

Spintronics, Multiferroics and magnetic chirality



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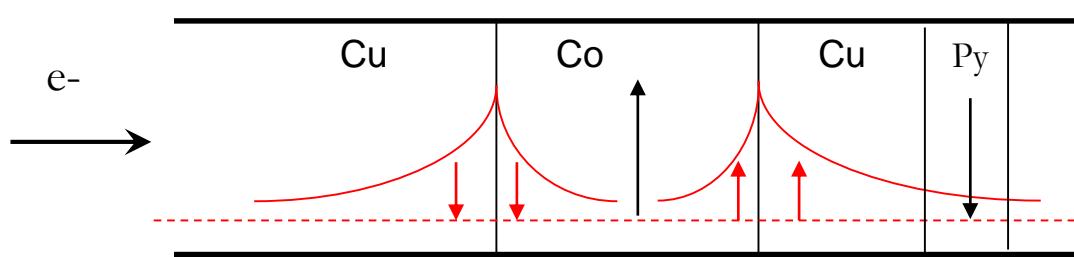
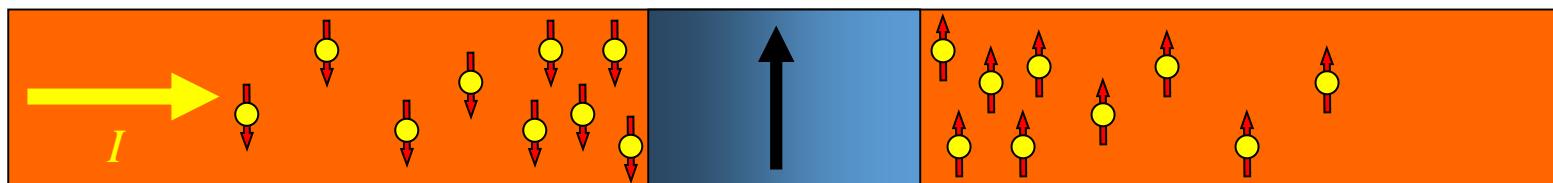
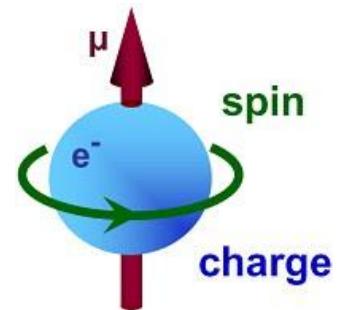
P. Thibaudeau, J. Tranchida: CEA-DAM

Spintronics

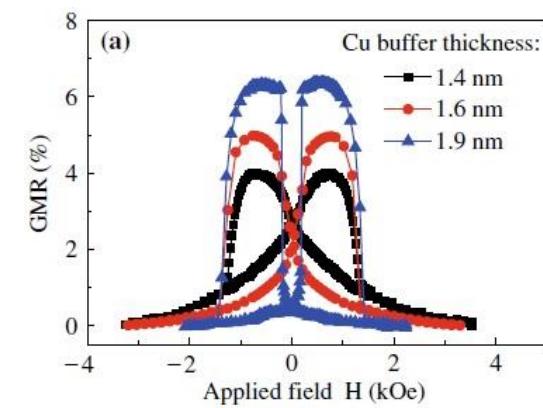
« Classical electronics has forgotten the spin of the electron »

→ SPINTRONICS

Based on spin effects and born with the spin valve (A. Fert, P. Grunberg, late 80's)



→ Magneto-resistance in multilayers



From discovery to applications in record time: Computer read head, MRAMs

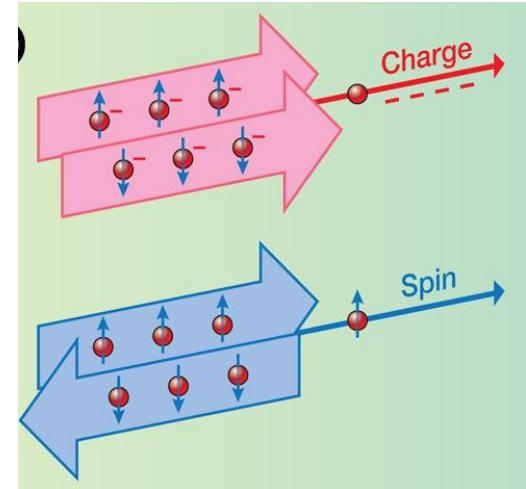
Present issues

Major present problem: **dissipation**...!

Do we really need moving charges?

→ **Pure spin currents**

→ **Electric field control of magnetism**



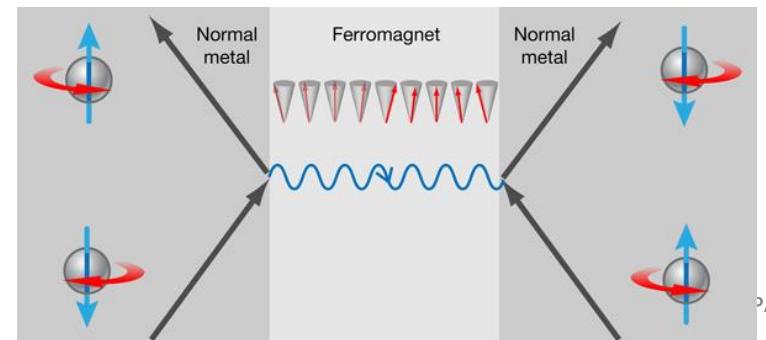
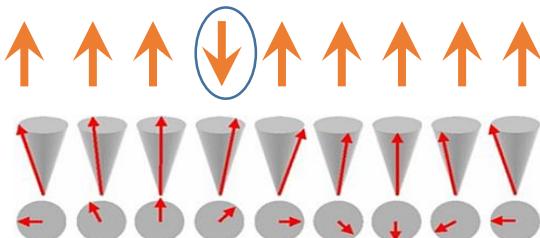
Elements needed to design devices:

- Spin sources → generate spins
- Spin conductors → propagate spins
- Spin sensors → transform spin into voltage/current

Fundamental research on pure spin transport, spin/charge conversion, multiferroics

Introduce insulators in spintronics ex:

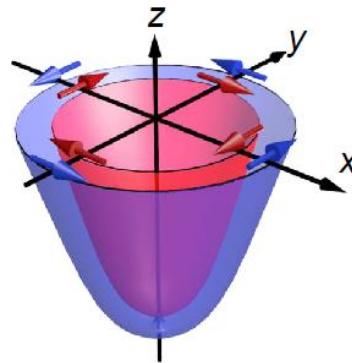
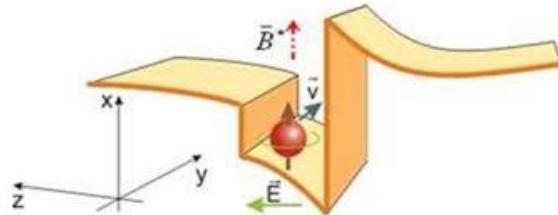
Propagating spins using the elementary magnetic excitations: the magnons



Spin orbitronics

Spin-orbit effects converting spin and charge

Rashba-Edelstein Effect :

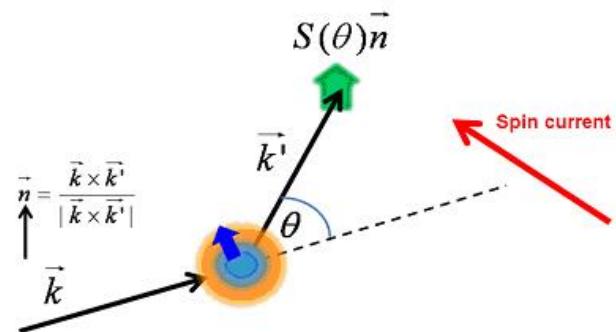


Built-in electric field at interfaces couples to the spin

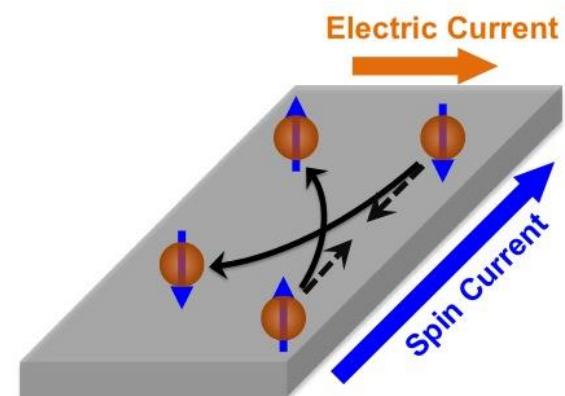
$$H = \frac{\mathbf{p}^2}{2m} - \mathbf{b} \cdot \boldsymbol{\sigma}$$

Rashba internal field: Spin

$$\mathbf{b} = \alpha \mathbf{e}_z \times \mathbf{p}$$



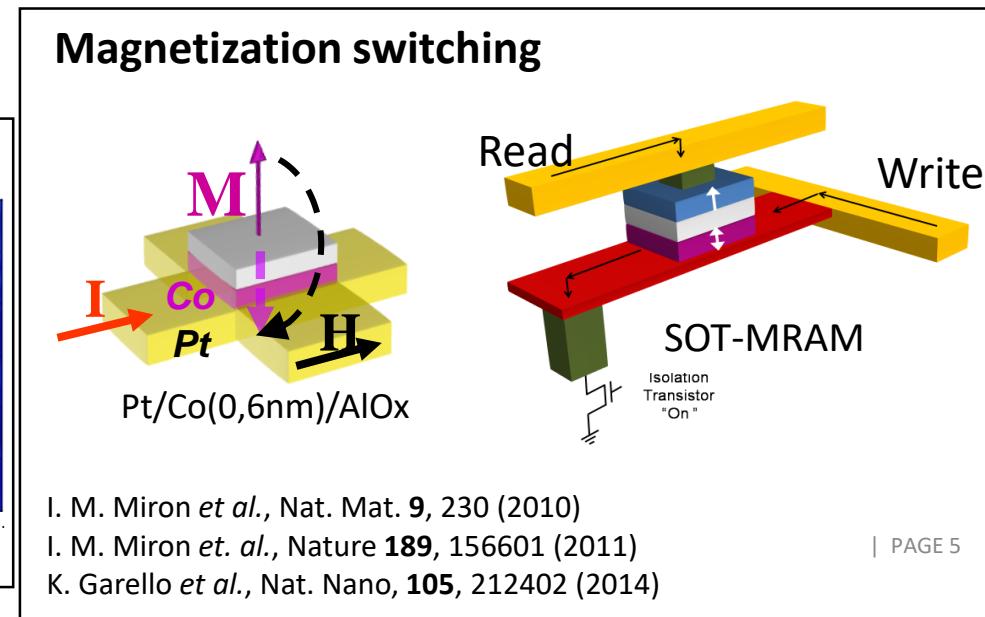
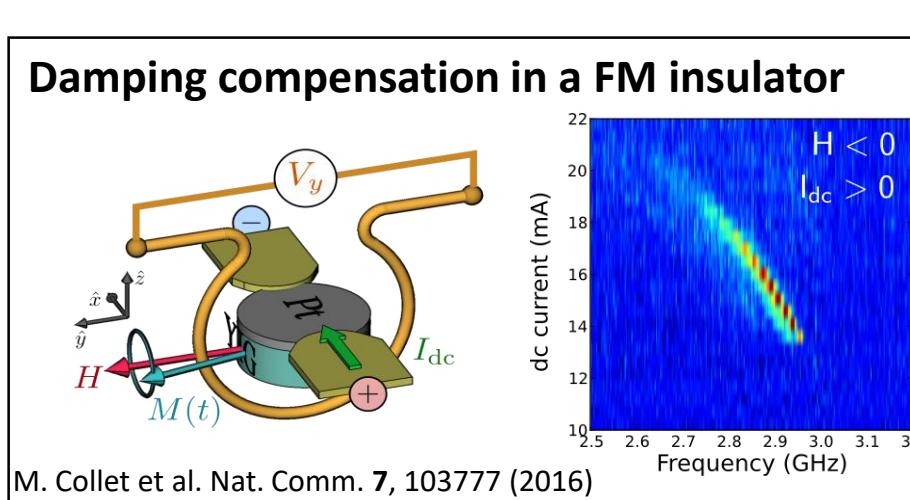
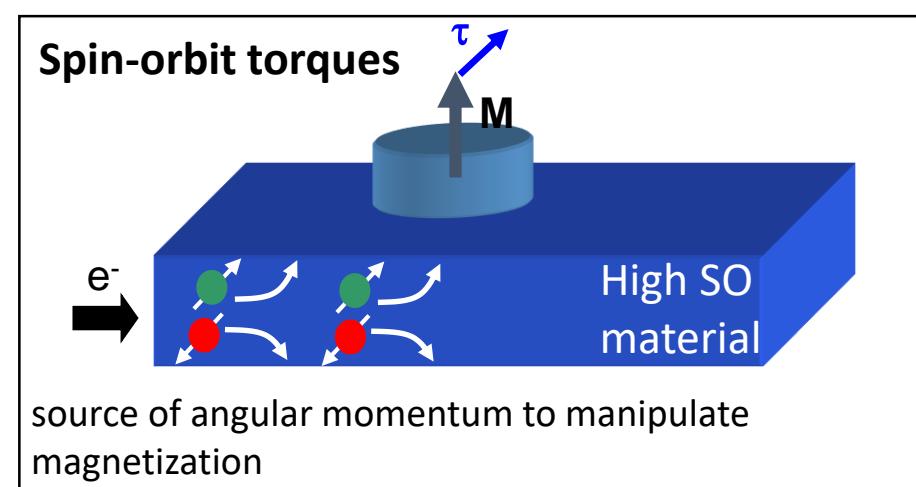
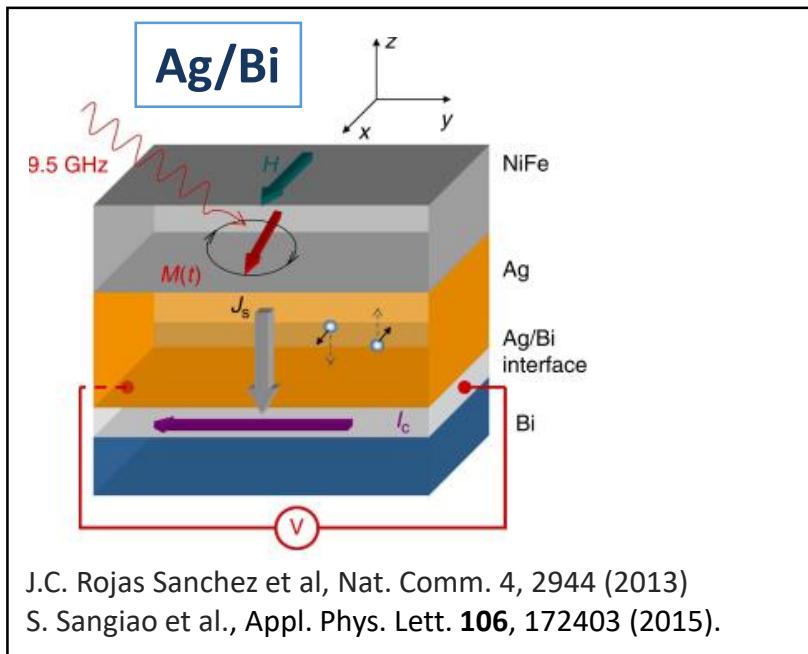
Spin Hall Effect :



→ *Spin current generation and measurement*

Spin Orbitronics

Giant spin-to-charge conversion at Rashba interfaces:



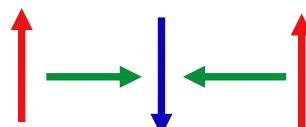
Magnetic Skyrmions

Exchange interaction



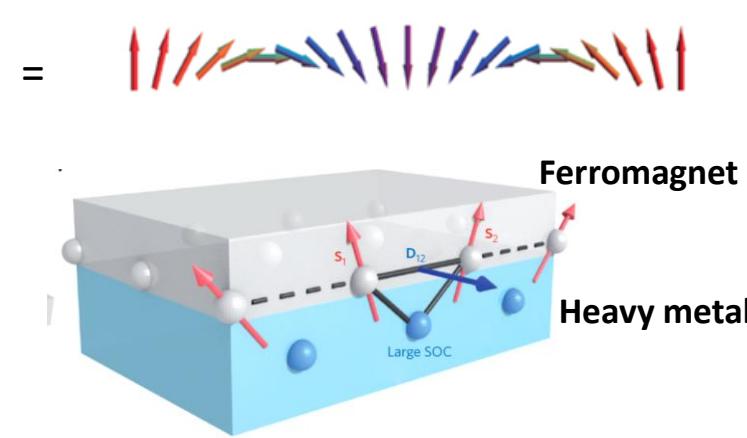
$$E_{ex} = -J \vec{S}_1 \cdot \vec{S}_2$$

Chiral interaction



$$E_{Ch} = -\vec{D} \cdot \vec{S}_1 \times \vec{S}_2$$

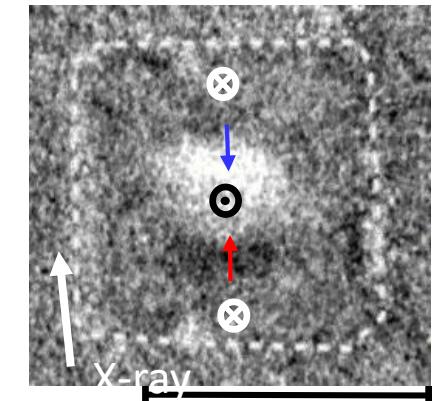
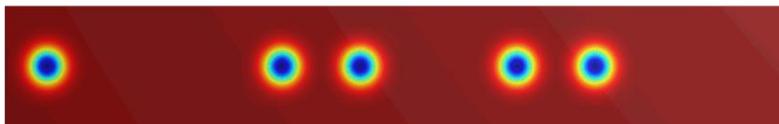
- Large spin-orbit coupling
- Structural inversion asymmetry



• Magnetic Skyrmions



- Nanometer size
- Homochiral, topologically protected
- Can be moved by small in-plane current
- → Logic and memory applications



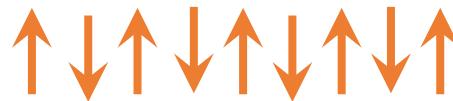
→ Room temperature skyrmions at zero field in Pt/Co/MgO nanostructures

Boule et al., Nat. Nano, 11, 449 (2016)

But: so far quite slow motion + ‘directional’ problems

Future directions: Antiferromagnets

Antiferromagnetic Spintronics

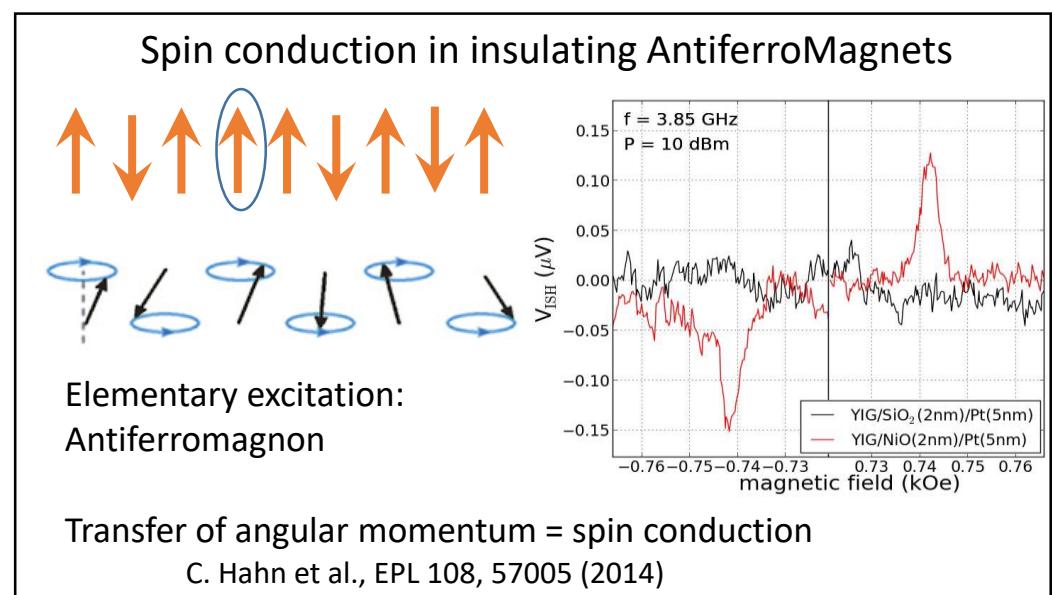


Antiferromagnetic materials could represent the future of spintronics

- immune to magnetic fields: no more unwanted data erasure
- produce no parasitic stray fields: secure ‘invisible’ data
- large magneto-transport effects: write/read with a spin current
- ultrafast TeraHertz dynamics: picosecond writing vs nanosecond today.

Addressing using:

- Spin currents
- Electrical polarization in Multiferroics

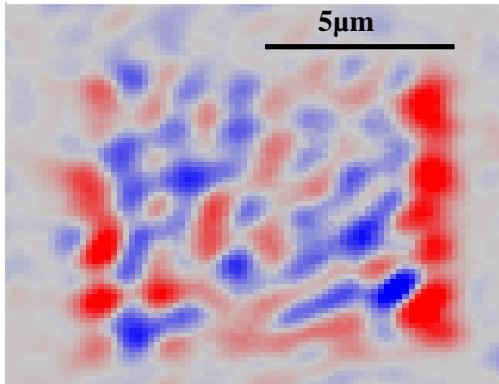
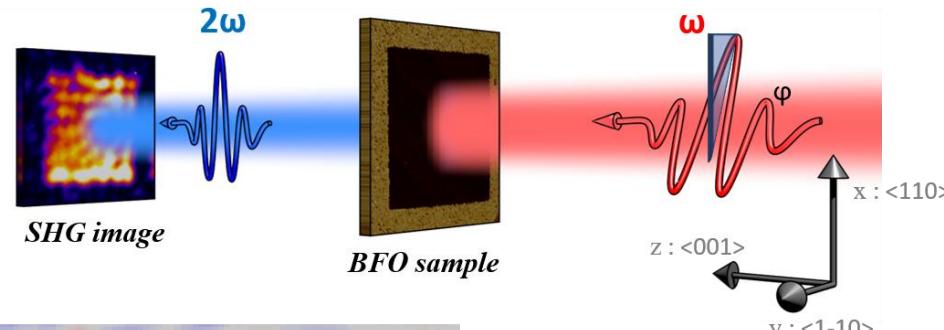


Imaging the Antiferromagnetic order

Second harmonic generation induced by :

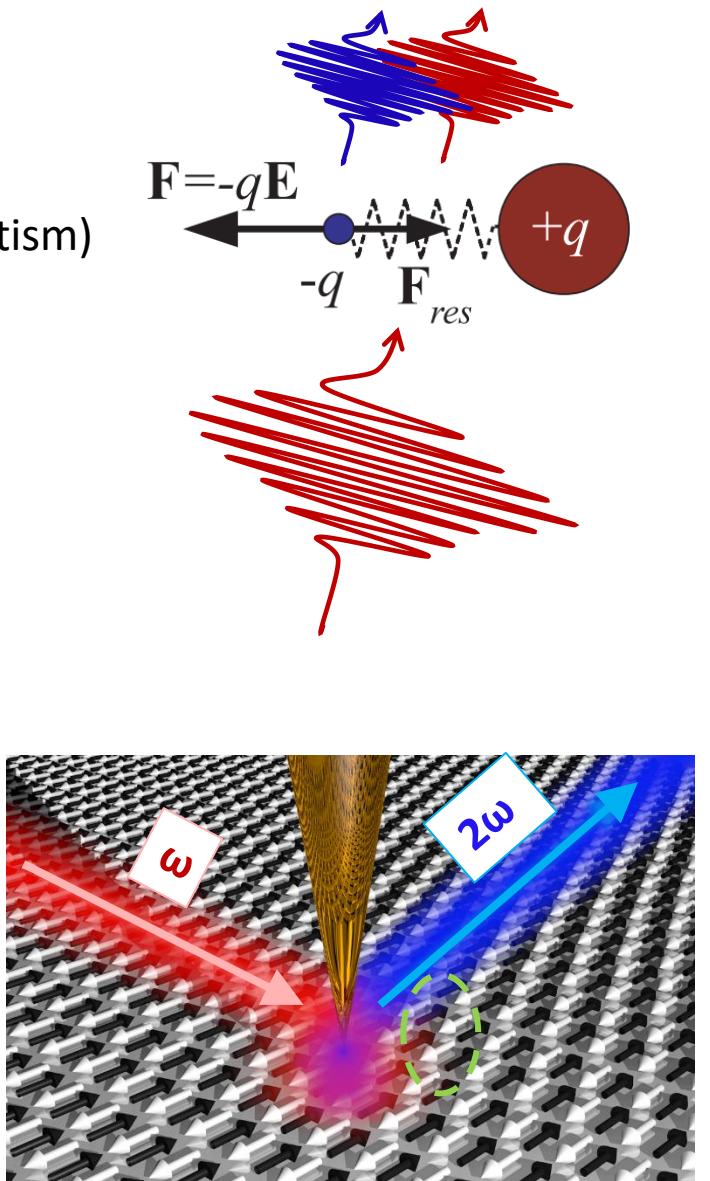
Fiebig, Pavlov & Pisarev J. Opt. Soc. Am. B **22** 96 (2005)

- Spatial centro-symmetry breaking (ferroelectricity)
- Time-inversion symmetry breaking ((anti)-ferromagnetism)



AF domains in BFO by second harmonic generation
J.-Y Chauleau et al., Nat. Mat. **16**, 803 (2017)

Project:



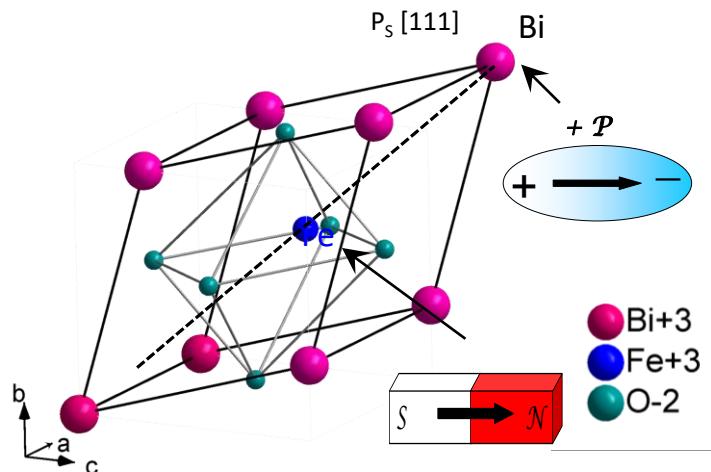
Tip-enhanced near-field light

AF multiferroics

BiFeO₃ : ferroelectric, ferroelastic and anti-ferromagnetic at 300K

Ferroelectric properties ($T_c \sim 1090$ K)

- Cubic perovskite structure → pseudo-cubic : rhombohedral distortion
- P_s due to Bi³⁺ and Fe³⁺ displacements along [111]



Kubel et al., Acta Cryst. B, 46 , 698 (1990)

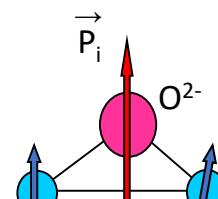
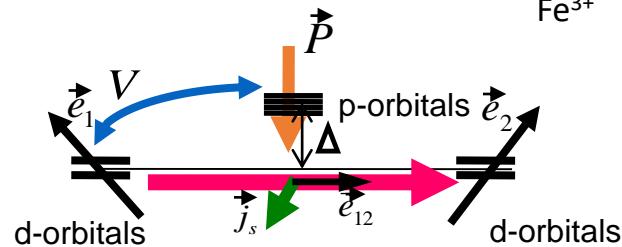
Large atomic displacements → large P_s

Magnetic properties: G type AF + magnetoelectric coupling

Magneto-electric coupling

In a solid with non-collinear magnetism (generalised Dzyaloshinski-Moriya interactions):

Typically oxydes with distorted crystallographic cells :



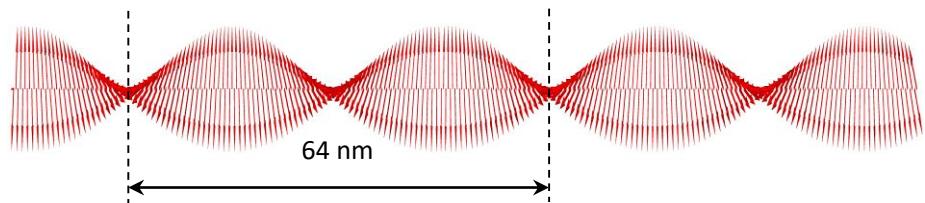
$$\vec{P} \propto \vec{e}_{12} \times \vec{j}_s$$

➤ P_S due to the magnetic structure

Katsura-Balatsky-Nagaosa PRL07

$$E_{ME} = \gamma_{ME} \sum \vec{P} \cdot (\vec{R}_{ij} \times (\vec{S}_i \times \vec{S}_j)) \quad \text{with } \gamma_{ME} \text{ the inhomogeneous ME coupling constant}$$

Leads to an AF cycloidal ordering :



AF skyrmions?

Dream: generate AF skyrmions in BiFeO_3 !

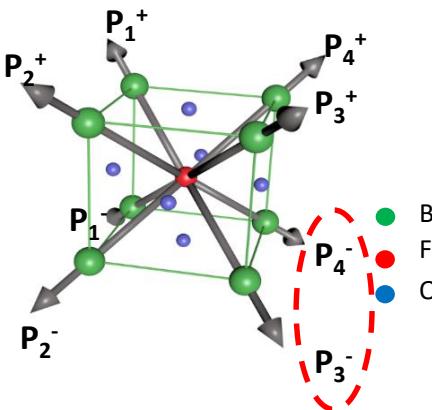
Would be fast, move in straight lines, addressable with an electric field...

But: very difficult to generate without the magnetic field 'knob'

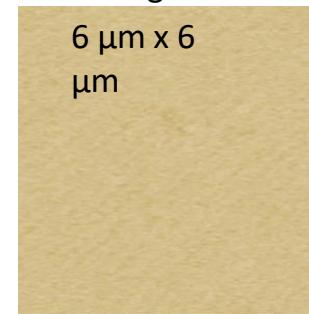
Idea: Use the frustration induced by a large density of FE domain walls

Sample : DSO/SRO/ BiFeO_3 (001) (≈ 110 nm)

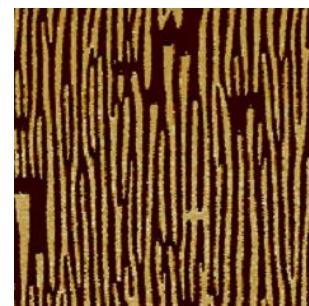
Ferroelectric configuration in thin epitaxial layers:



PFM images
6 μm x 6 μm



Out of plane P

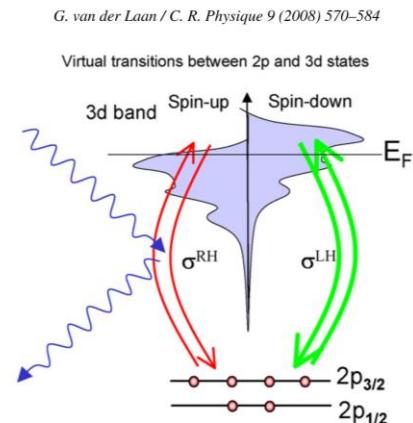
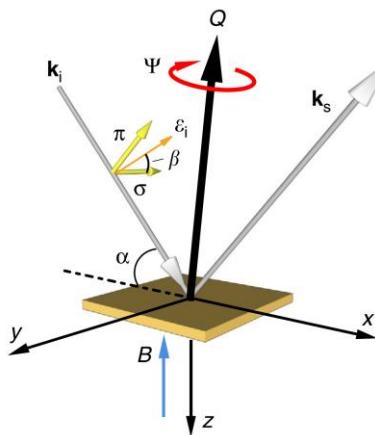


In-plane P

FE stripe domains

AF cycloids in Ferroelectric stripe pattern : XRMS

X-ray resonant magnetic diffraction:



Direct access to magnetic chirality:

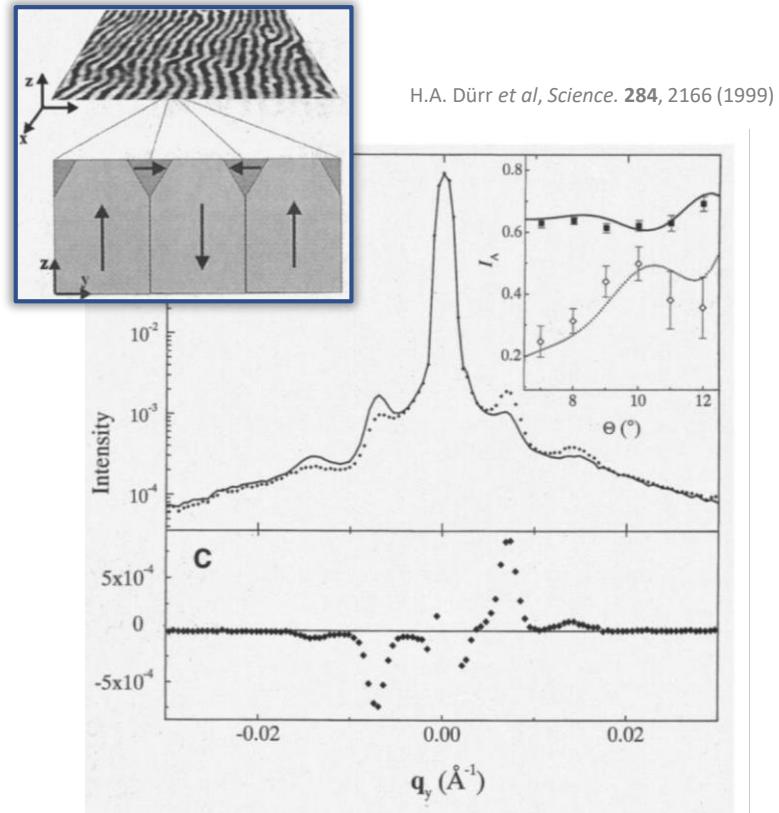


Table 1
Extensions to X-ray scattering possible with synchrotron radiation

Extensions to X-ray scattering	Allows one to probe
Tunable X-rays at resonance	Element, site and valence specificity
Polarized X-rays	Magnetic orbital and spin profile
Soft X-rays	Nanoscale sensitivity (down to 1 nm)
Coherent radiation	Local configuration
Pulsed radiation	Dynamics

AF cycloids in Ferroelectric stripe pattern : XRMS



XRMS + Open Port

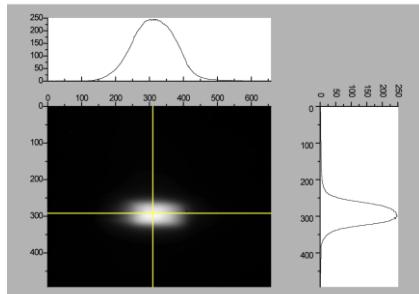
h-slit = 40µm / v-slit 10µm

Calculated @2.2m from M8

18µm * 7µm FMHM

Measured @ 2.2m from M8

17.3µm * 7.4µm

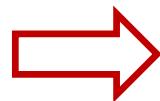


- **Energy range : 50-1700 eV** (optimized 70-1000 eV)
- **High flux : >5 10¹² ph/s** on the sample,
- **Resolving power E/ΔE ≥ 10000**

Variable polarization: linear and circular:

2 Apple-II undulators HU80 + HU44

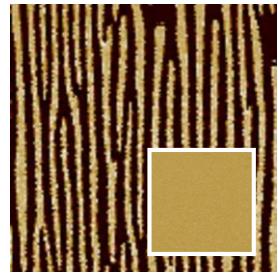
=> the whole energy range in first harmonic



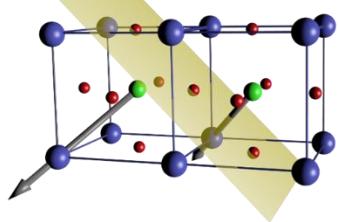
**Resonant elastic scattering
Resonant inelastic scattering
Holography**

XRMS in reflectivity

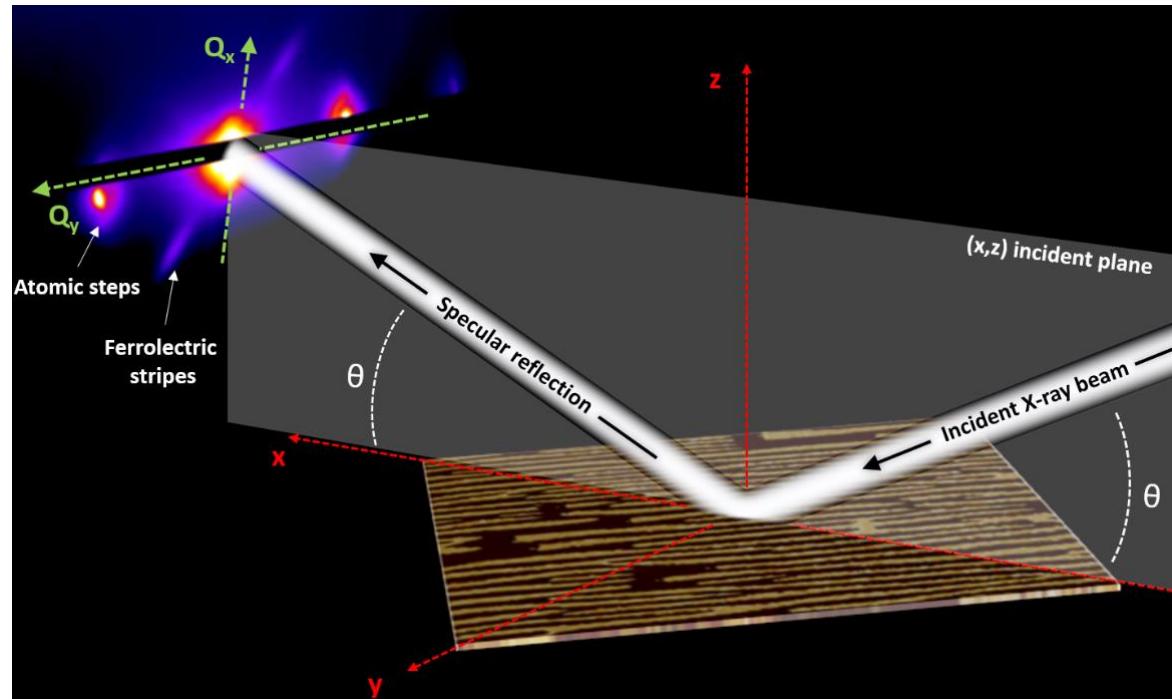
$\text{BiFeO}_3 / \text{SrRuO}_3 // \text{DyScO}_3$
71° walls



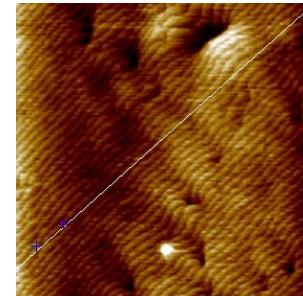
$4 \times 4 \mu\text{m}^2$ in-plane PFM
+ out of plane in inset



XRMS experiment in reflectivity configuration at O K edge

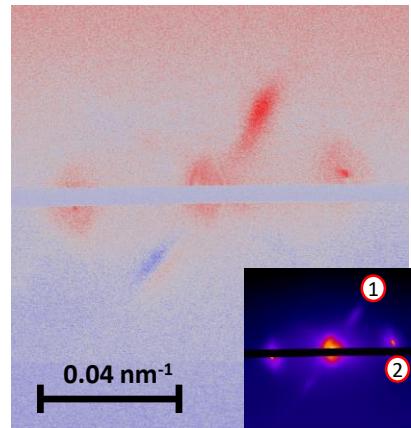


AFM image:

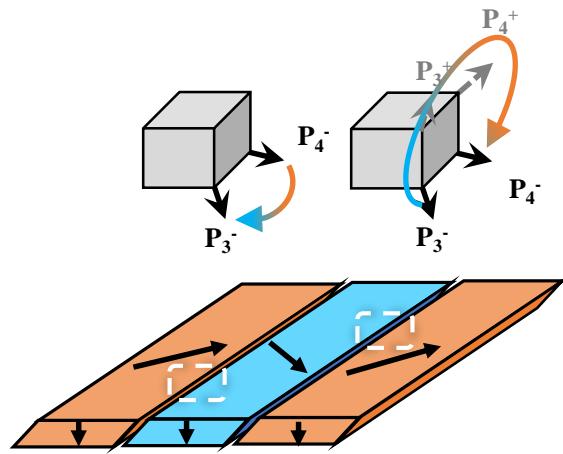


Chiral electrical polarization

Dichroism at the Oxygen K edge !



Chiral FE structures

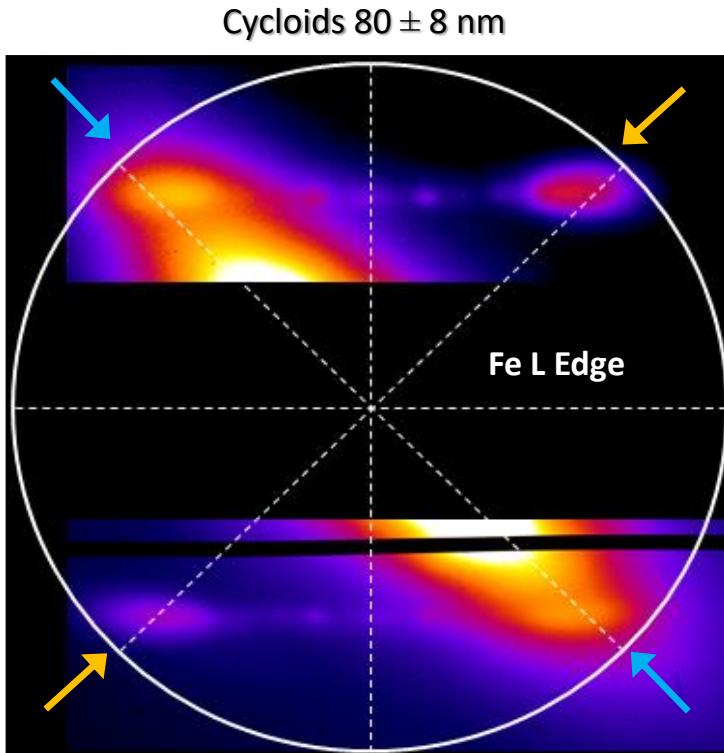


Very surprising because this is not the lowest energy state...

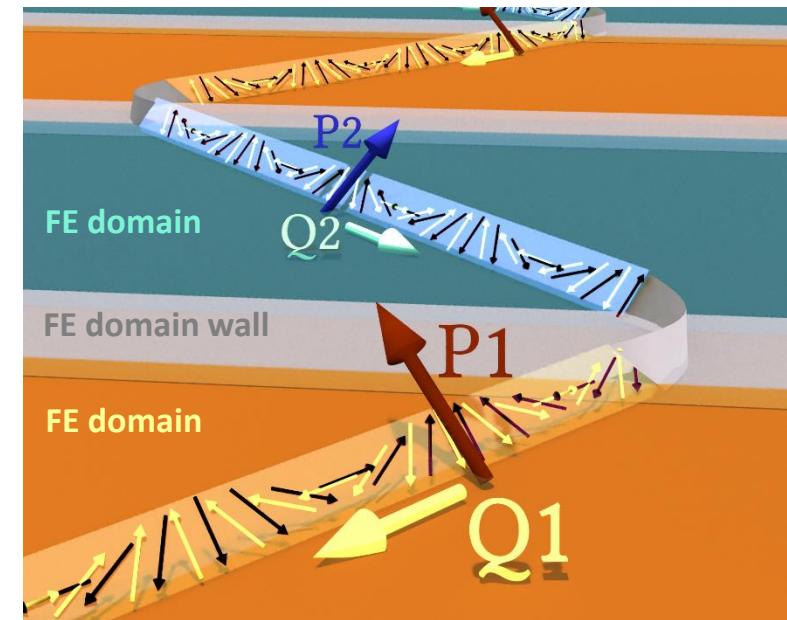
AF cycloids

What about the magnetic textures?

→ Fe L edge



scattering amplitude:



$$f_n^{\text{res}} = f_0(\hat{\epsilon} \cdot \hat{\epsilon}') - i f_1(\hat{\epsilon} \times \hat{\epsilon}') \cdot \mathbf{m}_n + f_2(\hat{\epsilon}' \cdot \mathbf{m}_n)(\hat{\epsilon} \cdot \mathbf{m}_n)$$

$\hat{\epsilon}$ and $\hat{\epsilon}'$: polarization states of incident and diffracted beams

\mathbf{m}_n : local magnetization vector

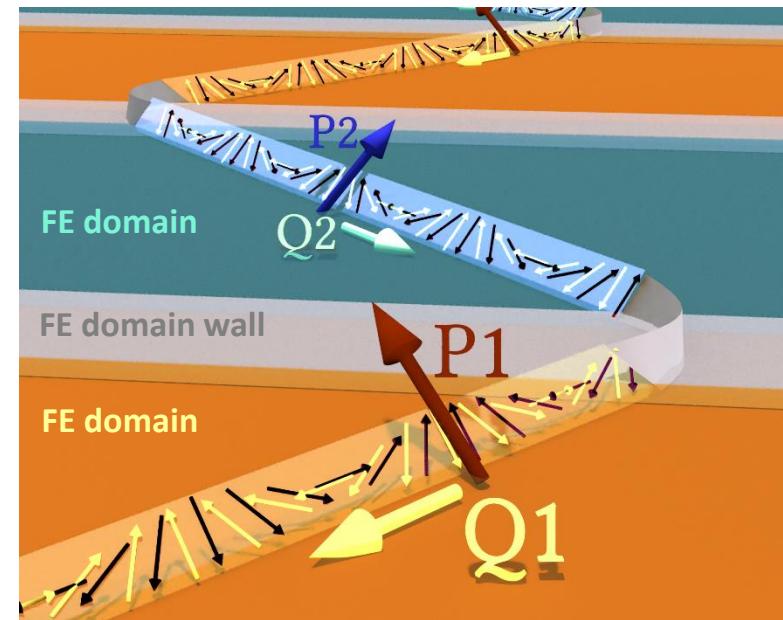
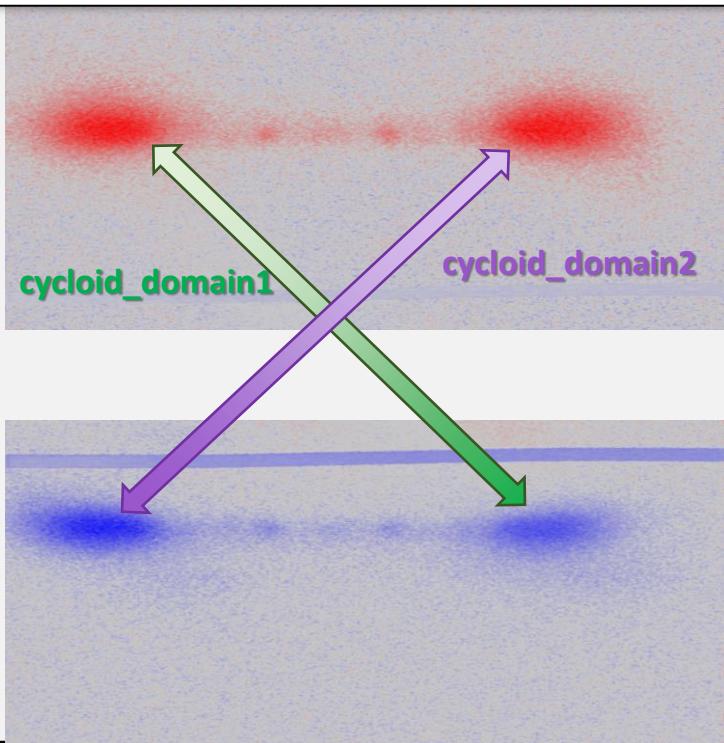
f_0, f_1, f_2 resonance factors for, the monopole, magnetic dipole and quadrupole parts of the scattering amplitude.

AF cycloids

What about the magnetic textures?

→ Fe L edge

Normalized circular dichroism : chiral structures

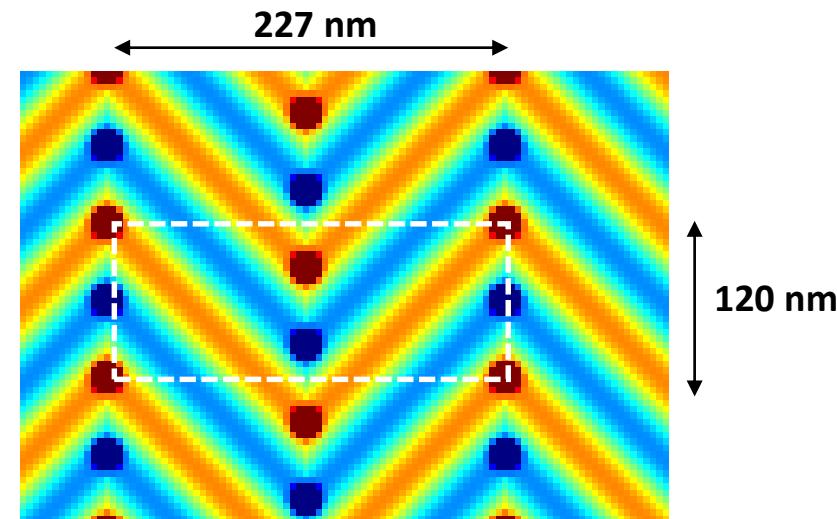
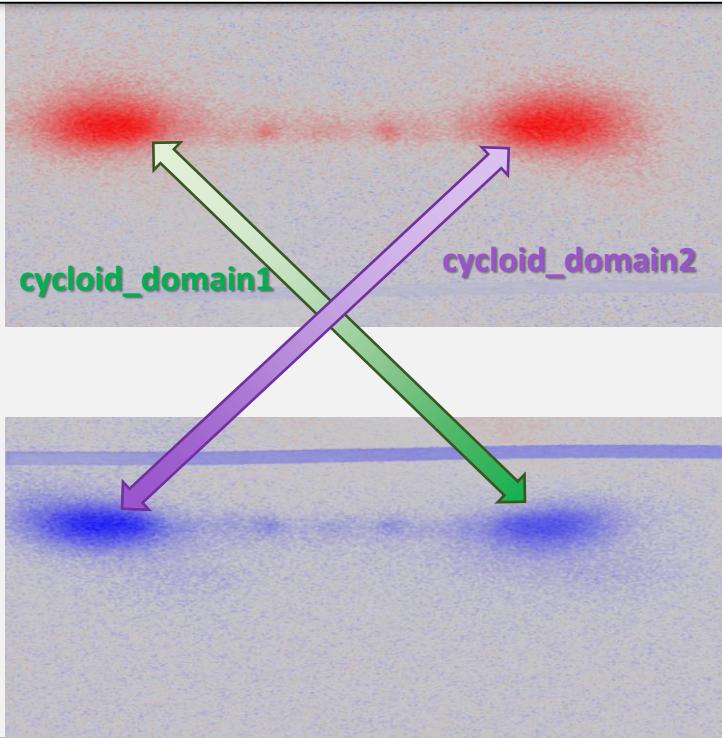


AF cycloids

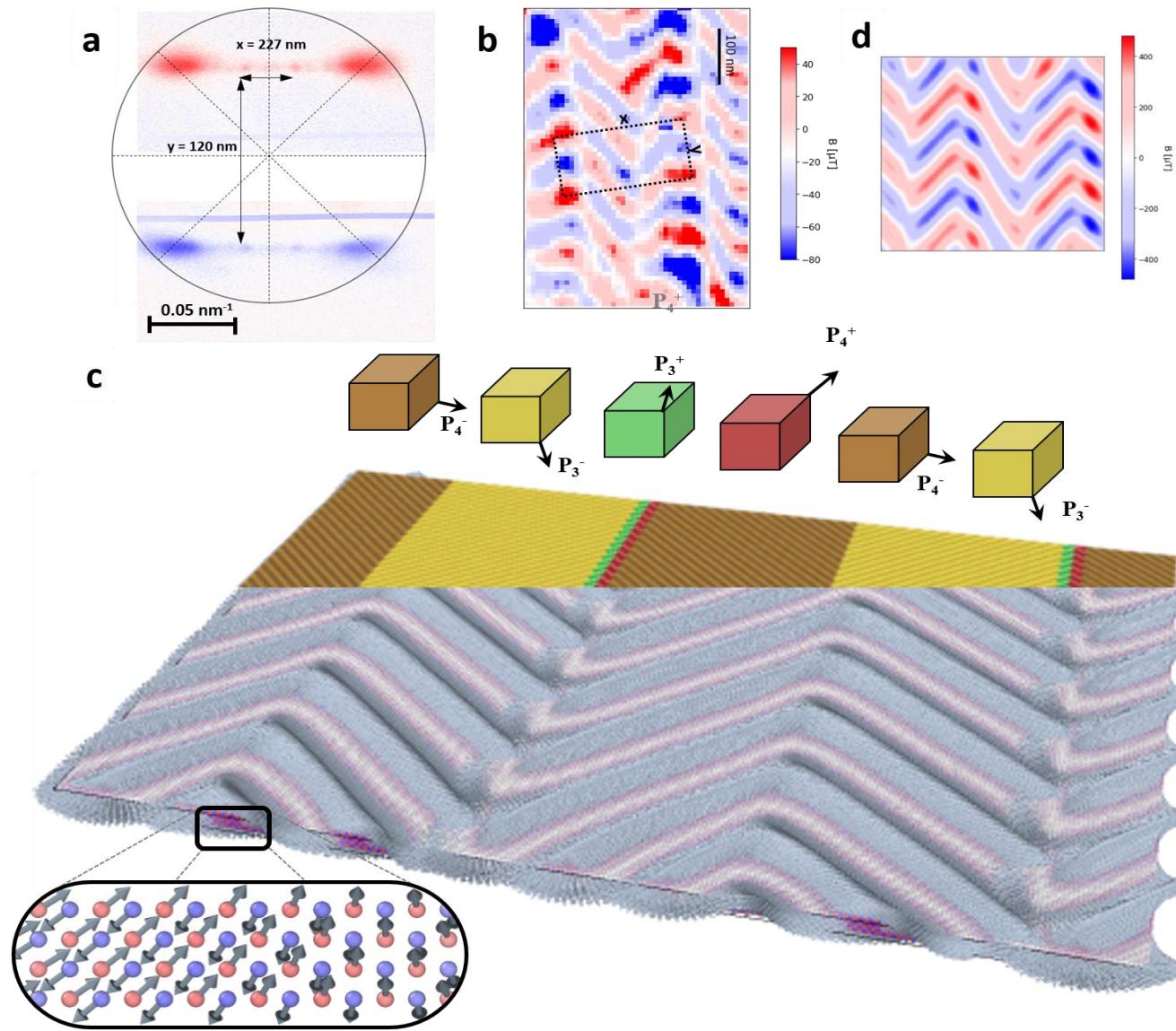
What about the magnetic textures?

→ Fe L edge

Normalized circular dichroism : chiral structures



Magnetic simulations



Conclusions / Perspectives

Present issues in *information technologies*:

Dissipation → replace charge by spin

Size → reduce stray fields (Antiferromagnets),
use topological protection (skyrmions)

Speed → Reduce friction and size (skyrmions),
Change magnetic interaction:
Antiferromagnets (THz)