

# Isospin transport in nuclear collisions studied by multidetectors INDRA-FAZIA

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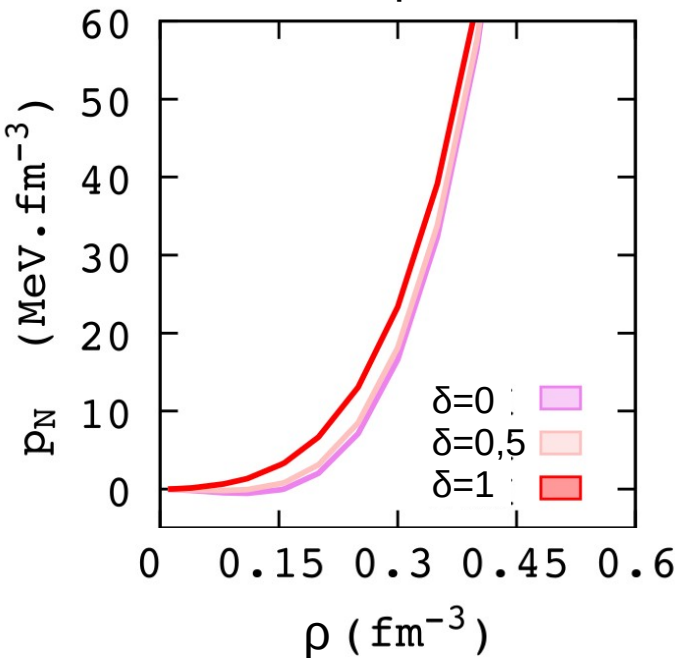
# Introduction : theoretical context

Def. Equation of State (EoS) : Pressure = f(Temperature, Density).

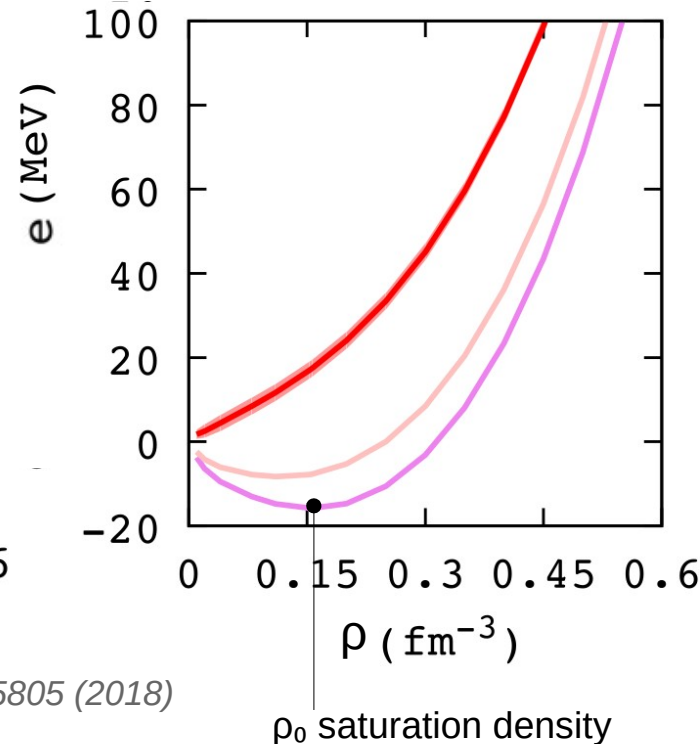
A well-known one is the perfect gas equation  $P = \rho RT$

In nuclear physic : **EoS of nuclear matter**  $P(T, \rho, \delta)$  , with  $\rho$  density of nuclear matter  
 More often written Energy( $T, \rho, \delta$ )

P vs  $\rho$



E vs  $\rho$



for  $\infty$  nuclear matter

$$\delta = \frac{\rho_n - \rho_p}{\rho}$$

for a nucleus

$$\delta = \frac{N - Z}{A}$$

Margueron et al. Physical Review C 97, 025805 (2018)

To get EoS's parameters we use a limited expansion

expanded in power of  $\delta$

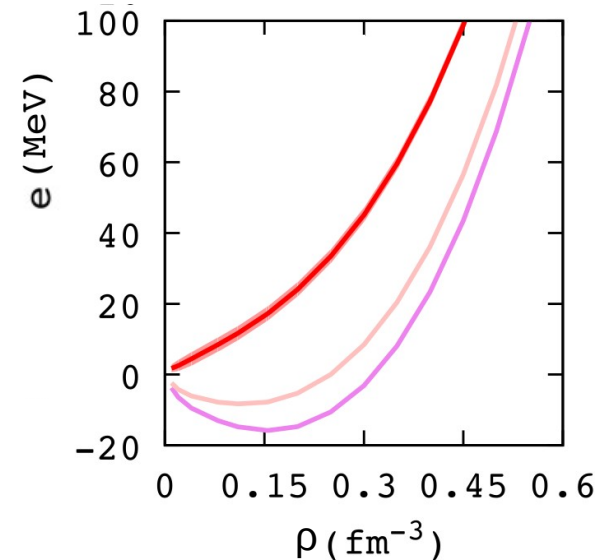
$$e(\rho, \delta) = e_{is}(\rho) + e_{iv}(\rho)\delta^2 + O(\delta^4)$$

expanded in power of  $x = (\rho - \rho_0)/(3\rho_0)$  around saturation density  $\rho = \rho_0$

$$e_{is}(x) = E_{sat} + \frac{1}{2}K_{sat}x^2 + O(x^4) \text{ well-known at first orders}$$

$$e_{iv}(x) = E_{sym} + L_{sym}x + \frac{1}{2}K_{sym}x^2 + O(x^4)$$

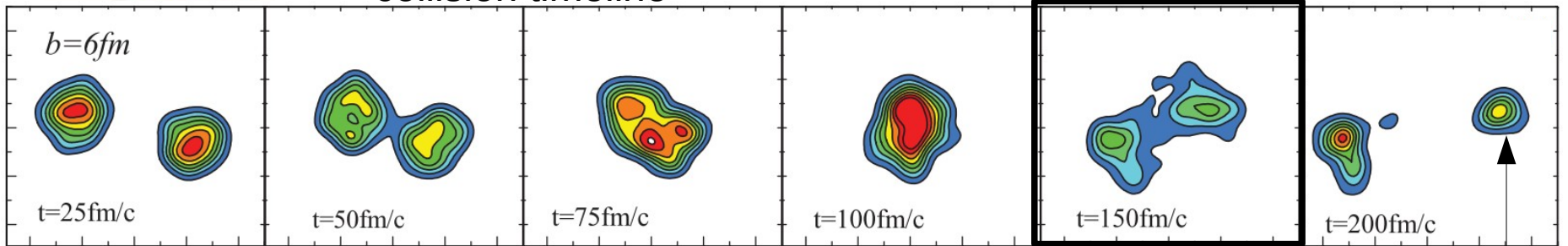
Parameter	Relative uncertainty
$E_{sym}$	5 %
$L_{sym}$	30 %
$K_{sym}$	100 %





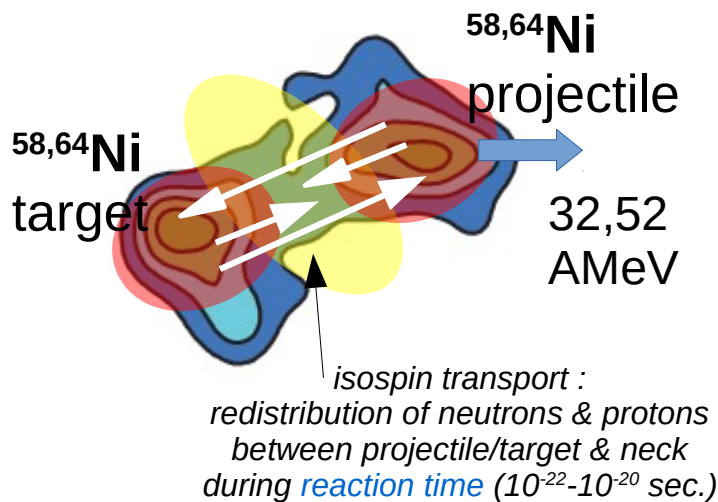
# Experimental way to constrain these parameters

collision timeline



Yingxun Zhang, et al. *Physical Review C*, 85(2)

Quasi-Projectile



Requirements :  
**impact parameter  $b$** , deduced from fragments multiplicity  
**reaction time**, deduced from  $b$  and  $E(\text{projectile})$

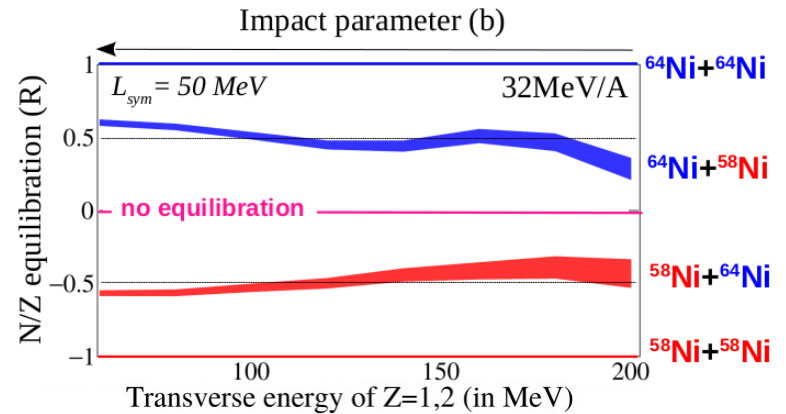
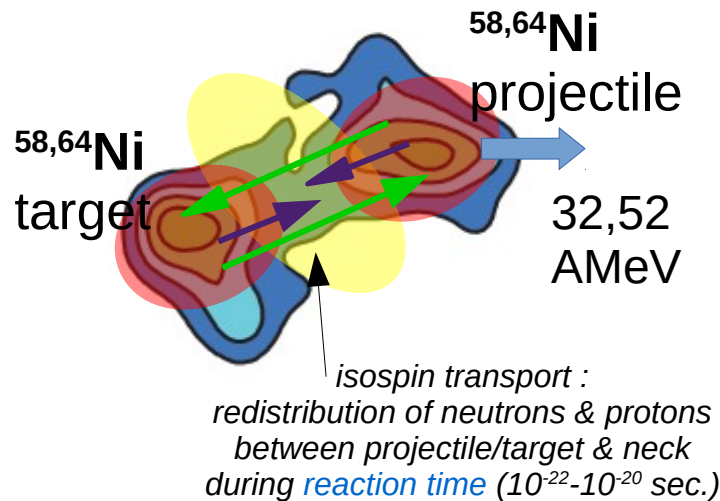
neutrons stream and protons stream difference  $|\vec{j}_n - \vec{j}_p|$ ,

deduced from  $\left(\frac{N}{Z}\right)_{ini} - \left(\frac{N}{Z}\right)_{final}$

$$|\vec{j}_n - \vec{j}_p| \propto \overbrace{L_{sym} \vec{\nabla} \rho}^{\text{migration}} - \overbrace{e_{iv}(\rho) \vec{\nabla} \delta}^{\text{diffusion}}$$

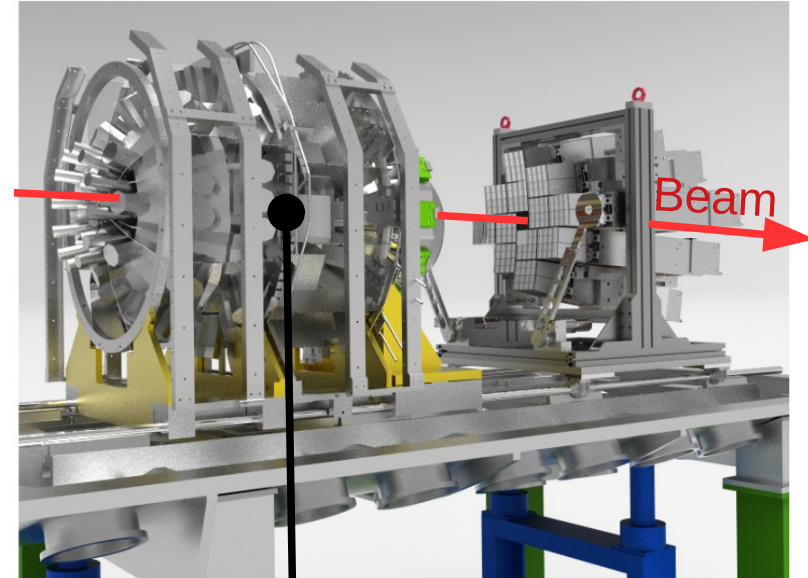
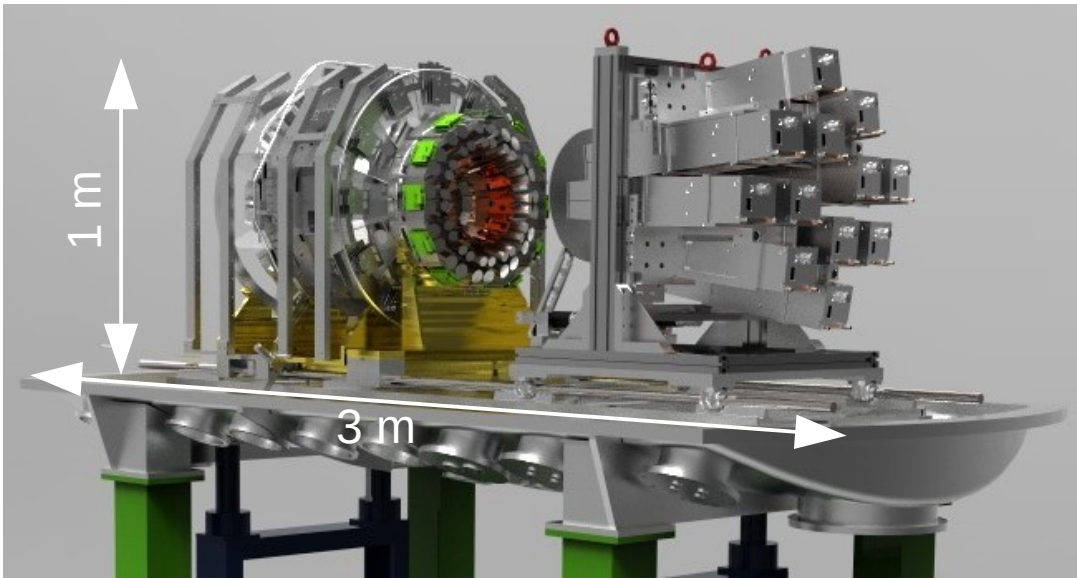
neutrons enrichment of the neck  
change with the couple chosen target/projectile couple

equilibration isotopic  
Change with the collision energy



Lopez and Piantelli, E789 proposal

# The multidetector : INDRA-FAZIA



**Target**

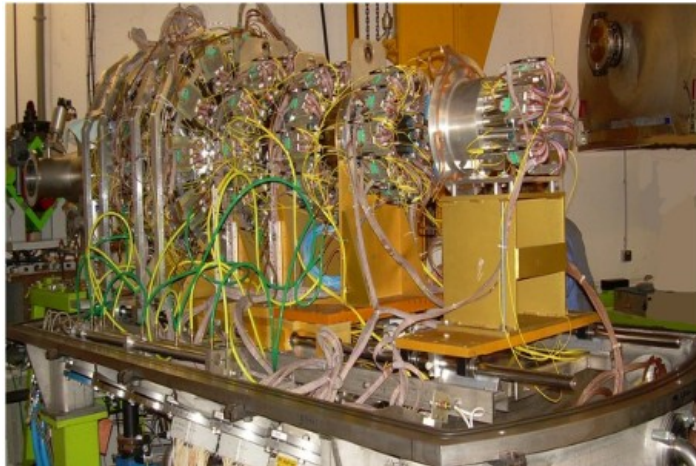
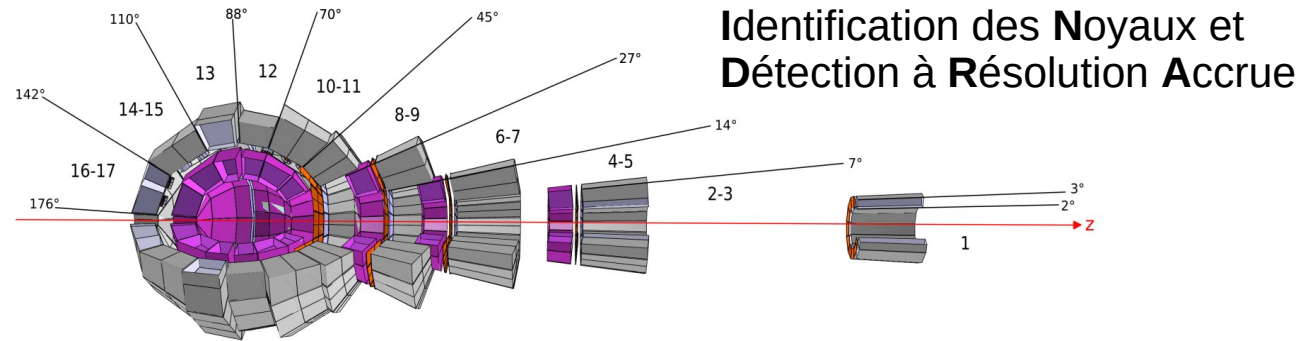


almost  $4 \pi$  angular coverage

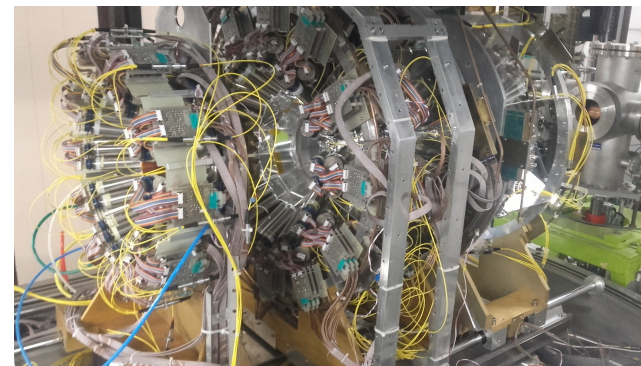
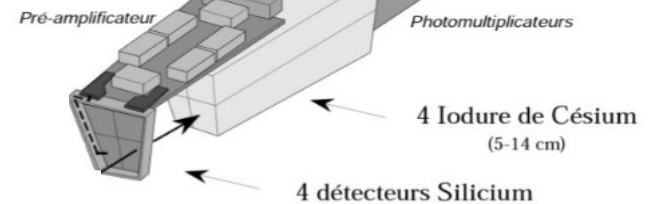
A & Z Quasi-Projectile well-measured

# The multidetectors : INDRA

25 years  
analogic electronic



## Module NDRA



12 rings left = 240 telescopes  
Angular coverage 14° to 176°

Z identification up to  $Z = 92$   
Isotopic identification up to  $Z = 6$

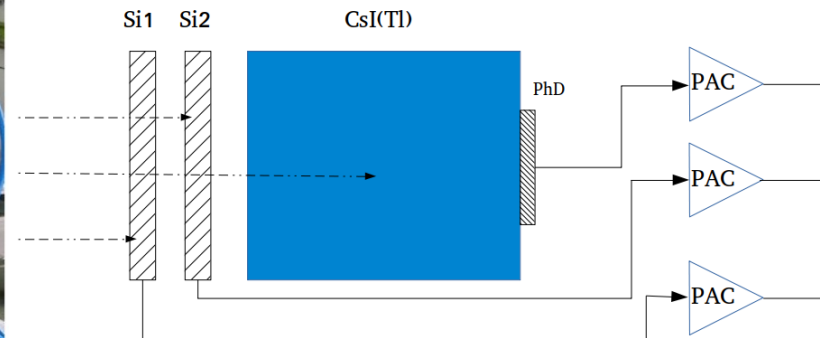


## Forward **A** and **Z** Identification **A**rray



12 blocks x 16 telescopes =  
192 telescopes = 576 detectors

Angular coverage from  $1,5^\circ$  to  $14^\circ$

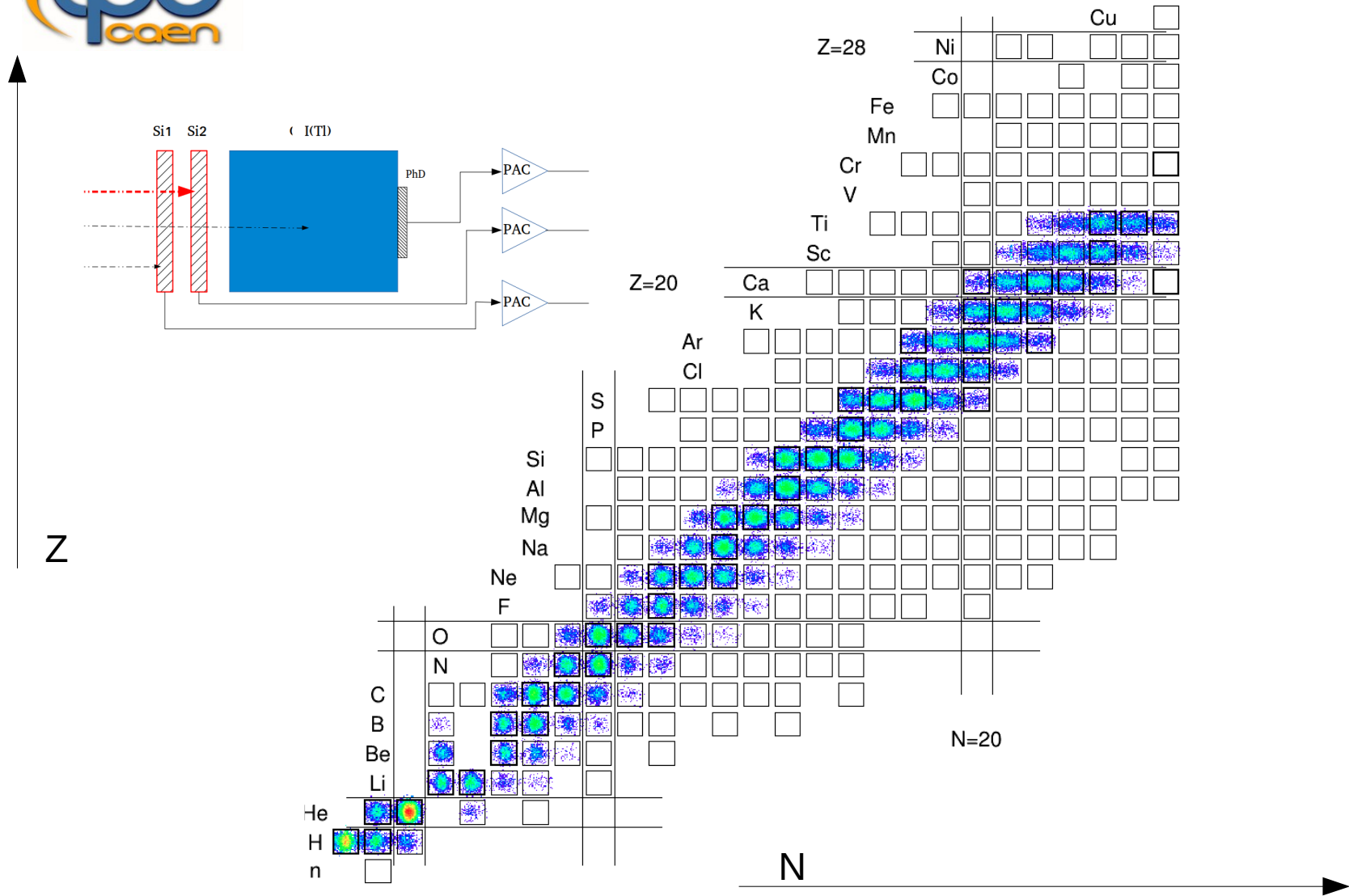


- ✓ numeric electronic = better signal treatment
- ✓ Si new generation optimized for mass detection (homogeneity)
- ✓ Good angular resolution (small size of telescopes)

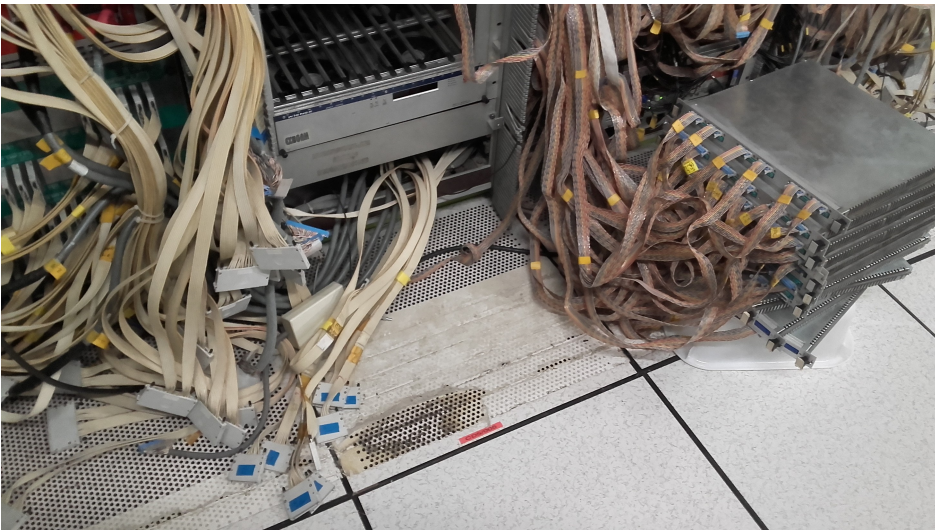
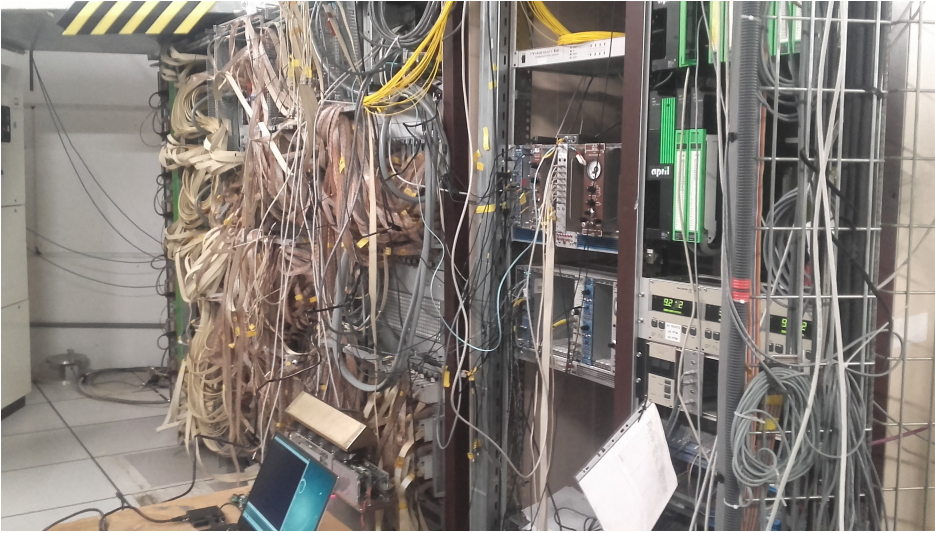
Z identification up to  $Z = 92$

Isotopic identification up to  $Z = 25$  !

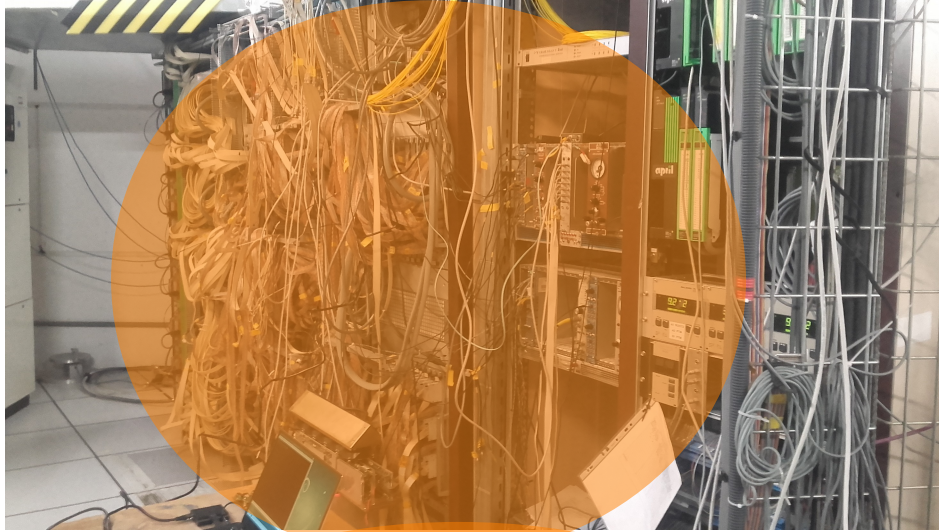
# The multidetectors : FAZIA



Identification map







**INDRA**

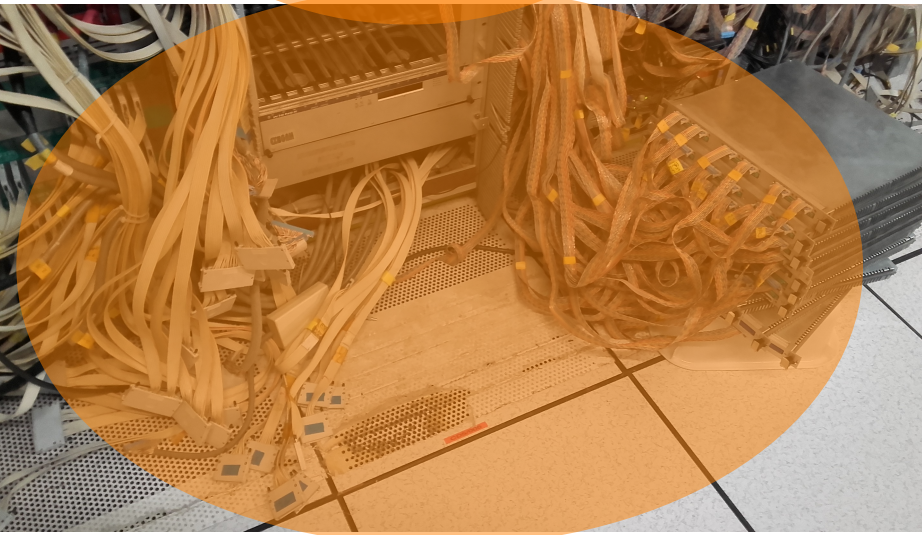
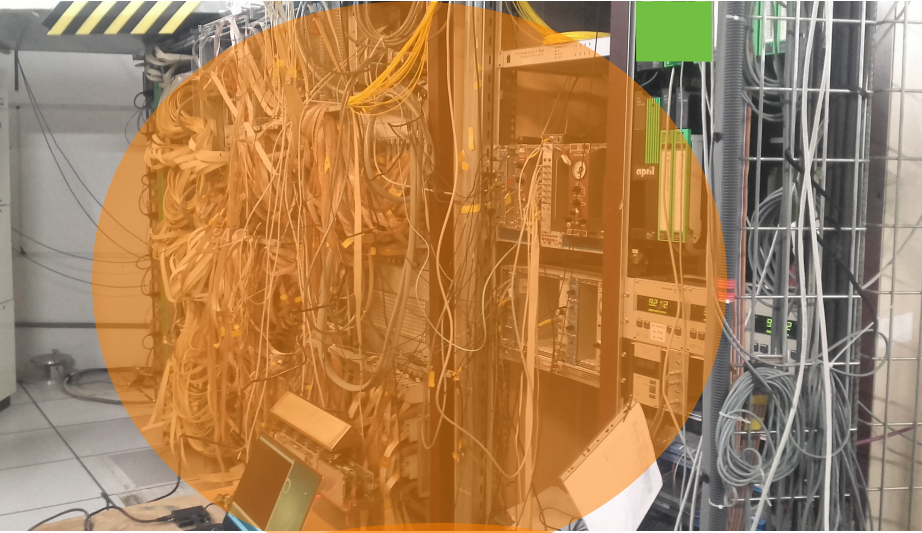




# The multidetectors : comparison

**FAZIA**

**INDRA**



# The multidetectors : comparison



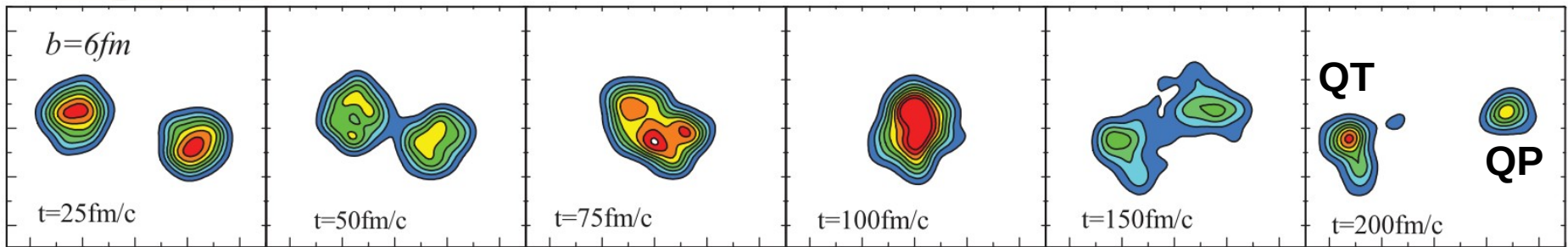
**FAZIA**



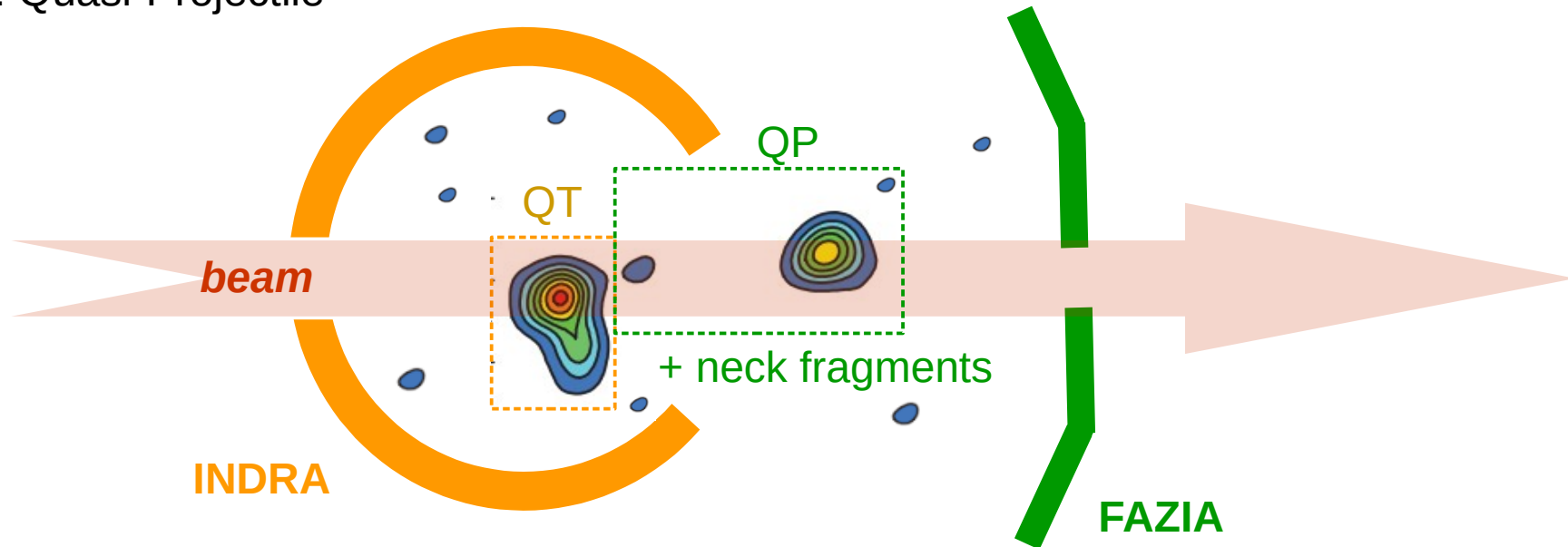
**INDRA**







QT : Quasi-Target  
 QP : Quasi-Projectile





## Mounting E789 INDRA-FAZIA (december 2018 – march 2019)

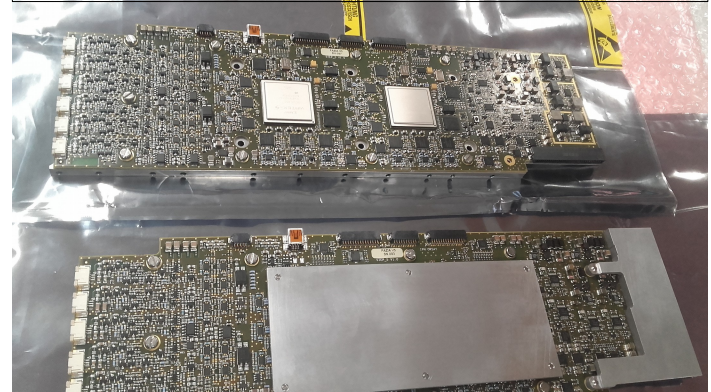
# Mounting E789 : Mounting FAZIA *Block's electronic*

Cooling plate



*Working in vacuum (6h de pumping)*

Telescopes electronic cards



PS card, Half Bridge, Block card



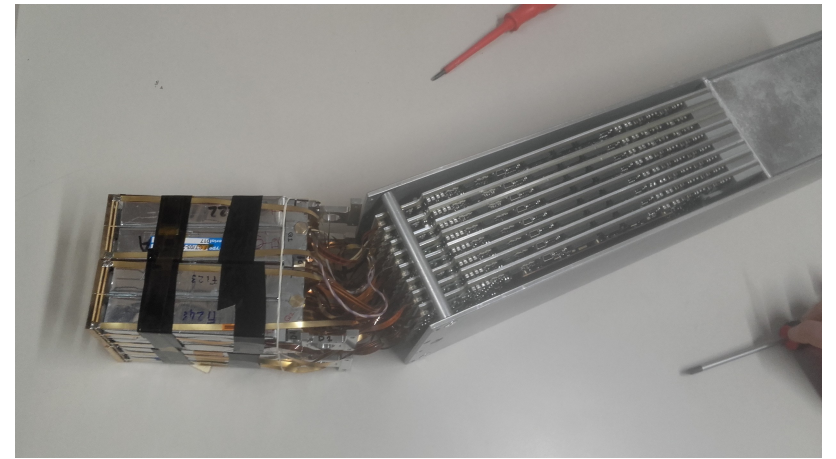
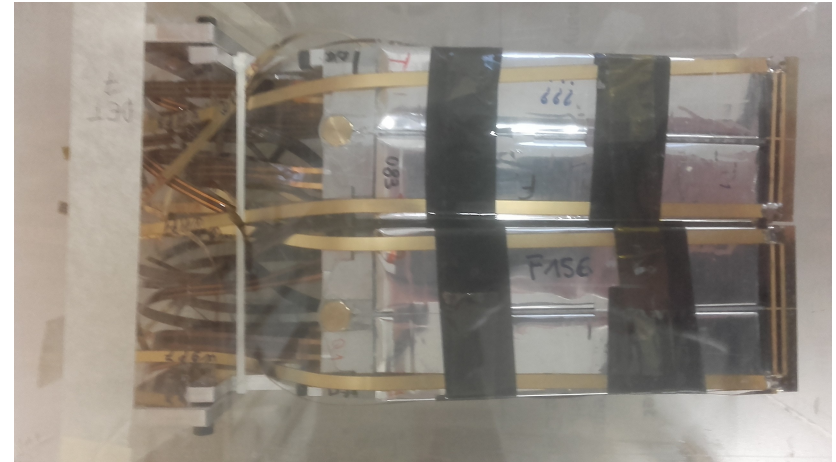
All mounted in the shelter





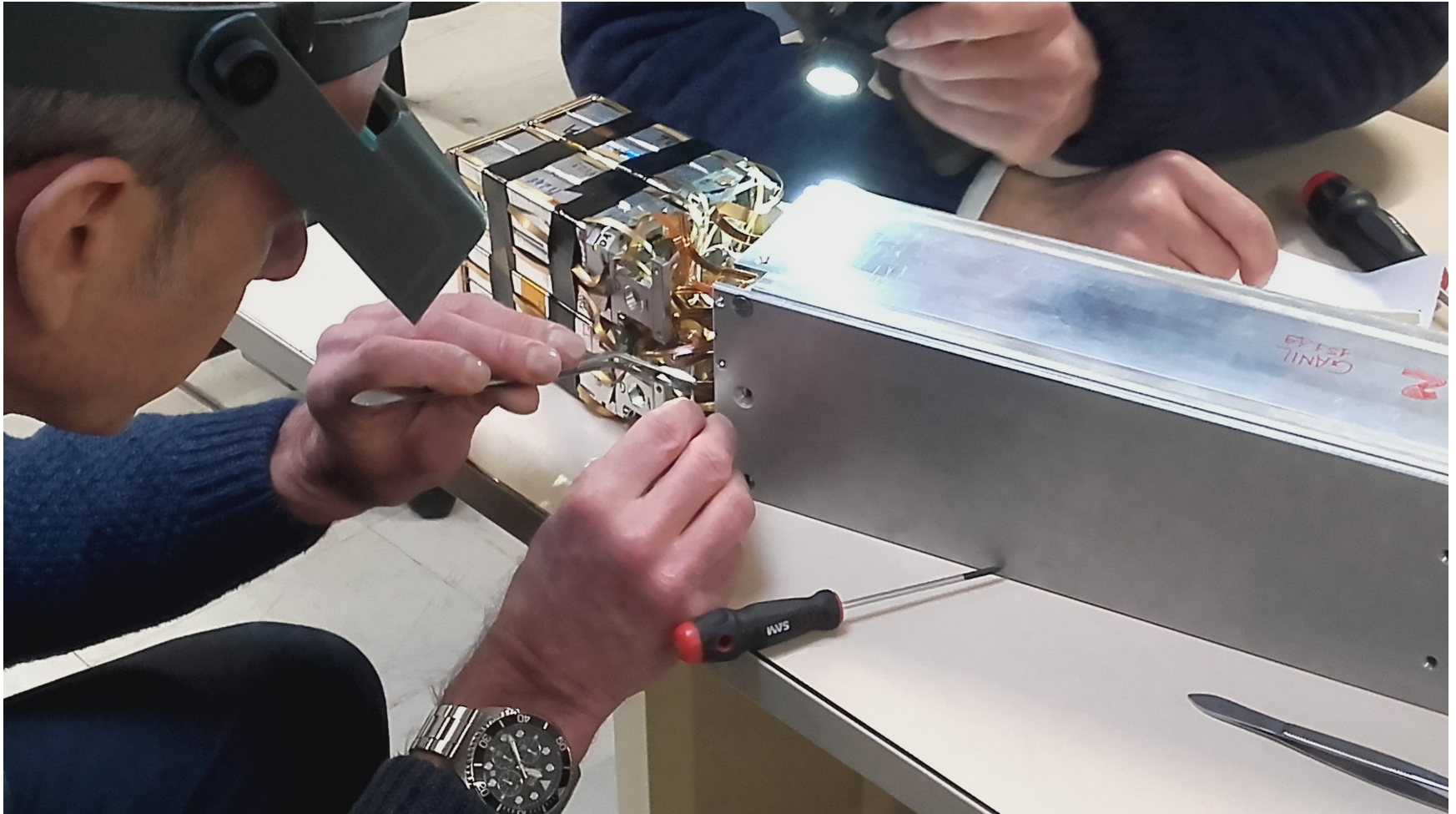
# Mounting E789 : Mounting FAZIA

*Detection's head*



1 cable Kapton by detector = 48 cables by head of detection

# Mounting E789 : Mounting FAZIA *Detector's head*



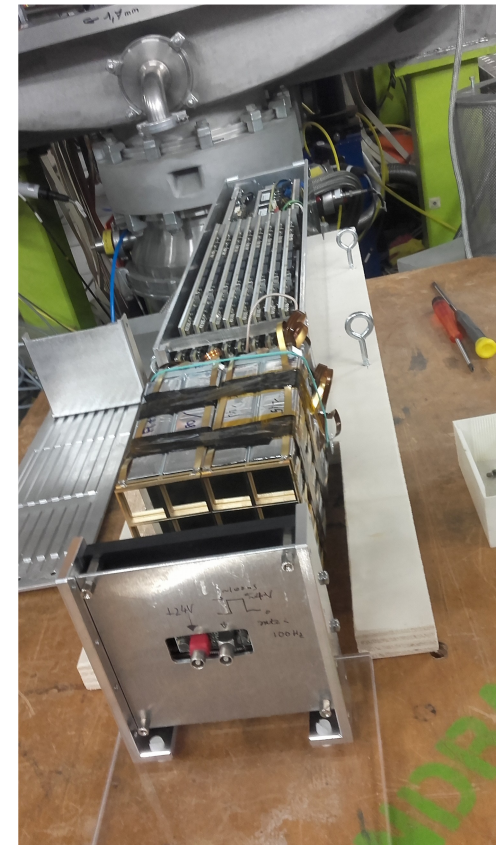
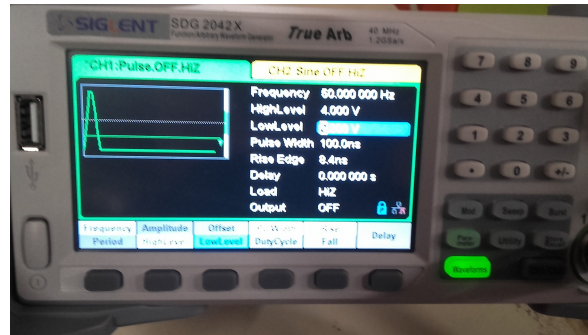
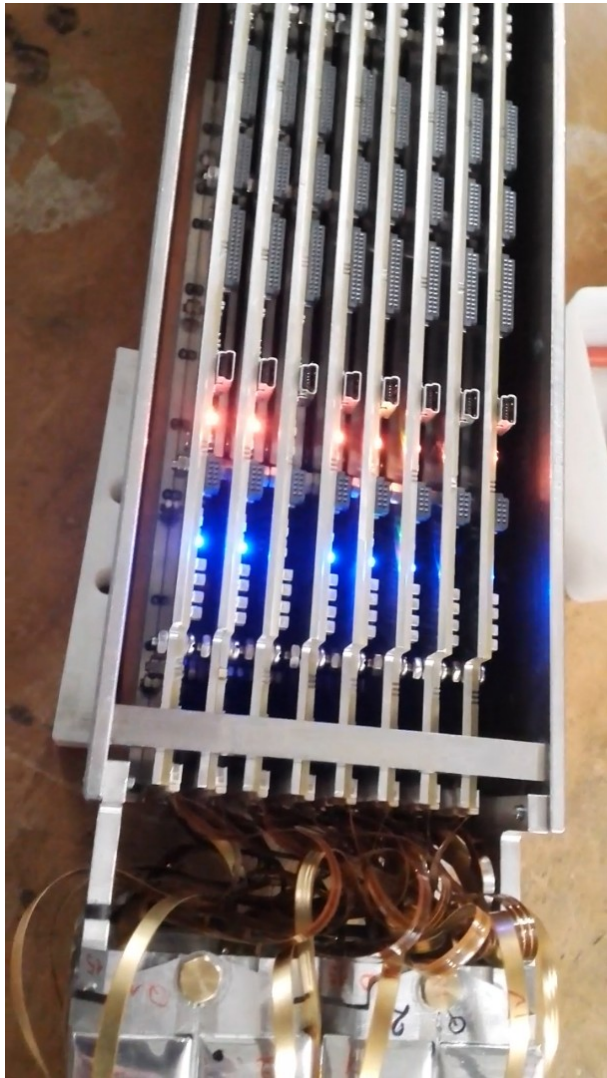
*Giovanni Casini connecting the detector's head to the electronic cards*



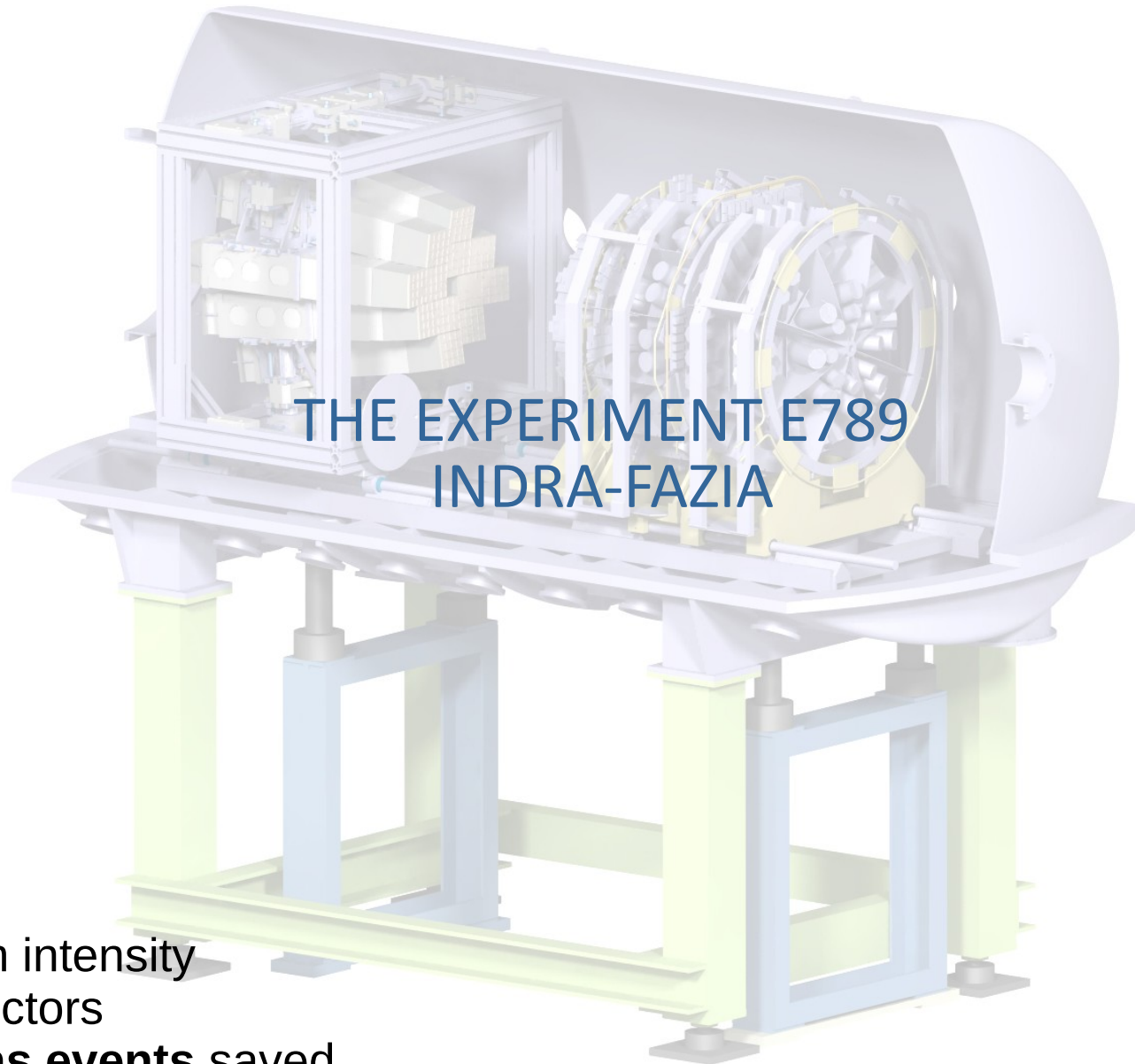
# Mounting E789 : Mounting FAZIA

Electronic tests  
Pulser tests  
Infrared LED tests  
Cosmics/alphas tests

*Each block is tested alone  
before mounting in the  
vacuum chamber*



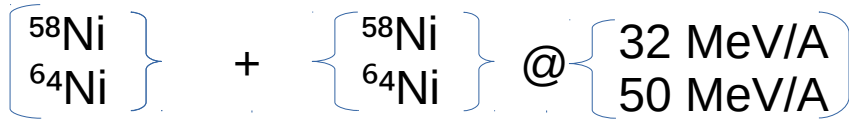




- ✓ Good beam intensity
- ✓ Stable detectors
- ✓ **240 millions events saved**

# The experiment : the beam

2 projectiles X 2 targets X 2 energies = 8 systems



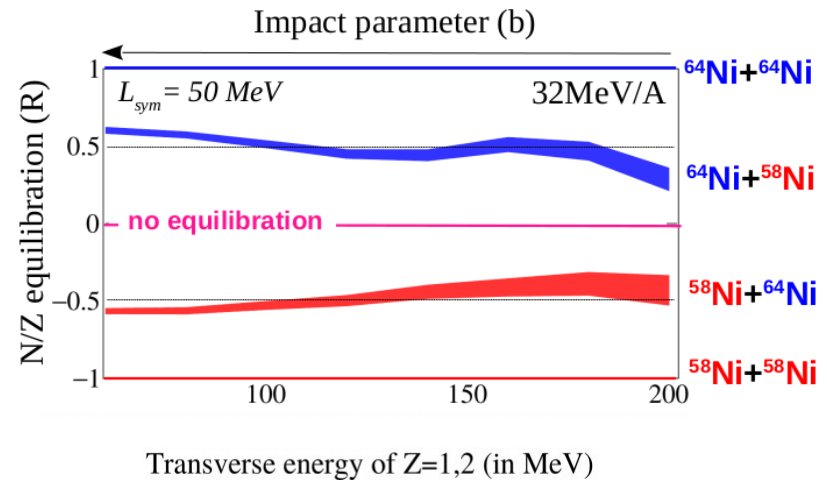
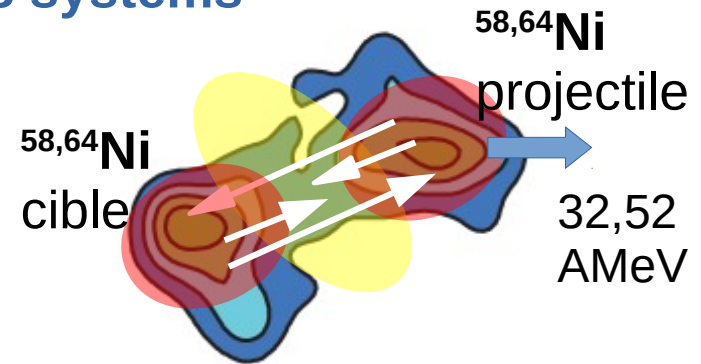
isospin gradient

density gradient  
interaction time

diffusion

migration

Study symmetric systems avoid experimental influence

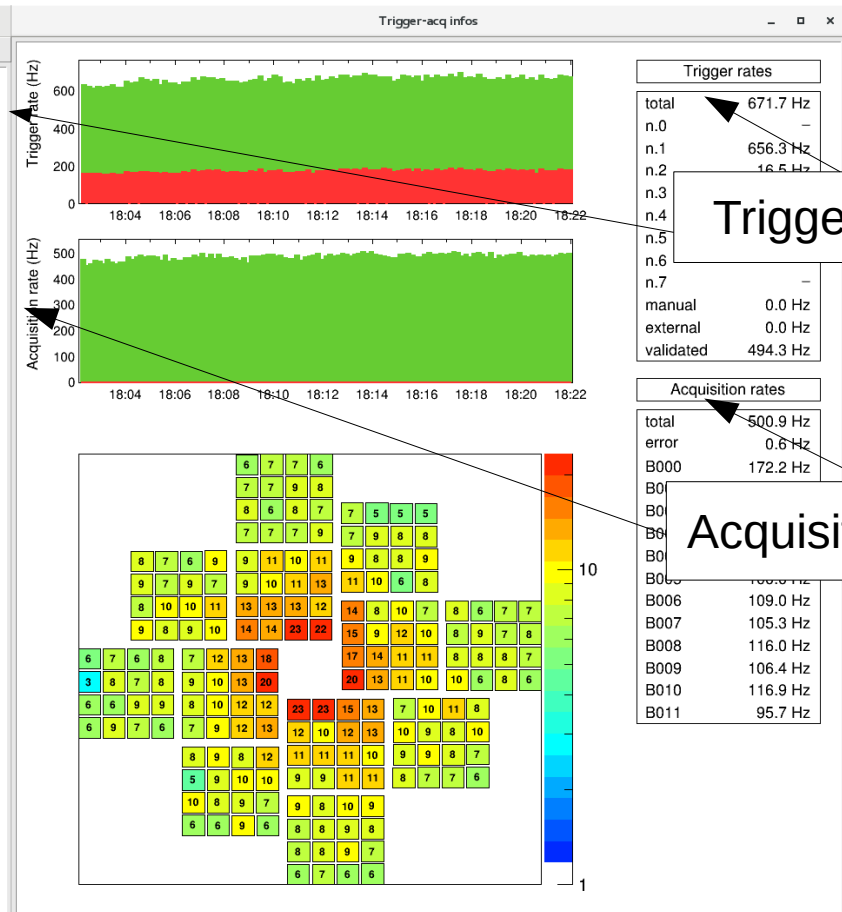
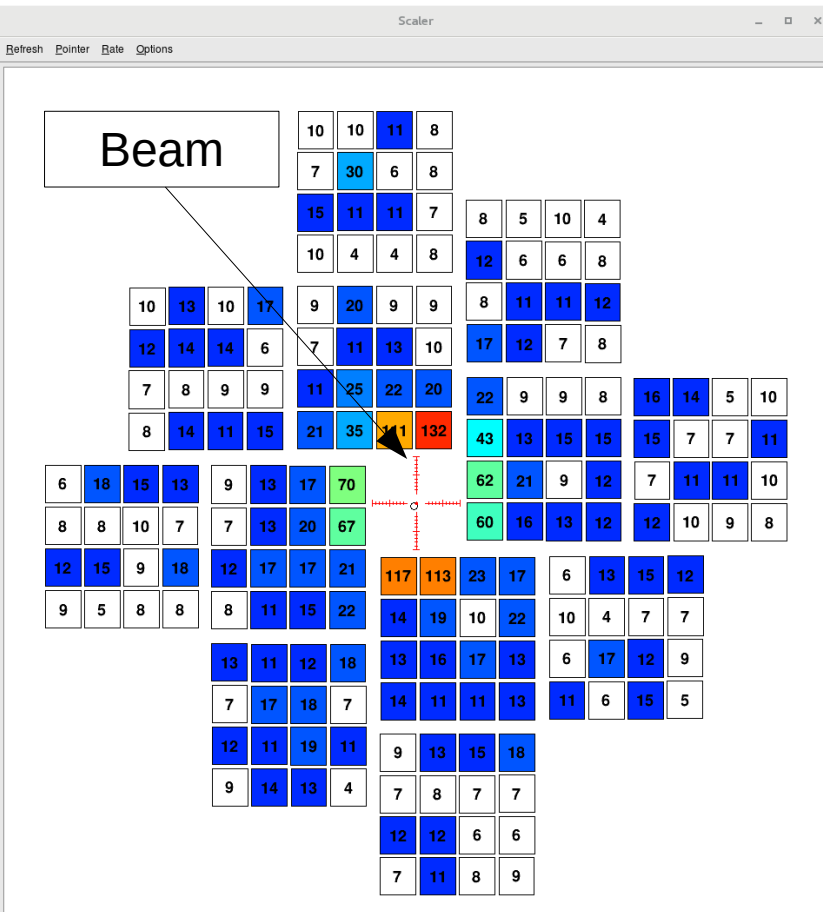




# The experiment: monitoring

Temperature of FAZIA cards  
 Temperature of INDRA electronic  
 Vacuum quality      Electronic stability

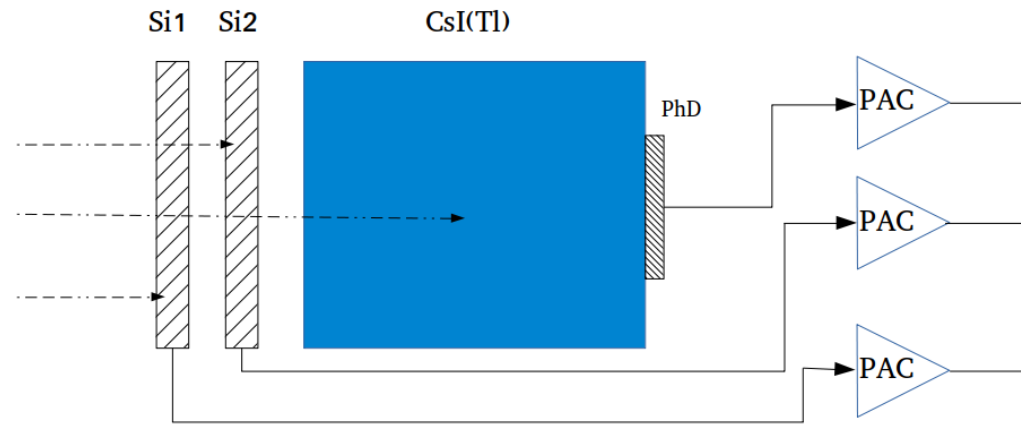
Beam intensity  
 Beam focus  
 Telescopes counting rate





# DATA REDUCTION

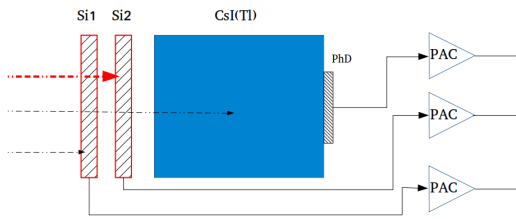
# Data reduction : identification methods



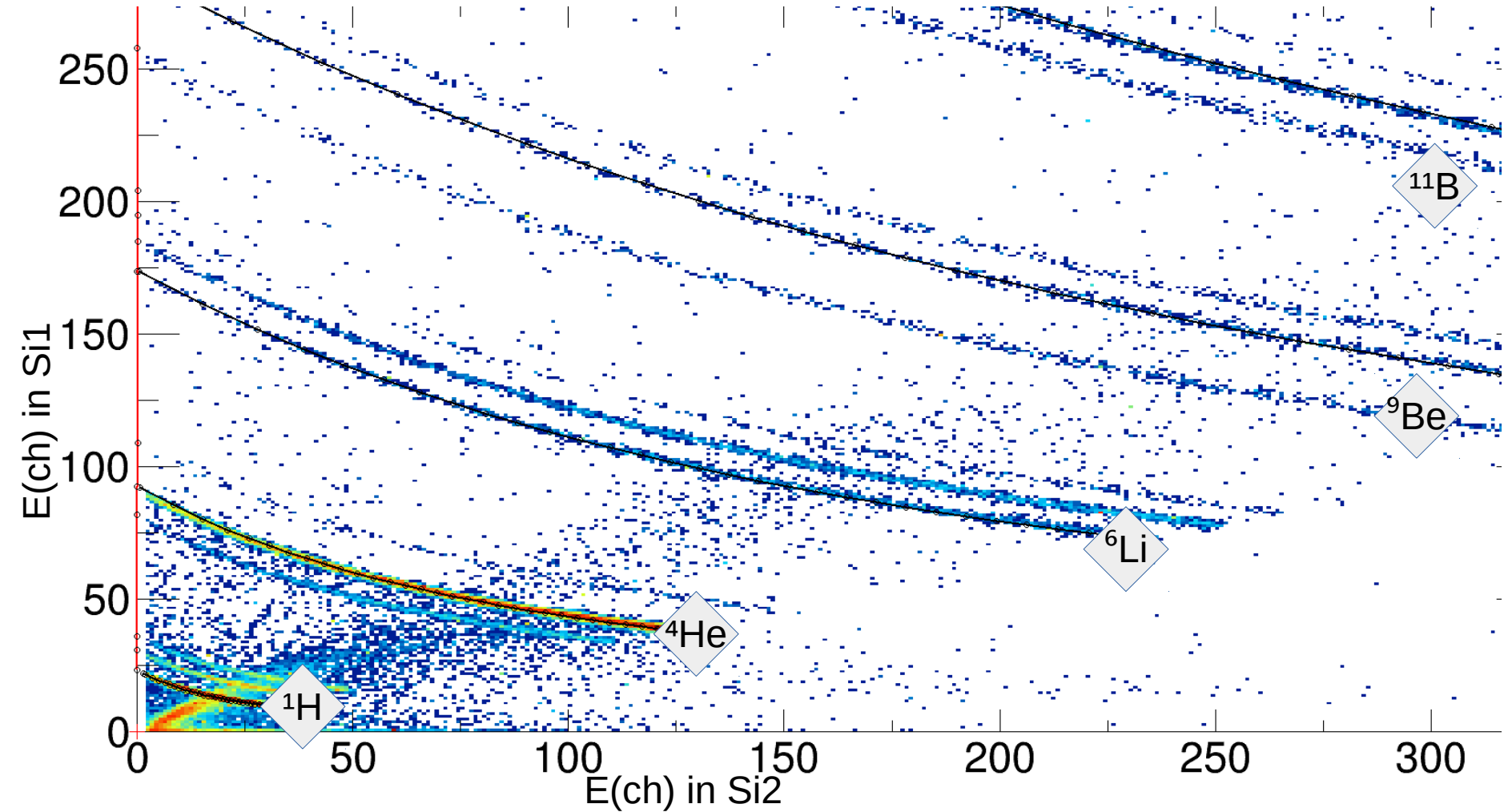
- The **Pulse Shape Analysis** (PSA) method in the first, second silicon or cesium iodure scintillator (Si1, Si2, CsI(Tl))
- The  $\Delta E - E$  method between the first and the second silicon (Si1 - Si2) but also between the second silicon and the CsI (TI) (Si2 - CsI)

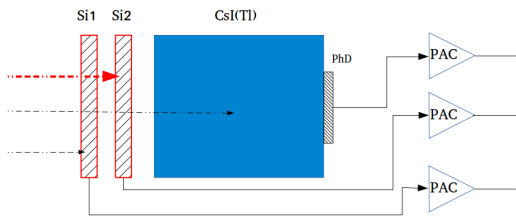


$$\text{Bethe - Bloch} \quad \Delta E \propto \frac{AZ^2}{E}$$

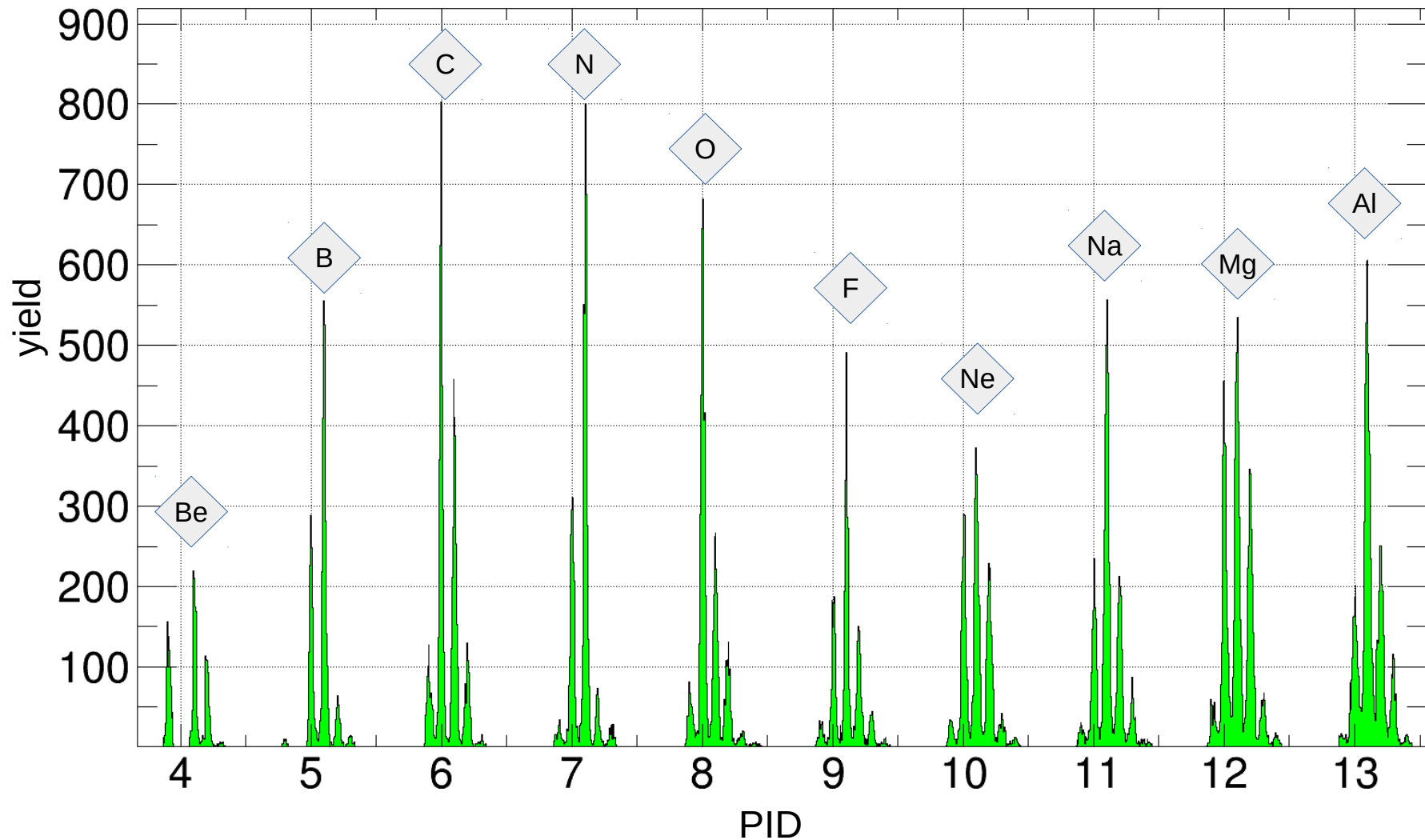


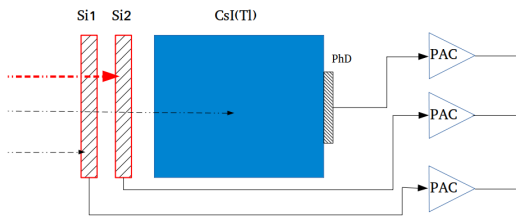
# Data reduction : Identification Si1-Si2



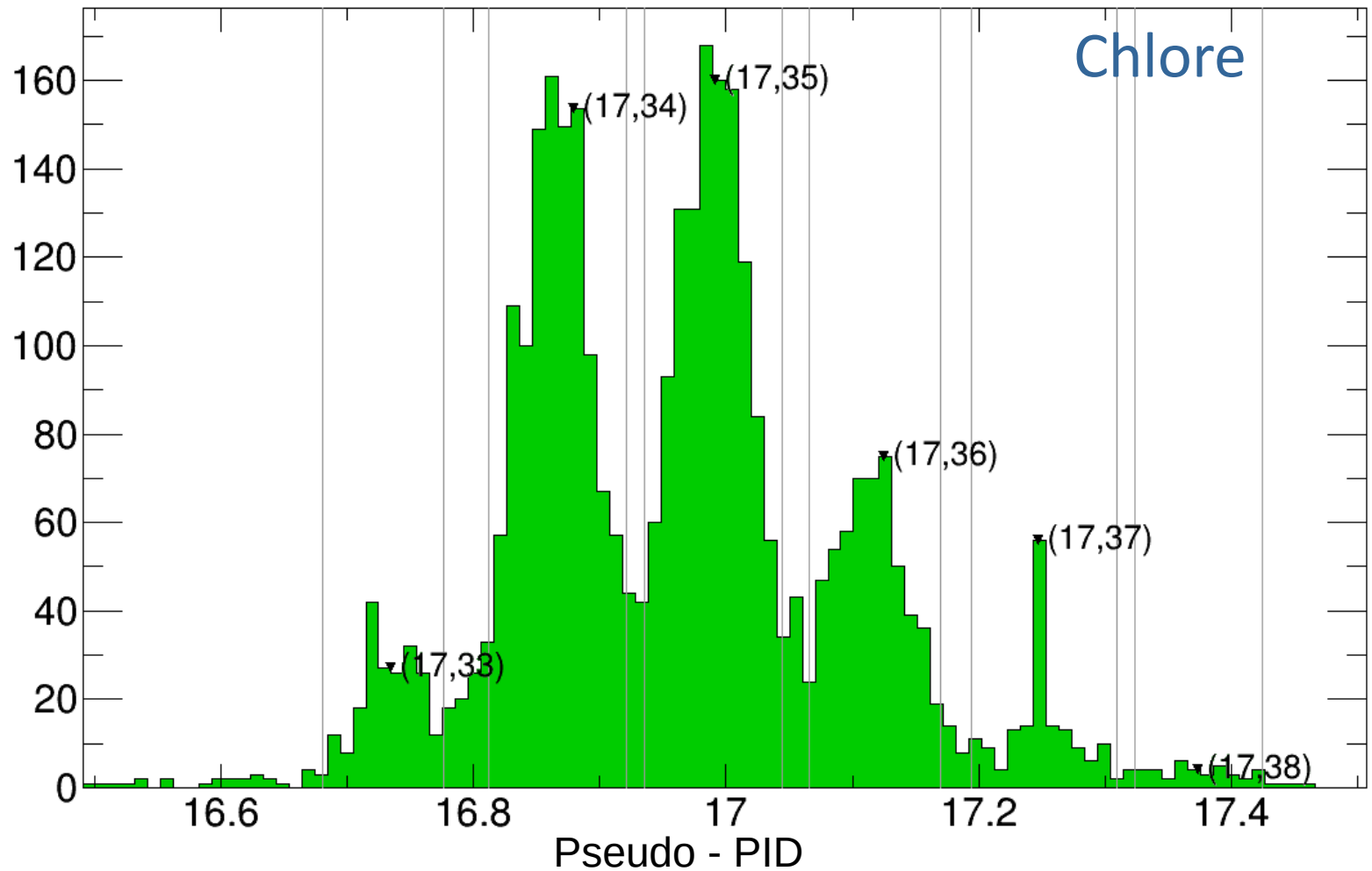


# Data reduction : Identification Si1-Si2

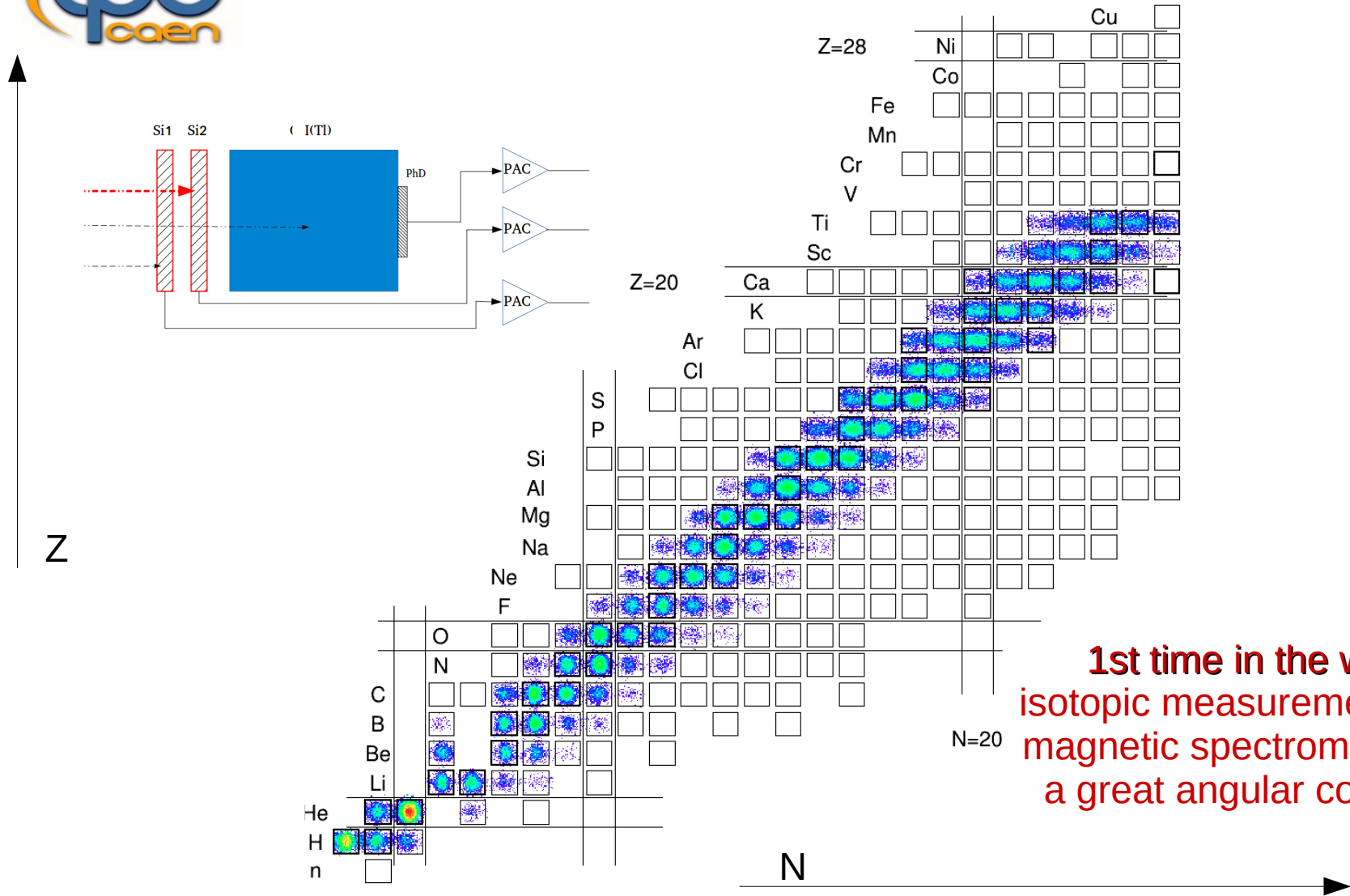




# Data reduction : Identification Si1-Si2

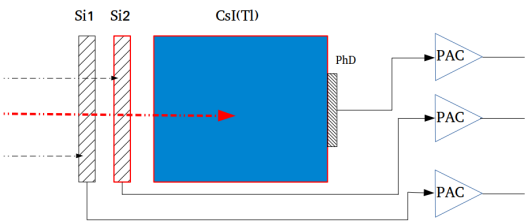




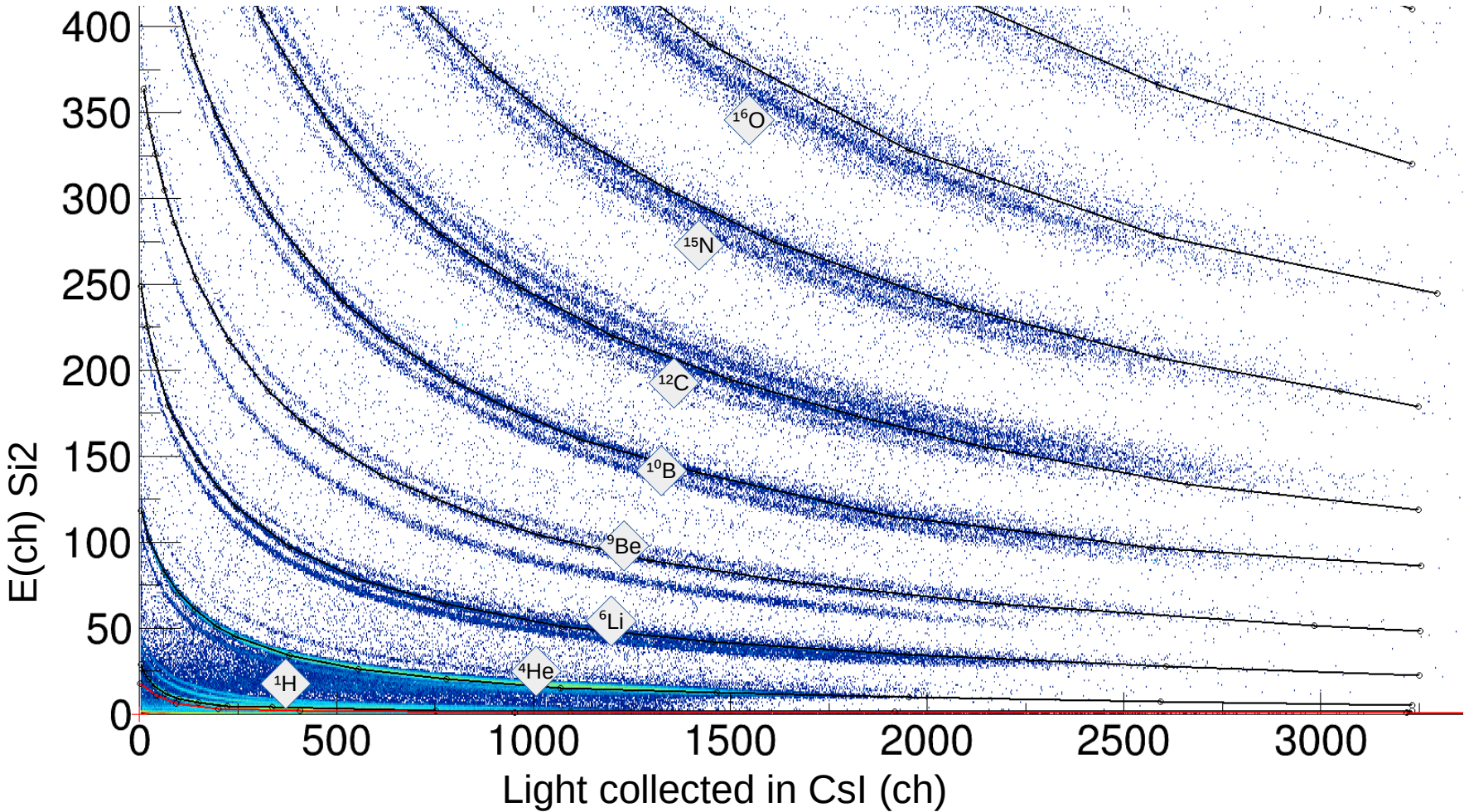


**1st time in the world  
isotopic measurement like a  
magnetic spectrometer with  
a great angular coverage**

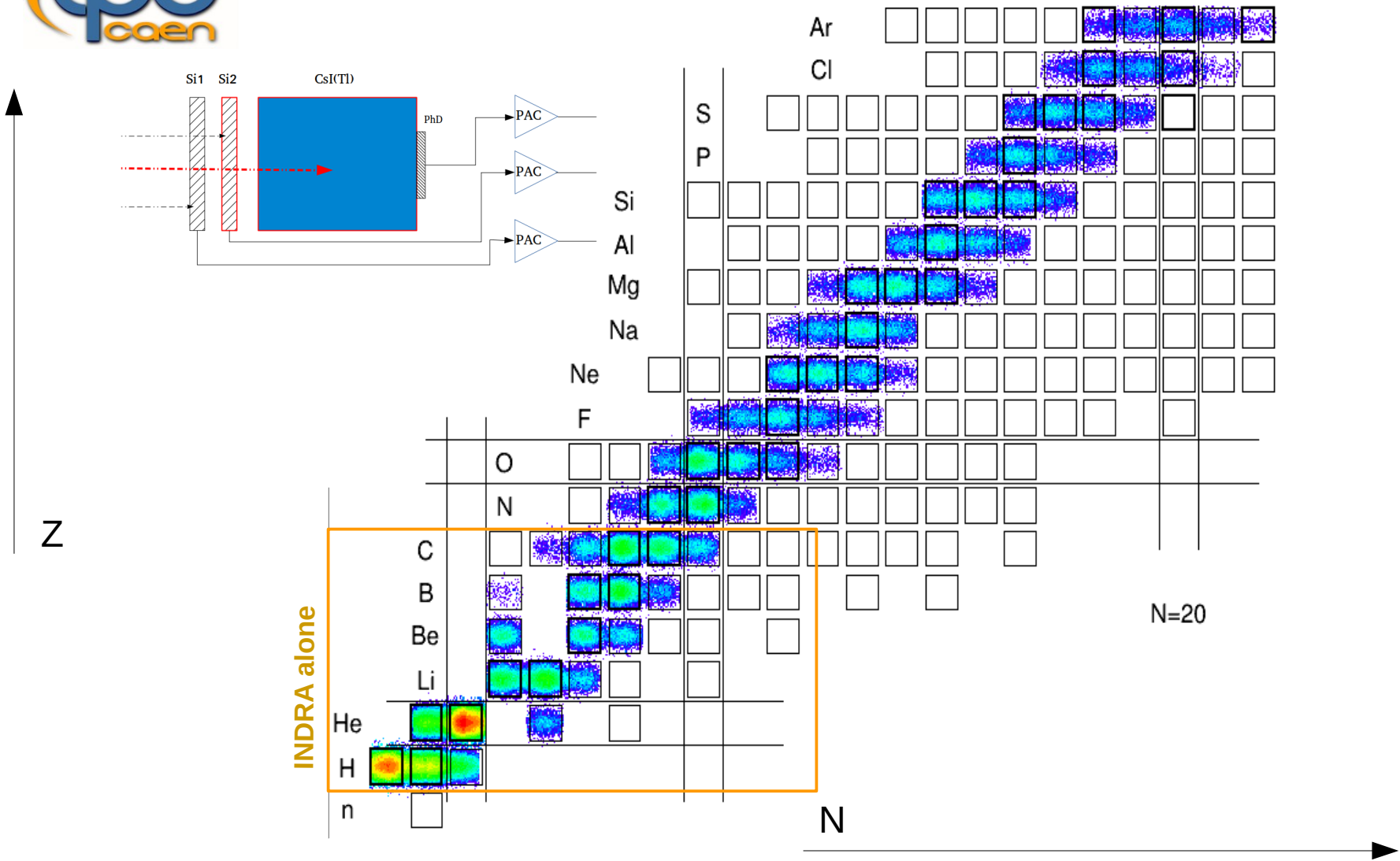
Identification map



# Data reduction: Identification Si2-CsI

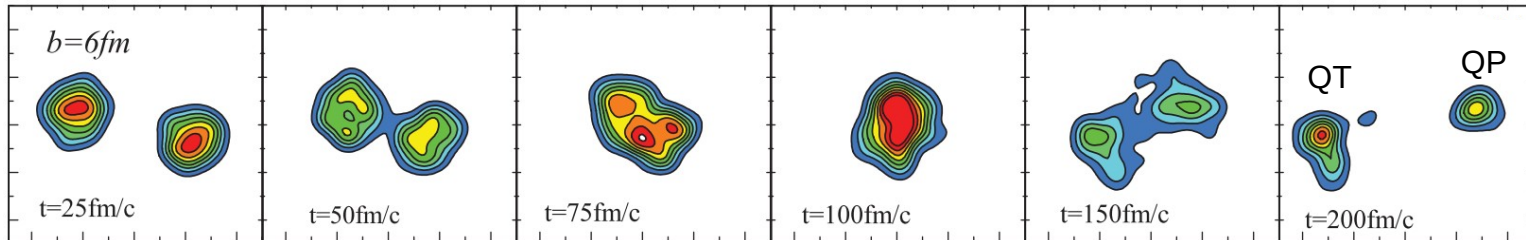


# Data reduction : Identification Si2-CsI



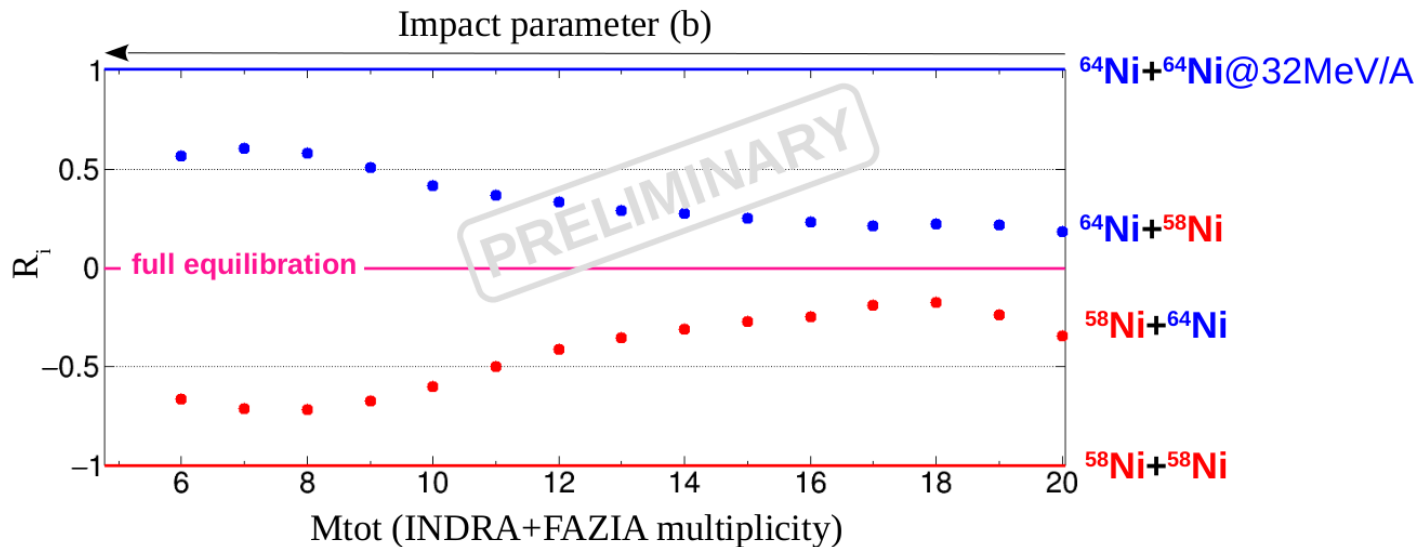
Identification map, Z vs N

# Preliminary results : isospin equilibration

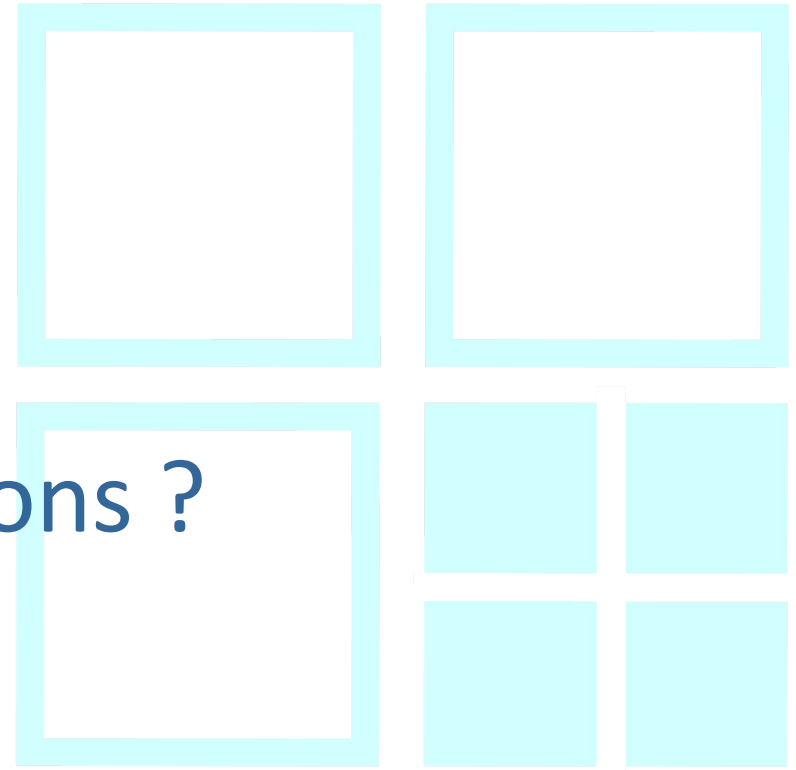


$$R_i = \left\langle \frac{N}{Z} \right\rangle_{QP}$$

$R_i = +1$  ( $-1$ ) : no N/Z equilibration  
 $R_i = 0$  : full N/Z equilibration



Smaller b  
 ↓  
 Nuclei more mixed  
 ↓  
 N/Z équilibré



Questions ?

INDRA FAZIA



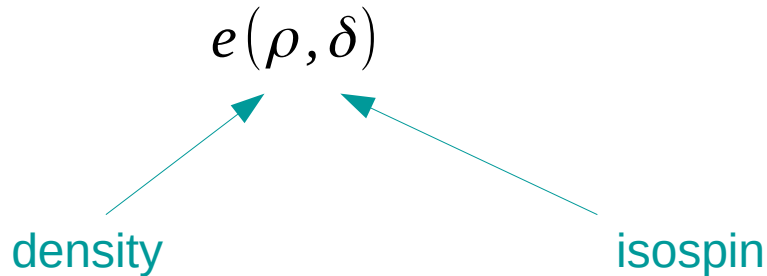


# BACKUP

B-W

$$E_L(N, Z) = \underbrace{-a_V \cdot A}_{\text{Volume}} + \underbrace{a_S \cdot A^{2/3}}_{\text{Surface}} + \underbrace{(a_V^{sym} A + a_S^{sym} A^{2/3}) \delta^2}_{\text{Symétrie}} + \underbrace{a_C \cdot \frac{Z(Z-1)}{A^{1/3}}}_{\text{Coulomb}}$$

- Historical formula of Bethe-Weizsacker, good starting point
- System at **T = 0** and **without  $\rho$  dependence**
- System more realistic = function of protons and neutrons density :





# Les collisions et l'isospin

$$\vec{j}_n = D_n^\rho \vec{\nabla} \rho - D_n^\delta \vec{\nabla} \delta \quad : \text{courant de neutrons}$$

$$\vec{j}_p = D_p^\rho \vec{\nabla} \rho - D_p^\delta \vec{\nabla} \delta \quad : \text{courant de protons}$$

$D^\rho$  : coefficient de transport en densité  
 $D^\delta$  : coefficient de transport en isospin

$$|\vec{j}_n - \vec{j}_p| = \underbrace{(D_n^\rho - D_p^\rho) \vec{\nabla} \rho}_{\text{migration}} - \underbrace{(D_n^\delta - D_p^\delta) \vec{\nabla} \delta}_{\text{diffusion}}$$

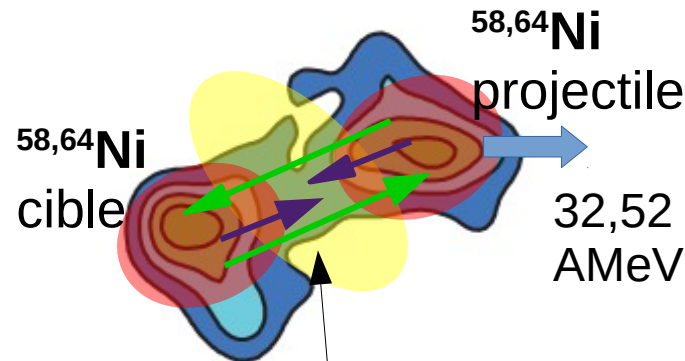
$$(D_n^\rho - D_p^\rho) \propto 4\delta \frac{\partial e_{iv}}{\partial \rho} \qquad (D_n^\delta - D_p^\delta) \propto 4\rho e_{iv} (\rho = \rho_0)$$

$\vec{\nabla} \rho$  : gradient d'isospin

$\vec{\nabla} \delta$  : gradient de densité

Varie avec le couple cible/projectile choisi

Varie avec l'énergie de collision



transport d'isospin :  
redistribution de neutrons & protons  
entre projectile/cible & col  
pendant le temps de réaction ( $10^{-22}$ - $10^{-20}$  sec.)

J. Frankland (GANIL, cnrs)

# Rapport de Rami

rapport isobarique ou  
isotopique des particules  
détectées pour le **système  
riche en neutrons**

rapport isobarique ou  
isotopique des particules  
détectées pour le **système  
riche en protons**

rapport isobarique ou  
isotopique des  
particules détectées

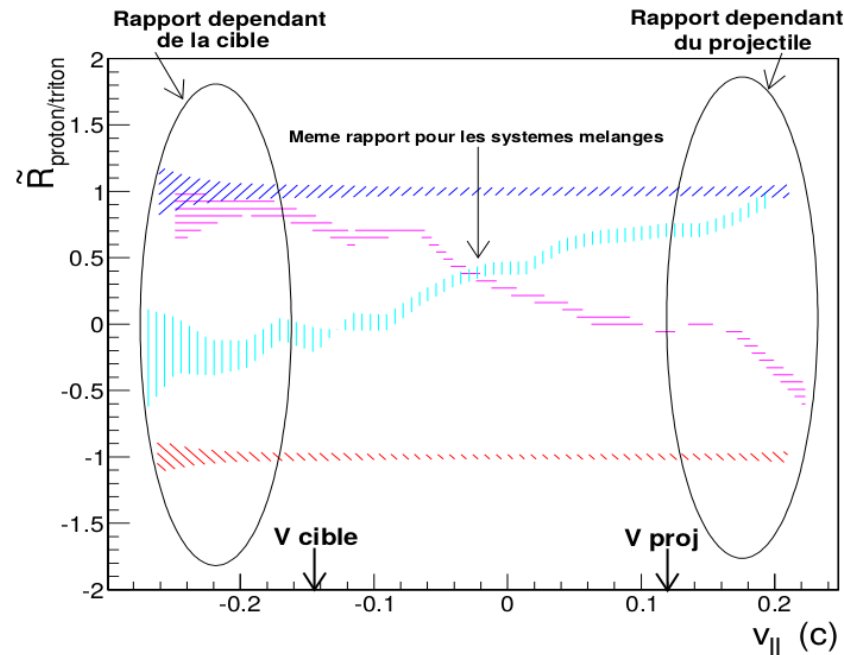
Rapport de Rami

$$\tilde{R}_{p/n} = \frac{2R_{p/n} - R^N_{p/n} - R^P_{p/n}}{R^N_{p/n} - R^P_{p/n}}$$

- Marqueur d'isospin : Une fois normalisé, on retrouve l'isospin du QP ou de la QT à l'avant du centre de masse ou non
- Elimine le biais expérimental

Pour la détermination du rapport de Rami, il faut se normaliser par rapport aux deux systèmes les plus extrêmes étudiés

- 1 pour système très riche en neutrons (correspondant à l'exposant N)
- -1 pour système très riche en protons (exposant P)



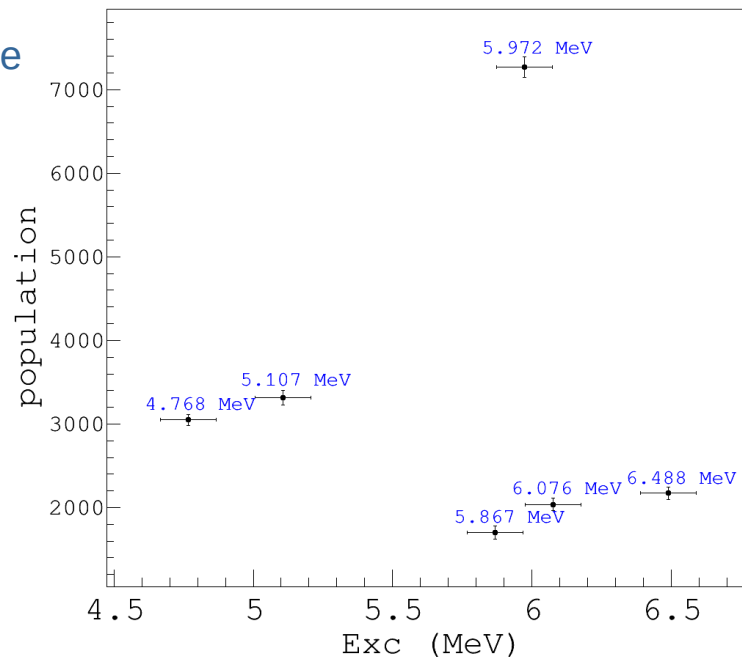
E. Legouée, Ph.D Thesis, Université de Caen-Basse Normandie (2013)

Hypothèses : Source de 10B à l'équilibre thermodynamique à une température donnée

$$\left[ \frac{dN}{dE} \right]_{\alpha} = A \times \exp\left(-\frac{E}{\langle T \rangle}\right) \sum_i (2J_i + 1) \times \frac{\Gamma_i/2\pi}{(E - E_i)^2 + \Gamma_i^2/4} \frac{\Gamma_{\alpha,i}}{\Gamma_i}$$

Constante de normalisation  $\uparrow$  Energie d'excitation  $\uparrow$  Moment angulaire total du niveau  $i$   $\uparrow$  Largeur intrinsèque du niveau  $i$   $\uparrow$  Proportion de décroissance  $\alpha$  du niveau  $\approx 1$   
 $\downarrow$  Nombre de décroissances  $\alpha$   $\downarrow$  Température moyenne du système  $\downarrow$  Terme de résonance

- Normalisation par  $(2J+1)$  de la population de chaque niveau
- Décroissance exponentielle non-observée, modèle thermique non-valide ici
- Noyaux produits dans différentes conditions de **température** car il n'y a pas eu de sélection selon le paramètre d'impact ou la température





# Les simulations, quelques ingrédients

## AMD-QMD

### Antisymetrized Molecular Dynamic – Quantum Molecular Dynamic

- Paquets d'onde gaussiens
- Potentiel local de Skyrme (dépend de la densité locale)
- Potentiel de Yukawa décrivant la surface
- Potentiel dépendant des vitesses
- Algorithme stochastique de collisions à deux corps
- Principe de Pauli plus ou moins bien traité

*Mouvement de toutes les particules du système sous l'action de leurs interactions mutuelles*

## HIPSE

### Heavy-Ion Phase Space Exploration

- Phase d'approche de la collision basée sur l'équation du mouvement à deux corps
- Réarrangement des nucléons en plusieurs clusters de particules légères selon le paramètre d'impact de la réaction. Suit les lois de coalescence dans l'espace des impulsions et des positions
- After-burner, la partition est propagée en prenant en compte les effets de reaggregation dus aux interactions nucléaires fortes et coulombiennes
- Les décroissances secondaires sont prises en compte à l'aide d'un code d'évaporation

*Très phénoménologique  
Interaction de Skyrme simplifiée  
Approximation soudaine*

