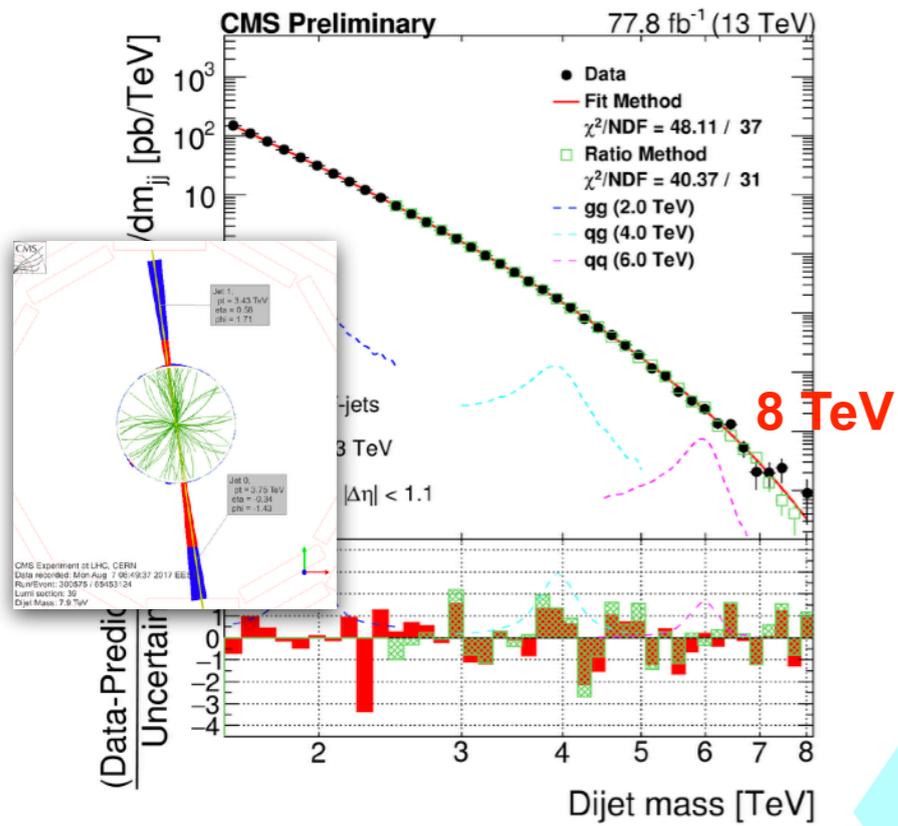


Les expériences du futur : potentiel de physique
Recherche de nouvelle physique



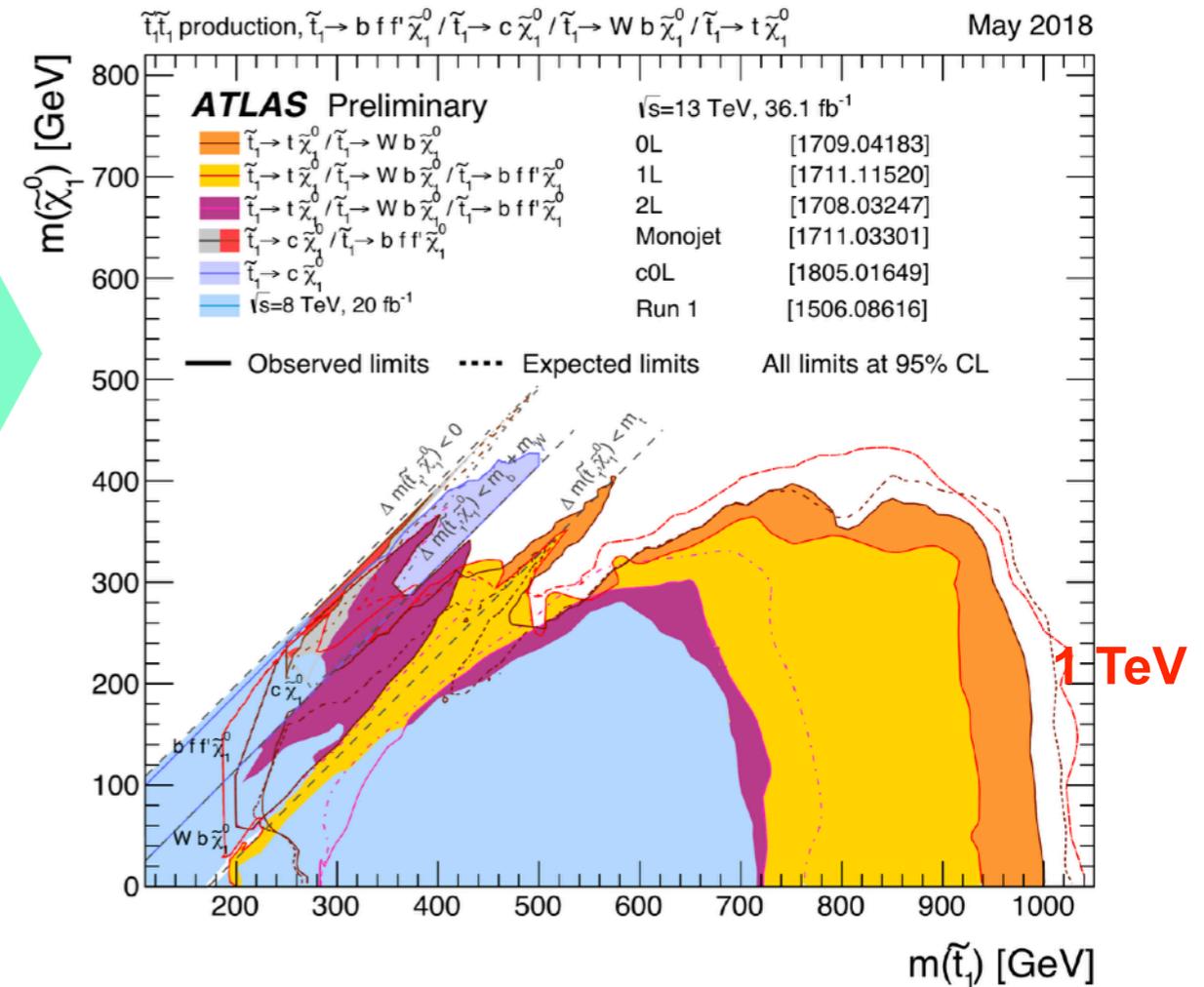
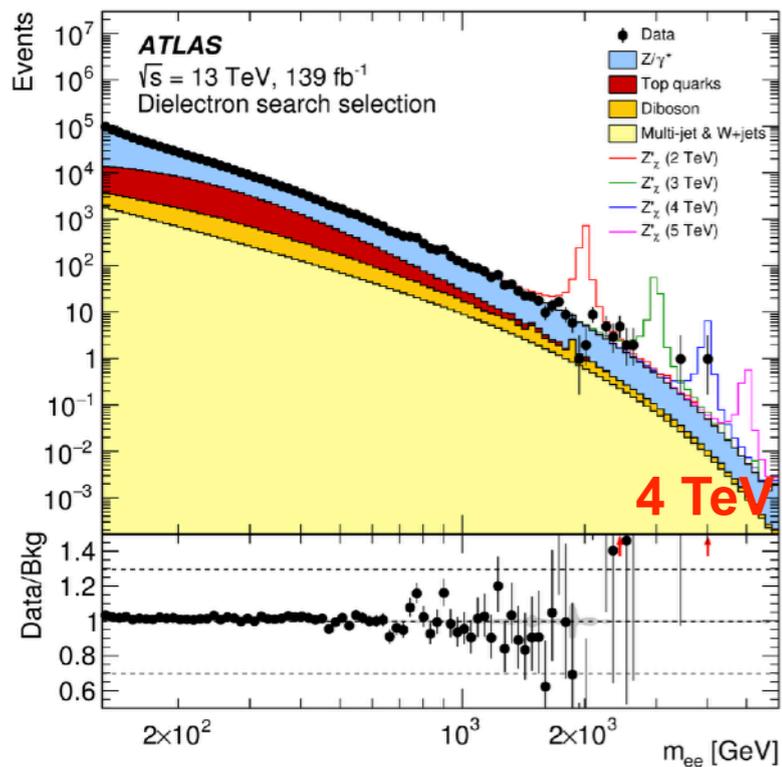
École de Gif 2019
*Questions ouvertes en physique
des particules*
2-6 septembre 2019
École polytechnique, Palaiseau

Search for Physics Beyond the SM (BSM)



search for Supersymmetry (here : stops)

search for new resonances



Résultat le plus marquant du LHC (après le Higgs)
 so far searches for new particles have not succeeded at unraveling new physics at the LHC

no evidence for SUSY at the TeV scale at the LHC

50% of the 900+ and 8 out of the last 10 CMS publications have title starting with "Search for" !

The Big Questions

For most physicists, search for New Physics (NP) is the most compelling motivation for future high-energy colliders, with the goal of providing elements of response to the big questions:

➡ **Is the Higgs boson a portal to new physics?**

- deviations from SM predictions (couplings, properties)
- Higgs potential and baryon asymmetry
- rare decay modes

➡ **Is the Higgs boson fundamental or composite?**

- composite Higgs
- new particles associated with EWK symmetry breaking
- heavy gauge bosons

➡ **Are there new particles or interactions around at the TeV scale?**

- supersymmetry (SUSY)
- extended Higgs sectors

➡ **Can we close the search for thermal relic WIMPs at future colliders?**

- Dark Matter (DM)

➡ **Can we probe the feebly-interacting sector?**

- feebly-interacting particles (axions, ALPs, dark photons)
- right-handed neutrinos

for these lectures, only a few chosen examples of searches at colliders, for illustration

Models with 1st Order Phase Transition

Question: is there BSM physics coupled to the Higgs capable of modifying the Higgs potential and allow for a 1st order EW phase transition as required for BAU?

Below, an example of minimal extension from the SM potential

➔ the SM Higgs potential

$$V_0^{\text{SM}}(H) = -\mu^2 |H|^2 + \lambda |H|^4$$

➔ the Real Scalar Singlet class of models

- add a single real singlet S (1 ddl) with coupling to the Higgs field (4 ddl)

$$V_0(H) = V_0^{\text{SM}}(H) - \frac{1}{2}\mu_S^2 S^2 + \frac{1}{4}\lambda_S S^4 + \lambda_{HS} |H|^2 S^2$$

(simplified)

with

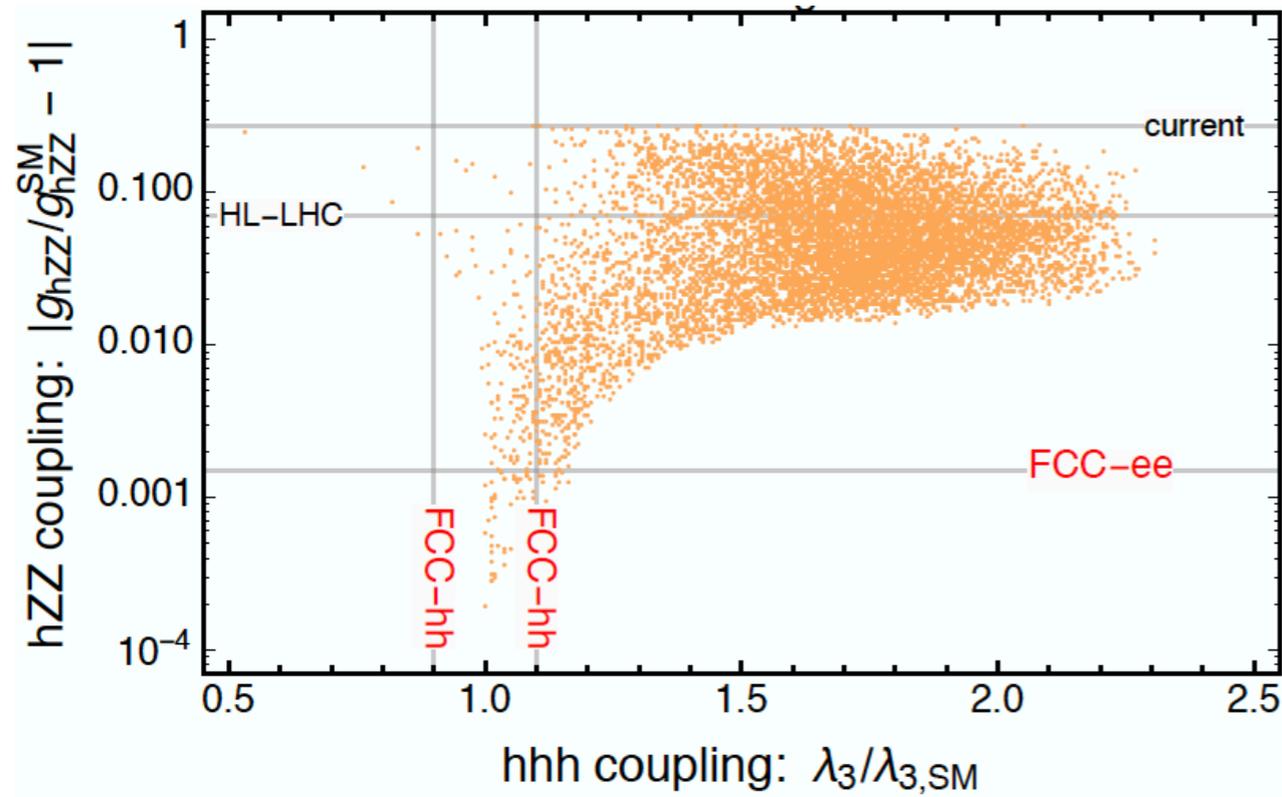
$$\begin{aligned} h &= h_0 \cos \gamma + S \sin \gamma \\ \phi &= -h_0 \sin \gamma + S \cos \gamma \end{aligned}$$

Equivalence theorem:

$$B(\phi \rightarrow hh) = B(\phi \rightarrow ZZ) = 25\%$$

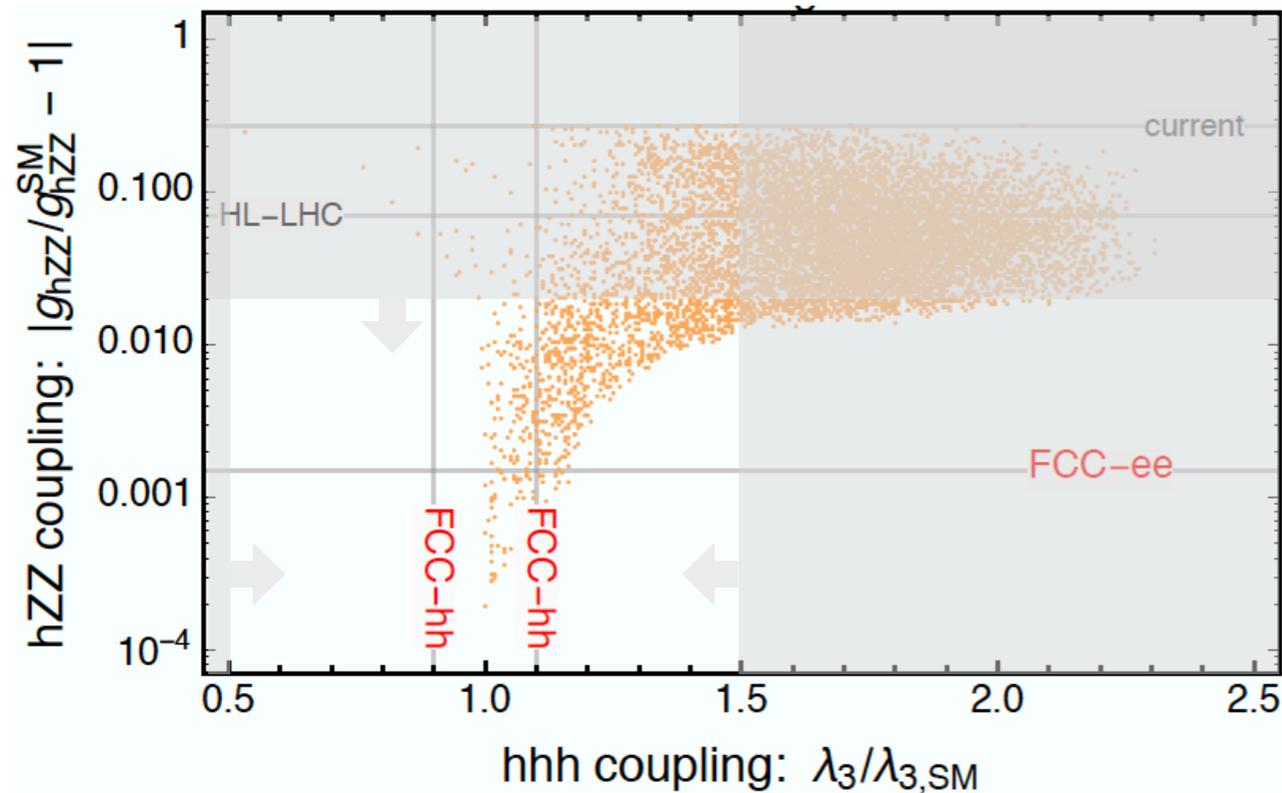
Real Scalar Singlet Model

Parameter space scan: the points indicate a first order phase transition



Real Scalar Singlet Model

Parameter space scan: the points indicate a first order phase transition

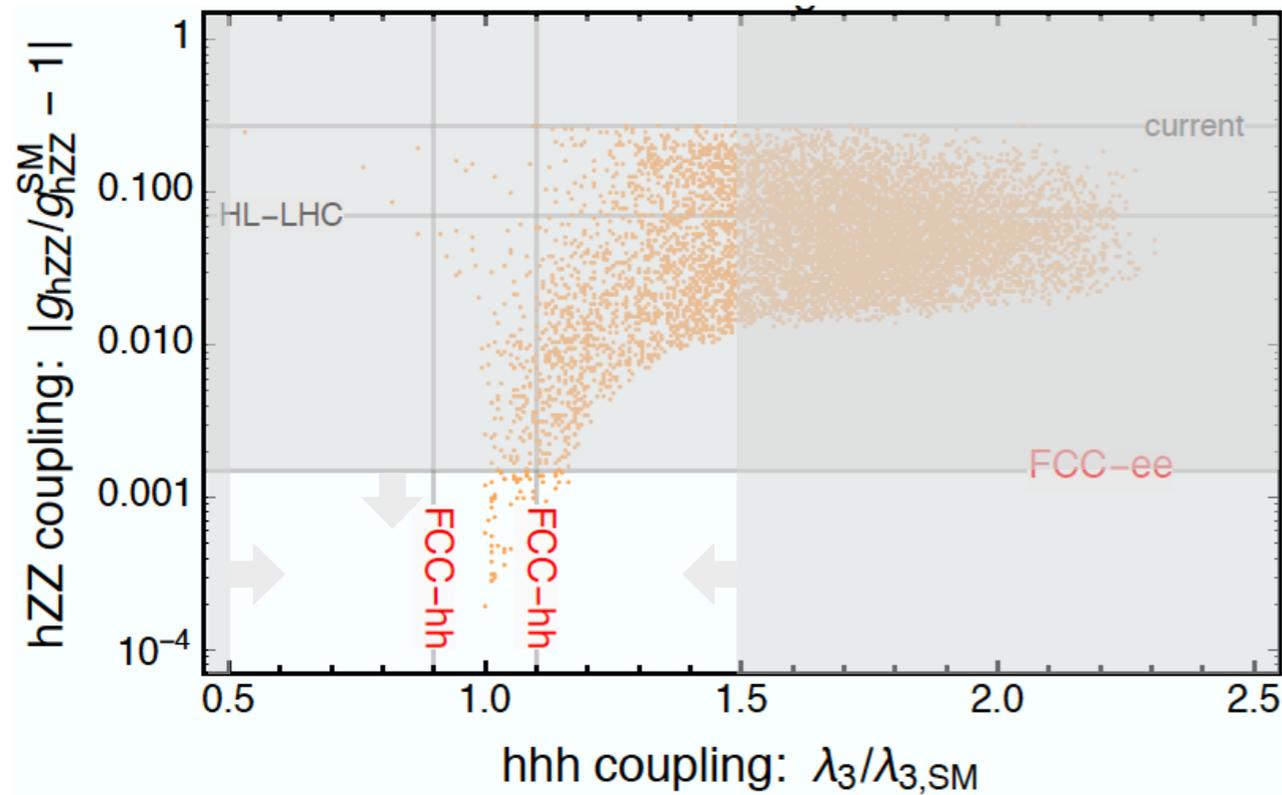


HL-LHC

► g_{HZZ} to 0.02 and λ to 50%
probe a good portion of parameter space, but not all

Real Scalar Singlet Model

Parameter space scan: the points indicate a first order phase transition

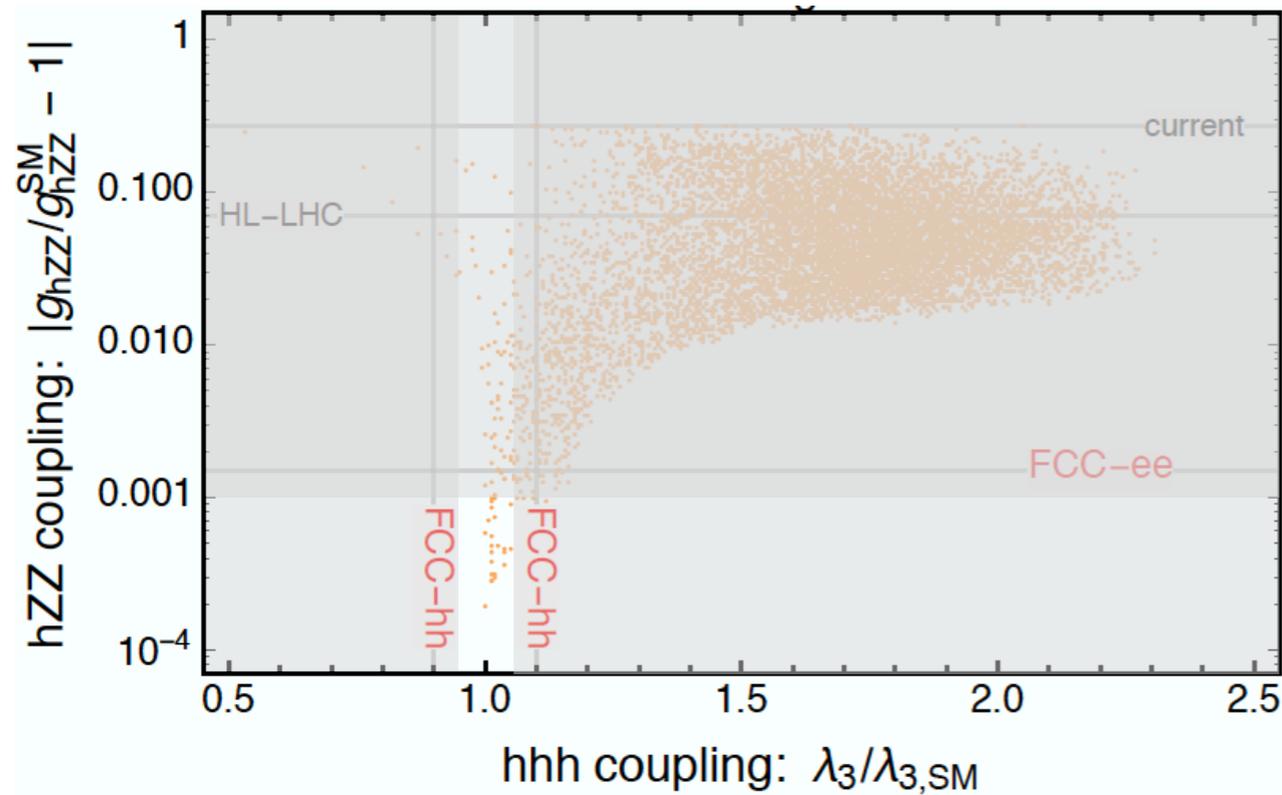


FCC-ee

→ g_{HZZ} to 1.5×10^{-3} allows to probe most of the parameter space

Real Scalar Singlet Model

Combined constraints from precision Higgs measurements at FCC-ee and FCC-hh

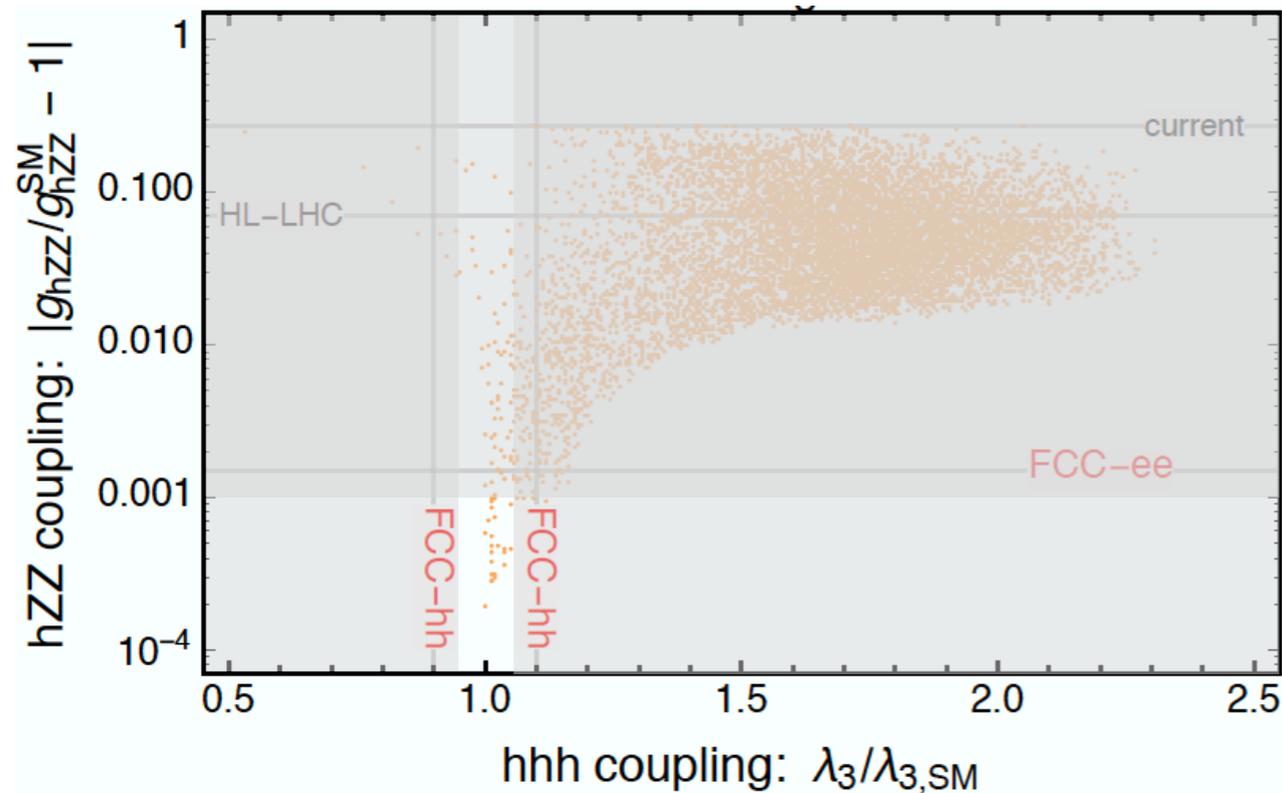


FCC-ee/hh

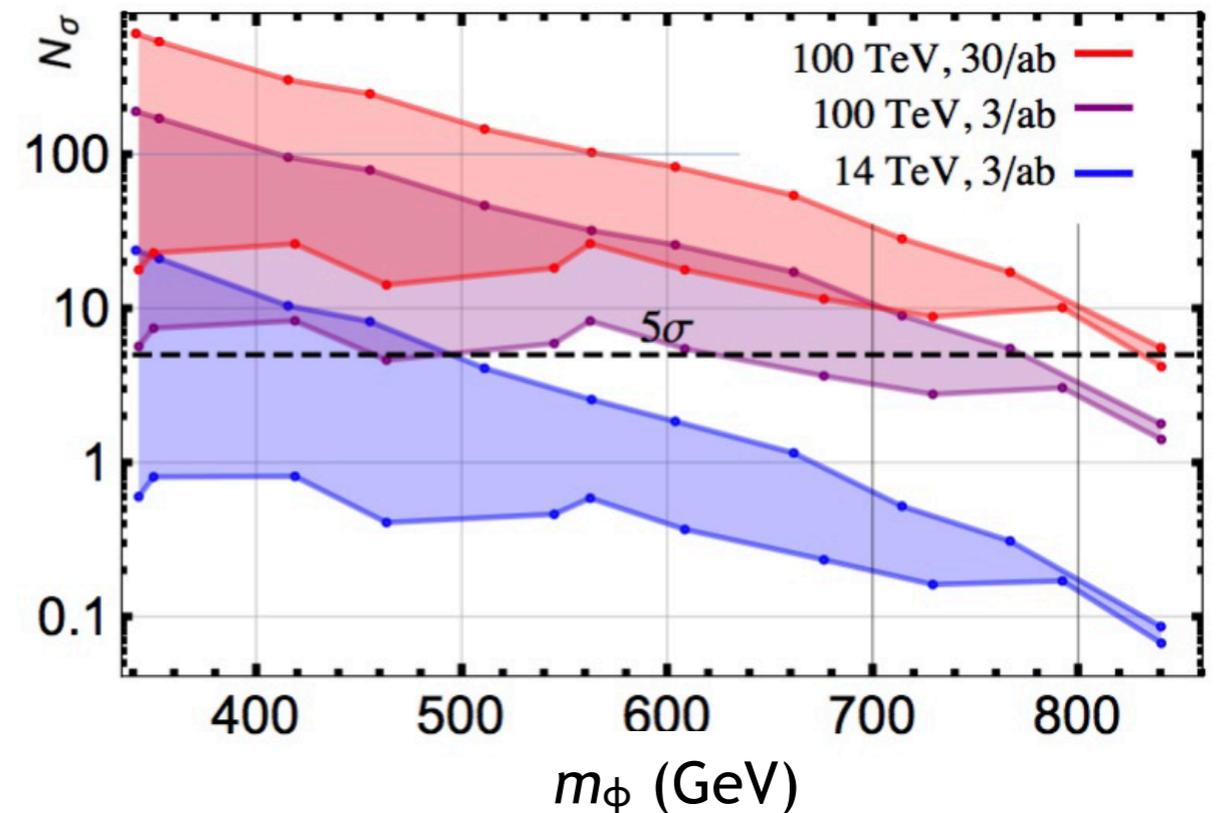
→ g_{HZZ} to 1×10^{-3} and λ to 6%
leaves only a very small portion
of the parameter space

Real Scalar Singlet Model

Combined constraints from precision Higgs measurements at FCC-ee and FCC-hh



Direct detection of extra heavy scalar states at FCC-hh



$$\phi \rightarrow hh \rightarrow b\bar{b}\gamma\gamma/\tau^+\tau^-$$

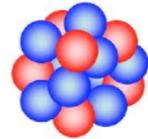
Operators that lead to visible deviations in the Higgs self coupling are likely to first manifest themselves through deviations of single-Higgs couplings to gauge bosons

Self-coupling measurements remain important to interpret any SM departure in single-Higgs couplings to understand the origin of these deviations

Higgs Compositeness

Question: is the Higgs point-like? If not, how big is it?

- can compositeness explain the smallness of the electroweak scale and solve the Naturalness problem?



Basic idea of Composite Higgs (CH) models

- all the degrees of freedom of the SM apart from the Higgs are elementary
- the Higgs instead arises as a bound state from a strong dynamics
- such dynamics is roughly described by two parameters:
the overall mass scale m_* and its overall coupling strength g_*

$$\mathcal{L}_U^{d=6} = c_H \frac{g_*^2}{m_*^2} \mathcal{O}_H$$

Scaling of Higgs couplings by a common factor

$$+ \frac{1}{g_*^2 m_*^2} [c_{2W} g^2 \mathcal{O}_{2W} + c_{2B} g'^2 \mathcal{O}_{2B}]$$

4-fermion contact interactions
W', Z' resonances

$$+ \frac{1}{m_*^2} [c_W \mathcal{O}_W + c_B \mathcal{O}_B]$$

2-fermion 2-boson contact interactions

$$+ c_{y_t} \frac{g_*^2}{m_*^2} \mathcal{O}_{y_t} + c_{y_b} \frac{g_*^2}{m_*^2} \mathcal{O}_{y_b}$$

Modified top and bottom
Yukawa couplings

+ ...

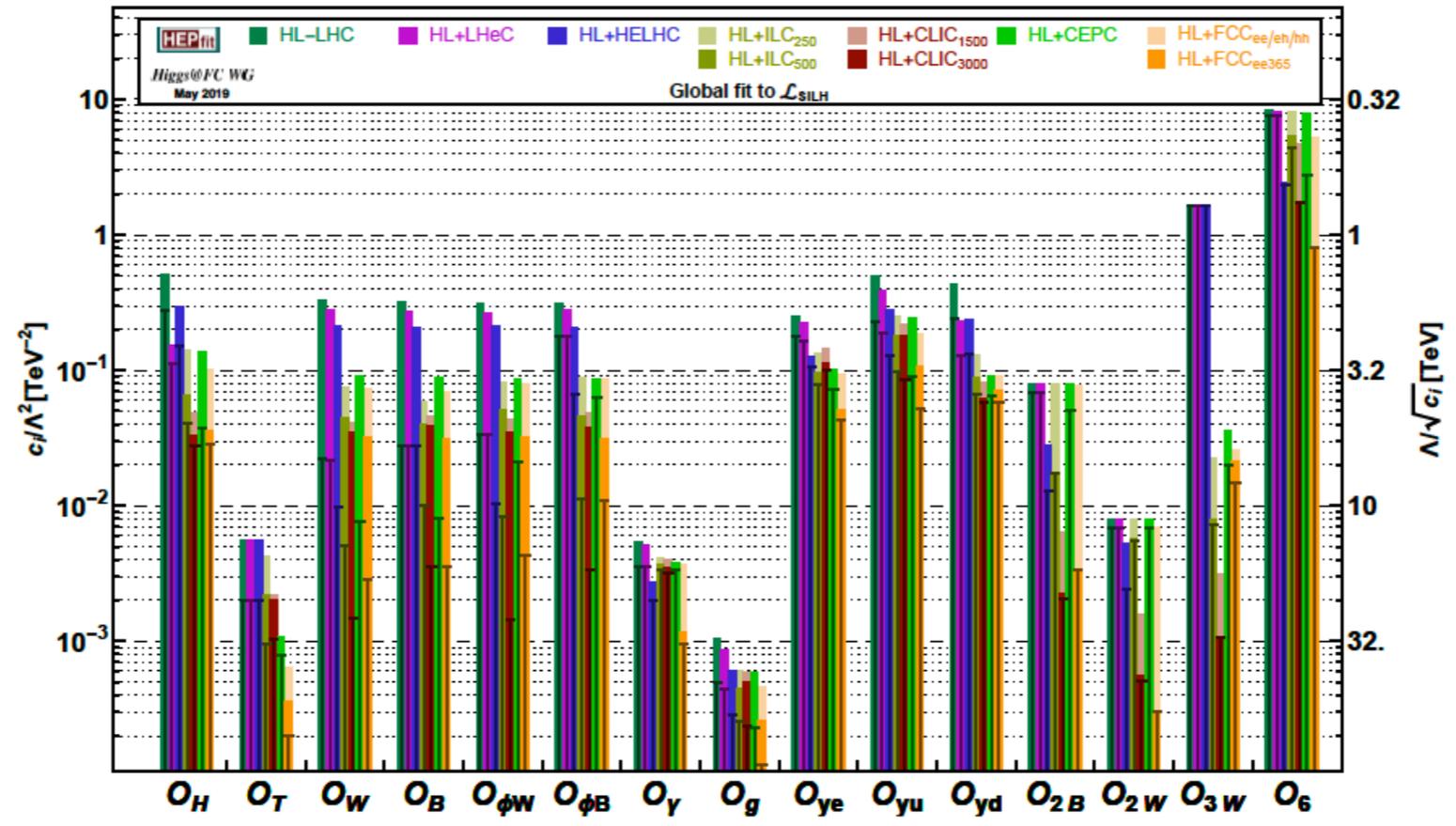
and more

this physics is encoded in
Strongly-Interacting Light
Higgs (SILH)
dimension-6 operators

see backup

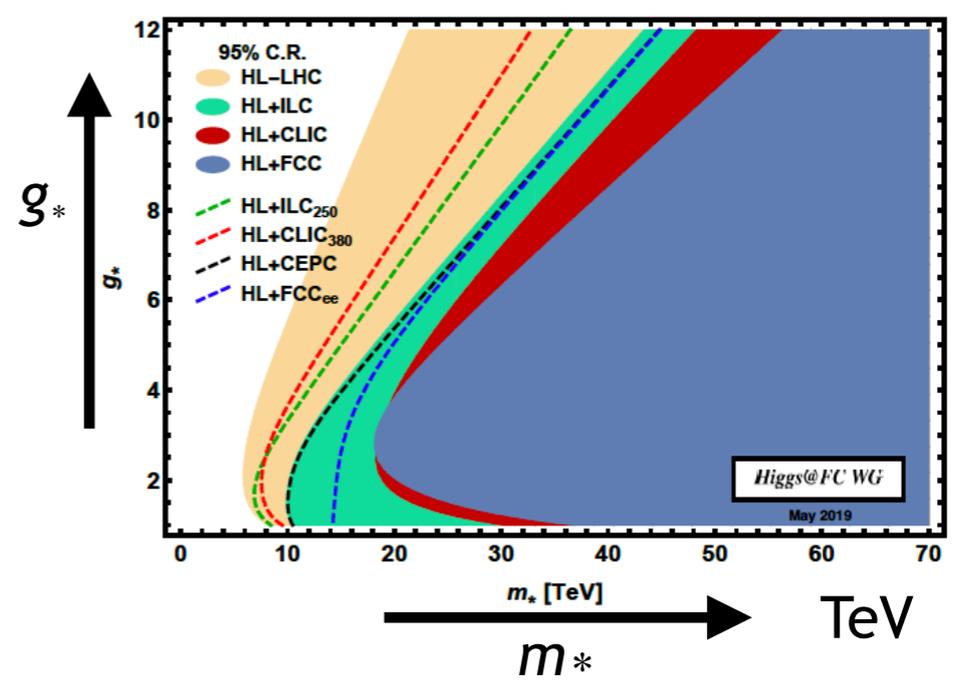
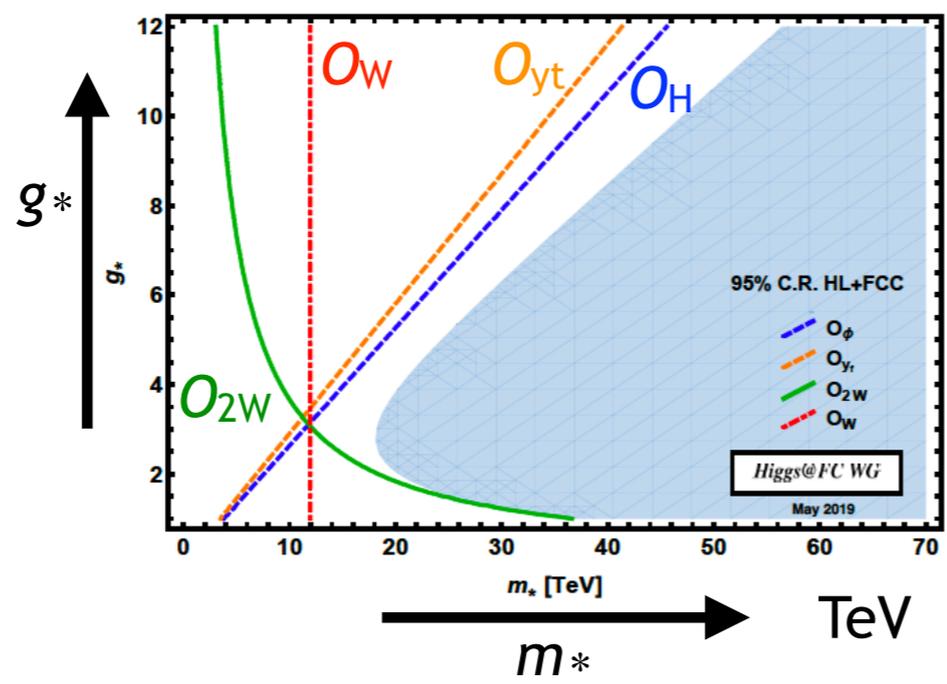
Bounds on SILH Lagrangian Interactions

result of the SILH fit on Higgs and EW precision observables for various machines

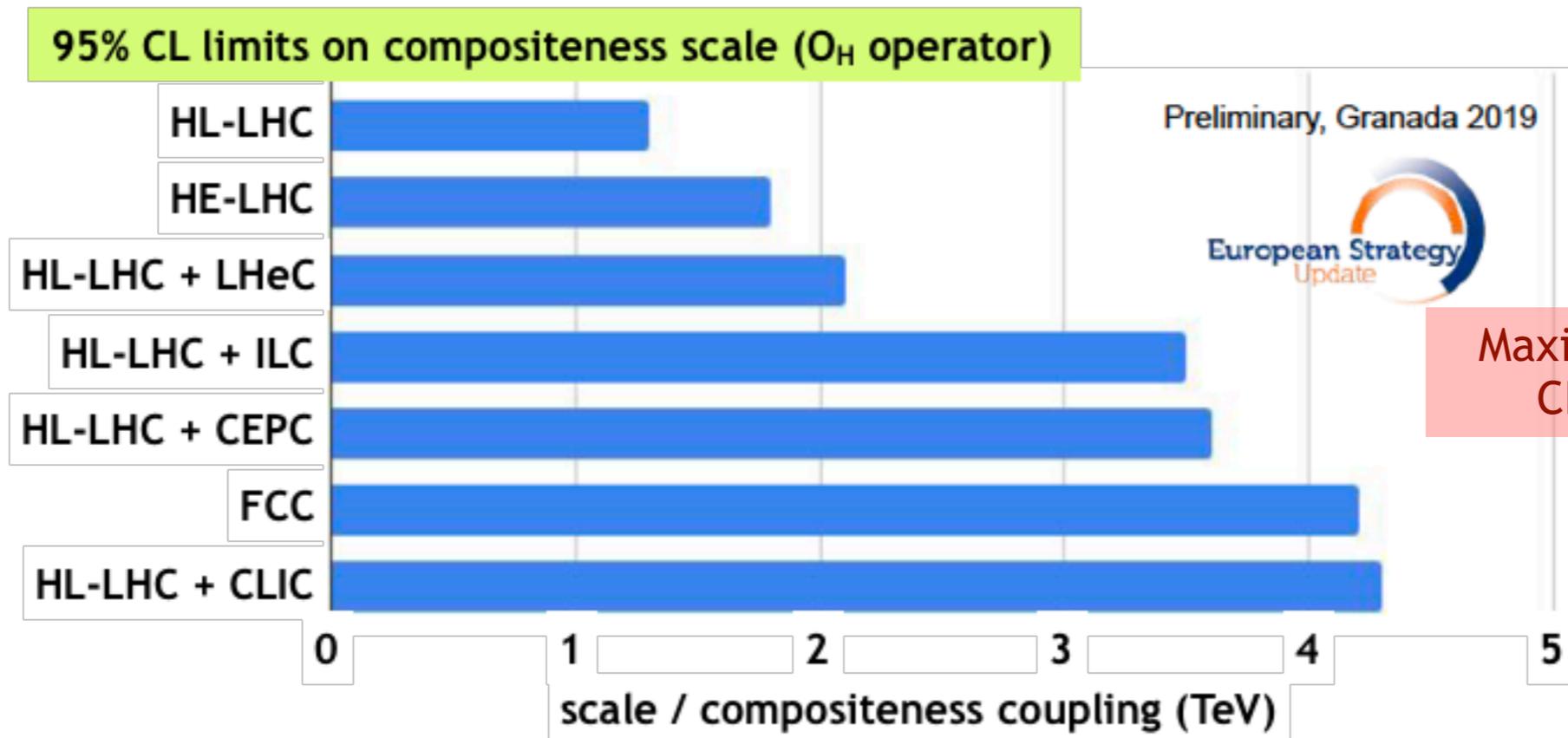
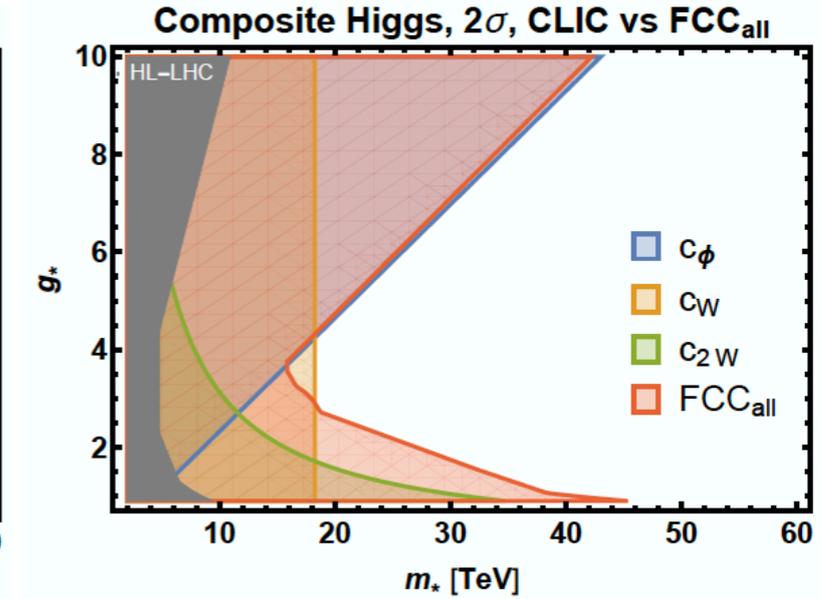
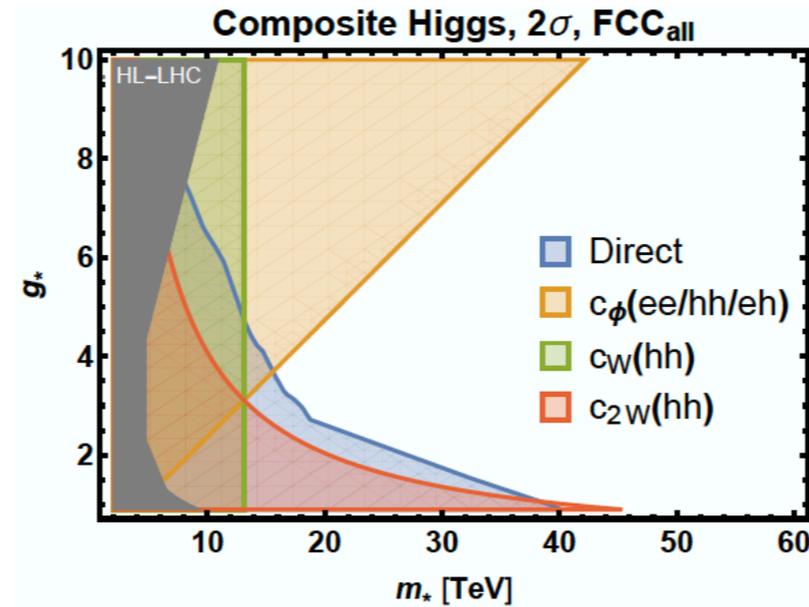
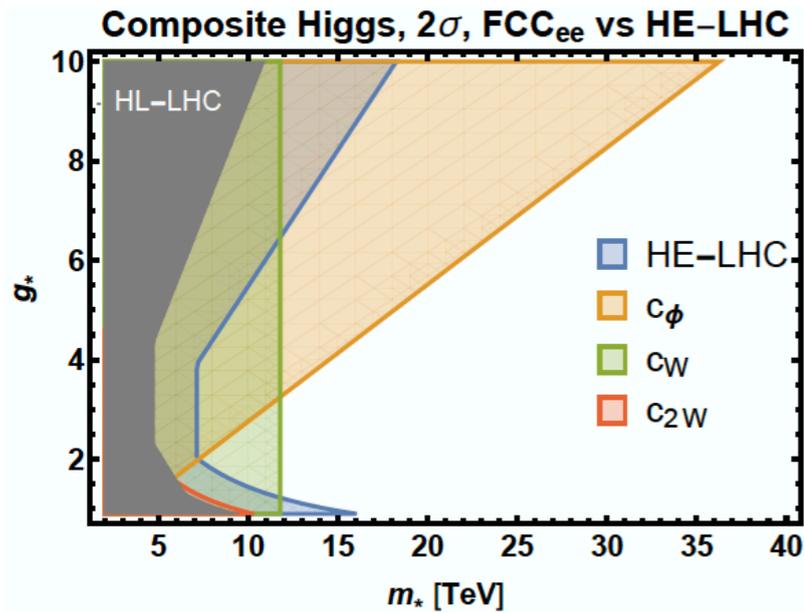


Bottom line: important BSM information from precision measurements

constraints on compositeness scale

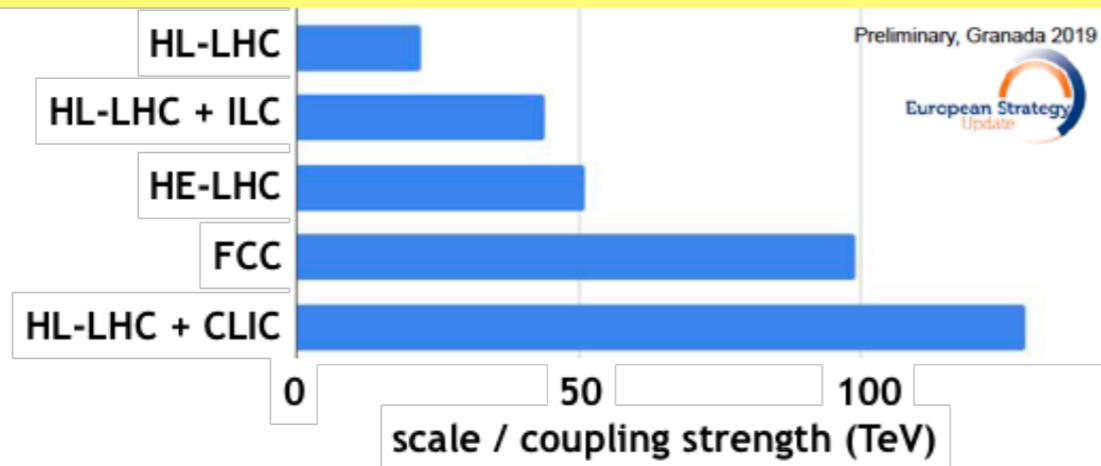


Higgs Compositeness

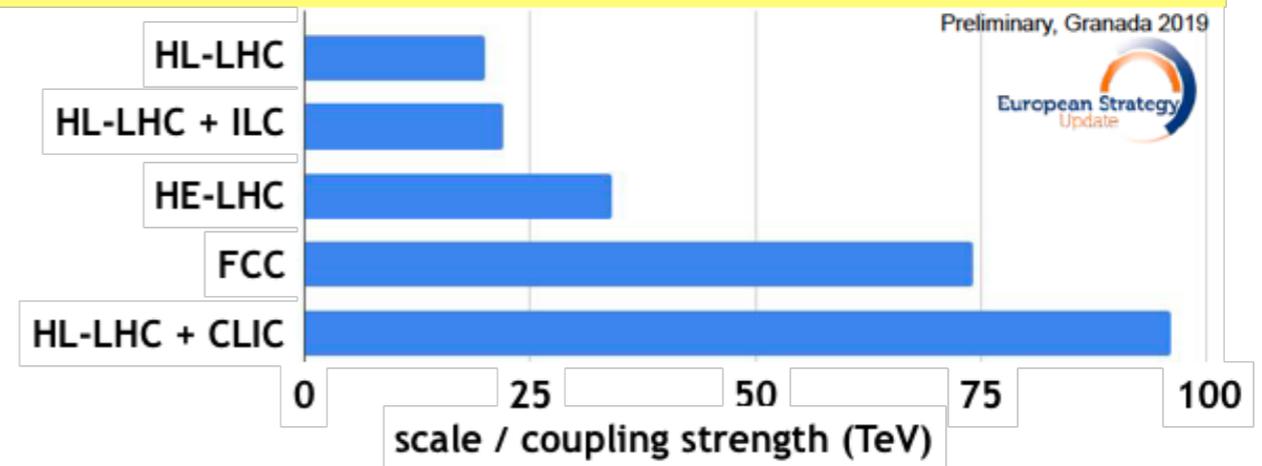


Compared SILH Reach at Future Colliders

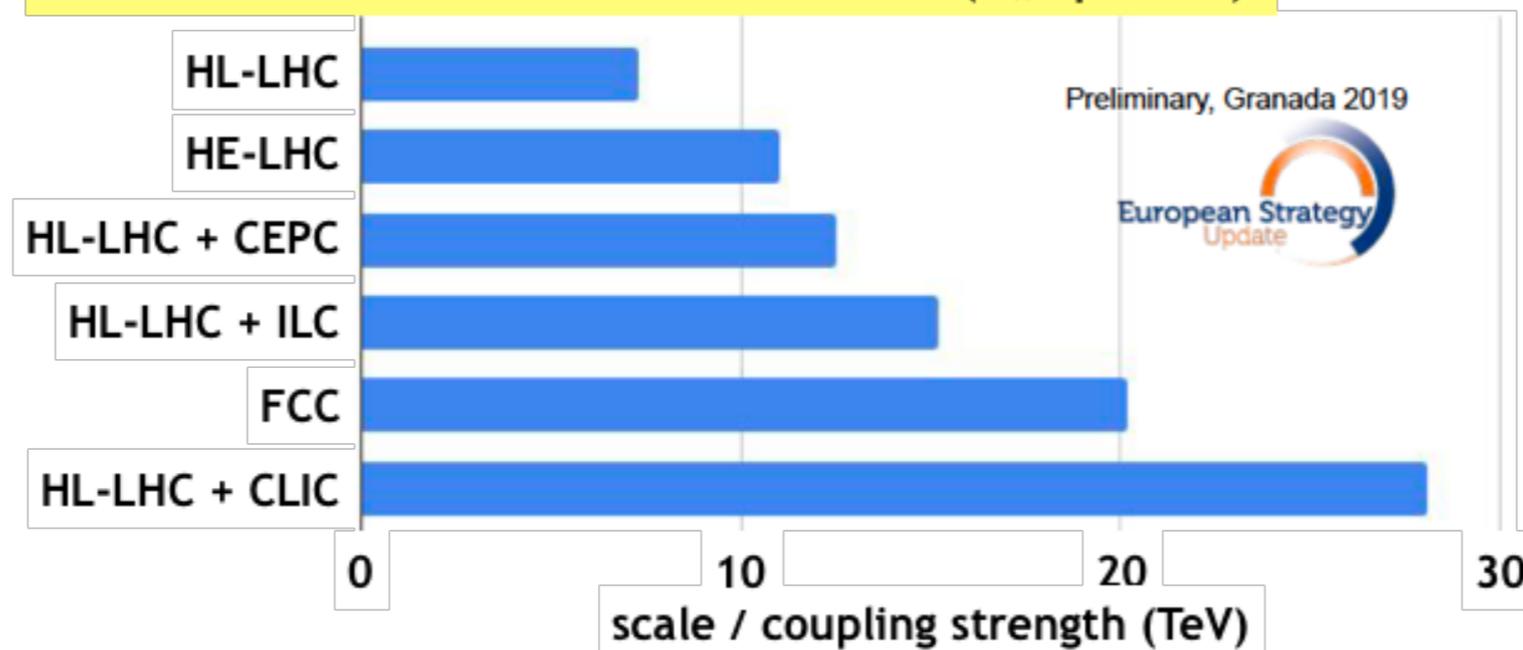
95% CL scale limits on 4-fermion contact interactions (Y couplings)



95% CL scale limits on 4-fermion contact interactions (W couplings)



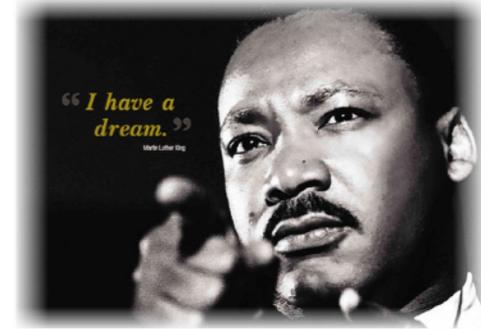
95% CL scale limits on contact interactions (O_W operator)



Bottom line:
FCC-hh and CLIC
always win...

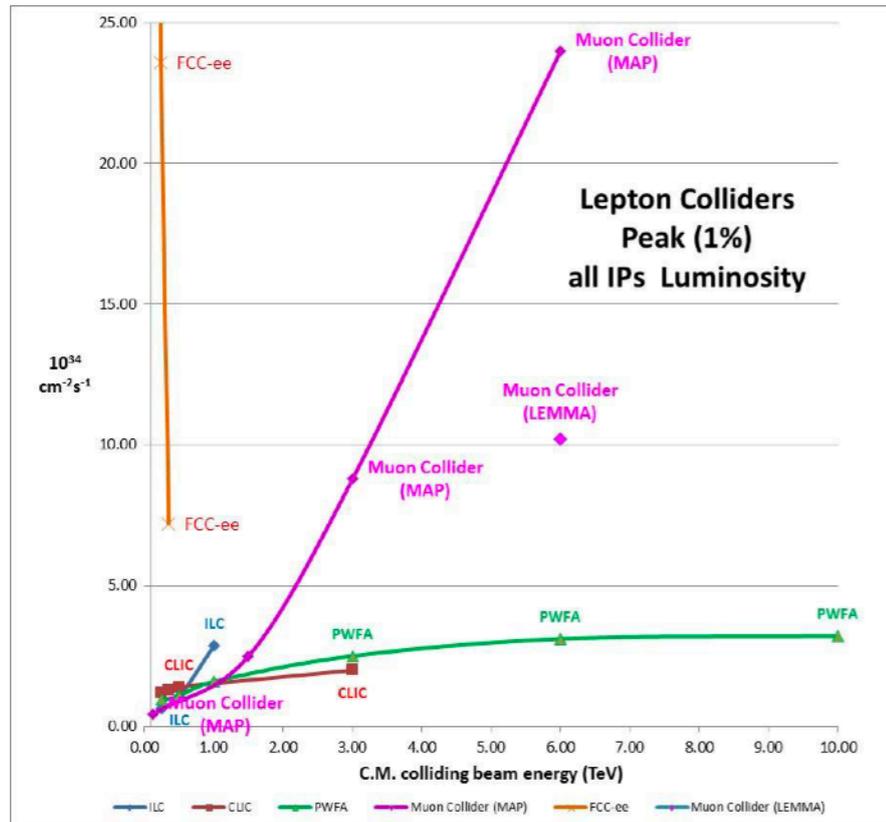
What about muon colliders?

A Very High-Energy Muon Collider?



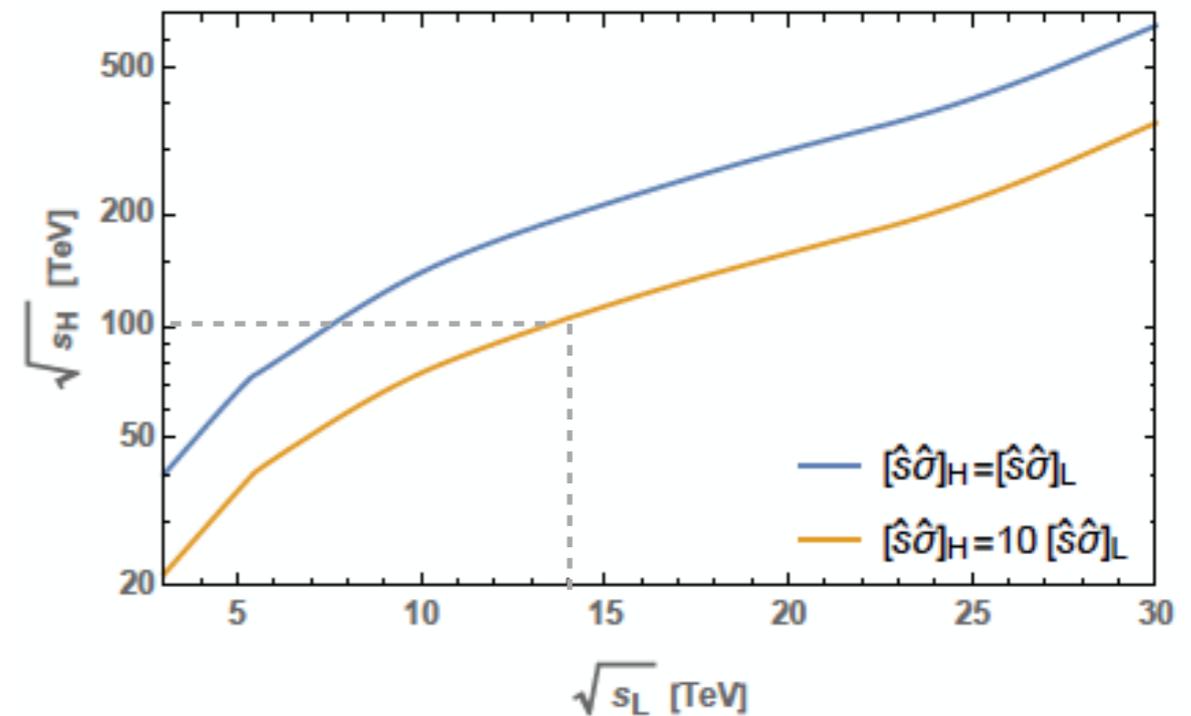
Let's have a dream...

... Imagine a muon collider at $\sqrt{s} = 14$ TeV running for 5 years at $\mathcal{L} = 2 \times 10^{35} \text{ cm}^{-2}\text{s}^{-1}$



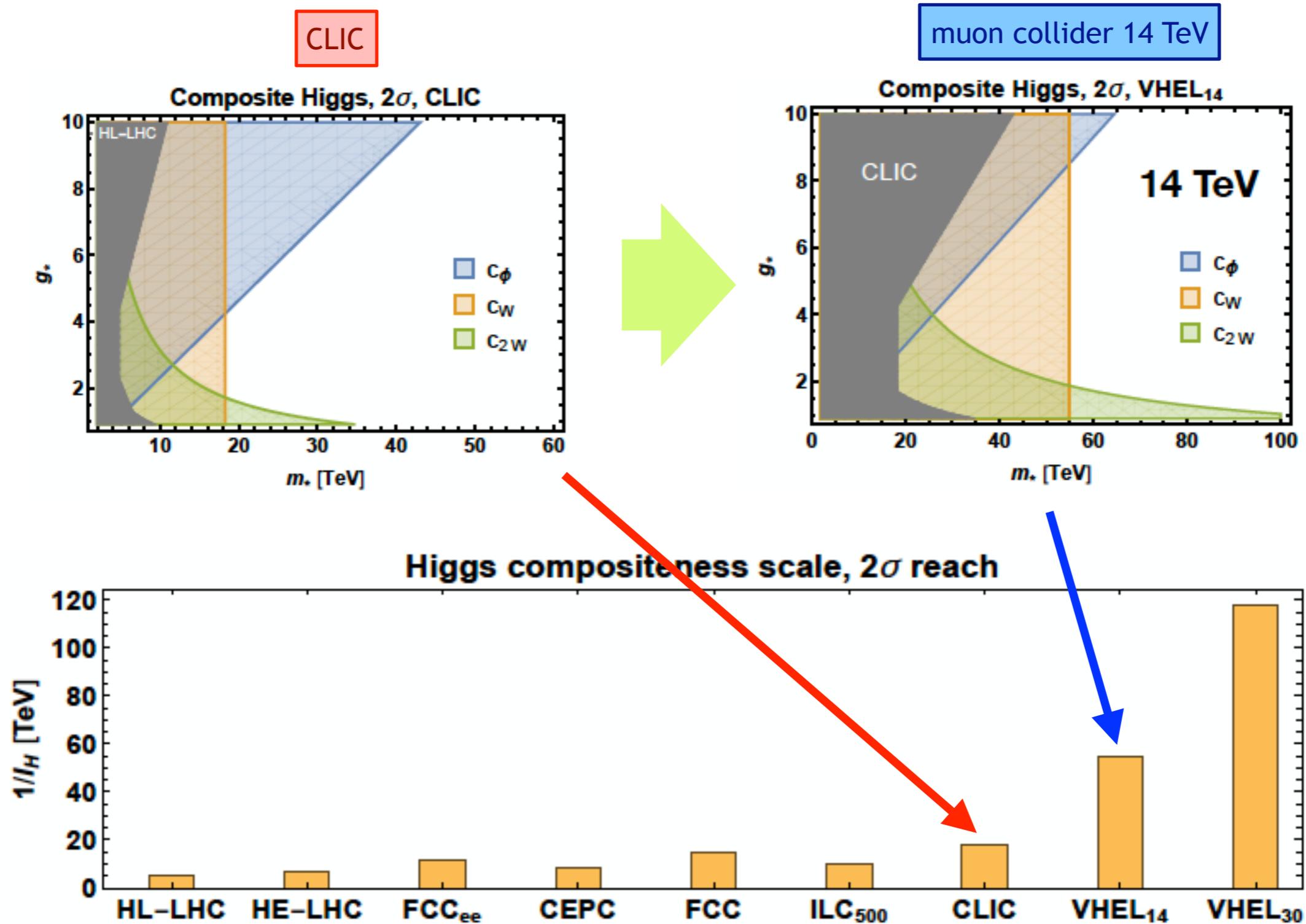
J.P. Delahaye, DPhP seminar

- muon colliders are the only reasonable way to imagine high-energy high-luminosity lepton colliders
- in terms of BSM physics reach, a 14 TeV muon collider compares with a 100 TeV pp collider... but much cleaner!



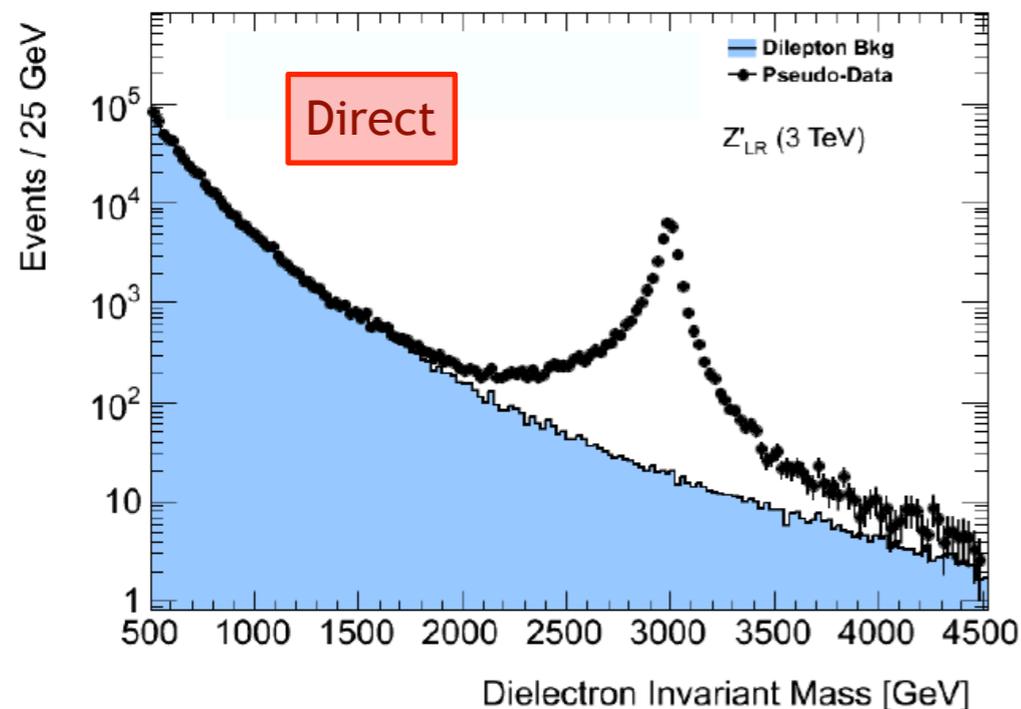
like CLIC... just ten times better!

Higgs Compositeness at Muon Collider



New Resonances, Particles or Forces

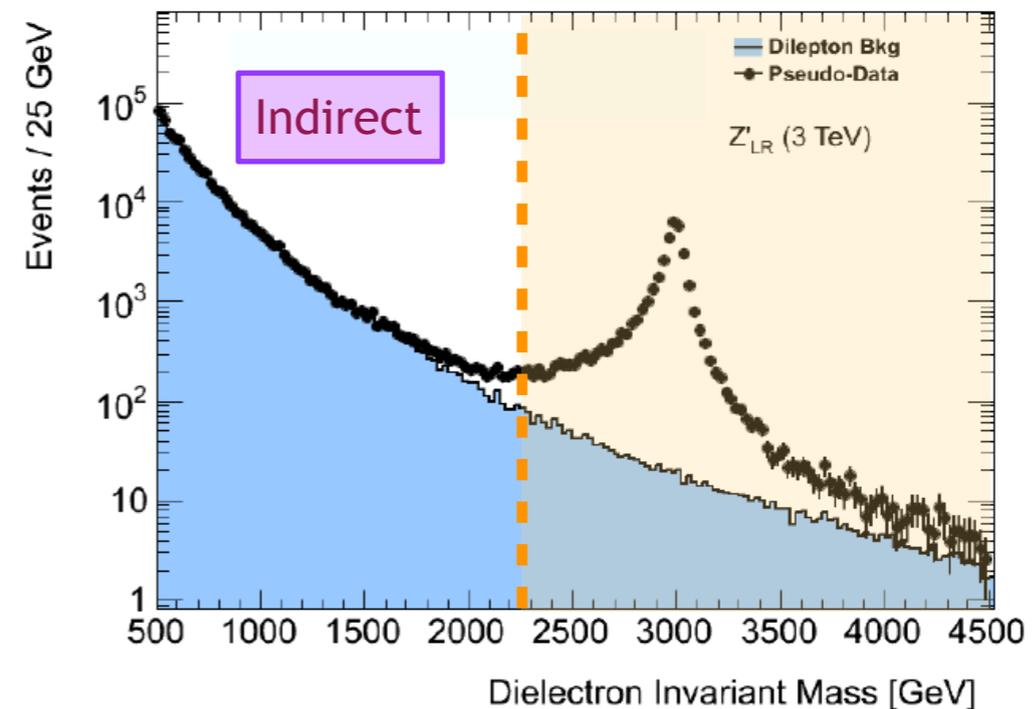
Question: is there a new very-short distance force on top of the electroweak and strong forces?



Z' resonance seen in the dilepton spectrum

Seeking a peak

- $M < \sqrt{s}$ for lepton colliders
- $M < 0.3-0.5 \sqrt{s}$ for hadron colliders, assuming “weak” couplings



no resonance seen but *deviations* wrt SM

Deviation in high-mass tails

- lepton colliders: sensitive to ratio $[\text{mass}/\text{coupling}] \gg \sqrt{s}$
- hadron colliders: only relevant if $g_{Z'} > g_{SM}$
 $\Rightarrow [\text{mass}/\text{coupling}] \gg 0.5\sqrt{s}$

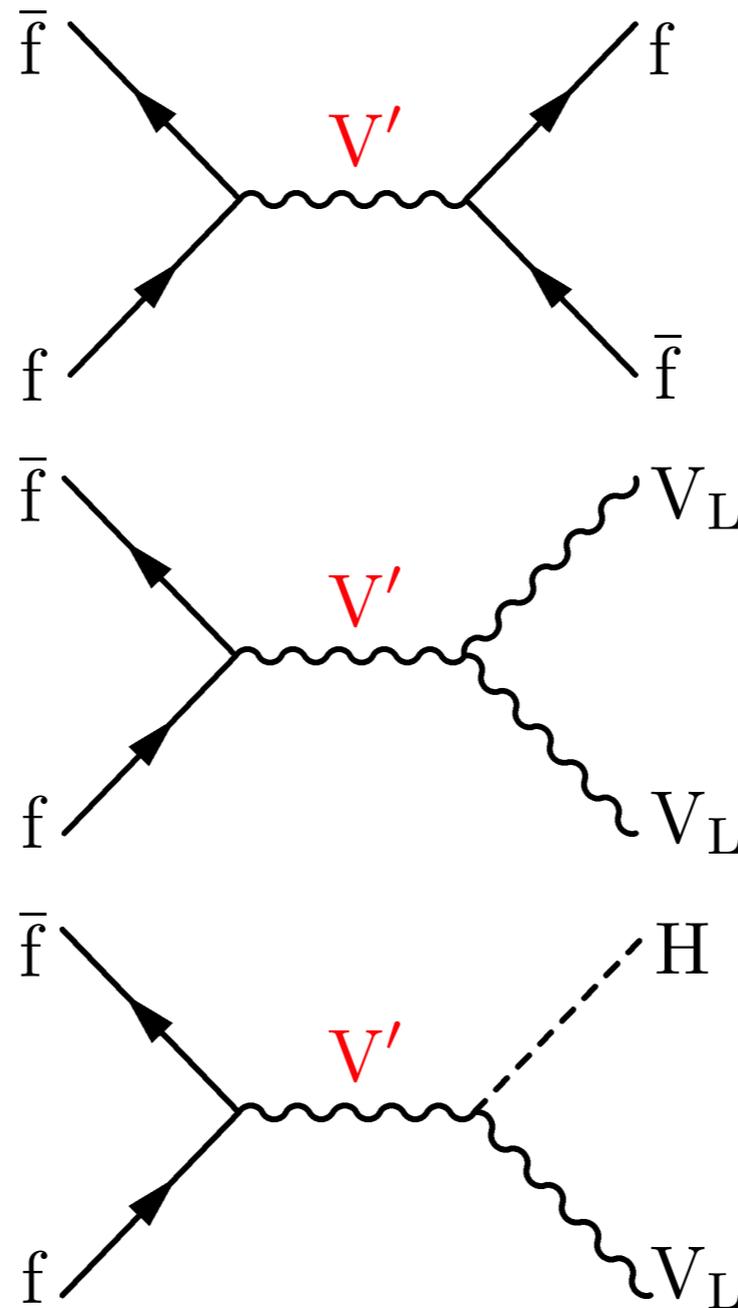
High-Energy Probes

example:
new W' , Z' vector bosons

at low-energy typically accounted for by dimension-6 operators interfering with SM processes

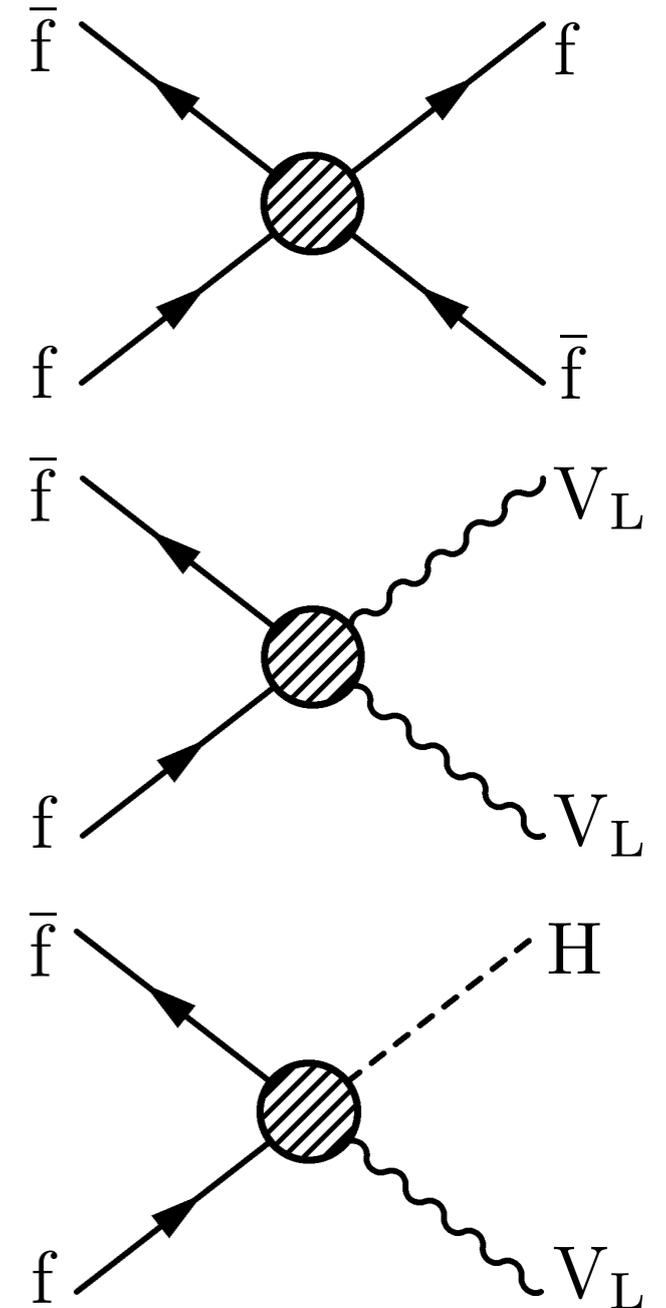
effects grow like center-of-mass energy squared of the collision

$$\propto \frac{s}{\Lambda^2}$$



new resonance in the corresponding spectra

$q^2 \ll M_{V'}$



excites the corresponding contact interaction operators in EFT fit

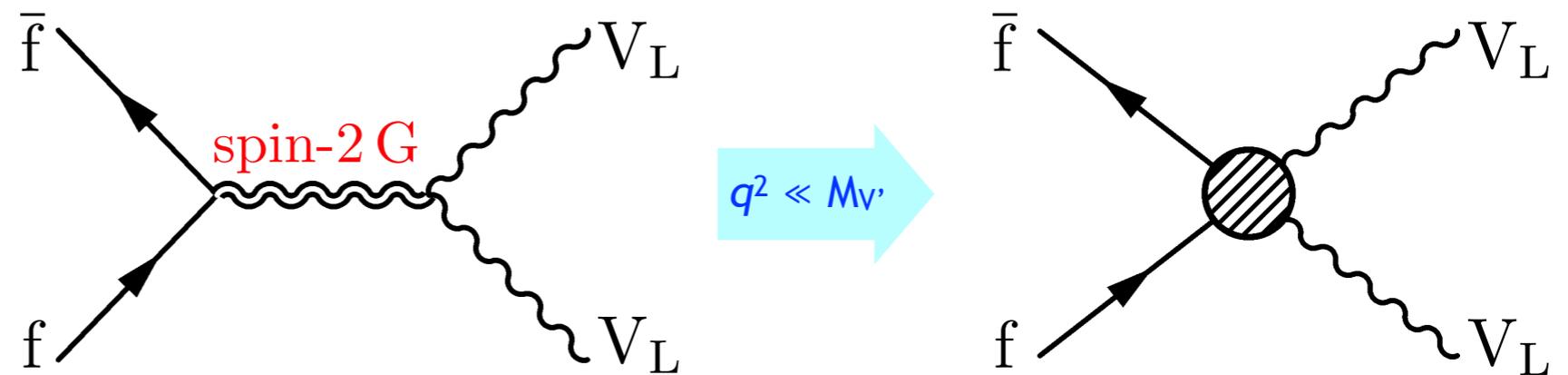
High-Energy Probes

➤ another example:
spin-2 graviton

➤ at low-energy:
dimension-8
operators

➤ grow with
center-of-mass
energy as

$$\propto \frac{s^2}{\Lambda^4}$$



due to huge growth with energy,
can show up at high-energy
colliders even in the absence of
deviations at low energy

Heavy Bosons Benchmark Models

➤ Sequential Standard Model (SSM)

- simple model as benchmark to compare sensitivities of experiments
- massive W' and Z' bosons with couplings to fermions identical to SM W and Z bosons

➤ Y-Universal Z' model

- another “simple” model with a new Heavy Dark Photon Z'
- massive U(1) gauge symmetry
- charges equal to SM hypercharge

$$\mathcal{L} = -\frac{1}{4g_{Z'}^2} \left(Z'_{\mu\nu} \right)^2 - \frac{1}{4g'^2} \left(B_{\mu\nu} - Z'_{\mu\nu} \right)^2 + \frac{M_{Z'}^2}{2g_{Z'}^2} \left(Z'_\mu \right)^2$$

B field:
generator of
U(1) $_{\gamma}$ SM
symmetry
(coupling = g')

Direct searches at high-energy

- most sensitive: dilepton spectra

Indirect searches at low energy:

- one single (“universal”) operator, known as Y-Universal

$$\frac{g_{Z'}^2}{2g'^2 M_{Z'}^2} \left(\partial^\mu B_{\mu\nu} \right)^2$$

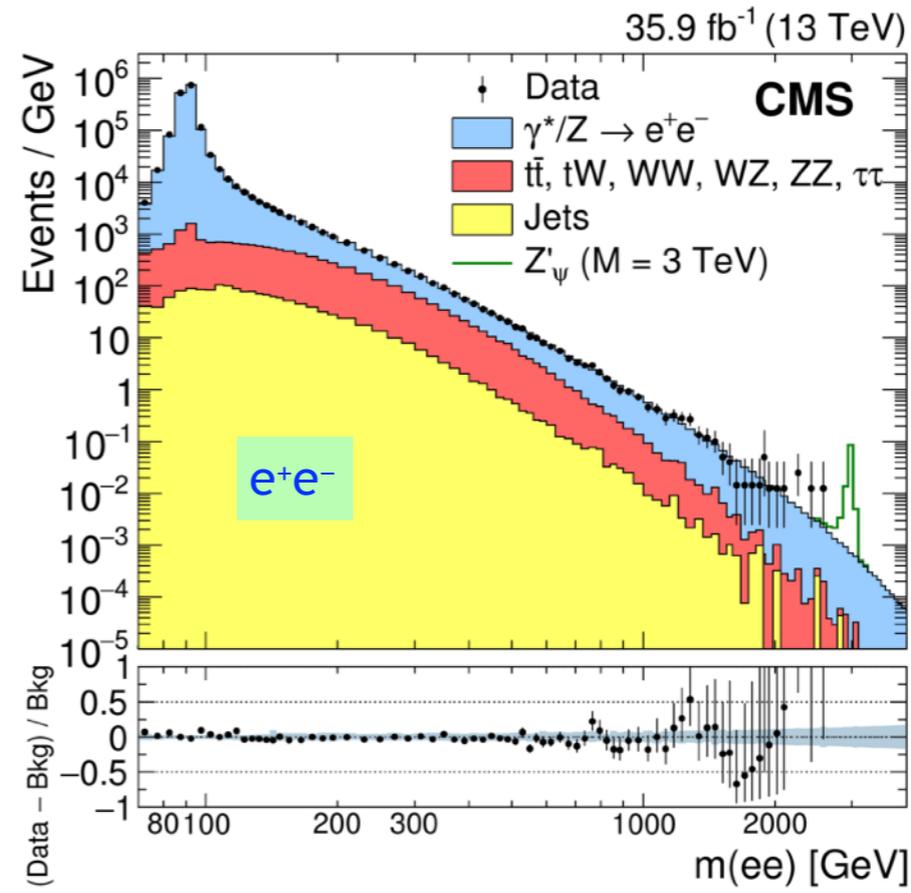
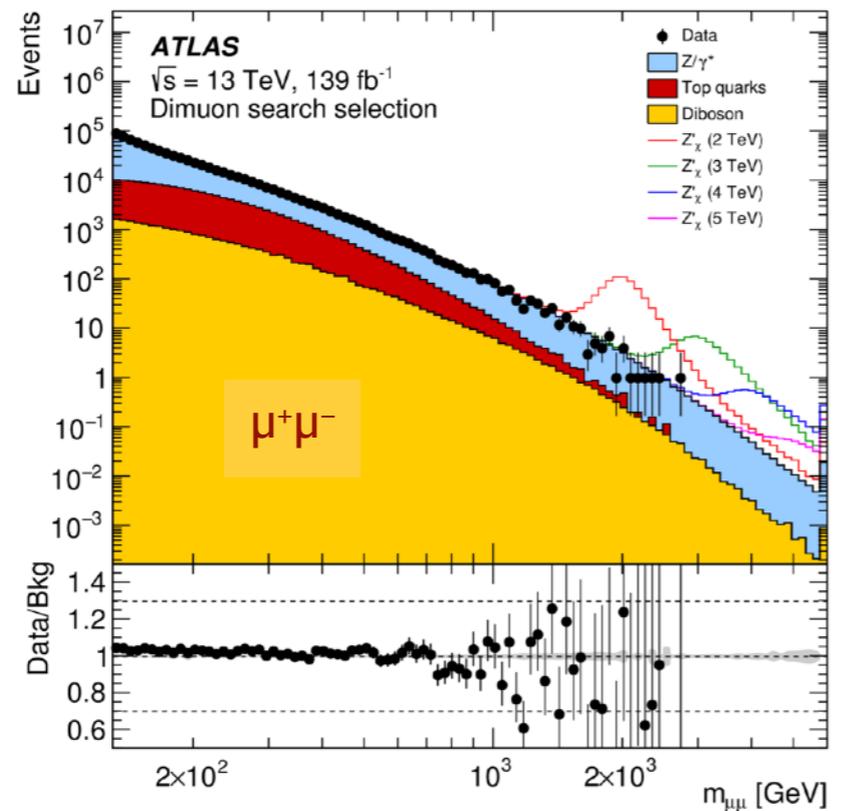
also known as c_{2B}

$$\frac{c_{2B}}{\Lambda^2} = \frac{g_{Z'}^2}{g'^4 M_{Z'}^2} = \frac{Y}{g'^2 m_W^2}$$

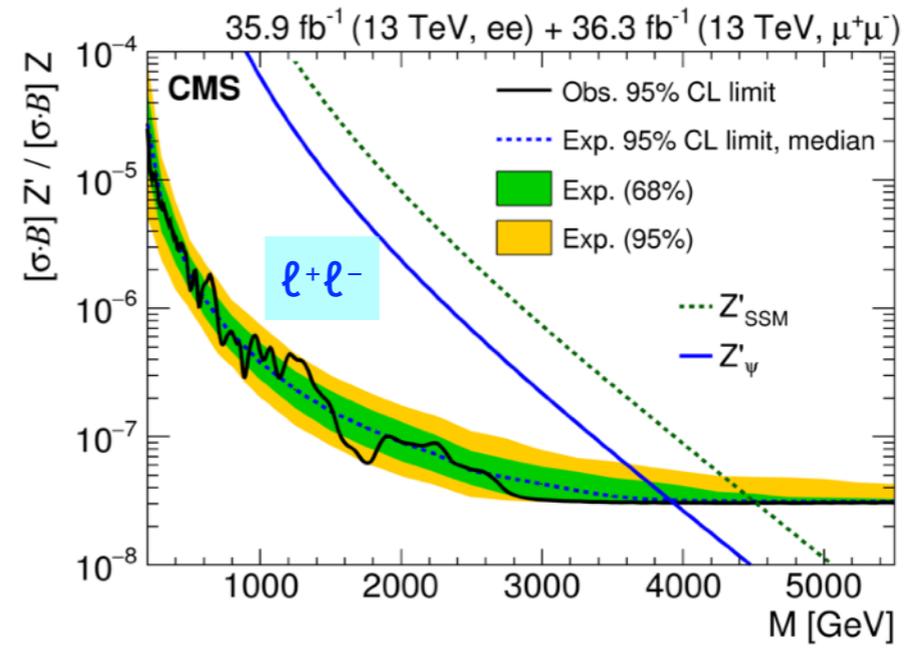
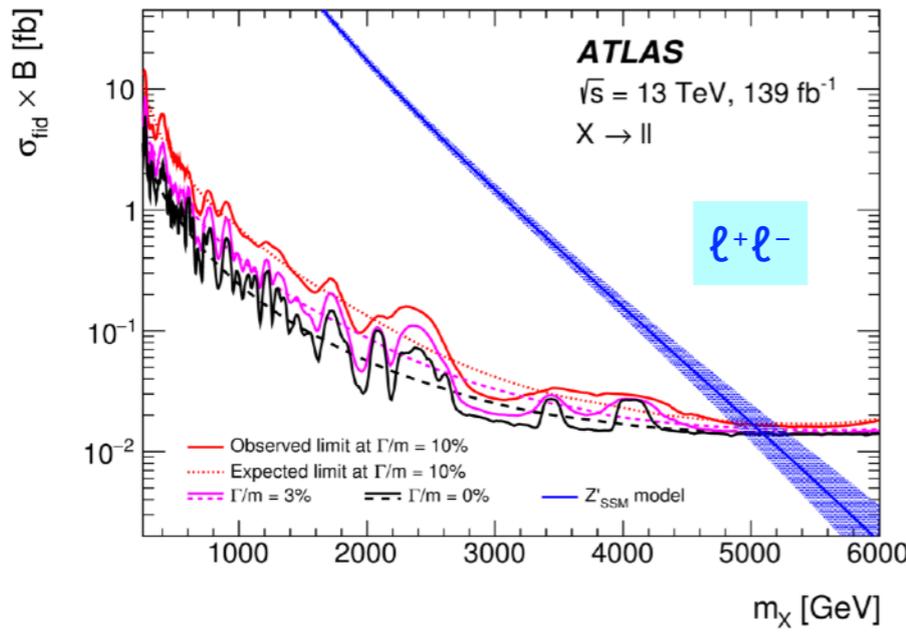
many other
models with
flavour
dependence,
etc.

Direct Search for Resonances at LHC

ATLAS Run-2
13 TeV
L = 139 fb⁻¹



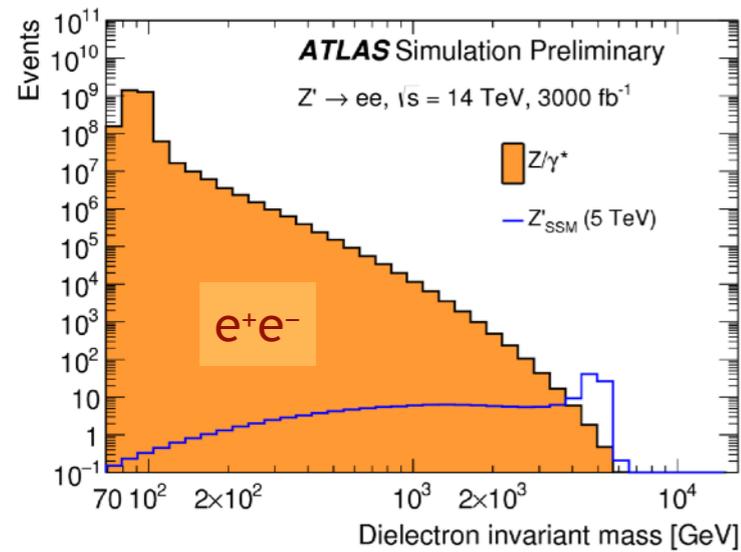
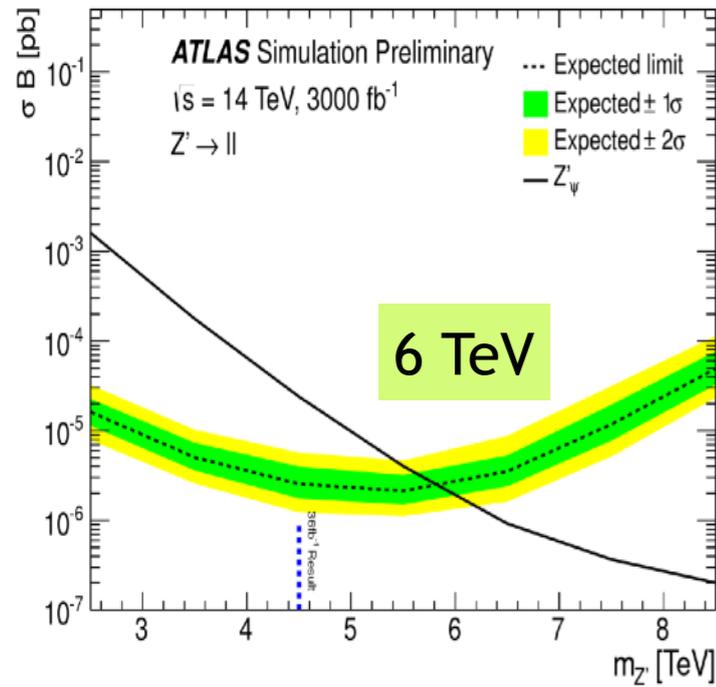
CMS 2016
13 TeV
L = 36 fb⁻¹



Typical
Run-2 limits
on Z' : 5 TeV

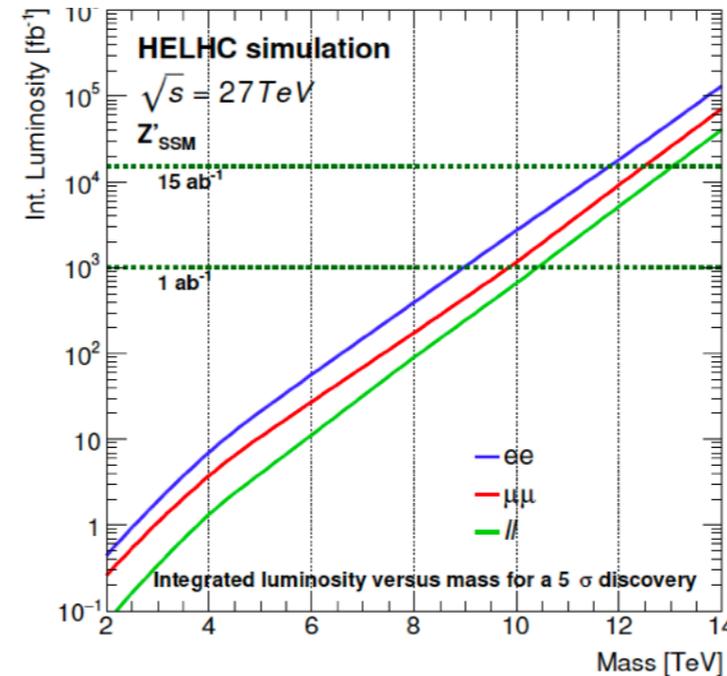
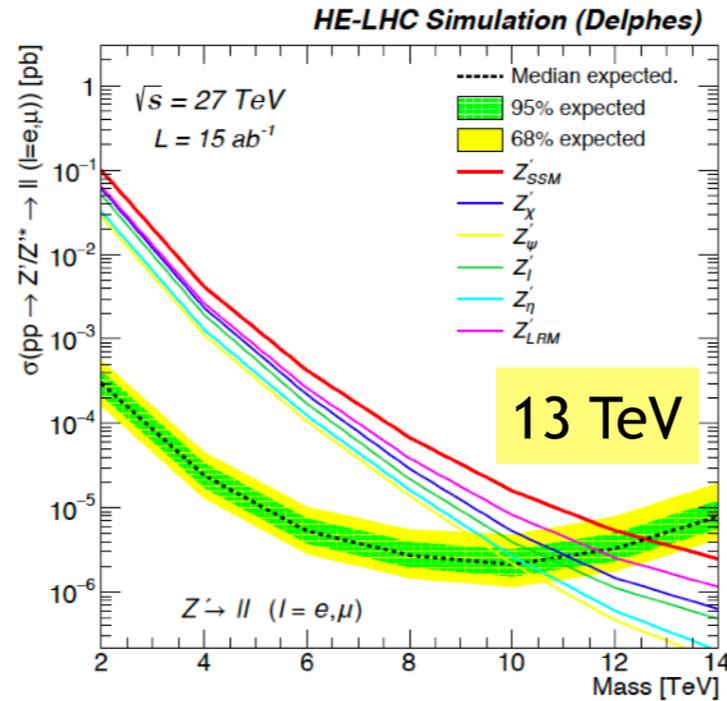
Direct Search for Resonances LHC → FCC

HL-LHC



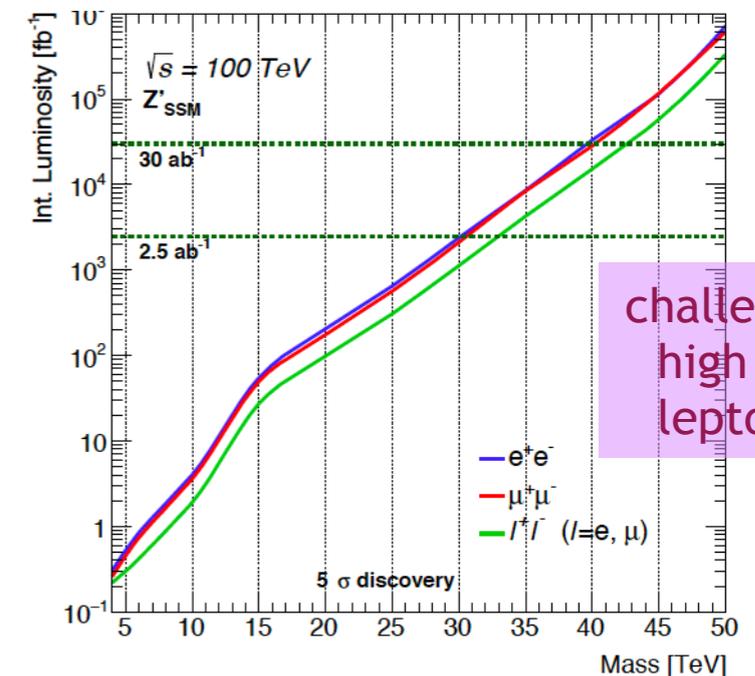
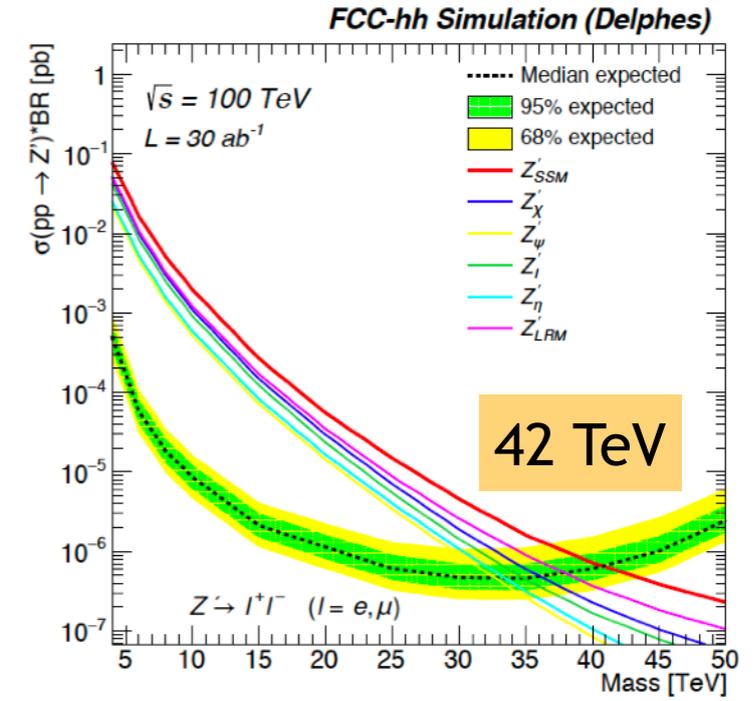
ATL-PHYS-PUB-2018-044

HE-LHC



HL/HE-LHC (2018) BSM
 arxiv:1812.07831

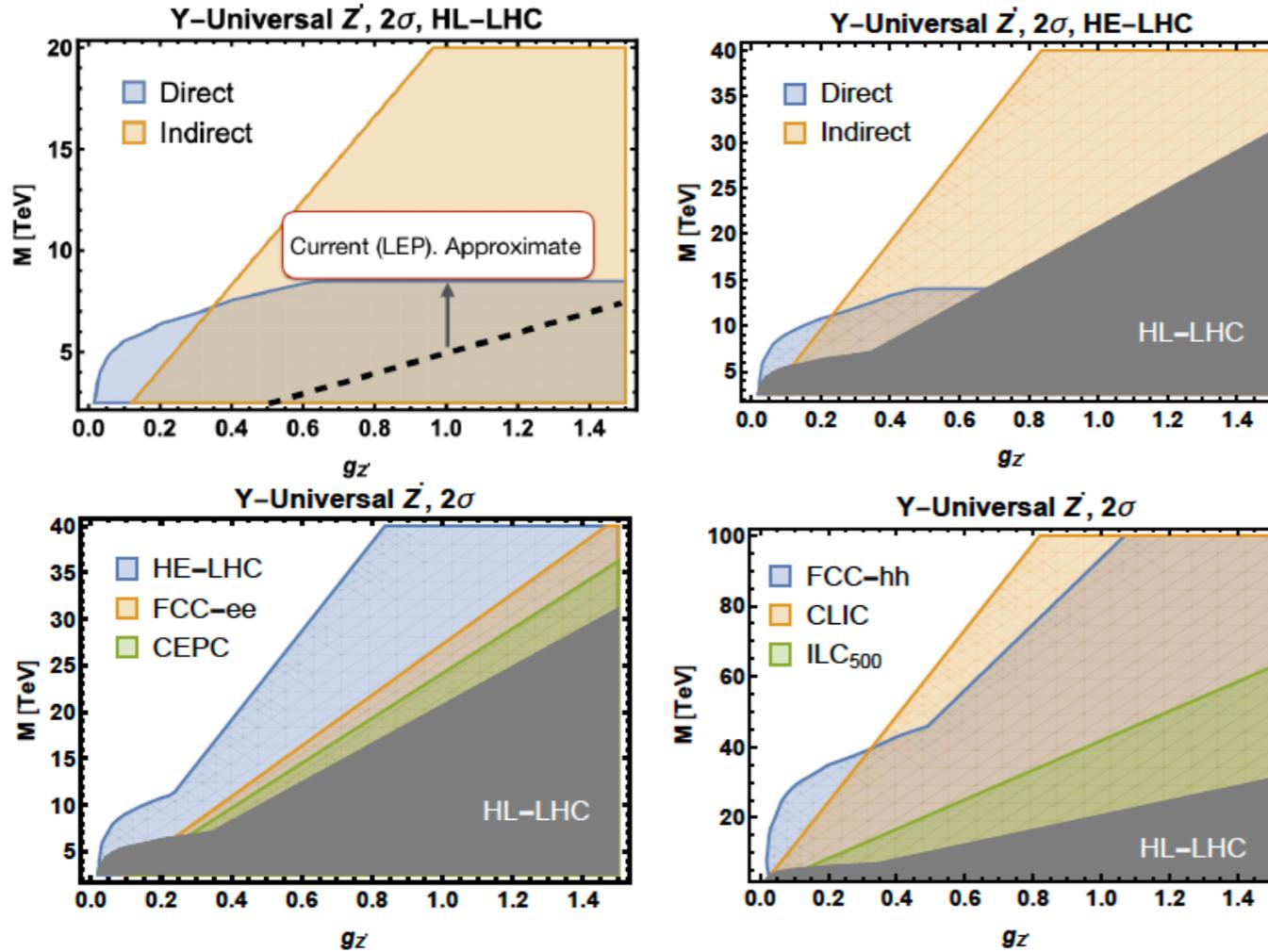
FCC-hh



challenge:
 high pT
 leptons

C. Helsen et al. (2019)
 arxiv:1902.11217

Z' Constraints at Future Colliders



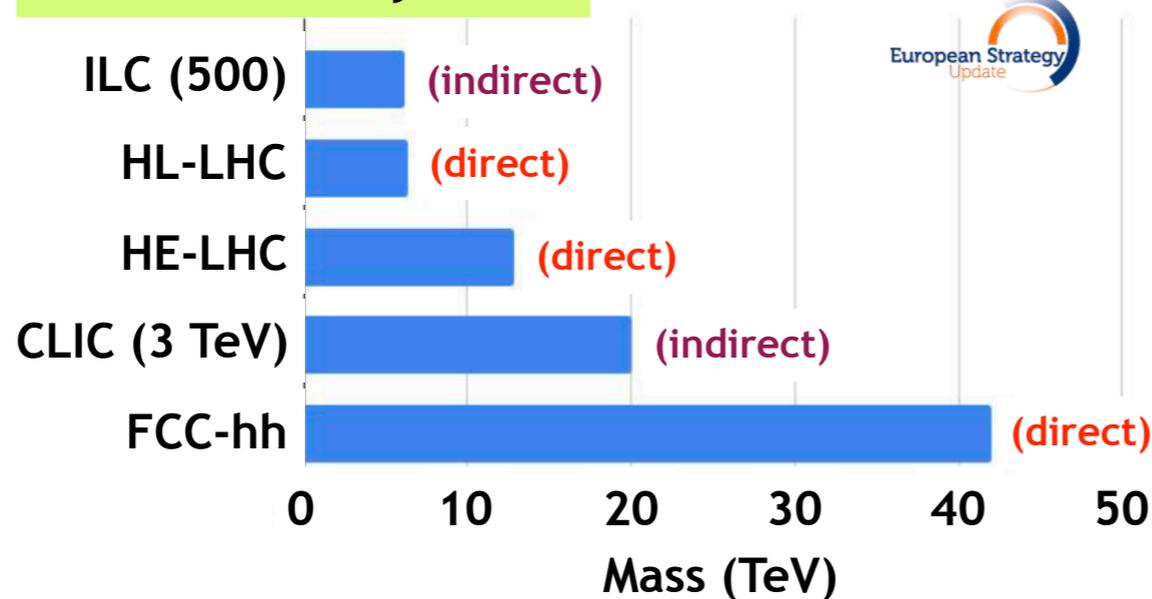
Best reach

- direct: FCC-hh
- indirect: CLIC 3-TeV

This is an example where high-energy makes a difference even for indirect constraints

- certain operators vary like v^2/Λ^2 , precision probes large Λ (e.g. anomalous Higgs couplings)
- other operators vary like Q^2/Λ^2 , those probe large Λ even with moderate precision

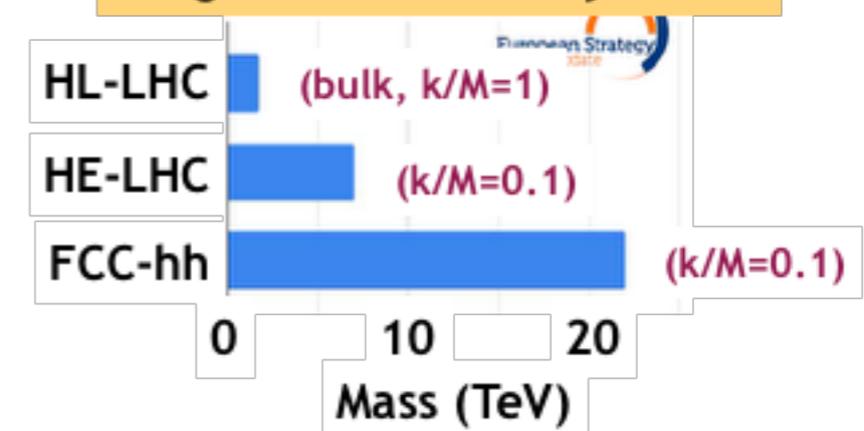
Z' SSM discovery reach



Preliminary, Granada 2019



RS graviton discovery reach

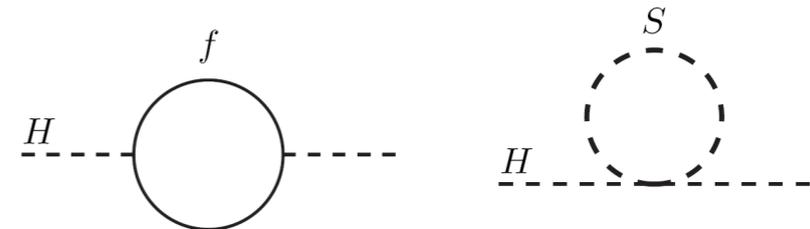


Supersymmetry

Supersymmetry: a family of **weakly-coupled** theoretical models that give solutions to the **Naturalness** problem, realise **unification at the GUT scale**, and, in certain realisations, provide a candidate of **Dark Matter** (R-parity conserving SUSY)

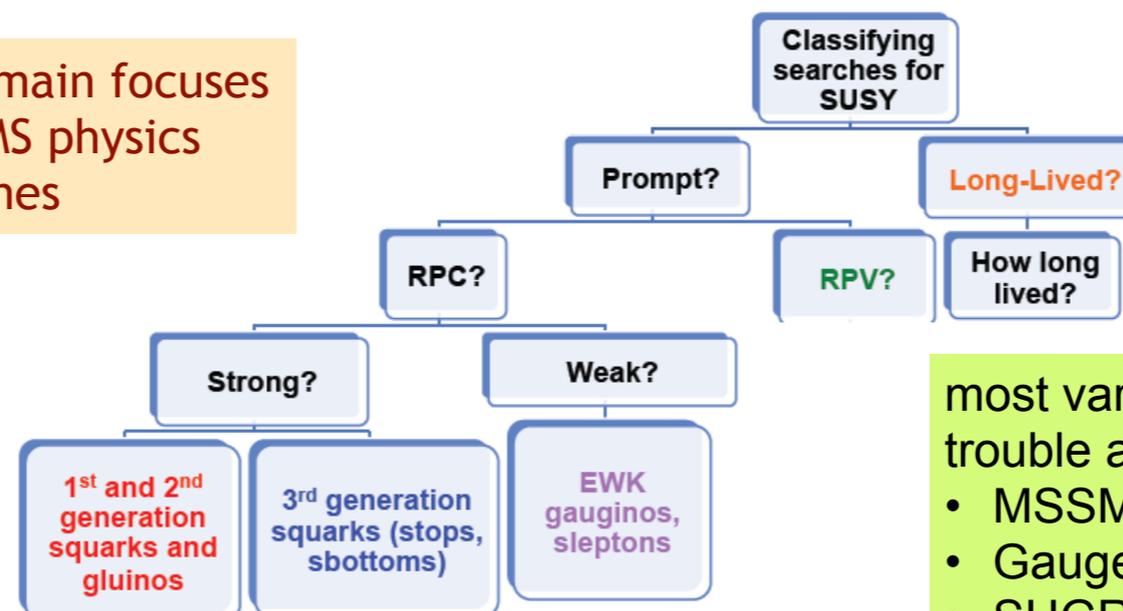
Spin 0	Spin 1/2	Spin 1	Spin 3/2	Spin 2
sleptons	leptons		gravitino	graviton
squarks	quarks			
Higgs	Higgsino			
	photino	photon		
	Zino	Z		
	Winos	W+ W-		
	gluinos	gluons		

SUSY is a fundamental space-time symmetry between fermions and bosons that regularises the Higgs boson mass



- each SM particle has a SUSY partner
- at which scale is SUSY broken?

SUSY is one of the main focuses of ATLAS and CMS physics programmes



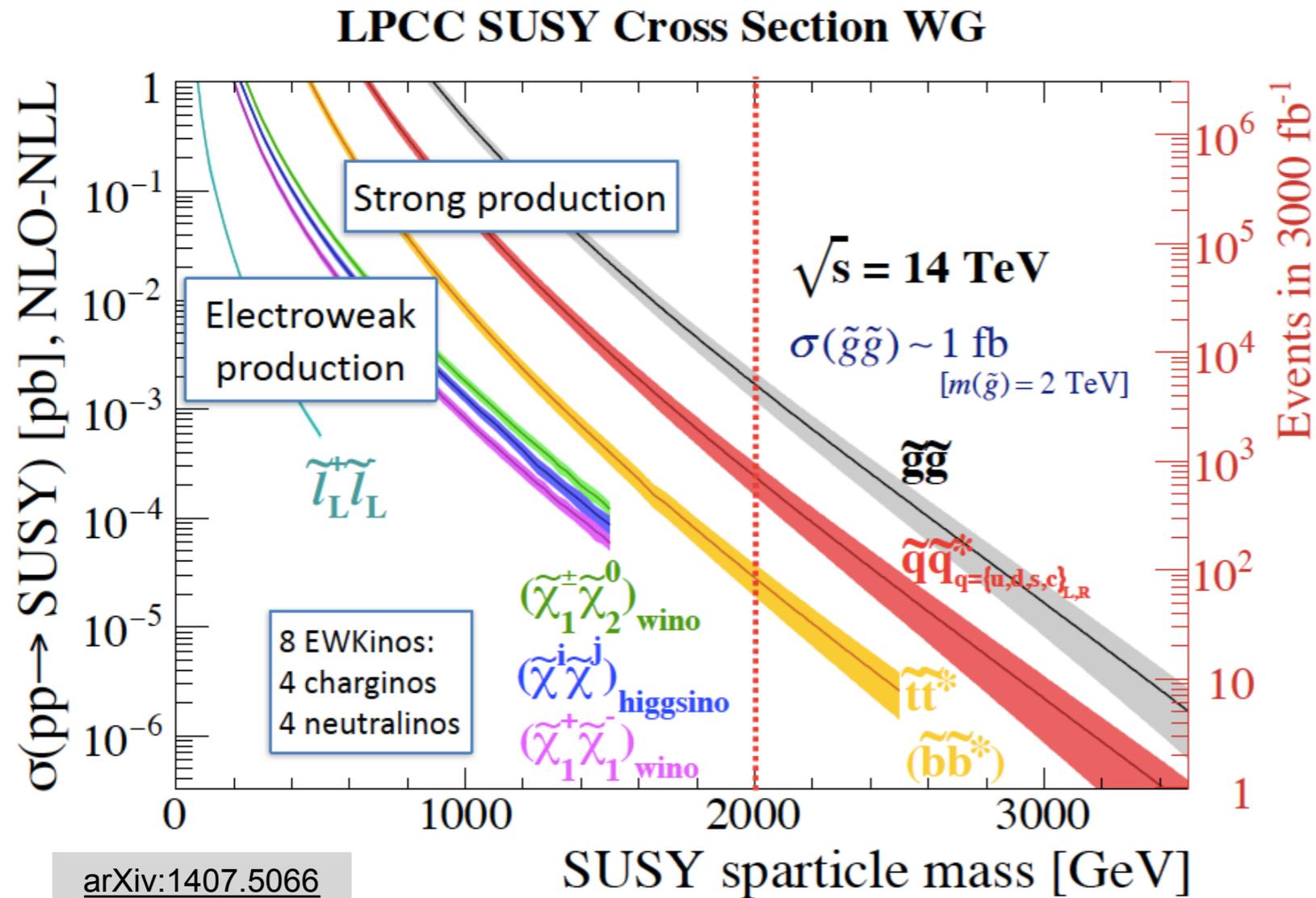
SUSY has an extended Higgs sector, for instance 2HDM as in MSSM

- CP even: h and H
- CP odd: A
- charged: H⁺ and H⁻

most variants of SUSY are in trouble after LHC Run-1 and-2

- MSSM, NMSSM, ...
- Gauge Mediation
- SUGRA
- Stealth SUSY,
- etc.

SUSY Production



Strong production

- gluinos
- squarks of 1st and 2d generations
- top and bottom squarks

Electroweak production

- neutralinos
- charginos
- sleptons

arXiv:1407.5066

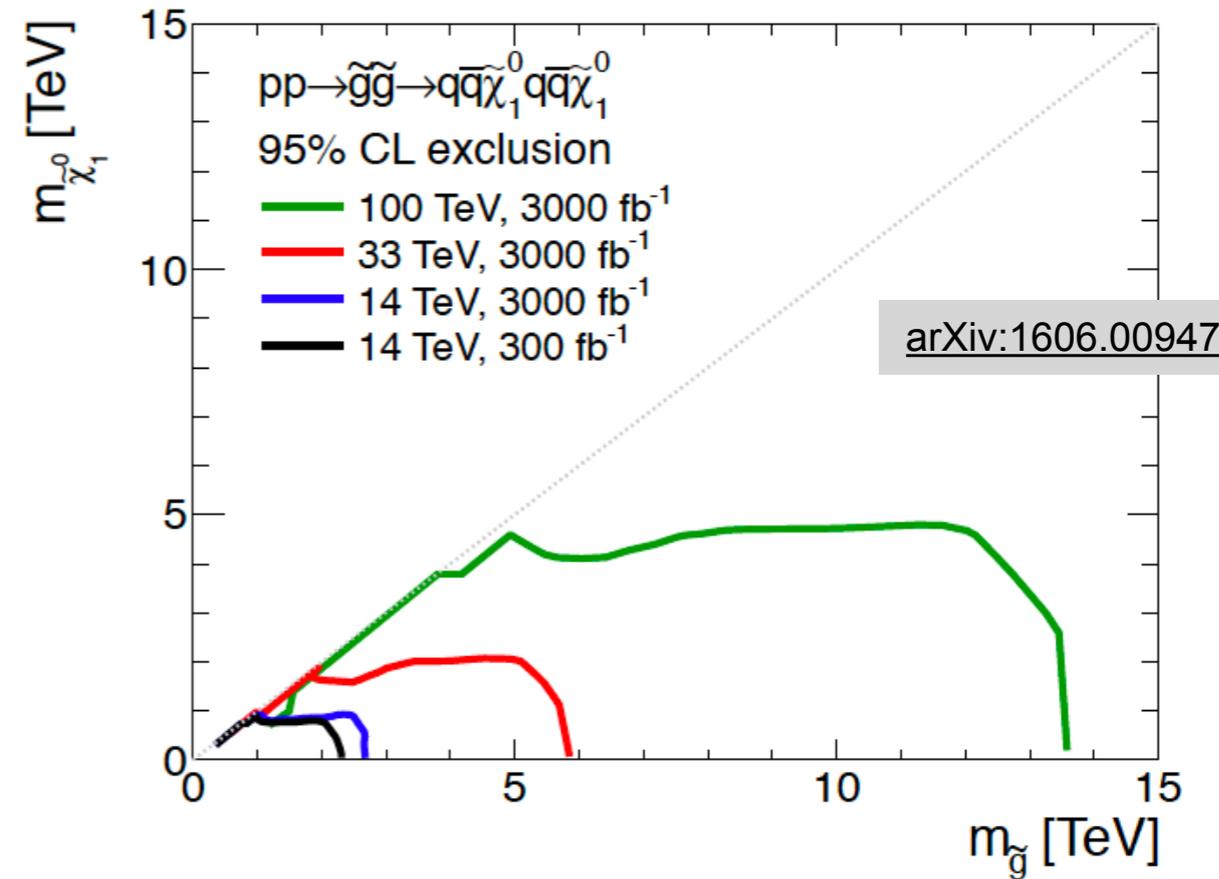
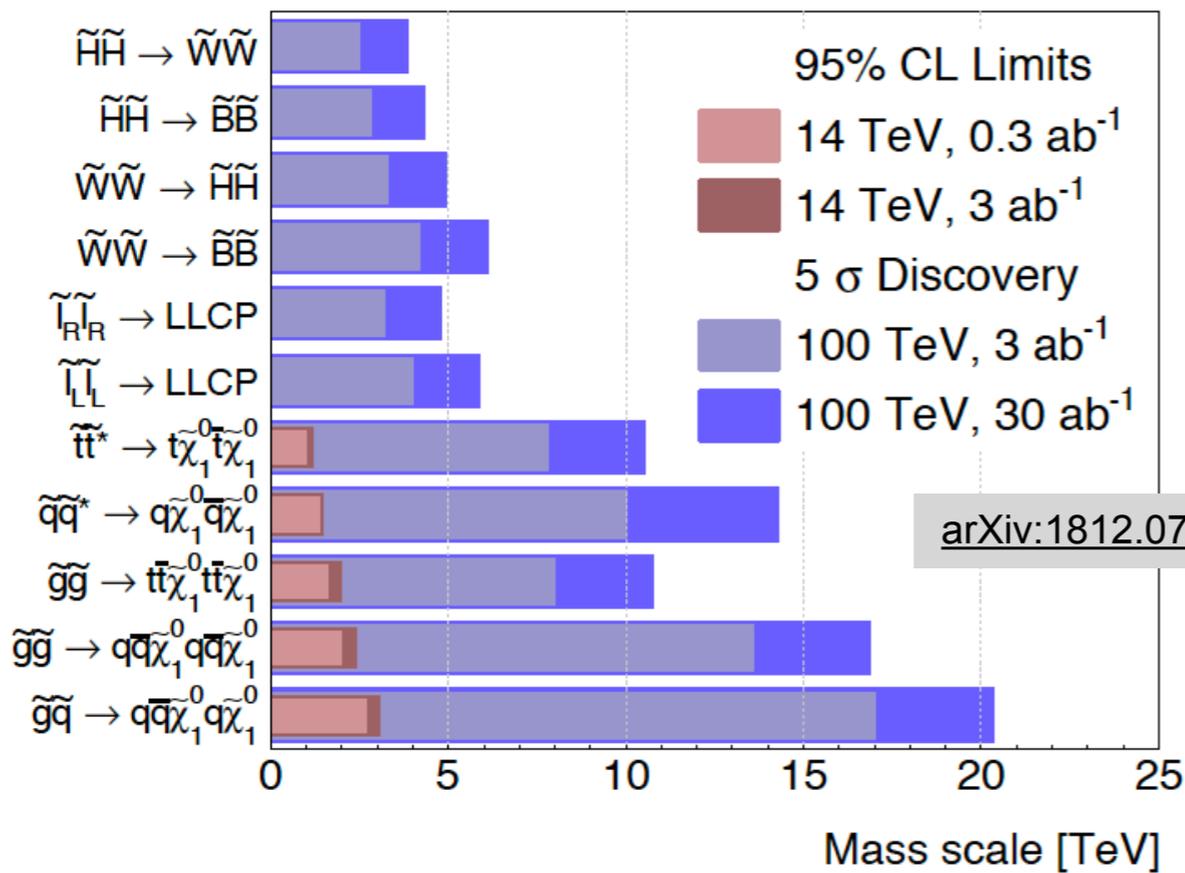
SUSY signatures

- R-parity conserving
 - large transverse missing energy
- R-parity violating
 - feebly interacting or non-prompt

If you believe in SUSY:

- direct searches are more powerful than precision constraints because SUSY is weakly coupled
- lepton colliders will provide limited improvement
- high-energy proton colliders bring significant improvement in direct coverage

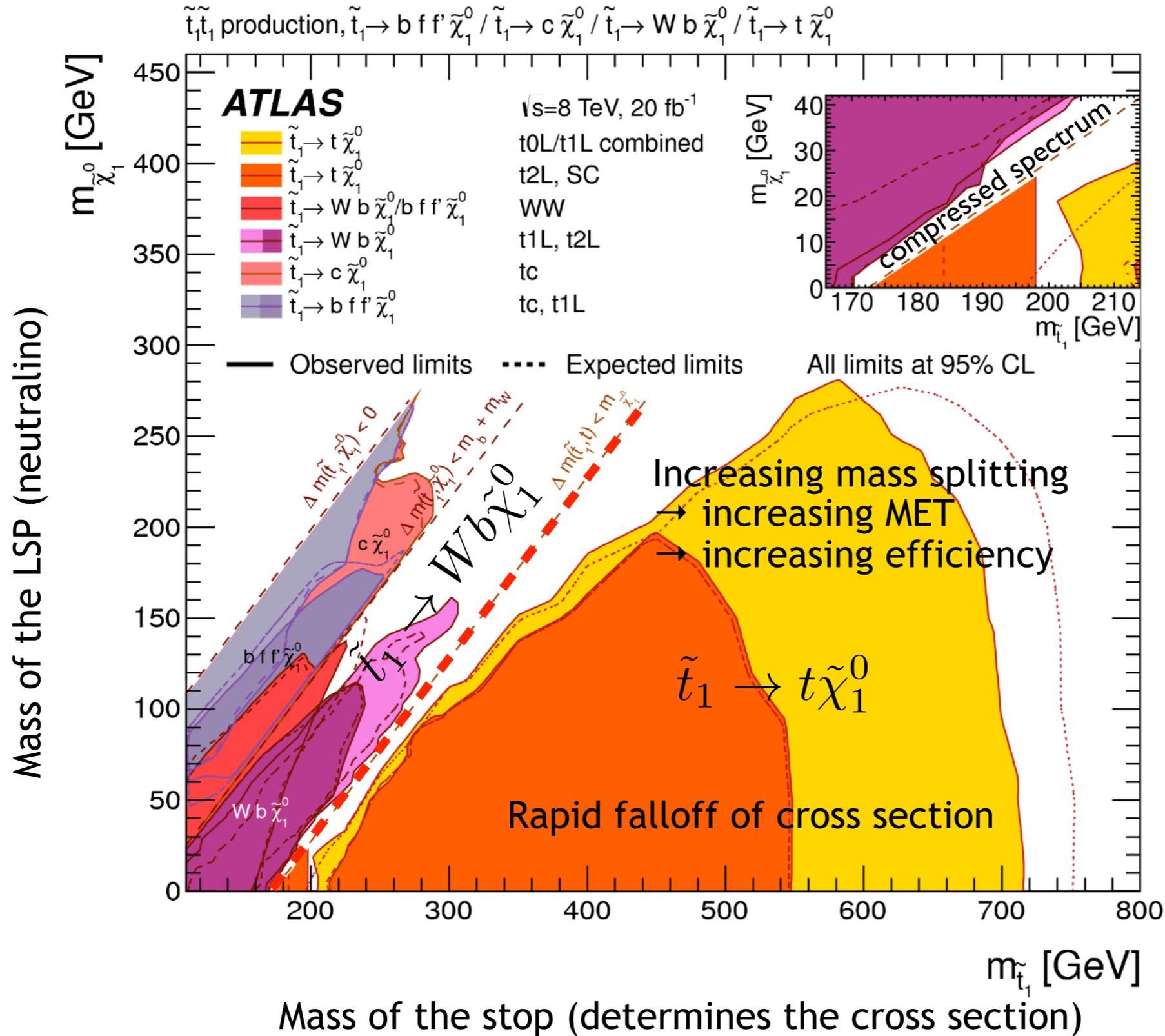
Glauino Searches



	Model	$\int \mathcal{L} dt [\text{ab}^{-1}]$	\sqrt{s} [TeV]	Mass limit (95% CL exclusion)	Conditions
HL-LHC	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\bar{q}\tilde{\chi}_1^0$	3	14	3.2 TeV	$m(\tilde{\chi}_1^0)=0$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\bar{q}\tilde{\chi}_1^0$	3	14	1.5 TeV	$m(\tilde{g}) \sim m(\tilde{\chi}_1^0)+10 \text{ GeV}$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$	3	14	2.5 TeV	$m(\tilde{\chi}_1^0)=0$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$	3	14	2.6 TeV	$m(\tilde{\chi}_1^0)=500 \text{ GeV}$
HE-LHC	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\bar{q}\tilde{\chi}_1^0$	15	27	5.7 TeV	$m(\tilde{\chi}_1^0)=0$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\bar{q}\tilde{\chi}_1^0$	15	27	2.6 TeV	$m(\tilde{g}) \sim m(\tilde{\chi}_1^0)+10 \text{ GeV}$
	NUHM2, $\tilde{g} \rightarrow t\bar{t}$	15	27	5.9 TeV	$m(\tilde{\chi}_1^0)=0$
FCC-hh	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\bar{q}\tilde{\chi}_1^0$	30	100	17.0 TeV	$m(\tilde{\chi}_1^0)=0$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\bar{q}\tilde{\chi}_1^0$	30	100	7.5 TeV	$m(\tilde{g}) \sim m(\tilde{\chi}_1^0)+10 \text{ GeV}$ (*)
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$	30	100	11.0 TeV	$m(\tilde{\chi}_1^0)=0$

FCC-hh max reach: 15-20 TeV

Example of Stop Exclusion Plot

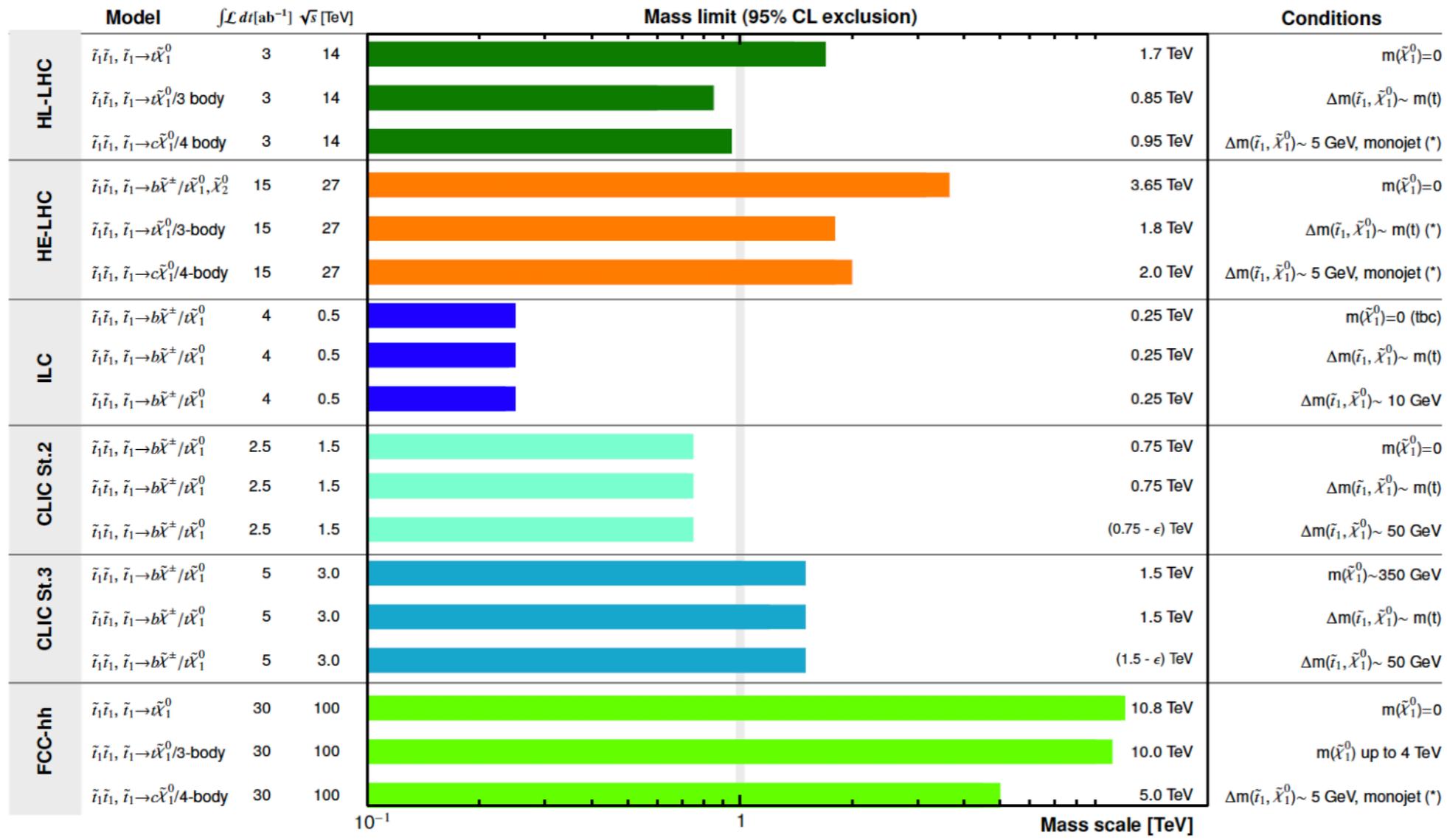


Limits extend to $m(\text{stop}) \approx 700$ GeV for low-mass LSP

No limit on $m(\text{stop})$ for $m(\text{LSP}) > 275$ GeV

Stop Searches

discovery potential



HL/HE-LHC
up to 1.4/3.2 TeV

ILC/CLIC
up to $\sqrt{s}/2$

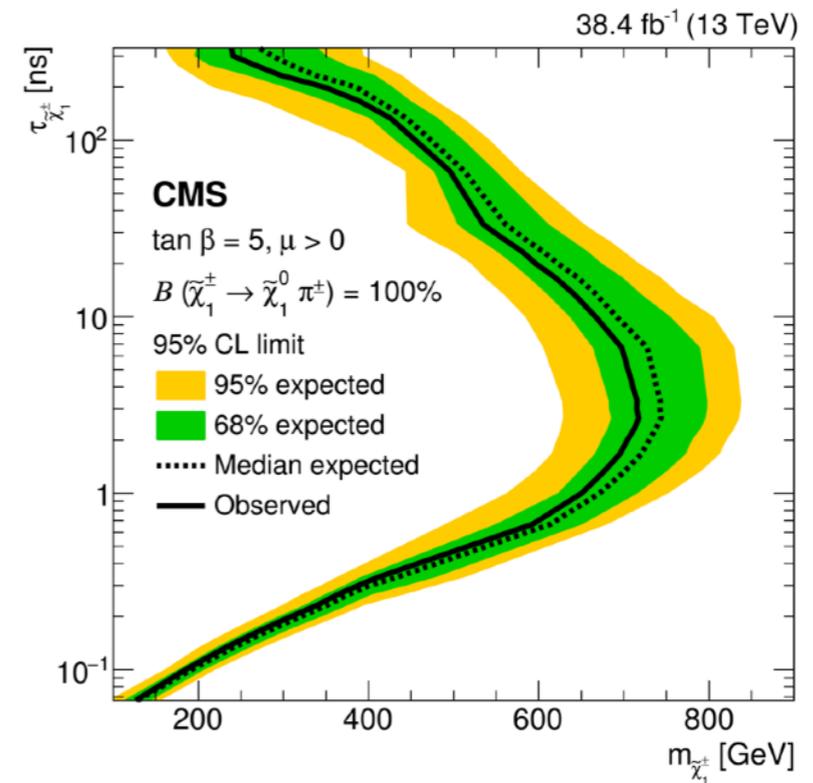
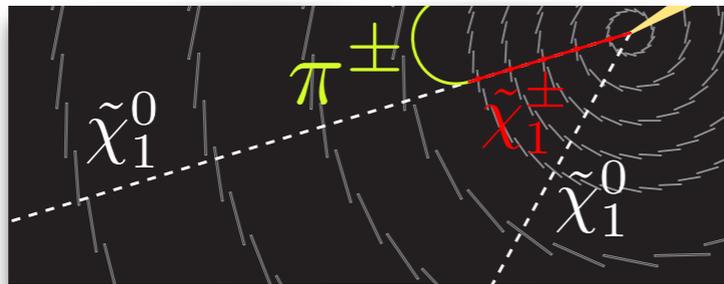
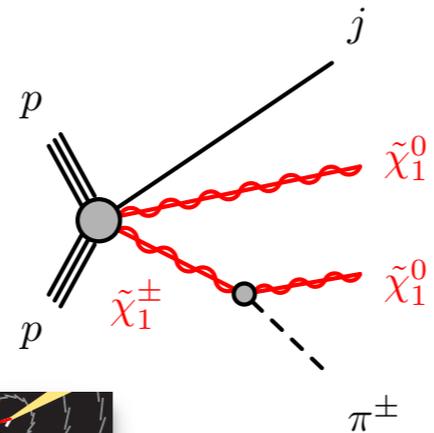
FCC-hh
up to 8 TeV

SUSY can/will never die... but if no stop is found at future colliders, SUSY cannot serve its main purpose, which is to solve the Naturalness problem

Disappearing Tracks

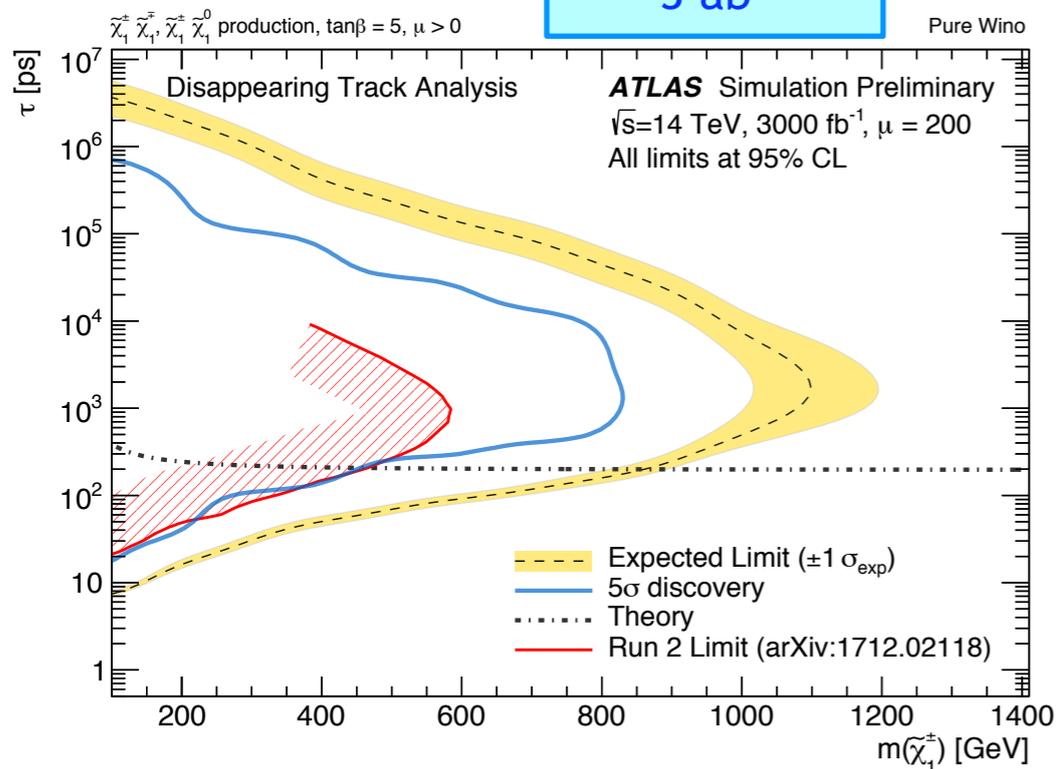
Example: Wino-like DM

$$\chi = \begin{pmatrix} \chi^+ \\ \chi^0 \\ \chi^- \end{pmatrix} \quad \text{dark matter}$$

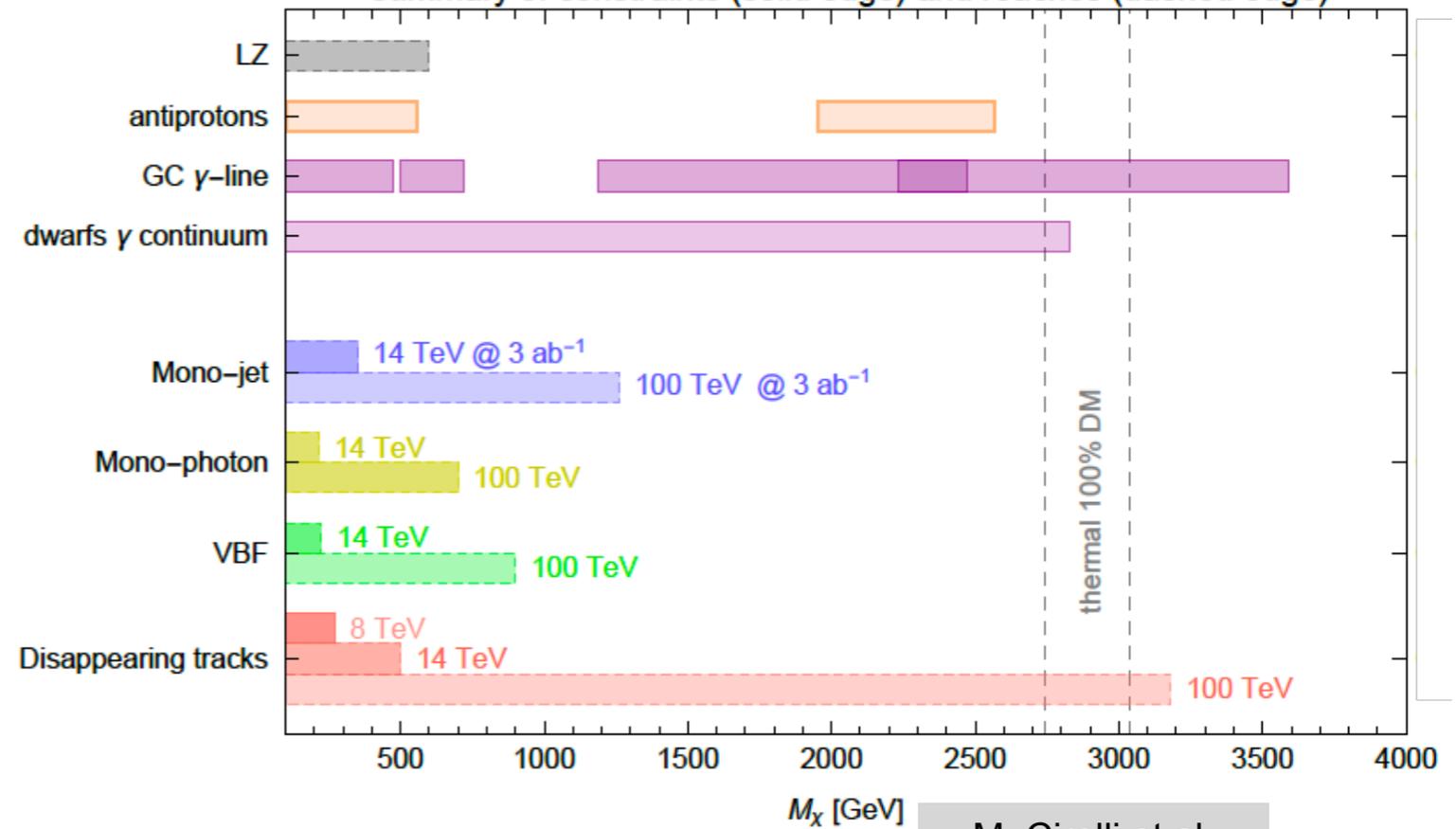


CMS 2016
13 TeV
L = 36 fb⁻¹

ATLAS HL-LHC
3 ab⁻¹



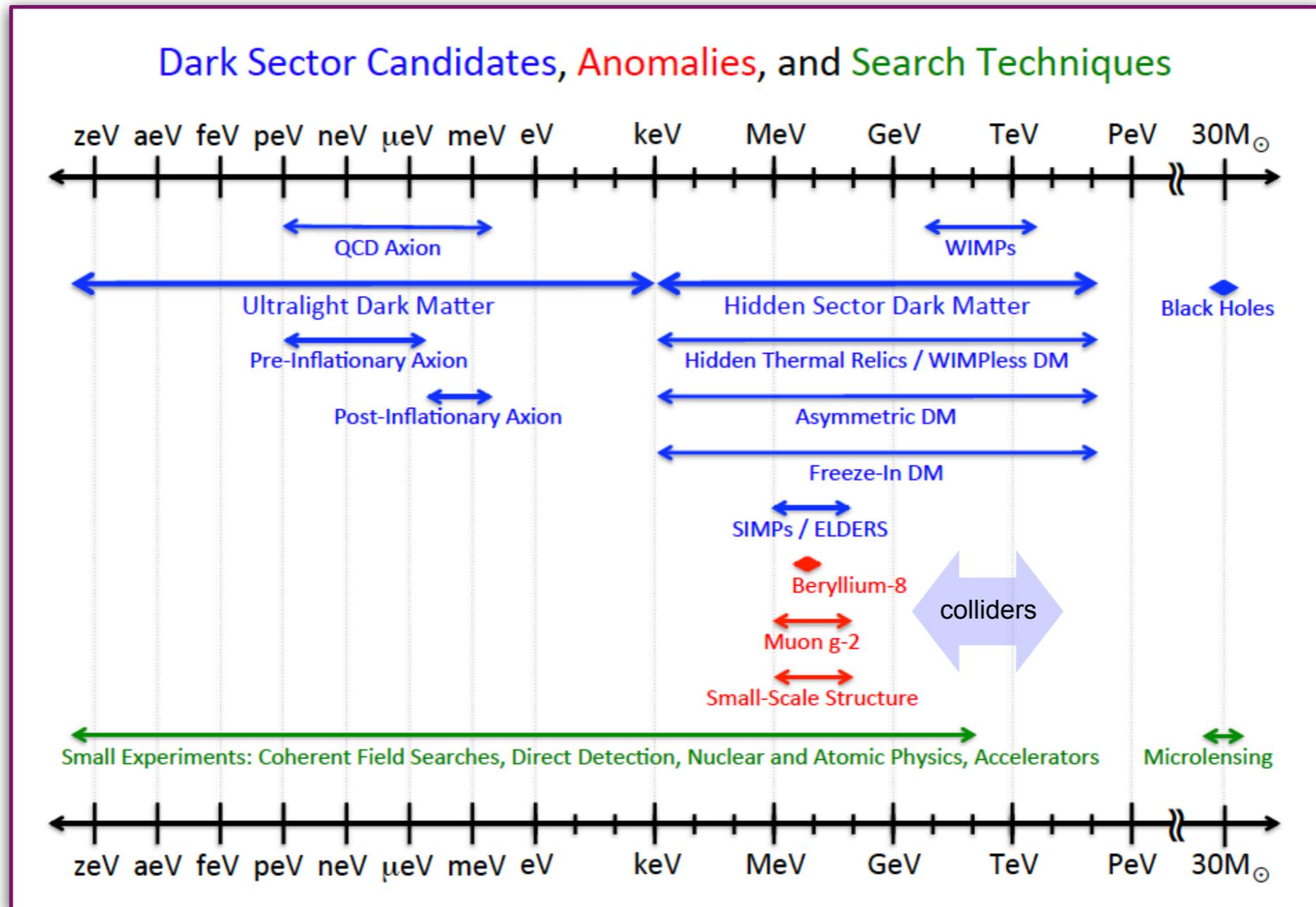
Wino-like (minimal 3plet) Dark Matter:
summary of constraints (solid edge) and reaches (dashed edge)



M. Cirelli et al.
arXiv:1407.7058

Dark Matter: The Log Crisis

Our ignorance of Dark Matter is logarithmic!



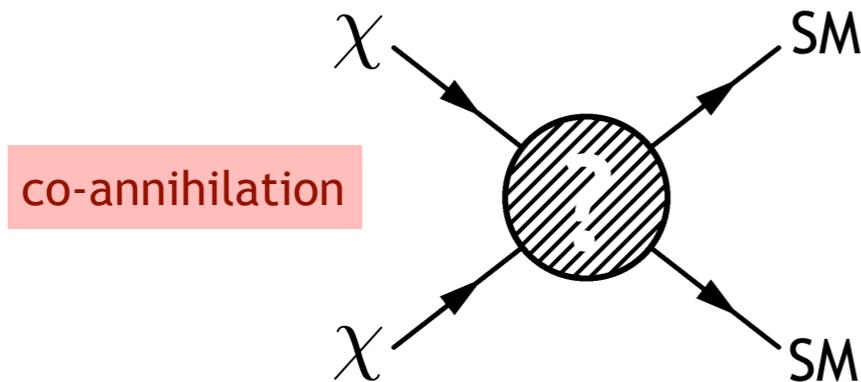
Mass ranges for dark matter and mediator particle candidates, experimental anomalies, and search techniques

Thermal WIMPs

WIMP = weakly-interacting massive particle

WIMP scenario: the relic density of Dark Matter is set by non-relativistic annihilation (freeze-out) to SM particles, either through

- a SM portal (e.g., the Higgs boson)
- a new mediator (consider scalar, pseudoscalar, vector or axialvector mediators)



co-annihilation

the WIMP miracle?

g_{eff} = effective coupling strength

$$\langle \sigma v \rangle \sim g_{\text{eff}}^4 / M_{\text{DM}}^2$$

arXiv:1606.00947

From cosmology: thermal freeze-out

$$\Omega_{\text{DM}} h^2 \sim 0.12 \times \left(\frac{M_{\text{DM}}}{2 \text{ TeV}} \right)^2 \left(\frac{0.3}{g_{\text{eff}}} \right)^4$$

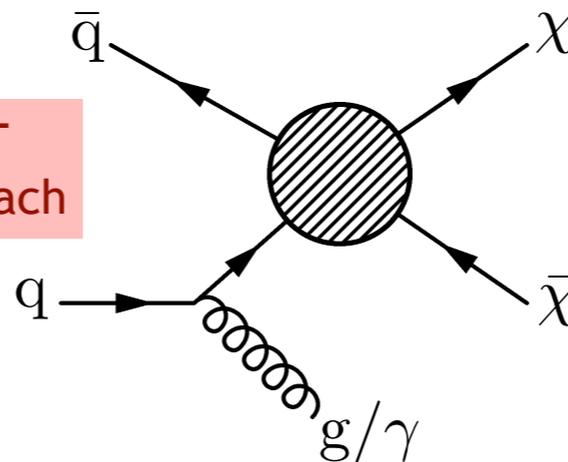
WIMPs can account for relic abundance for masses in the few GeV to few TeV range

At colliders: the recoil method

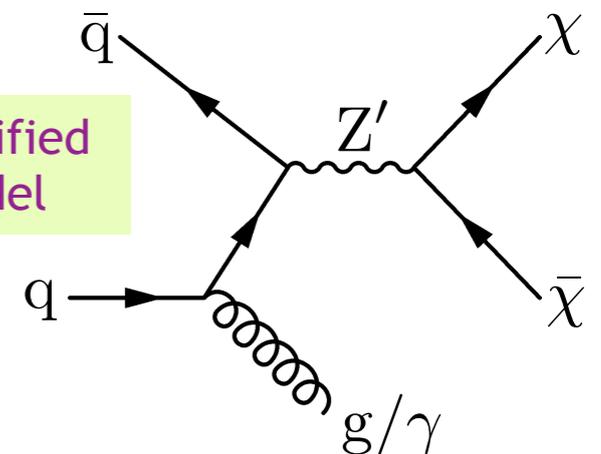
- the DM particles are pair-produced with a radiated SM particle
- main background: $Z \rightarrow \nu\bar{\nu}$

and 2 approaches

EFT approach



Simplified Model

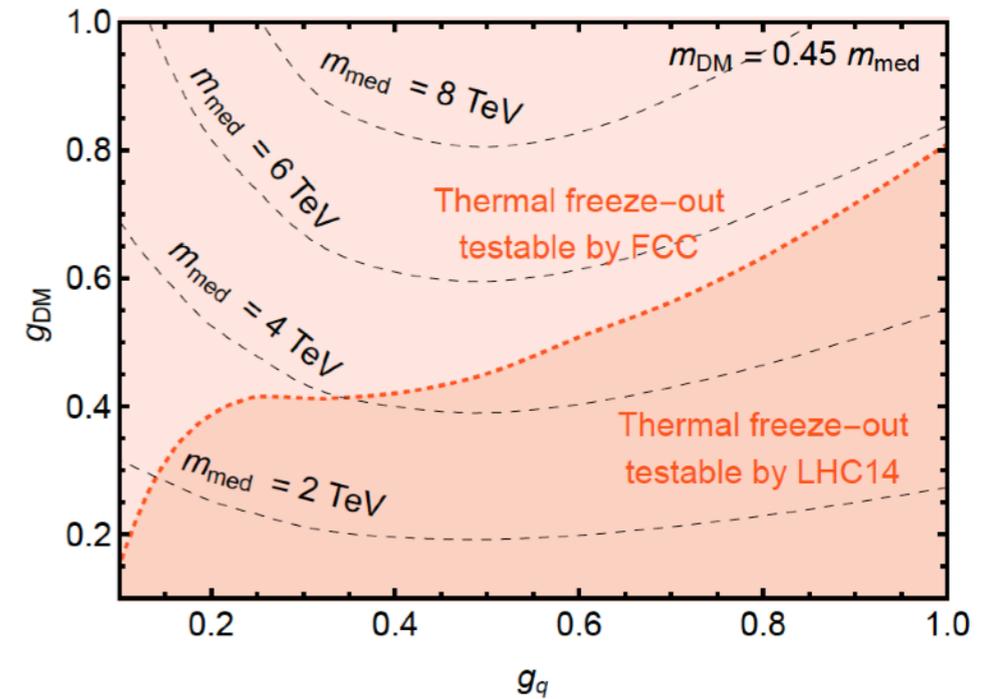


A Simplified Dark Matter Model

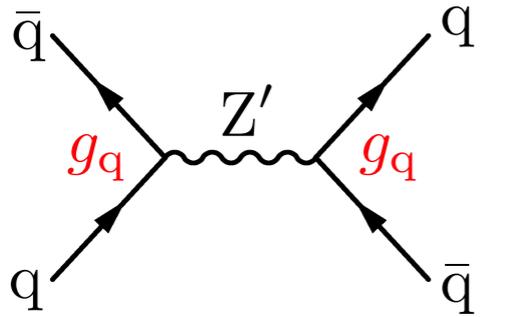
Consider a simple model where Dark Matter interacts via a new Z' boson

- the WIMP is a Majorana fermion
- the mediator Z' has only vector couplings to the SM quarks

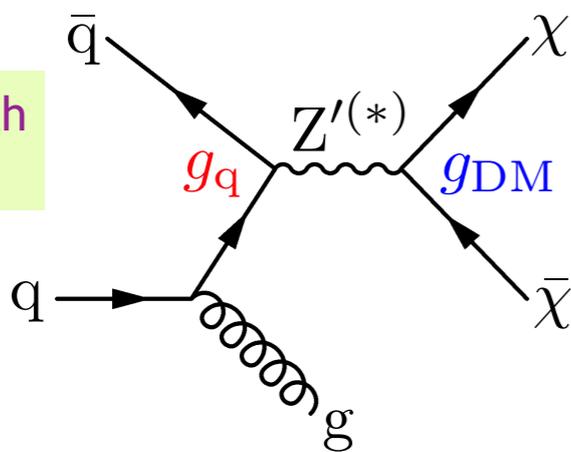
no constraints from EWPO & dilepton resonance searches
only relevant if no direct DM detection



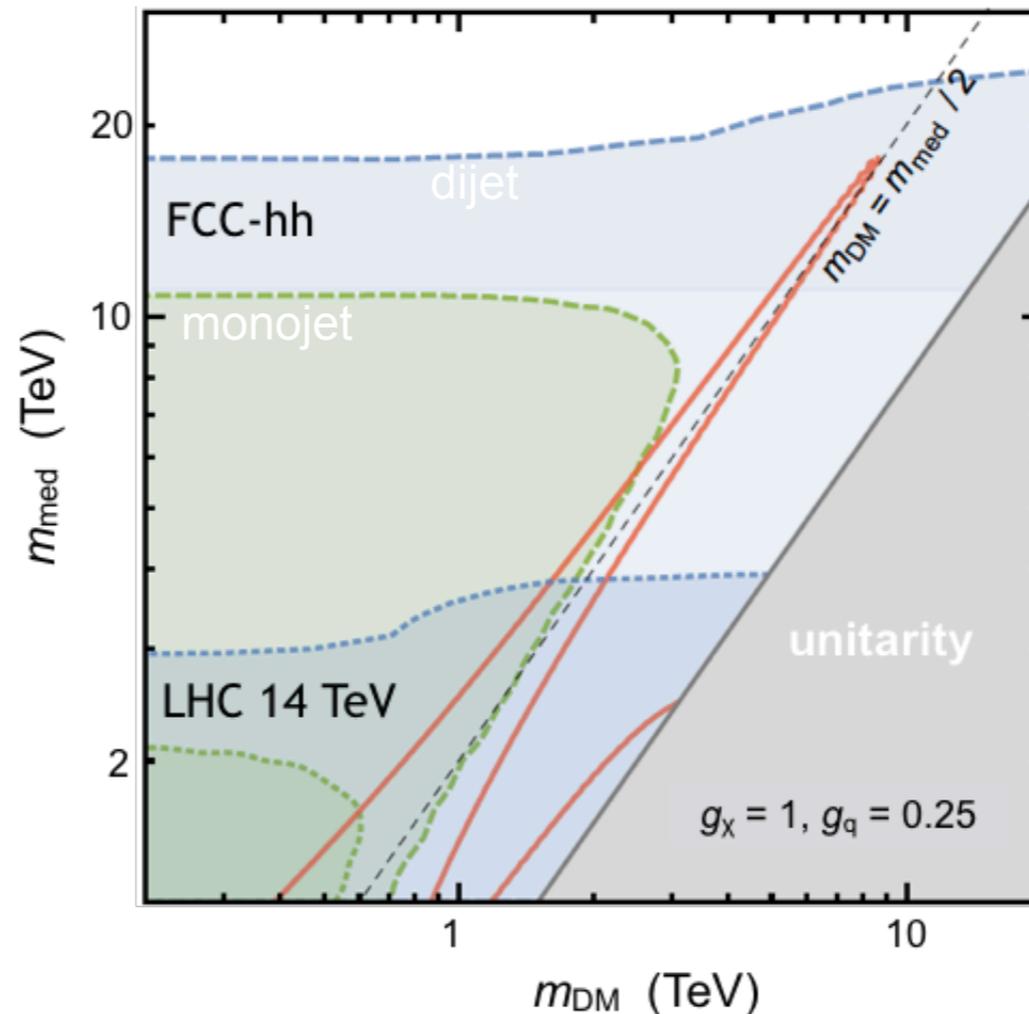
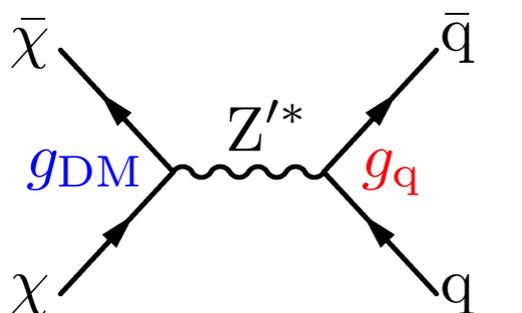
di-jet resonance



mono-jet with large $E_{T,miss}$



relic density



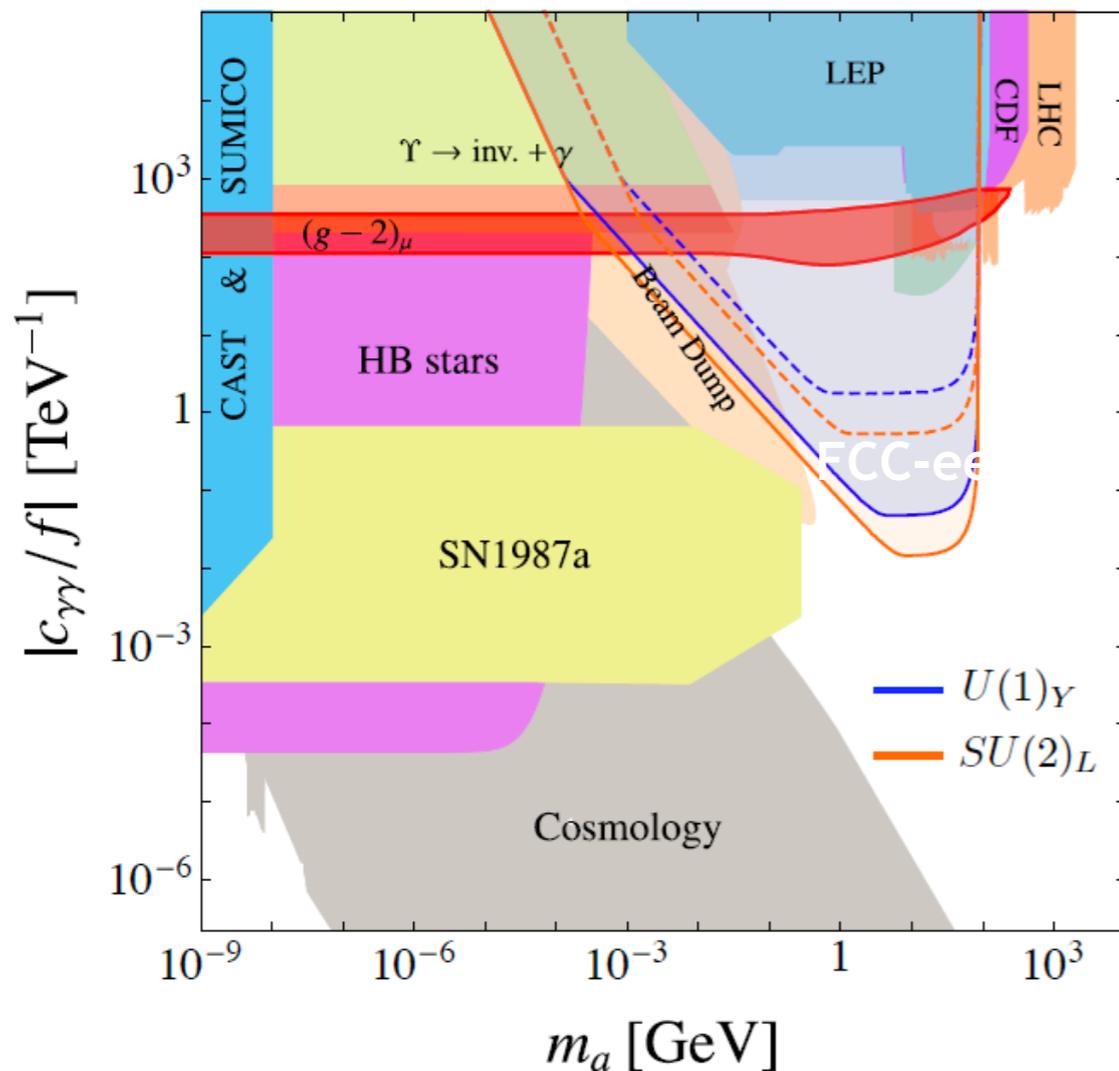
for illustration, fix the ratio between DM and mediator mass and pick mass that gives the right relic density

FCC-hh covers a much larger parameter space with reasonable couplings

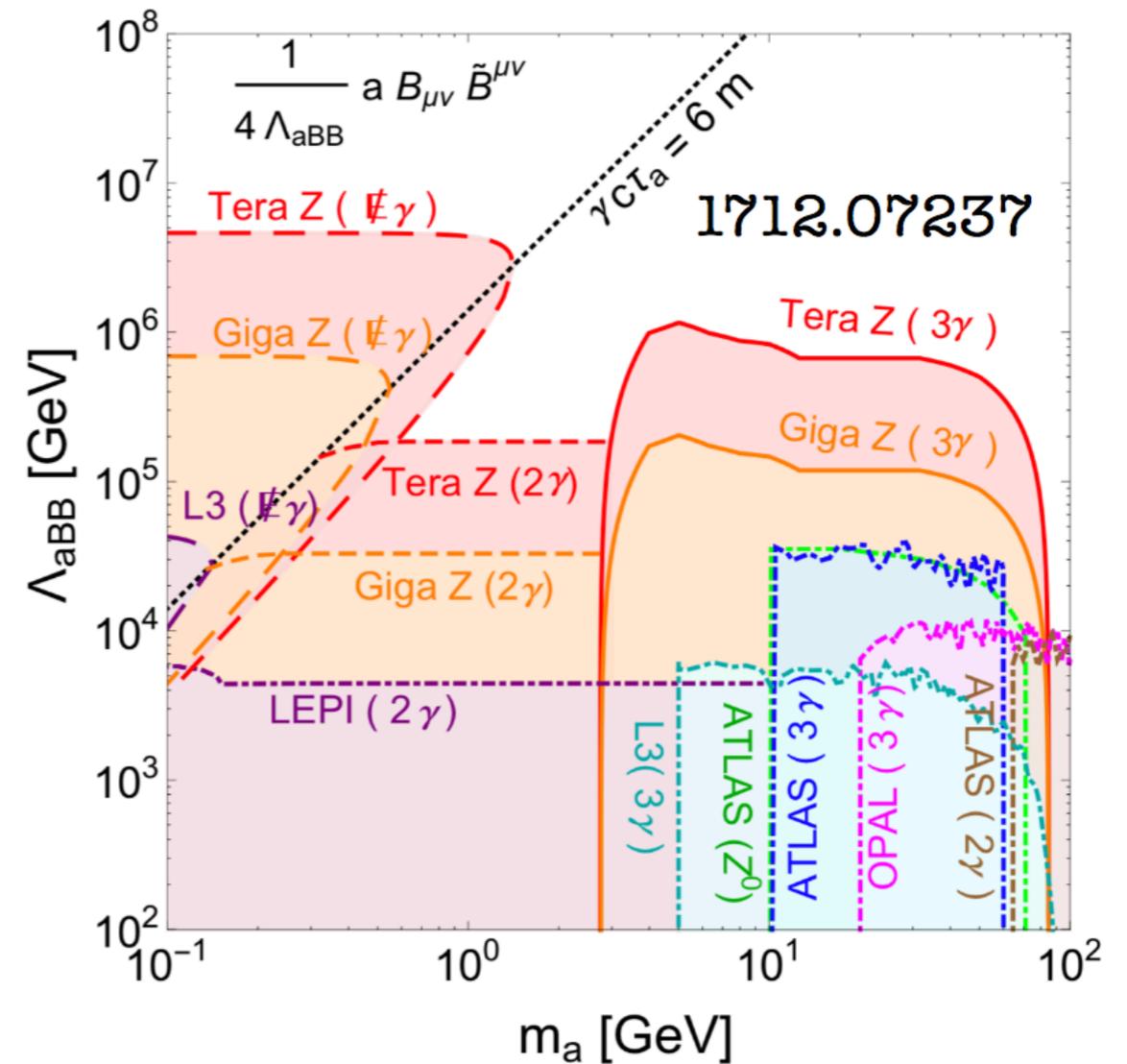
New Physics at the Z-Pole

A weekly coupled window on the Dark Sector

Search axion-like particles (ALP)



$$Z \rightarrow \gamma a, a \rightarrow \gamma\gamma$$



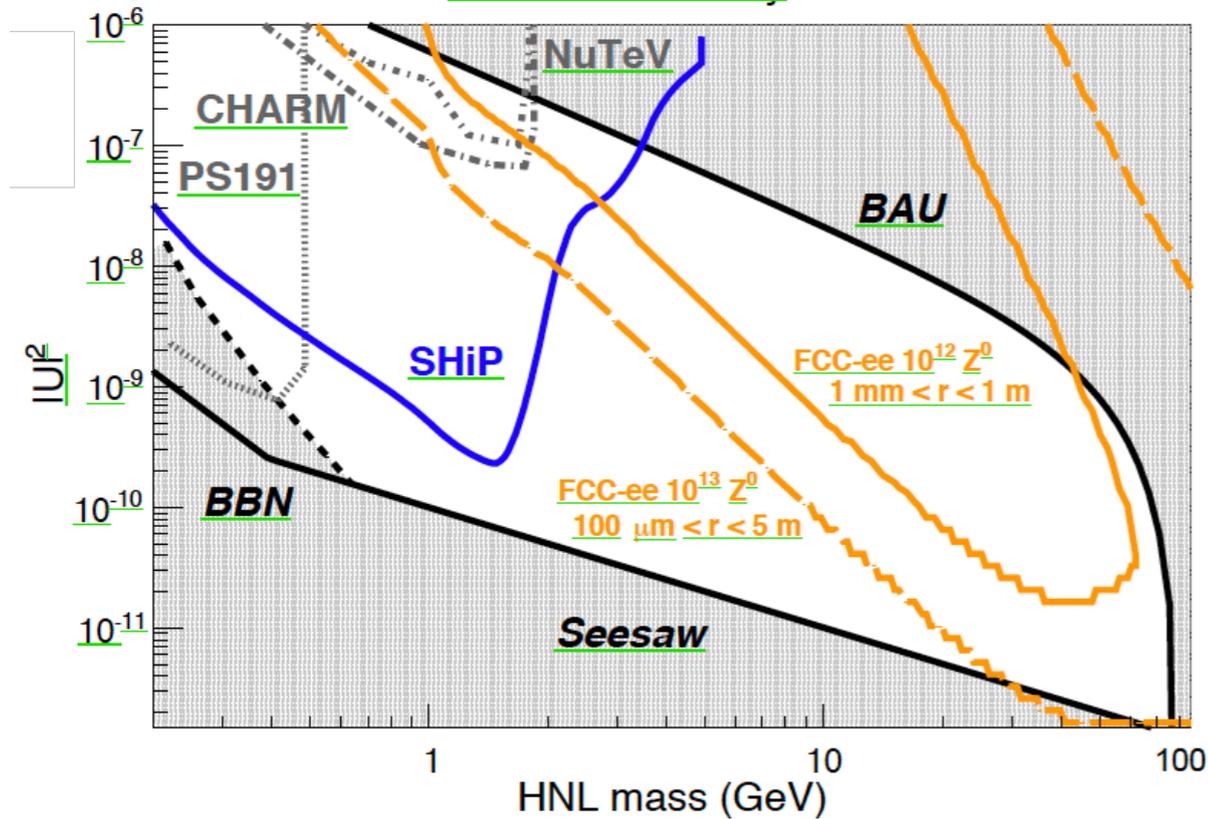
- very light: $Z \rightarrow \gamma +$ missing energy
- light: $Z \rightarrow \gamma\gamma$
- heavier: $Z \rightarrow \gamma\gamma\gamma$

New Physics at the Z-Pole

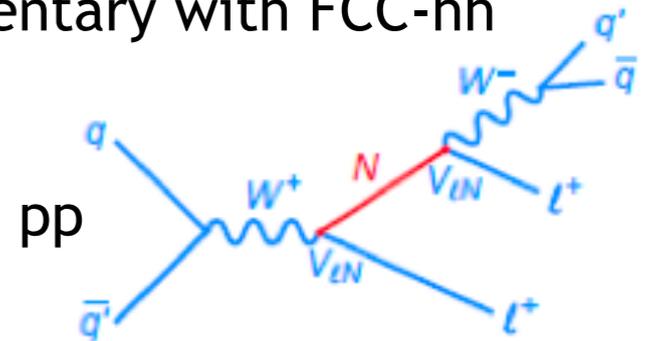
Huge production of neutrinos: $2 \times n_Z \times 20\%$!

Search for $Z \rightarrow \nu N$, $N \rightarrow W \ell / Z \nu$

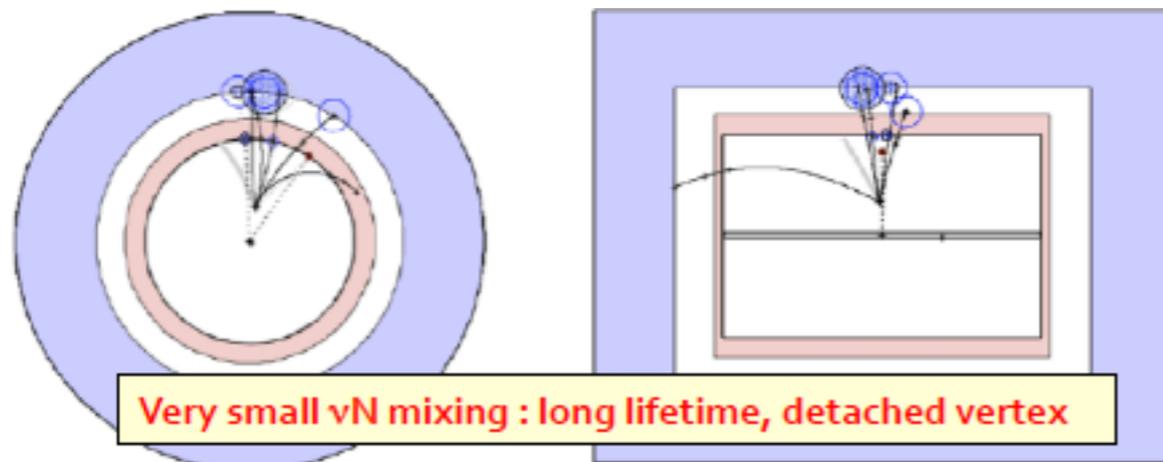
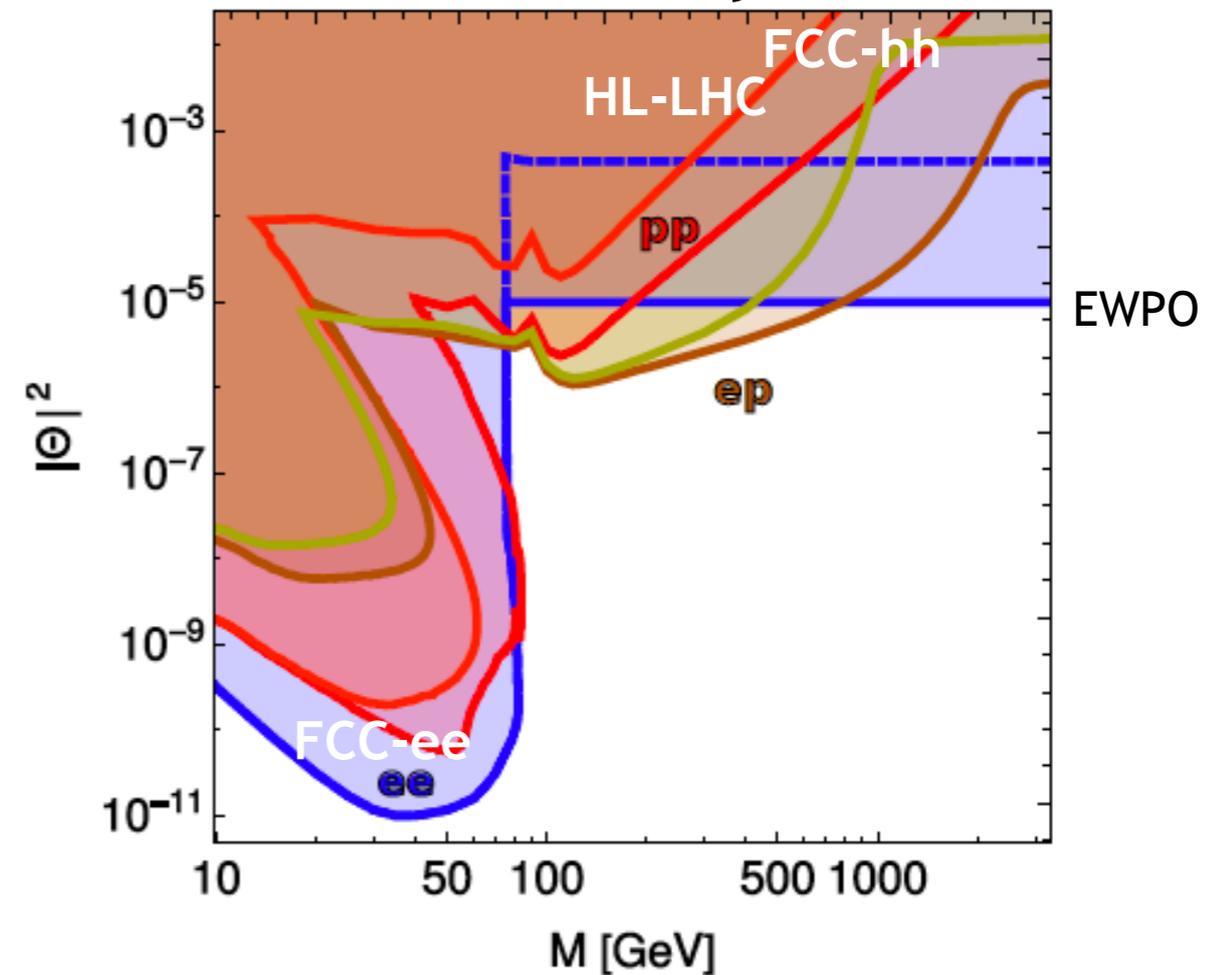
Inverted hierarchy

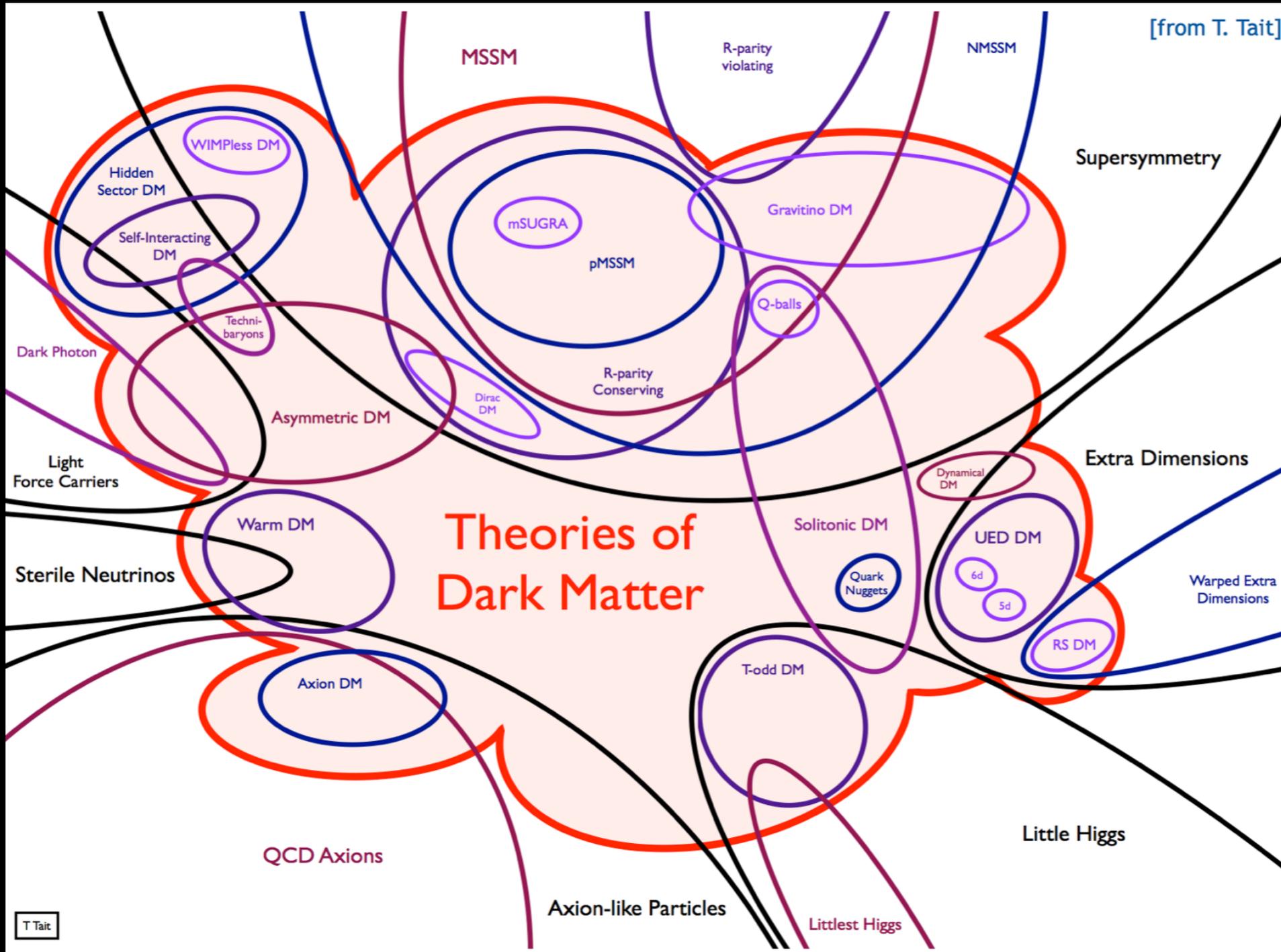


complementary with FCC-hh



Search for heavy neutrinos





Gautier Hamel de Monchenault