

# NMSSM signatures at the LHC

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# Why the NMSSM?

- Relatively heavy Higgs @ LHC  $\Rightarrow$  fine-tuning in the MSSM
- $\mu$ -problem of the MSSM: SUSY conserving mass term for  $H_u, H_d$ 
  - $\mu = 0$   $\Rightarrow$  experimentally excluded ( $\mu \geq 100$  GeV)
  - $\mu = M_{\text{Planck}}$   $\Rightarrow$  hierarchy problem ( $\mu \sim M_{\text{susy}}$ )
- Solution: add a gauge singlet  $S$  coupled to  $H_u, H_d$

$$W_{\text{MSSM}} = \mu H_u H_d + \dots \quad \rightarrow \quad W_{\text{NMSSM}} = \lambda S H_u H_d + \frac{\kappa}{3} S^3 + \dots$$

After minimisation of the potential  $\mu_{\text{eff}} \equiv \lambda \langle S \rangle \sim M_{\text{susy}}$

- The NMSSM is the simplest SUSY extension of the SM where the electroweak scale originates from the SUSY breaking **only**
  - $\lambda \sim \kappa \rightarrow 0, \mu_{\text{eff}} \neq 0$   $\Rightarrow$  effective MSSM + decoupled singlet sector
- $\Rightarrow$  the parameter space of the NMSSM includes the MSSM **and more**

# What's the NMSSM?

## ■ Particle content:

- $\tilde{S}$ : one more neutralino  $\Rightarrow \tilde{\chi}_{i=1..5}^0$
- $S_R$ : one more neutral CP-even Higgs  $\Rightarrow h_{i=1..3}$
- $S_I$ : one more neutral CP-odd Higgs  $\Rightarrow a_{i=1..2}$

## ■ Parameters: $V_{\text{Higgs}} = V_F + V_D + V_{\text{soft}}$

$$V_{\text{soft}} = \left( \lambda A_\lambda H_u H_d S + \frac{\kappa}{3} A_\kappa S^3 + \text{h.c.} \right) + m_{H_u}^2 |H_u|^2 + m_{H_d}^2 |H_d|^2 + m_S^2 |S|^2$$

+ 3 minimization equations:

$$\mu_{\text{eff}} = \lambda \langle S \rangle, \quad \tan\beta = \frac{\langle H_u^0 \rangle}{\langle H_d^0 \rangle}, \quad M_Z^2 = g^2 \left( \langle H_u^0 \rangle^2 + \langle H_d^0 \rangle^2 \right)$$

$\Rightarrow$  6 free parameters in the Higgs sector:  $\lambda, \kappa, A_\lambda, A_\kappa, \mu_{\text{eff}}, \tan\beta$

Recall: in the MSSM, 2 free parameters in the Higgs sector ( $m_A, \tan\beta$ )

NB:  $A_\lambda, A_\kappa \leftrightarrow m_A, m_P$  (diagonal elements of the CP-odd mass matrix)

# Parameters at the weak scale

- If  $A_\lambda = A_\kappa = 0$ , the Higgs sector is invariant under a  $U(1)$  R-symmetry explicitly broken by gaugino masses  $\Rightarrow a_1$  is a **light pseudo-Goldstone**  
Light pseudoscalar also with PQ symmetry when  $\kappa = 0$  (nMSSM)  
In this case the singlino is also very light ( $m_{a_1} \simeq 2 m_{\tilde{S}}$  for correct DM)

- In general:  $\Delta W = \mu H_u H_d + \mu' S^2 + \xi_F S$ ,  $\Delta V_{\text{soft}} = (m_3^2 H_u H_d + m_S'^2 S^2 + \xi_S S + \text{h.c.})$

One can always put  $\mu = 0$  by the field redefinition  $S \rightarrow S - \mu/\lambda$

However assuming  $\xi_F = \xi_S = \mu' = m_S' = m_3 = 0 \Leftrightarrow Z_3$  symmetry

$\Rightarrow$  Domain wall problem: anisotropies in CMB + wrong BBN

$\Rightarrow$  Assume tiny  **$Z_3$ -breaking** interactions (no impact on spectrum)

$\Rightarrow$  Tadpole problem:  $\xi_F = \Lambda M_{\text{susy}}$ ,  $\xi_S = \Lambda M_{\text{susy}}^2$  (hierarchy pb)

$\Rightarrow$  Assume new **discrete R-symmetries** at the high scale such that tadpole terms are generated radiatively with small coefficients

NB: In nMSSM (where  $\kappa = 0$ ), one needs  $\xi_F, \xi_S \neq 0$



# Parameters at the high scale

- mSUGRA:  $M_{1/2}, m_0, A_0$  (at  $M_{\text{GUT}}$ ) +  $\lambda, \kappa, \tan\beta, \text{sgn}(\mu_{\text{eff}})$  (at  $M_{\text{weak}}$ ) ?  
⇒ Z<sub>3</sub>-NMSSM: 1 free parameter ( $\mu_{\text{eff}}$ ) for 3 min. equations ( $M_{\text{weak}}$ )  
↳ Solution: assume semi-universal singlet soft terms at  $M_{\text{GUT}}$ 
  - Parameters:  $\lambda, \tan\beta, \text{sgn}(\mu_{\text{eff}}), M_{1/2}, m_0, A_0, [A_\lambda, A_\kappa]$
  - Minimization conditions ⇒  $\mu_{\text{eff}}, \kappa, m_S$  (at  $M_{\text{weak}}$ )
    - ✦ Guess  $M_{\text{GUT}}$  and  $\kappa, m_S$  at this scale
    - ✦ Run the RGEs down to  $M_{\text{weak}}$ , compute  $\mu_{\text{eff}}, \kappa, m_S$  ←
    - ✦ Run the RGEs up to  $M_{\text{GUT}}$ , distance from universality →
  - Similar algorithm for GMSB (checking for correct  $m_S$ )
- cNMSSM: take  $A_\lambda = A_\kappa = A_0$ , keep only  $m_S = m_0$  (very constrained)
- NUH-NMSSM: Non-Universal Higgs sector
  - Parameters:  $M_{1/2}, m_0, A_0, A_\lambda, A_\kappa, \lambda, \kappa, \tan\beta, \mu_{\text{eff}}$
  - Minimization conditions ⇒  $m_{H_u}, m_{H_d}, m_S$  (at  $M_{\text{weak}}$ )
  - ↳ Used to study the nMSSM ( $\kappa = A_\kappa = 0$  but  $\xi_F, \xi_S \neq 0$ )

# The NMSSMTools package

## ■ Package that contains 4 programs:

- **NMHDECAY** for general NMSSM with soft terms at  $M_{\text{weak}}$
- **NMHDECAYCPV** for  $Z_3$ -NMSSM with CPV soft terms at  $M_{\text{weak}}$
- **NMSPEC** for mSUGRA (non-universal) models with inputs at  $M_{\text{GUT}}$
- **NMGMSB** for GMSB (extended) models with inputs at  $M_{\text{mes}}$

Each in 4 versions: 1 point, random scan, grid scan or MCMC scan

## ■ For a given set of input parameters, it computes:

- (2 loop) RGEs for all parameters  $M_{\text{weak}} \leftrightarrow M_{\text{GUT}}$
- Sparticle/Higgs masses and mixings (2 loop LL or *à la* Slavich)
- Higgs decay widths in 2/3 particles (QCD corrections as in **HDECAY**)
- DM relic density + (in)direct detection (interface with **micrOMEGAs**)
- B physics:  $b \rightarrow s\gamma$ ,  $B_s \rightarrow \mu^+\mu^-$ ,  $B \rightarrow X_s l^+l^-$ ,  $B \rightarrow \tau\nu$ ,  $\Delta M_{s(d)}$ ,  $Y \rightarrow a\gamma$ ,  $(g-2)_\mu \dots$

## ■ I/O files in SLHA2 conventions + script: `run $PATH/$P{inp}$S`

➡ `$PATH/$P{spectr}$S`, `$PATH/$P{omega}$S` (1 point)

➡ `$PATH/$P{out}$S`, `$PATH/$P{err}$S` (scan versions)

# Input files

```
inp.dat
# Input file for NMSSMTools
# Based on SUSY LES HOUCHEs ACCORD II

BLOCK MODSEL
  3      1      # NMSSM particle content
  1      0      # IMOD (0=general NMSSM, 1=mSUGRA, 2=GMSB)
  10     0      # ISCAN (0=no scan, 1=grid scan, 2=random scan)
  9      0      # Call micrOmegas (default 0=no, 1=relic density only,
#              2=dir. det. rate, 3=indir. det. rate, 4=both det. rates)
  8      0      # Precision for Higgs masses (default 0: as before,
#              1: full 1 loop + full 2 loop from top/bot Yukawas
#              2: as 1 + pole masses - 1&2 by courtesy of P. Slavich)

BLOCK SMINPUTS
  1      127.9200 # ALPHA_EM^-1(MZ)
  2      1.166390-5 # GF
  3      .117200  # ALPHA_S(MZ)
  4      91.18700  # MZ
  5      4.21400   # MB(MB) (running mass)
  6      171.400   # MTOP (pole mass)
  7      1.77700   # MTAU

BLOCK MINPAR
# 0      1000.00   # MSUSY (If /= SQRT(2*MQ1+MU1+MD1)/2)
  3      2.00      # TANB

BLOCK EXTPAR
# 1      100.00    # M1 (If /= M2/2)
  2      200.00    # M2
# 3      600.00    # M3 (If /= 3*M2)
  11     -2500.00  # AU3
  12     -2500.00  # AD3
  13     -2500.00  # AE3
# 16     -2500.00  # AE2 = AE1 (If /= AE3)
  33     200.00    # ML3
# 32     200.00    # ML2 = ML1 (If /= ML3)
  36     200.00    # ME3
# 35     200.00    # ME2 = ME1 (If /= ME3)
  43     1000.00   # MQ3
# 42     1000.00   # MQ2 = MQ1 (If /= MQ3)
  46     1000.00   # MU3
# 45     1000.00   # MU2 = MU1 (If /= MU3)
  49     1000.00   # MD3
# 48     1000.00   # MD2 = MD1 (If /= MD3)
  61     .700      # LAMBDA
  62     .0500     # KAPPA
  63     1200.     # ALAMBDA (If MA is not an input)
  64     0.00      # AKAPPA
  65     530.00    # MUEFF
# 124    1321.00   # MA (If ALAMBDA is not an input)
```

```
randinpsp.dat
# Input file for NMSSMTools
# Based on SUSY LES HOUCHEs ACCORD II

BLOCK MODSEL
  3      1      # NMSSM particle content
  1      1      # IMOD (0=general NMSSM, 1=mSUGRA, 2=GMSB)
  10     2      # ISCAN (0=no scan, 1=grid scan, 2=random scan)
  9      0      # Call micrOmegas (default 0=no, 1=relic density only,
#              2=dir. det. rate, 3=indir. det. rate, 4=both det. rates)
  8      0      # Precision for Higgs masses (default 0: as before,
#              1: full 1 loop + full 2 loop from top/bot Yukawas
#              2: as 1 + pole masses - 1&2 by courtesy of P. Slavich)

BLOCK SMINPUTS
  1      127.9200 # ALPHA_EM^-1(MZ)
  2      1.166390-5 # GF
  3      .117200  # ALPHA_S(MZ)
  4      91.18700  # MZ
  5      4.21400   # MB(MB) (running mass)
  6      171.400   # MTOP (pole mass)
  7      1.77700   # MTAU

BLOCK MINPAR
# 0      1000.00   # MSUSY (If /= SQRT(2*MQ1+MU1+MD1)/2)
  17     0.00      # M0_min
  18     1000.00   # M0_max (If /= min)
  27     100.00    # M12_min
  28     1000.00   # M12_max (If /= min)
  37     1.00      # TANB_min at MSUSY
  38     50.00     # TANB_max at MSUSY (If /= min)
  4      1.00      # SIGMU
  57     -1000.00  # A0_min
  58     1000.00   # A0_max (If /= min)

BLOCK EXTPAR
# 107    0.00      # M1_min at MGUT (If /= M12)
# 108    0.00      # M1_max at MGUT (If /= min)
# 207    0.00      # M2_min at MGUT (If /= M12)
# 208    0.00      # M2_max at MGUT (If /= min)
# 307    0.00      # M3_min at MGUT (If /= M12)
# 308    0.00      # M3_max at MGUT (If /= min)
# 217    0.00      # MHD^2_min at MGUT (If /= M0)
# 218    0.00      # MHD^2_max at MGUT (If /= min)
# 227    0.00      # MHU^2_min at MGUT (If /= M0)
# 228    0.00      # MHU^2_max at MGUT (If /= min)
  617    1.0-2     # LAMBDA_min at MSUSY
  618    7.0-1     # LAMBDA_max at MSUSY (If /= min)
# 637    0.00      # ALAMBDA_min at MGUT (If /= A0)
# 638    0.00      # ALAMBDA_max at MGUT (If /= min)
  647    -1000.00  # AKAPPA_min at MGUT (If /= A0)
  648    1000.00   # AKAPPA_max at MGUT (If /= min)

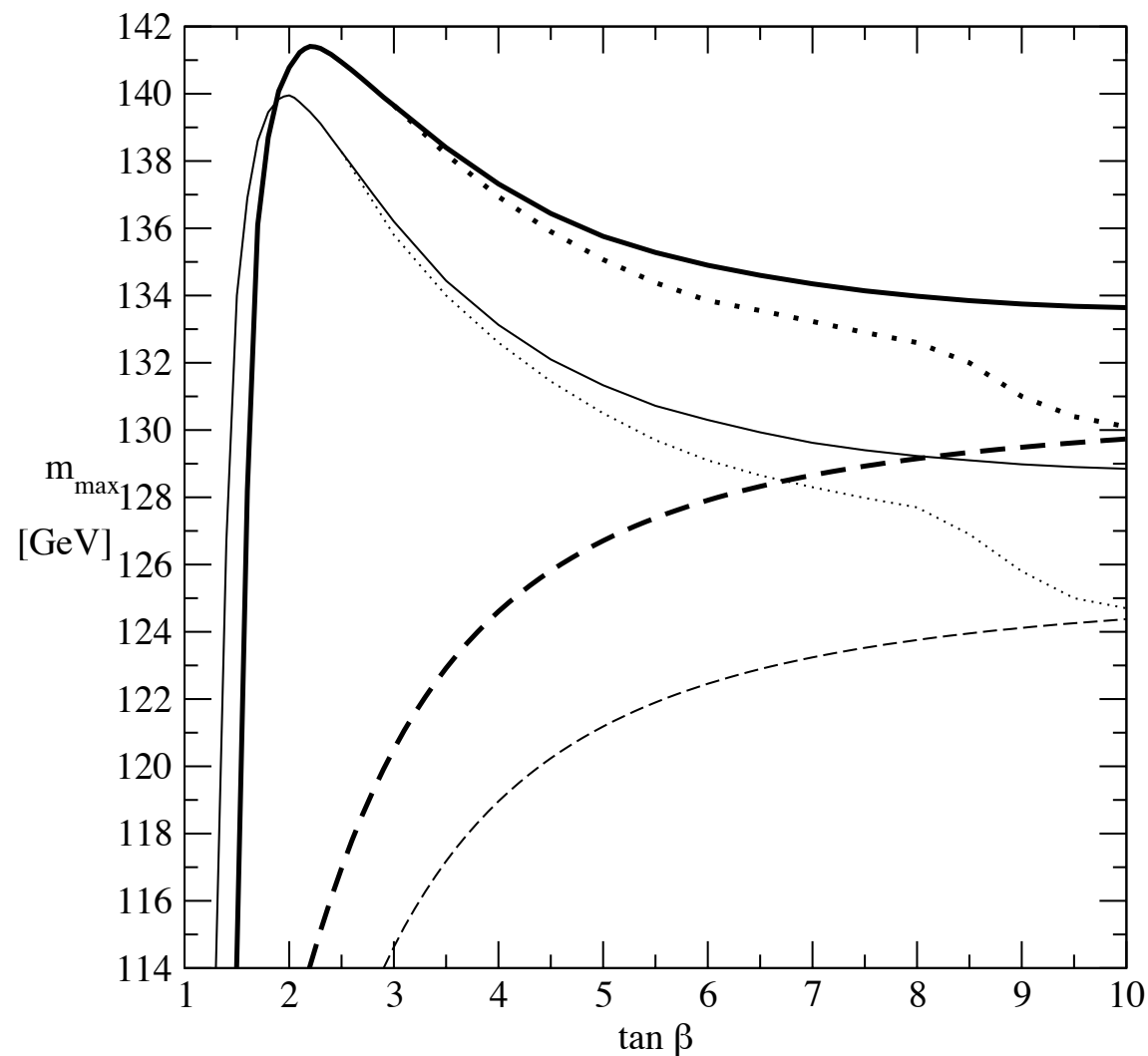
BLOCK STEPS
  0      100      # NTOT
  1      1        # ISEED
```

# Radiative corrections to the Higgs mass

- At tree level:  $m_h^2 = M_Z^2 \left( \cos^2 2\beta + \frac{\lambda^2}{g^2} \sin^2 2\beta \right) \sim (g^2 + \lambda^2)v^2$
- Radiative corrections considered:
  - Full one loop top/stop, bottom/sbottom:  $h_{t/b}^4 v^2 \left( \log \left( \frac{M_{\text{susy}}^2}{m_{\text{top}}^2} \right) + \text{finite} \right)$   
+ stop/sbottom D-term couplings:  $g^2 h_{t/b}^2 v^2 \left( \log \left( \frac{M_{\text{susy}}^2}{m_{\text{top}}^2} \right) + \text{finite} \right)$   
+ leading logs at two loops:  $(h_{t/b}^6 + h_{t/b}^4 \alpha_s) v^2 \log^2 \left( \frac{M_{\text{susy}}^2}{m_{\text{top}}^2} \right)$
  - EW leading logs at one loop:  $(g^2 + \lambda^2 + \kappa^2)^2 v^2 \log \left( \frac{M_{\text{susy}}^2}{m_{\text{top}}^2} \right)$
  - Higgs pole mass corrections:  $g^2 h_{t/b}^2 \left( \log \left( \frac{M_{\text{susy}}^2}{m_{\text{top}}^2} \right) + \text{finite} \right)$
- Terms *à la* Slavich:
  - full two loop stop/sbottom:  $h_{t/b}^4 \alpha_s v^2 \left( \log \left( \frac{M_{\text{susy}}^2}{m_{\text{top}}^2} \right) + \text{finite} \right)$
  - full one loop EW:  $(g^2 + \lambda^2 + \kappa^2)^2 v^2 \times \text{finite}$
- Still missing:  $h_{t/b}^6 v^2 \left( \log \left( \frac{M_{\text{susy}}^2}{m_{\text{top}}^2} \right) + \text{finite} \right)$

# Interlude: how to cook a 125 GeV Higgs

- The large  $\lambda$  small  $\tan\beta$  solution



- ➡ Requires heavy singlet sector disfavoured in  $\tilde{S}$  LSP scenario

- Large  $M_{\text{susy}}$  (as in the MSSM)  $\Rightarrow$  large fine-tuning, defined as

$$\Delta = \max \left( \left| \frac{\partial \ln(M_Z)}{\partial \ln(p_i)} \right| \right)$$

- The push-up mechanism

$$\mathcal{M}_H^2 = \begin{bmatrix} m_S^2 & \lambda(\dots) & \dots \\ \lambda(\dots) & m_h^2 & \dots \\ \dots & \dots & m_A^2 \end{bmatrix}$$

$$\underline{m_h = 110 \text{ GeV}}$$

$$\underline{m_S = 105 \text{ GeV}}$$

mixing

$$\underline{m_h = 125 \text{ GeV}}$$

$$\underline{m_S = 90 \text{ GeV}}$$

- ➡ Requires  $\lambda \sim .1$

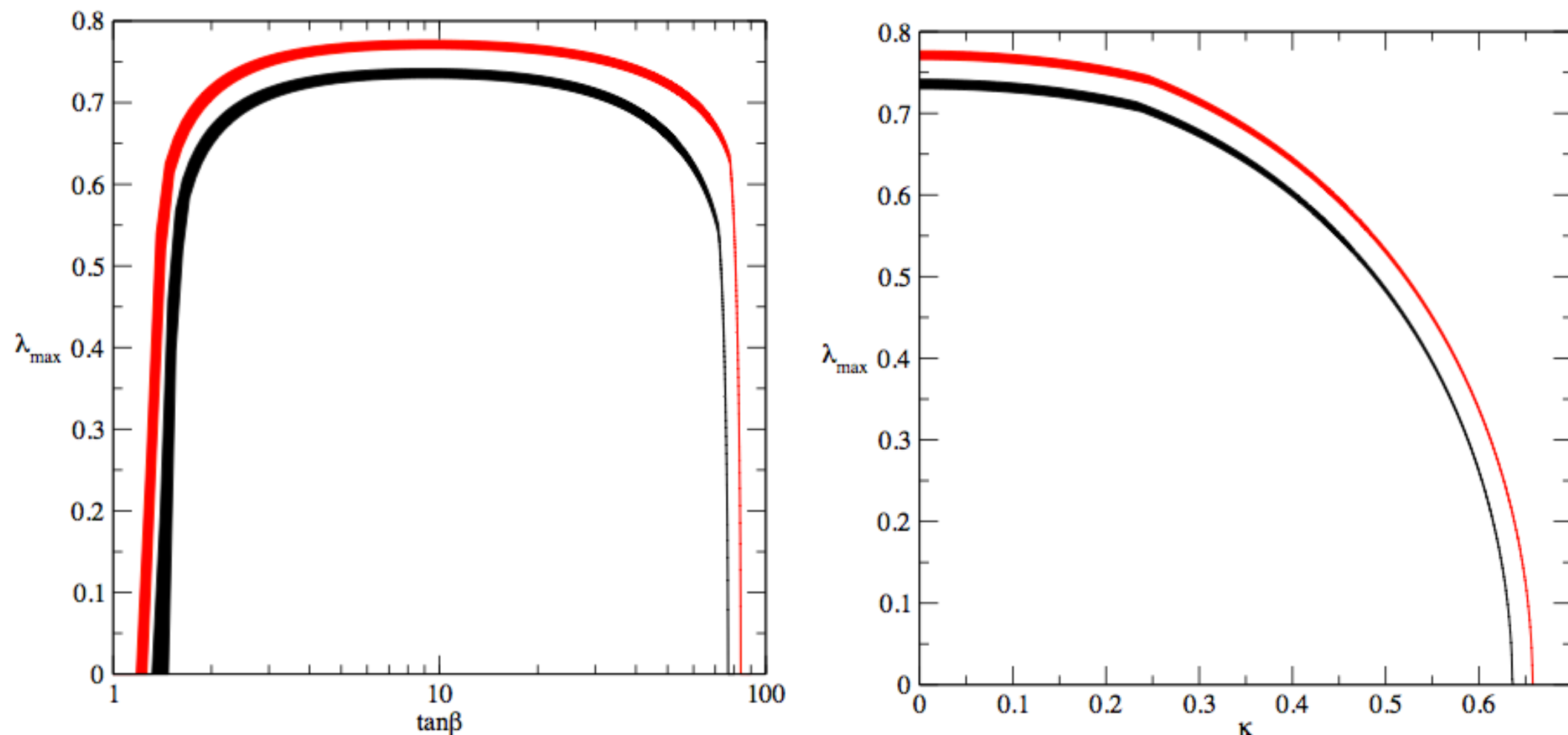


# Theoretical constraints

## ■ No unphysical minimum of the scalar potential:

- $\langle H_u^0 \rangle, \langle H_d^0 \rangle, \langle S \rangle \neq 0 \Rightarrow \kappa \leq \lambda$
- $\langle H_u^+ \rangle, \langle H_d^- \rangle = 0 \Rightarrow \lambda \leq g$

## ■ No Landau pole for $\lambda, \kappa, h_t$ below the GUT scale:



Left panel: upper bound on  $\lambda$  ( $\lambda_{\max}$ ) as a function of  $\tan\beta$  for fixed  $\kappa = 0.01$ . Right panel:  $\lambda_{\max}$  as a function of  $\kappa$  for fixed  $\tan\beta = 10$ . Black (lower) bands: light spectrum, red (upper) bands: heavy spectrum. Inside the bands the top quark mass is  $171.2 \pm 2.1$  GeV.

# Experimental constraints

After computing all observables, NMSSMTools checks:

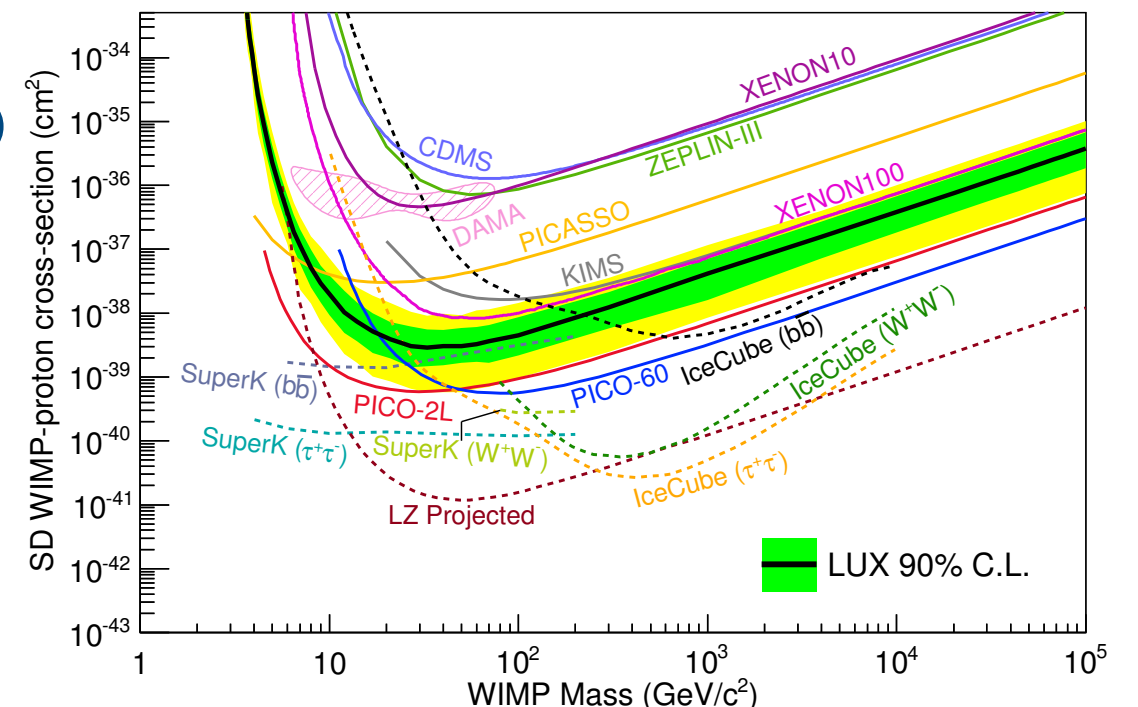
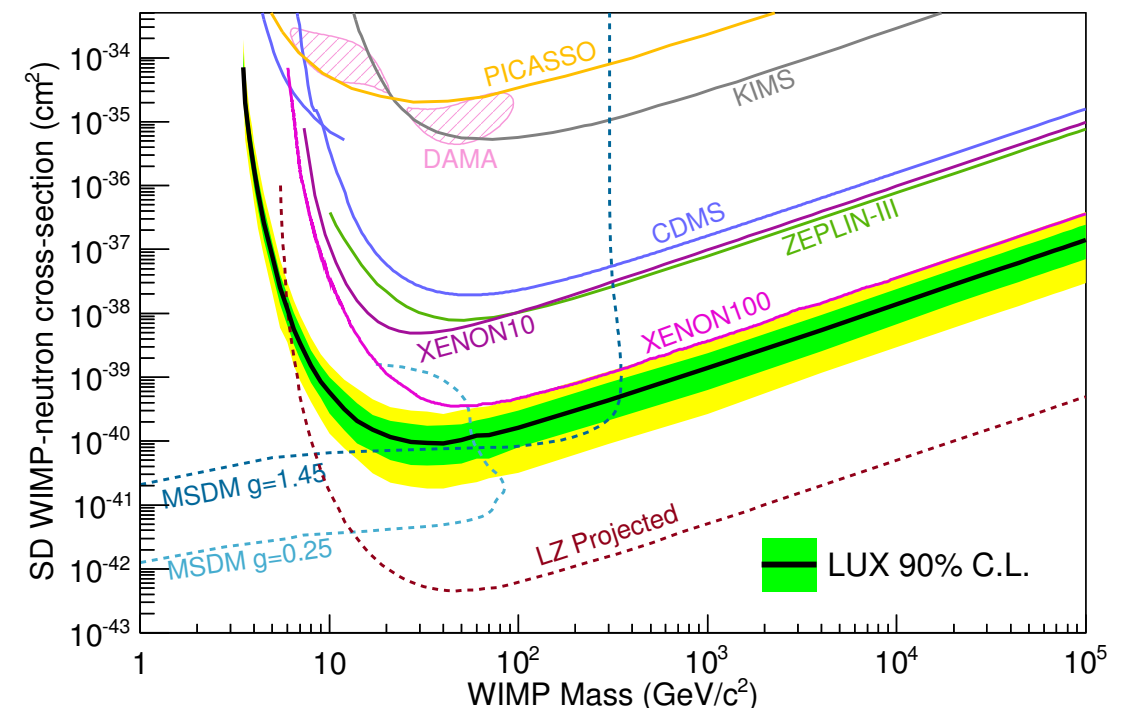
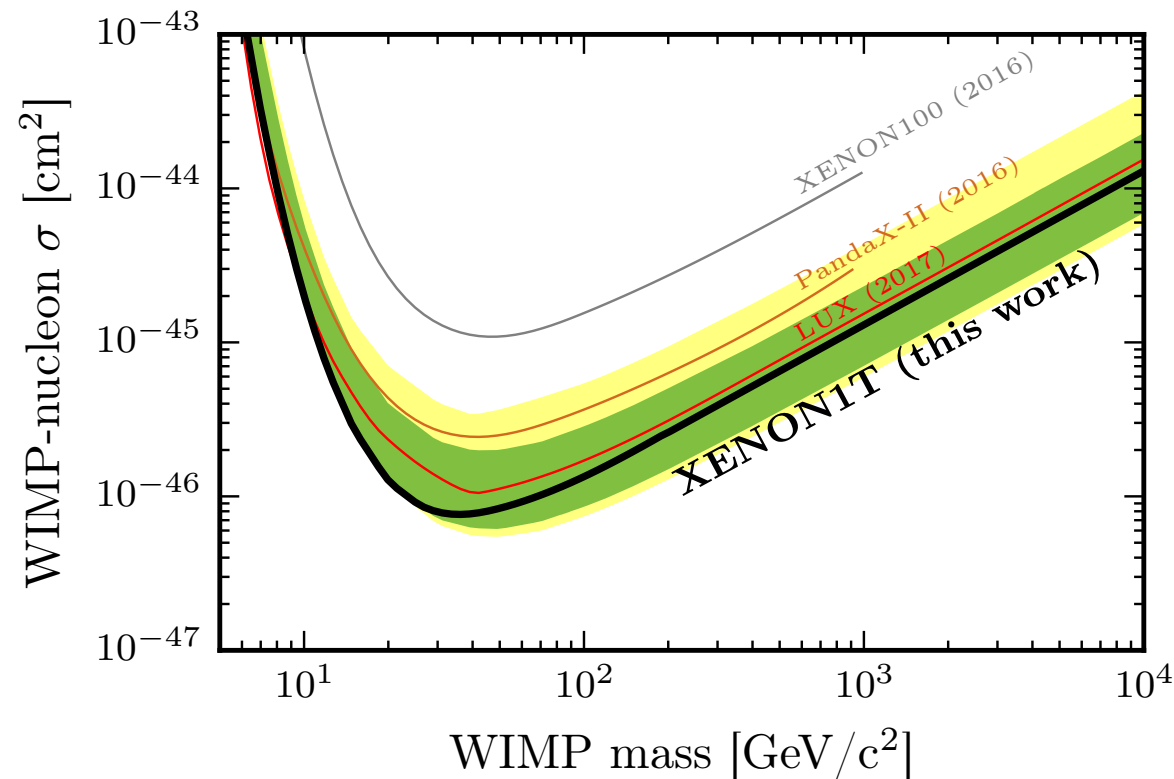
- $\tilde{\chi}_1^0$  (or gravitino) is the LSP
- LEP limits on  $\tilde{\chi}^\pm$  and  $\tilde{\chi}^0$  (direct search +  $\Gamma_{\text{inv}}(Z)$ )
- Tevatron + LEP constraints on gluinos/sfermions
- LEP limit on the charged Higgs mass  $m_{h^\pm}$
- LEP constraints from neutral Higgs searches:
  - $e^+e^- \rightarrow hZ$  with  $h \rightarrow bb, \tau\tau, jj, \gamma\gamma$ , invisible, "any" ( $Z$  recoil mass)
  - $e^+e^- \rightarrow hZ$  with  $h \rightarrow aa$  and  $a \rightarrow bb$  or  $\tau\tau$
  - $e^+e^- \rightarrow ha$  with  $h \rightarrow aa$  and  $a \rightarrow bb$  or  $\tau\tau$
  - $e^+e^- \rightarrow ha$  with  $h/a \rightarrow bb$  or  $\tau\tau$
- Constraints on very light Higgs states/mixing with light mesons
- CDF/D0, BABAR, BELLE, CLEO and LHCb limits on  $B$ -physics
- BNL constraint on  $(g-2)_\mu$  from  $e^+e^-$  data ( **$3\sigma$  away from SM**)



# Dark Matter constraints

■ Planck constraints:  $\Omega h^2 = 0.1187 \pm 10\%$  (or just upper bound)

■ SI limits from LUX, PandaX, XENON-1T

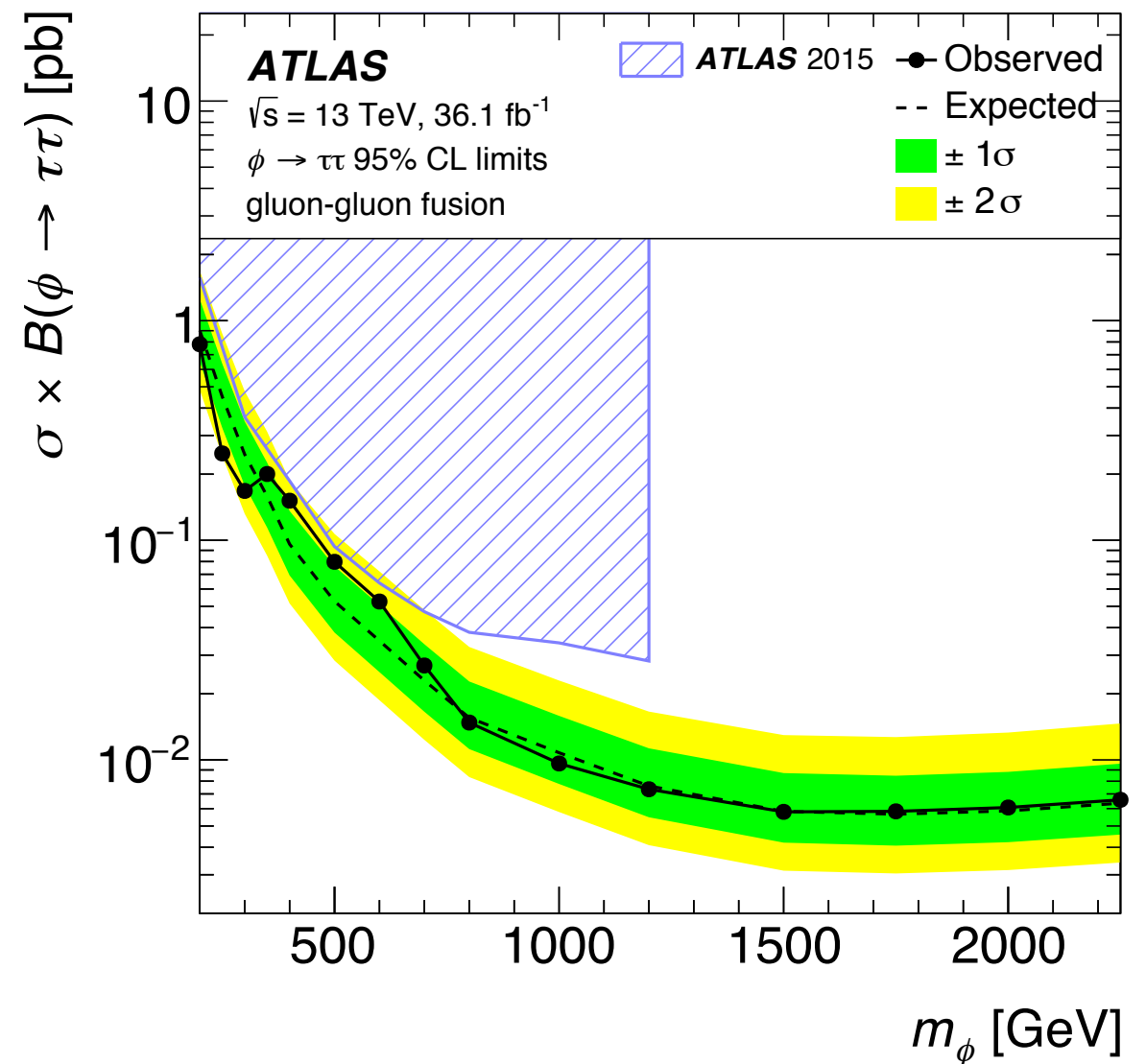
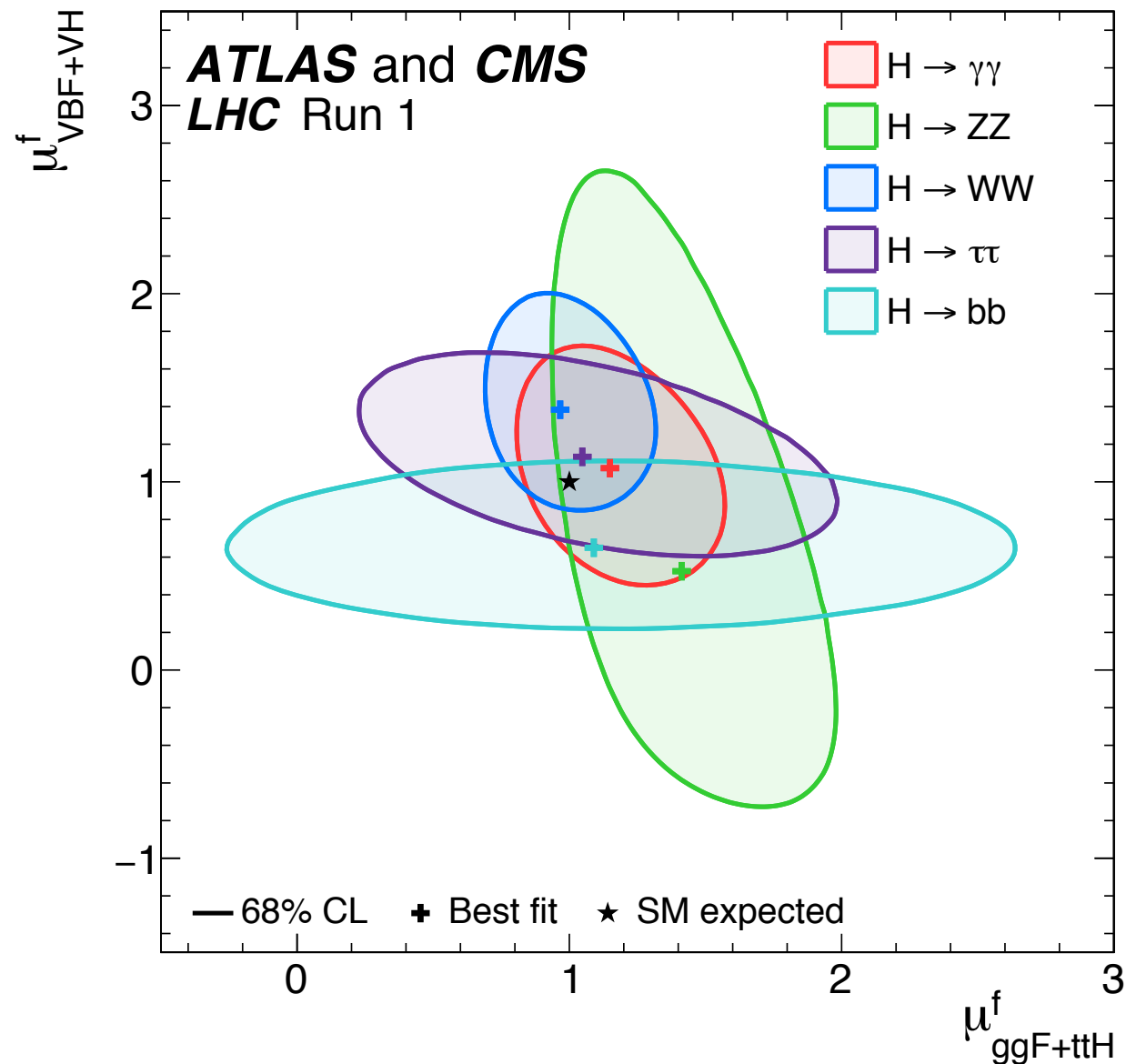


■ SD limits from LUX, PandaX, PICO  
(from neutron and proton Xsections)

■ Relic density and direct detection  
cross sections from **micrOMEGAs**

# LHC constraints

- Limit on the gluino mass:  $m_g > 1.6 \text{ TeV}$
- Mass of the observed (SM-like) Higgs state:  $m_h = 125.1 \pm 3 \text{ GeV}$
- Couplings of the observed Higgs:
- Constraints on heavy Higgs states  $H/A \rightarrow bb$  or  $\tau\tau$



# Light singlino/higgsino scenario

## ■ Light higgsinos are natural:

- Mass given by  $\mu$  (MSSM) or  $\mu_{\text{eff}}$  (NMSSM)
- EWSB requires fine-tuning if  $\mu_{\text{eff}} \gg M_Z$

## ■ But are not good dark matter candidates:

- Relic density too small unless  $|\mu_{\text{eff}}| > 1 \text{ TeV}$
- Large Direct Detection rates (SD via Z-exchange)

## ■ In the MSSM:

- bino LSP: relic density too large
- bino-higgsino mixing: DD rates too large unless  $m_{\text{LSP}} > 1 \text{ TeV}$

## ■ In the NMSSM: singlino LSP!

## ■ Mass relation in singlet sector: $m_{\tilde{S}}^2 = m_S^2 + \frac{1}{3}m_P^2 - \frac{4}{3} \left( \lambda^2 \frac{A_\lambda}{\mu_{\text{eff}}} + \kappa \right) v_u v_d$

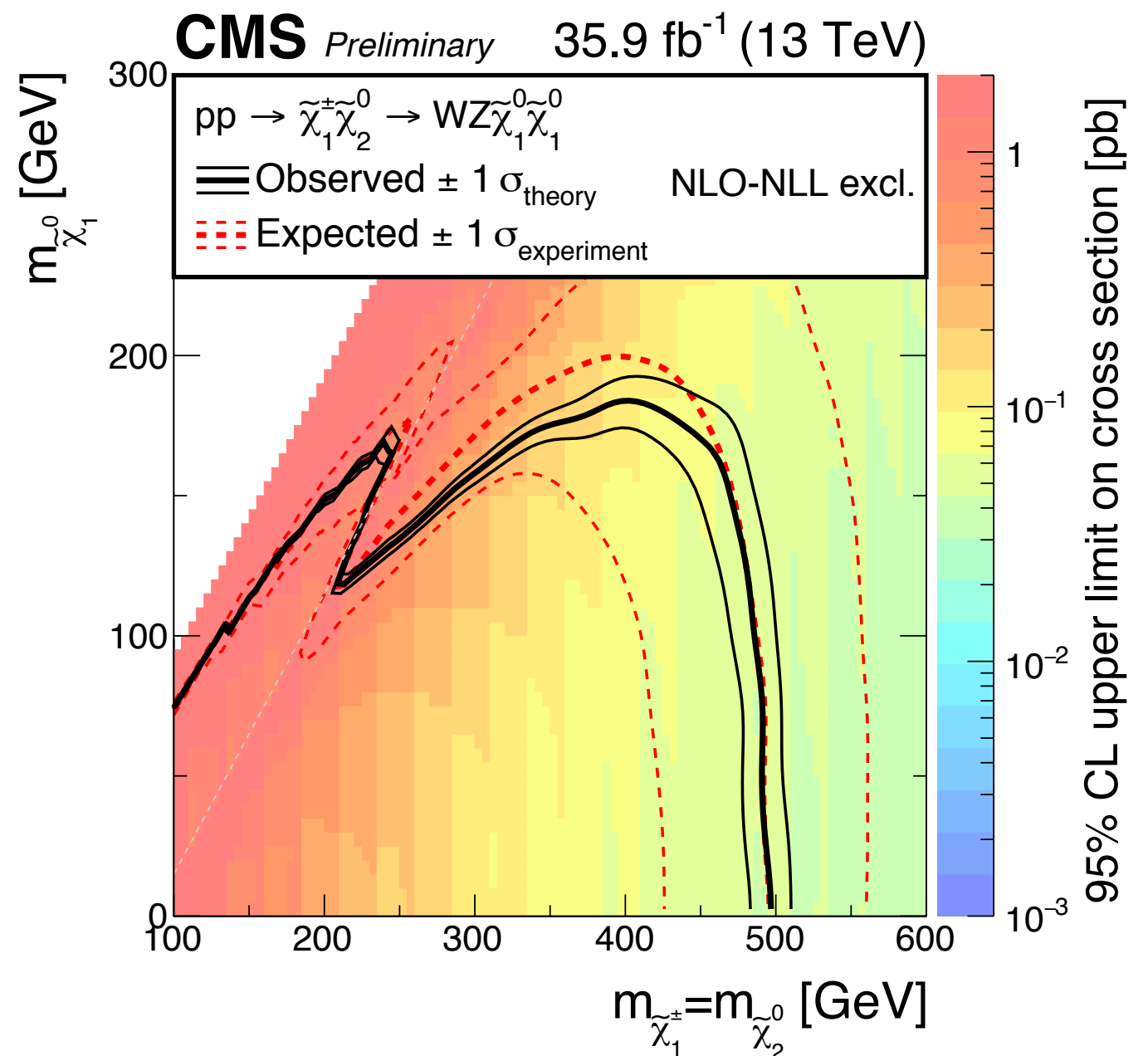
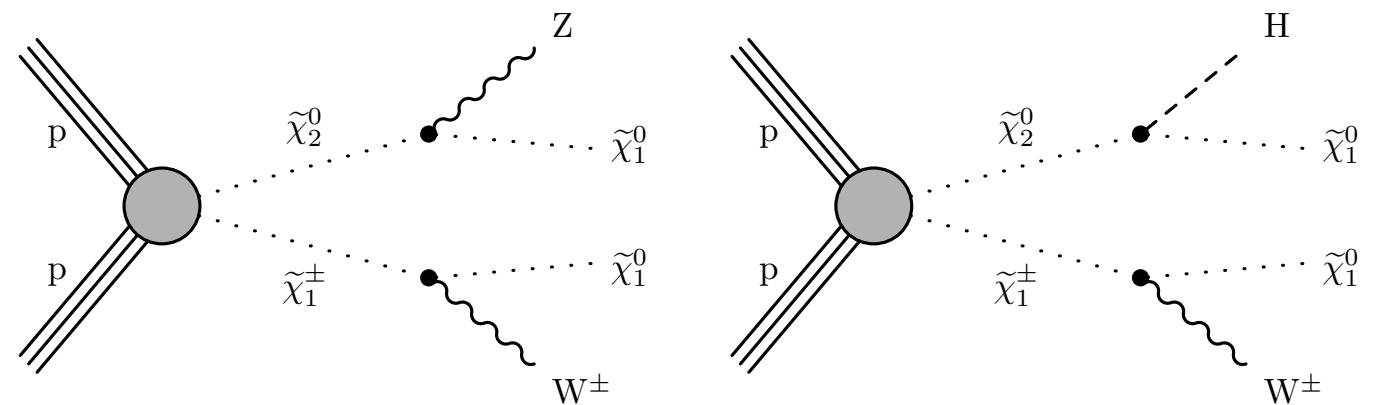
Last term negligible in decoupling limit ( $\lambda, \kappa \sim 0$ ) or large  $\tan\beta$

# Singlino dark matter

- Efficient processes to get correct relic density:
  - annihilation via the pseudo-scalar singlet  $P$  in the s-channel  
requires  $m_P \simeq 2m_{\tilde{S}} \Rightarrow 3m_{\tilde{S}}^2 + m_S^2 = 4 \left( \lambda^2 \frac{A_\lambda}{\mu_{\text{eff}}} + \kappa \right) v_u v_d$   
hence  $\lambda \gtrsim .2$  and  $\tan\beta$  not too large unless  $m_{\tilde{S}} \lesssim 20 \text{ GeV}$
  - $h$  or  $Z$  exchange in the s-channel if  $m_h \simeq 2m_{\tilde{S}}$  or  $M_Z \simeq 2m_{\tilde{S}}$
  - co-annihilation with higgsinos if  $m_{\tilde{S}} \simeq |\mu_{\text{eff}}|$
- Note: singlino/singlet coupling  $\sim \kappa \implies$  no higgsino mixing needed
- Still: LSP acquires higgsino component unless mass splitting large
- Direct Detection limits affect the present scenario if  $m_{\tilde{S}} \gtrsim 5 \text{ GeV}$ 
  - SI: singlet Higgs exchange in the t-channel, negative interference with the SM-like Higgs  $\implies$  cross section possibly below v floor
  - SD: weaker constraint but t-channel  $Z$  exchange more efficient  
 $Z/\text{LSP}$  coupling from higgsino component  $\implies$  upper bound on  $\lambda$   
For light LSP: similar upper bound on  $\lambda$  from  $Z$  invisible width

# LHC constraints on light higgsinos

- Search for  $W^* \rightarrow \tilde{\chi}_1^\pm + \tilde{\chi}_2^0$ ,  
with  $m_{\tilde{\chi}_1^\pm} = m_{\tilde{\chi}_2^0}$  (winos)  
and  $\tilde{\chi}_1^0 = \text{bino LSP}$
- H is the  $\sim 125$  GeV Higgs  
Best sensitivity in  $3l$  ( $W, Z$ )
- Result: limit on cross section times branching ratios in the plane  $m_{\tilde{\chi}_1^\pm} = m_{\tilde{\chi}_2^0}$ ,  $m_{\tilde{\chi}_1^0}$
- Line: assuming pure winos and 100% branching ratios
- How does this translate for light singlino/higgsinos?





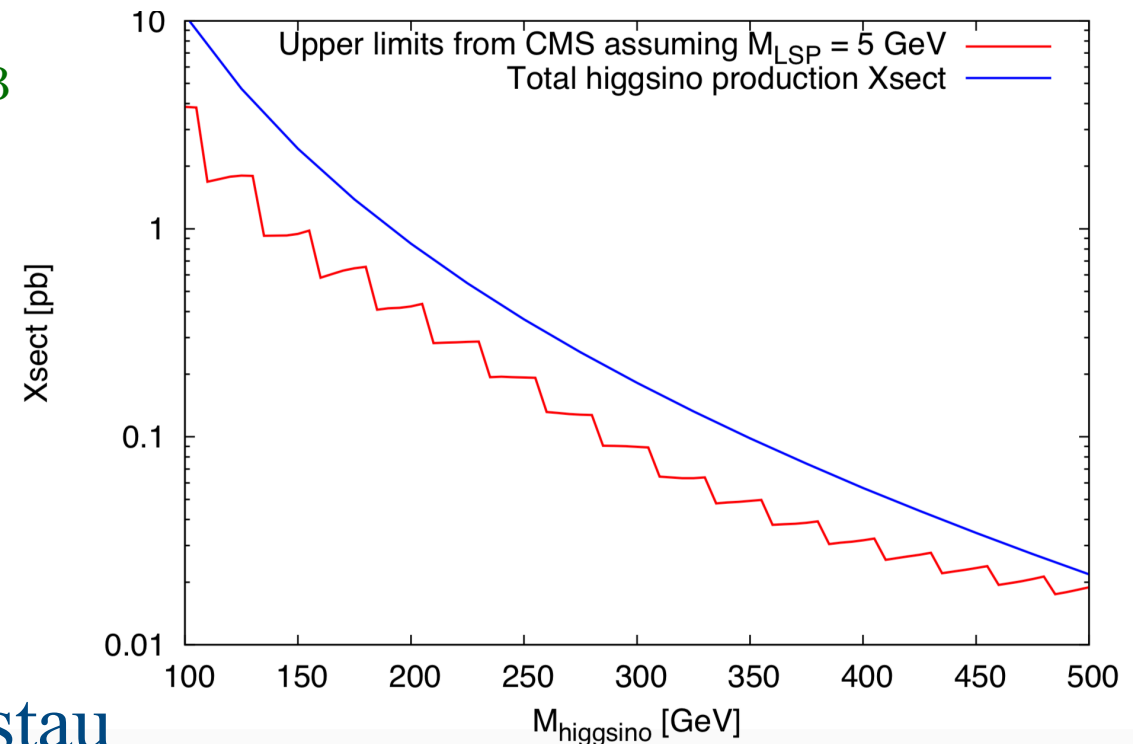
# And how to avoid them

- 1 charged higgsino  $\chi_1^\pm$ , 2 neutrals  $\chi_{2,3}^0$

Assuming mass degenerate states

total cross section  $\sim$  half of winos

For  $m_{\text{LSP}} = 5$  GeV: not excluded if  
branching ratios to  $W, Z < 100\%$



- $\chi_1^\pm \rightarrow W^{\pm(*)} + \chi_1^0$  unless light bino or stau

- $\chi_{2,3}^0 \rightarrow Z^{(*)} + \chi_1^0$  and  $\chi_{2,3}^0 \rightarrow h^{(*)} + \chi_1^0$  are of order  $\lambda^2$ : comparable BRs

$\chi_{2,3}^0 \rightarrow h^{(*)} + \chi_1^0$  limits much weaker  $\Rightarrow$  consider only  $\chi_{2,3}^0 \rightarrow Z^{(*)} + \chi_1^0$

Alleviates constraints, notably for  $M_1 \sim |\mu_{\text{eff}}|$  (bino component of  $\chi_{2,3}^0$ )

- $\chi_{2,3}^0 \rightarrow S/P + \chi_1^0$  (if kinematically allowed) suppressed compared to  $Z/h$

However if the latter are kinematically forbidden, BRs to  $S/P = 100\%$

- Other loophole: light bino (still singlino LSP) then higgsinos  $\chi_{3,4}^0 \rightarrow \chi_2^0$   
considered (conservatively) not to contribute to the signal



# Recasting procedure and results

- $\chi_1^\pm$  and  $\chi_{2,3}^0$  not mass degenerate ( $m_{\chi_3^0} - m_{\chi_1^\pm} \simeq m_{\chi_1^\pm} - m_{\chi_2^0} \lesssim 20$  GeV)
  - ➔ find the  $\chi_1^\pm \chi_i^0$  production cross sections (from an array)
  - ➔ weigh by  $X_i/(X_1+X_2)$  where  $X_i = \chi_1^\pm \chi_i^0 W$  reduced coupling squared (production) times branching ratios to  $W, Z$  (decay)
  - ➔ Find a degenerate triplet (mass  $m_T$ ) with same cross section
  - ➔ Compare to the cross section excluded by CMS at  $m_T, m_{\chi_1^0}$

- Results in the pNMSSM

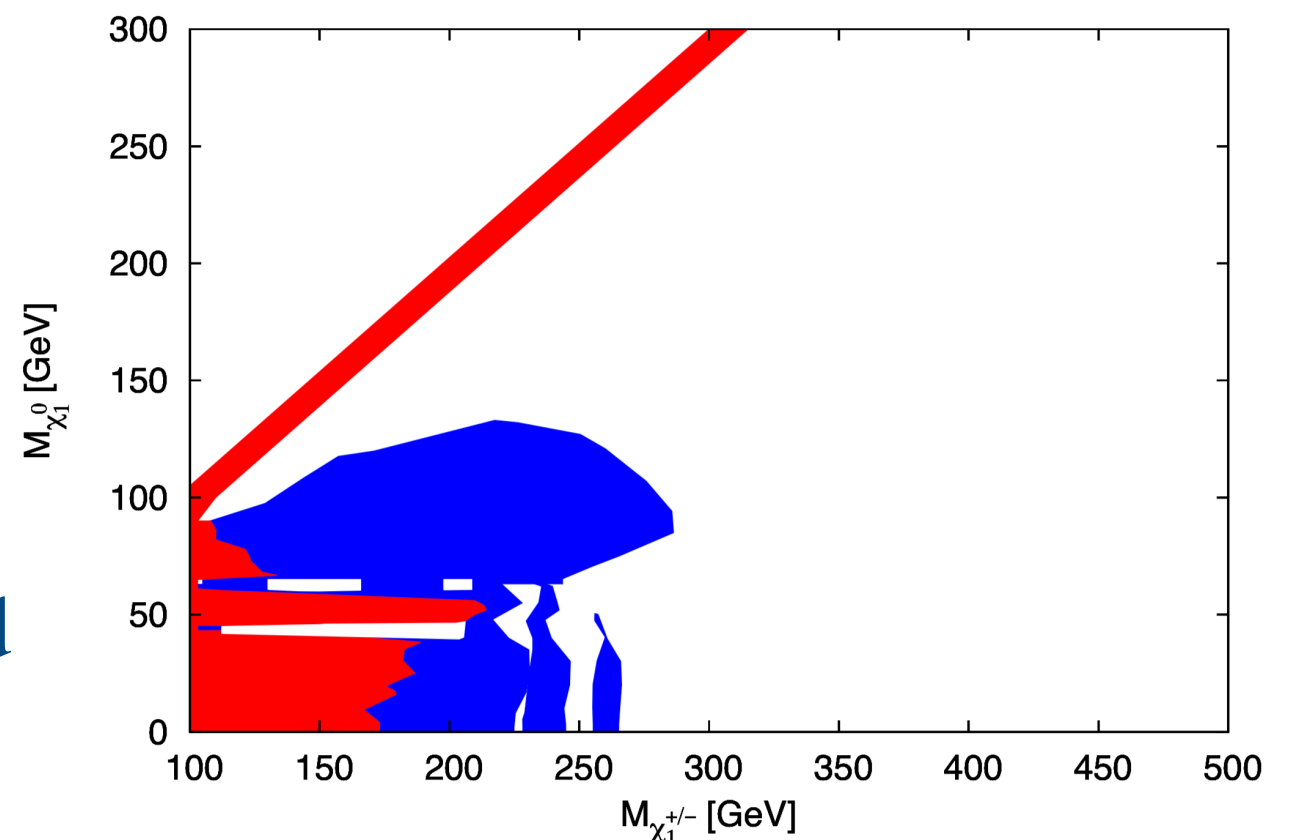
- ➔ red: excluded for arbitrary  $M_1$

- ➔ blue: for  $M_1 > 300$  GeV

Diagonal band: coannihilation

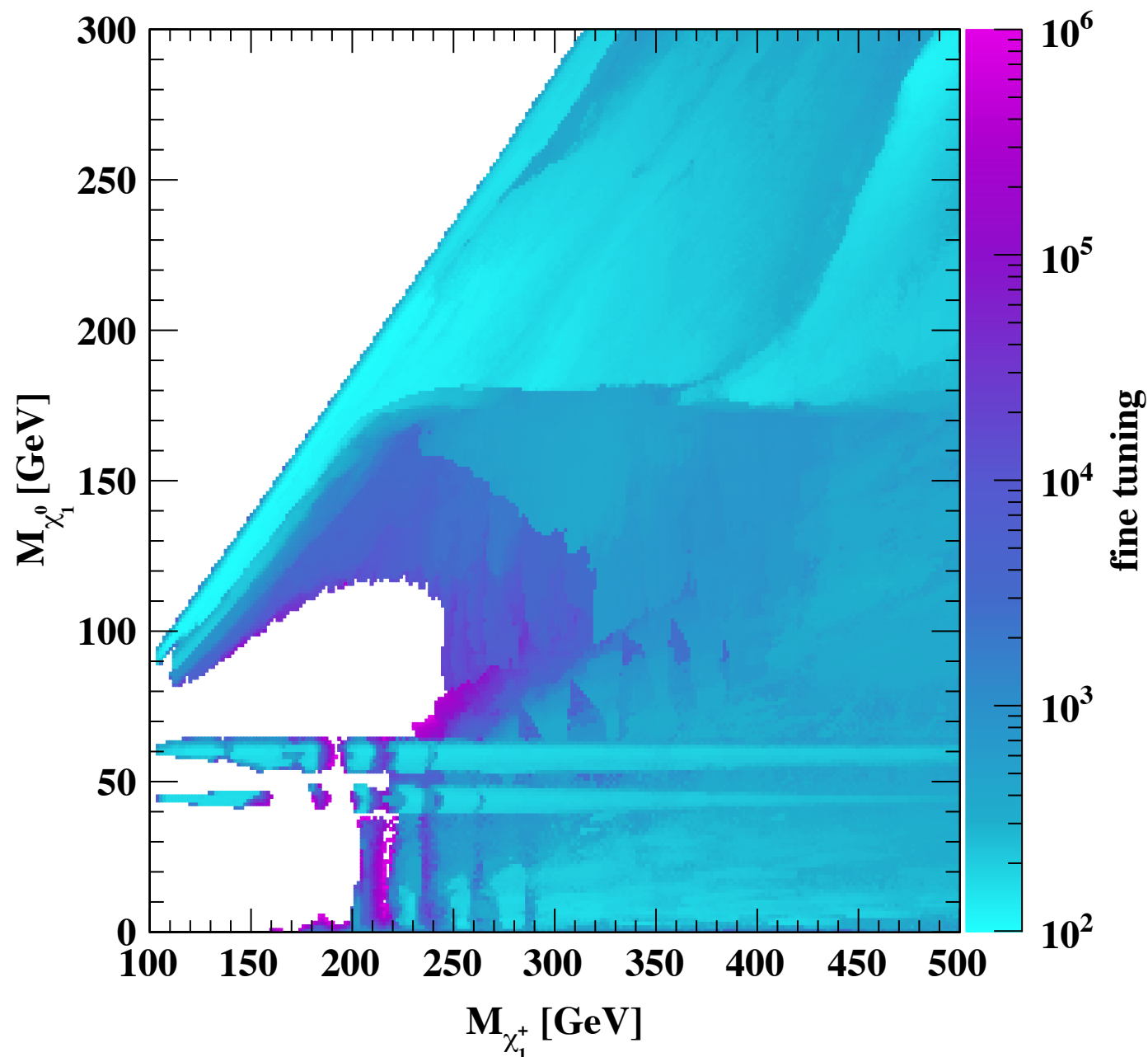
Near  $Z/h$  funnel:  $S/P$  not needed

for annihilation:  $\chi_{2,3}^0 \rightarrow S/P + \chi_1^0$



# Results in the NUH-NMSSM

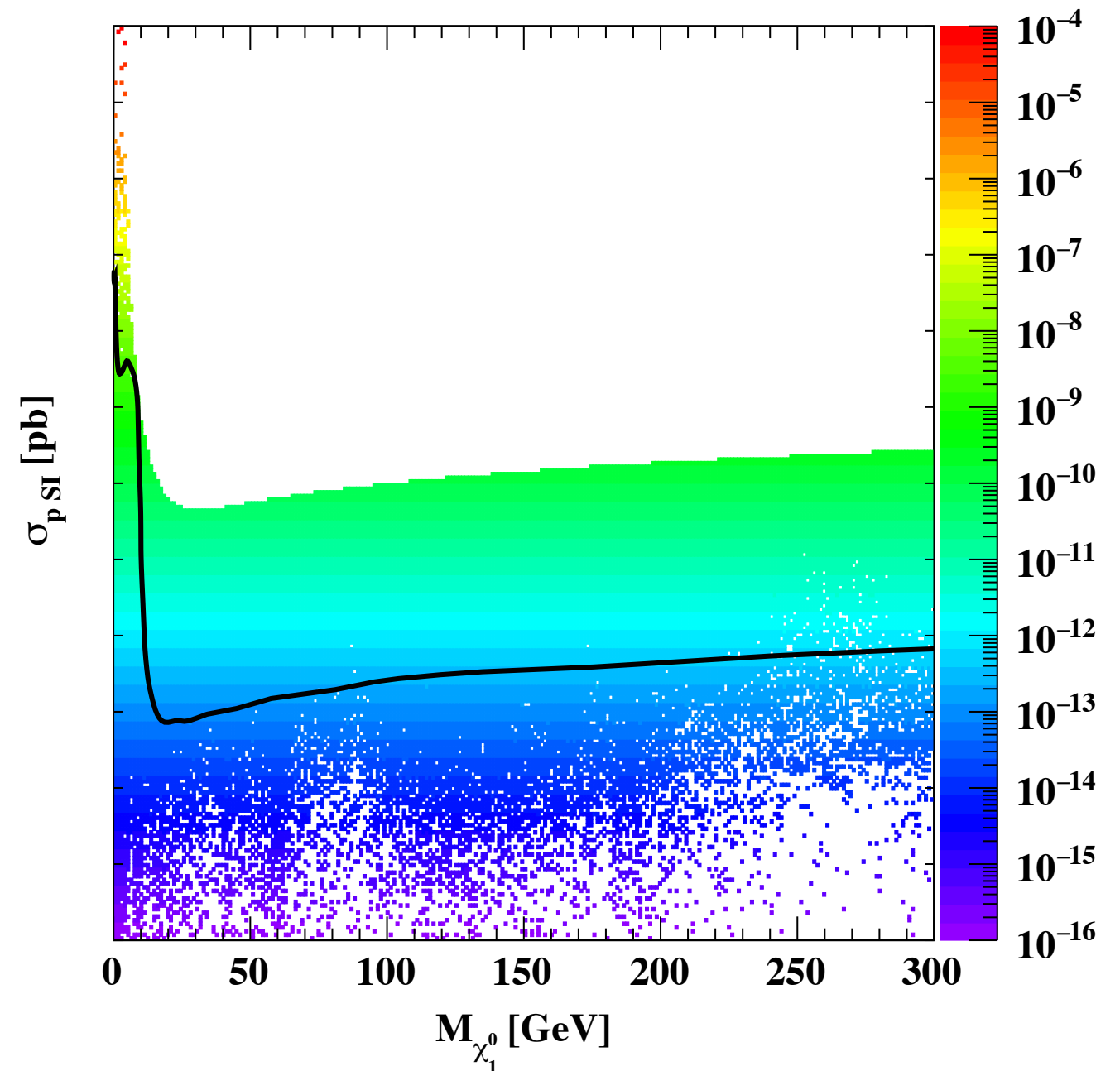
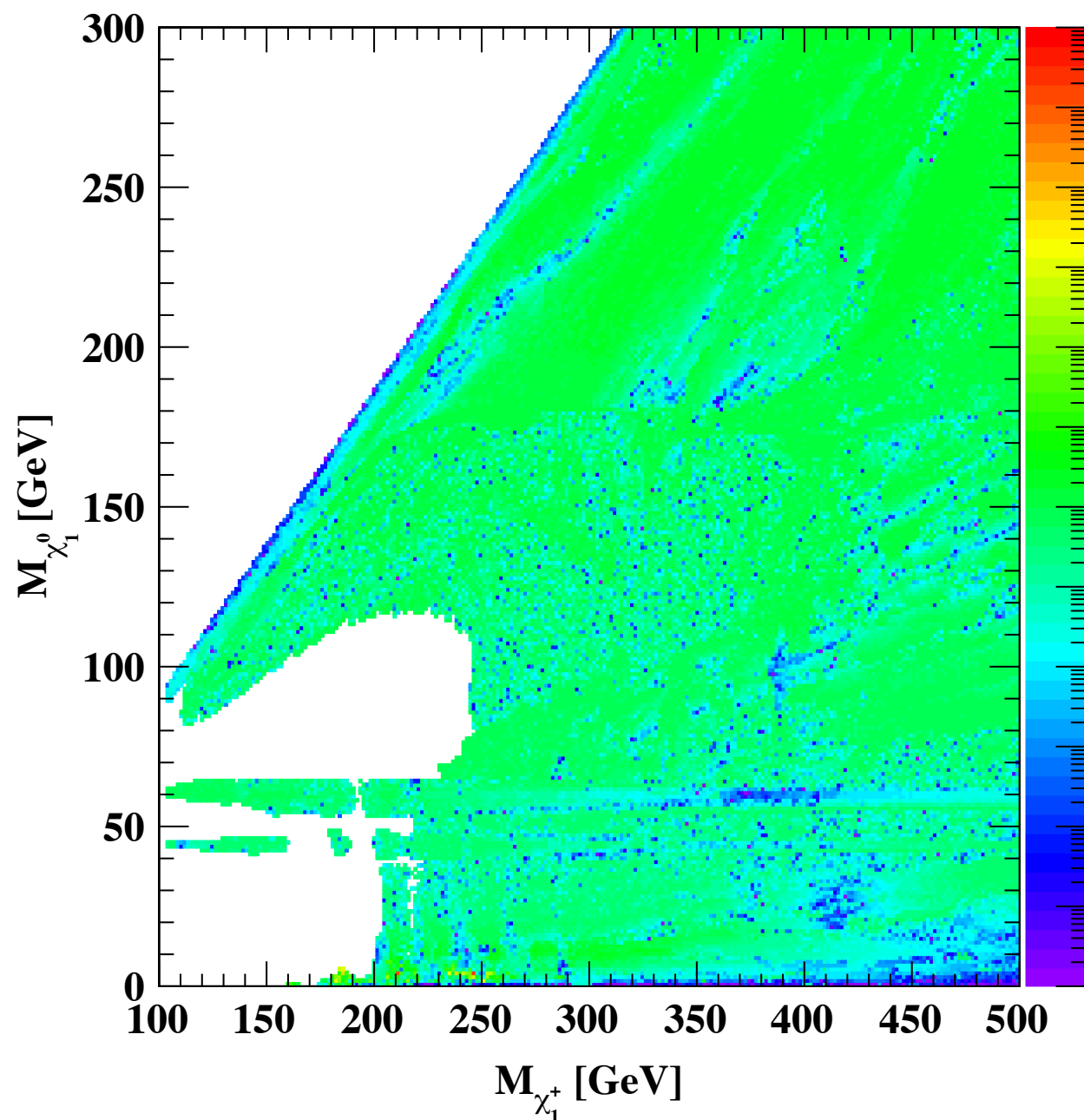
- Large MCMC scans of NUH-NMSSM parameter space ( $\sim 10^{12}$  pts)
- ➡ Fine-tuning map in the  $m_{\chi_1^\pm}, m_{\chi_1^0}$  plane divided in a 2D histogram
- ➡ In each bin: keep the point with the lowest fine-tuning



- Excluded (white) region = red+part of blue in pNMSSM
- Low FT regions (light blue):
  - close to  $Z/h$  funnels
  - coannihilation region
  - light singlino ( $m_{\chi_1^\pm} > 20$  GeV)
  - $P$  + off-shell  $Z$  ( $m_{\chi_1^\pm} > m_{top}$ )
- Large FT in regions where CMS constraints are strong

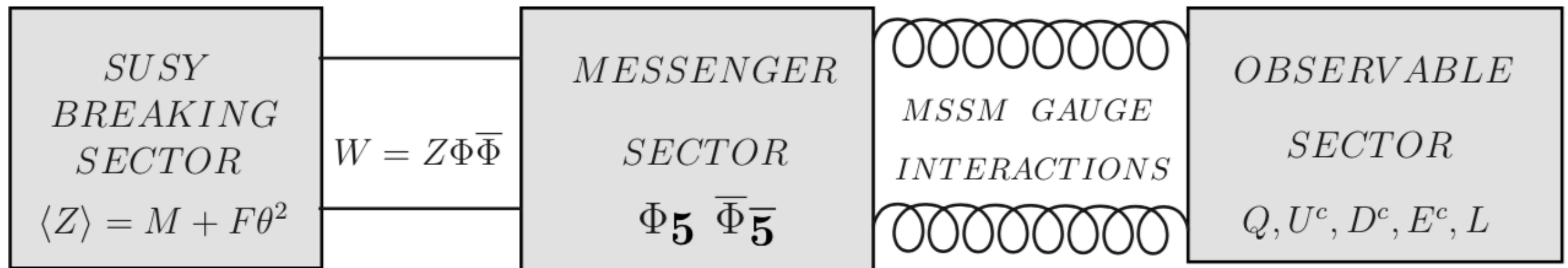
# Direct detection rates

- SI Direct Detection rates can fall below the  $\nu$  floor
  - in regions where LSP is very singlino-like
  - elsewhere due to  $S/h$  negative interference



# Gauge mediated SUSY breaking

- SUSY has too many free parameters  $\Rightarrow$  predictive scenarios?



- Minimal GMSB is a very predictive setup
- It solves easily the SUSY flavour problem
- Problem: it is hard to obtain a 125 GeV Higgs
- Solutions:
  1. New tree level  $\Delta m_h$  at messenger scale, possible in the **NMSSM** with extra Higgs singlet + singlet/messenger coupling
  2. Large  $A_t$  at messenger scale from matter/messenger coupling

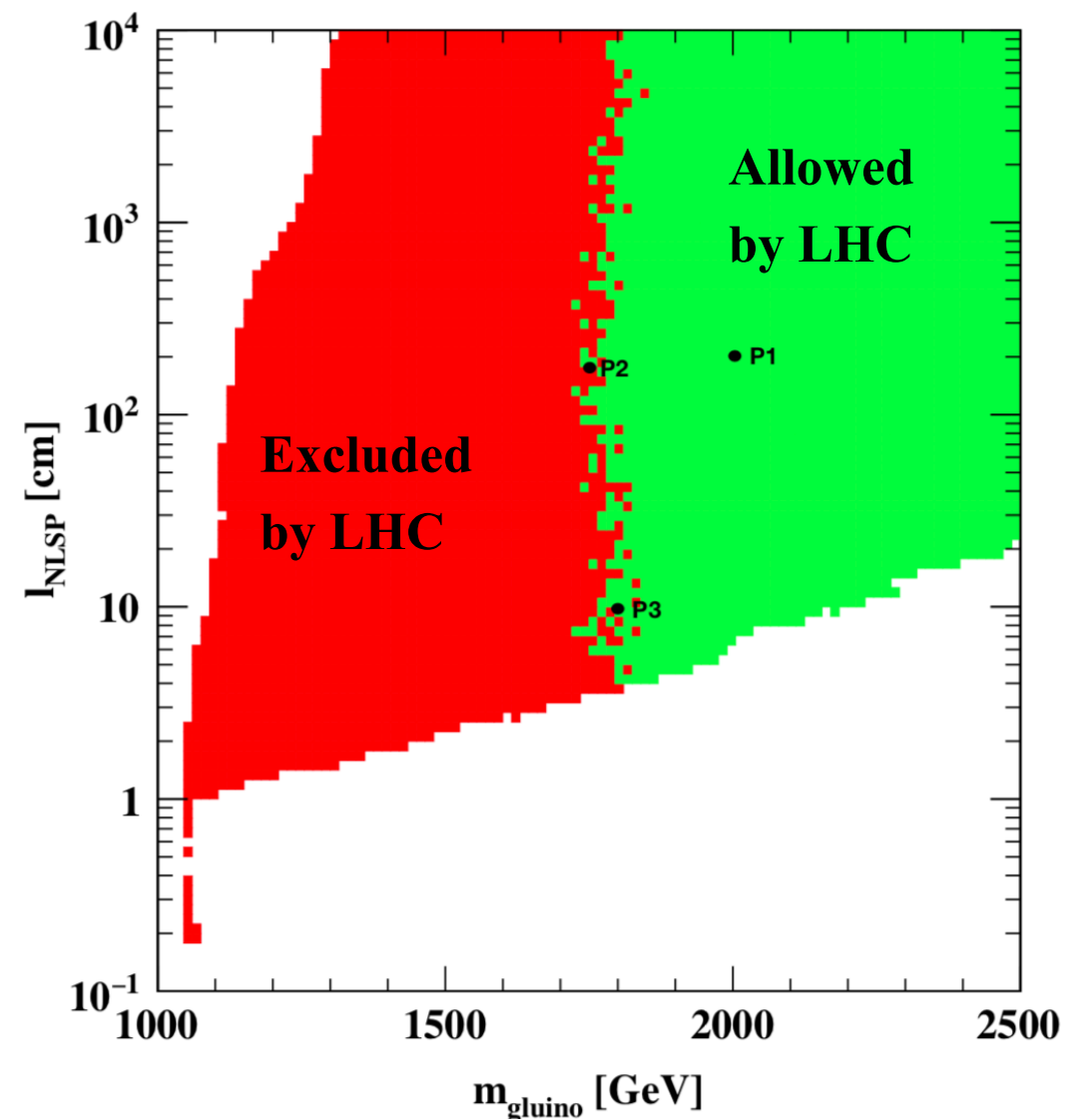
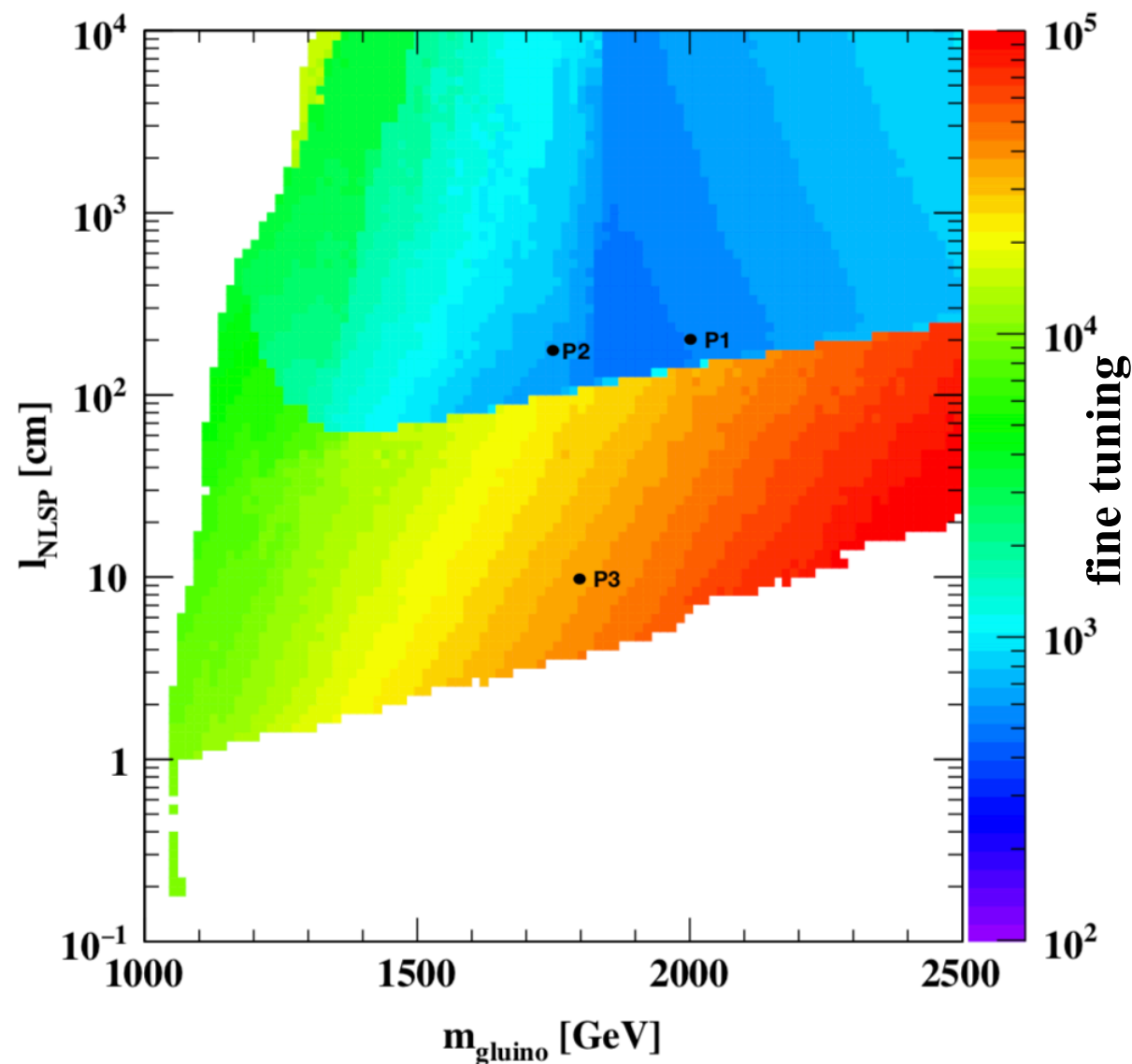
# Solution 1: the DGS model

- 2 messengers  $\Phi_u, \Phi_d$  coupled to the singlet [Delgado, Giudice, Slavich '07]
- Only 4 parameters (all other fixed by correct EWSB):
  - $\lambda$ : Higgs singlet/doublet coupling  $H_u H_d S$
  - $\xi$ : singlet/messenger coupling  $\Phi_u \Phi_d S$
  - $M$ : messenger scale
  - $\Lambda$ : SUSY scale
- **The push-up scenario** [Allanach, Badziak, Hugonie, Ziegler '15]:
  - $m_{h1} \sim 90$  GeV (singlet-like),  $m_{h2} \sim 125$  GeV (SM-like)  $\Rightarrow \lambda, \xi \sim 10^{-2}$
  - Additional light pseudoscalar Higgs with  $m_a \sim 25$  GeV
  - Gluino/squarks as light as  $\sim 1.5$  TeV (depending on  $\Lambda$ )
  - Singlino NLSP  $\sim 100$  GeV decaying to gravitino + pseudoscalar
- Possibly displaced singlino to gravitino decay (depending on  $M$ )
  - First studied in [Allanach, Badziak, Cottin, Desai, Hugonie, Ziegler '16]
  - Updated in [Badziak, Desai, Hugonie, Ziegler '18]: new constraints on the Heavy Higgs sector push the gluino mass beyond the LHC reach...



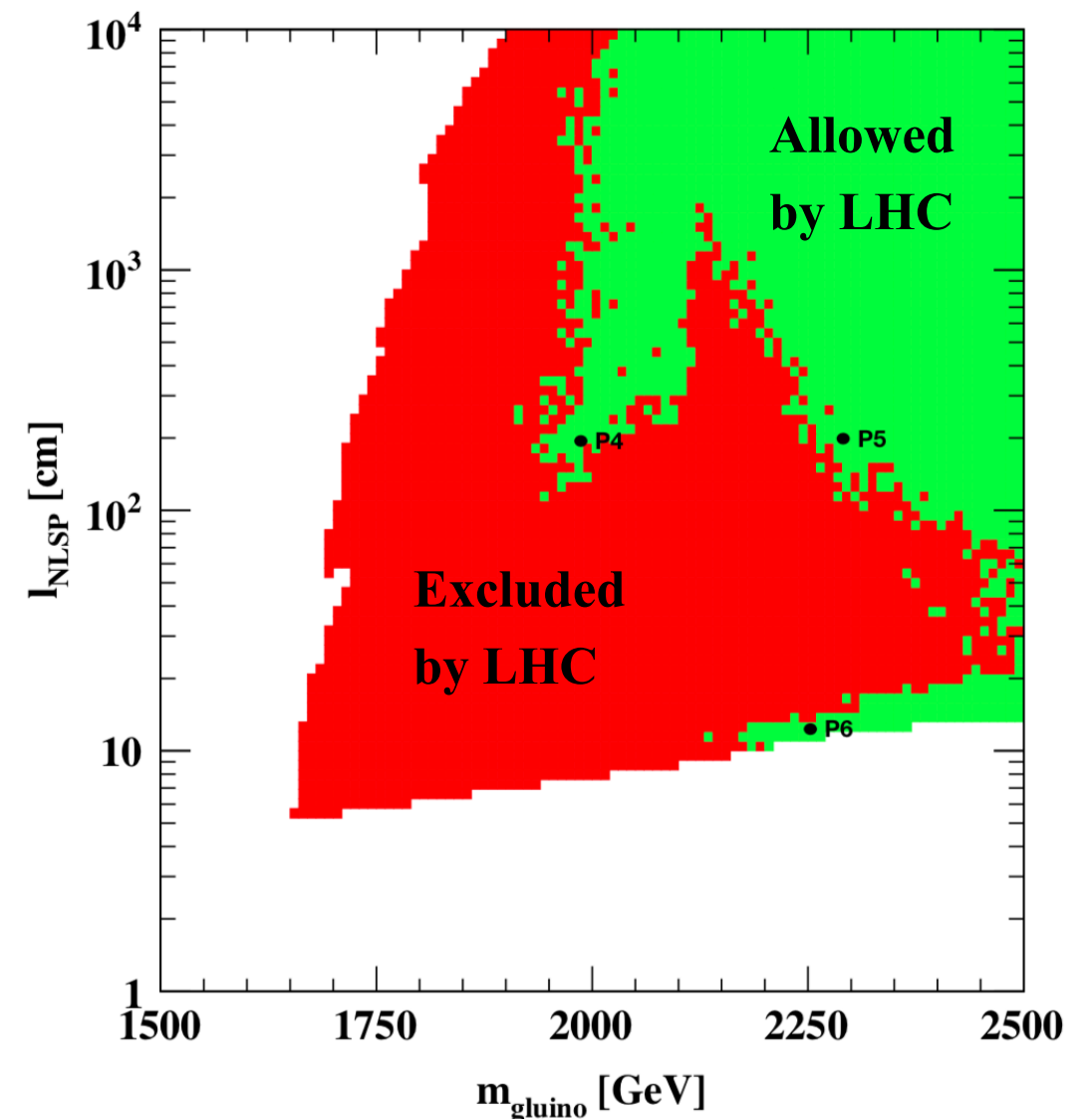
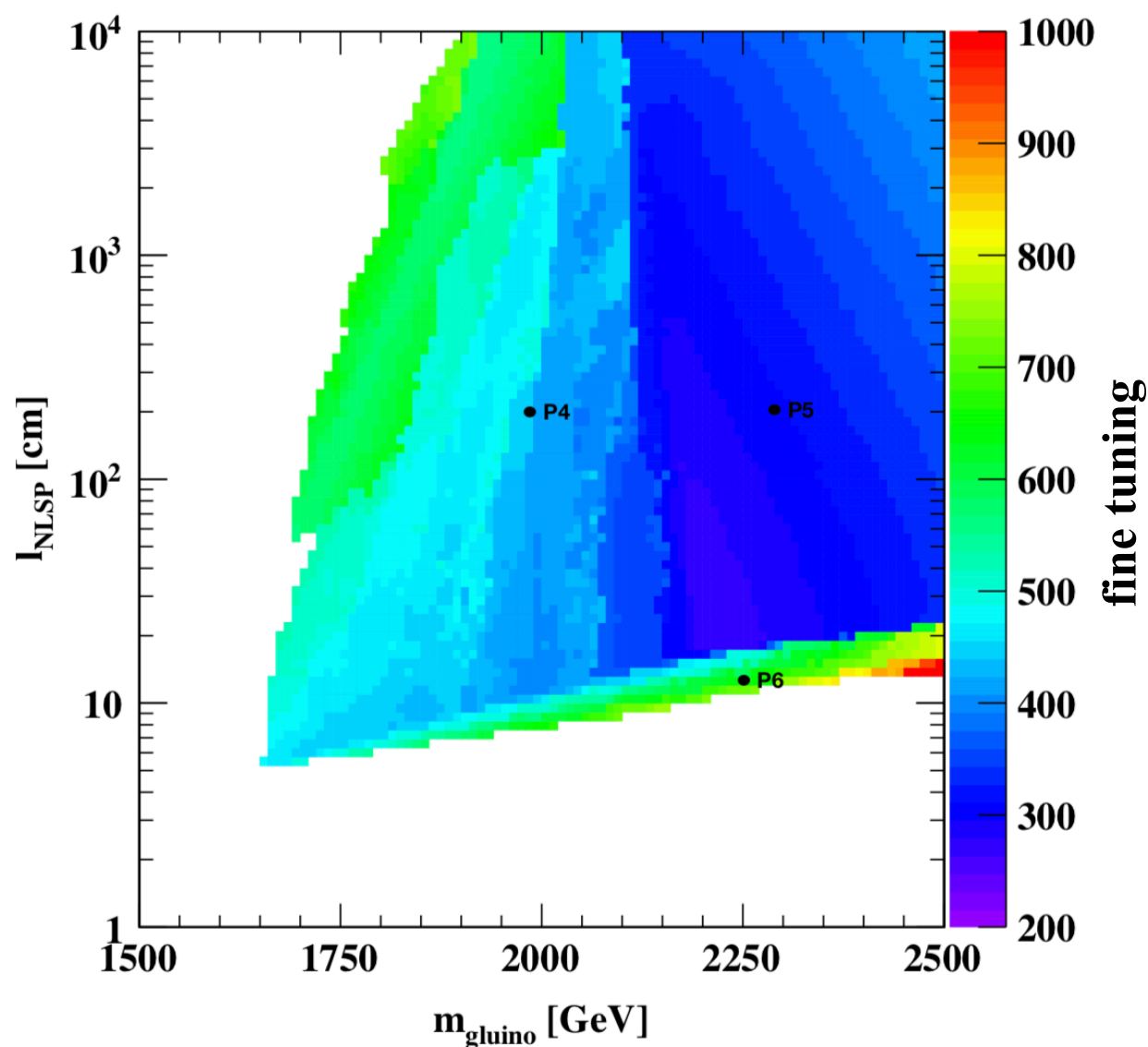
# Solution 2: the U model

- 1 messenger  $\Phi_u$  coupled to matter and Higgs sector
- Now 5 free parameters:
  - $\lambda_t$ :  $\Phi_u Q U$  coupling
  - $M$ : messenger scale
  - $\lambda$ :  $H_u H_d S$  coupling
  - $\lambda_S$ :  $\Phi_u H_d S$  coupling
  - $\Lambda$ : SUSY scale
- Possibly light gluino with displaced singlino NLSP to gravitino LSP decay!



# Solutions 1 & 2: the DGSU model

- Mixture of the 2 previous models with 2 messengers  $\Phi_u, \Phi_d$
- Still 5 free parameters:
  - $\xi$ :  $\Phi_u \Phi_d S$  coupling
  - $M$ : messenger scale
  - $\lambda$ :  $H_u H_d S$  coupling
  - $\lambda_t$ :  $\Phi_u Q U + \Phi_u S H_d$
  - $\Lambda$ : SUSY scale
- Again light gluino with displaced singlino NLSP to gravitino LSP decay



# Conclusions

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- Light singlino/higgsino scenario still alive and kicking
- Could be challenging to test
  - Small DD rates (possibly below the  $\nu$  floor)
  - $\chi_{2,3}^0 \rightarrow S/P + \chi_1^0$  difficult channel at the LHC
- Benchmark points given in [arXiv:1806.09478]
  - ➡ Test ATLAS sensitivity to non degenerate higgsinos
- Displaced singlino NLSP to gravitino LSP vertices in GMSB models could also be challenging signals [arXiv:1810.05618]