

Strangeness measurements up to Tevatron and predictions for LHC

Ingrid Kraus



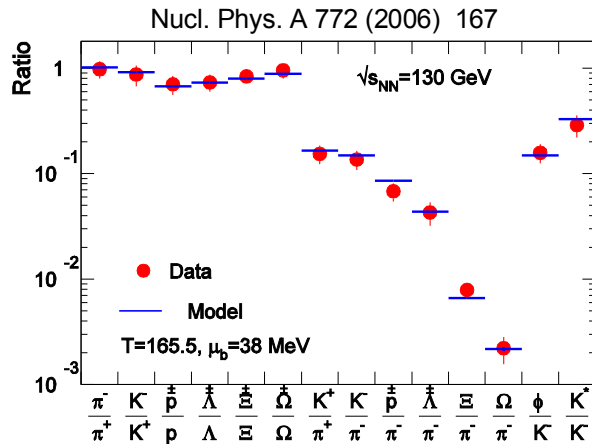
A Large Ion Collider Experiment

European Organisation for Nuclear Research

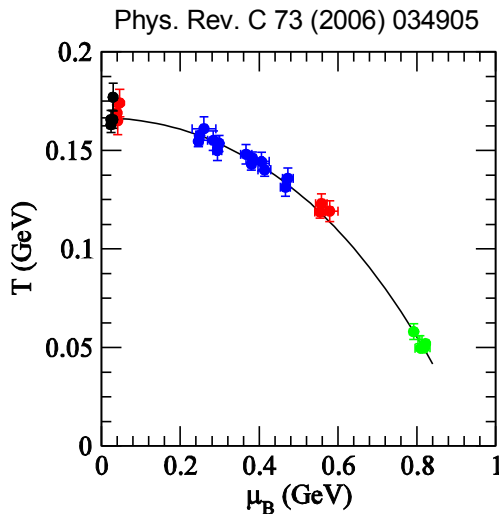


heavy-ion collisions

statistical model: grand-canonical ensemble



→ 2 parameters
 T, μ_B

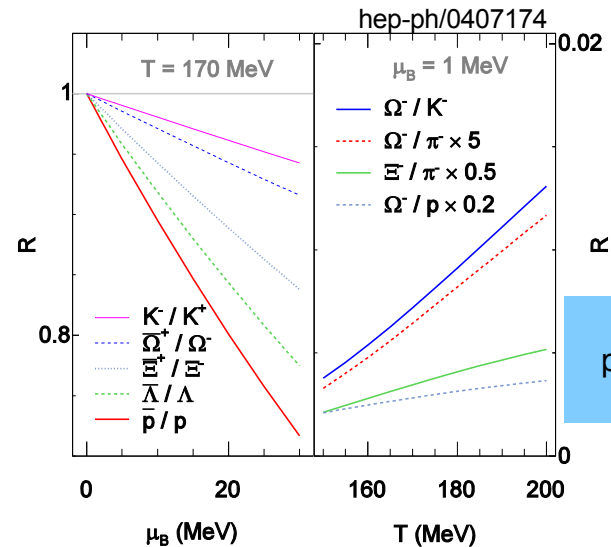


freeze-out curve
→ extrapolate to
LHC

Phys. Rev. C 74 (2006) 034903

\bar{h}/h Ratio		mixed Ratio	
π^+/π^-	$0.9998^{+0.0002}_{-0.0010}$	K^+/π^+	$0.180^{+0.001}_{-0.001}$
K^+/K^-	$1.002^{+0.008}_{-0.002}$	K^-/π^-	$0.179^{+0.001}_{-0.001}$
\bar{p}/p	$0.989^{+0.011}_{-0.045}$	p/π^-	$0.091^{+0.009}_{-0.007}$
$\bar{\Lambda}/\Lambda$	$0.992^{+0.009}_{-0.036}$	Λ/p	$0.473^{+0.004}_{-0.006}$
$\bar{\Xi}^+/\Xi^-$	$0.994^{+0.006}_{-0.026}$	Ξ^-/Λ	$0.160^{+0.002}_{-0.003}$
$\bar{\Omega}^+/\Omega^-$	$0.997^{+0.003}_{-0.015}$	Ω^-/Ξ^-	$0.186^{+0.008}_{-0.009}$

→ calculate predictions for LHC



→ extract
parameters
from data

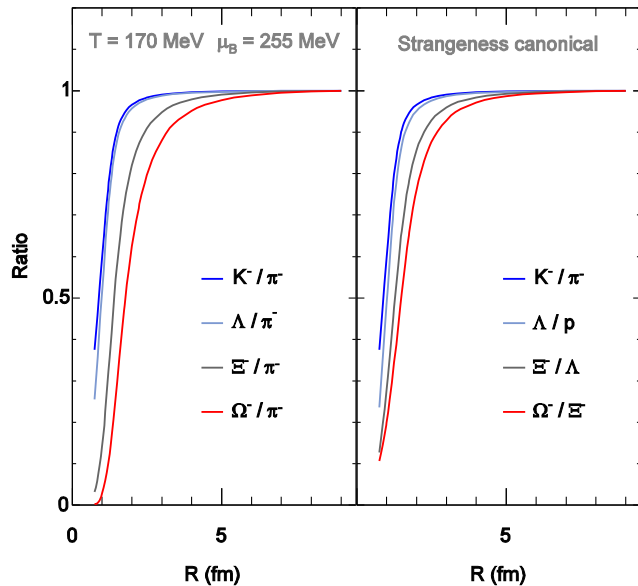
small colliding systems

canonical ensemble

small systems / peripheral collisions,
low energies

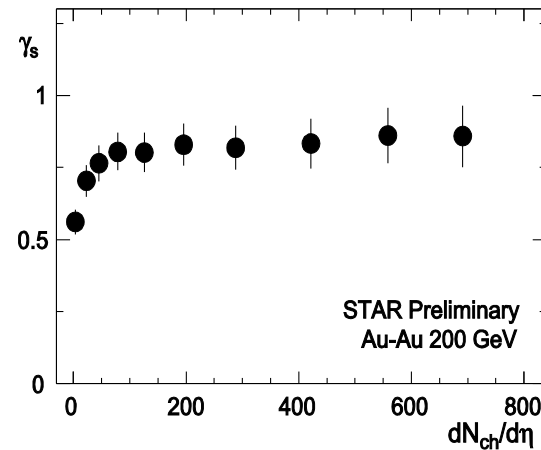
suppressed phase-space for particles
related to conserved charge

$$n_i^{\text{canonical}} \approx n_i^{\text{grand-canonical}} \frac{I_S(x)}{I_0(x)}$$



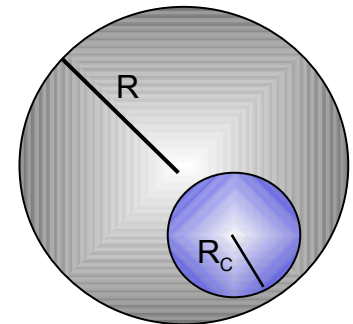
suppression beyond canonical expectation

J. Phys. G 31(2005) S101



deviations:
strangeness
understauration
factor γ_s

alternative:
small cluster (R_C)
in fireball (R):
 $R_C \leq R$
→ **chemical
equilibrium**
in subvolumes

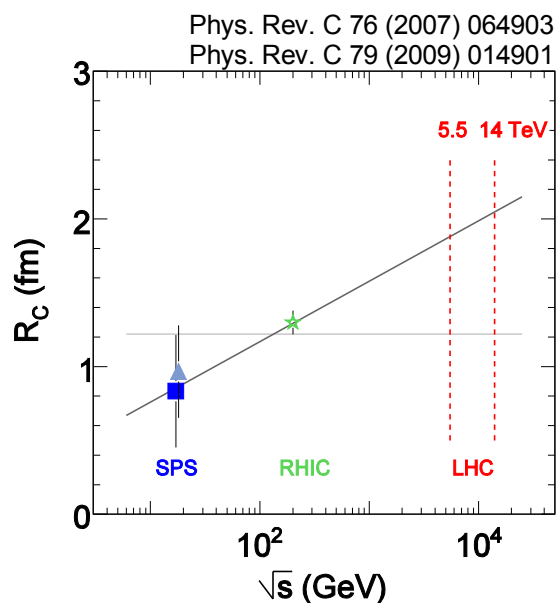


predictions for pp interactions at LHC

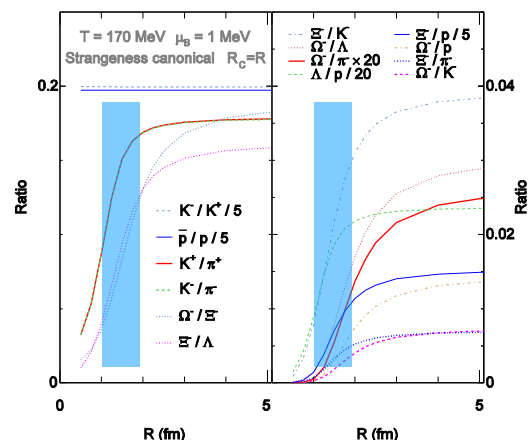
extrapolation of cluster size

what defines R_C in p+p?

- initial size of p+p system: R_C const
- final state of large number of produced hadrons:
→ increase with \sqrt{s} and multiplicity

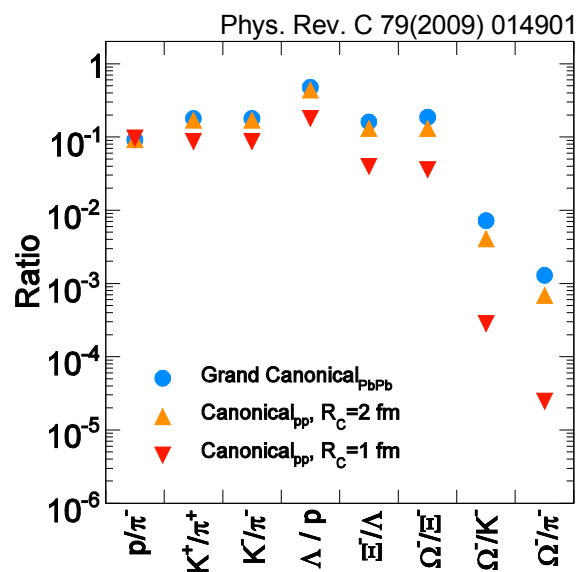


measurement of cluster size



significant increase of ratios at $R_C \approx 1.5$ fm

→ determine R_C from data



sensitivity increases with strangeness difference

→ R_C from Ω / π

what do we know about pp ?

soft physics

- multiplicity distributions
- spectra and mean-pt
- strangeness production

event characterisation / scaling behaviour

- c.m. energy \sqrt{s}
- multiplicity
- hard vs soft events

	<i>ISR</i>	<i>SppS</i>	<i>Tevatron</i>
\sqrt{s} (GeV)	~ 10	~ 100	~ 1000

multiplicity distribution

ISR Phys Rev D 30 (1984) 528 $\sqrt{s} = 30 - 62 \text{ GeV}$

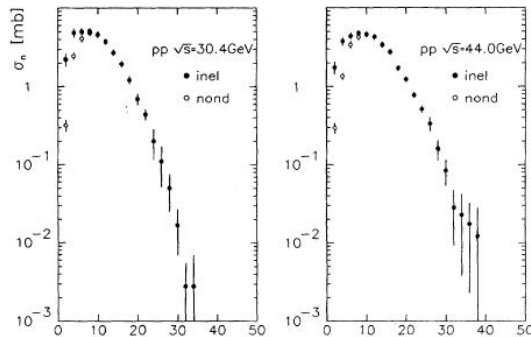
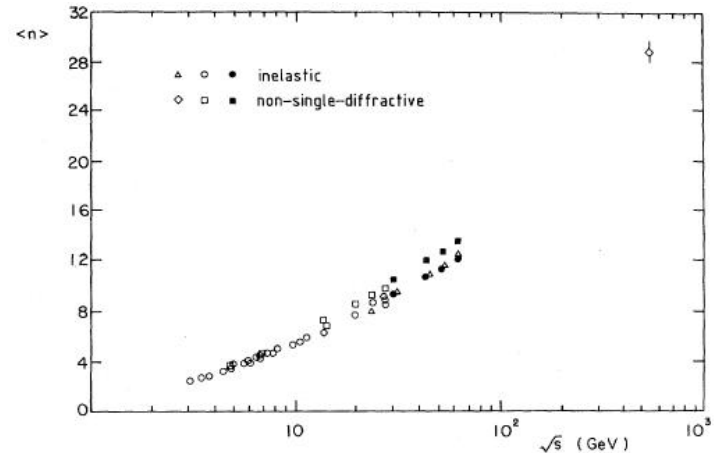
4 π detector

p+p

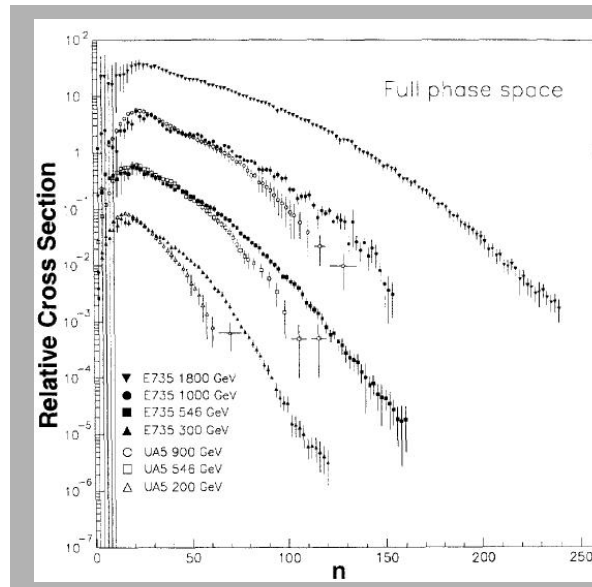
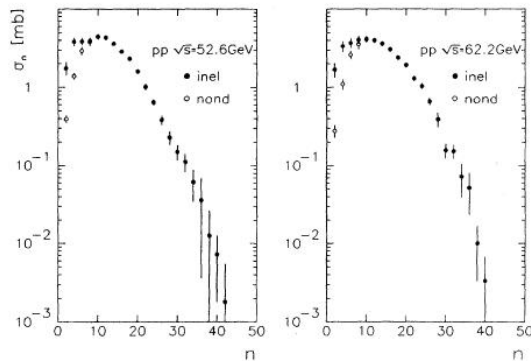
energy dependence of $\langle n \rangle$

Feynman scaling and *limiting fragmentation* predict $\langle n \rangle \sim \ln s$
in data $\langle n \rangle$ rises faster

$$\rightarrow \langle n \rangle = A + B \ln s + C \ln^2 s$$



non single
diffractive
events:
both p
destroyed



E735

Phys Lett B 435 (1998) 453

$\sqrt{s} = 300 - 1800 \text{ GeV}$

UA5

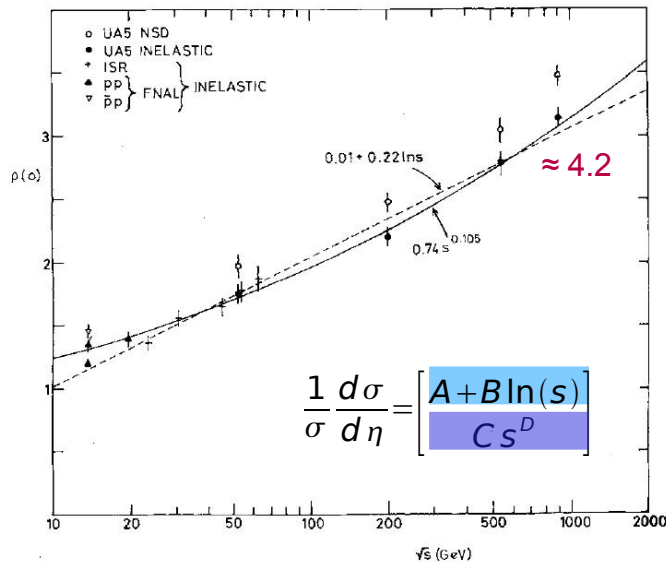
data

$\sqrt{s} = 200 - 900 \text{ GeV}$

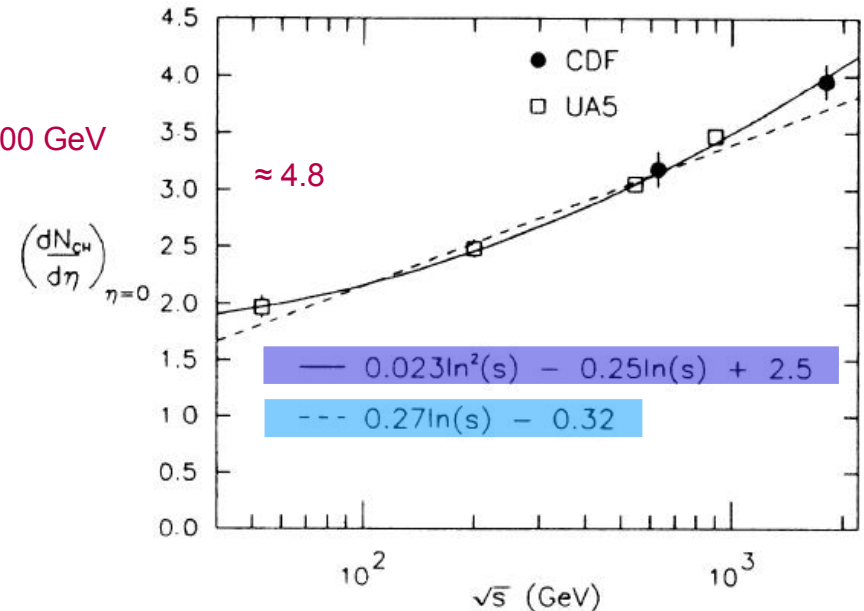
multiplicity density vs \sqrt{s}

UA5 Z Phys C 33 (1986) 1 $\sqrt{s} = 200, 900 \text{ GeV}$

CDF Phys Rev D 41 (1990) 2330 $\sqrt{s} = 630, 1800 \text{ GeV}$
 $p_t > 50 \text{ MeV}/c$



at $\sqrt{s} = 14,000 \text{ GeV}$



multi-particle production in

- statistical hydrodynamical models
determined by initial energy density
- parton-parton interactions with string fragmentation models
described by multiple parton-parton scattering
or number of strings

charged particle midrapidity density

- scales ~ with $\ln(s)$ at lower $\sqrt{s} \rightarrow$ available energy
- increases at SppS and Tevatron faster than $\ln s$

KNO scaling

Koba, Nielsen, Olesen Nucl Phys B 40 (1971) 317

scaling of multiplicity distribution in high energy hadron collisions

"the normalised multiplicity distribution keeps its form independently of the beam energy and just scales up as $\ln s$ "

$$\begin{array}{c} \text{multiplicity} \\ \text{distribution} \end{array} P_n(s) = \frac{\begin{array}{c} \text{cross} \\ \text{section for} \\ \text{the} \\ \text{multiplicity} \\ \text{being } n \text{ at} \\ \text{the c.m.} \\ \text{energy } \sqrt{s} \end{array} \sigma_n(s)}{\sigma_{tot}(s)} = \frac{1}{\langle n \rangle} \begin{array}{c} \text{independent of } \sqrt{s} \text{ except} \\ \text{through } n / \langle n \rangle \end{array} \psi(n/\langle n \rangle)$$

$\langle n \rangle$ average multiplicity at c.m. energy \sqrt{s}

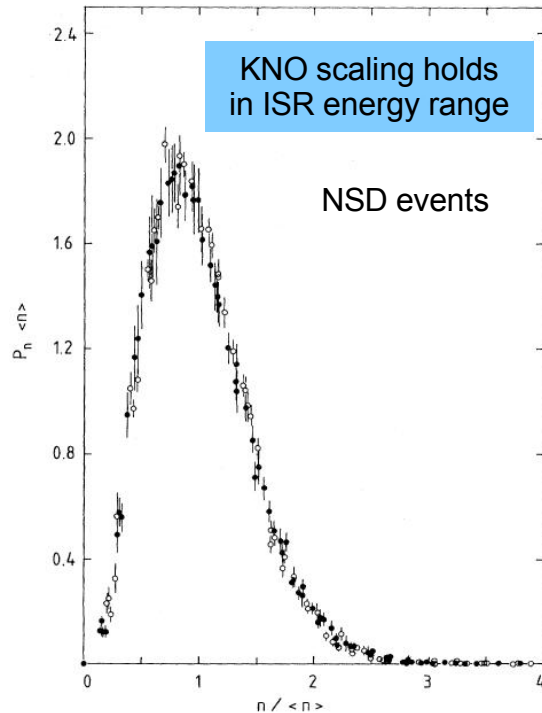
$\langle n \rangle P_n(s)$ is only a function of $n / \langle n \rangle$

multiplicity distribution

ISR Phys Rev D 30 (1984) 528 $\sqrt{s} = 30 - 62 \text{ GeV}$

4 π detector

p+p

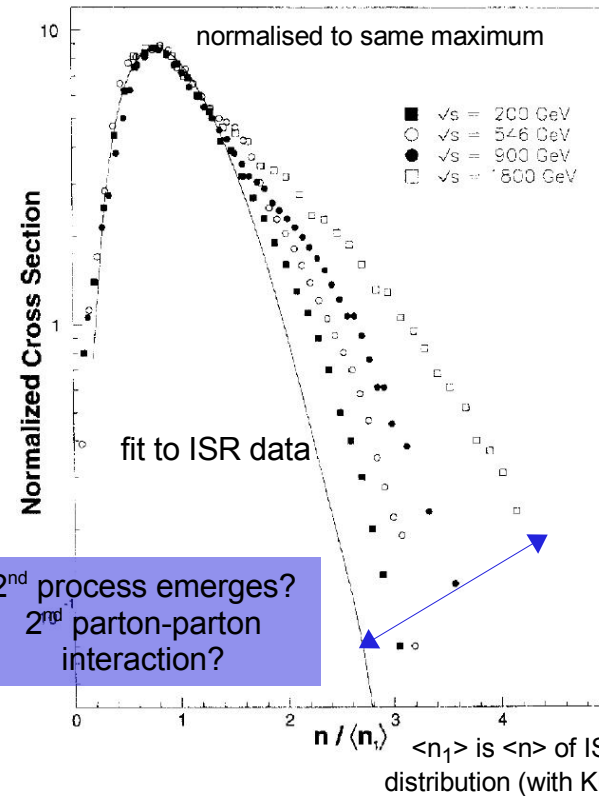


UA5 Z Phys C 43 (1989) 357 $\sqrt{s} = 200 - 900 \text{ GeV}$

KNO scaling breaks at full phase space

negative binomial distribution describes data
(except 900 GeV full phase space)

E735 Phys Lett B 435 (1998) 453 $\sqrt{s} = 300 - 1800 \text{ GeV}$



1st component
KNO-like scaling with
const cross section
(32mb)

2nd component
cross section
increases with \sqrt{s}
(18mb at 1800 GeV)

CDF Phys Rev D 56 (1997) 3811 $\sqrt{s} = 1800 \text{ GeV}$

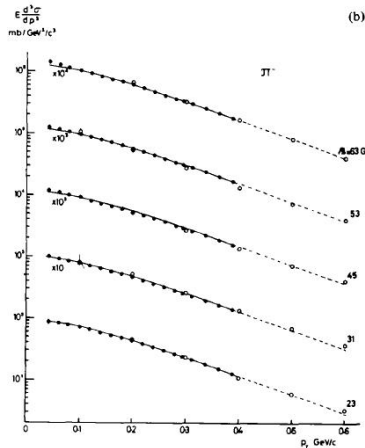
multiple jet events:

kinematic independent double parton interactions
cross section $\sim 15 \text{ mb}$

shape of pt spectra

ISR Phys Lett 64B (1976) 111 $\sqrt{s} = 23 - 63 \text{ GeV}$

π^- : $40 < pt < 400 \text{ MeV/c}$ **p+p**



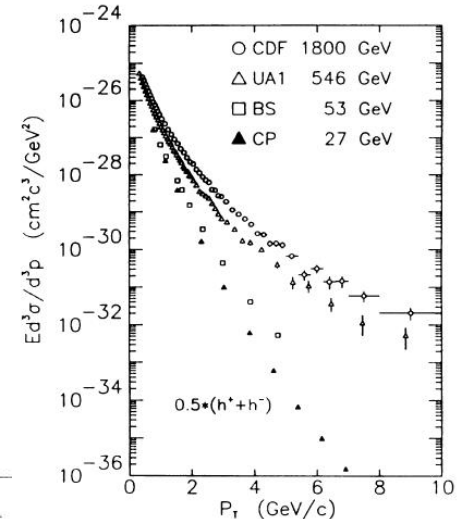
$\sqrt{s} = 63 \text{ GeV}$

53
45
31
23

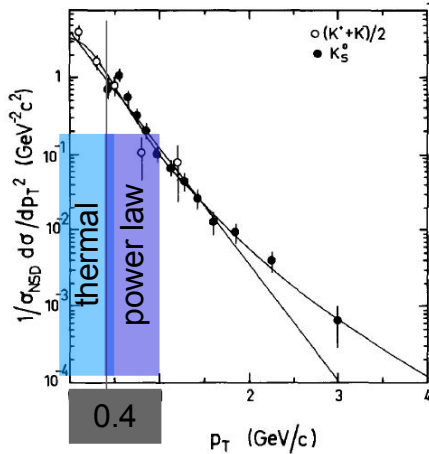
turnover at low pt
exp(pt) fails
exp(mt) better
→ **shape dictated by transverse energy**

CDF Phys rev Lett 61 (1988) 1819 $\sqrt{s} = 630, 1800 \text{ GeV}$

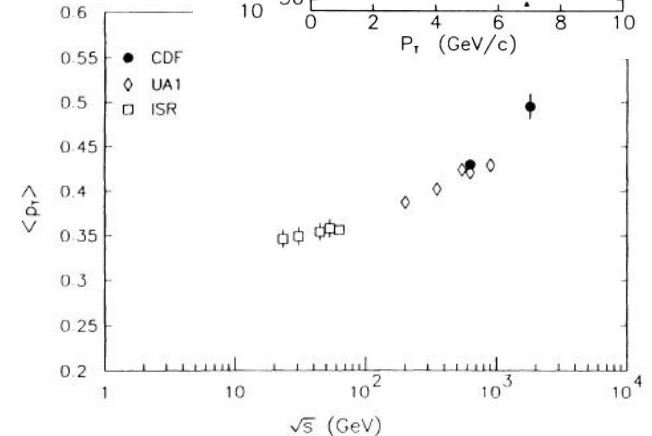
$pt > 0.4 \text{ GeV/c}$
power law in measured range



UA5 Nucl Phys B 258 (1985) 505 $\sqrt{s} = 540 \text{ GeV}$



$$E \frac{d^3 \sigma}{d^3 p} = \left[\frac{A \exp(-b m_t)}{C \frac{p_0^n}{(p_0 + p_t)^n}} \right]$$



mean-pt vs \sqrt{s} and multiplicity

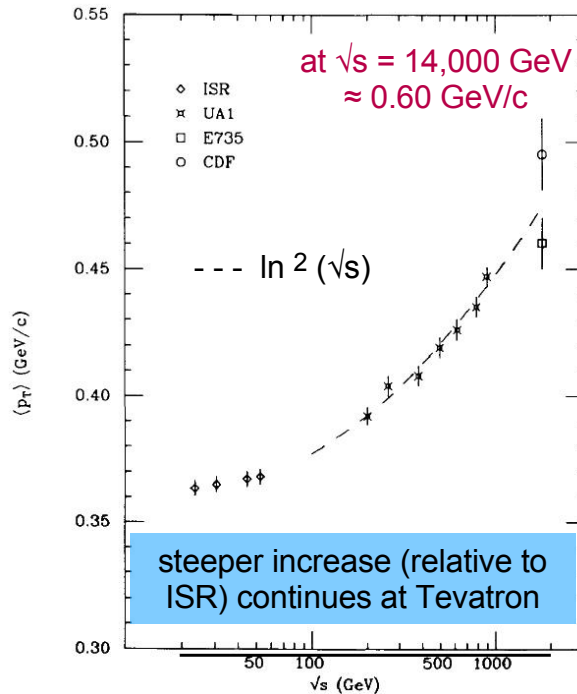
UA1 Nucl Phys B 335 (1990) 261 $\sqrt{s} = 0.2 - 0.9 \text{ TeV}$

magnetic and calorimetric analysis

$p_t > 0.25 \text{ GeV/c}$

pt spectra extrapolated with power law

thermal distribution results in 6% higher $\langle p_t \rangle$



CDF 0904.1098 [hep-ex] $\sqrt{s} = 1.96 \text{ TeV}$

$p_t > 0.4 \text{ GeV/c}$

- power law @ $p_t < 10 \text{ GeV/c}$

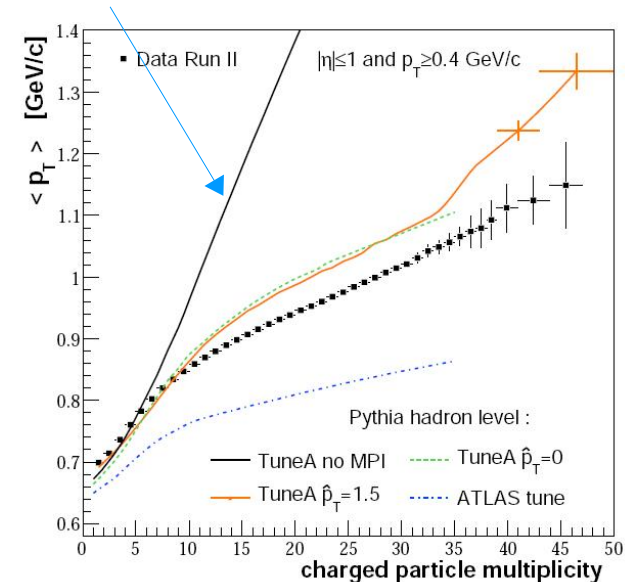
- Pythia tune A @ $p_t < 20 \text{ GeV/c}$

- "more sophisticated parametrisation" above

→ *min bias collisions are mixture of hard and soft processes*

$$f = A \left(\frac{p_0}{p_t + p_0} \right)^n + B \left(\frac{1}{p_t} \right)^s$$

2 ↔ 2 interactions
 hard collisions ↔ large Nch
 → rises too fast



underlying event

CDF Phys. Rev. D 65 (2002) 092002 $\sqrt{s} = 1.8 \text{ TeV}$

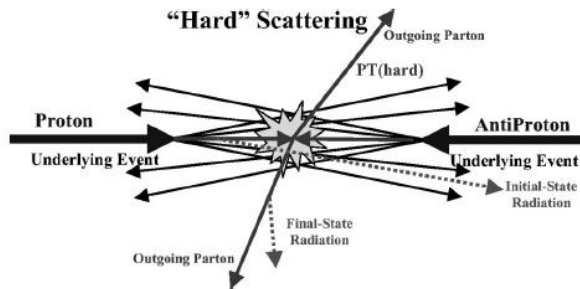
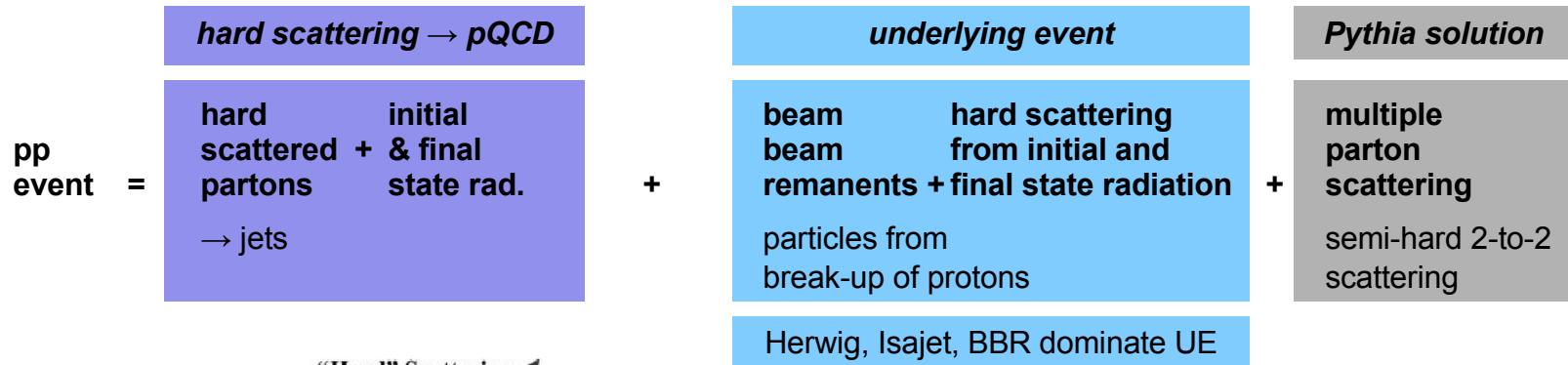


FIG. 1. Illustration of the way the QCD Monte Carlo models simulate a proton-antiproton collision in which a hard 2-to-2 parton scattering with transverse momentum, $p_T(\text{hard})$, has occurred. The resulting event contains particles that originate from the two outgoing partons (plus initial and final-state radiation) and particles that come from the breakup of the proton and antiproton ("beam-beam remnants"). The "hard scattering" component consists of the outgoing two "jets" plus initial and final-state radiation. The "underlying event" is everything except the two outgoing hard scattered "jets" and consists of the "beam-beam remnants" plus possible contributions from the "hard scattering" arising from initial and final-state radiation.

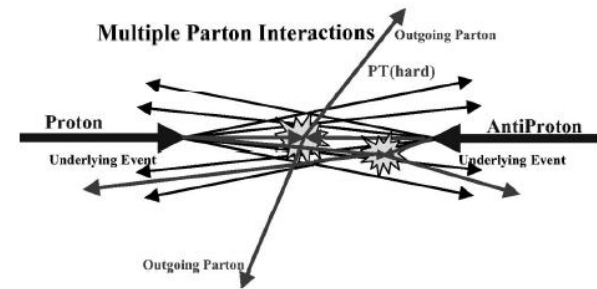
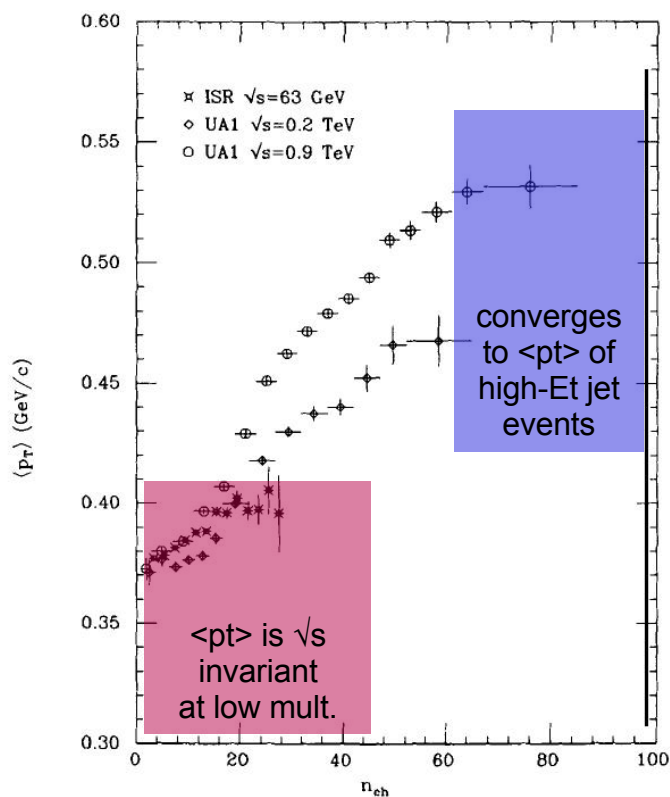


FIG. 2. Illustration of the way PYTHIA models the "underlying event" in proton-antiproton collision by including multiple parton interactions. In addition to the hard 2-to-2 parton-parton scattering with transverse momentum, $p_T(\text{hard})$, there is a second "semi-hard" 2-to-2 parton-parton scattering that contributes particles to the "underlying event."

mean-pt vs multiplicity

UA1 Nucl Phys B 335 (1990) 261 $\sqrt{s} = 0.2 - 0.9 \text{ TeV}$

$p_T > 0.25 \text{ GeV/c}$

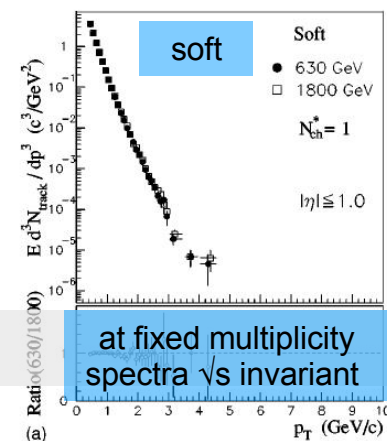


CDF Phys. Rev. D 65 (2002) 072005 $\sqrt{s} = 1800, 630 \text{ GeV}$

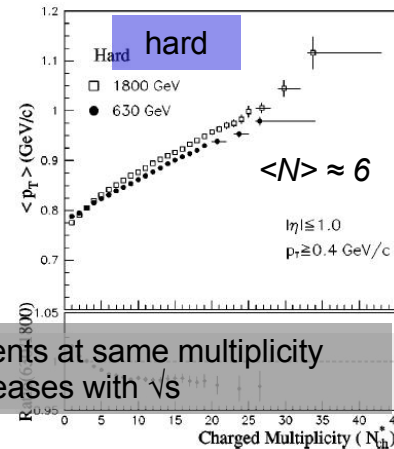
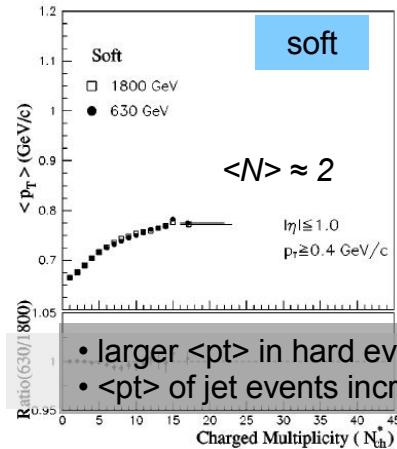
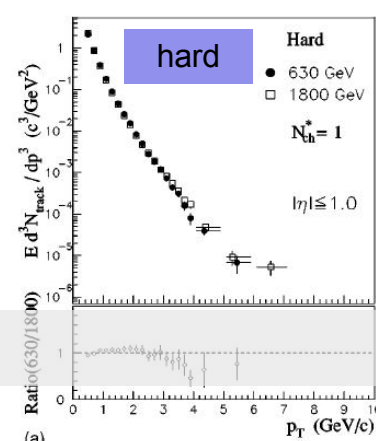
hard events: E_t or $p_T > 1.1 \text{ GeV}$ in jet cone of $R = 0.7$

soft events: no high-pt jet

pt distribution



$p_T > 0.4 \text{ GeV/c}$



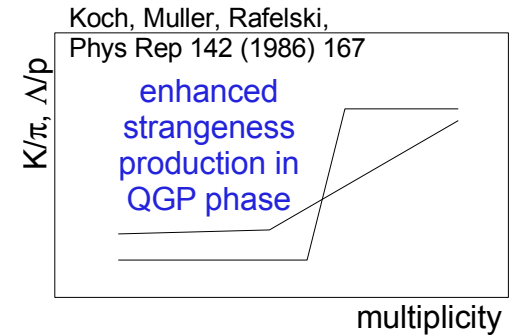
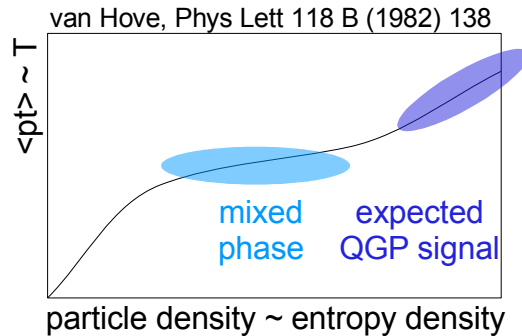
- larger $\langle p_T \rangle$ in hard events at same multiplicity
- $\langle p_T \rangle$ of jet events increases with \sqrt{s}

QPG search in $\bar{p}p$ at $\sqrt{s} = 1.8 \text{ TeV}$

E 735 at Tevatron C0 intersection region of the Fermilab collider

Observables

- $\langle p_t \rangle$ vs multiplicity
- strange / non-strange vs mult.
 π, p, γ
 K, Λ, Ξ, ϕ



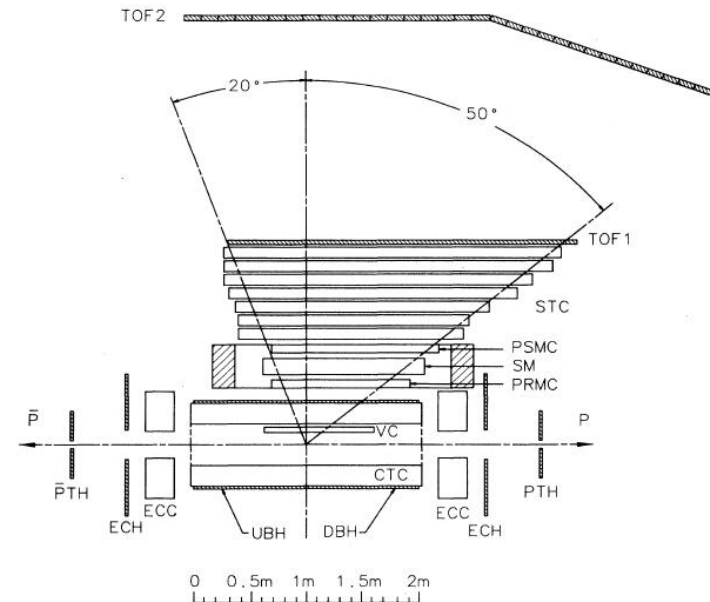
→ ruled out by SppS and Tevatron data

at \sqrt{s} energy

1800 GeV, 300, 546, 1000 GeV
 compare to SppS at 200, 546, 900 GeV

detector set-up

- multiplicity hodoscope $|\eta| < 3.25 \rightarrow$ Nch
- magnetic spectrometer $-0.36 < \eta < 1.0$
 18 degree azimuthal acceptance, $p > 150 \text{ MeV}/c$
 pre- and post-magnet drift chambers \rightarrow tracking
- ToF \rightarrow pid



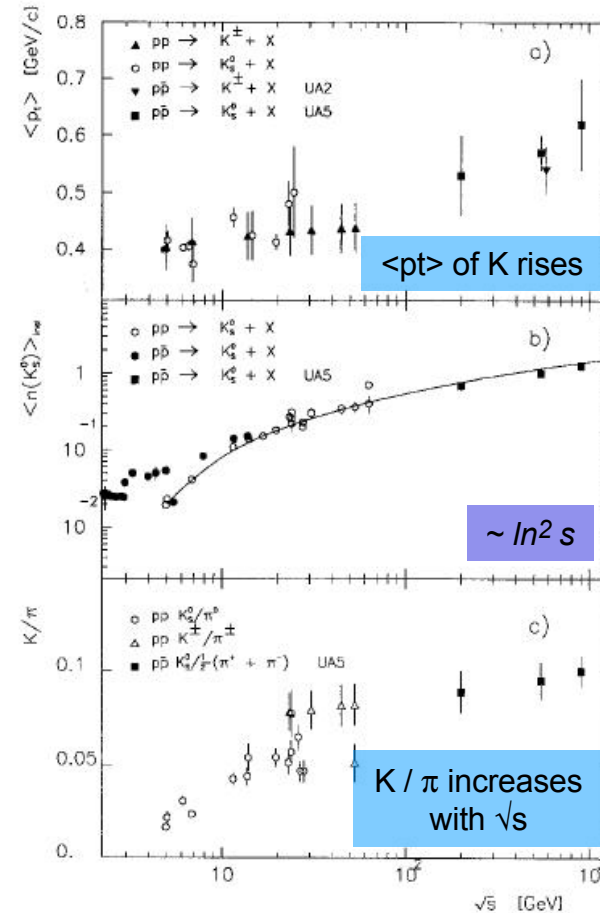
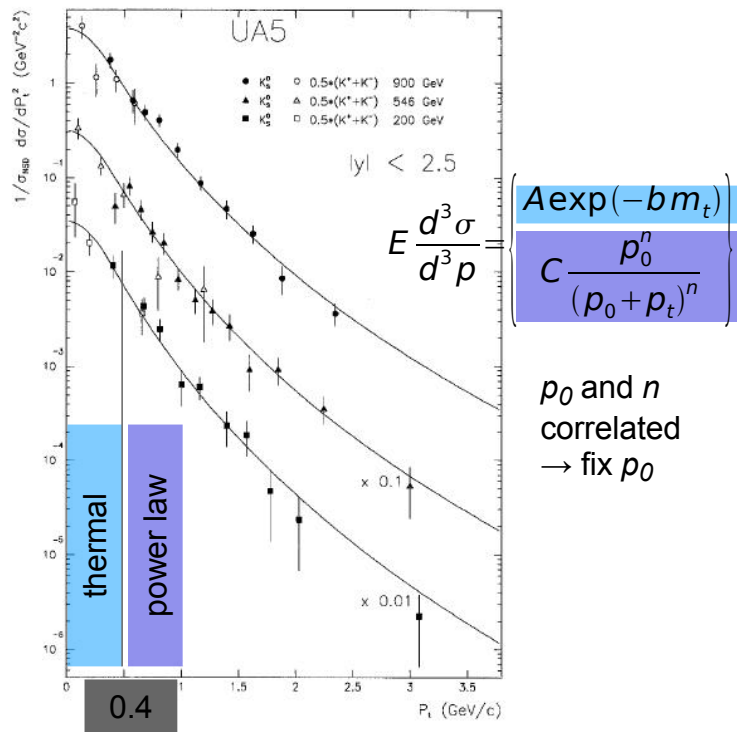
K/π ratio vs \sqrt{s} at $Sp\bar{p}S$

UA5 Z Phys C 41 (1988) 179 $\sqrt{s} = 0.2 - 0.9$ TeV
Phys Lett B 199 (1987) 311

no B field → PID by decay topology

V0's: K0s, momentum from decay kinematics

kinks: K in 3 prong decay, **momentum!**



faster than expected from ISR

$\langle p_t \rangle$ might increase with multiplicity

like charged particles, width and height (dn/dy) scale with $\ln s$

canonical suppression diminishes as $\langle K \rangle$ increases

V0 yield vs multiplicity

CDF Phys. Rev. D 72 (2005) 052001 $\sqrt{s} = 1800, 630 \text{ GeV}$

comparative study of event structure

- as function of multiplicity

- as function of E_t

hard events: E_t or $p_t > 1.1 \text{ GeV}$ in jet cone of $R = 0.7$

soft events: no high-pt jet

pt distribution

pt (K0s) $> 0.8 \text{ GeV/c}$

pt (Λ) $> 1.1 \text{ GeV/c}$

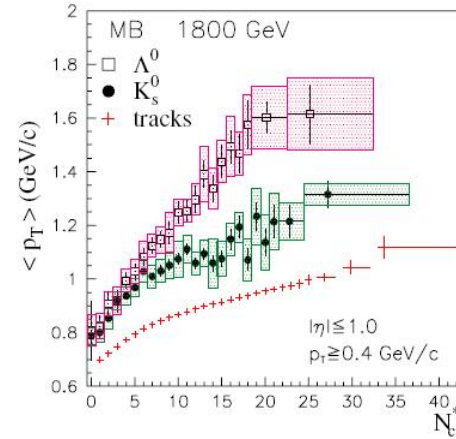
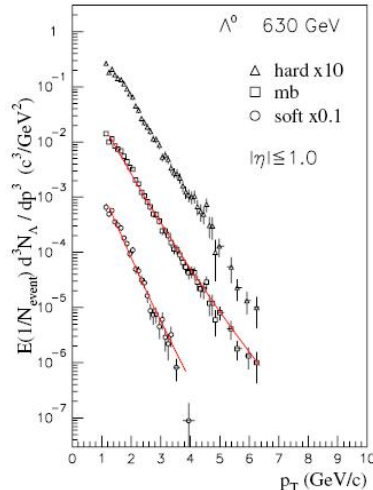
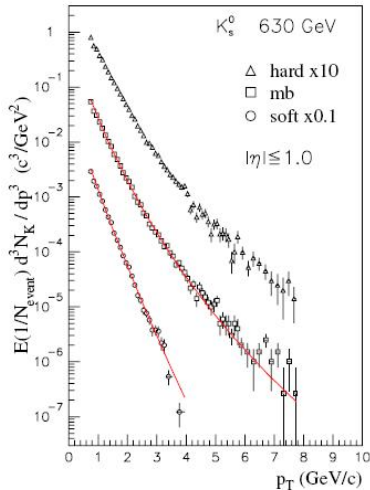
extrapolate with power law

down to $p_t = 0.4 \text{ GeV/c}$

part below missed

→ only high-pt measured → **pQCD part**

→ **soft physics missed ?**



mean-pt
increases with mass

at low multiplicity:

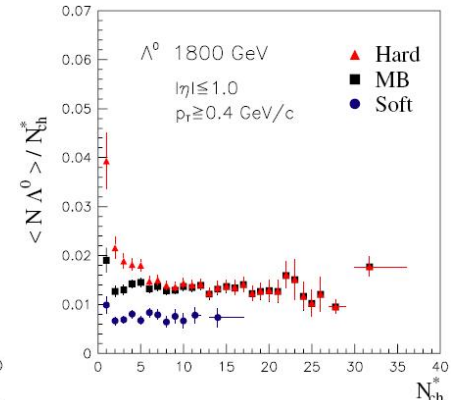
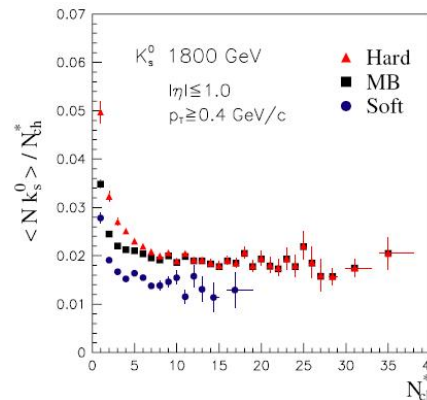
\sqrt{s} invariant

V0 / Nch vs multiplicity

const

yield dominantly in extrapolated part

→ **might be artifact of extrapolation**



→ hard events favour larger K/ π and larger $\langle p_t \rangle$ at same mult.

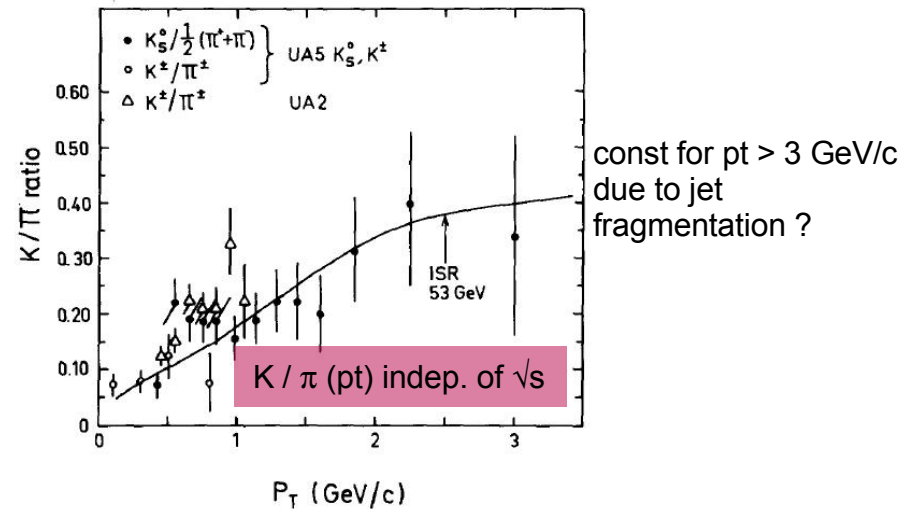
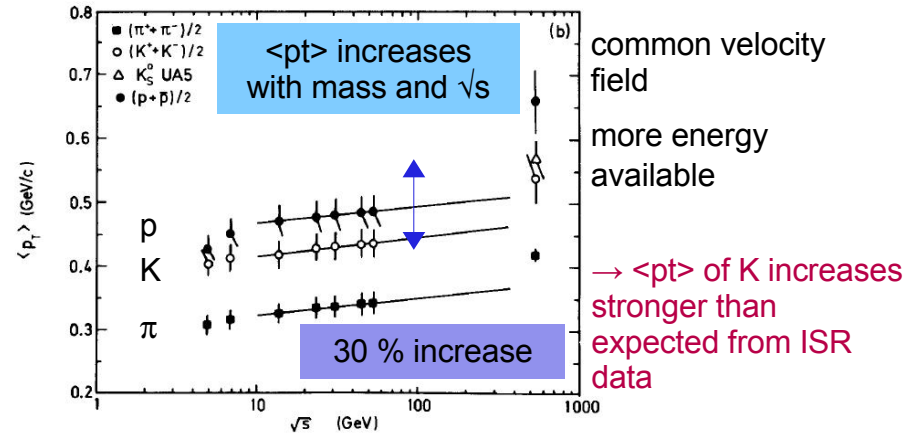
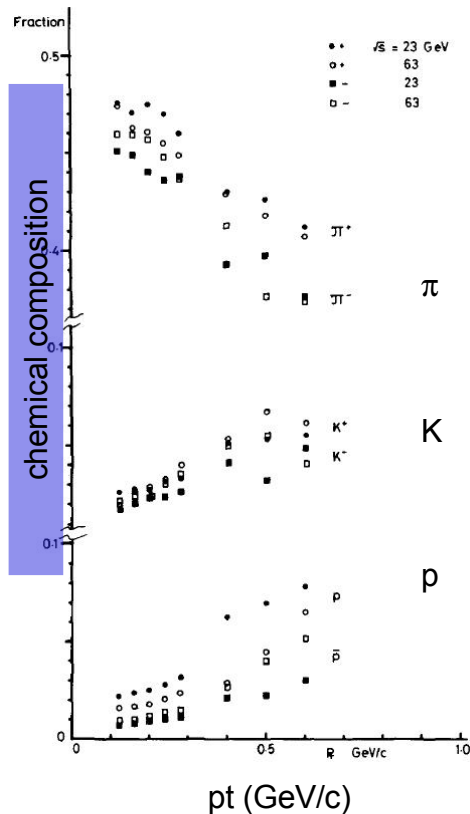
K/ π ratio vs pt

ISR Phys Lett 64B (1976) 111 $\sqrt{s} = 23 - 63$ GeV
Nucl Phys B 116 (1976) 77

π : $40 < p_t < 400$ MeV/c
K: $100 < p_t < 300$ MeV/c
p: $100 < p_t < 500$ MeV/c

p+p

UA5 Nucl Phys B 258 (1985) 505 $\sqrt{s} = 540$ GeV



strangeness vs multiplicity

E 735 Phys Rev Lett 64 (1990) 991 $\sqrt{s} = 1800$ GeV

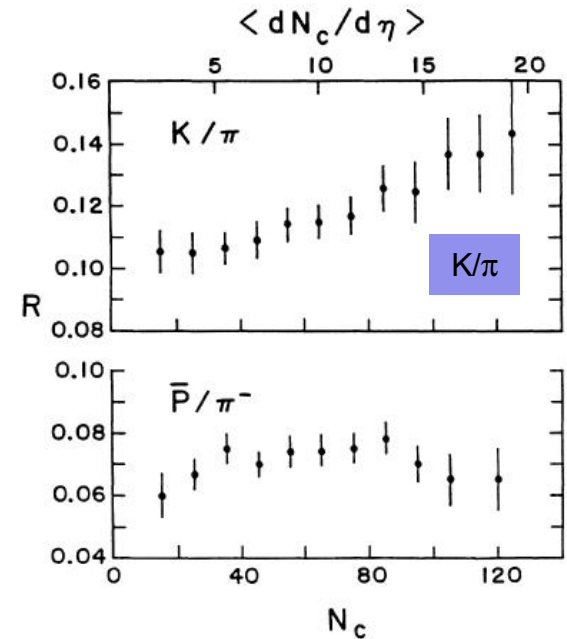
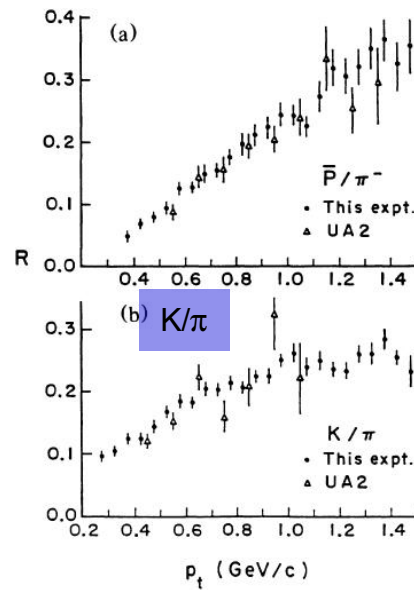
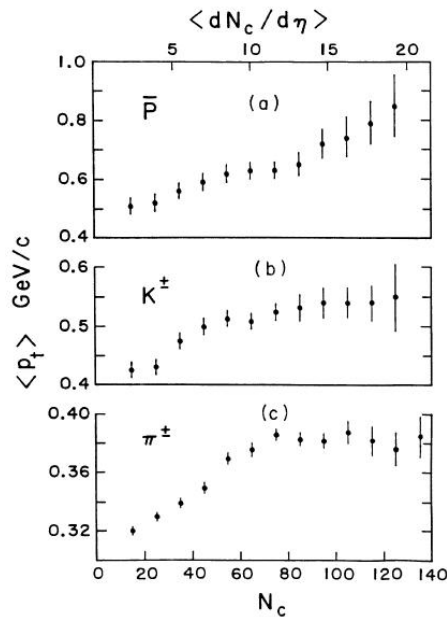
Run I data

K/ π

$0.25 < p_t < 1.5$ GeV/c

$$\frac{dN/dp_t^2}{B(p_t + p_0)^{-n}} = \frac{A \exp(-a p_t)}{\pi} \frac{K}{\pi}$$

- plateau in $\langle p_t \rangle$ of π and K?
- K/π (p_t) is \sqrt{s} invariant
- K/π increases with multiplicity



strangeness vs \sqrt{s} and multiplicity

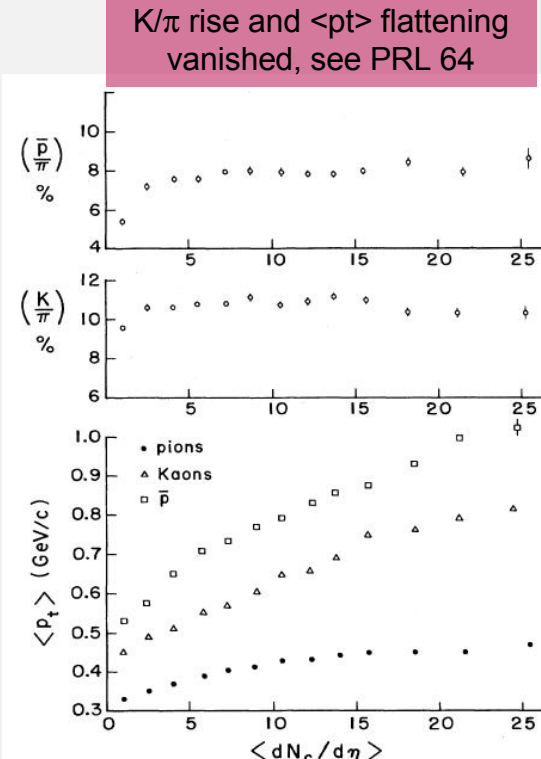
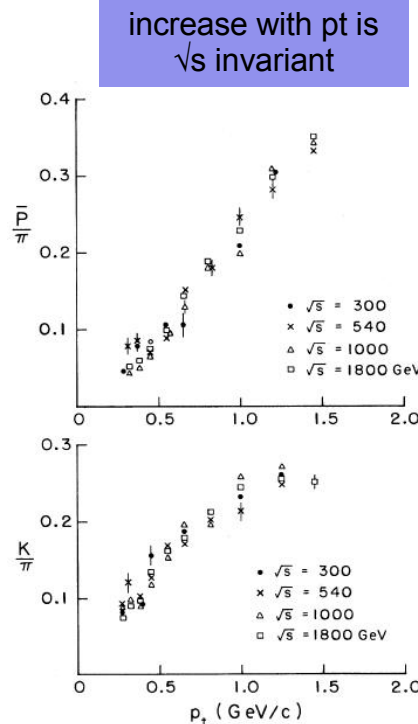
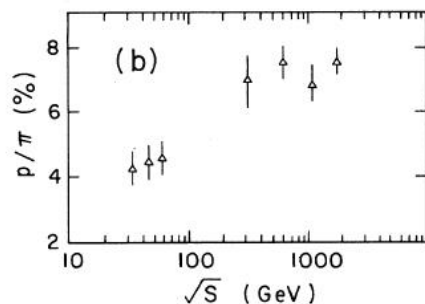
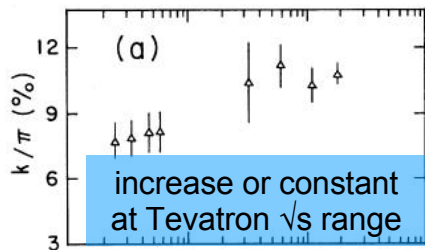
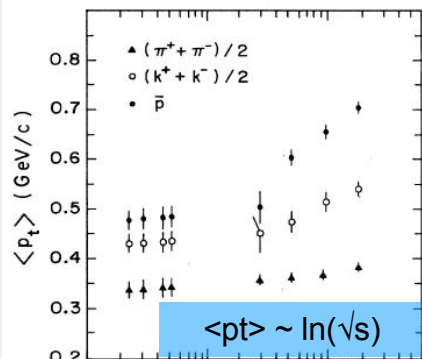
E 735 Phys. Rev. D 48 (1993) 984 $\sqrt{s} = 300, 540, 1000, 1800$ GeV

Run II data

K/ π

$0.25 < p_t < 1.5$ GeV/c

$$\frac{dN}{dp_t^2} = \frac{A \exp(-a p_t)}{B (p_t + p_0)^{-n}} \quad \begin{array}{l} K \quad \langle p_t \rangle = 2/a \\ \pi \quad \langle p_t \rangle = 2p_0 / (n-3) \end{array}$$



hyperon production vs multiplicity

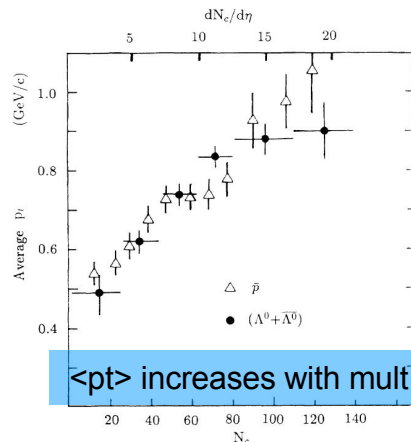
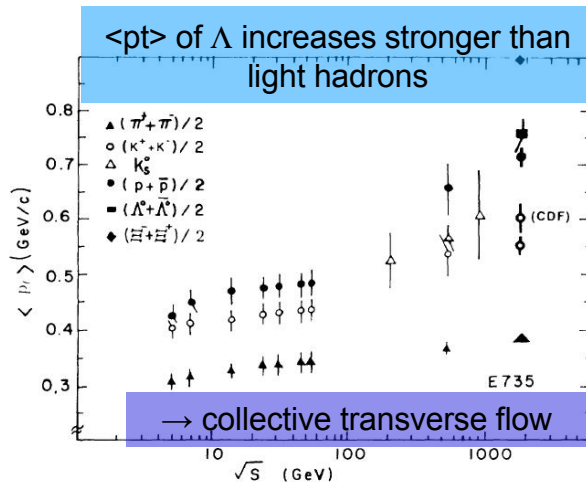
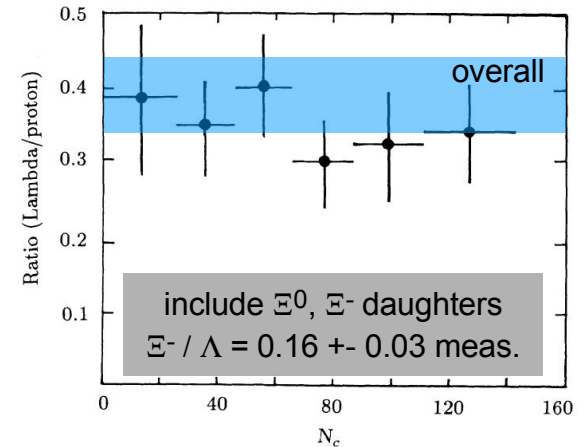
E 735 Phys. Rev. D 46 (1992) 2773 $\sqrt{s} = 1.8 \text{ TeV}$
 Nucl. Phys. A 544 (1992) 343c
 Nucl. Phys. A 525 (1991) 165c

Λ / N_{ch} , Λ / p $0.5 < p_t < 1.5 \text{ GeV}/c$

$$p = p + \Lambda + \Sigma^+$$

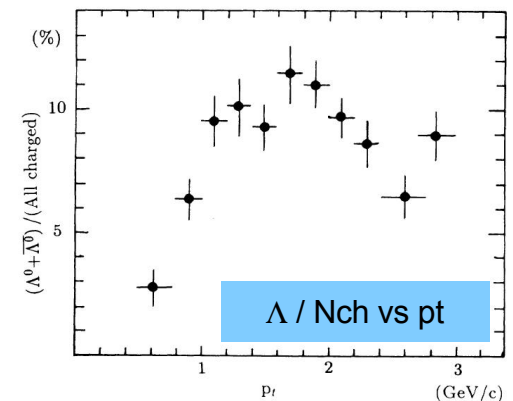
$$\Lambda = \Lambda + \Sigma^0 + \Xi^0 + \Xi^-$$

$$\frac{dn}{p_t dp_t} = A \exp(-b p_t)$$



Λ / p vs multiplicity is constant

p / π is const $\rightarrow \Lambda / \pi$ flat too

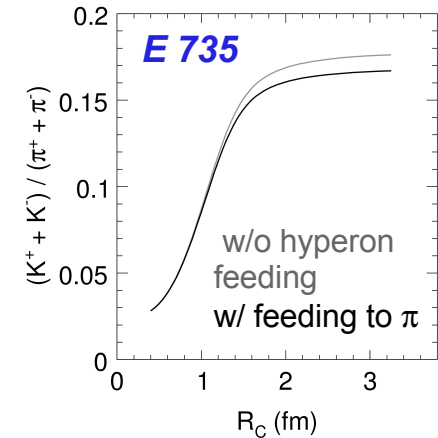


comparison statistical model vs data

Statistical model hep-ph/0407174

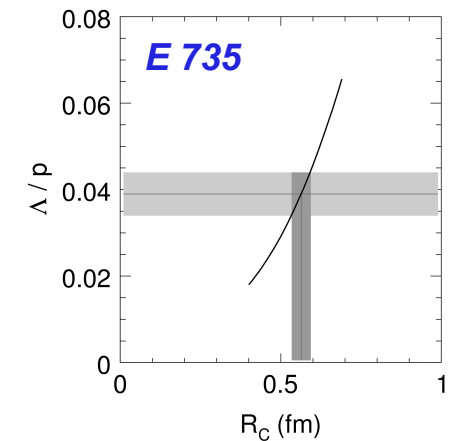
K/ π vs multiplicity

- overall $K/\pi \approx 0.1$
- K/π (multiplicity) ≈ 0.1 *E735*
- CDF K/π not comparable to model due to pt cut-off
- consistent with statistical model with significant canonical suppression of strange-particle phase-space



Λ/p vs multiplicity

- overall $\Lambda/p \approx 4\%$
- Λ/p vs multiplicity const
- CDF Λ/p not comparable to model due to pt cut-off
- consistent with statistical model
- canonical suppression seems stronger than in K data



→ **more data** – on multi-strange hyperons – needed to check whether model applies in jet-dominated regime

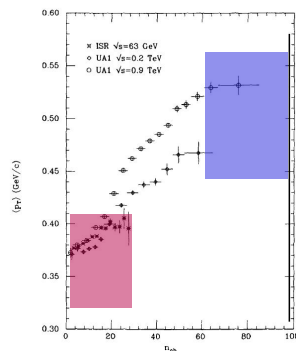
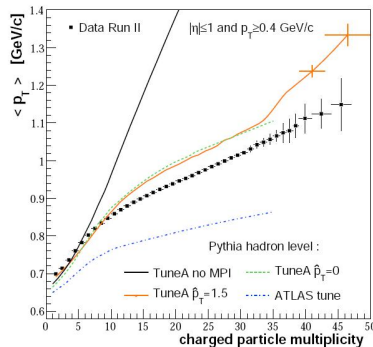
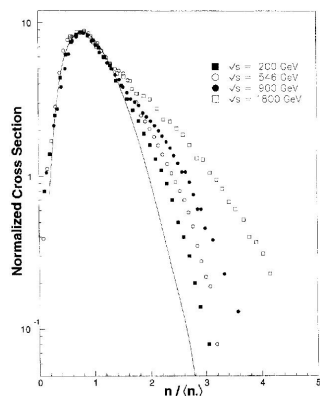
strangeness canonical
 $T = 170$ MeV, $\mu_B = \mu_Q = 0$

summary

multiplicity and $\langle pt \rangle$

- mean multiplicity rises stronger than $\ln s$ at all energies
- KNO type distribution exhibits shoulder
- spectra consist of two components, thermal and power law
- mean- pt vs mult. rises moderately
- complex event structure in p+p
- no QPG signature

- low multiplicity: mean- $pt \sim \sqrt{s}$ invariant



strangeness production

- $\langle pt \rangle$ increases with mass, multiplicity and \sqrt{s}
- K/π increases with \sqrt{s}
- K/π vs multiplicity is constant → no QPG signature
- data show significant canonical suppression

