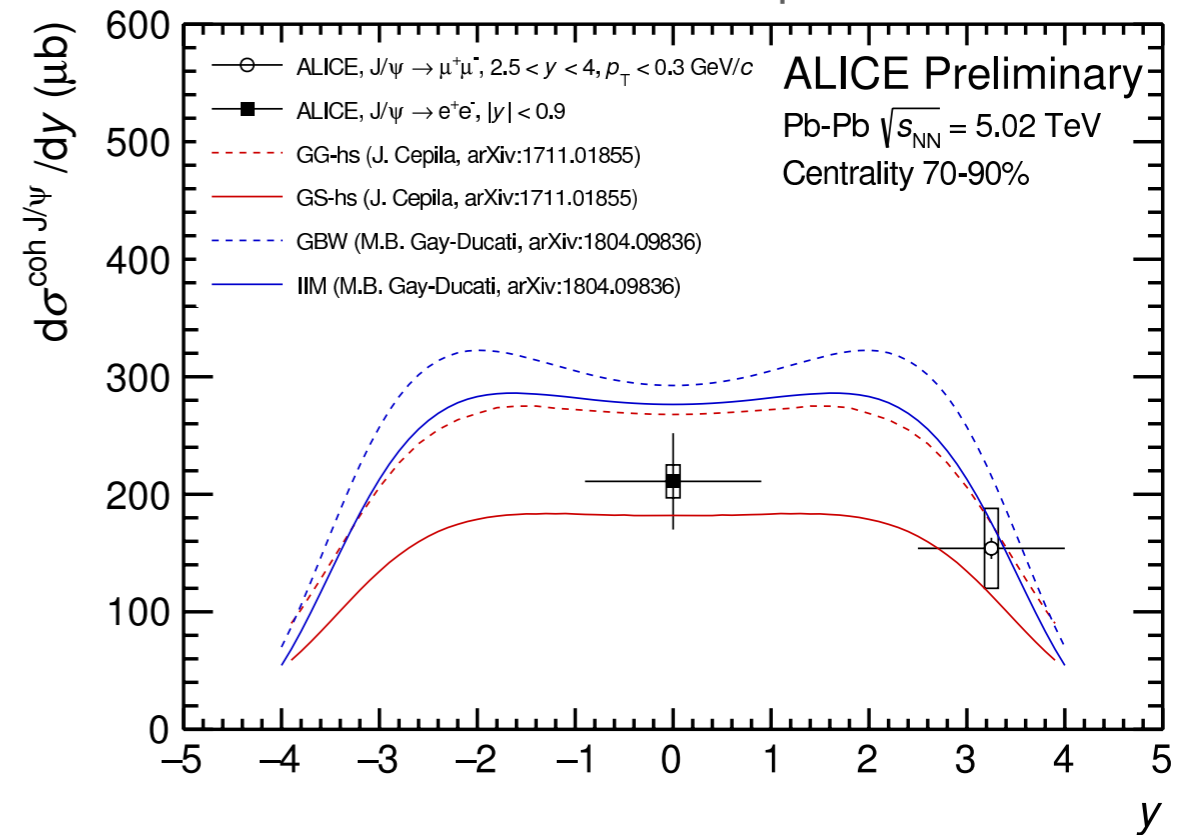
Branching ration = 5.93%

Acceptance x efficiency for coherent J/ψ

Integrated luminosity

Rapidity width

2015 data sample



ALI-PREL-309948

PoS (HardProbes2018) 181

- ▶ Systematics uncertainties determination

Muon tracking efficiency

- ❖ Estimation of the systematics on the tracking efficiency

Three main parts :

I. Cluster map Data/Monte Carlo comparison

- Spot the unexpected detector issues
- Include them in simulations

II. Tuning the kinematics distributions

- Parametrization of the kinematics distributions
- Make the most realistic simulations

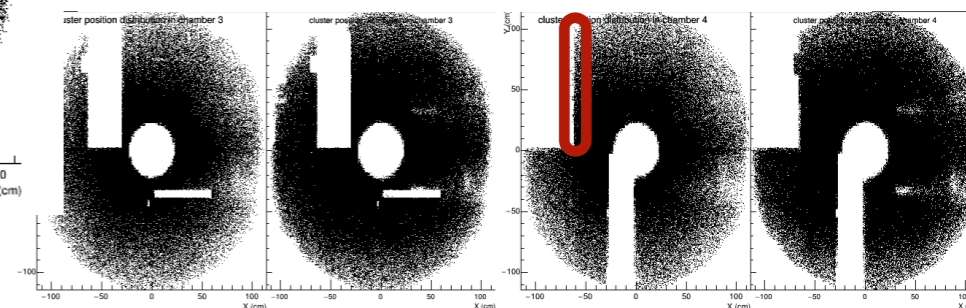
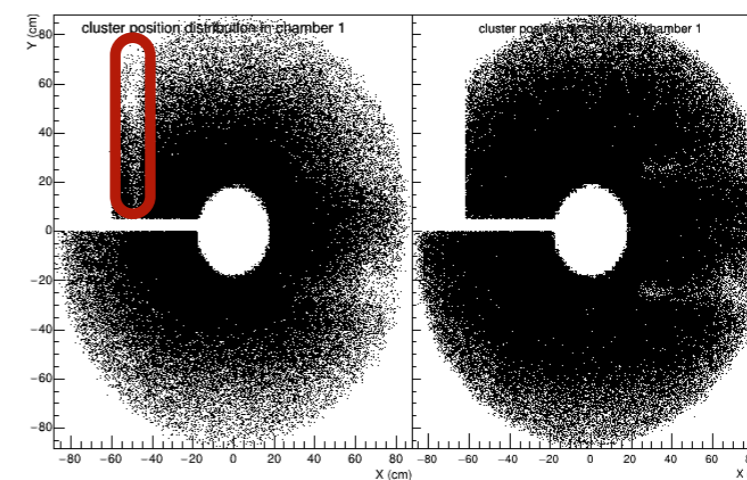
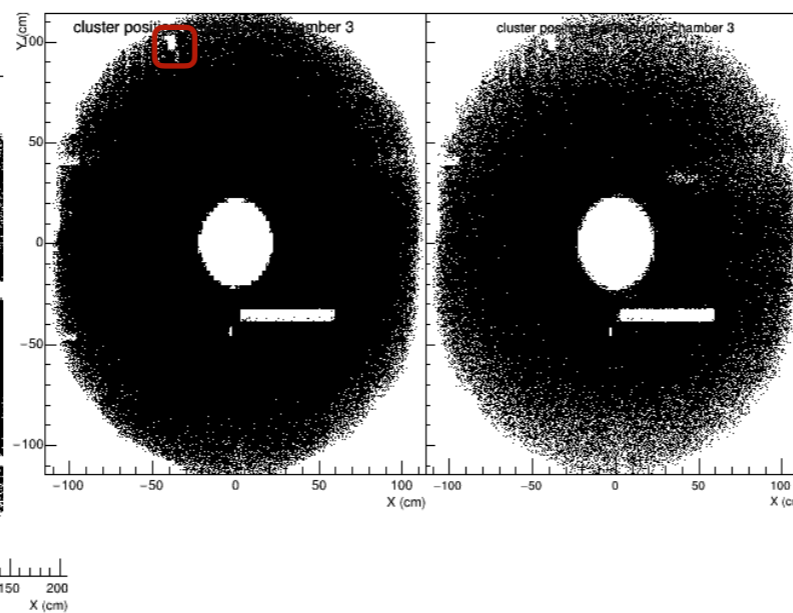
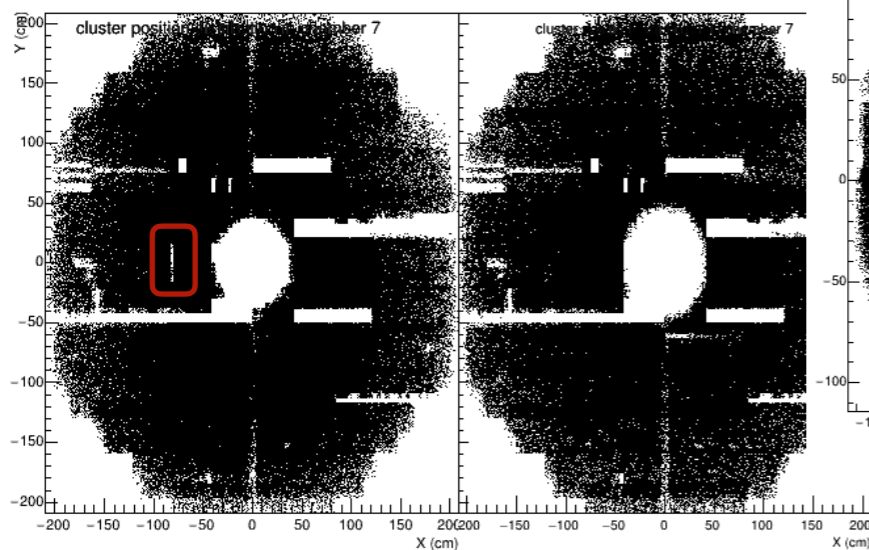
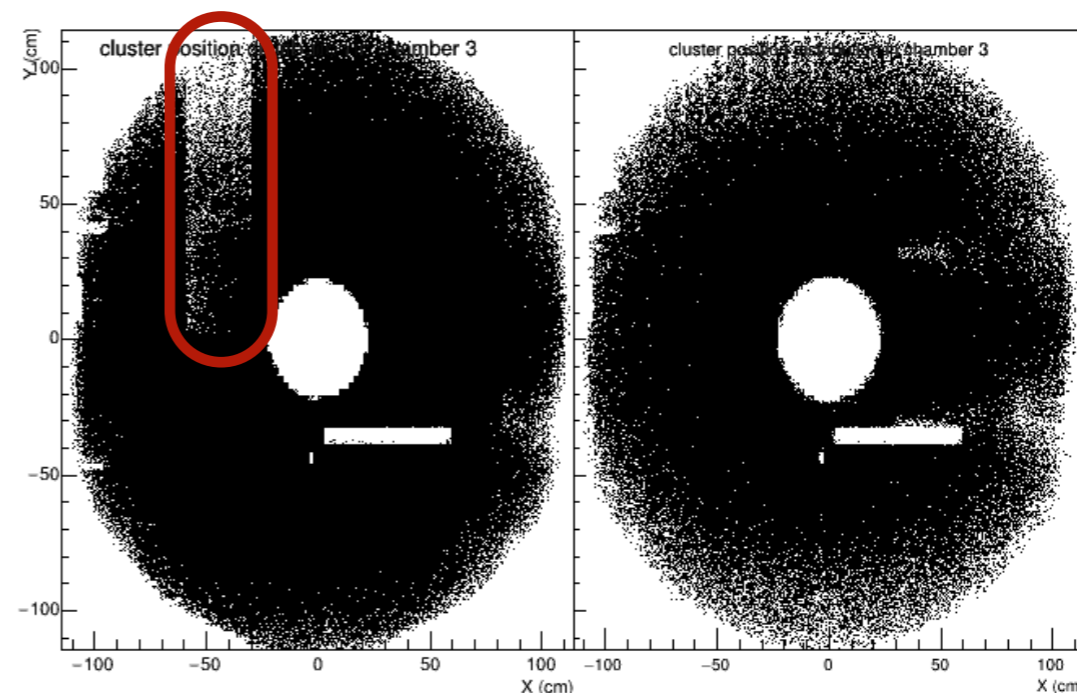
III. Acceptance x Efficiency and systematics

- Compute the Acceptance x Efficiency factor
- Quantify the residual data/MC discrepancies to be used as systematics

Cluster map data/MC comparison

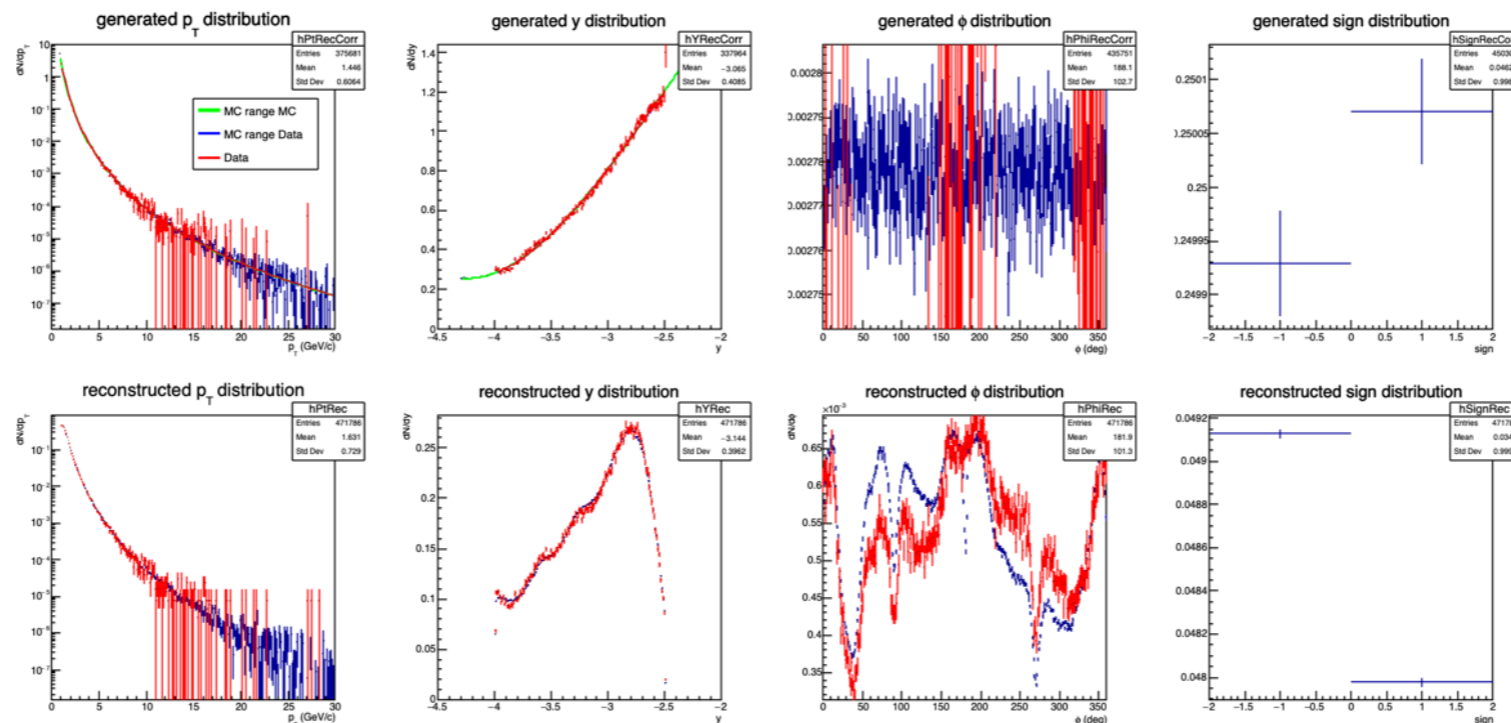
► Method

- Perform single muon simulations
- Compare data and simulations
- Identify **detector issues** not included status map
- **Fix** the problems :
 - detector status and/or offline rejection algorithm
 - add the dead detector elements to the rejectList



Tuning the kinematics distributions : the data driven method

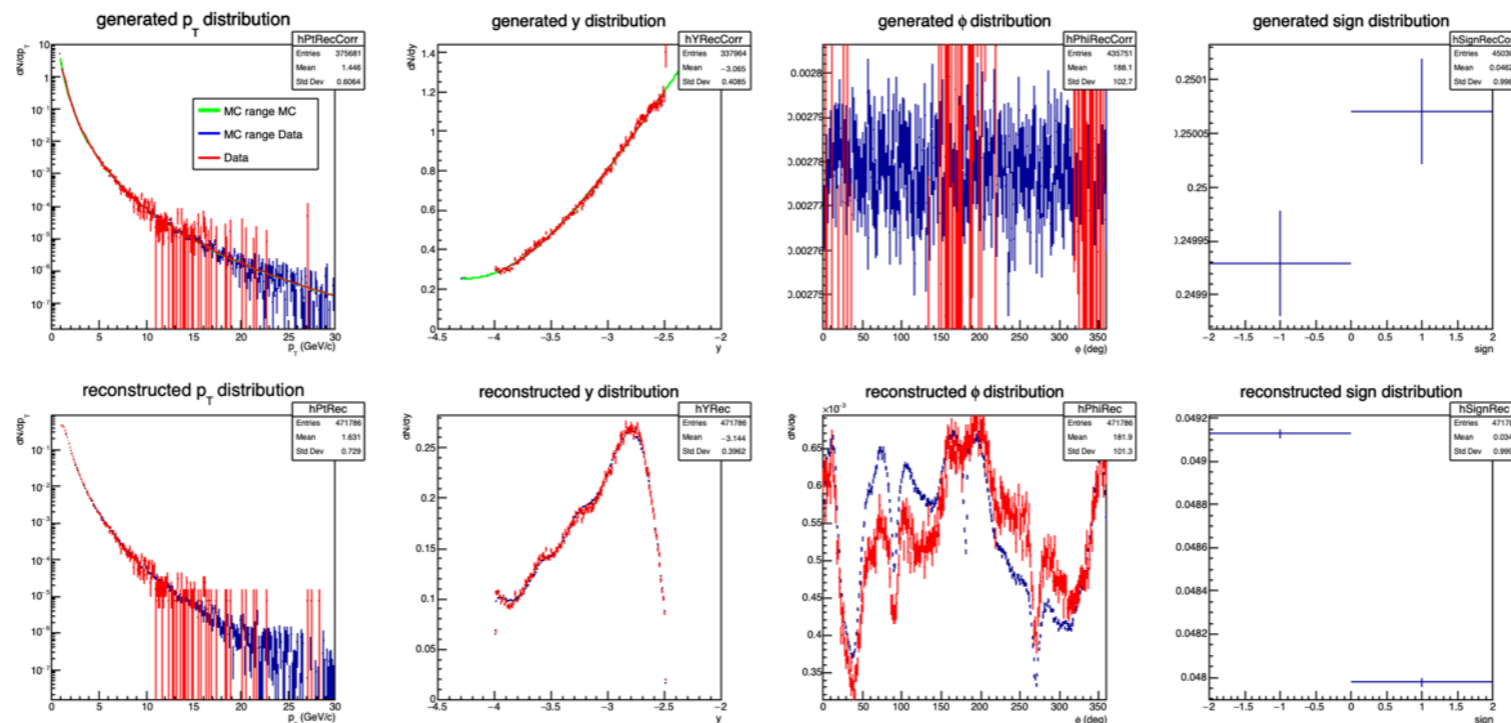
- ▶ Step 0 :
single-muon simulations with parameters from previous period for the kinematic distributions
- ▶ Step $n > 0$:
 - $\mathcal{A} \times \epsilon$ is measured from MC simulation and applied to the data reconstructed distribution
→ equivalent generated distribution
 - Fitting the new corrected data distributions
 - Compute the **weight** :
ratio of the new function to the original generated MC distribution
→ Re-weight the generated muons



Tuning the kinematics distributions : the data driven method

- ▶ Step 0 :
single-muon simulations with parameters from previous period for the kinematic distributions
- ▶ Step $n > 0$:
 - $\mathcal{A} \times \epsilon$ is measured from MC simulation and applied to the data reconstructed distribution
→ **equivalent generated distribution**
 - Fitting the new corrected data distributions
 - Compute the **weight** :
ratio of the new function to the original generated MC distribution
→ Re-weight the generated muons

The procedure is done **iteratively** until parameters for generated MC distributions **converged**



Acceptance times efficiency

► Measure **intrinsic efficiency** :

- Use the reconstructed tracks
- Use the redundancy between chambers
- Assuming the efficiency of one chamber is independent on the others

$$\epsilon_{ch\ i} = \frac{N_{i-j}}{N_{i-j} + N_{0-j}}$$

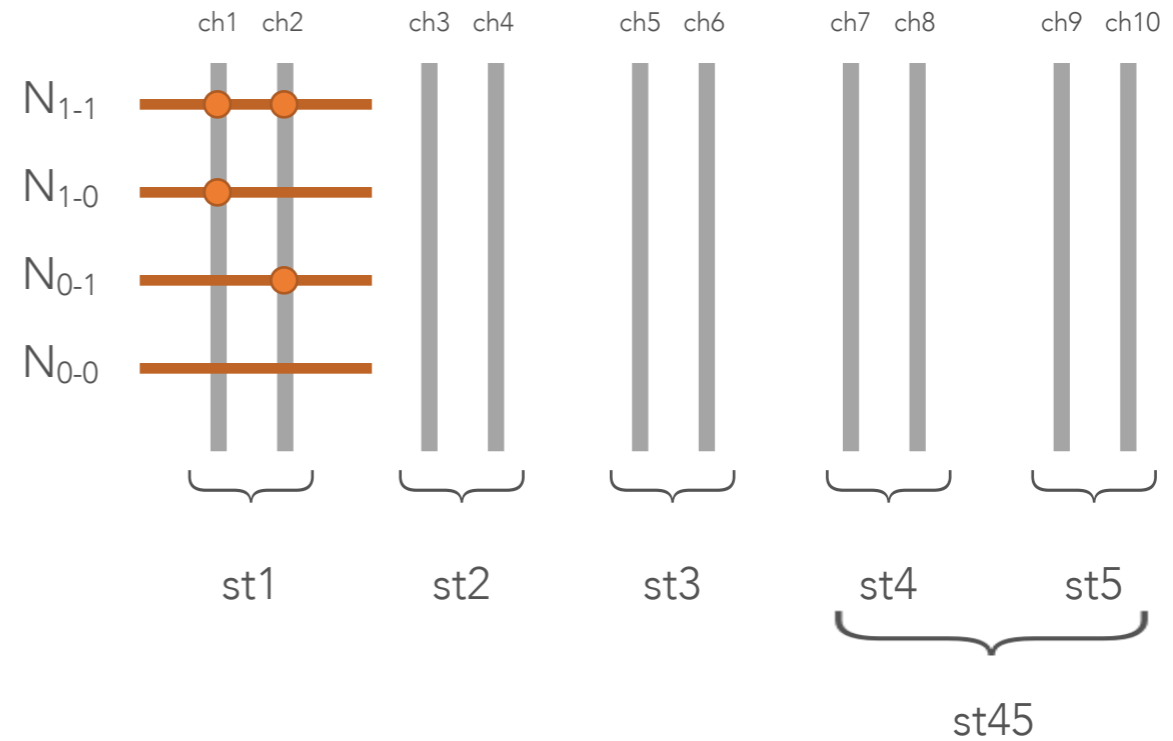
► Method based on the tracking algorithm properties

- For stations 1, 2 and 3 :
at least 1 cluster per station

$$\epsilon_{st\ 1(2)(3)} = 1 - \left(1 - \epsilon_{ch\ 1(3)(5)}\right) \left(1 - \epsilon_{ch\ 2(4)(6)}\right)$$

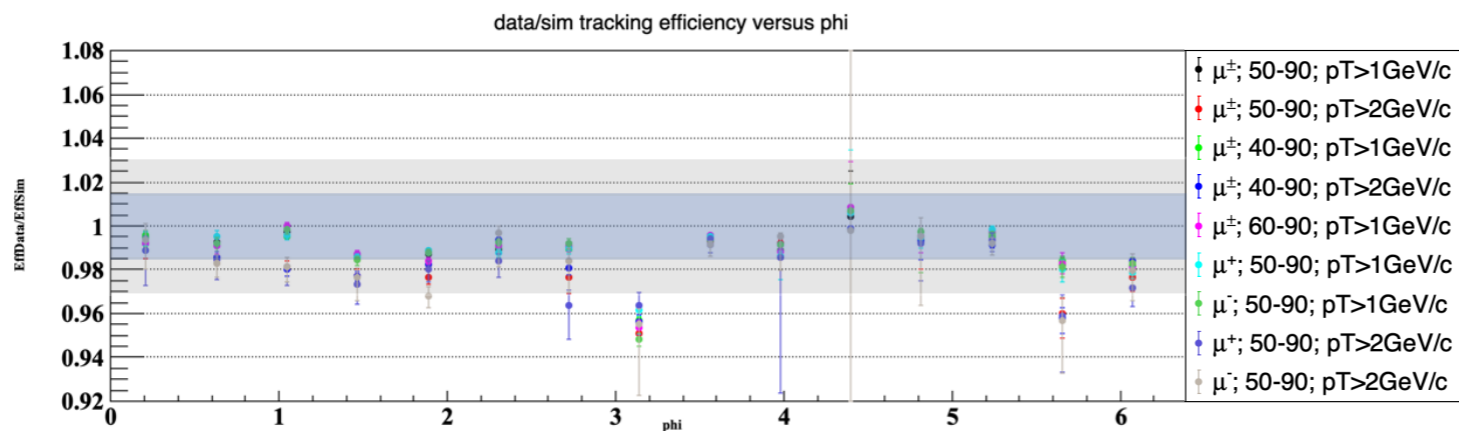
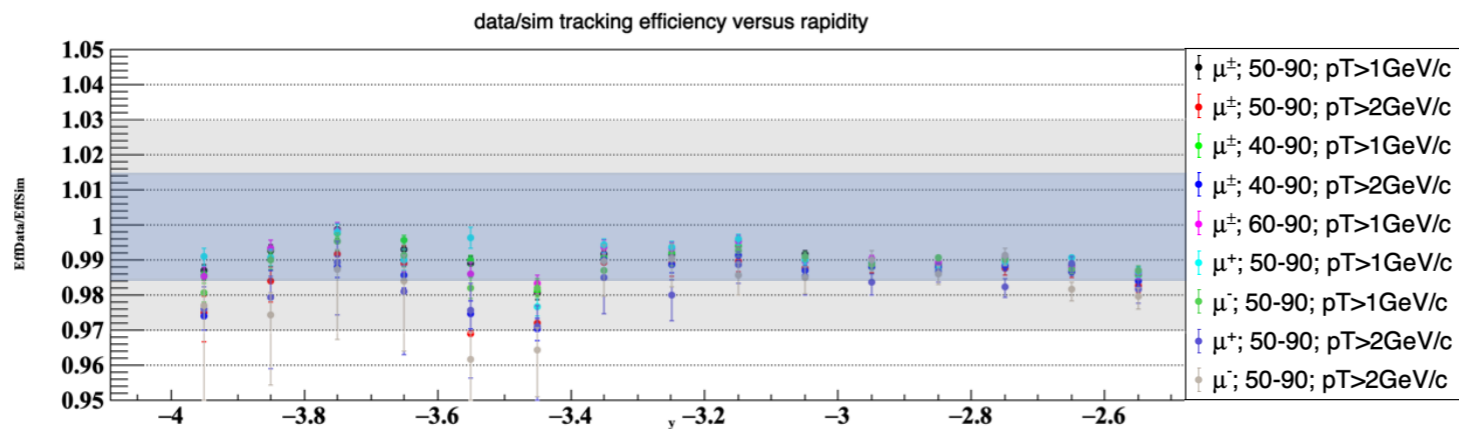
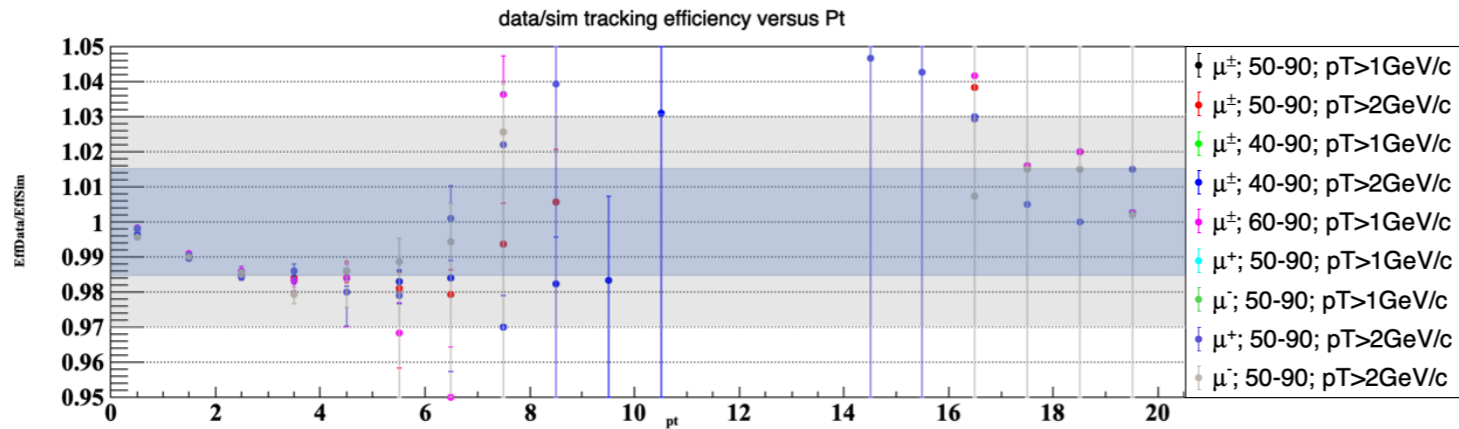
- For stations 4 and 5 :
at least 3 cluster in the last 4 chambers

$$\epsilon_{st\ 45} = \prod_{i=7}^{10} \epsilon_{ch\ i} + \sum_{i=7}^{10} \left((1 - \epsilon_{ch\ i}) \prod_{j=7; j \neq i}^{10} \epsilon_{ch\ j} \right)$$



$$\epsilon_{tracking} = \epsilon_{st1} \epsilon_{st2} \epsilon_{st3} \epsilon_{st45}$$

Tracking systematics



► If we take a systematic of 1.5% all data/MC differences are within $2\sigma - 3\sigma$

→ systematic of 3% for dimuons

Conclusion and outlooks

- ❖ Excess in the yield of J/ψ at very low p_T confirmed in peripheral collisions at 5.02 TeV
- ❖ Excess is expected to be more significant with respect to 2.76 TeV

Analysis in progress

- ▶ Study of the $\langle p_T \rangle$ of the J/ψ excess
- ▶ Will be extended to more central collisions
 - Centrality dependence of the $\langle p_T \rangle$
 - Coherence condition
- ▶ Will be performed for more bins in rapidity
 - Rapidity dependence of the cross-section
 - Combine results with UPC analysis and extract $\sigma_{\text{YpB}}(+y)$ and $\sigma_{\text{YpB}}(-y)$
 - Probe the gluonic content for different Bjorken- x