

BACKUP

**Top quark physics  
at the LHC.  
Selected highlights**

*Laboratory Seminar  
LPNHE  
Sorbonne Université  
Paris  
24th June 2019*

*Francesco Spanò*



ROYAL  
HOLLOWAY  
UNIVERSITY  
OF LONDON

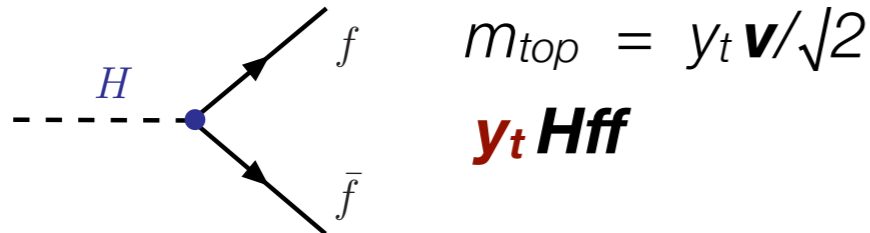
**BACK-UP**

# General top quark features

# Why heavy (quarks) at multi-purpose colliders? (I)

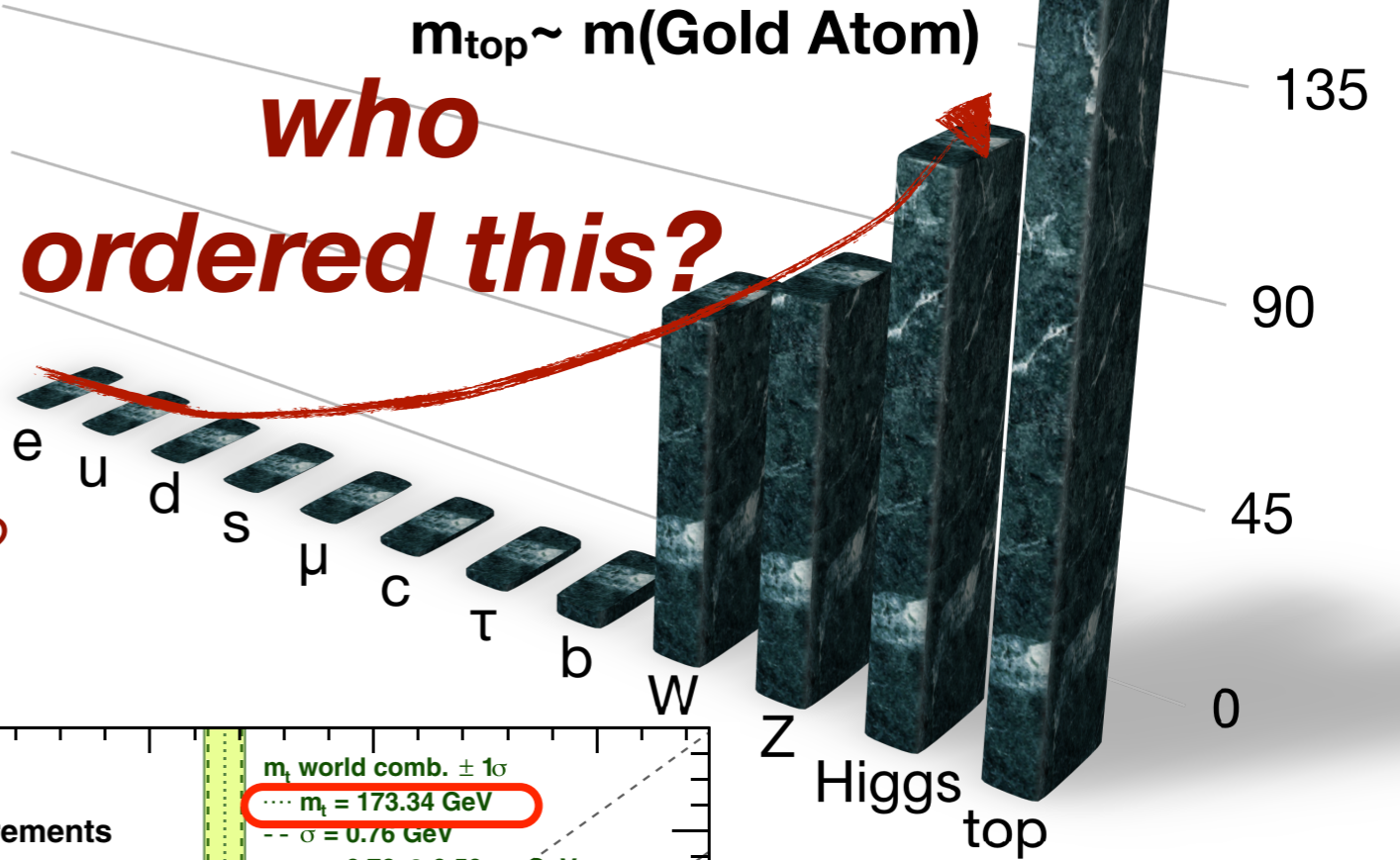
**“Special” reasons**

**top & b are the most massive known quarks**



**largest (unmeasured) coupling  $y_t$  to Higgs boson** → **Is  $y_t \sim 1$ ?**

■ Masses of known fundamental particles



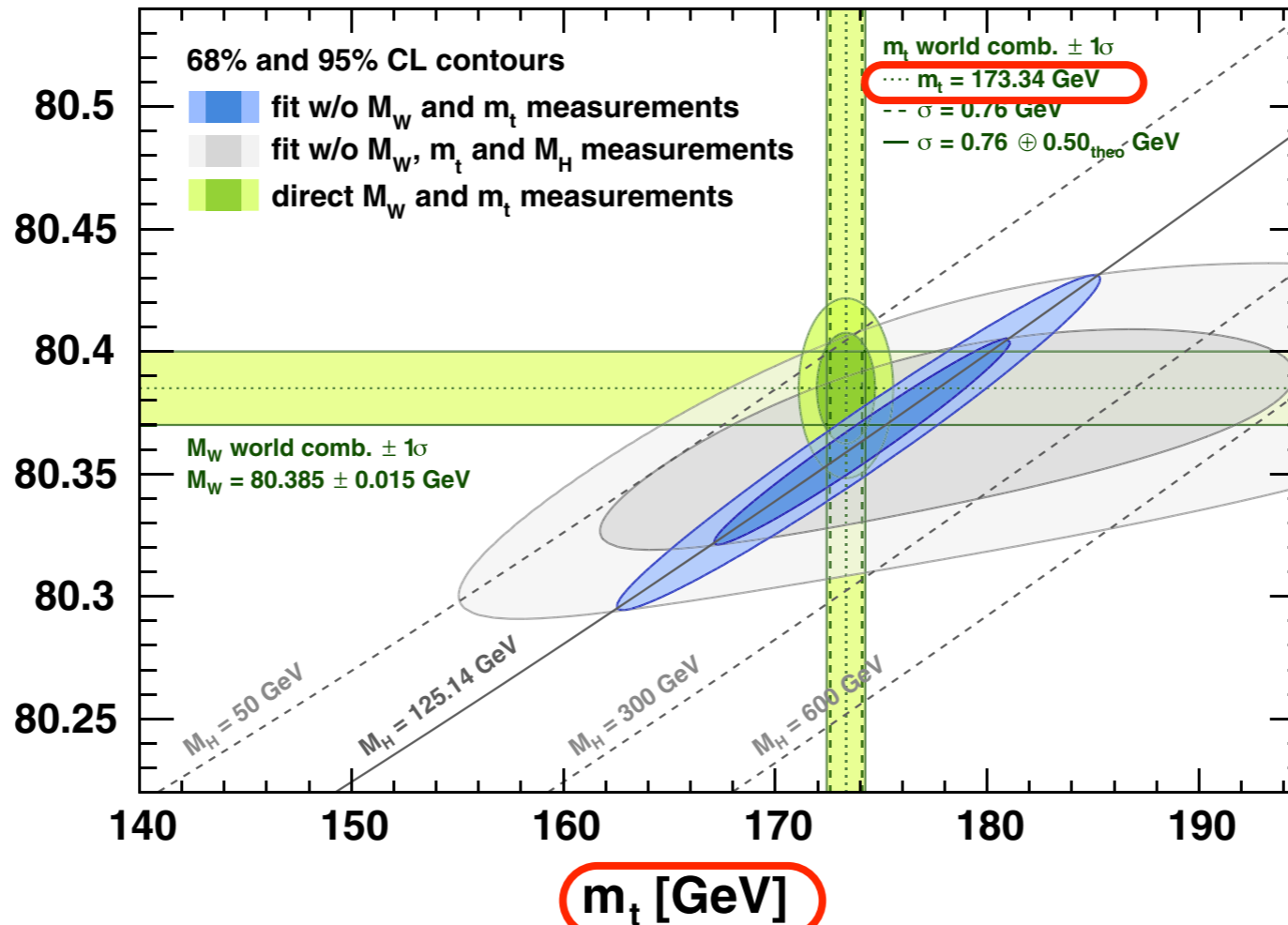
## New W mass by ATLAS

$m_W = 80370 \pm 7$  (stat)  $\pm 11$  (exp syst)  $\pm 14$  (mod sys) MeV =  **$80370 \pm 19$  MeV**  
 7 TeV pp-collision data

## Tevatron indirect W mass

$m_W = 80351 \pm 15$  (stat)  $\pm 10$  (syst) =  **$80351 \pm 18$  MeV**  
 using 1.9 TeV p-antip collision

**$M_W$  [GeV]**



The GFitter Group, Eur. Phys. J. C 74:3046 (2014)

*The reason(s) for mass hierarchy is (are) still a mystery*

# top quark production

# The SM Lagrangian (2019)

reminder(II)

The Standard Model

$$\phi = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + \mathbf{H} \end{pmatrix}$$

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i \bar{\Psi} \not{D} \Psi + h.c. + |\mathcal{D}_\mu \phi|^2 - V(\phi) + \bar{\Psi}_i y_{ij} \Psi_j \phi + h.c.$$

$\mathbf{H}$  = Higgs particle

$$\mathcal{L}_W \sim g^2 (v + H)^2 W^+_\mu W^{-\mu}$$

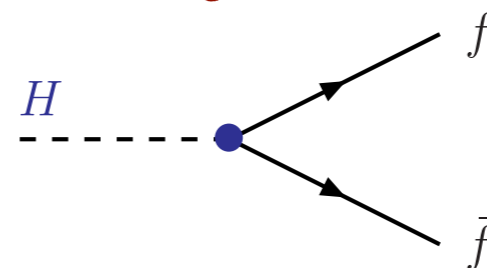
$$M_W = \frac{1}{2} g_2 v = \left( \frac{\sqrt{2} g^2}{8 G_\mu} \right)^{1/2}$$

$$v = \frac{1}{(\sqrt{2} G_\mu)^{1/2}} \rightarrow v \sim 246 \text{ GeV}$$

$$\mathcal{L}_f = m_f \bar{f}_L f_R + y_f H \bar{f}_L f_R / \sqrt{2} + h.c.$$

mass term                      interaction term

$$m_f = y_f v / \sqrt{2}$$



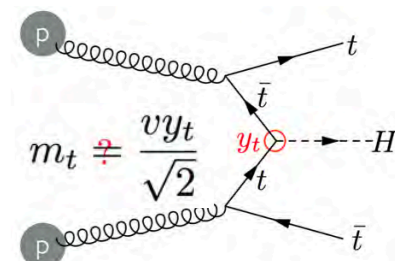
$$m_{b,obs} \sim 4 \text{ GeV}$$

$$m_t \sim 173 \text{ GeV}$$

**top: largest (unmeasured) coupling  $y_t$  to Higgs boson**

replacing values gives  $y_{top} = \sqrt{2} m_{top} / v \sim \sqrt{2} 173 / 246 \sim 0.99$

**Is  $y_t \sim 1$ ? Is it SM Higgs?**



# The SM Lagrangian (2019)

reminder(I)

The Standard Model: QF theory invariant under  $SU(2) \times U(1) \times SU(3)$

observed

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i \bar{\Psi} \not{D} \Psi + h.c. + |\not{D}_\mu \phi|^2 - V(\phi) + \bar{\Psi}_i y_{ij} \Psi_j \phi + h.c.$$

Gauge fields are in kinetic terms and cov. derivative  $D_\mu$

**Spontaneous symmetry breaking:** the Lagrangian shows the **possibility that at a given energy scale, the symmetry of the observed physical states is different from the symmetry of the Lagrangian interactions**, by realising one of multiple asymmetric configurations (minimum potential energy state)

**spin-1 W&Z bosons emerge as massive** by coupling to H **photon remains massless**

$g_2$   $SU(2)$  gauge coupling

$$\mathcal{L}_{W,Z} \sim g_2^2 (v+H)^2 W_\mu^+ W^{-\mu} + g_2^2 / 8 \cos^2 \theta_W Z_\mu Z^\mu$$

$$M_W = \frac{1}{2} g_2 v = \left( \frac{\sqrt{2} g^2}{8 G_\mu} \right)^{1/2} M_Z \cos \theta_W = M_W$$

$$v = \frac{1}{(\sqrt{2} G_\mu)^{1/2}} \rightarrow v \sim 246 \text{ GeV}$$

**spin-0 H=Higgs boson emerges**

$$\phi = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v+H \end{pmatrix}$$

**spin 1/2 fermions (u,c,d,s,b,...)**

- negative chirality ( $F_L$ ) state **couple to W, Z** by covariant derivative
- obtain mass from **assuming gauge invariant coupling terms** (Yukawa coupling) **to  $\phi$** :  $F_L \phi f_R = y (v+H) f_L f_R$   
 $\mathcal{L}_Y = -Y_{ij}^d \bar{Q}_{Li}^I \phi d_{Rj}^I - Y_{ij}^u \bar{Q}_{Li}^I \epsilon \phi^* u_{Rj}^I + h.c.$

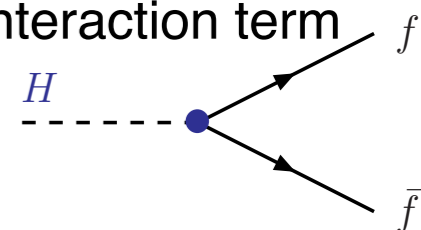
Diagonalise Y and replace  $\phi$  from SSB

$$\mathcal{L}_f = m_f \bar{f}_L f_R + y_f H \bar{f}_L f_R / \sqrt{2} + h.c.$$

mass term

$$m_f = y_f v / \sqrt{2}$$

fermion-Higgs interaction term



# Top production @ LHC: differential growth

(Campbell et al, Rept.Prog.Phys.70:892007)

$$\sigma = \sum_{i,j} \int_0^1 dx_1 dx_2 f_i(x_1, \mu) f_j(x_2, \mu) \hat{\sigma}_{ij} = \sum_{i,j} \int \left( \frac{d\hat{s}}{\hat{s}} dy \right) \left( \frac{dL_{ij}}{d\hat{s} dy} \right) (\hat{s} \hat{\sigma}_{ij}) \sim \sum_{i,j} \frac{\Delta\hat{s}}{\hat{s}} \left( \frac{dL_{ij}}{d\hat{s}} \right) (\hat{s} \hat{\sigma}_{ij})$$

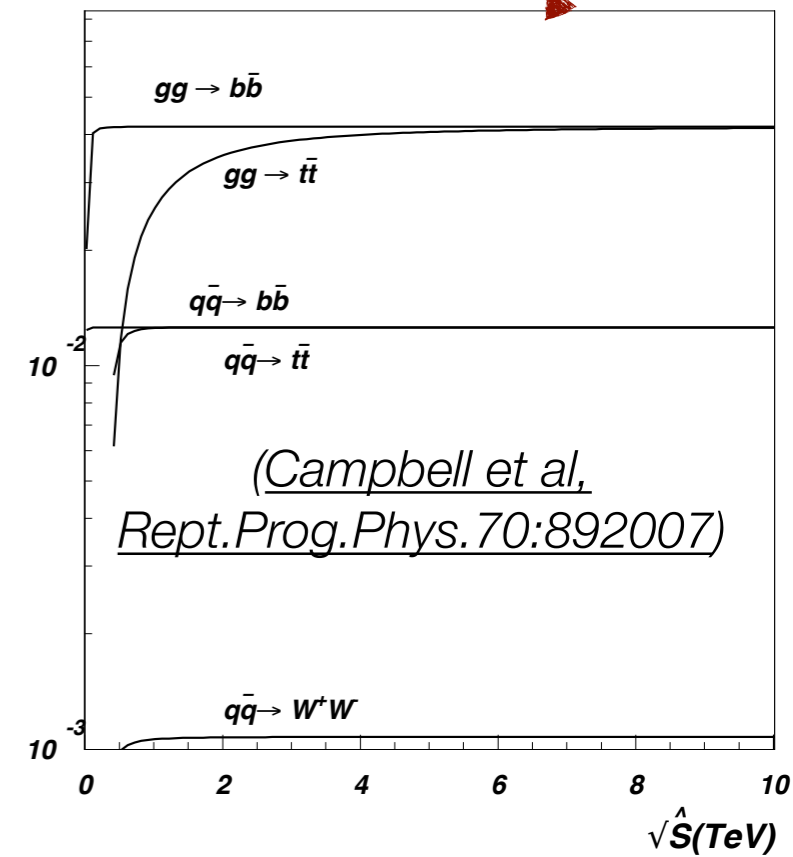
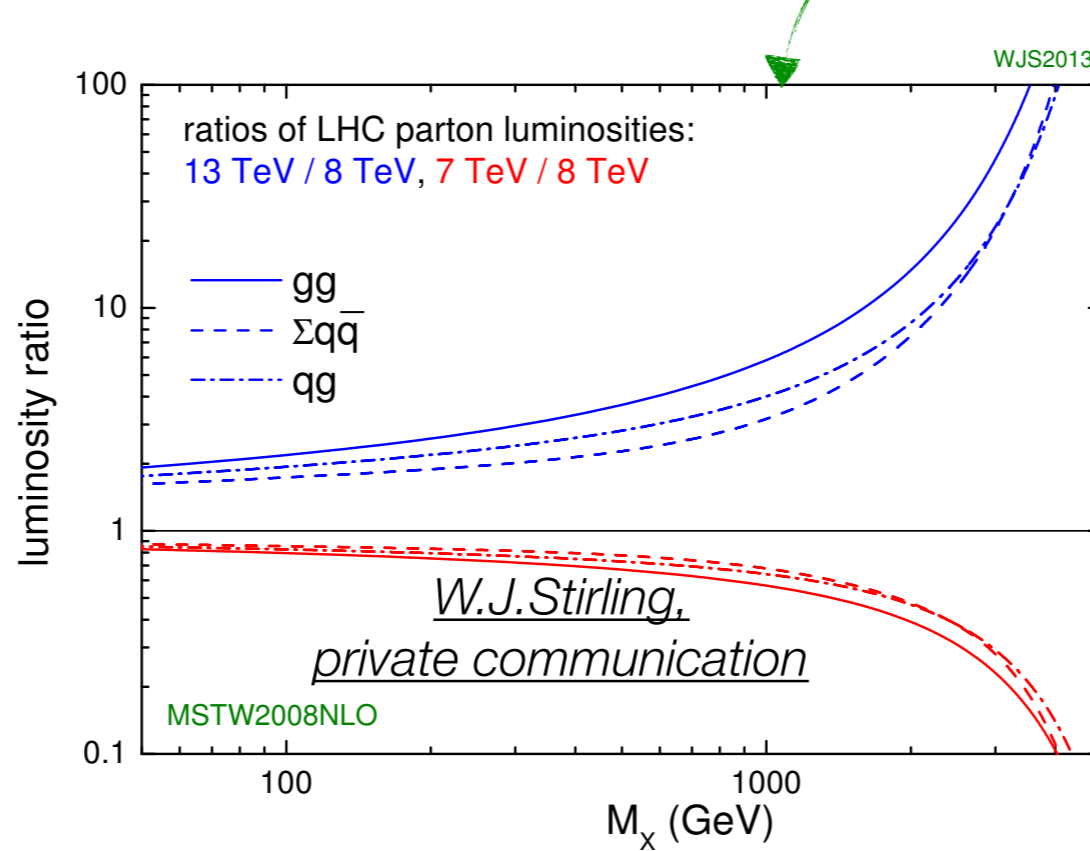
$$\left( \frac{dL_{ij}}{d\hat{s} dy} \right) = \frac{1}{s} \frac{1}{1 + \delta_{ij}} [f_i(x_1, \mu) f_j(x_2, \mu) + (1 \leftrightarrow 2)]$$

- Different **x-range** and **center of mass** dependence incorporated in **parton luminosities** →

▶ **gg** → **X** dominated processes **grow more than qq** → **X** ones

▶ **larger gains at high multi-TeV masses** ~up to **O(100)**

- **Cross sections in “tails” increase more rapidly than** inclusive value



Magano, Rojo, JHEP(1208),2012:10

$R^{th, nnpdf} = 14\text{TeV to } 8\text{ TeV xsec ratios}$

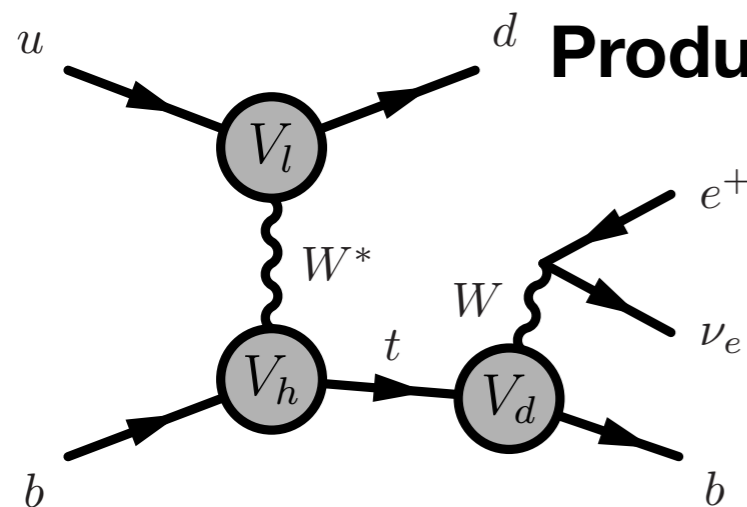
Cross Section	$R^{th, nnpdf}$	$\delta_{PDF}(\%)$	$\delta_{\alpha_s}(\%)$	$\delta_{scales}(\%)$
$t\bar{t}/Z$	2.12	± 1.3	-0.8 - 0.8	-0.4 - 1.1
$t\bar{t}$	3.90	± 1.1	-0.5 - 0.7	-0.4 - 1.1
$t\bar{t}(M_{t\bar{t}} \geq 1\text{ TeV})$	8.18	± 2.5	-1.3 - 1.1	-1.6 - 2.1
$t\bar{t}(M_{t\bar{t}} \geq 2\text{ TeV})$	24.9	± 6.3	-0.0 - 0.3	-3.0 - 1.1



# Top quark predictions@ LHC- the NNLO revolution : single top

## Production and decay @ NNLO, decay in narrow width approx

*Phys. Rev. D 94, 071501 (2016)*



@  $\sqrt{s} = 8 \text{ TeV}$

$V_{l,h,q}$  = three corrections to light quark line, heavy quark line and decay including **corrections from two loop, one loop+1 real emission, two real emission**

inclusive [pb]	LO	NLO	NNLO
$t$ quark	$143.7^{+8.1\%}_{-10\%}$	$138.0^{+2.9\%}_{-1.7\%}$	$134.3^{+1.0\%}_{-0.5\%}$
$\bar{t}$ quark	$85.8^{+8.3\%}_{-10\%}$	$81.8^{+3.0\%}_{-1.6\%}$	$79.3^{+1.0\%}_{-0.6\%}$

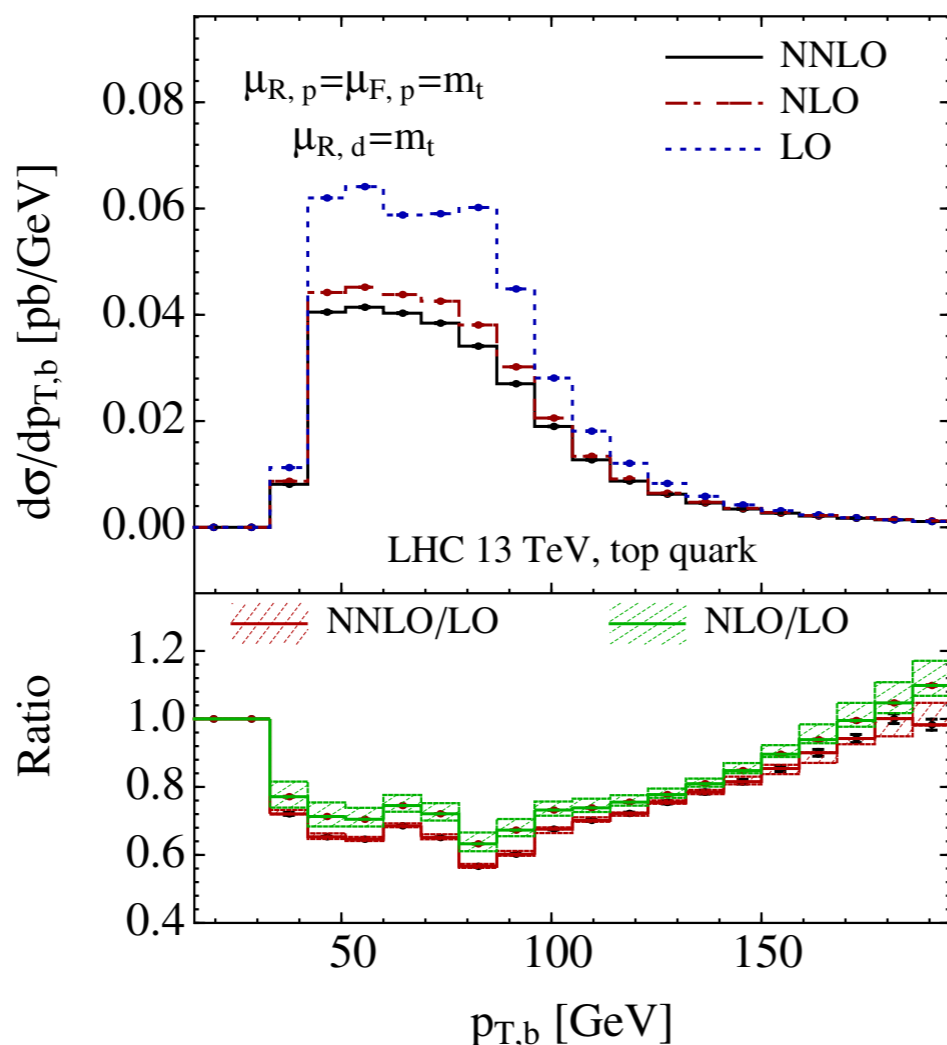
2jets, 1b-jet antik $_T$  R=0.5  
 $p_{T,jet} > 40 \text{ GeV}$ ,  $|\eta_{jet}| < 5$  (2.4 for b-jet)  
 $p_{T,lep} > 30 \text{ GeV}$  |  $|\eta_{lep}| < 2.4$



fiducial [pb]		LO	NLO	NNLO
$t$ quark	total	$4.07^{+7.6\%}_{-9.8\%}$	$2.95^{+4.1\%}_{-2.2\%}$	$2.70^{+1.2\%}_{-0.7\%}$
	corr. in pro.		-0.79	-0.24
	corr. in dec.		-0.33	-0.13
$\bar{t}$ quark	total	$2.45^{+7.8\%}_{-10\%}$	$1.78^{+3.9\%}_{-2.0\%}$	$1.62^{+1.2\%}_{-0.8\%}$
	corr. in pro.		-0.46	-0.15
	corr. in dec.		-0.21	-0.08

*stable values*  
*reduced uncertainties*

Differential distributions also available @NNLO

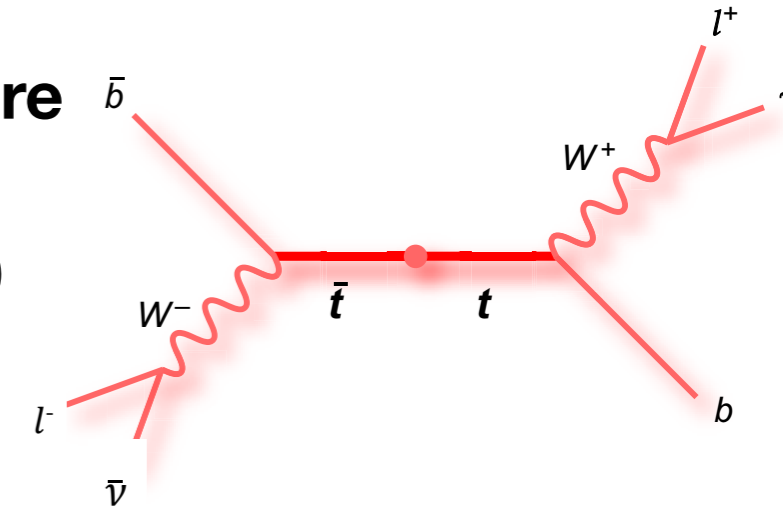


# Standard reasons: Extreme test of SM: $d\sigma_{tt}/dp_T$ "saga" - dilepton

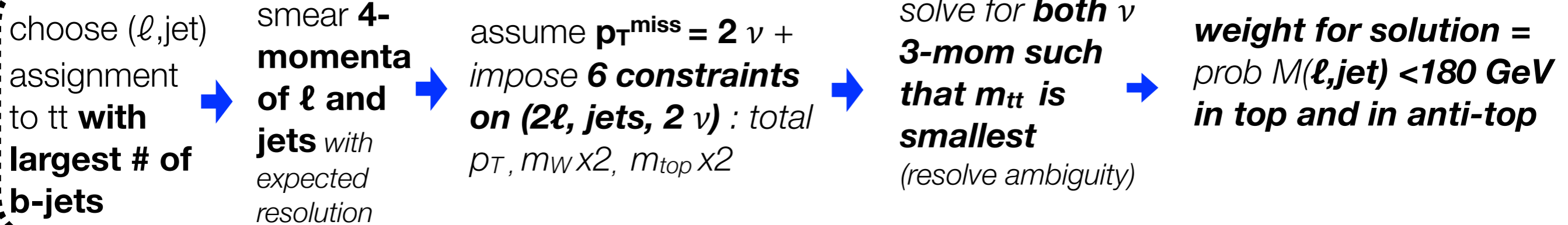
$\sqrt{s}=13\text{TeV}$

JHEP 02 (2019) 149

- Particle flow  $\rightarrow$  individual particles using all CMS subdet  $\rightarrow$  **Require**
  - **2 opposite sign  $\ell$  ( $e, \mu$ ),  $\geq 2$  jets,  $\geq 1$  b-tag**
  - **$m(\ell^+\ell^-) > 20$  GeV and  $\neq M_Z$  (15 GeV window), large  $p_T^{\text{miss}} (>40$  GeV)**
- **Bkg:** data-driven Z+jets, simulated tW,W/Z jets, other tt



- **Reconstruct tt system with kinematic reco**

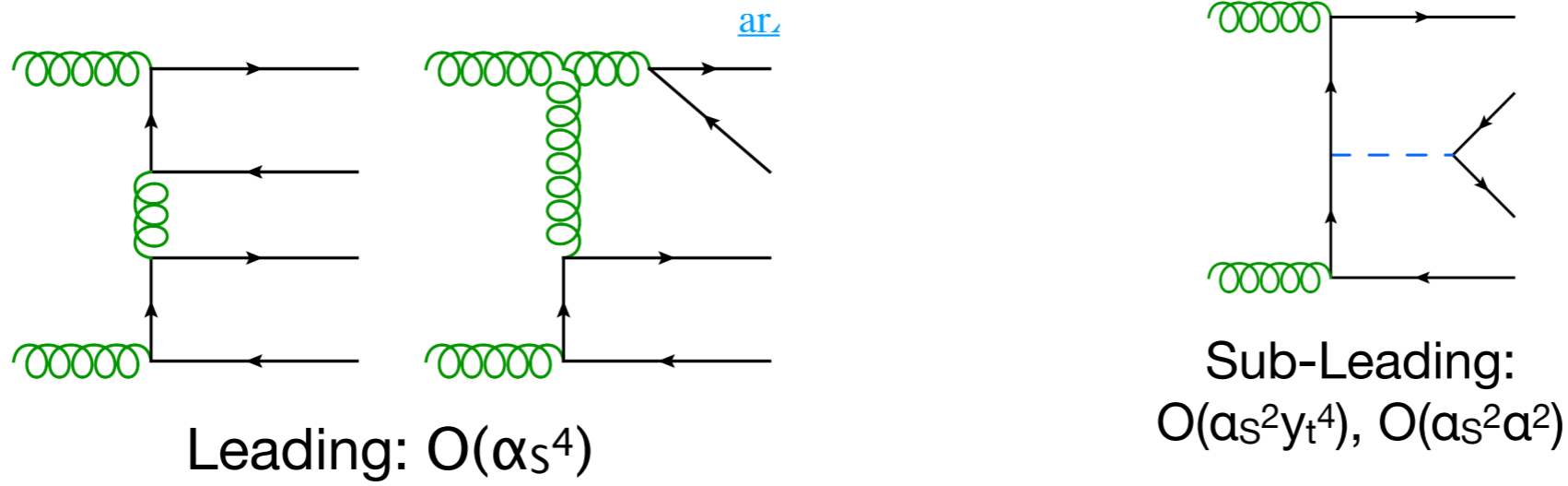


**two separate  $\nu$  3-momenta for given assignment = weighted average of 100 smeared repetitions**

**keep  $(\ell, \text{jet})$  assignment with maximum  $\Sigma$  weights**

- **Bkg-subtract & Unfold to parton and particle level  $\rightarrow d\sigma_{tt}/dX$**   $\rightarrow$

# Status of Search for observation of 4 top quarks



$$\sigma^{\text{NLO}}(\text{tttt}) = 11.97 \text{ fb at NLO QCD + NLO QED} \quad 13 \text{ TeV}$$

Significance obs. (exp.) [ $\sigma$ ]	ATLAS 36 fb <sup>-1</sup>	CMS 36 fb <sup>-1</sup>	CMS 139 fb <sup>-1</sup>
<b>SS/ML</b>	3.0 (0.8) <u>1</u>	1.6 (1.0) <u>3</u>	2.6 (2.7) <u>6</u>
<b>1L/OS</b>	1.0 (0.6) <u>2</u>	0.0 (0.4) <u>4</u>	-
<b>Combination</b>	2.8 (1.0) <u>2</u>	1.4 (1.1) <u>4</u>	-

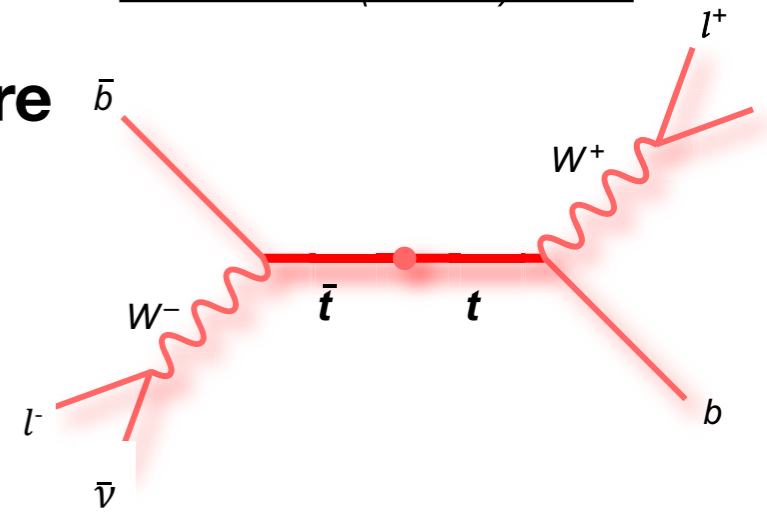
(table and diagrams by Nedaa Alexandra Asbah)

# Standard reasons: Extreme test of SM: $d\sigma_{tt}/dp_T$ "saga" - dilepton

$\sqrt{s}=13\text{TeV}$

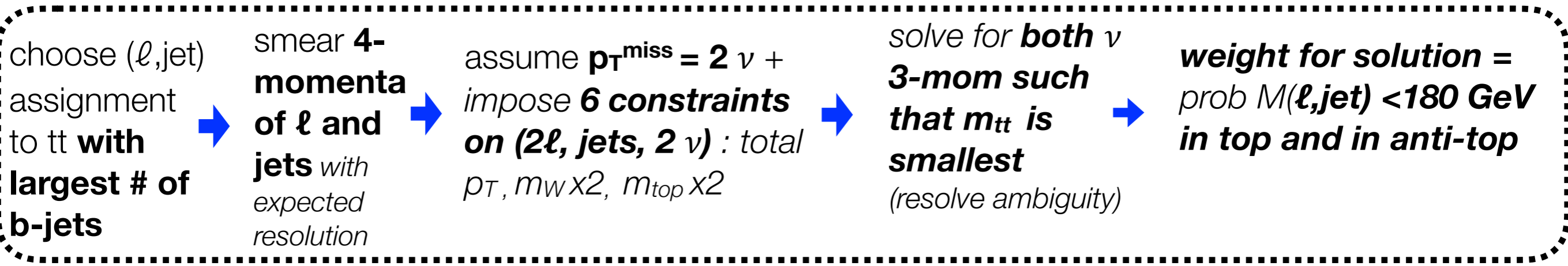
JHEP 02 (2019) 149

- Particle flow  $\rightarrow$  individual particles using all CMS subdet  $\rightarrow$  **Require**
  - 2 opposite sign  $\ell$  (e, $\mu$ ),  $\geq 2$  jets,  $\geq 1$  b-tag
  - $m(\ell^+\ell^-) > 20$  GeV and  $\neq M_Z$  (15 GeV window), large  $p_T^{\text{miss}} (>40$  GeV)



• **Bkg:** data-driven Z+jets, simulated tW,W/Z jets, other tt

## Reconstruct tt system with kinematic reco



90% efficient on tt

two separate  $\nu$  3-momenta for given assignment = weighted average of 100 smeared repetitions

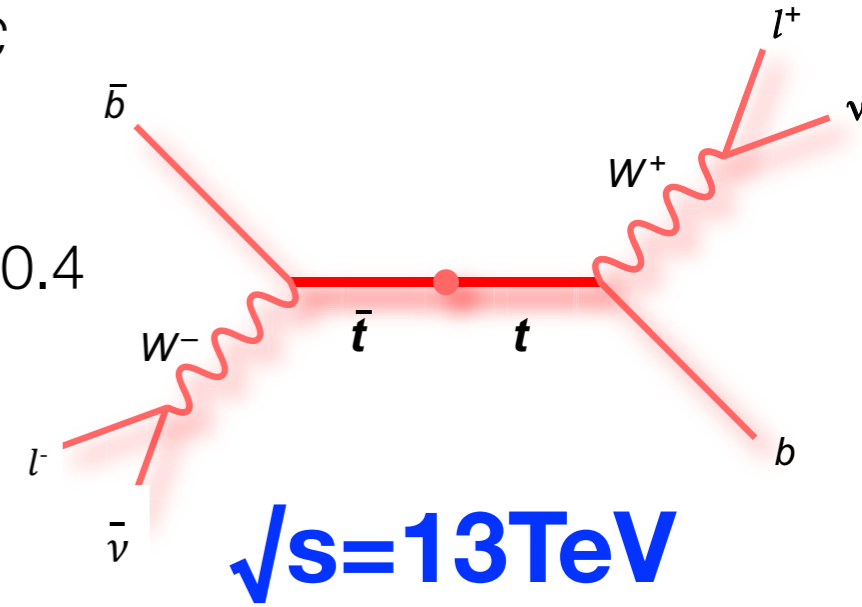
keep ( $\ell$ ,jet) assignment with maximum  $\Sigma$  weights

• **Bkg-subtract & Unfold to parton and particle level  $\rightarrow d\sigma_{tt}/dX$**

# Extreme test of SM: double and triple diffxsec - dilepton+jets

*CMS-TOP-18-004, submitted to Eur. Phys J. C*

- 13 TeV Dilepton selection as *JHEP 02 (2019) 149*
- **Extra jets:** central, high  $p_T$  jets with  $\Delta R(\text{e-jet}, \text{lep}) = \Delta R(\text{e-jet}, \text{b-jet}) > 0.4$
- **Bkg:** data-driven Z+jets, simulated tW, W/Z jets, other tt
- **Reconstruct tt system with dilepton kinematic reco**

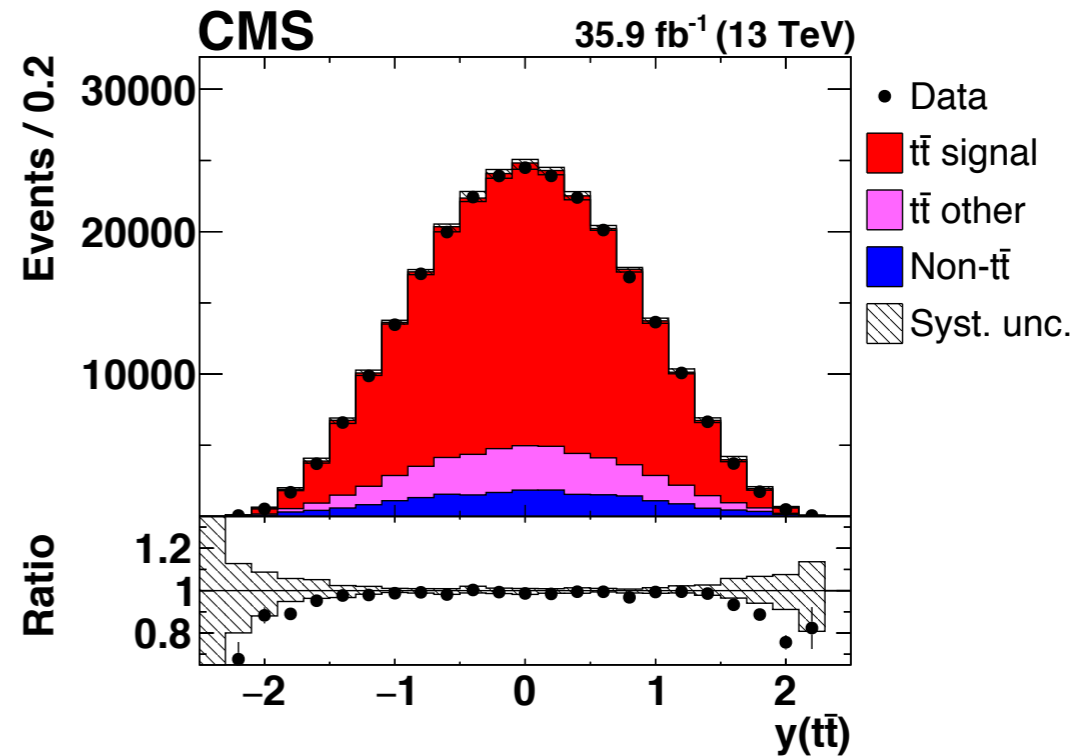
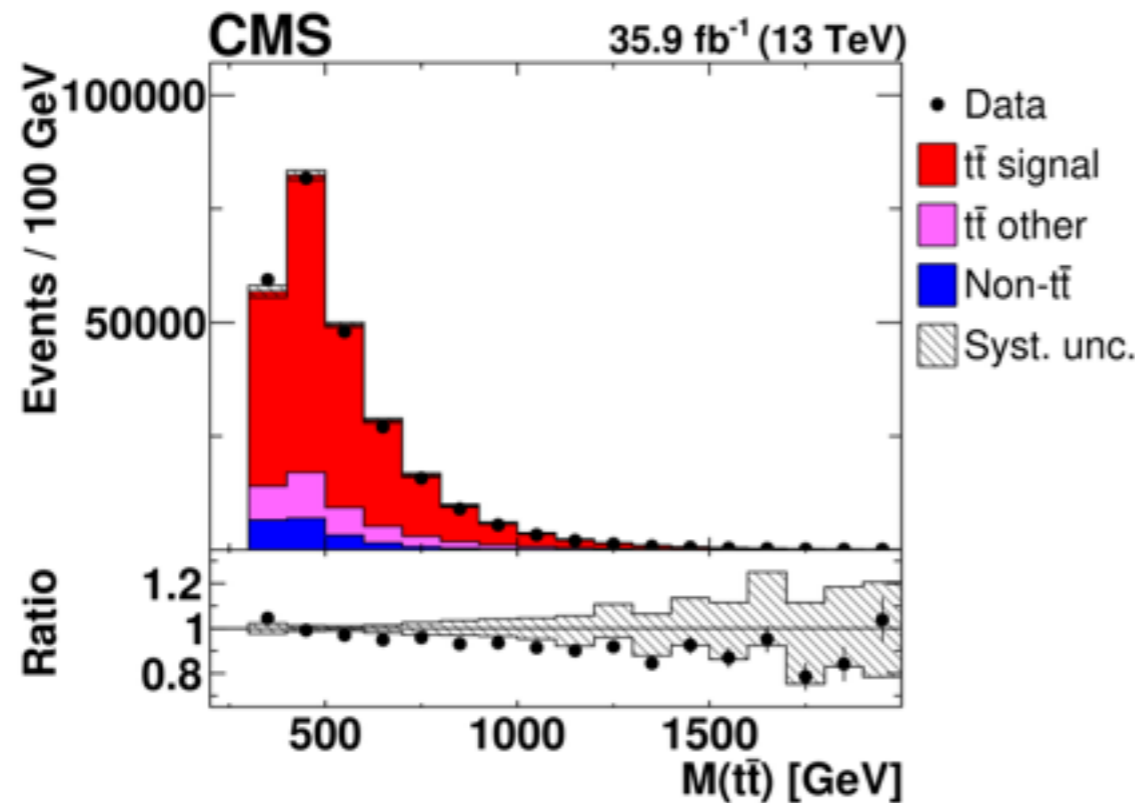


standard for 2d distributions

- $M_{tt}$  vs  $\{p_{T, \text{top}} | y_{\text{top}} | |y_{tt}| \Delta\eta(t, \bar{t}), \Delta\phi(t, \bar{t}), p_{T, tt}\}$
- $[|y_{\text{top}}|, p_{T, \text{top}}]$

loose: for 3d distributions

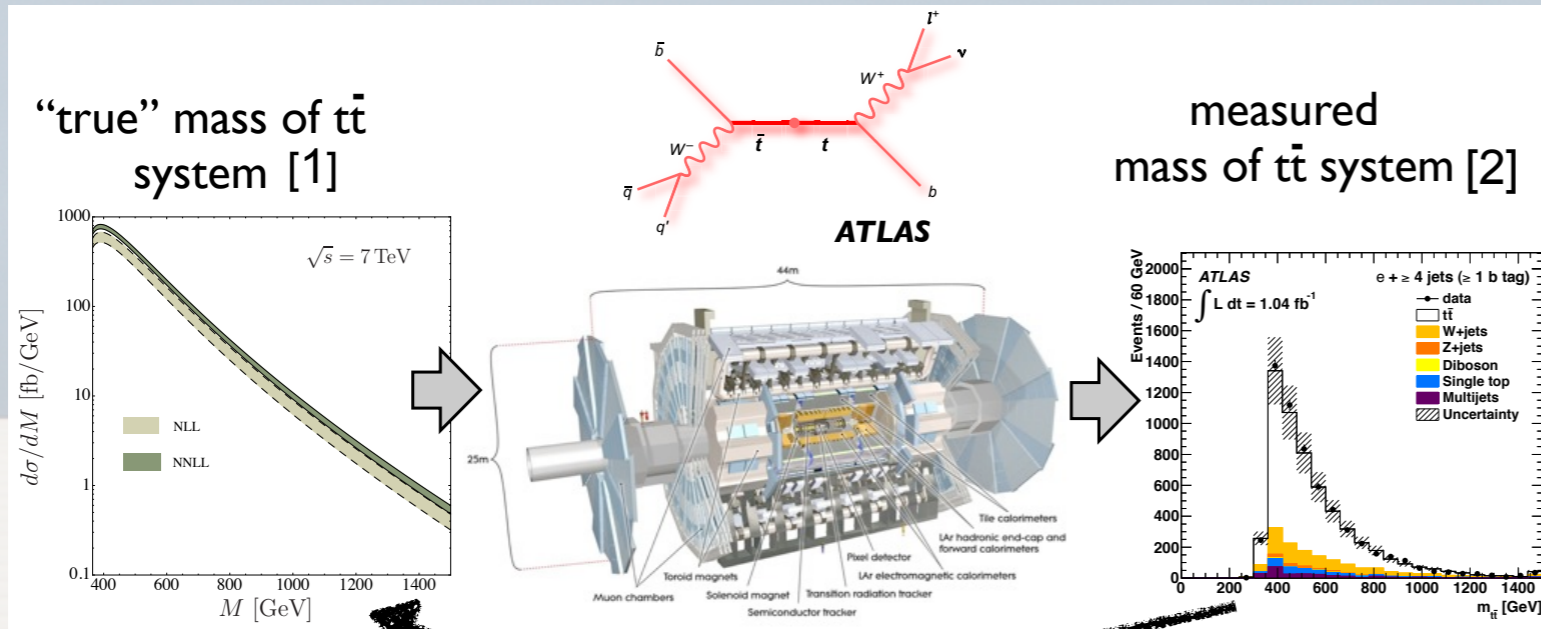
- assume  $\mathbf{p}_T^{\text{miss}} = 2 \nu + p_z^{\text{miss}} = \mathbf{p}(\ell^+ \ell^-), \mathbf{E}(\ell^+ \ell^-)$ ,  $m(2 \nu) > 0, m(WW) \geq 2 M_W$
- keep  $(\ell, \text{jet})$  assignment with maximum  $p_T$  jets
- $[M_{tt}, y_{tt}, N_{\text{extra jets}}]$  2 bins (0,1) and 3 bins (0,1,2)



# Unfolding Foundations

• Unfolding: **infer an unknown distribution  $f(y)$**  for a variable  $y$  **from the measured distribution  $g(s)$**  by using knowledge and/or assumptions on the probability distribution that links the observation to the “true” value.

$$g(s) = \int_{\Omega} K(s, y) f(y) dy + b(s)$$



# Unfolding

**FS**  
**EPJ Web of conference**  
**vol 55 (2013) 03002**

$$\nu_i = \int_{s_{i-1}}^{s_i} g(s) ds$$

$$\mu_j = \int_{y_{j-1}}^{y_j} f(y) dy$$

$$E[n_i] = \nu_i = \sum_{j=1}^M R_{i,j} \mu_j + \beta_i$$

$$E[\mathbf{n}] = \boldsymbol{\nu} = R\boldsymbol{\mu} + \boldsymbol{\beta}$$

- Ingredients**
- response matrix  $R$ : prob true in bin  $i$  is reco in bin  $j$
  - number of observed events  $n$
  - reconstruction efficiency
  - estimate of expected bkg  $\boldsymbol{\beta}$

the Max LKL solution

$$E[\mathbf{n}] = \boldsymbol{\nu} = R\boldsymbol{\mu} + \boldsymbol{\beta}$$

$$\mathcal{L} = \prod_{i=1}^N \nu_i^{n_i} \frac{e^{-\nu_i}}{n_i!}$$

$$\partial \log \mathcal{L}(\mu_i) / \partial \mu_i = 0 \quad \forall i$$

**Small changes in input (can) lead to large changes** in the ML estimate.

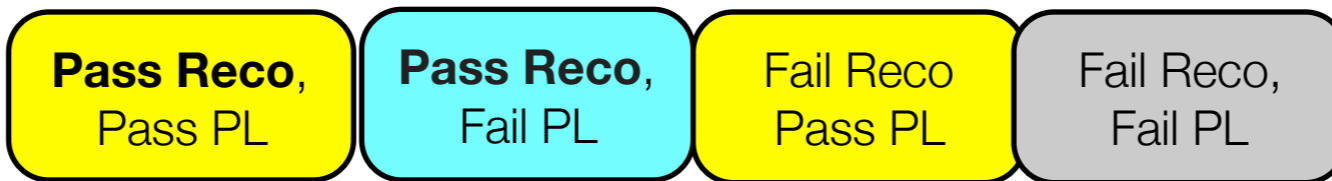
- **Regularize** = Reduce impact of high frequency while keeping info of high significance, stable components → **reduction in variance** w.r.t. ML estimator

# Example $d\sigma_{t\bar{t}}/dp_{T,top(-jet)}$ l+jets @ particle level (PL)- $\sqrt{s} = 8$ TeV

*Phys. Rev. D 93, 032009 (2016)*

**Reco**

*selected out*



*similar (simpler) formulas for parton level*

*to correct to*

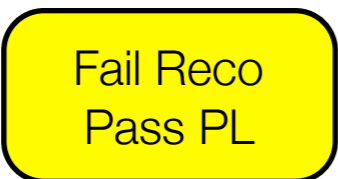
$$\frac{d\sigma_{t\bar{t}}}{dp_{T,ptcl}}(p_{T,ptcl}^i) = \frac{1}{\Delta p_{T,ptcl}^i \mathcal{L}_{ptcl!reco}^{f^i}} \cdot \sum_j M_{ij}^{-1} f_{reco!ptcl}^j f_{t\bar{t},l+jets}(N_{reco}^j - N_{reco,bgnd}^j)$$

*events that pass particle & reco*  
*events that pass particle*

*events that pass reco & part*  
*events that pass reco*

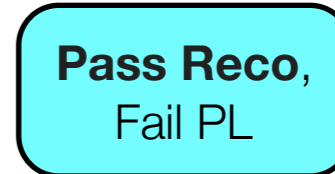
**Unfold** at level of

**restore**

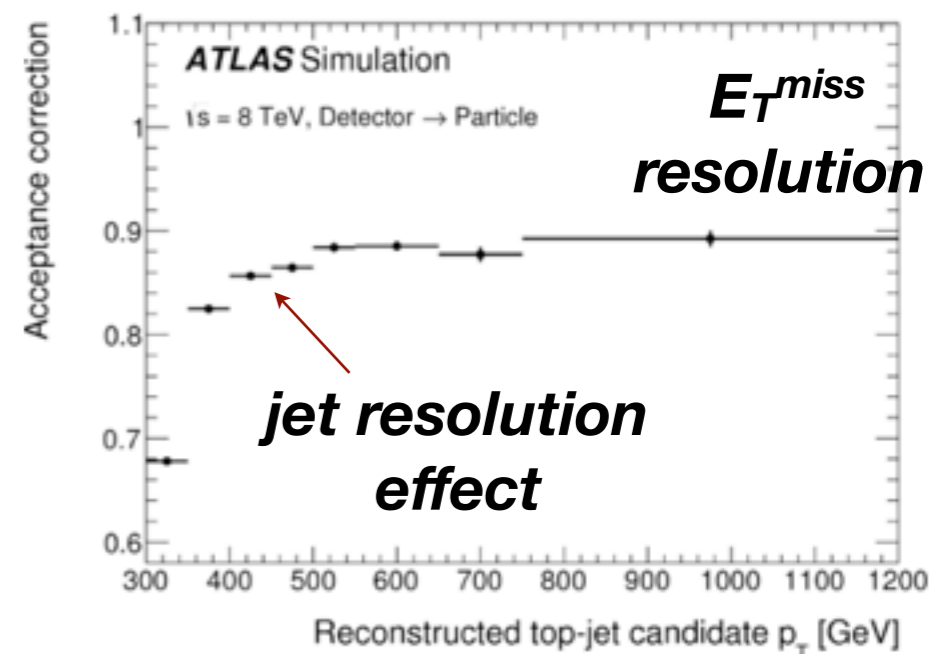
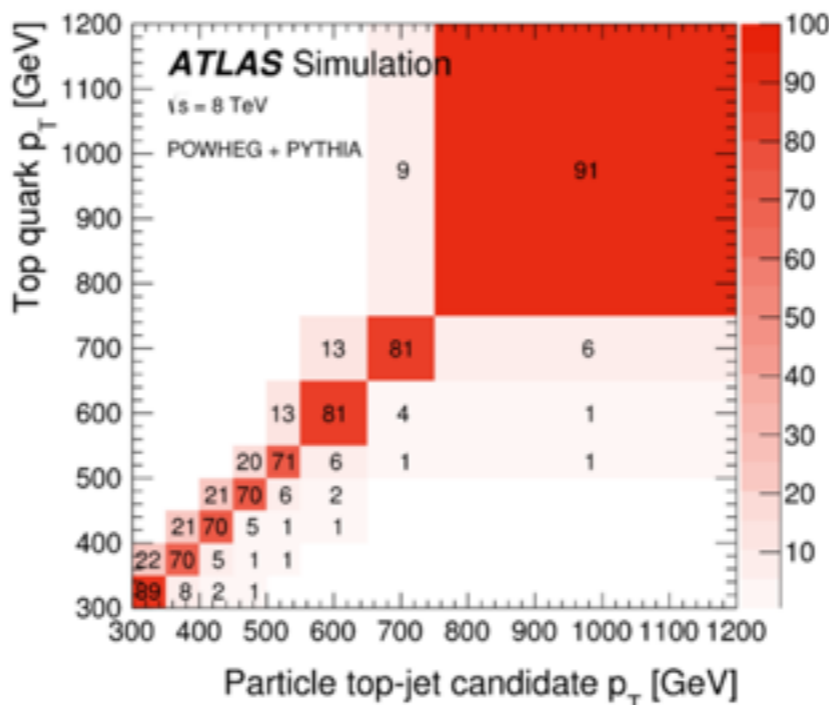
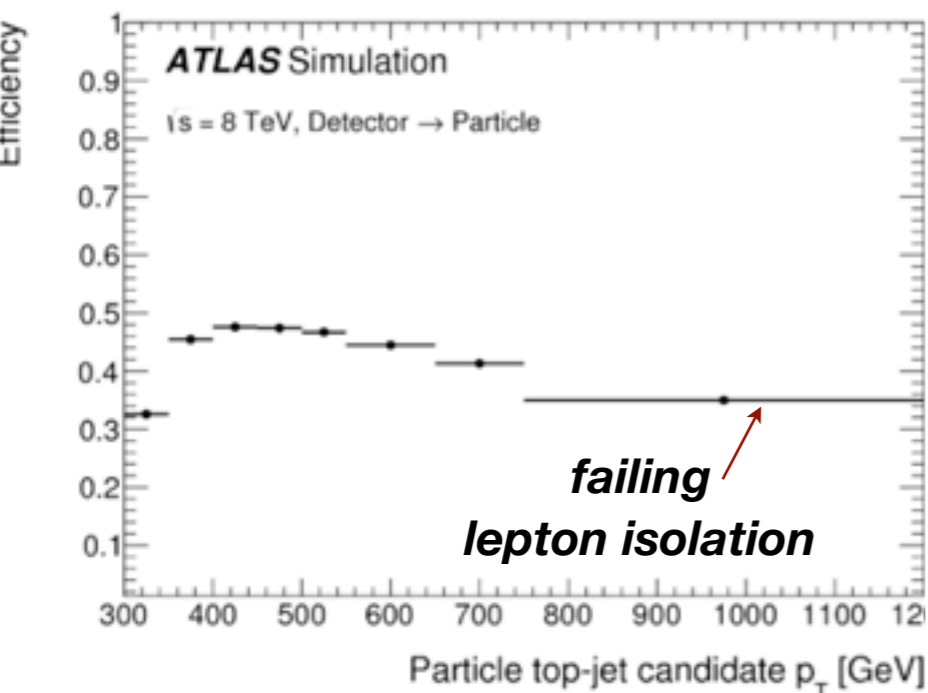


**Pass Reco, Pass PL**

**take**



**away**

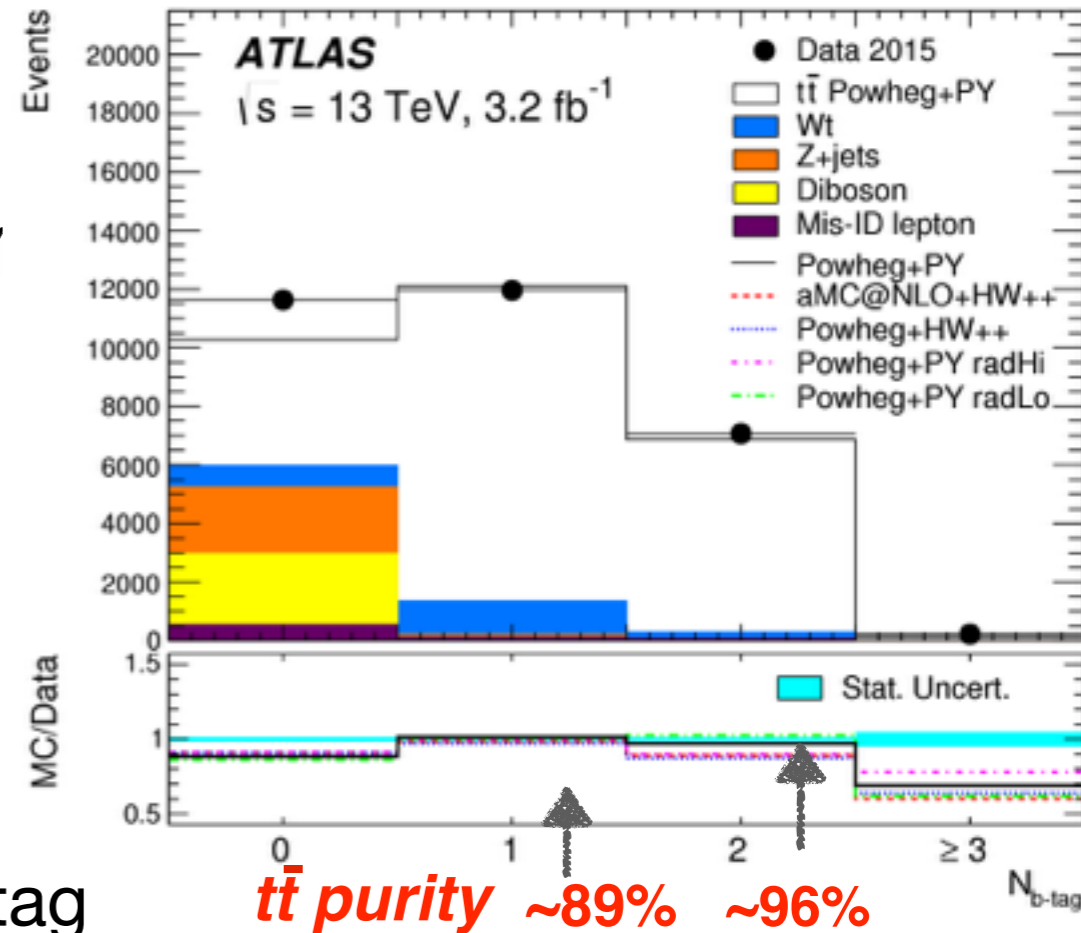
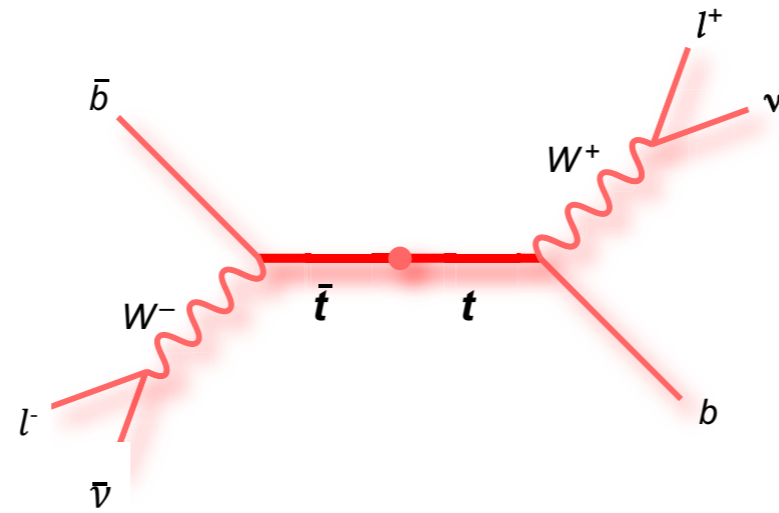


# Special reasons: measure $\sigma_{t\bar{t}}$ - dilepton @ $\sqrt{s} = 13$ TeV

Phys. Lett. B761 (2016) 136

## example from ATLAS

- Select **opposite sign  $e\mu$** , minimal use of jet/ $E_T^{miss}$  info
- **Bkg: mostly single top (Wt)** data-driven fake leptons Z +jets



- Simultaneously **fit for  $\sigma_{t\bar{t}}$  and  $\epsilon_b$** , efficiency to select, reconstruct and recognize b-jet in 1-b-tag and 2-b-tag samples  $\rightarrow$  **minimize jet & b-tag syst**

from simulation

$$N_1 = \mathcal{L} \sigma_{t\bar{t}} \epsilon_{e\mu} 2\epsilon_b (1 - C_b \epsilon_b) + N_1^{bkg}$$

$$N_2 = \mathcal{L} \sigma_{t\bar{t}} \epsilon_{e\mu} C_b \epsilon_b^2 + N_2^{bkg}$$

- Dominated by “External” Systematic effects: Luminosity (1.5%), Beam (2.3%) energy
- tt Modelling uncertainty reduced from parton (~3%) to particle (~2%)

- Extrapolate to **particle (called fid) & parton level**

$$\sigma_{t\bar{t}} = 818 \pm 8 \text{ (stat)} \pm 27 \text{ (syst)} \pm 19 \text{ (lumi)} \pm 12 \text{ (beam)} \text{ pb} \quad \delta\sigma_{t\bar{t}}/\sigma_{t\bar{t}} \sim 4.4 \%$$

$$\sigma_{t\bar{t}}^{\text{fid}} = 11.32 \pm 0.10 \text{ (stat)} \pm 0.29 \text{ (syst)} \pm 0.26 \text{ (lumi)}$$

$$\delta\sigma_{t\bar{t}}/\sigma_{t\bar{t}} \sim 3.9$$

NNLO+NNL prediction : 832 +40-46 pb



Table 1: The  $\chi^2$  values (taking into account data uncertainties and ignoring theoretical uncertainties) and dof of the measured cross sections with respect to the predictions of various MC generators.

Cross section variables	dof	$\chi^2$		
		'POW+PYT'	'POW+HER'	'MG5+PYT'
$[y(t), p_T(t)]$	15	57	18	35
$[M(t\bar{t}), y(t)]$	15	26	18	36
$[M(t\bar{t}), y(t\bar{t})]$	15	28	17	23
$[M(t\bar{t}), \Delta\eta(t, \bar{t})]$	11	66	68	124
$[M(t\bar{t}), \Delta\phi(t, \bar{t})]$	15	14	18	10
$[M(t\bar{t}), p_T(t\bar{t})]$	15	21	22	29
$[M(t\bar{t}), p_T(t)]$	15	77	34	68
$[N_{\text{jet}}^{0,1+}, M(t\bar{t}), y(t\bar{t})]$	23	34	31	34
$[N_{\text{jet}}^{0,1,2+}, M(t\bar{t}), y(t\bar{t})]$	35	50	66	63

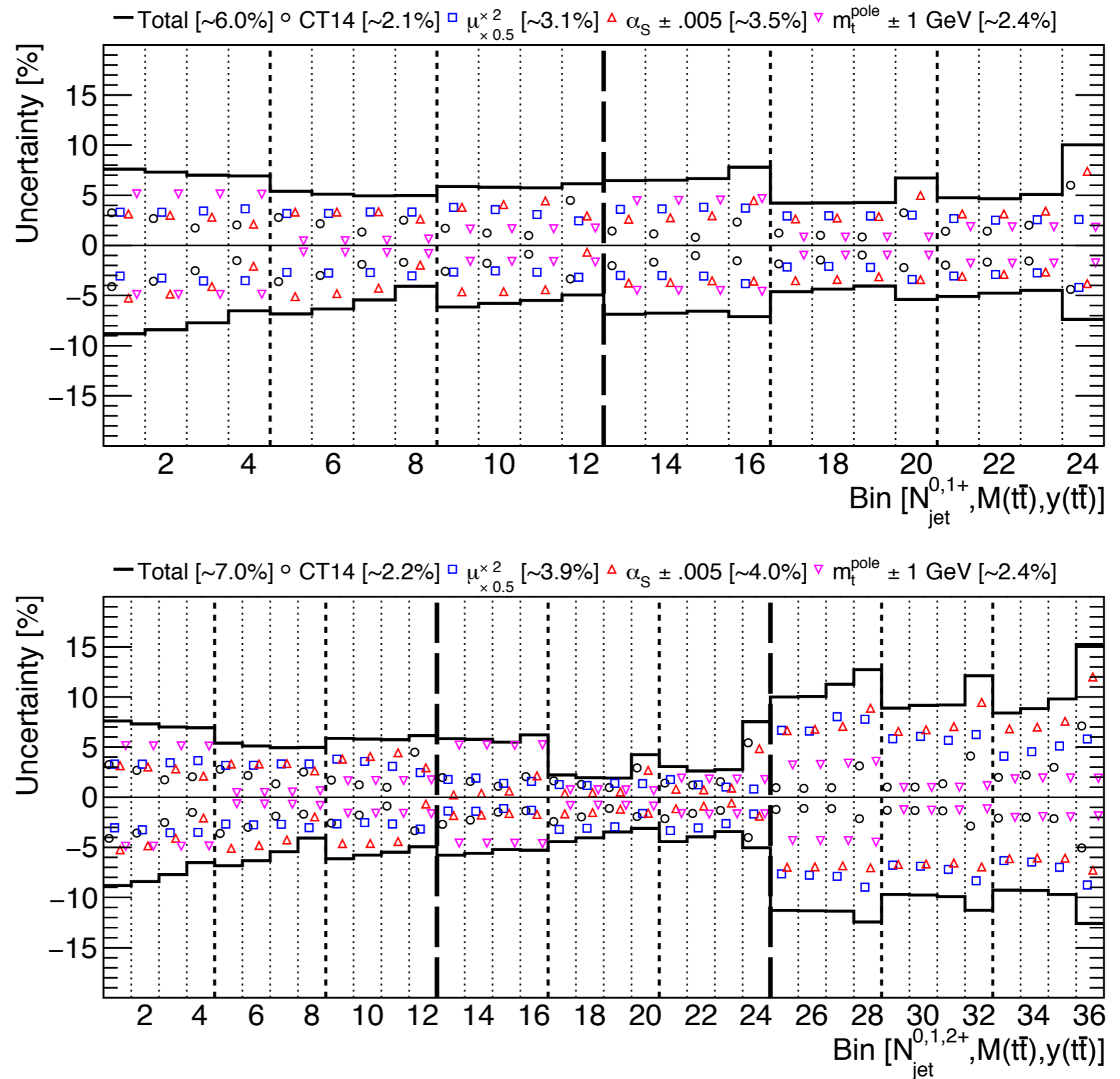
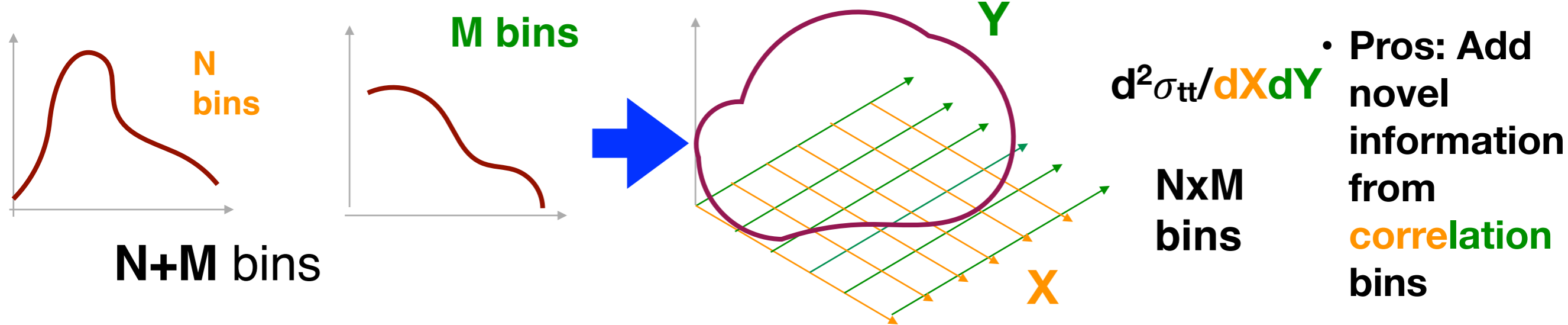


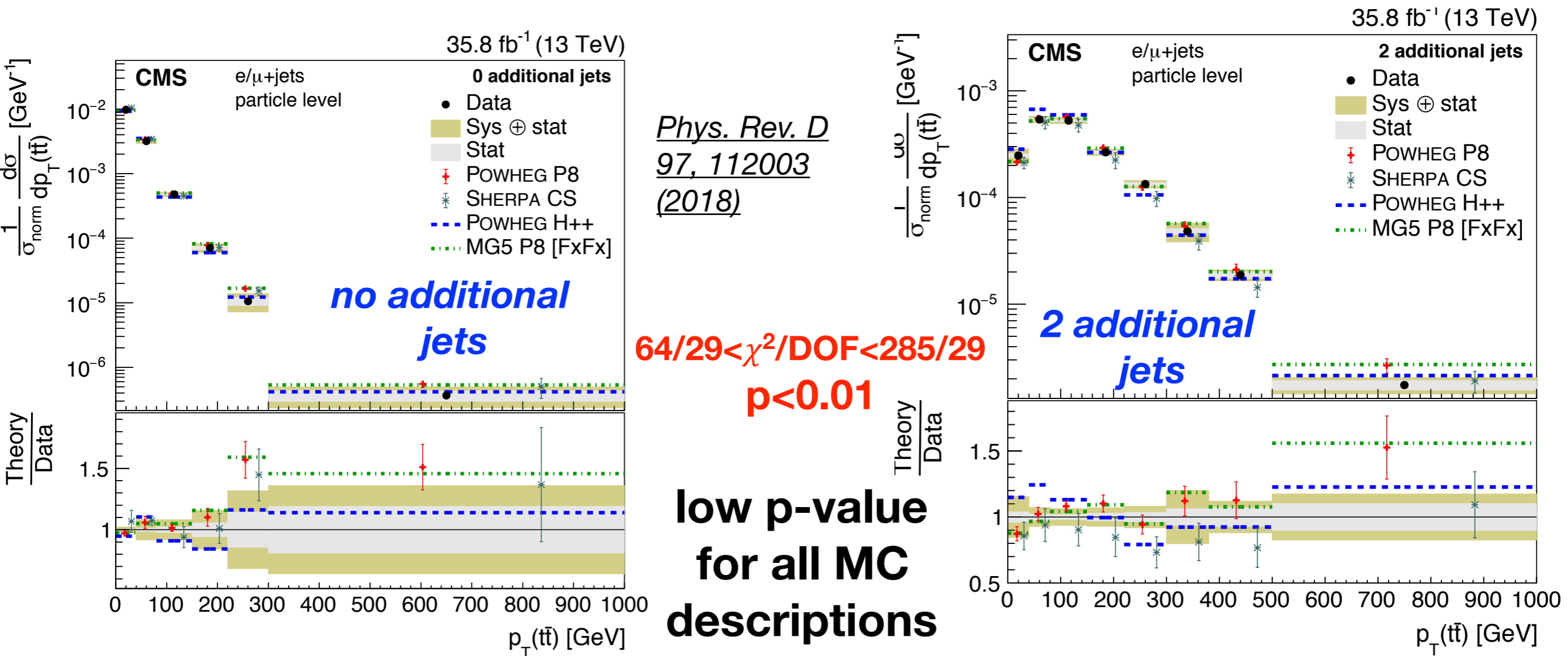
Figure 13: The theoretical uncertainties for  $[N_{\text{jet}}^{0,1+}, M(t\bar{t}), y(t\bar{t})]$  (upper) and  $[N_{\text{jet}}^{0,1,2+}, M(t\bar{t}), y(t\bar{t})]$  (lower) cross sections, arising from PDF,  $\alpha_S(m_Z)$ , and  $m_t^{\text{pole}}$  variations, as well as the total theoretical uncertainties, with their bin-averaged values shown in brackets. The bins are the same as in Figs. 10 and 11.

# 1 Extreme test of SM: Double differential cross section @13 TeV

“Multiply” number of new regions!



global inconsistency in  $p_{T,tt}$  in bins of  $N_{jet}$  @ particle level vs NLO+PS

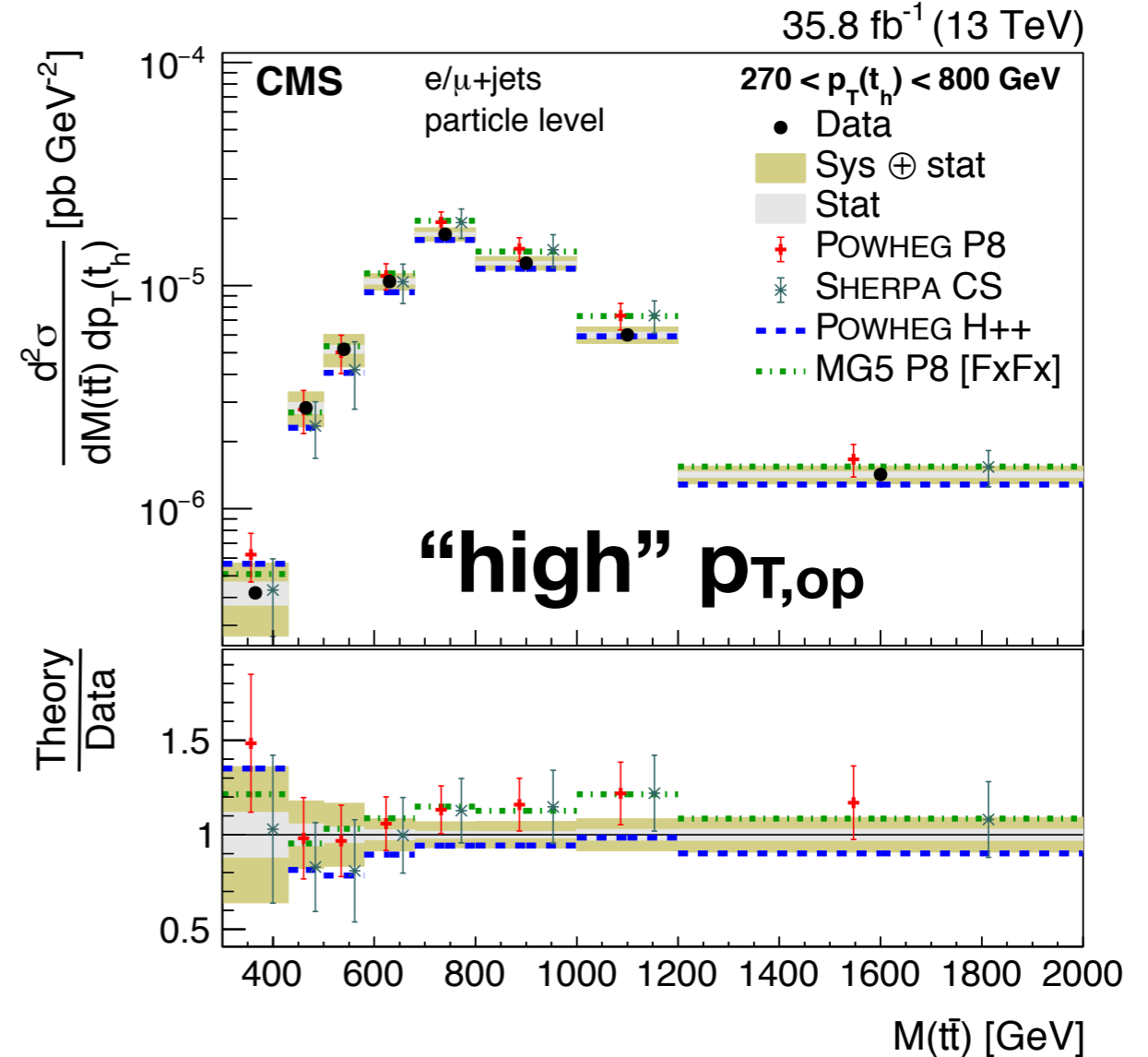
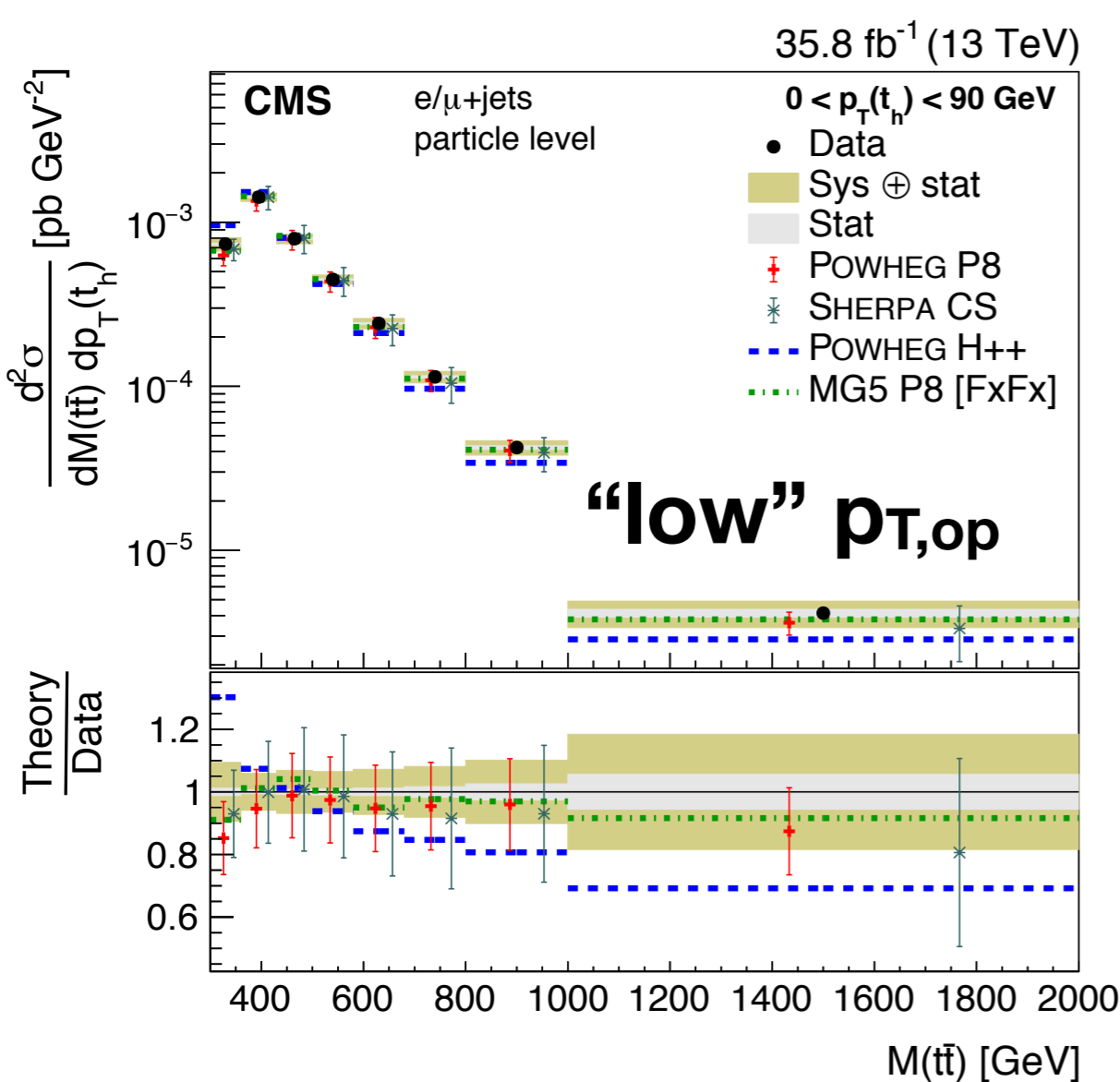


# 1 Extreme test of SM: Double differential cross section @13 TeV

@ particle level

Phys. Rev. D 97, 112003 (2018)

check  $m_{t\bar{t}}$  consistency as Lorentz boost increases (larger  $p_{T,op}$ )



- “Inconsistency” reduced to “tension” by including theory uncertainties
- same @ parton

N <sub>dof</sub> = 32	no theory unc.		with theory unc.	
	χ <sup>2</sup>	p-val	χ <sup>2</sup>	p-val
NLO+PS				
PW+PY8	73.2	<0.01	47.4	0.039
SHERPA	66.5	<0.01	57.2	<0.01
PW+H++	152	<0.01		
MG5+PY	48.9	0.028		

*Eur. Phys. J. C 77 (2017) 804*

**NLO predictions**  
(MCFM +App/GRID)  
**with free gluon PDF parameters**

## Measurements

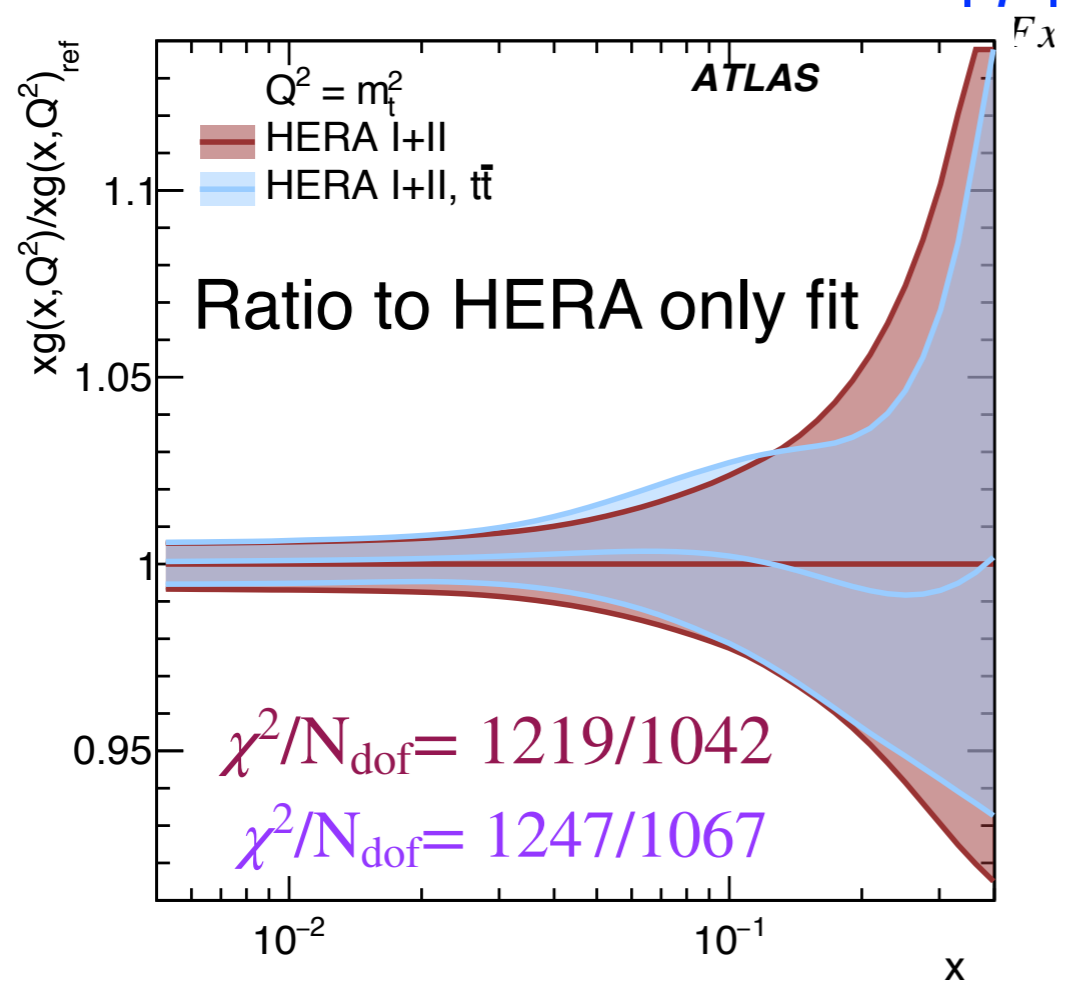
- HERA DIS
- LHC measurements in NNPDF/CT14

with & w/o ATLAS  
 $1/\sigma_{tt} d\sigma_{tt}/dX$  with  $X = |\eta^\ell|, y^{e\mu}, E^{e+} E^\mu$

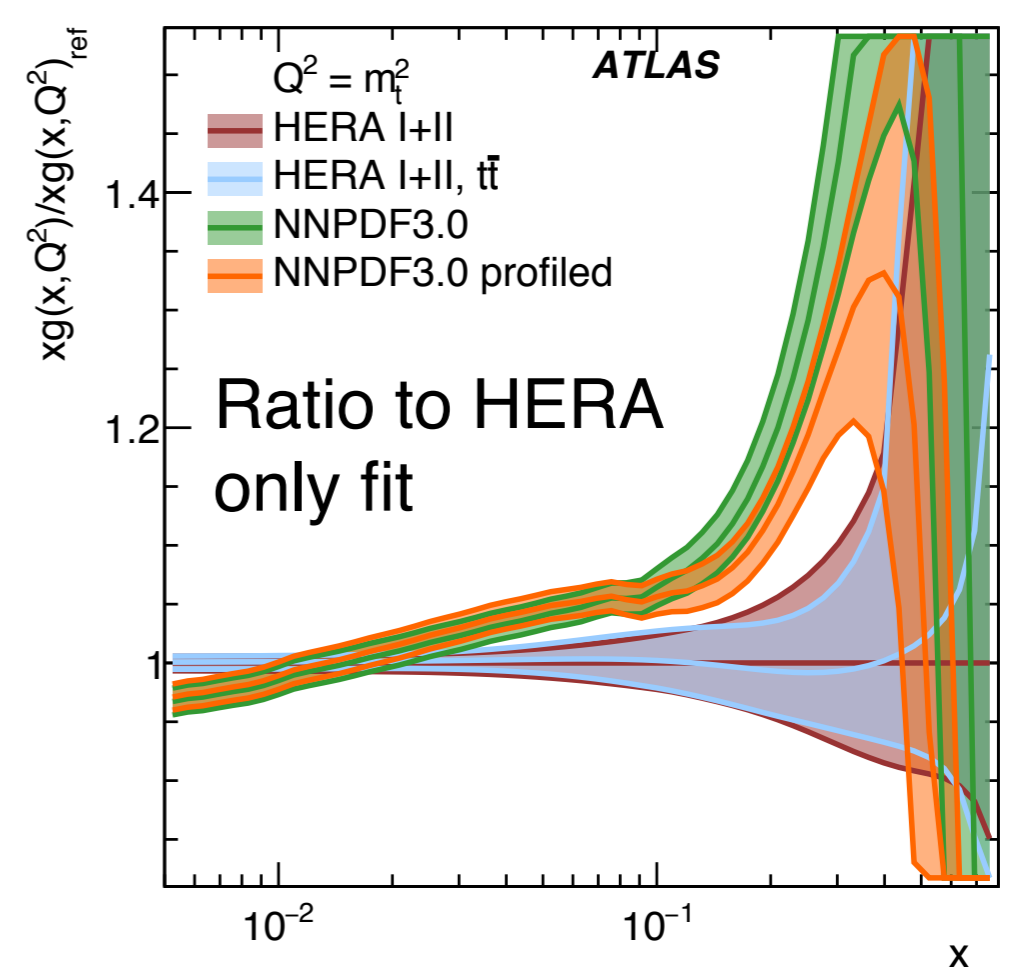


**extract gluon PDF parameters**

•  $\chi^2$  Fit (xFitter)



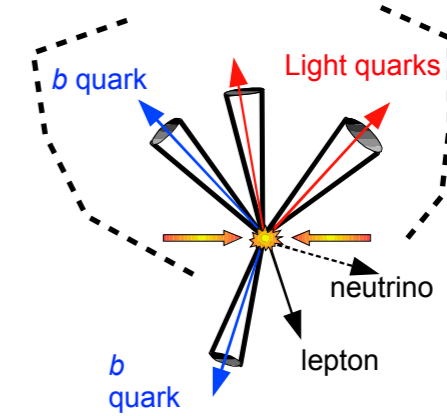
**Adding tt info to HERA reduces uncertainty by 10 to 25%**



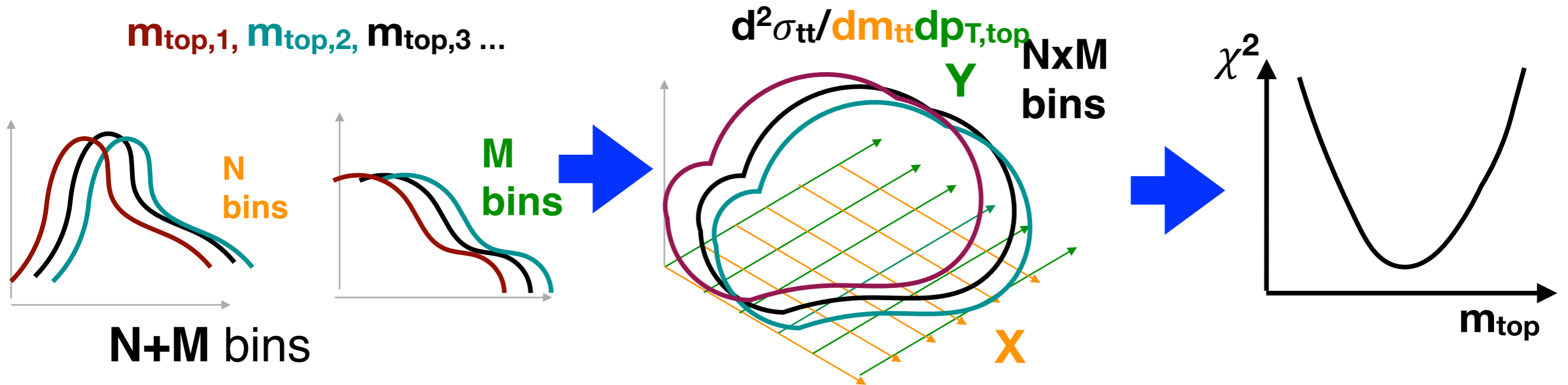
**Including tt info in NNPDF/CT14 lowers gluon pdf at high x**

# Looking at the future: top quark mass from $d\sigma_{tt}/dXdY$

Measure  $d\sigma_{tt}/dXdY$  in  $l+jets$



- $\sigma_{tt}$  depends on top quark mass ( $m_{top}$ ) → Measure  $m_{top}$  using  $d\sigma_{tt}/dm_{tt}dp_{T,top}$



- **Pros:** Add novel information from **correlation bins** + use most precise predictions
- **Cons(?):** ...requires more events, larger sensitivity to modelling?
- Minimise  $\chi^2$  of  $d^2\sigma_{tt}/dm_{tt}dp_{T,top}$  VS **NNLO  $m_{top}$ -dependent prediction**

# Extreme test of SM:double differential cross sections

# Extreme test of SM: Double differential cross sections

*Phys. Rev. D 97, 112003 (2018)*

## Absolute @ particle level

kine vs jet multiplicities and final state objects

- **FxFx** has up to 2 partons @NLO: it is **more consistent with 2d diffxsec** even without theory uncertainties

- $p_{T,tt}$  is only observable with global inconsistency

Distribution	$\chi^2/\text{dof}$	p-value	$\chi^2/\text{dof}$	p-value	$\chi^2/\text{dof}$	p-value
	POWHEG+P8 with unc.		SHERPA with unc.		POWHEG+P8	
Additional jets	1.52/6	0.958	27.3/6	<0.01	10.1/6	0.121
Additional jets vs. $p_T(t_h)$	35.1/44	0.830	64.6/44	0.023	71.6/44	<0.01
Additional jets vs. $M(t\bar{t})$	27.5/36	0.845	68.9/36	<0.01	38.8/36	0.345
Additional jets vs. $p_T(t\bar{t})$	64.6/29	<0.01	181/29	<0.01	175/29	<0.01
$p_T(\text{jet})$	70.2/47	0.016	374/47	<0.01	133/47	<0.01
$ \eta(\text{jet}) $	120/70	<0.01	174/70	<0.01	171/70	<0.01
$\Delta R_{jt}$	60.9/66	0.655	215/66	<0.01	168/66	<0.01
$\Delta R_t$	64.0/62	0.405	229/62	<0.01	121/62	<0.01
	SHERPA		POWHEG+H++		MG5_aMC@NLO+P8 FxFx	
Additional jets	63.0/6	<0.01	34.1/6	<0.01	11.1/6	0.086
Additional jets vs. $p_T(t_h)$	88.5/44	<0.01	230/44	<0.01	53.4/44	0.156
Additional jets vs. $M(t\bar{t})$	112/36	<0.01	300/36	<0.01	55.1/36	0.022
Additional jets vs. $p_T(t\bar{t})$	285/29	<0.01	223/29	<0.01	122/29	<0.01
$p_T(\text{jet})$	768/47	<0.01	624/47	<0.01	111/47	<0.01
$ \eta(\text{jet}) $	214/70	<0.01	259/70	<0.01	133/70	<0.01
$\Delta R_{jt}$	334/66	<0.01	959/66	<0.01	67.0/66	0.441
$\Delta R_t$	316/62	<0.01	483/62	<0.01	78.9/62	0.073

theo uncertainties available only for POWHEG+PY & SHERPA

Additional jets = jet multiplicities up to 5 additional jets with  $p_T > 30$  GeV



# Extreme test of SM: Double differential cross sections

*Phys. Rev. D 97, 112003 (2018)*

## Normalised @ particle level

kine vs jet multiplicities and final state objects

- **FxFx** has up to 2 partons @NLO: it is **more consistent with 2d diffxsec** even without theory uncertainties

- $p_{T,t\bar{t}}$  is only observable with global inconsistency

Distribution	$\chi^2/\text{dof}$	$p$ -value	$\chi^2/\text{dof}$	$p$ -value	$\chi^2/\text{dof}$	$p$ -value
	POWHEG+P8 with unc.		SHERPA with unc.		POWHEG+P8	
Additional jets	2.20/5	0.820	26.4/5	<0.01	12.5/5	0.029
Additional jets vs. $p_T(t_h)$	28.6/43	0.955	35.8/43	0.773	69.7/43	<0.01
Additional jets vs. $M(t\bar{t})$	24.5/35	0.908	46.1/35	0.100	38.9/35	0.298
Additional jets vs. $p_T(t\bar{t})$	73.3/28	<0.01	122/28	<0.01	164/28	<0.01
$p_T(\text{jet})$	75.3/46	<0.01	184/46	<0.01	134/46	<0.01
$ \eta(\text{jet}) $	141/69	<0.01	162/69	<0.01	160/69	<0.01
$\Delta R_{jt}$	69.9/65	0.317	157/65	<0.01	173/65	<0.01
$\Delta R_t$	82.2/61	0.036	163/61	<0.01	126/61	<0.01
	SHERPA		POWHEG+H++		MG5_aMC@NLO+P8 FxFx	
Additional jets	62.4/5	<0.01	35.4/5	<0.01	9.31/5	0.097
Additional jets vs. $p_T(t_h)$	79.8/43	<0.01	194/43	<0.01	51.4/43	0.178
Additional jets vs. $M(t\bar{t})$	86.3/35	<0.01	287/35	<0.01	48.2/35	0.068
Additional jets vs. $p_T(t\bar{t})$	282/28	<0.01	232/28	<0.01	112/28	<0.01
$p_T(\text{jet})$	692/46	<0.01	623/46	<0.01	112/46	<0.01
$ \eta(\text{jet}) $	213/69	<0.01	255/69	<0.01	121/69	<0.01
$\Delta R_{jt}$	301/65	<0.01	976/65	<0.01	65.2/65	0.469
$\Delta R_t$	325/61	<0.01	506/61	<0.01	74.7/61	0.112

theo uncertainties available only for POWHEG+PY & SHERPA

Additional jets = jet multiplicities up to 5 additional jets with  $p_T > 30$  GeV

# Extreme test of SM: Double differential cross section @13 TeV

*Phys. Rev. D 97, 112003 (2018)*

**Absolute  
@ parton  
level**

• **POWHEG +PY8  
inconsistency is  
reabsorbed if  
theory  
uncertainties are  
included**

Distribution	$\chi^2/\text{dof}$	$p\text{-value}$	$\chi^2/\text{dof}$	$p\text{-value}$	$\chi^2/\text{dof}$	$p\text{-value}$
	POWHEG+P8 with unc.		POWHEG+P8		NNLO QCD+NLO EW	
$p_T(t_{\text{high}})$	16.4/12	0.173	27.4/12	<0.01		
$p_T(t_{\text{low}})$	22.4/12	0.033	42.7/12	<0.01		
$p_T(t_h)$	16.4/12	0.175	24.0/12	0.020	5.13/12	0.953
$ y(t_h) $	1.28/11	1.000	1.41/11	1.000	2.27/11	0.997
$p_T(t_\ell)$	22.2/12	0.035	38.3/12	<0.01	9.56/12	0.654
$ y(t_\ell) $	2.04/11	0.998	2.42/11	0.996	8.14/11	0.700
$M(t\bar{t})$	7.67/10	0.661	11.6/10	0.314	24.7/10	<0.01
$p_T(t\bar{t})$	5.38/8	0.717	46.5/8	<0.01		
$ y(t\bar{t}) $	3.98/10	0.948	5.66/10	0.843	9.26/10	0.507
$ y(t_h) $ vs. $p_T(t_h)$	23.6/44	0.995	41.6/44	0.577		
$M(t\bar{t})$ vs. $ y(t\bar{t}) $	20.6/35	0.975	35.0/35	0.469		
$p_T(t_h)$ vs. $M(t\bar{t})$	38.9/32	0.188	59.3/32	<0.01		
	POWHEG+H++		MG5_aMC@NLO+P8 FxFx		—	
$p_T(t_{\text{high}})$	6.60/12	0.883	16.3/12	0.180		
$p_T(t_{\text{low}})$	28.5/12	<0.01	15.3/12	0.225		
$p_T(t_h)$	5.09/12	0.955	11.0/12	0.530		
$ y(t_h) $	2.39/11	0.997	2.21/11	0.998		
$p_T(t_\ell)$	6.55/12	0.886	17.4/12	0.136		
$ y(t_\ell) $	2.54/11	0.995	3.99/11	0.970		
$M(t\bar{t})$	4.16/10	0.940	12.1/10	0.275		
$p_T(t\bar{t})$	55.0/8	<0.01	26.8/8	<0.01		
$ y(t\bar{t}) $	11.9/10	0.292	8.92/10	0.540		
$ y(t_h) $ vs. $p_T(t_h)$	57.9/44	0.077	40.2/44	0.634		
$M(t\bar{t})$ vs. $ y(t\bar{t}) $	40.8/35	0.229	58.7/35	<0.01		
$p_T(t_h)$ vs. $M(t\bar{t})$	93.0/32	<0.01	166/32	<0.01		

theory uncertainties  
available only for  
POWHEG+PY8

# Extreme test of SM: Double differential cross section @13 TeV

*Phys. Rev. D 97, 112003 (2018)*

**Normalised  
@ parton  
level**

• **POWHEG +PY8  
inconsistency is  
reabsorbed if  
theory  
uncertainties are  
included**

**consistent  
with absolute**

Distribution	$\chi^2/\text{dof}$	$p\text{-value}$	$\chi^2/\text{dof}$	$p\text{-value}$	$\chi^2/\text{dof}$	$p\text{-value}$
	POWHEG+P8 with unc.		POWHEG+P8		NNLO QCD+NLO EW	
$p_T(t_{\text{high}})$	18.4/11	0.073	24.4/11	0.011		
$p_T(t_{\text{low}})$	16.6/11	0.120	40.0/11	<0.01		
$p_T(t_h)$	16.1/11	0.138	22.9/11	0.018	4.99/11	0.932
$ y(t_h) $	1.25/10	1.000	1.33/10	0.999	2.23/10	0.994
$p_T(t_\ell)$	23.6/11	0.014	33.0/11	<0.01	8.67/11	0.652
$ y(t_\ell) $	2.03/10	0.996	2.29/10	0.994	8.18/10	0.611
$M(t\bar{t})$	7.78/9	0.556	11.3/9	0.259	24.4/9	<0.01
$p_T(t\bar{t})$	5.52/7	0.597	40.9/7	<0.01		
$ y(t\bar{t}) $	3.89/9	0.919	5.36/9	0.802	9.29/9	0.411
$ y(t_h) $ vs. $p_T(t_h)$	22.7/43	0.995	38.8/43	0.654		
$M(t\bar{t})$ vs. $ y(t\bar{t}) $	20.2/34	0.970	33.2/34	0.507		
$p_T(t_h)$ vs. $M(t\bar{t})$	34.4/31	0.309	57.4/31	<0.01		
	POWHEG+H++		MG5_aMC@NLO+P8 FxFx		—	
$p_T(t_{\text{high}})$	4.10/11	0.967	13.2/11	0.283		
$p_T(t_{\text{low}})$	17.4/11	0.096	11.9/11	0.370		
$p_T(t_h)$	3.61/11	0.980	9.95/11	0.535		
$ y(t_h) $	1.63/10	0.998	1.11/10	1.000		
$p_T(t_\ell)$	8.36/11	0.680	16.4/11	0.128		
$ y(t_\ell) $	1.57/10	0.999	2.48/10	0.991		
$M(t\bar{t})$	3.57/9	0.937	7.61/9	0.574		
$p_T(t\bar{t})$	43.4/7	<0.01	20.5/7	<0.01		
$ y(t\bar{t}) $	5.94/9	0.746	4.65/9	0.864		
$ y(t_h) $ vs. $p_T(t_h)$	32.6/43	0.877	27.8/43	0.965		
$M(t\bar{t})$ vs. $ y(t\bar{t}) $	27.2/34	0.788	40.2/34	0.214		
$p_T(t_h)$ vs. $M(t\bar{t})$	67.9/31	<0.01	77.9/31	<0.01		

theory uncertainties  
available only for  
POWHEG+PY8

Distribution	$\chi^2/\text{dof}$	$p\text{-value}$	$\chi^2/\text{dof}$	$p\text{-value}$	$\chi^2/\text{dof}$	$p\text{-value}$
	POWHEG+P8 with unc.		SHERPA with unc.		POWHEG+P8	
$p_T(t_h)$	15.9/12	0.197	7.21/12	0.844	29.5/12	<0.01
$ y(t_h) $	1.96/11	0.999	1.48/11	1.000	2.23/11	0.997
$p_T(t_\ell)$	27.0/12	<0.01	22.3/12	0.034	80.2/12	<0.01
$ y(t_\ell) $	4.55/11	0.951	5.07/11	0.928	4.99/11	0.932
$M(t\bar{t})$	5.83/10	0.829	2.40/10	0.992	9.07/10	0.525
$p_T(t\bar{t})$	4.96/8	0.761	28.9/8	<0.01	41.2/8	<0.01
$ y(t\bar{t}) $	5.93/10	0.821	6.63/10	0.760	8.61/10	0.570
$ y(t_h) $ vs. $p_T(t_h)$	35.7/44	0.810	29.6/44	0.953	64.1/44	0.025
$M(t\bar{t})$ vs. $ y(t\bar{t}) $	25.9/35	0.867	24.2/35	0.914	56.2/35	0.013
$p_T(t_h)$ vs. $M(t\bar{t})$	47.4/32	0.039	57.2/32	<0.01	73.2/32	<0.01
	SHERPA		POWHEG+H++		MG5_aMC@NLO+P8 FxFx	
$p_T(t_h)$	13.5/12	0.335	32.1/12	<0.01	17.4/12	0.137
$ y(t_h) $	2.32/11	0.997	4.89/11	0.936	3.16/11	0.988
$p_T(t_\ell)$	39.4/12	<0.01	21.8/12	0.040	47.7/12	<0.01
$ y(t_\ell) $	5.54/11	0.902	4.04/11	0.969	7.22/11	0.781
$M(t\bar{t})$	2.86/10	0.985	52.8/10	<0.01	5.45/10	0.859
$p_T(t\bar{t})$	68.7/8	<0.01	46.8/8	<0.01	21.3/8	<0.01
$ y(t\bar{t}) $	12.1/10	0.276	18.6/10	0.046	8.13/10	0.616
$ y(t_h) $ vs. $p_T(t_h)$	48.3/44	0.305	116/44	<0.01	44.9/44	0.434
$M(t\bar{t})$ vs. $ y(t\bar{t}) $	41.5/35	0.208	219/35	<0.01	55.7/35	0.014
$p_T(t_h)$ vs. $M(t\bar{t})$	66.5/32	<0.01	152/32	<0.01	48.9/32	0.028

- Absolute @ particle level**
- **POWHEG +PY8 inconsistency is reabsorbed if theory uncertainties are included**
- **FxFx is the most consistent already without theory uncertainties**
- consistent with absolute**

## Normalised @ particle level

• POWHEG +PY8 inconsistency is reabsorbed if theory uncertainties are included

• FxFx is the most consistent already without theory uncertainties

consistent with absolute

Distribution	$\chi^2/\text{dof}$	p-value	$\chi^2/\text{dof}$	p-value	$\chi^2/\text{dof}$	p-value
	POWHEG+P8 with unc.		SHERPA with unc.		POWHEG+P8	
$p_T(t_h)$	14.9/11	0.186	6.99/11	0.800	29.4/11	<0.01
$ y(t_h) $	1.77/10	0.998	1.25/10	1.000	1.90/10	0.997
$p_T(t_\ell)$	25.3/11	<0.01	28.0/11	<0.01	74.0/11	<0.01
$ y(t_\ell) $	4.50/10	0.922	4.88/10	0.899	5.00/10	0.891
$M(t\bar{t})$	5.69/9	0.770	2.17/9	0.989	9.33/9	0.407
$p_T(t\bar{t})$	5.36/7	0.616	12.5/7	0.086	34.8/7	<0.01
$ y(t\bar{t}) $	5.79/9	0.761	6.68/9	0.671	8.48/9	0.486
$ y(t_h) $ vs. $p_T(t_h)$	27.6/43	0.967	32.7/43	0.872	53.8/43	0.126
$M(t\bar{t})$ vs. $ y(t\bar{t}) $	26.5/34	0.817	22.7/34	0.931	54.0/34	0.016
$p_T(t_h)$ vs. $M(t\bar{t})$	42.5/31	0.082	39.2/31	0.149	64.8/31	<0.01
	SHERPA		POWHEG+H++		MG5_aMC@NLO+P8 FxFx	
$p_T(t_h)$	13.9/11	0.238	34.1/11	<0.01	15.2/11	0.173
$ y(t_h) $	1.60/10	0.999	3.81/10	0.955	2.73/10	0.987
$p_T(t_\ell)$	37.3/11	<0.01	25.0/11	<0.01	40.5/11	<0.01
$ y(t_\ell) $	5.28/10	0.872	3.92/10	0.951	5.54/10	0.853
$M(t\bar{t})$	2.99/9	0.965	51.7/9	<0.01	4.98/9	0.836
$p_T(t\bar{t})$	59.4/7	<0.01	43.8/7	<0.01	17.9/7	0.013
$ y(t\bar{t}) $	11.3/9	0.253	18.2/9	0.033	8.37/9	0.498
$ y(t_h) $ vs. $p_T(t_h)$	47.7/43	0.287	108/43	<0.01	40.9/43	0.561
$M(t\bar{t})$ vs. $ y(t\bar{t}) $	37.6/34	0.308	234/34	<0.01	55.5/34	0.011
$p_T(t_h)$ vs. $M(t\bar{t})$	63.2/31	<0.01	126/31	<0.01	43.0/31	0.074

theory uncertainties available only for POWHEG+PY8 & SHERPA

## Parton distributions functions and top quark : the connection (II)

In c.m. frame of pp system, the parton momentum components are written

$$p_1^\mu = \frac{\sqrt{s}}{2}(x_1, 0, 0, x_1) \quad p_2^\mu = \frac{\sqrt{s}}{2}(x_2, 0, 0, -x_2) . \quad (\text{assume massless partons})$$

Using conservation of momentum & partons in the leading order picture for  $ff \rightarrow t\bar{t}$ , the rapidity of the top quark pair is

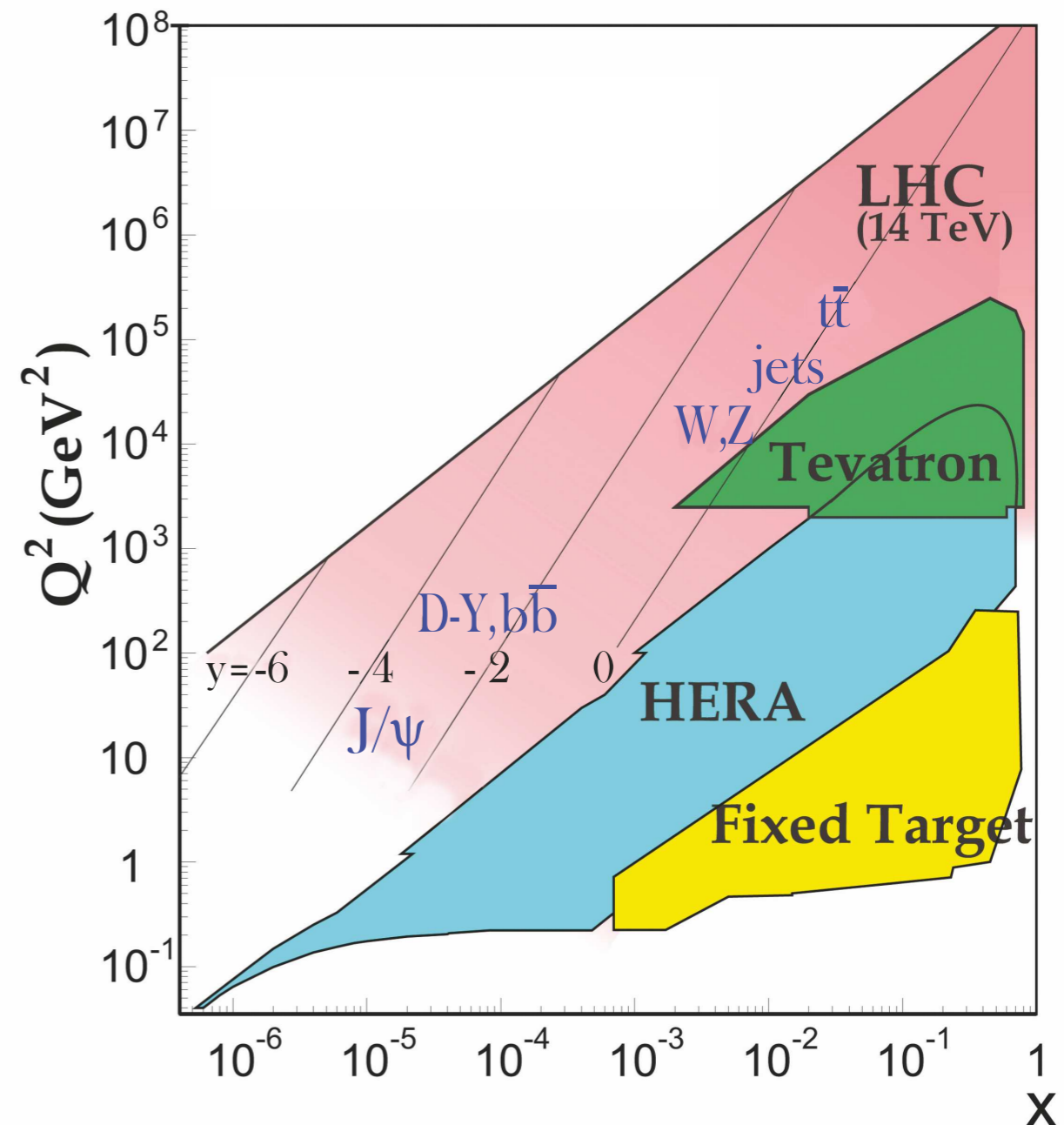
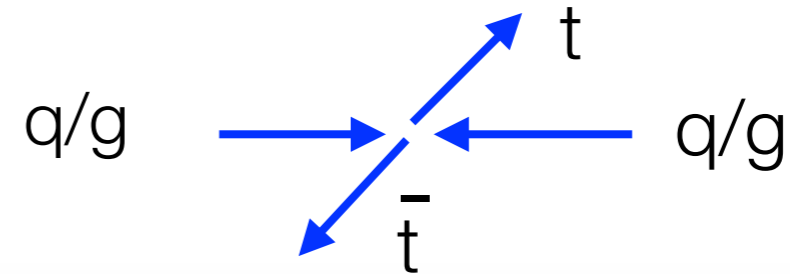
$$y(t\bar{t}) = \frac{1}{2} \ln \left( \frac{E(t\bar{t}) + p_z(t\bar{t})}{E(t\bar{t}) - p_z(t\bar{t})} \right) = \frac{1}{2} \ln \left( \frac{x_1}{x_2} \right)$$

$$\hat{s} = x_1 x_2 s = M(t\bar{t})^2$$



$$x_1 = \frac{M(t\bar{t})}{\sqrt{s}} e^y \quad x_2 = \frac{M(t\bar{t})}{\sqrt{s}} e^{-y}$$

- **maximum probed x at 13 TeV is ~0.25** given the bins of the rapidity and tt mass distributions ranges



# **The Top $p_T$ saga, including the new chapter @13 TeV**

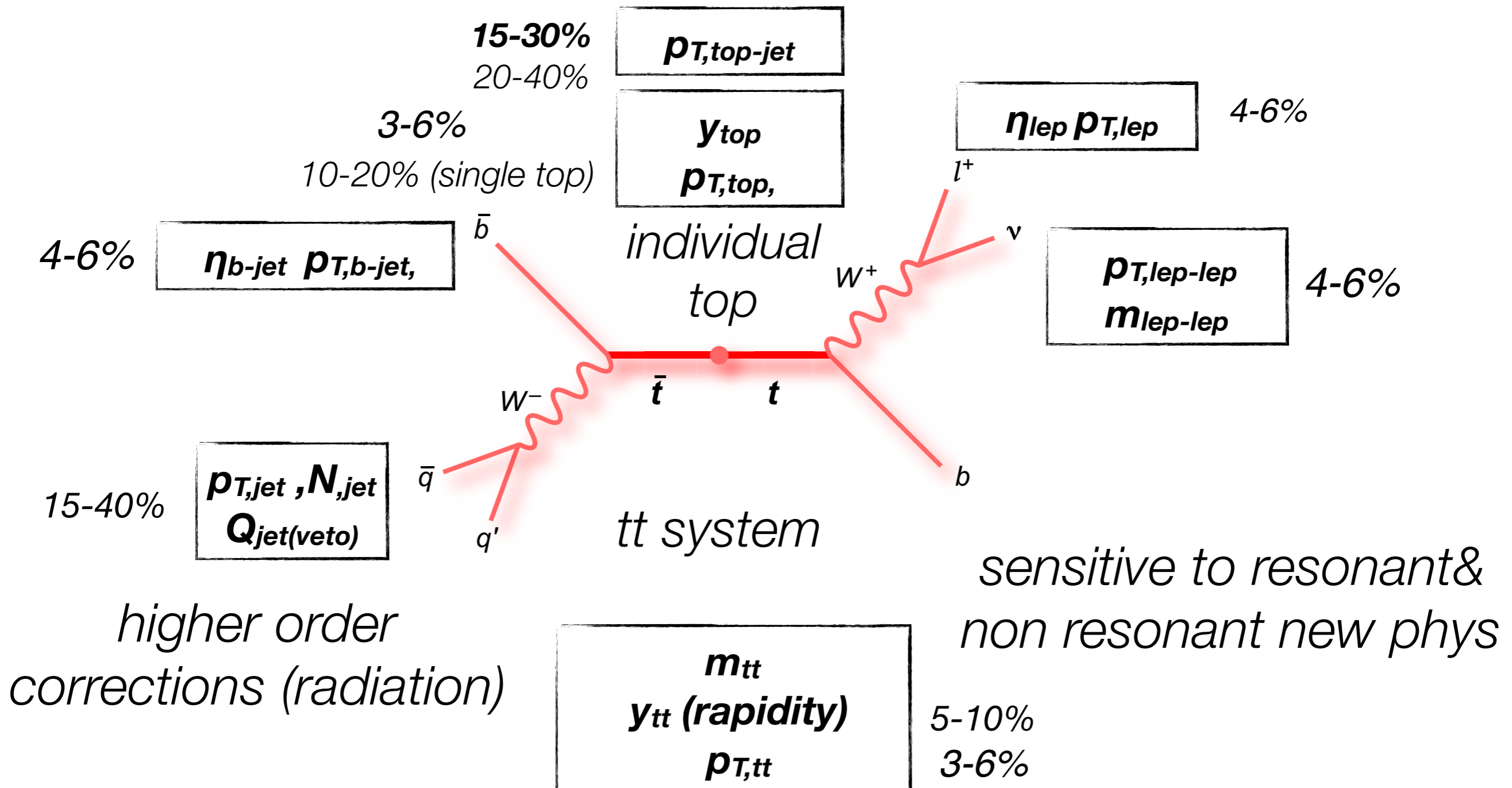
# What we measured at $\sqrt{s}=7,8$ TeV: **$X$ in $d\sigma_{(tt/t)}/dX$**



## Overview of current results at LHC



- Increasing variety of differential cross section results
  - More measurements in fiducial PS, exploiting particle-level object definition and pseudo-top
  - Pioneering results in boosted regime, first absolute differential cross sections appearing





parton level

# ATLAS vs CMS vs NNLO Theory: $1/\sigma_{tt} d\sigma_{tt}/dp_{T,top} @ \sqrt{s} = 8 \text{ TeV}$

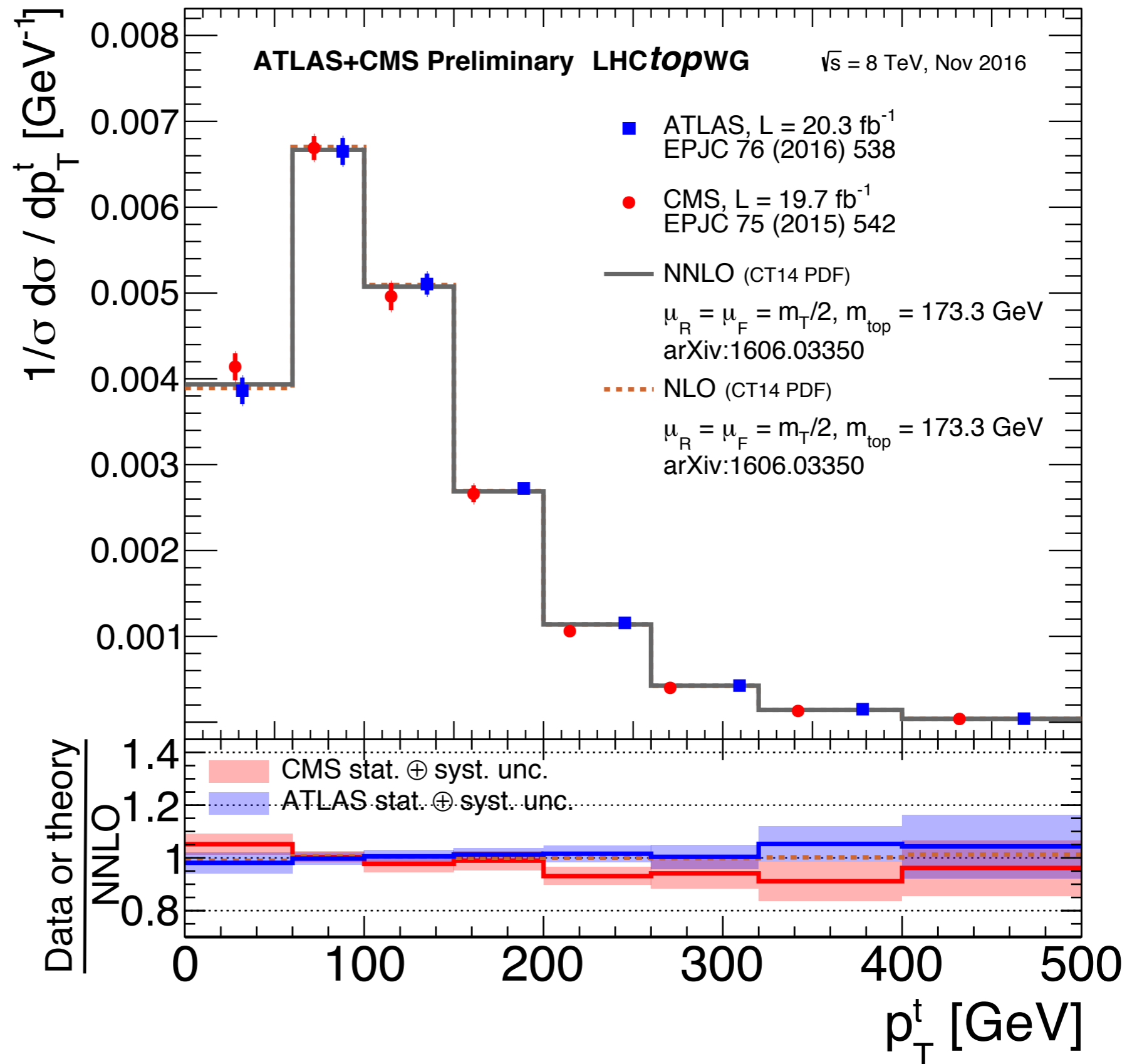
## *parton level*

- **ATLAS & CMS measurements are generally consistent with each other**

- CMS shows slight slope

- *Using latest predictions with dynamic factorisation & renormalization scale*

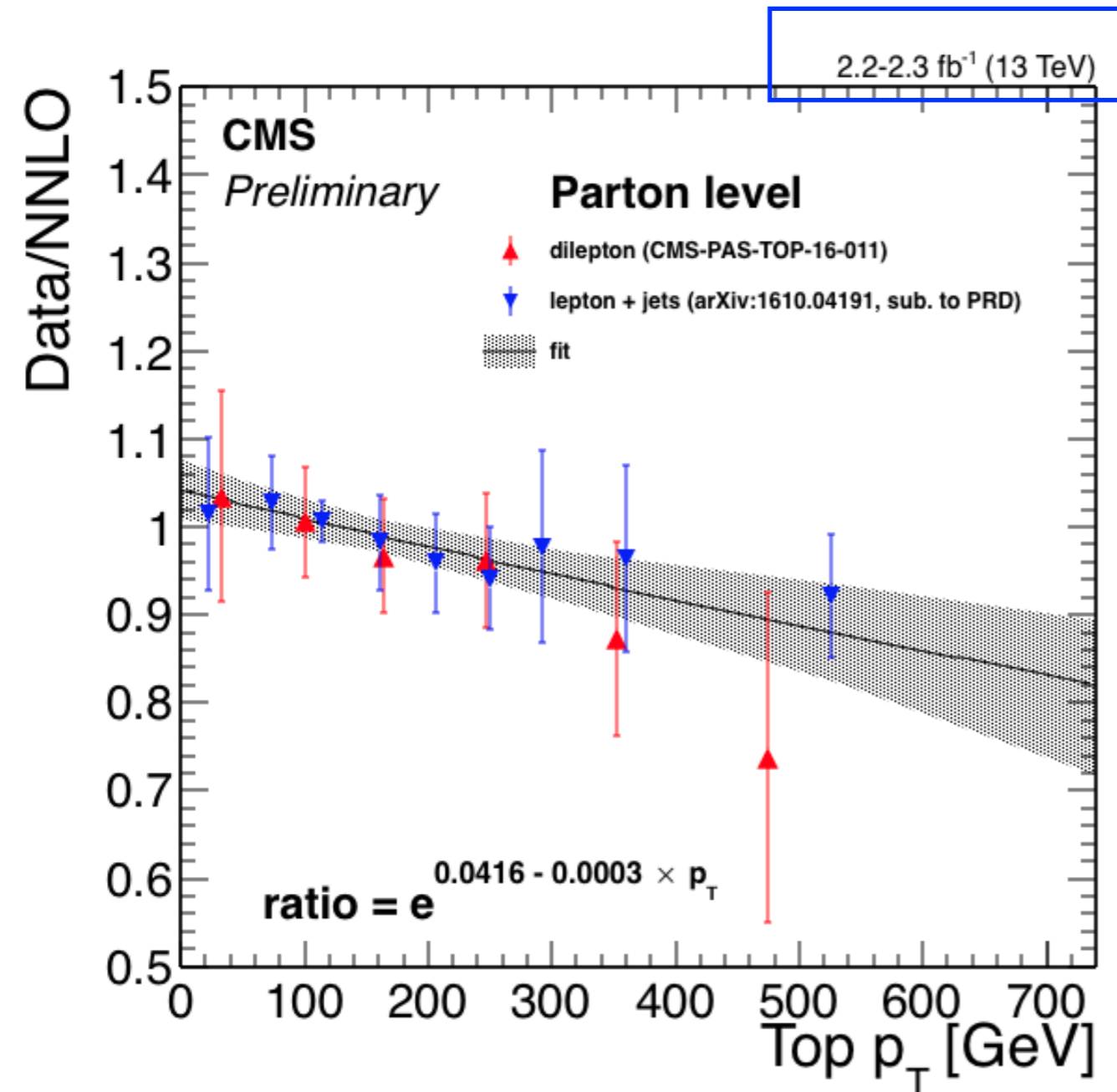
*Qualitative statement, no statistical test performed yet*



# The Top $p_T$ saga: NLO+PS @13 TeV

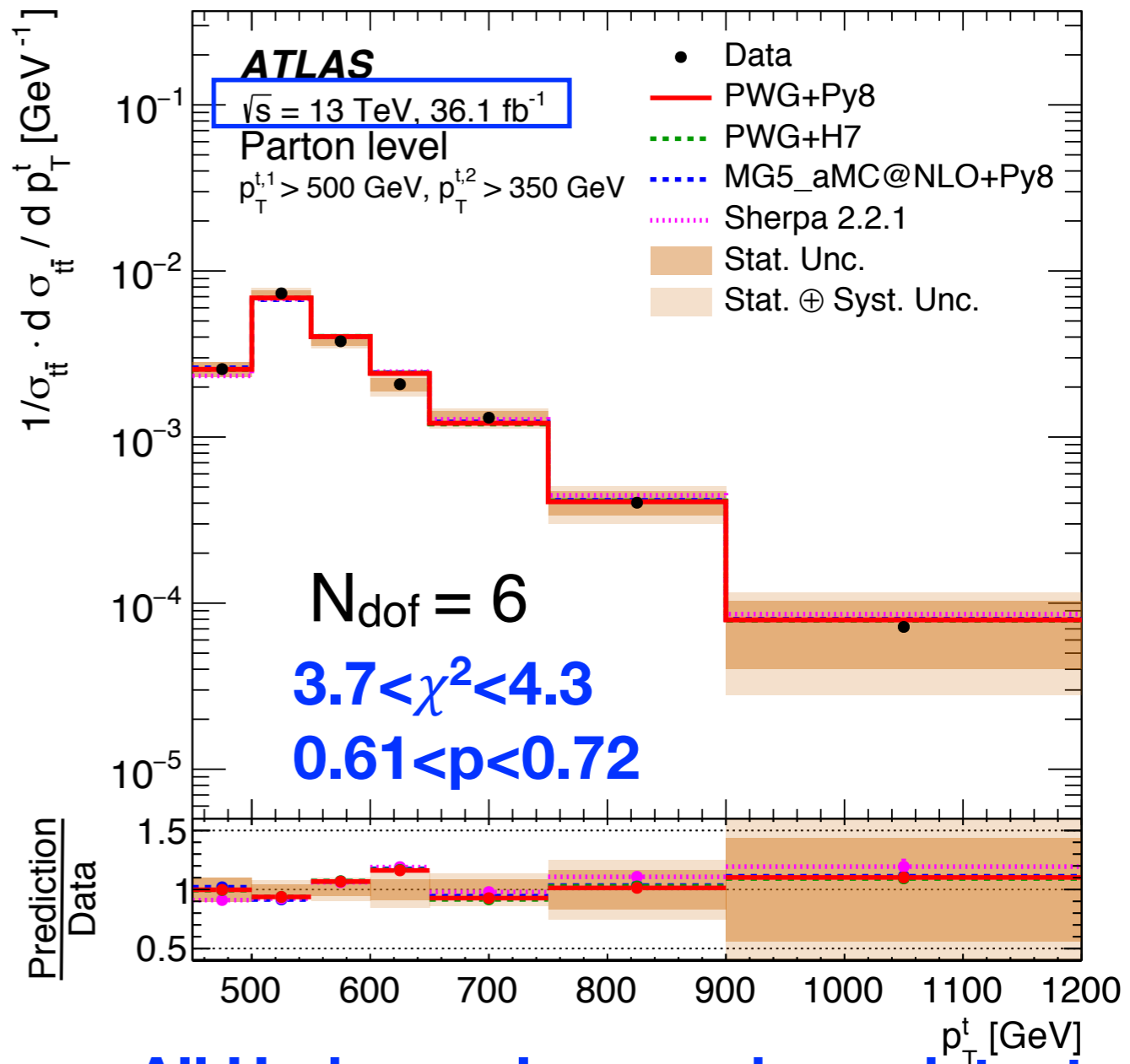
$p_{T,top}$

## Parton level vs NLO+PS



$\ell$ +jets and dilepton “resolved”

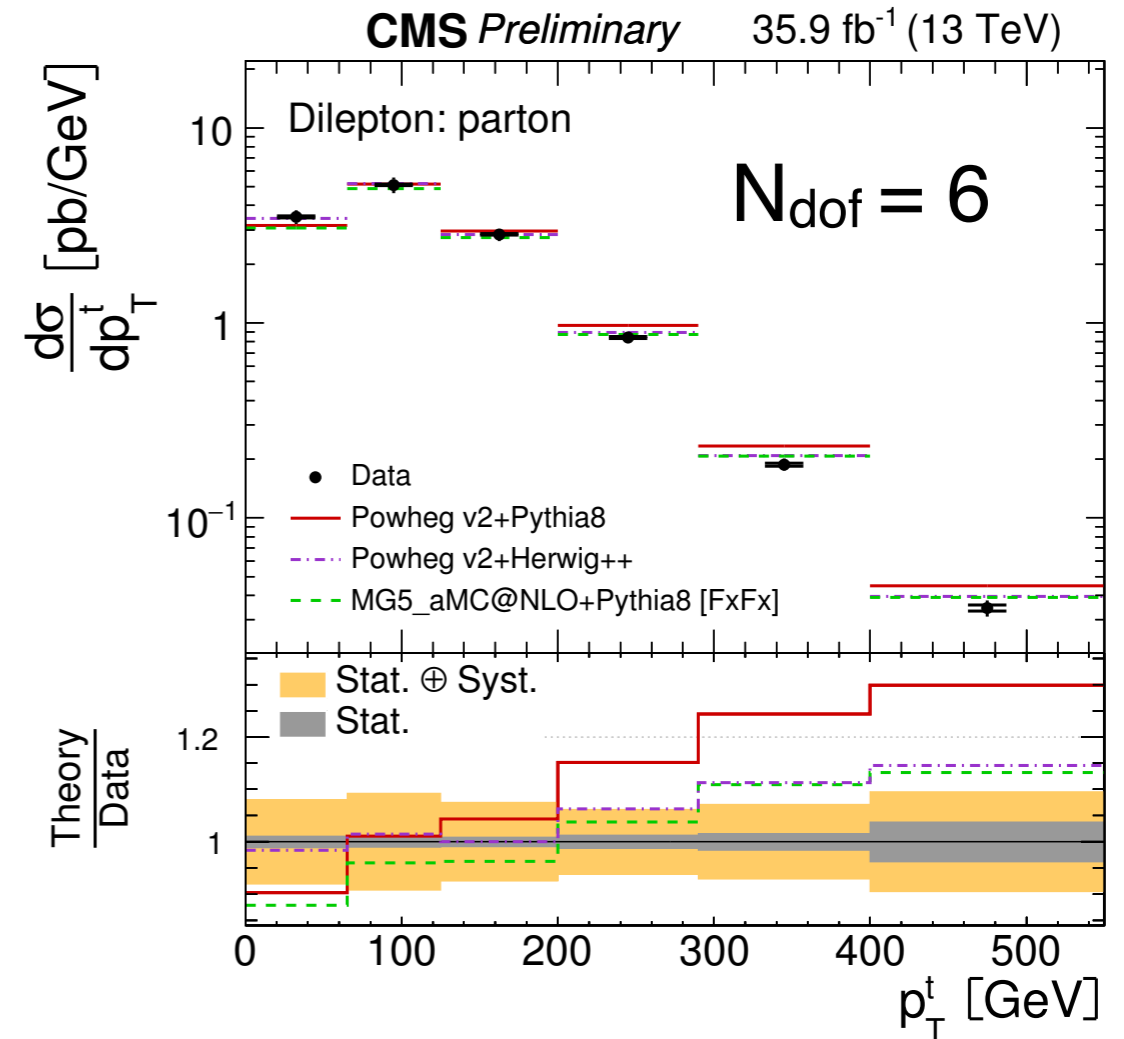
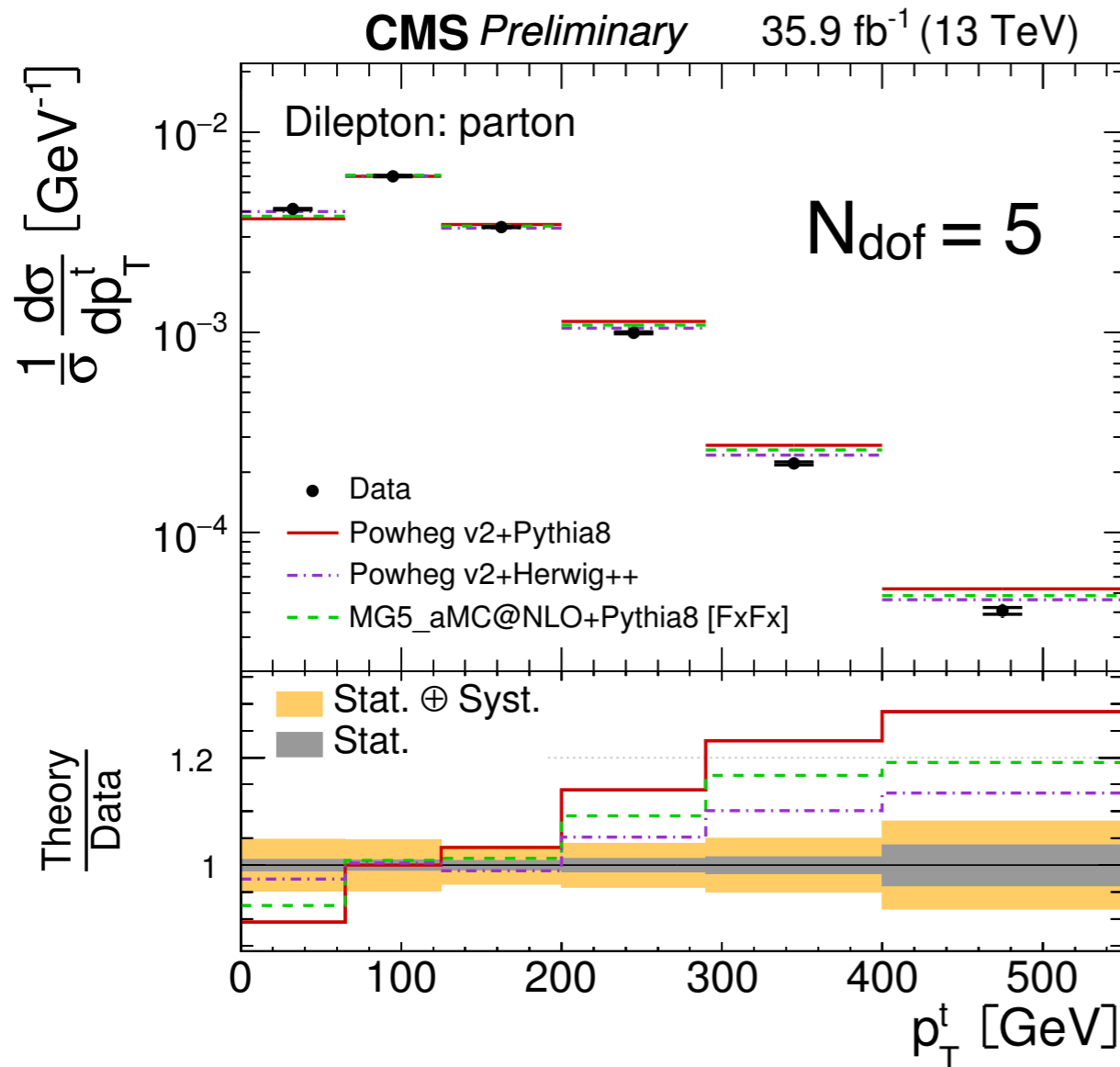
softer, but compatible with standard model within uncertainties  
measurement and theory.



**All Had : no slope and consistent  
larger stat uncertainties  
no theory uncertainty**

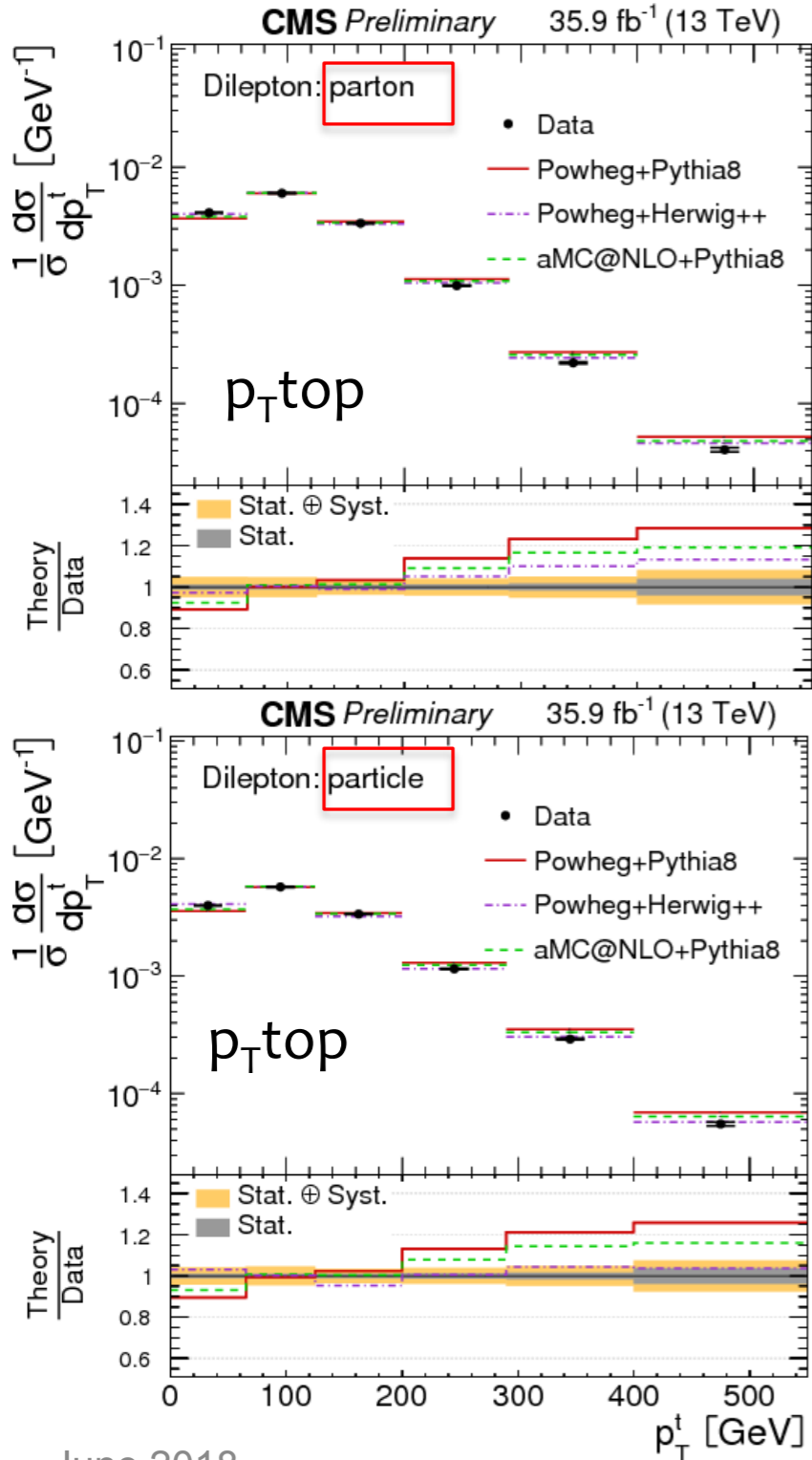
# The Top $p_T$ saga: dilepton NLO+PS @13 TeV

CMS-PAS-TOP-17-014



$N_{\text{dof}} = 5 \text{ o } 6$	normalised		absolute	
NLO+PS	$\chi^2$	p-val	$\chi^2$	p-val
PW+PY8	<b>63.5</b>	<b>&lt;10<sup>-3</sup></b>	<b>51</b>	<b>&lt;10<sup>-3</sup></b>
PW+H++	<b>10</b>	<b>0.087</b>	<b>8</b>	<b>0.239</b>
MG5+PY8	<b>28</b>	<b>&lt;10<sup>-3</sup></b>	<b>18</b>	<b>0.007</b>

**no theory  
uncertainties  
included**



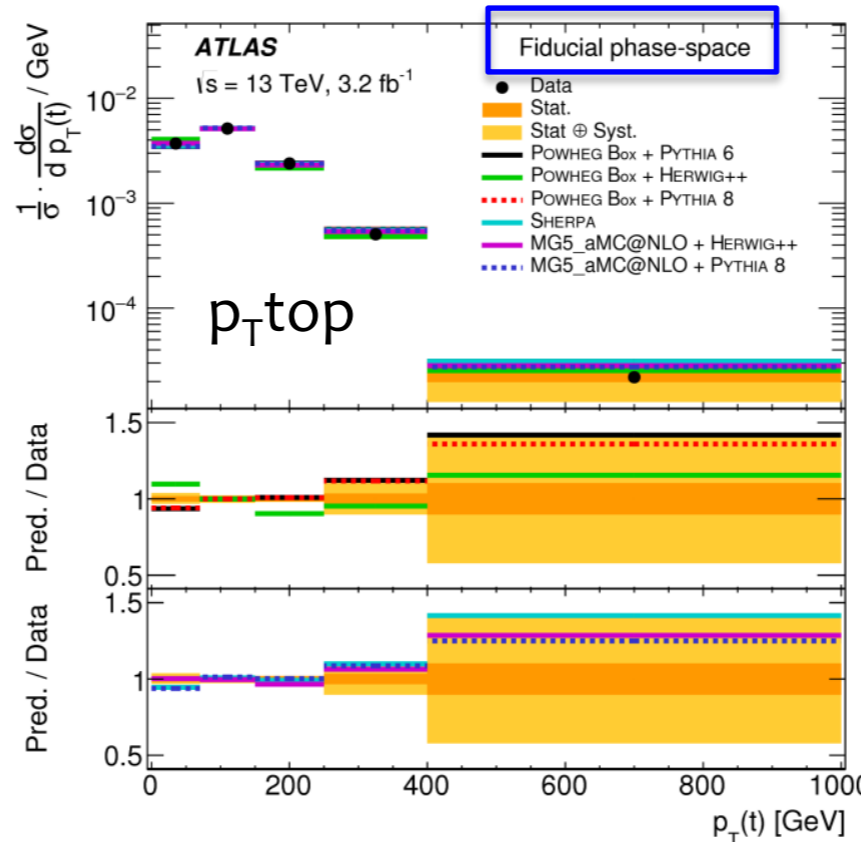
- 13TeV  $e\mu$  35.9fb<sup>-1</sup> (2016)
- Parton level in the full phase space
- Particle level, within a phase space close to experimental acceptance (fiducial phase-space)

**NEW**

- **PowhegV2+Pythia8** (NLO) chosen as default generator setup in ATLAS and CMS for Run2:

❖ Reasonable agreement except in top quark “direct” observables  $p_T$ ,  $p_T^{tt}$ ,  $m_{tt}$

❖ Need for full NNLO MC + PS predictions



Javier Fdez.

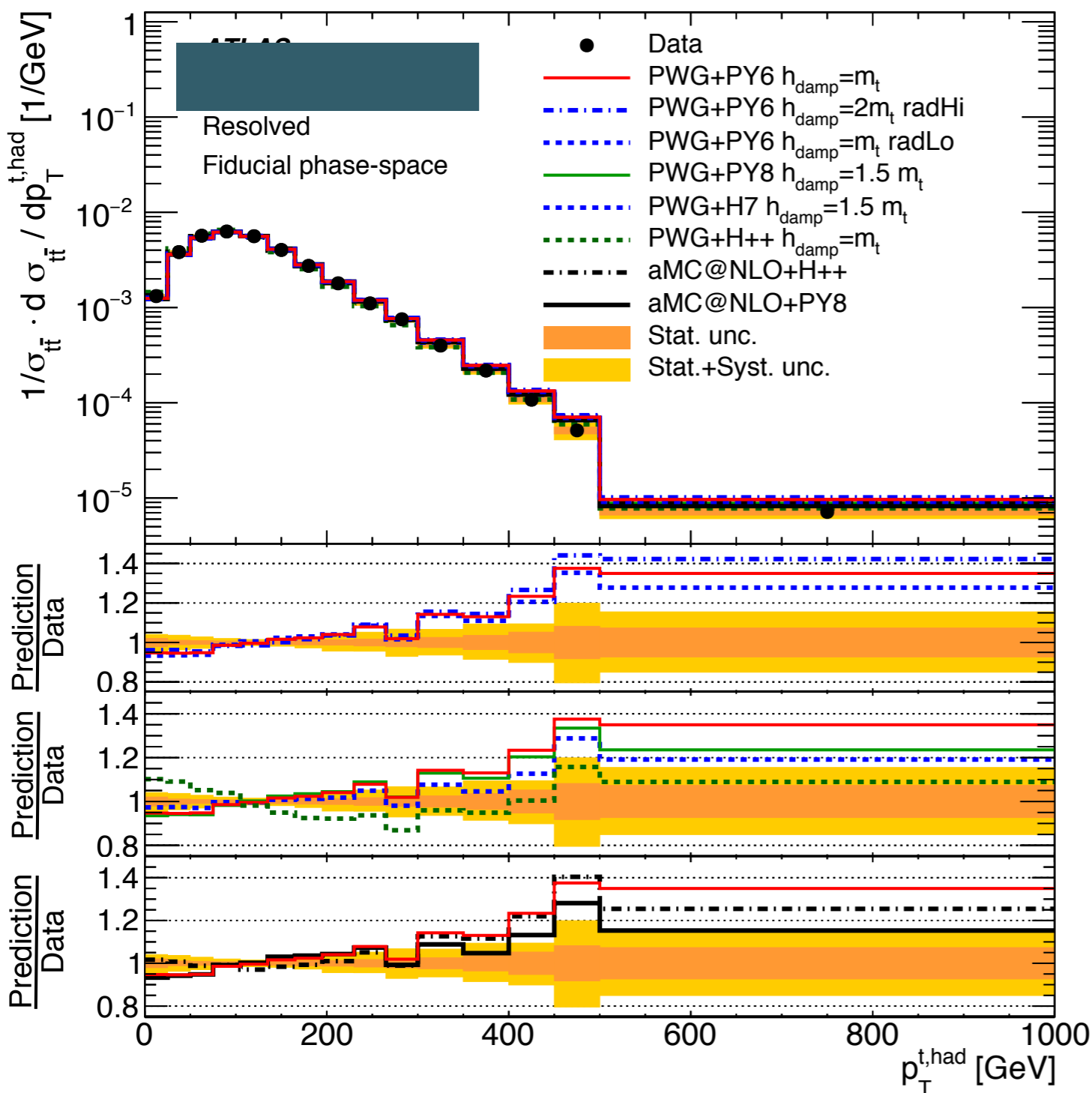
- Similar results in  $e\mu$  ATLAS: [Eur. Phys. J. C77 \(2017\) 299](#)

particle level

## Particle level vs NLO+PS

$1/\sigma_{tt} d\sigma_{tt}/dX$

*JHEP 11 (2017) 191*



$\ell$ +jets “resolved”

$N_{dof} = 14$

$21.5 < \chi^2 < 24.4$

$0.03 < p < 0.09$

Different radiation

Different hadronization

Different generator & hadronization

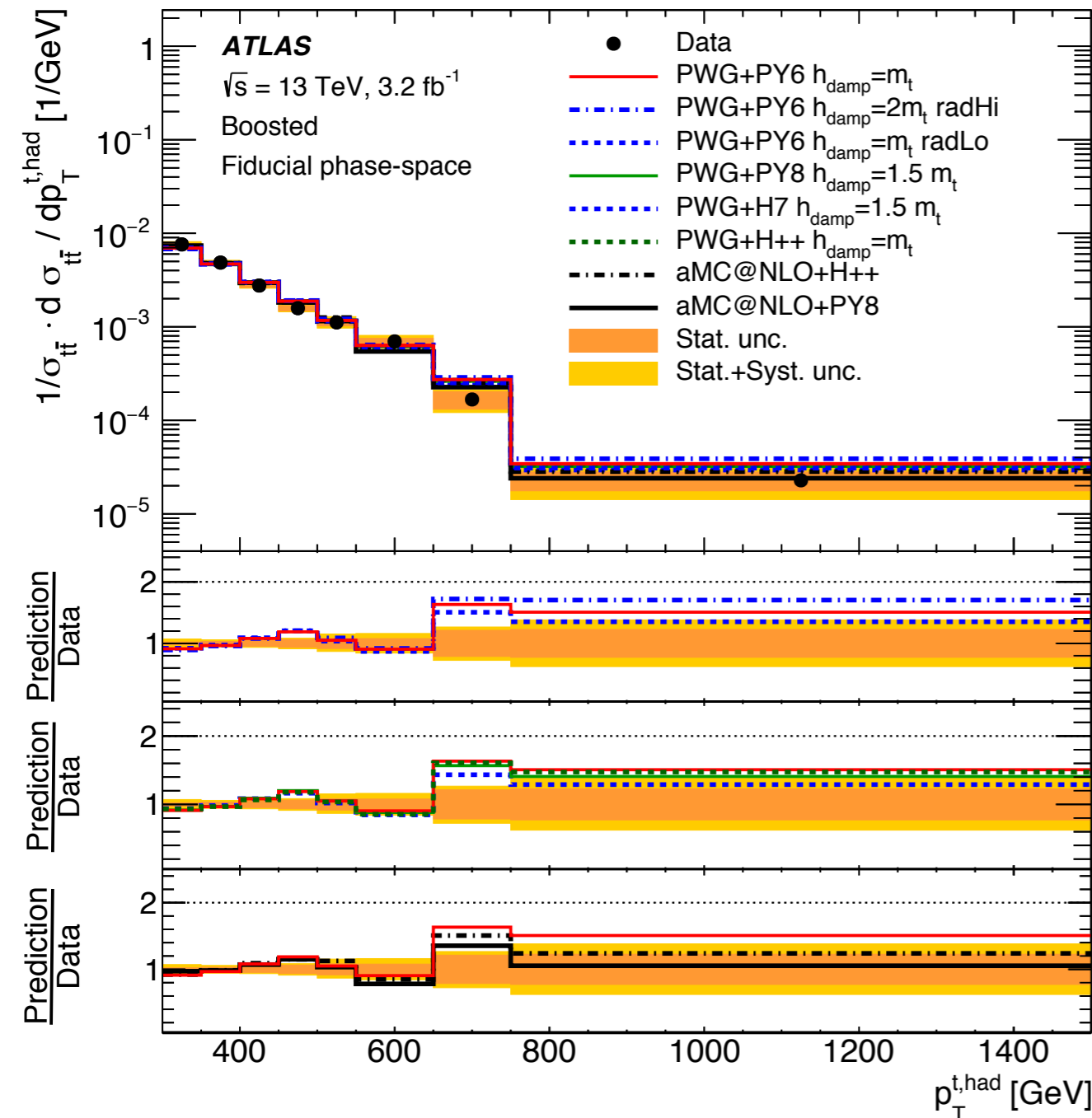
except for  
POWHEG+HERWIG

$\chi^2=15.4$   $p=0.35$  (-)

• Tension with most NLO+PS predictions: slope w.r.t measurement

## Particle level vs NLO+PS

$1/\sigma_{tt} d\sigma_{tt}/dX$  *JHEP 11 (2017) 191*



$\ell$ +jets “boosted”

Different  
Radiation

$$N_{\text{dof}} = 7$$

Different  
Hadronization

$$9.9 < \chi^2 < 11.5$$

$$0.12 < p < 0.20$$

Different  
generator  
& hadronization

- Tension with most NLO+PS predictions: slope w.r.t measurement.
- Less tension than in “resolved”, larger statistical uncertainties



# Extreme test of SM: the top $p_T$ “saga”

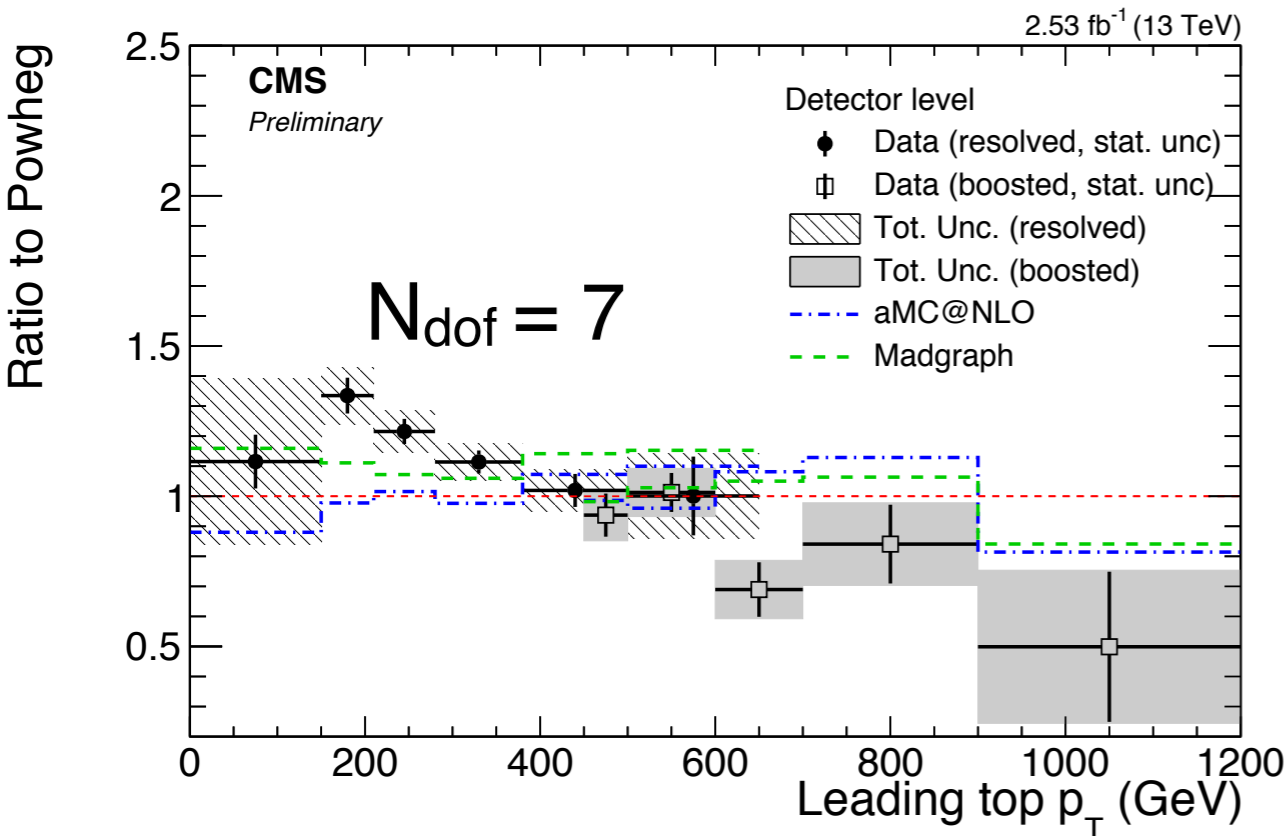
@13TeV

## Particle level vs NLO+PS

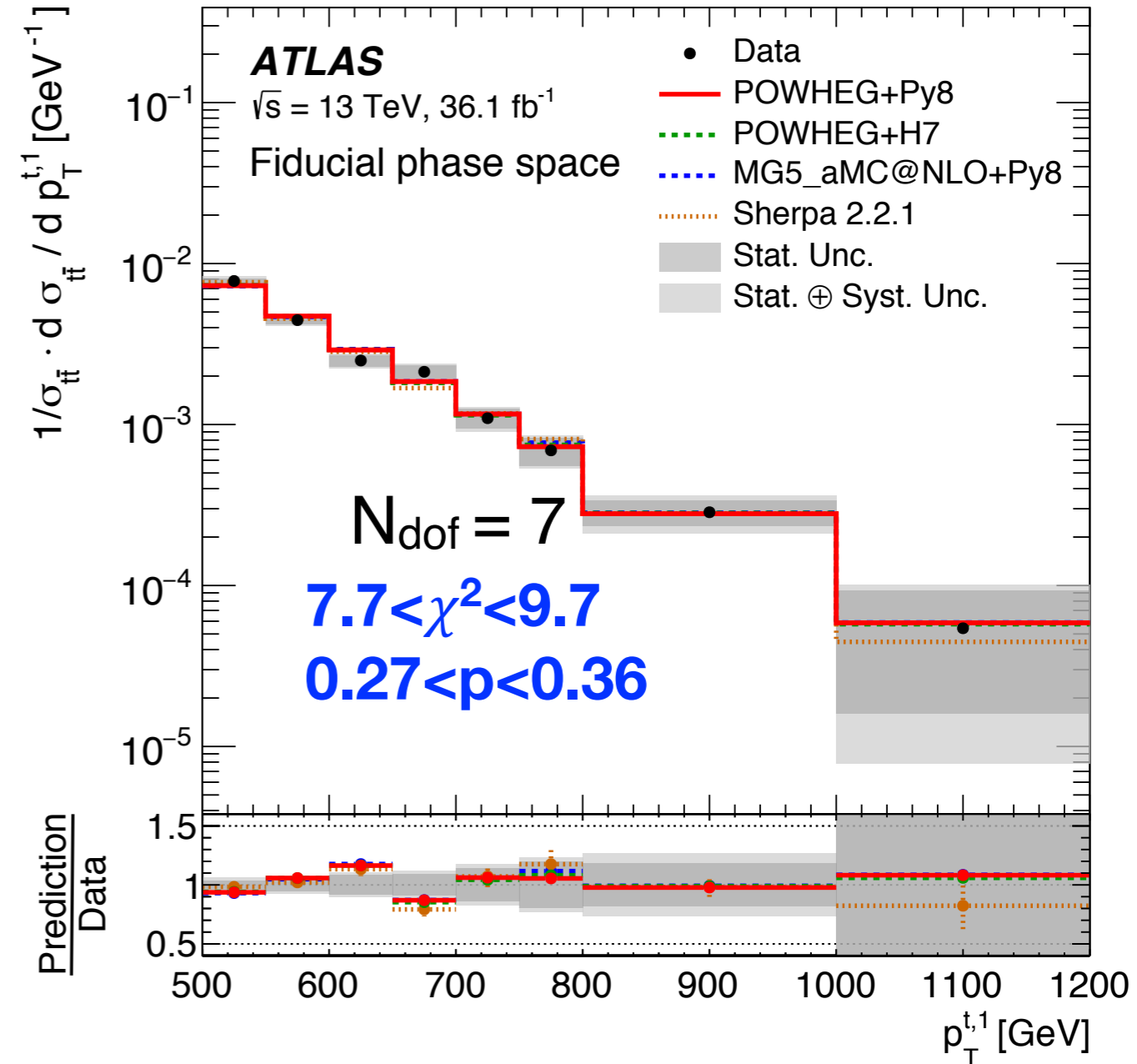
$p_{T,lead\ top}$

### Reco level

CMS-TOP-16-013,  $2.5\text{ fb}^{-1}$ , 13 TeV



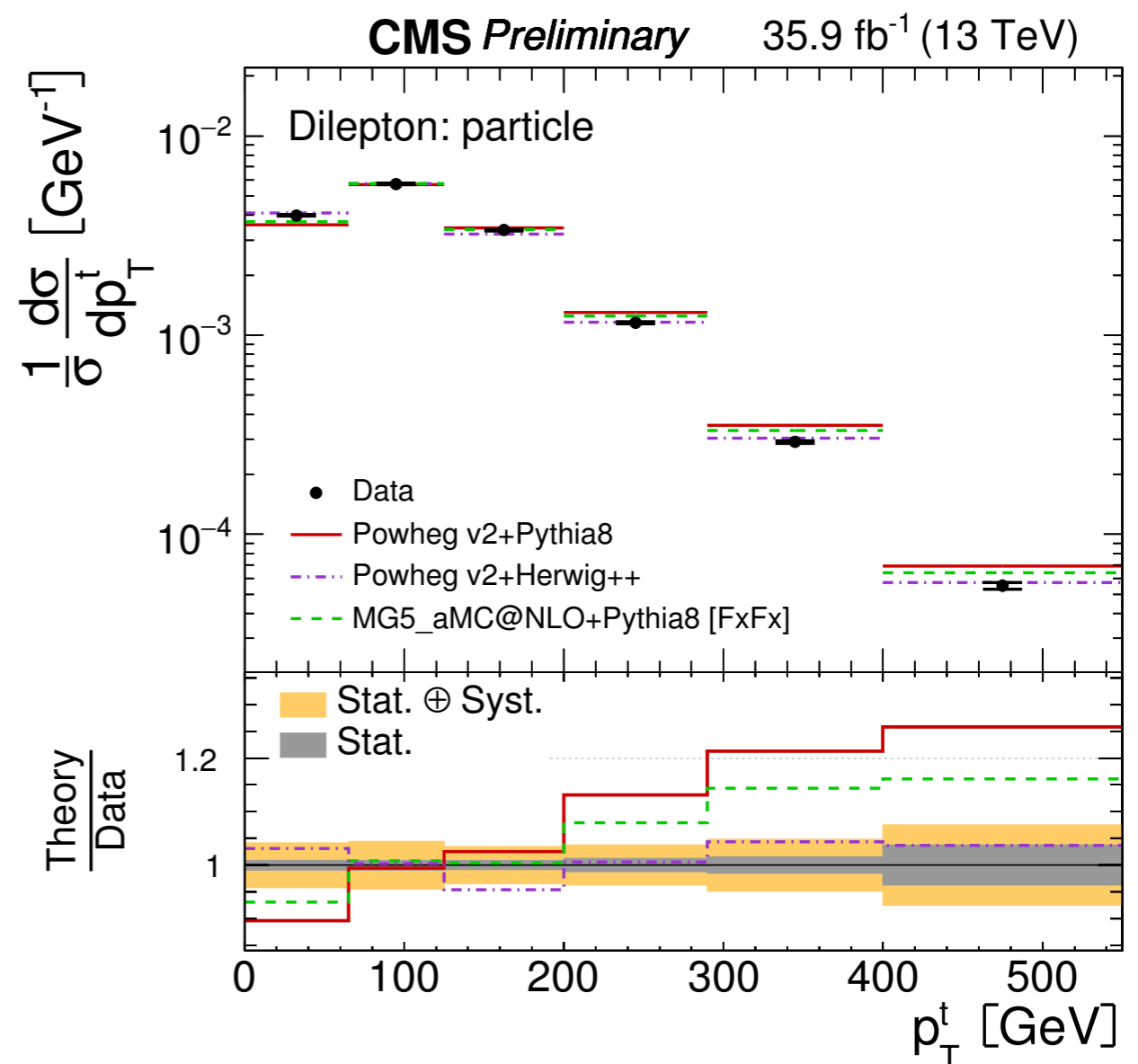
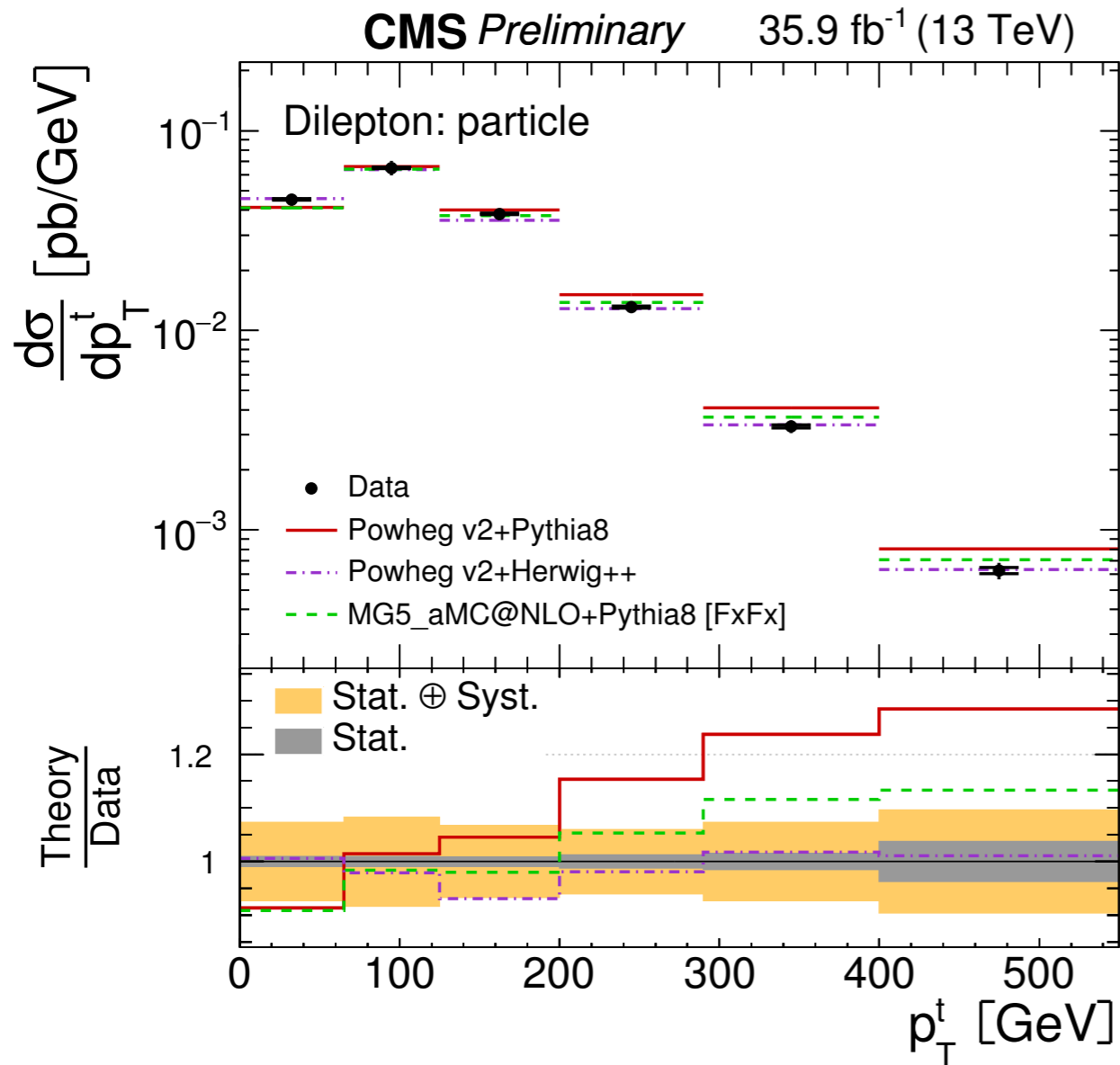
### Particle level vs NLO+PS



• Measurements agree with predictions

# The Top $p_T$ saga: dilepton NLO+PS @13 TeV

CMS-PAS-TOP-17-014



<b><math>N_{\text{dof}} = 5 \text{ or } 6</math></b>	normalised		absolute	
	$\chi^2$	p-val	$\chi^2$	p-val
NLO+PS				
PW+PY8	<b>128</b>	<b>&lt;10<sup>-3</sup></b>	<b>52</b>	<b>&lt;10<sup>-3</sup></b>
PW+H++	<b>6</b>	<b>0.306</b>	<b>3</b>	<b>0.830</b>
MG5+PY8	<b>45</b>	<b>&lt;10<sup>-3</sup></b>	<b>17</b>	<b>0.008</b>

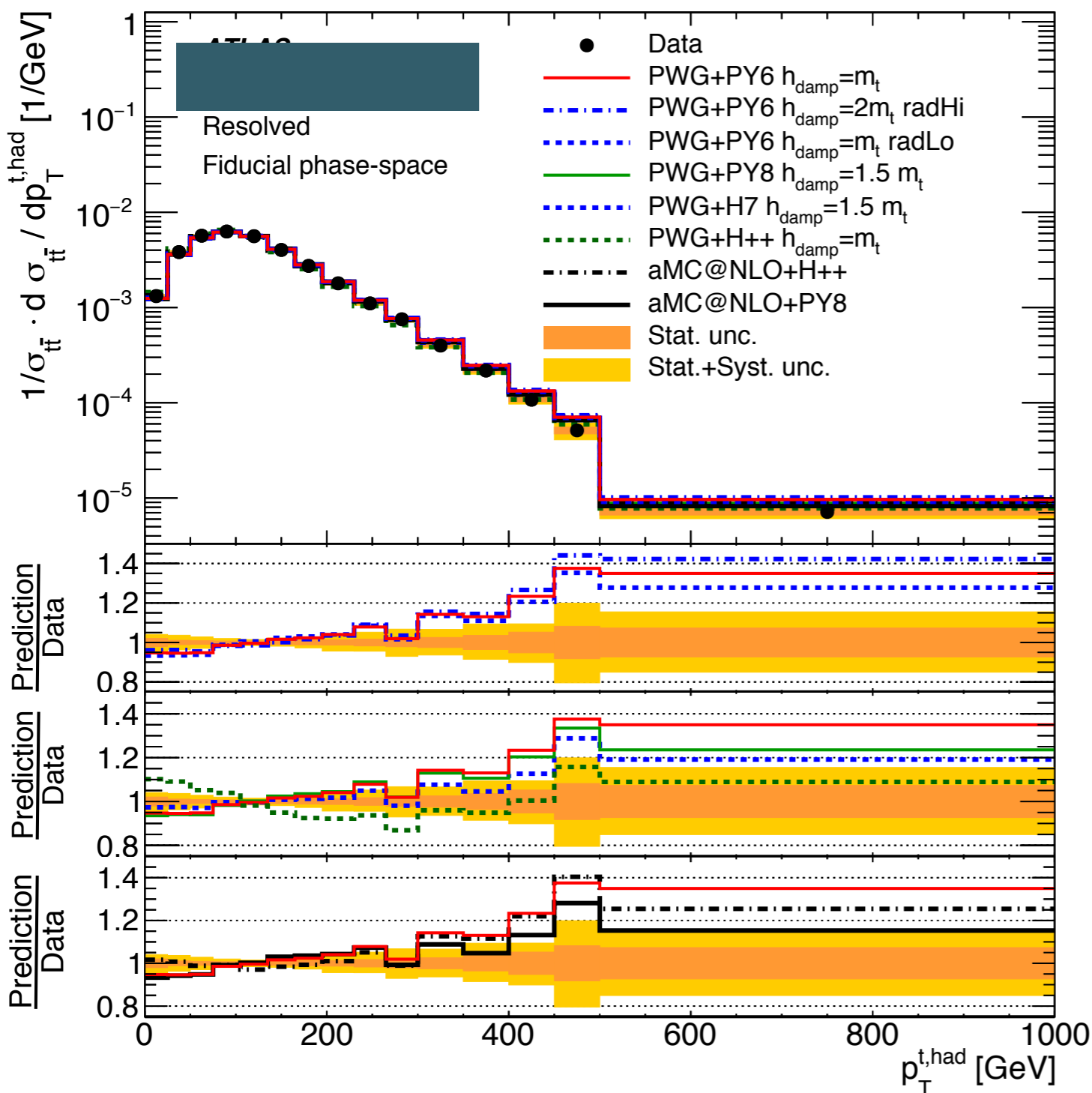
**no theory  
uncertainties  
included**

particle level

## Particle level vs NLO+PS

$1/\sigma_{tt} d\sigma_{tt}/dX$

*JHEP 11 (2017) 191*



$\ell$ +jets “resolved”

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Different radiation

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Different generator & hadronization

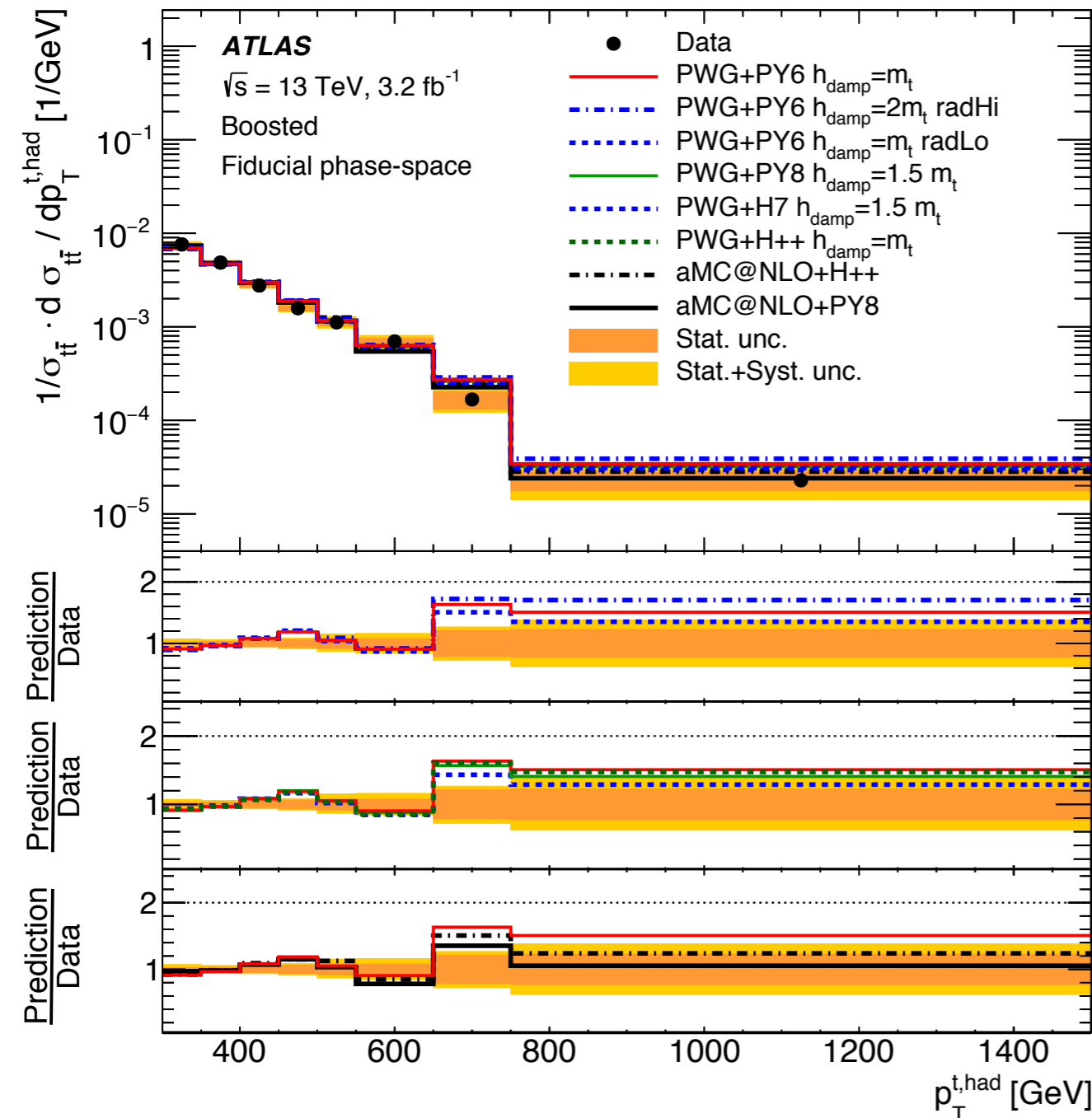
except for  
*POWHEG+HERWIG*

$\chi^2=15.4$   $p=0.35$  (-)

• Tension with most NLO+PS predictions: slope w.r.t measurement

## Particle level vs NLO+PS

$1/\sigma_{t\bar{t}} d\sigma_{t\bar{t}}/dX$  *JHEP 11 (2017) 191*



$\ell$ +jets “boosted”

Different  
Radiation

$$N_{\text{dof}} = 7$$

Different  
Hadronization

$$9.9 < \chi^2 < 11.5$$

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Different  
generator  
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- Tension with most NLO+PS predictions: slope w.r.t measurement.
- Less tension than in “resolved”, larger statistical uncertainties

# Extreme test of SM: the top $p_T$ “saga”

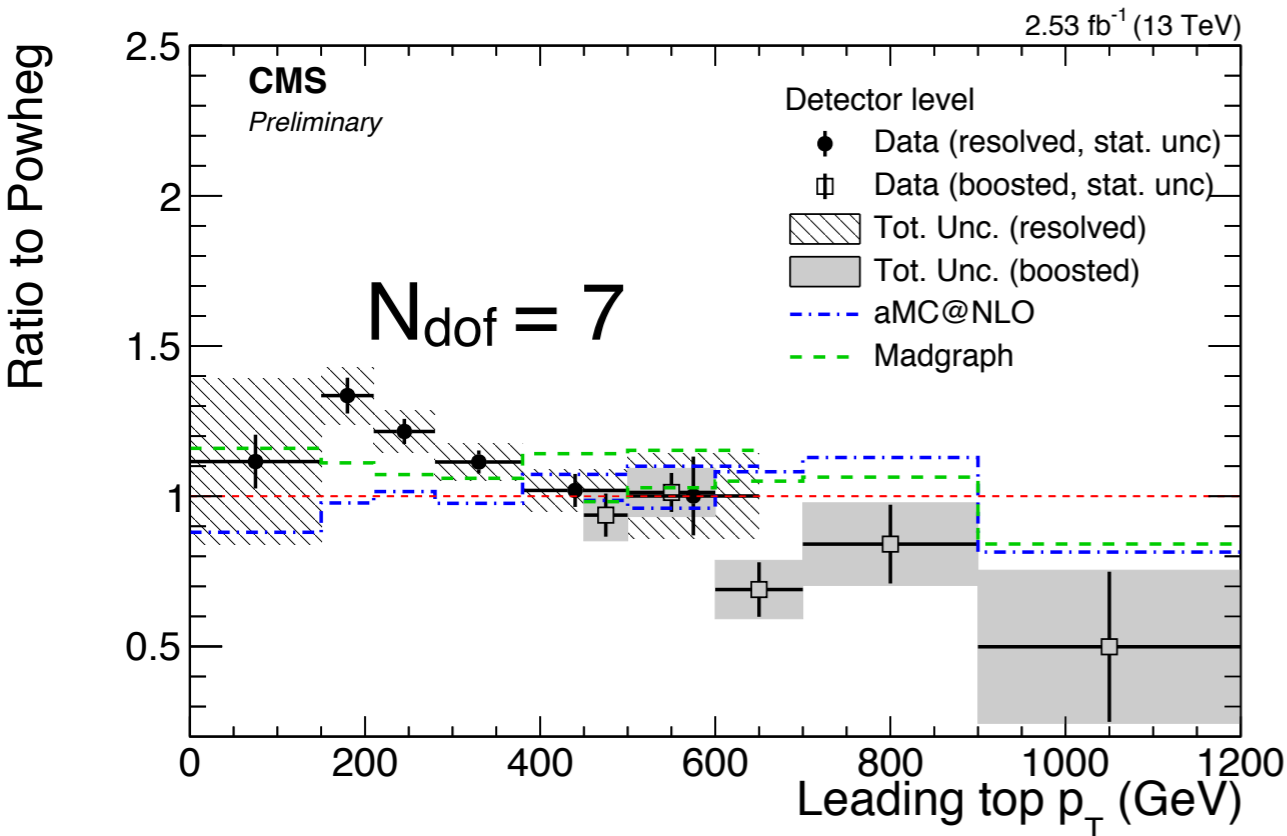
@13TeV

## Particle level vs NLO+PS

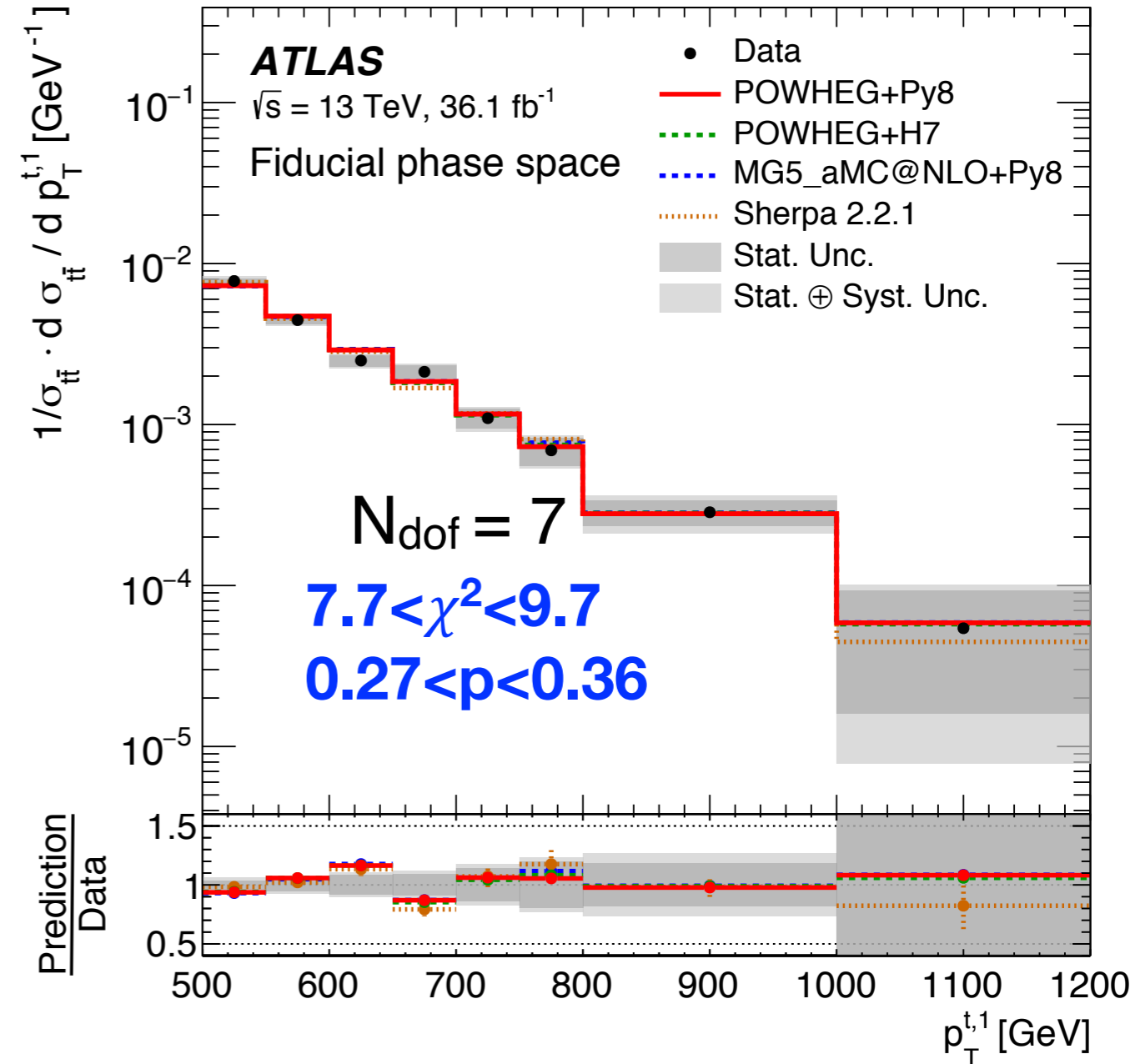
$p_{T,lead\ top}$

### Reco level

CMS-TOP-16-013,  $2.5\text{ fb}^{-1}$ , 13 TeV



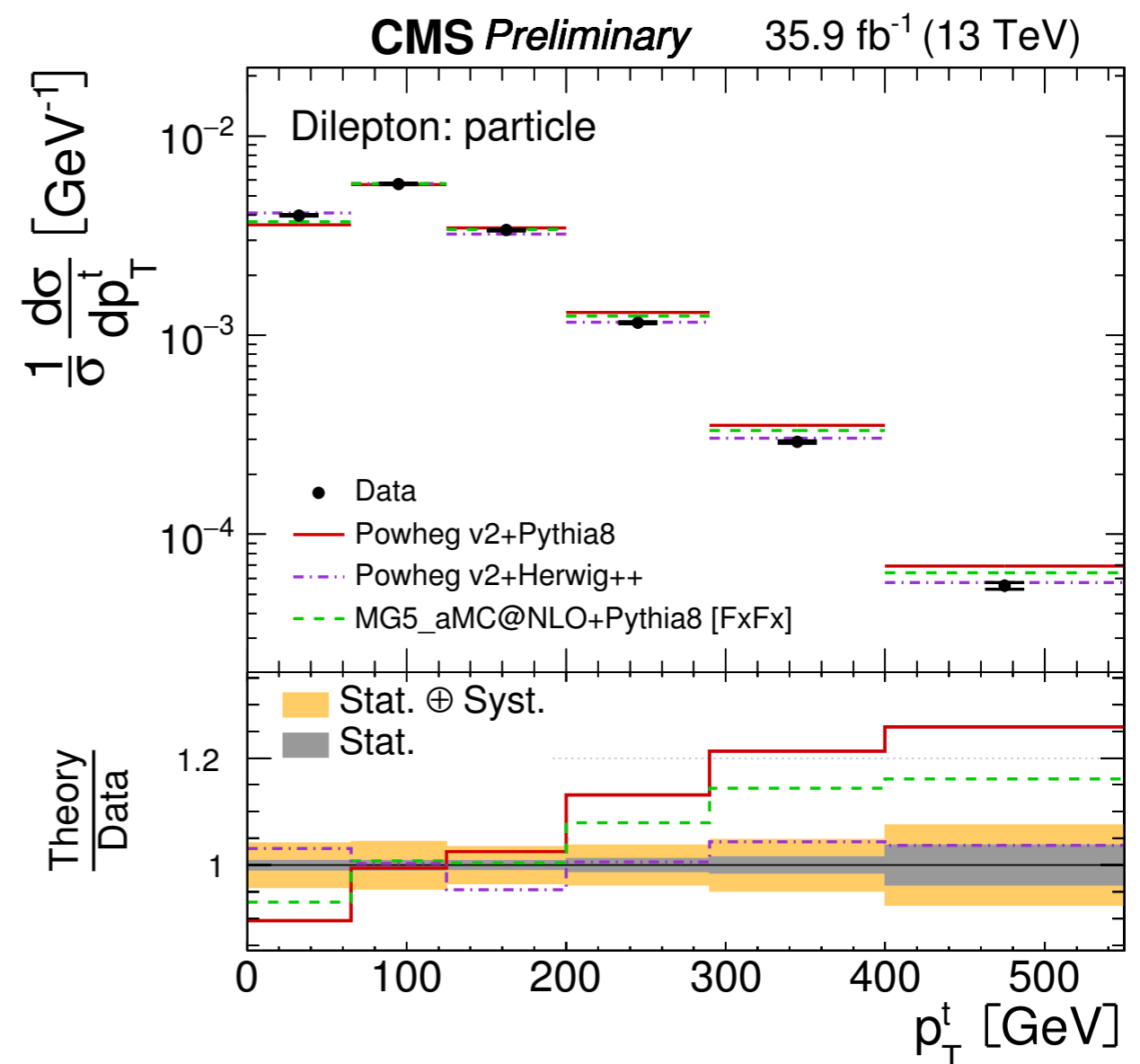
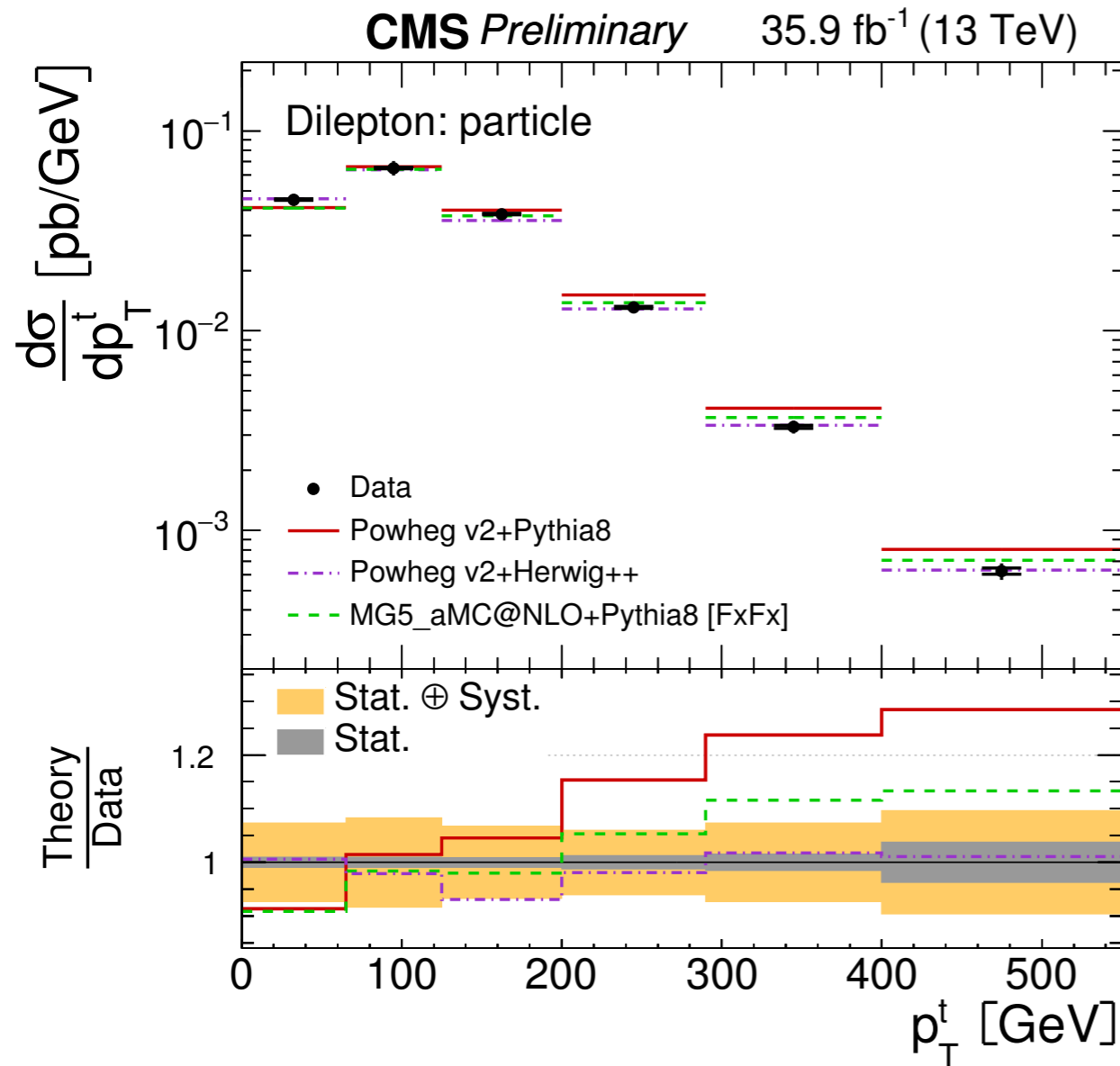
### Particle level vs NLO+PS



• Measurements agree with predictions

# The Top $p_T$ saga: dilepton NLO+PS @13 TeV

CMS-PAS-TOP-17-014



<b><math>N_{\text{dof}} = 5 \text{ or } 6</math></b>	normalised		absolute	
	$\chi^2$	p-val	$\chi^2$	p-val
NLO+PS				
PW+PY8	<b>128</b>	<b>&lt;10<sup>-3</sup></b>	<b>52</b>	<b>&lt;10<sup>-3</sup></b>
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**no theory  
uncertainties  
included**

CMS-TOP-16-013,  $2.5 \text{ fb}^{-1}$ , 13 TeV

In all-jets final state

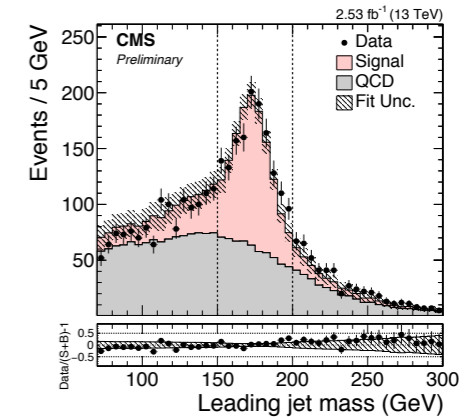
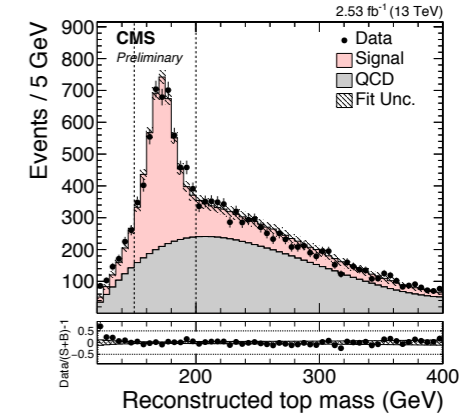
# ALL HADRONIC CMS @ 13 TeV (still preliminary)

## Resolved

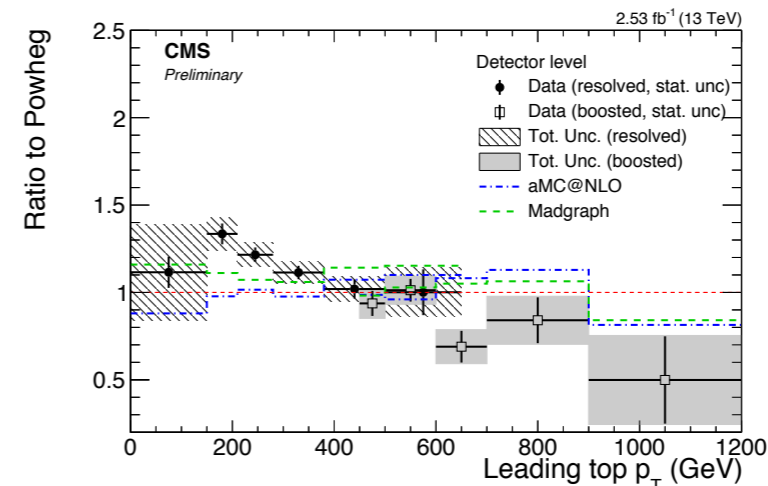
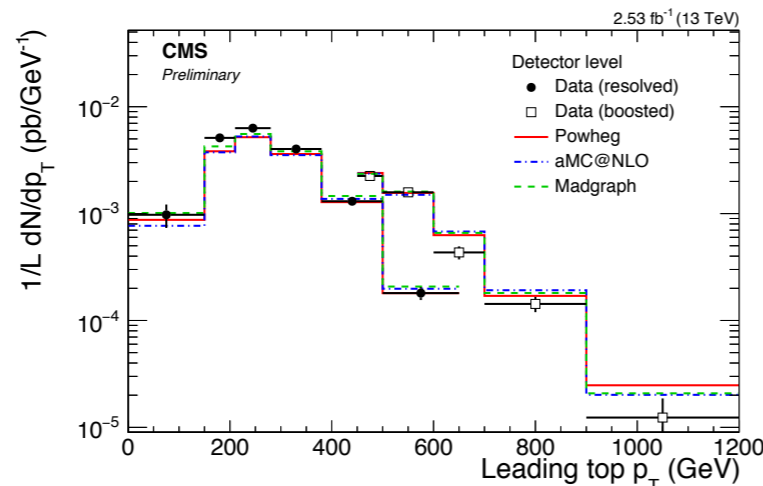
- Selection: at least 6 jets, 2 b tagged.
- Perform kinematic fit for  $t\bar{t}$  reconstruction (based on W and top mass constrains)
- Accept events with  $150 < m_t^{\text{fit}} < 200 \text{ GeV}$ , and fit probability greater than 0.02.

## Boosted

- 1 jet  $p_T > 200 \text{ GeV}$  and 1 jet  $p_T > 450 \text{ GeV}$
- each jet: softdrop mass  $> 50 \text{ GeV}$ , b tagged subjet, n-jettiness requirements.



Template Fit: Signal template from MC, background template from data by inverting b tagging.

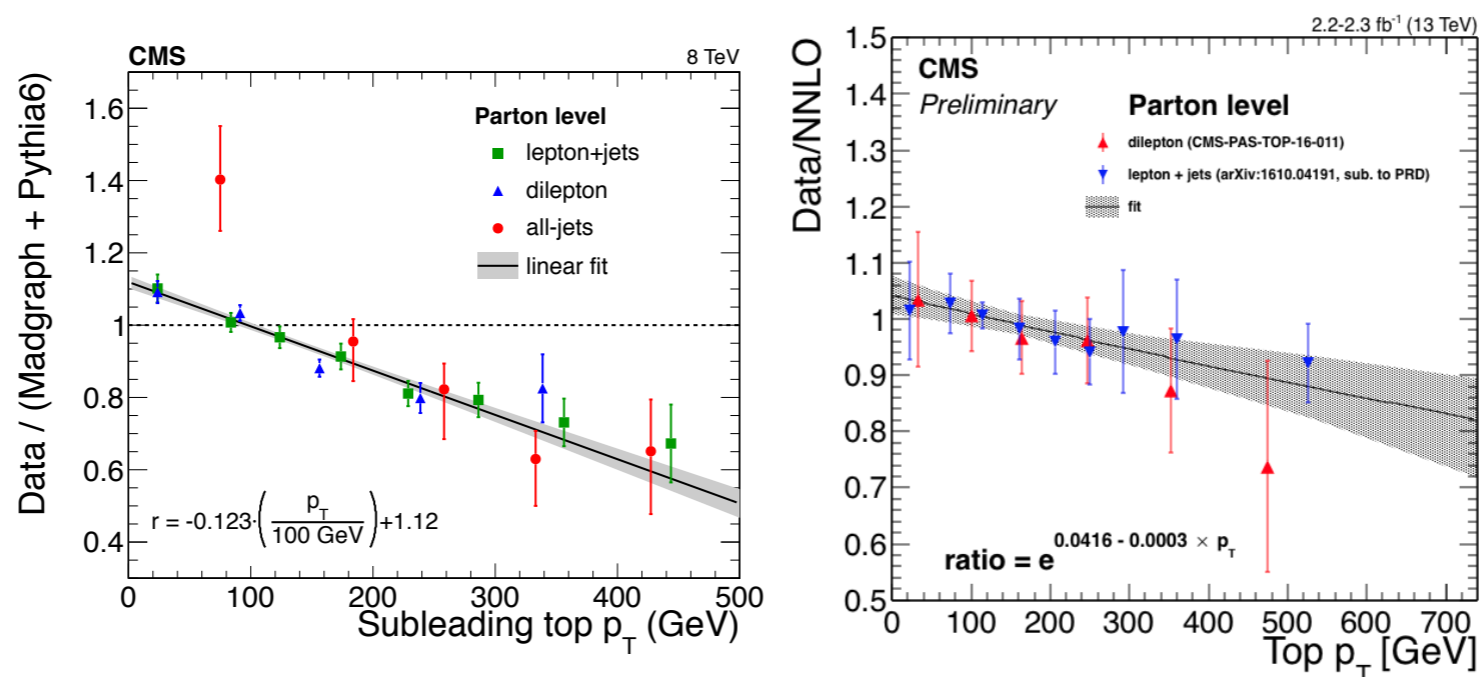


Soft  $p_T(t)$  confirmed in all-jets channel and persisting in boosted regime.



## Summary differential cross sections

- Measurements at 7, 8, and 13 TeV in various  $t\bar{t}$  decay channels.
- $p_T(t)$  observed softer, but compatible with standard model within uncertainties in measurement and theory.
- Persistent in boosted regime  $p_T(t) > 400$  GeV.



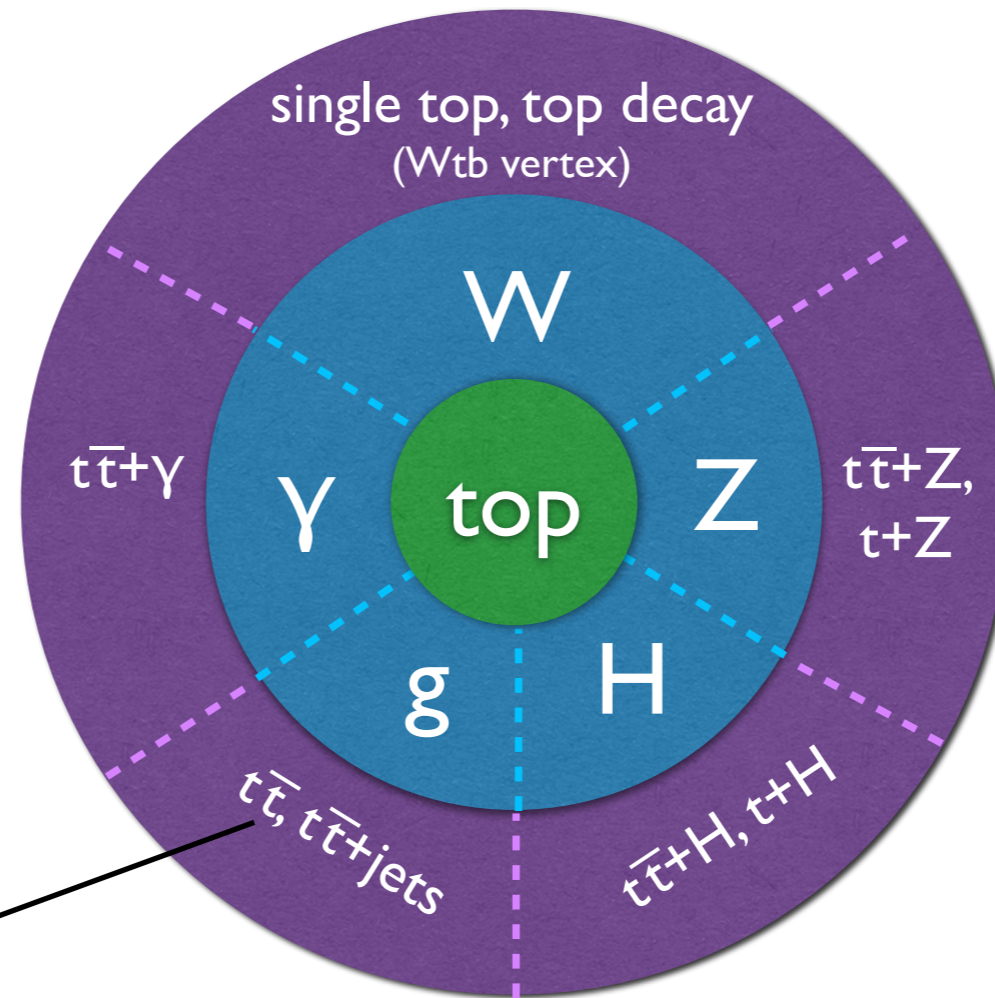
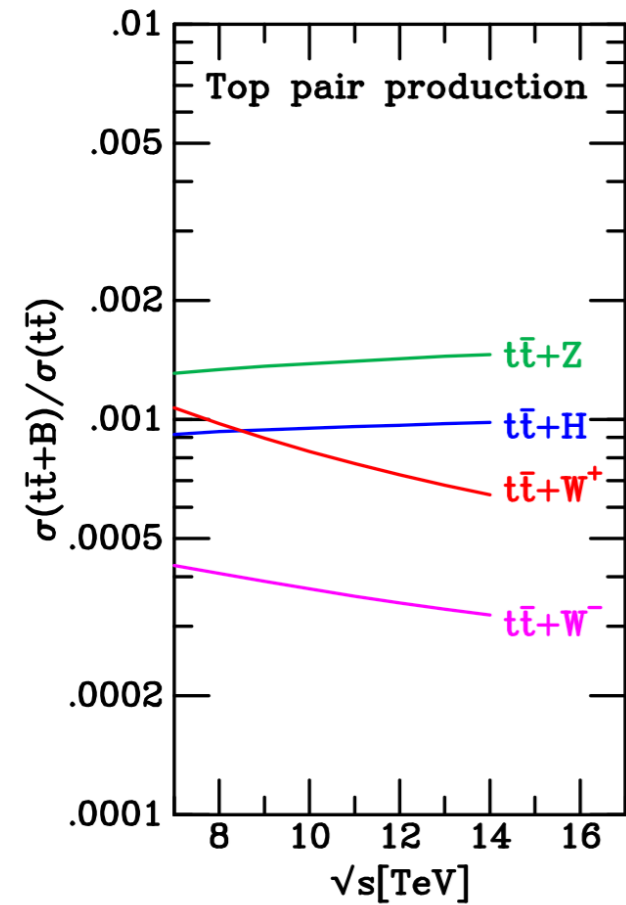
# top pair associated production

- The top quark couples to other SM fields through its **gauge** and **Yukawa interactions**
- $t \rightarrow Wb$  coupling measured already at the Tevatron

High statistics top physics at the LHC:  $t\bar{t}$  + bosons ( $\gamma$ , Z and H) becomes accessible!

First evidence on the coupling of the top quark to these particles from **production rate**

Important Standard Model test: **new physics** modifies the structure of the EW couplings

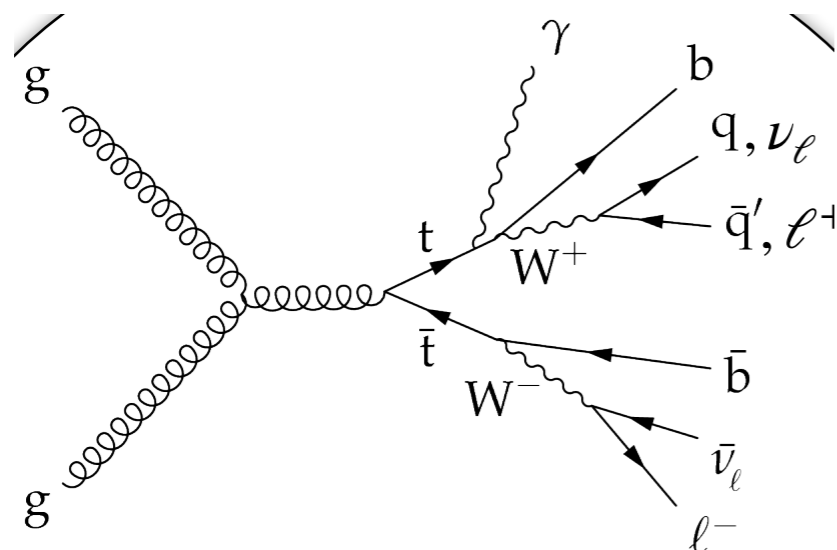


[T. Vazquez Schroder @ Top2014](#)

see *inclusive top cross section &  $d\sigma/dX$  with jets* (Lecture 5)

**Top** + **X** coupling, **how to measure it?**

# tt+photon: latest ATLAS



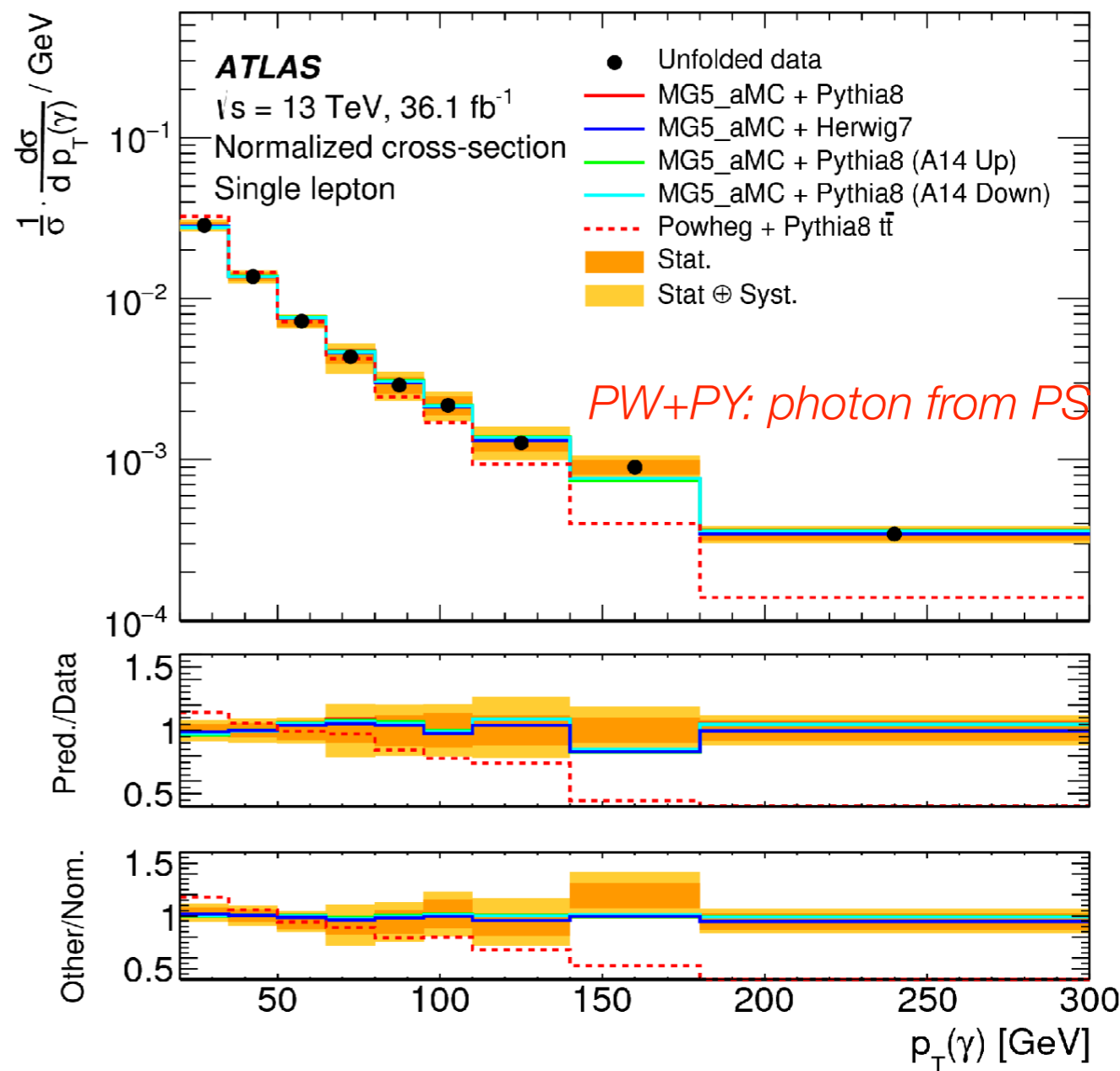
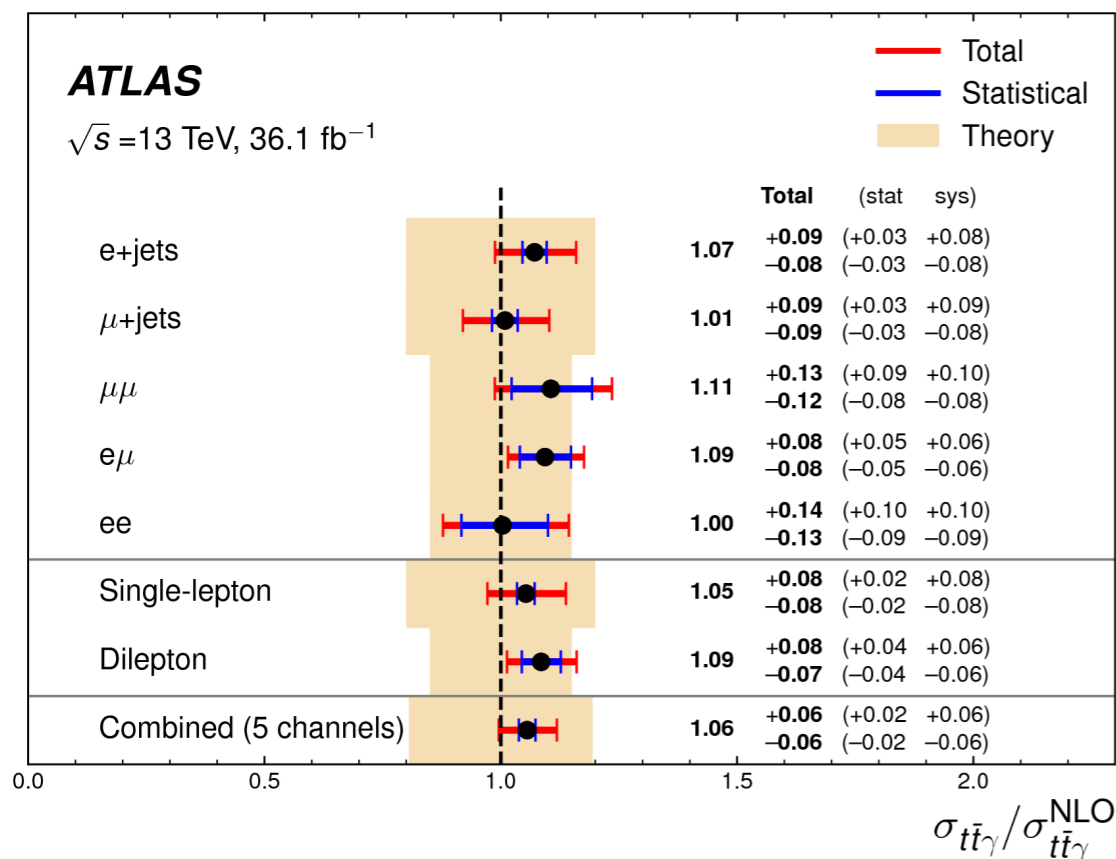
Eur. Phys. J. C 79 (2019) 382

$$\sigma_{\text{fid.}}^{\text{SL}} = 521 \pm 9 \text{ (stat.)} \pm 41 \text{ (syst.) fb}$$

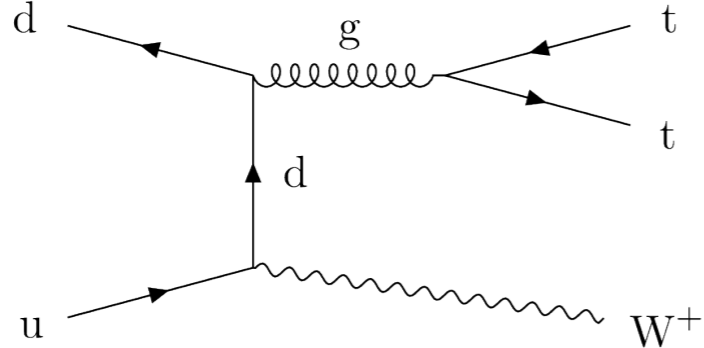
in agreement with NLO prediction:  $\sigma_{\text{fid.}}^{\text{pred}} = 495 \pm 99 \text{ fb}$

$$\sigma_{\text{fid.}}^{\text{DL}} = 69 \pm 3 \text{ (stat.)} \pm 4 \text{ (syst.) fb}$$

in agreement with NLO prediction:  $\sigma_{\text{fid.}}^{\text{pred}} = 63 \pm 9 \text{ fb}$



$tt+W/Z$

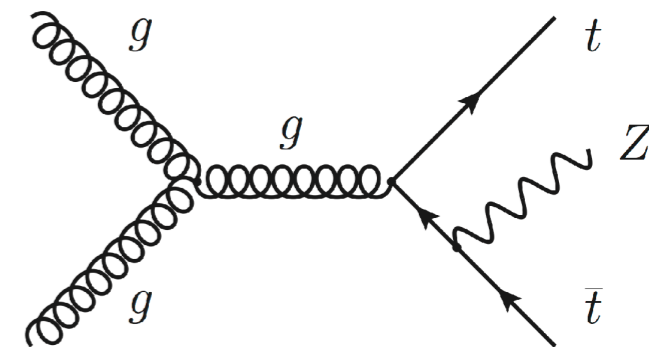


ATLAS (36/fb)

$$\sigma_{t\bar{t}Z} = 0.95 \pm 0.08 \text{ (stat.)} \pm 0.10 \text{ (syst.) pb}$$

$$\sigma_{t\bar{t}W} = 0.87 \pm 0.13 \text{ (stat.)} \pm 0.14 \text{ (syst.) pb}$$

multi-lepton ( $2\ell, 3\ell, 4\ell$ )

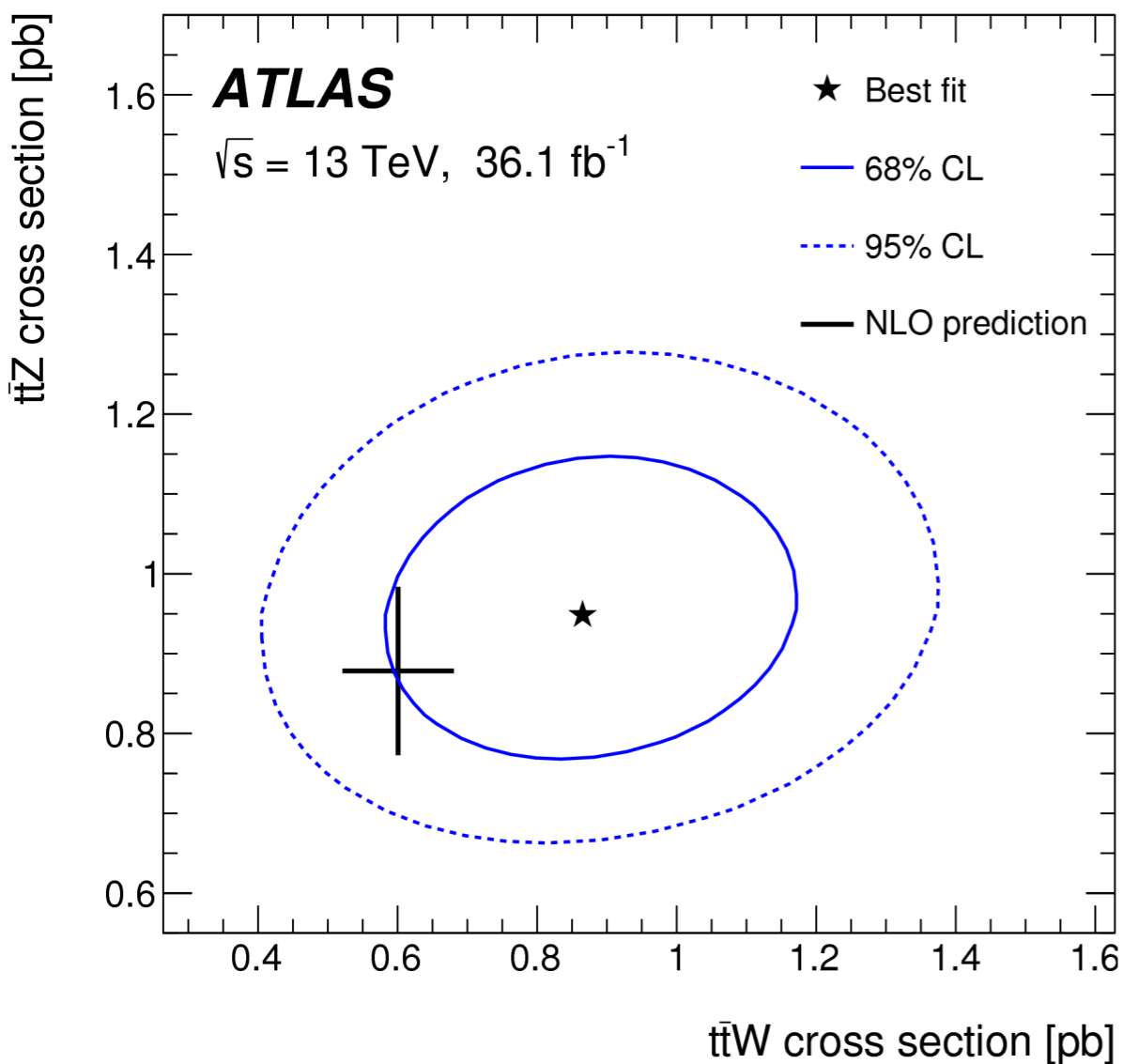


CMS (75/fb)

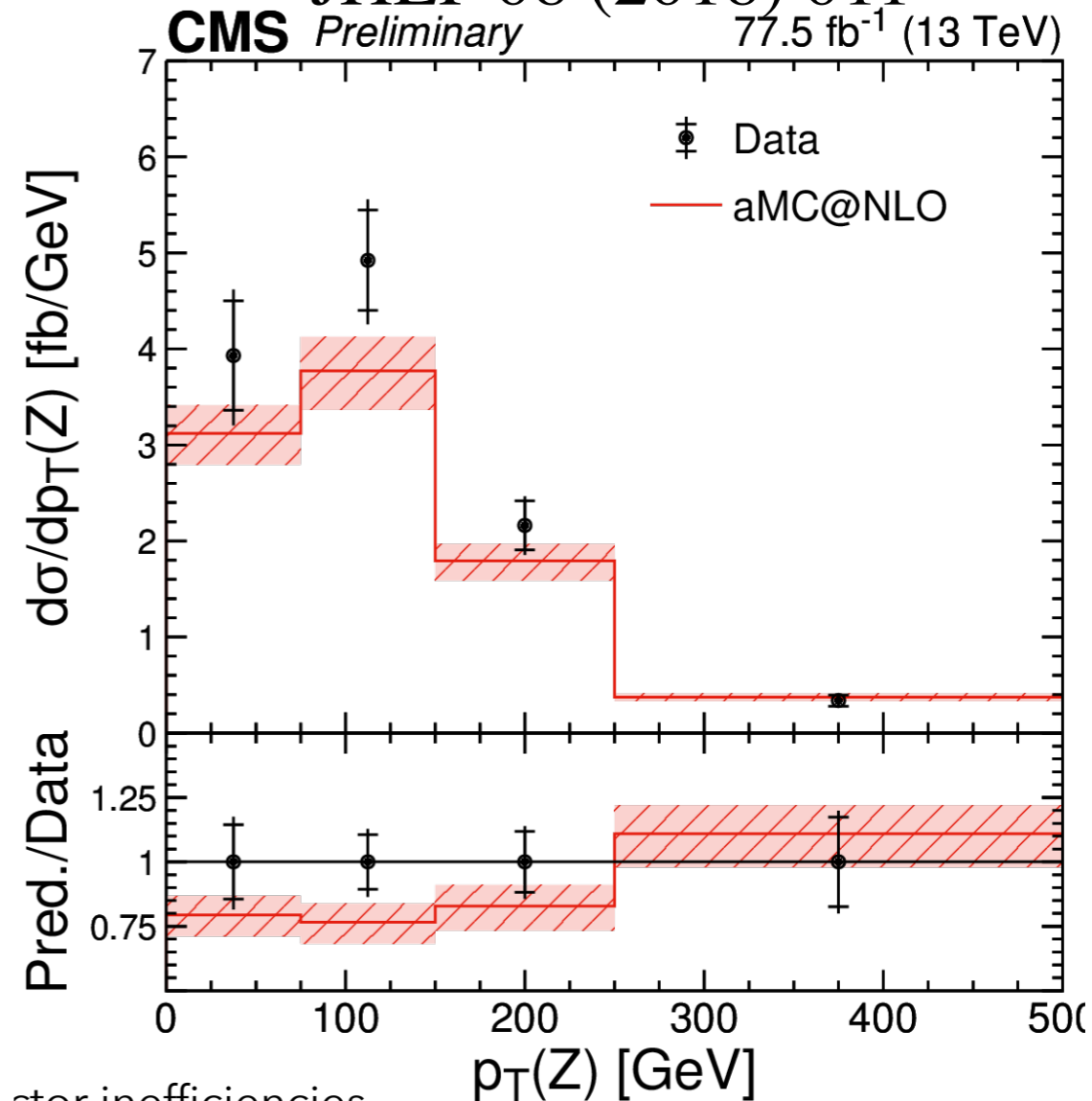
$$\sigma_{t\bar{t}W} = 0.77^{+0.12}_{-0.11} \text{ (stat.)}^{+0.13}_{-0.12} \text{ (syst.) pb}$$

$$\sigma_{t\bar{t}Z} = 0.99^{+0.09}_{-0.08} \text{ (stat.)}^{+0.12}_{-0.10} \text{ (syst.) pb}$$

Phys. Rev. D 99 (2019) 072009



JHEP 08 (2018) 011



ATLAS (36/fb)

Table 8: List of relative uncertainties in the measured cross sections of the  $t\bar{t}Z$  and  $t\bar{t}W$  processes from the fit, grouped in categories. All uncertainties are symmetrized. The sum in quadrature may not be equal to the total due to correlations between uncertainties introduced by the fit.

Uncertainty	$\sigma_{t\bar{t}Z}$	$\sigma_{t\bar{t}W}$
Luminosity	2.9%	4.5%
Simulated sample statistics	2.0%	5.3%
Data-driven background statistics	2.5%	6.3%
JES/JER	1.9%	4.1%
Flavor tagging	4.2%	3.7%
Other object-related	3.7%	2.5%
Data-driven background normalization	3.2%	3.9%
Modeling of backgrounds from simulation	5.3%	2.6%
Background cross sections	2.3%	4.9%
Fake leptons and charge misID	1.8%	5.7%
$t\bar{t}Z$ modeling	4.9%	0.7%
$t\bar{t}W$ modeling	0.3%	8.5%
Total systematic	10%	16%
Statistical	8.4%	15%
Total	13%	22%

# Top Properties : Spin correlations

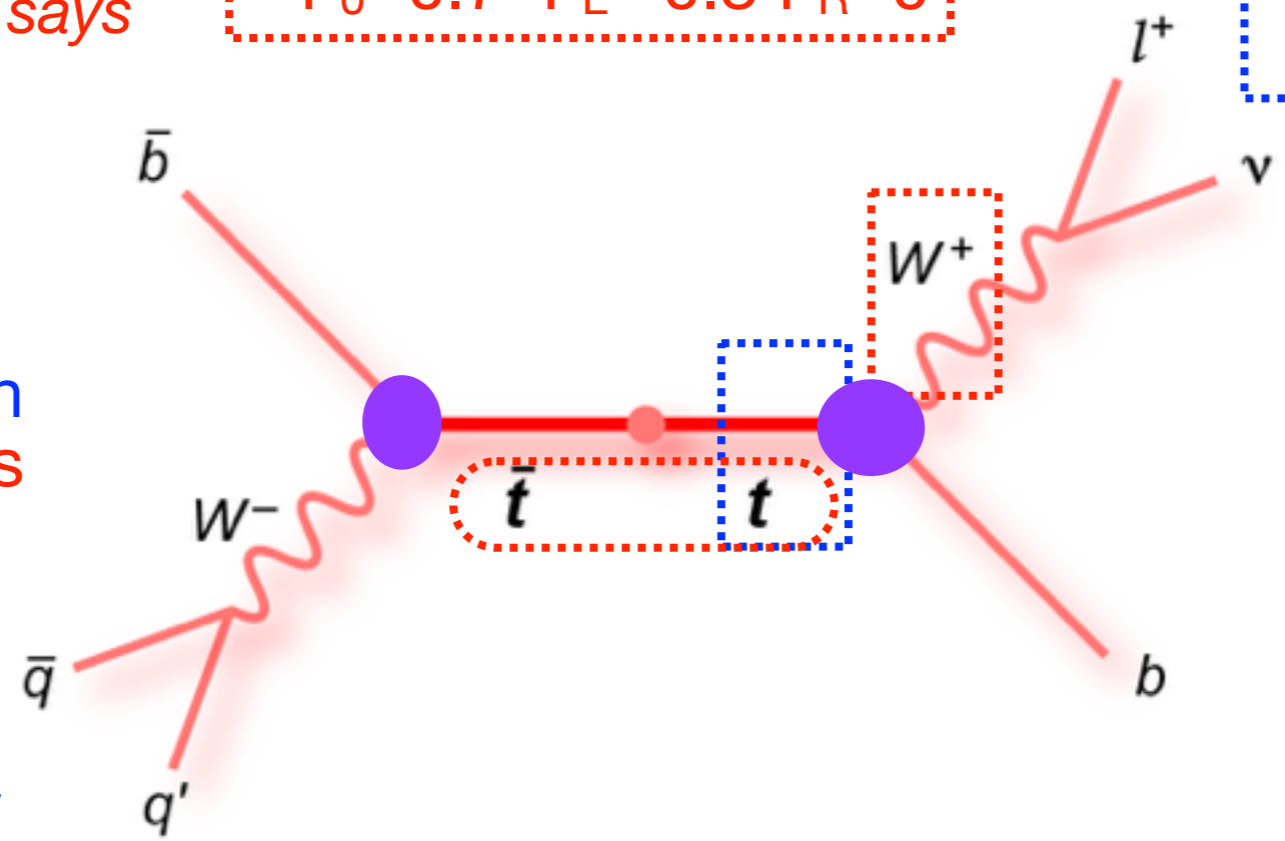
# Standard reasons: Top quark ( $t$ ) properties

**top quark decays before its spin flips** → Spin information is passed to decay products  
**undiluted by hadronization** → angles of  $t\bar{t}$  decay products are correlated



• W helicity in  $Wtb$  vertex? SM says  $F_0 \sim 0.7$   $F_L \sim 0.3$   $F_R \sim 0$

• is top quark unpolarised in  $pp \rightarrow t\bar{t}$  production? SM says yes



• Are  $t\bar{t}$  spins correlated in  $pp \rightarrow t\bar{t}$  production? SM says YES

• is top polarised in single top production? SM says YES

• are t and anti-t angular distribution different? SM says yes @NLO

• is CP violation visible in b-decay from  $t\bar{t}$ ? SM says yes at  $<10^{-2}$

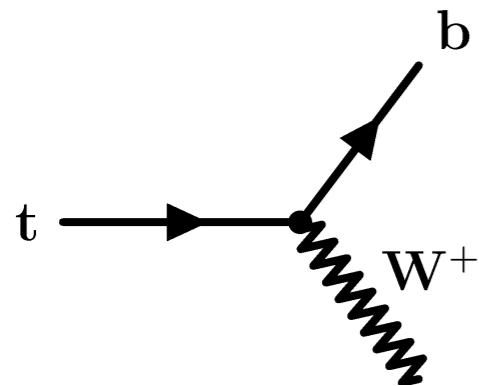
•  $Wtb$  vertex? SM says V-A: i.e. spin density matrix as foreseen in combination of  $t\bar{t}$  production and decay?

**All angular properties are found to be consistent with SM**



# The spin of the top quark

*see Liss, Maltoni, Quadt,  
PDG Review, 2016*



$$\frac{1}{m_t} < \frac{1}{\Gamma_t} < \frac{1}{\Lambda} < \frac{m_t}{\Lambda^2}$$

Production time < Lifetime < Hadronization time < Spin decorrelation time

individual  
top quark

- chiral coupling in SM  
Wtb vertex enhances  
specifically polarized  
W-boson and b-quark

- **Angular distribution of top quark decay products**  
follows the predictions of the  
top quark spin (*differently from b  
quark in which B meson decays  
isotropically*)
- Top quark polarization and  
(and consequently its spin) is  
directly observable by such  
angular distributions

Spin for

top quark  
pairs

single  
top

- in qq/gg tt interaction at pp  
collider
  - top and anti-top quarks are  
~unpolarised (as the initial  
g and q)
  - the spins of t and t are  
correlated

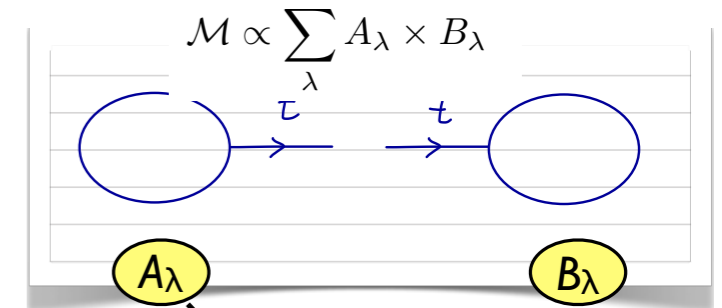
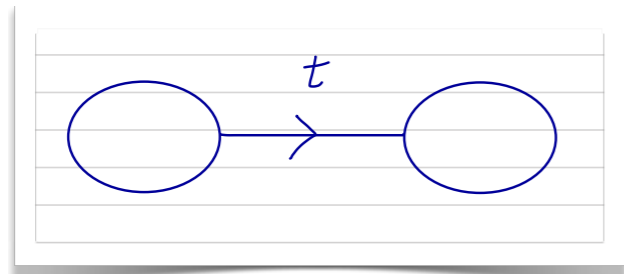
- **Observation of the top quark spin is strongly linked to its production and decay process**

# The spin of the top quark in $pp \rightarrow t\bar{t}$

B Lemmer, [arXiv 1410.1791](https://arxiv.org/abs/1410.1791)

J A Saavedra @ Top2014

- Top is narrow resonance:  $\Gamma_t/m_t \ll 1 \rightarrow$  the  $q\bar{q}/gg \rightarrow t\bar{t}$  production amplitude is factorizable in decay and production in the leading pole approx



$$\left( \frac{1}{3^2 \text{ or } 8^2} \sum_{\text{colours}} \right) \left( \frac{1}{2^2} \sum_{\text{spins}} \right) |\mathcal{M}(q\bar{q}/gg \rightarrow t\bar{t} \rightarrow (f_1 \bar{f}'_1 b) (f_2 \bar{f}'_2 \bar{b}))|^2 = \lambda_{ab} \rho_{ab, \bar{a}\bar{b}} \bar{\lambda}_{\bar{a}\bar{b}}$$

$a, b$ : top /anti-top quark spins,  $f$ : decay products  $\langle 0 \rangle = \text{tr}[\rho 0]$

$\rho_{ab, \bar{a}\bar{b}}$  Hermitian **production spin density matrix**  $\rightarrow$  **if  $\text{spin}_{\text{top}} = 1/2$ , Pauli matrices are a basis**

$$\rho_{ab, \bar{a}\bar{b}} \equiv \frac{1}{4} M^{00} \left( \delta_{ab} \delta_{\bar{a}\bar{b}} + P^i \sigma_{ab}^i \delta_{\bar{a}\bar{b}} + \bar{P}^{\bar{i}} \delta_{ab} \sigma_{\bar{a}\bar{b}}^{\bar{i}} + \hat{C}^{i\bar{i}} \sigma_{ab}^i \sigma_{\bar{a}\bar{b}}^{\bar{i}} \right)$$

$$P_i = \langle 2S_i \rangle \text{ top quark polarization} \quad S, \text{ top quark spin}$$

$$\hat{C}_{i\bar{i}} = \langle 4\hat{S}_i \hat{S}_{\bar{i}} \rangle \text{ top-anti-top spin correlation}$$

$\lambda_{ab}$  is the decay spin density matrix; if integrated over all decay product phase space, but one

$$\tilde{\lambda}(\vec{e})_{ab} \sim \delta_{ab} + \alpha_i \vec{e}_i \cdot \vec{\sigma}_{ab}$$

Including production and integrated decay density matrix in  $|\text{ME}|^2$  gives

$$\frac{d\sigma}{d^2\vec{e}_i d\vec{e}_j} \sim 1 + \alpha_i \vec{P} \vec{e}_i + \alpha_j \vec{P} \vec{e}_j + \alpha_i \alpha_j \vec{e}_i \hat{C} \vec{e}_j$$

# The spin of the top quark in pp collisions

B Lemmer, [arXiv 1410.1791](https://arxiv.org/abs/1410.1791)

[J A Saavedra @ Top2014](mailto:J.A.Saavedra@top2014)

- Going from generalized quantities to observable quantities requires the **choice of the spin quantization axis** : define spin axis as z direction and use associated polar coordinates

## in tt production

$$\frac{1}{\sigma} \frac{d^2\sigma}{d\cos\theta_i d\cos\theta_j} = \frac{1}{4} (1 + \underbrace{\alpha_i B_1 \cos\theta_i}_{\substack{\text{top quark} \\ \text{polarisation}}} + \underbrace{\alpha_j B_2 \cos\theta_j}_{\substack{\text{anti-top quark} \\ \text{polarisation}}} + \underbrace{\alpha_i \alpha_j C \cos\theta_i \cos\theta_j}_{\substack{\text{top-anti-top} \\ \text{spin correlation}}})$$

- **Degree of top quark polarization**

- single top : strong polarization
- $t\bar{t}$  production: B coefficients vanish at LO (pair invariance) :  $t\bar{t}$  almost unpolarized

- **Degree of top quark pair spin correlations:** depends on choice of spin quantization axis

*JHEP12(2015)026*

$$\langle |\text{MatrixEl}(t\bar{t})|^2 \rangle \rightarrow \frac{1}{\sigma} \frac{d^2\sigma}{d\cos\theta_+^a d\cos\theta_-^b} = \frac{1}{4} (1 + B_+^a \cos\theta_+^a + B_-^b \cos\theta_-^b - C(a,b) \cos\theta_+^a \cos\theta_-^b)$$

$a,b=\{k,n,r\}$

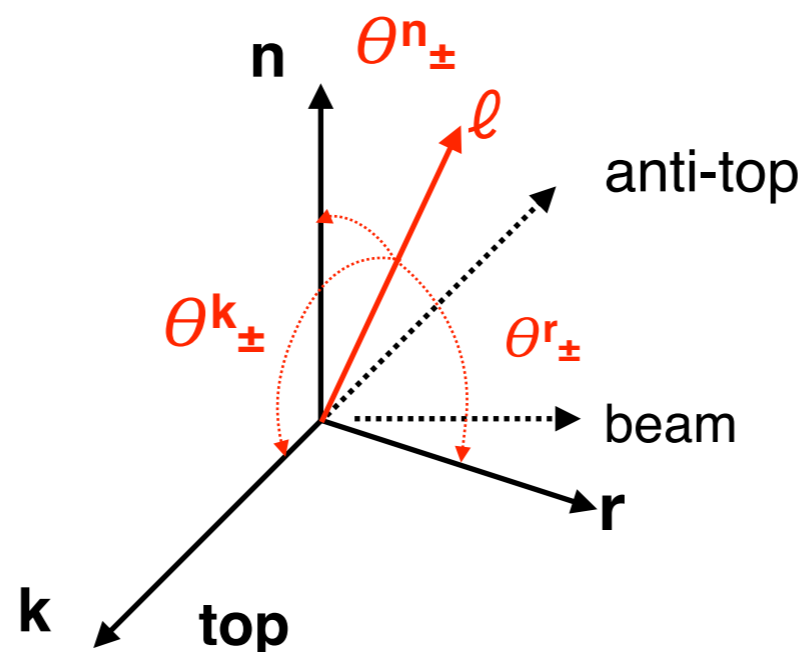
$+=\text{top decay}$   $-=\text{anti-top decay}$

$\cos\theta^{a,b}_{\pm}$ : **6 angles:**  
**2 lepton directions**  
in top/anti-top parent  
rest frame w.r.t  
**3 spin quantization**  
**axes** (a or b)

*dilepton events*

- ▶ **k**: top quark direction in  $t\bar{t}$  rest frame
- ▶ **n**:  $\perp$  to **k** & **laboratory beam direction**
- ▶ **r**:  $\perp$  to **k** and **n**

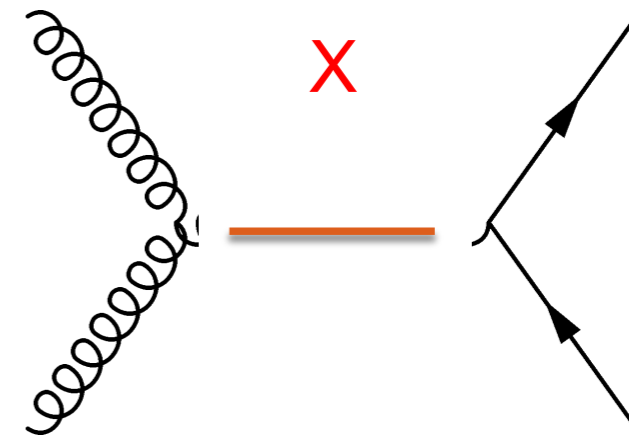
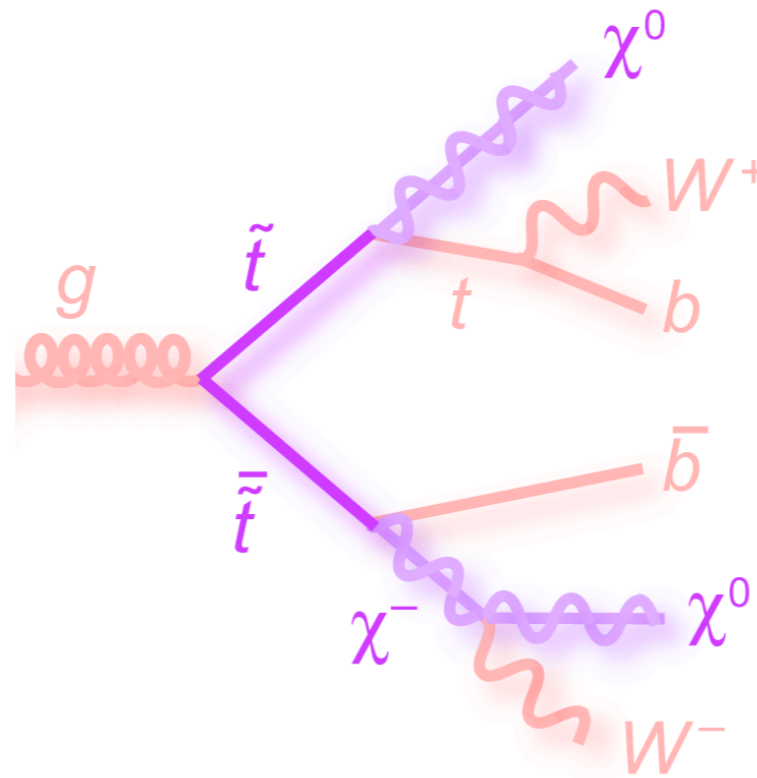
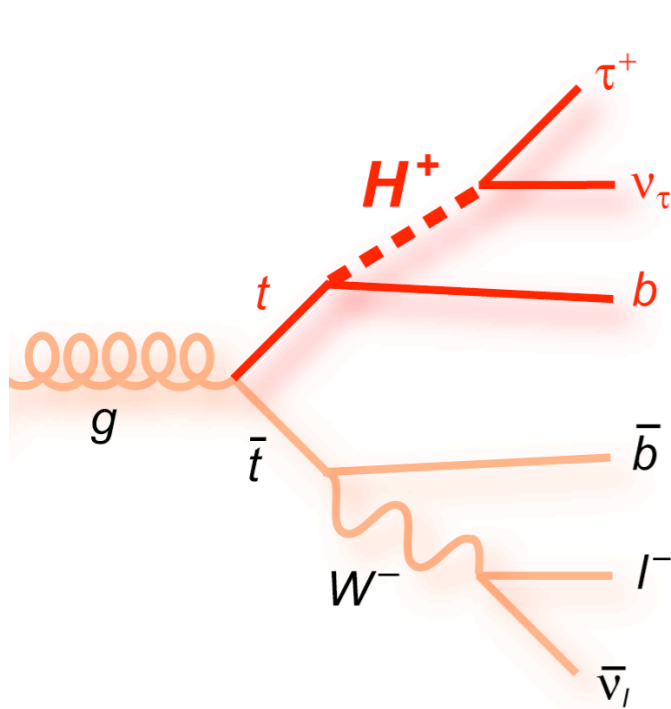
**15 spin density  
matrix elements**



- approximate CP symmetry of SM  $\rightarrow$ 
  - ▶ C (i,j) is symmetric
  - ▶ top and anti top quarks have same polarisation coefficients
- ▶ QCD invariant under  $P \rightarrow$  only P even and CP even coefficients are allowed

# Spin correlation: beyond the Standard Model

- Measured spin correlation can alter due to
  - Different decays
  - Different production
- Spin correlation: **test full chain from production to decay**



- Decays: charged Higgs,  $b'$ , ...

- Production: stop pairs, KK gravitons,  $Z'$ , Higgs...

Miriam Watson

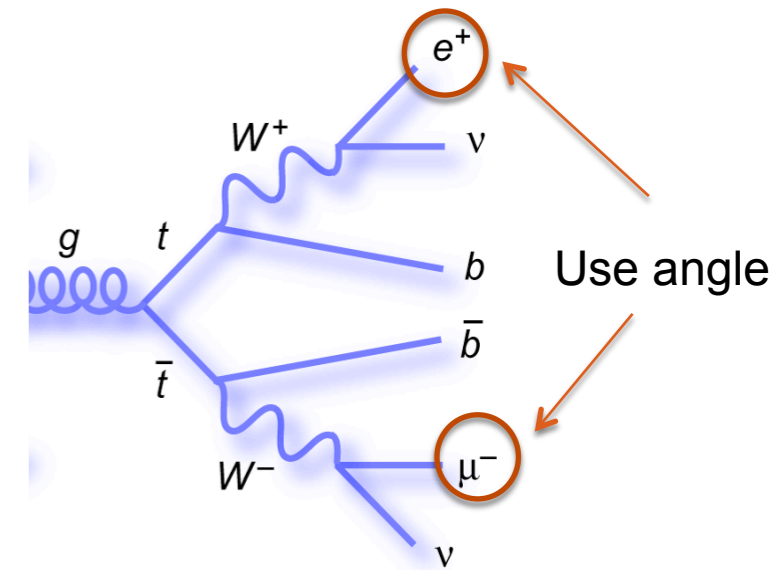
3

**Spin density matrix elements are consistent with SM predictions**

# Spin correlation with $\Delta\varphi$

$\sqrt{s}=13\text{TeV}$

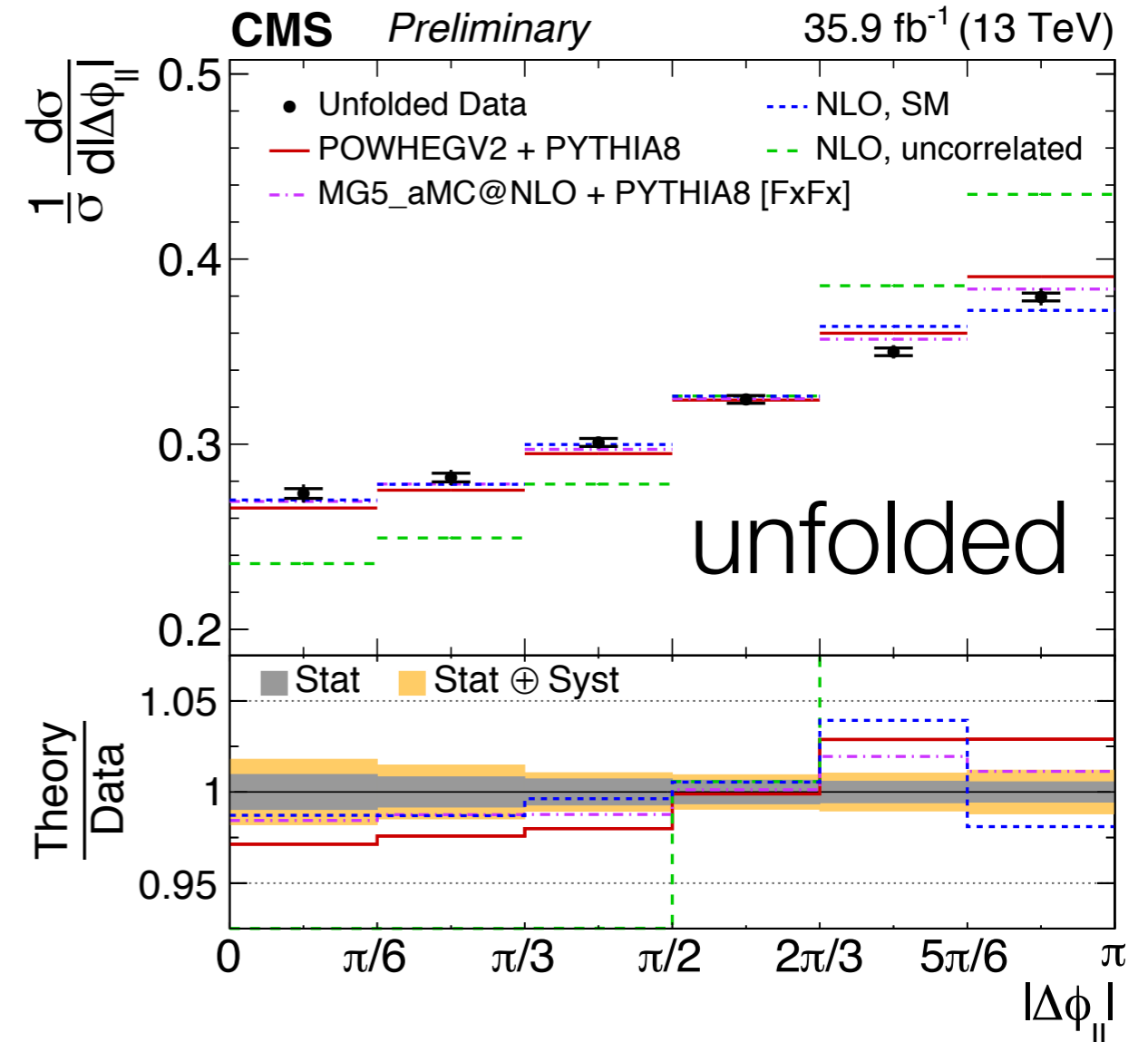
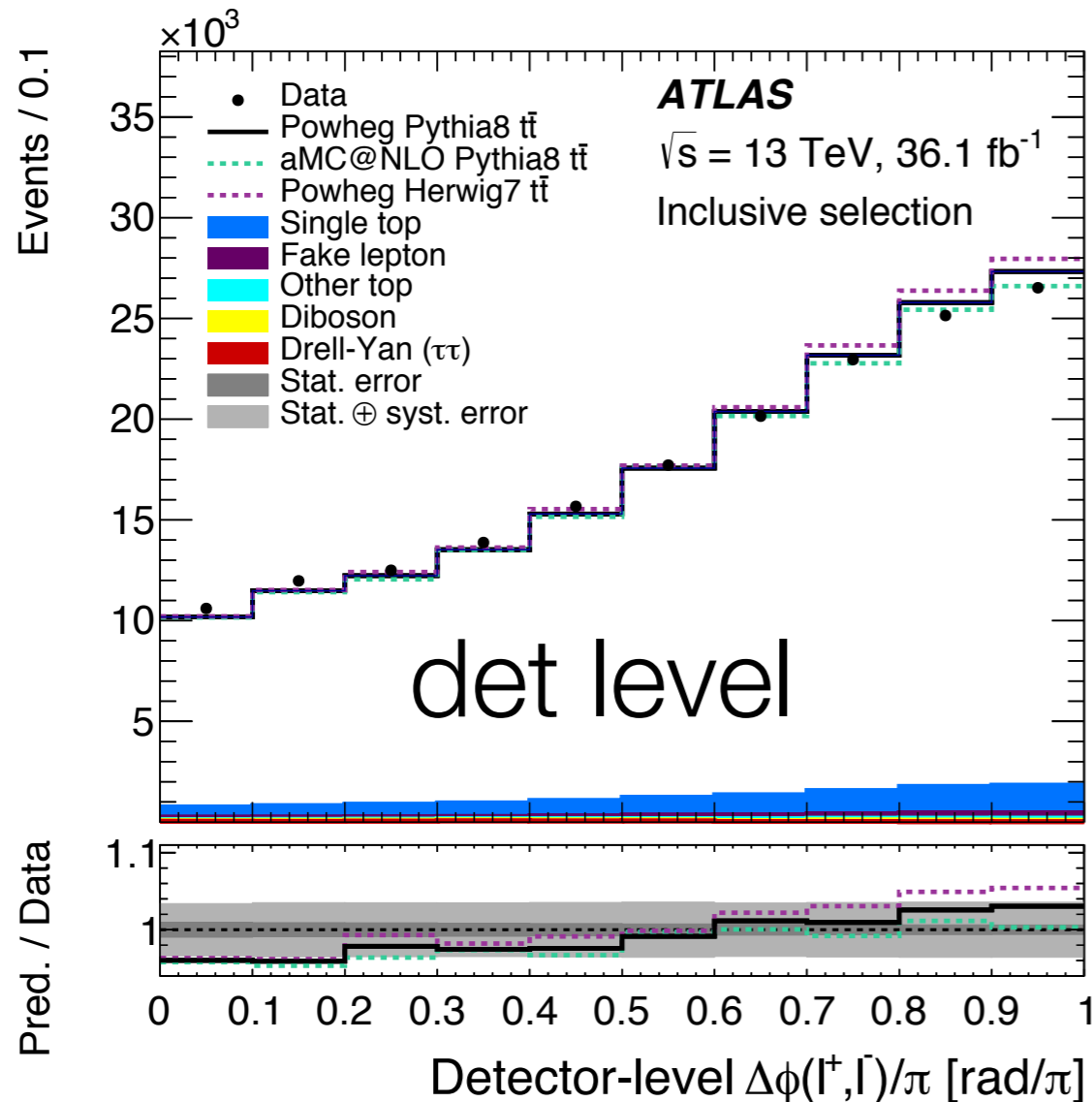
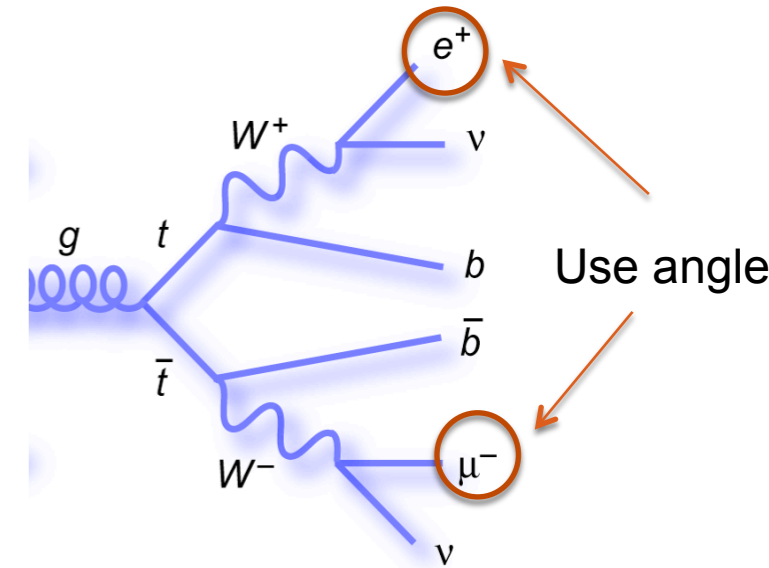
- Dilepton selection: ATLAS: **2 OS  $\ell$  ( $e\mu$  only),  $\geq 1$  b-jet,**
  - lepton has highest spin analysing power
- Derive  $\Delta\varphi$  difference in azimuthal angle between leptons in lab frame
  - no event reco, use lepton reco and resol
- Reconstruct tt final state :
  - constrains by  $m_W$  and  $m_{\text{top}}$ ,
  - test different  $\eta$  assumptions for  $2\nu$ : select assumption highest weight based on  $E_T^{\text{miss}}$  expected resolution
- Subtract bkg and unfold



# Spin correlation with $\Delta\phi$

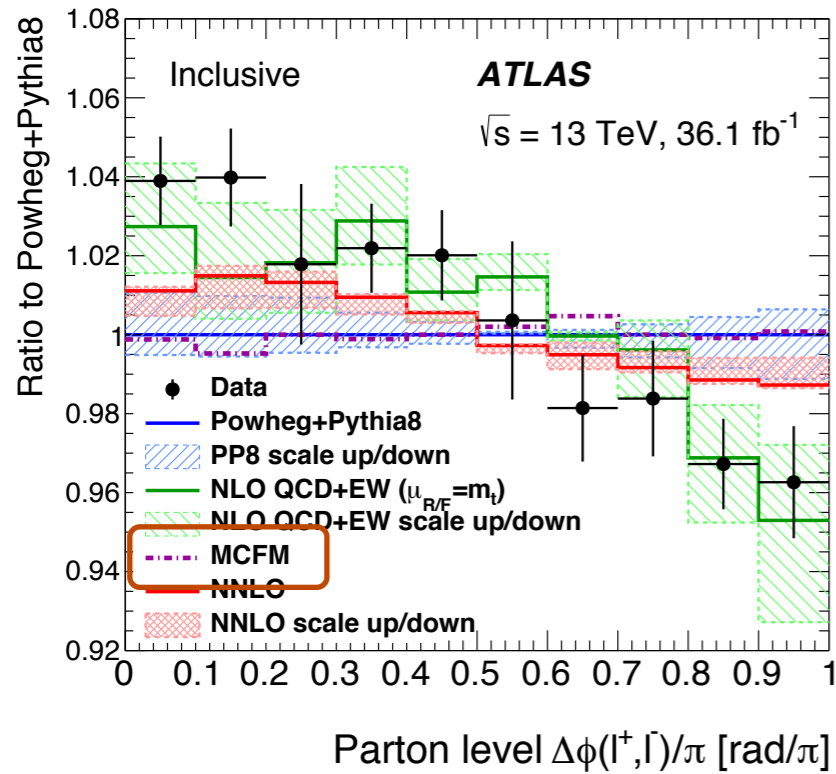
$\sqrt{s}=13\text{TeV}$

- Dilepton selection: ATLAS: **2 OS  $\ell$  ( $e\mu$  only),  $\geq 1$  b-jet,**
  - lepton has highest spin analysing power
- Derive  $\Delta\phi$  difference in azimuthal angle between leptons in lab frame
  - no event reco, use lepton reek and resol



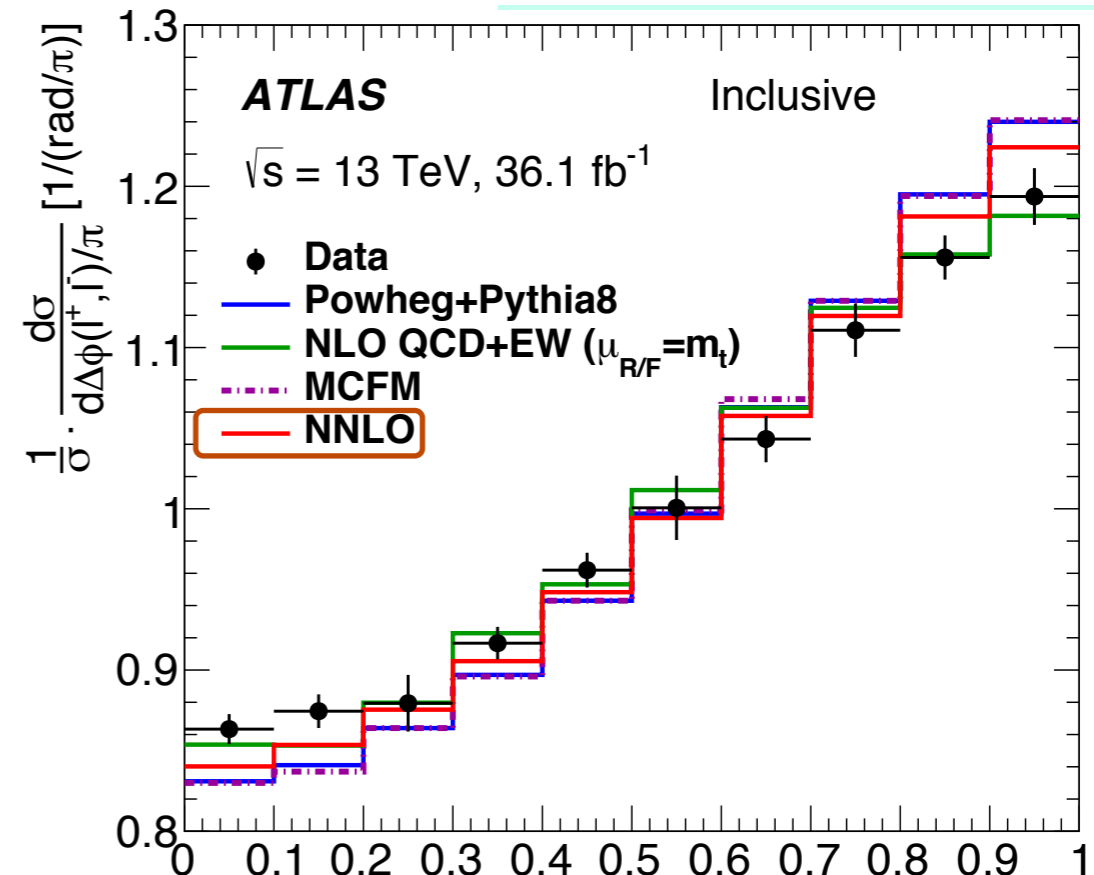
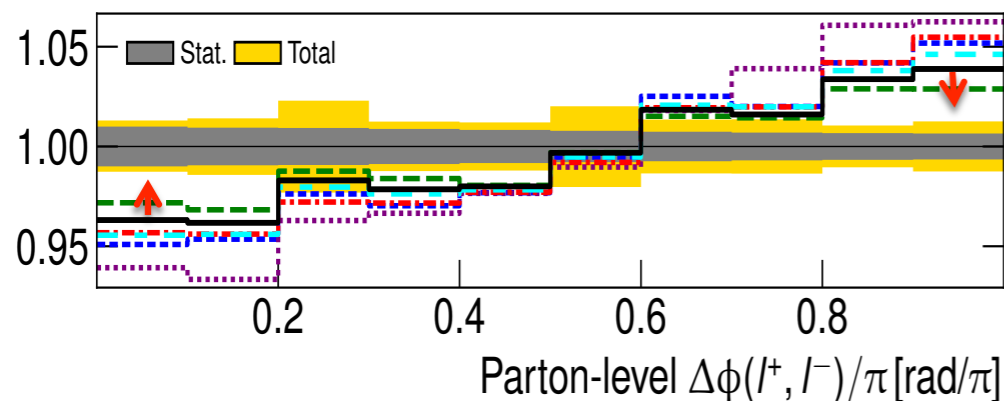
# ATLAS comparisons (I)

## MCFM with NLO in decay



## Add NNLO prediction

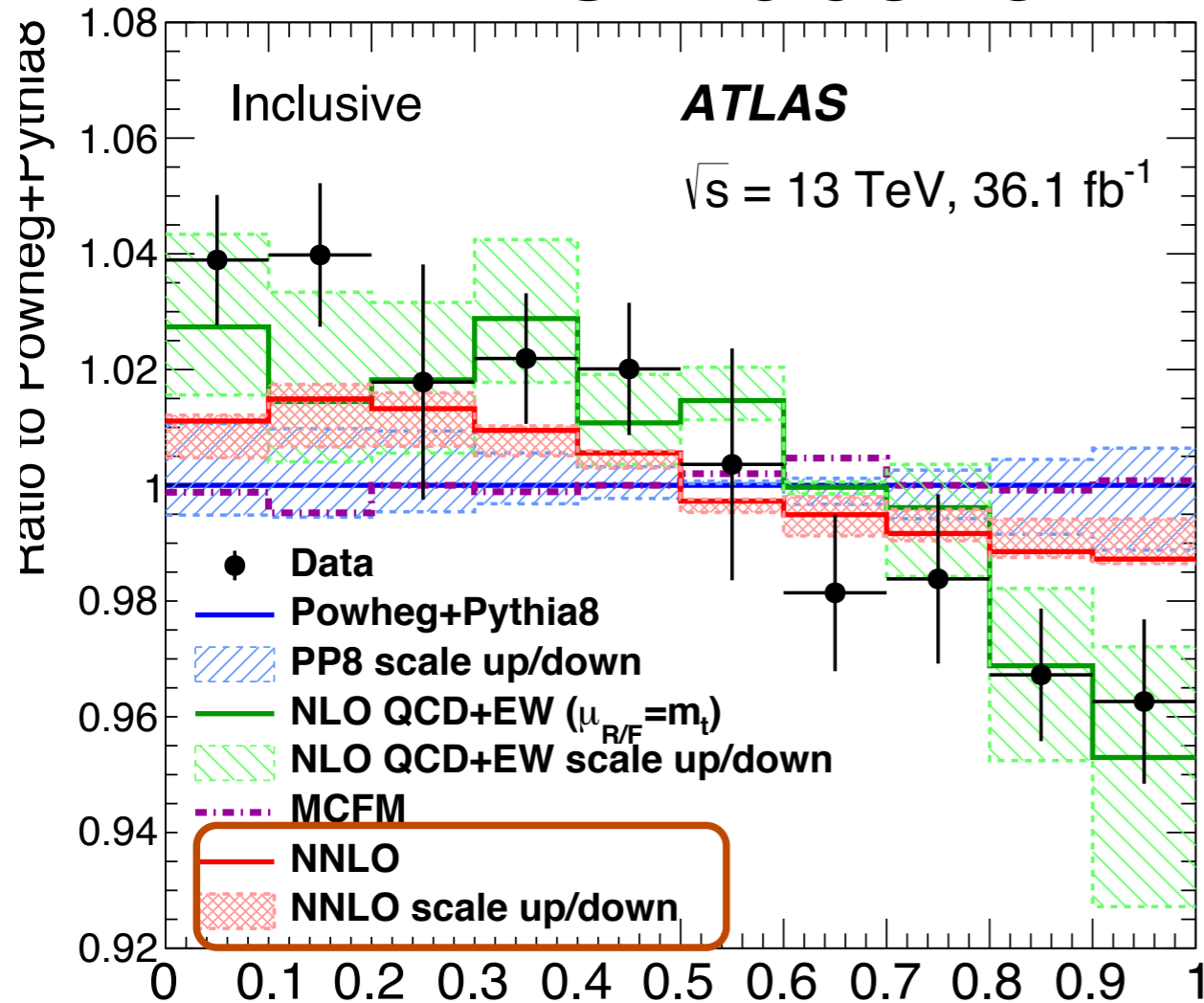
reweight top  $p_T$  to match NNLO



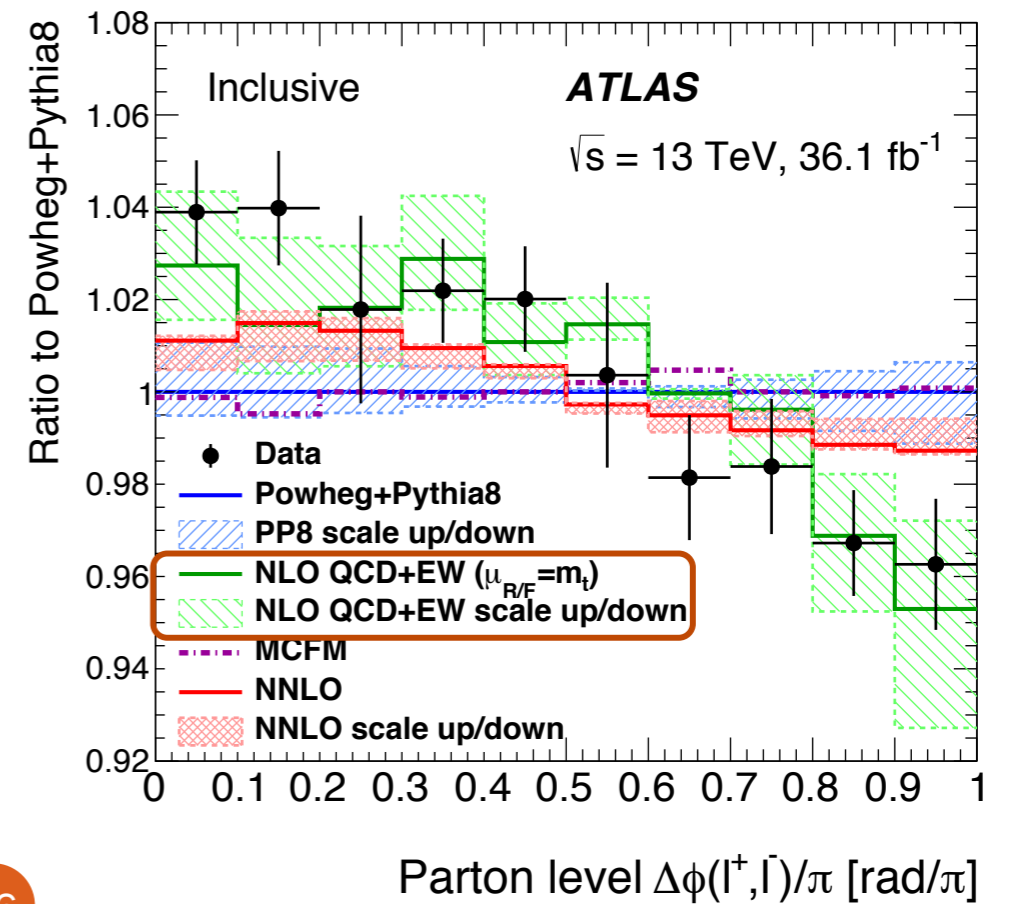


# ATLAS comparisons (II)

## NNLO inclusive



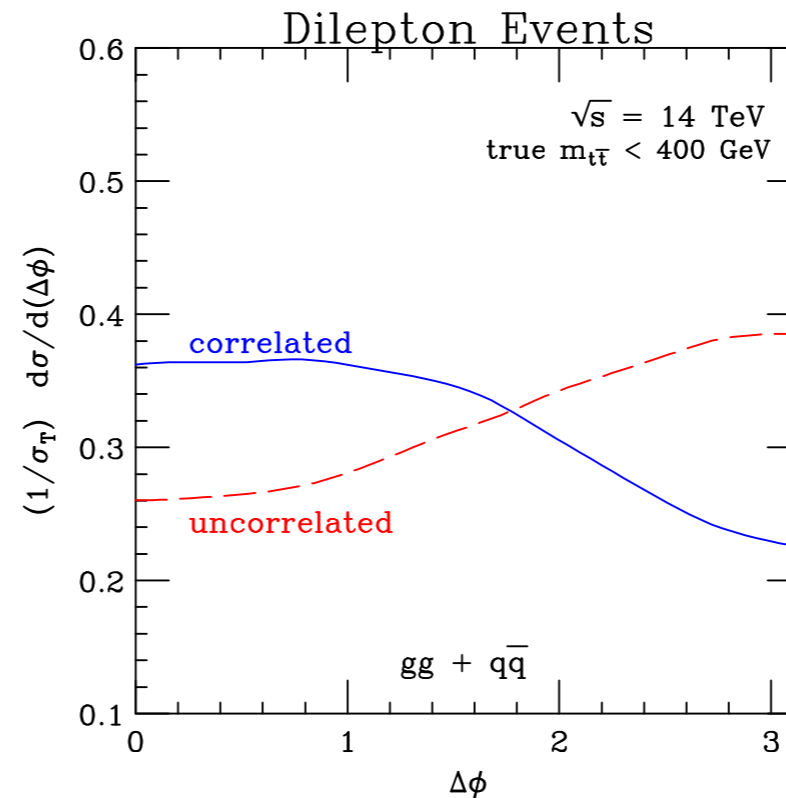
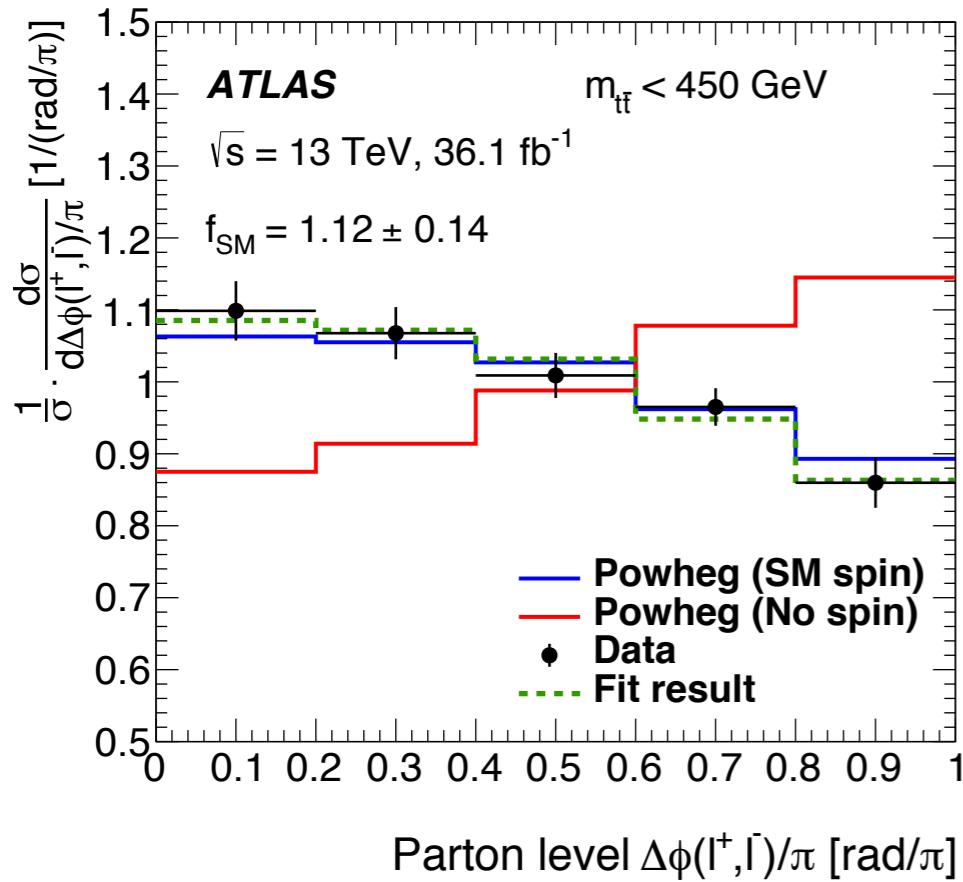
## NLO+EWK corr



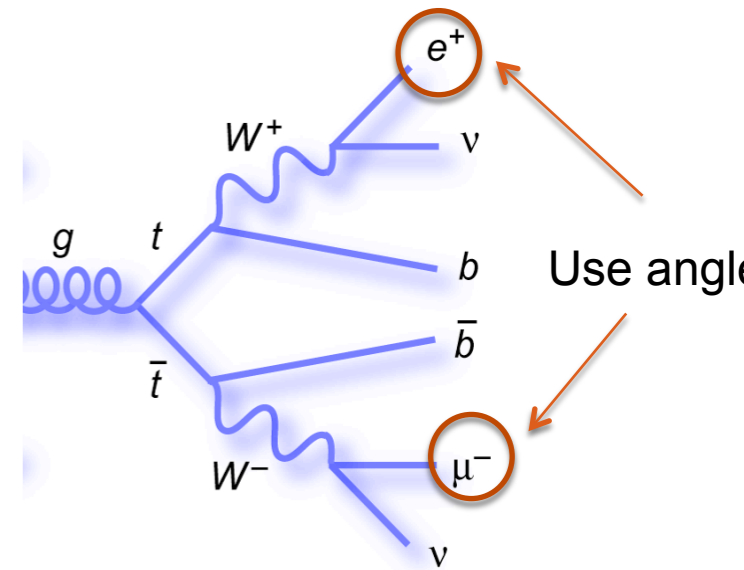
26

# Spin correlation with $\Delta\varphi$

- spin correlation sensitivity to  $\Delta\varphi$  is enhanced at low  $m_{t\bar{t}}$



Mahlon and Parke  
 Phys. Rev. D 81, 074024



- Reconstruct  $t\bar{t}$  final state : with constrains by mw and motion , test different eta assumptions for nus,select highest weight based on ETmiss expected resolution

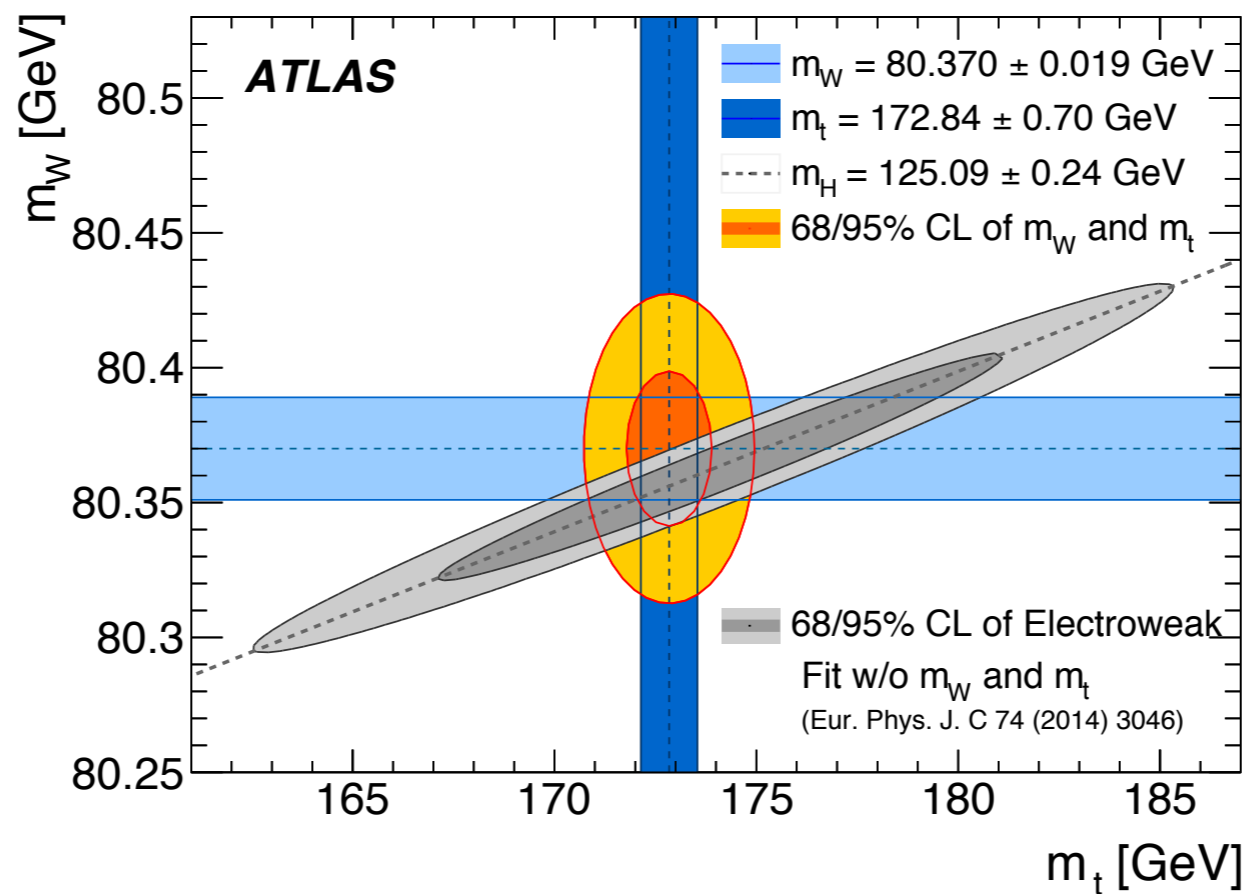
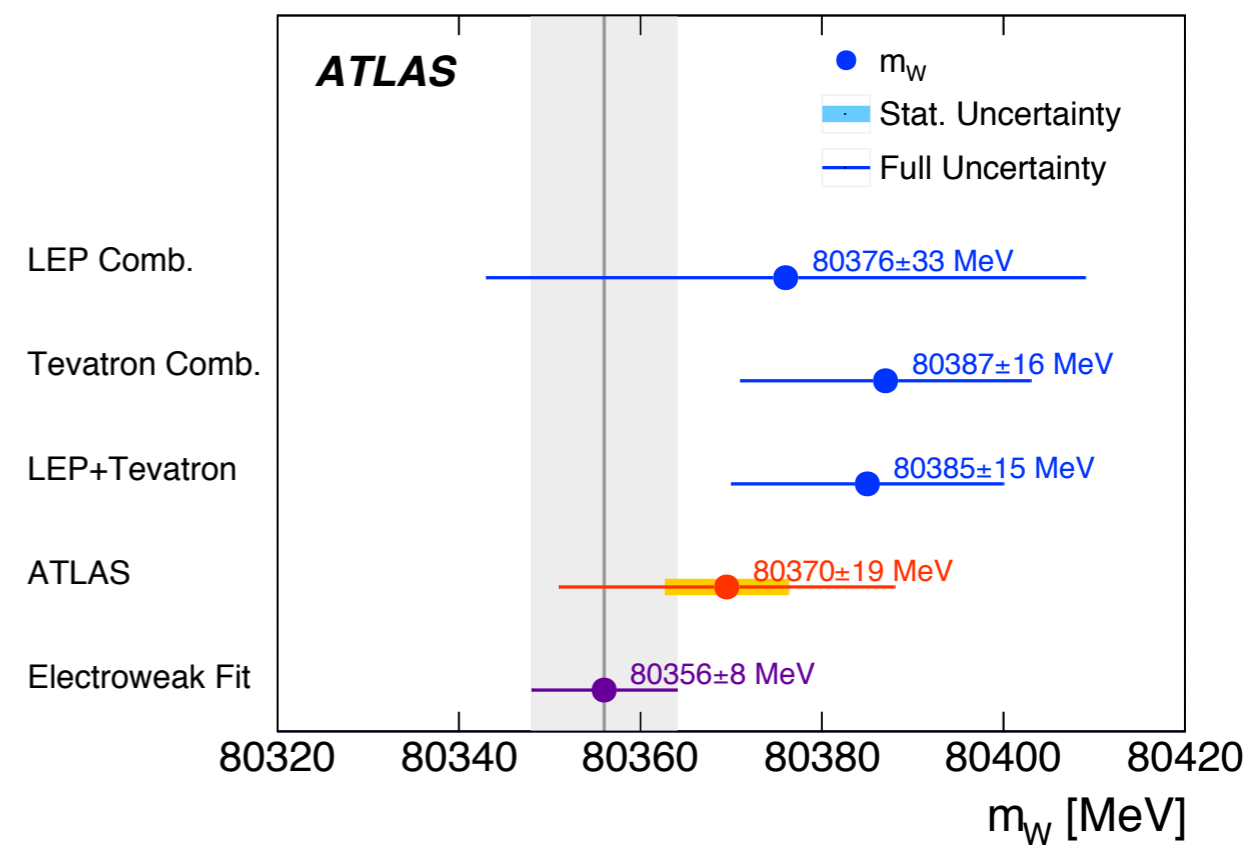
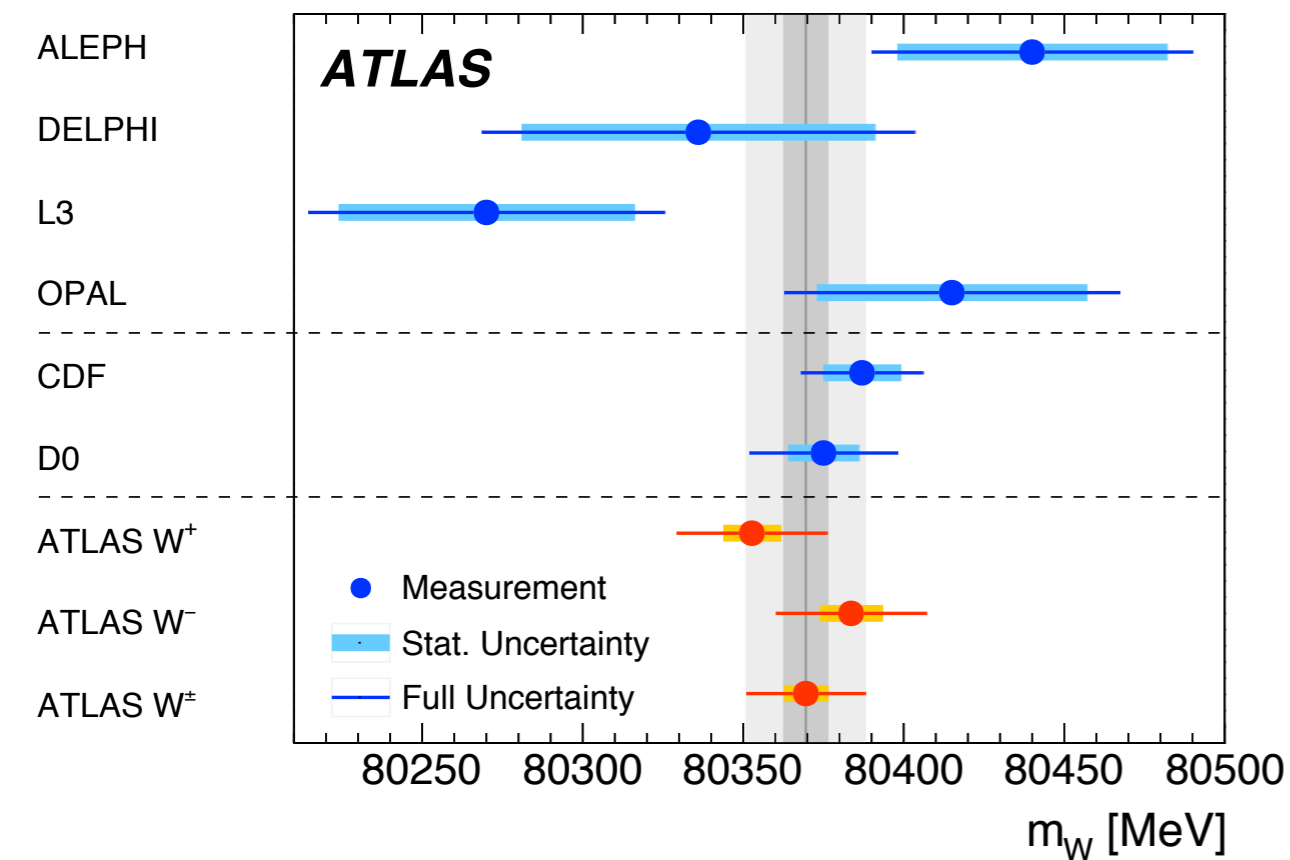
- Subtract bkg and unfold



# Top quark mass

# W mass from ATLAS

*arXiv:1701.07240*



*$m_W$  and  $m_{top}$  are from ATLAS measurements*

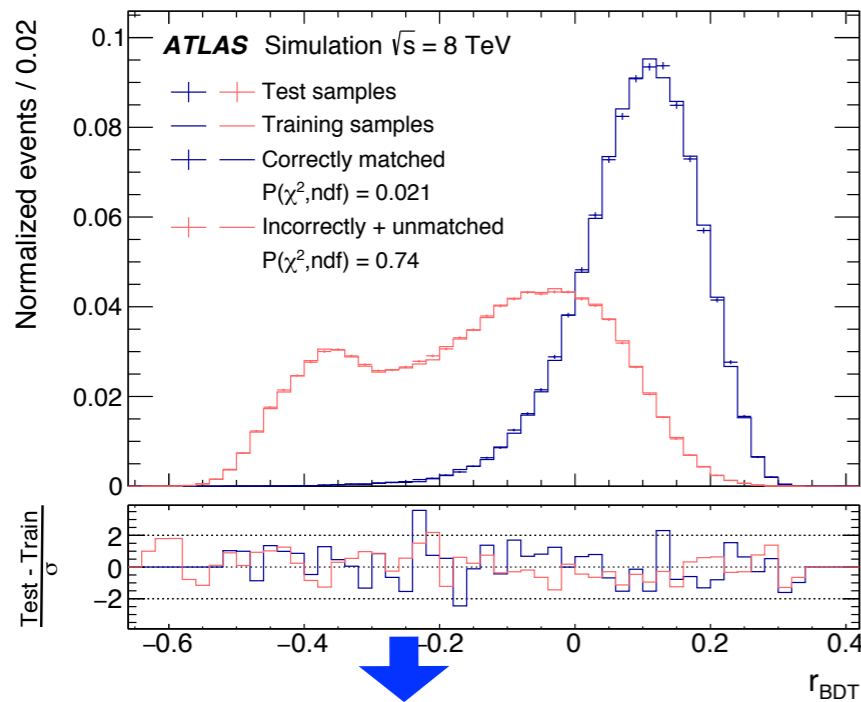
# Special reasons: Measure the top quark mass

most recent ATLAS from all data @  $\sqrt{s}=8$  TeV

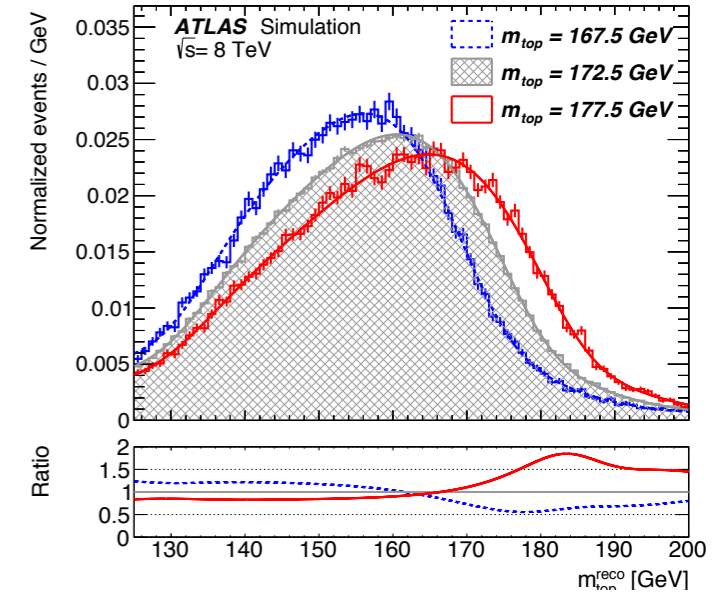
Eur. Phys. J. C79 (2019) 290

- Select  $\ell+\geq 4$  jets events (subtract  $W$ +jets, fakes, dibosons & single top top), require 2 b-jets
- Likelihood-based kinematic fit ( $m_W$  &  $m_{top}$  constr)
- Kine variables  $\rightarrow$  **Boosted Decision Tree**  $\rightarrow$  discriminant to select correct jet/lepton assignment.

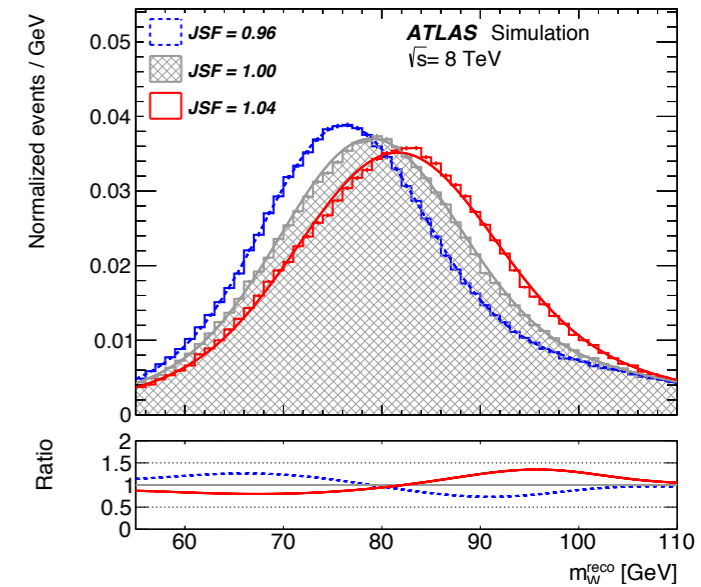
assign jets/ leptons to  $t\bar{t}$  decay products



•  $m_{top}^{reco}$  sensitive to  $m_{top}$ , JSF, bJSF

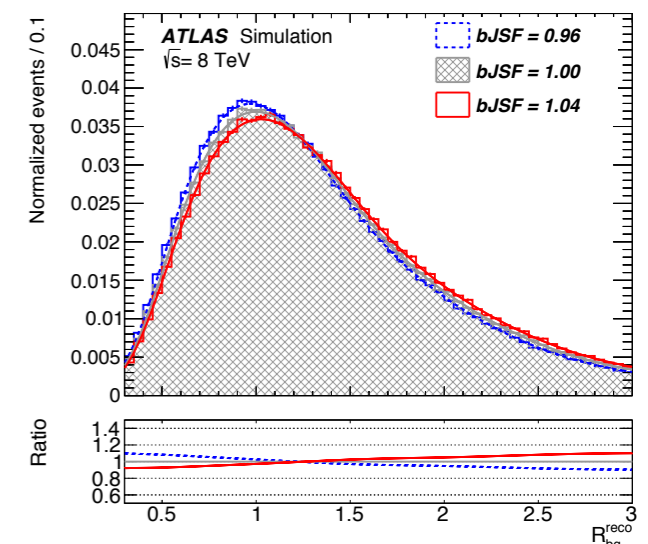


•  $m_W^{reco}$  sensitive to JSF



• likelihood-fit 3  $m_{top}$  sensitive variables to data

•  $R = \frac{\sum p_{T,b-jets}}{\sum p_{T,jets\_in\_W}}$  sensitive to bJSF

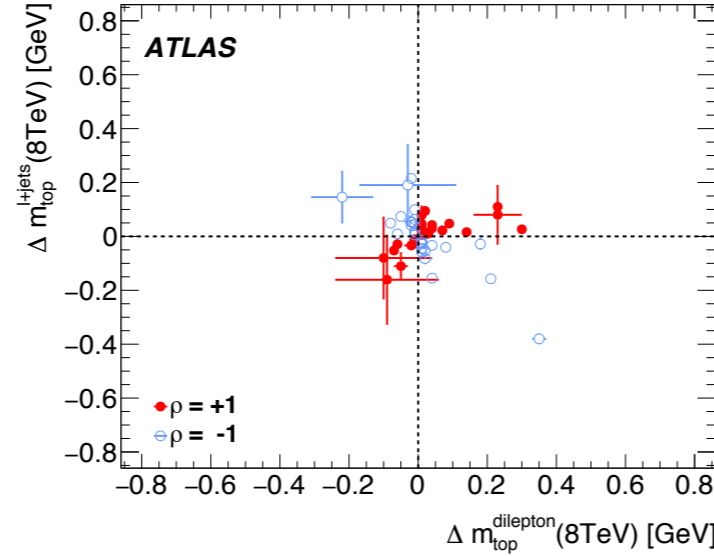
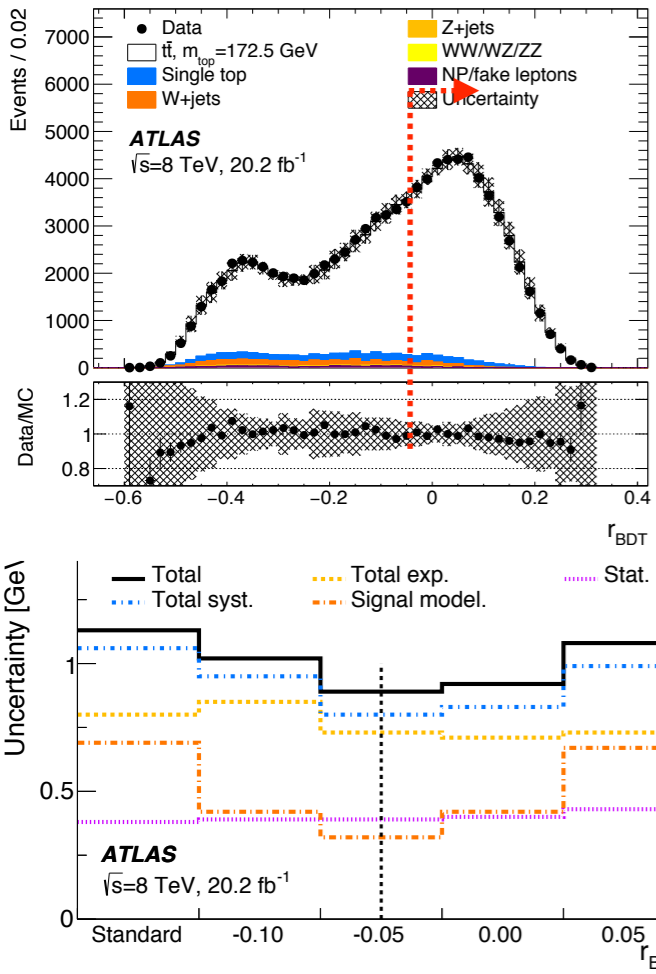


•  $m_{top}$  + 2 scale factors: jet and b-jet-to-light-jet energy  $\rightarrow$  reduce dominant jet & b-jet uncertainties

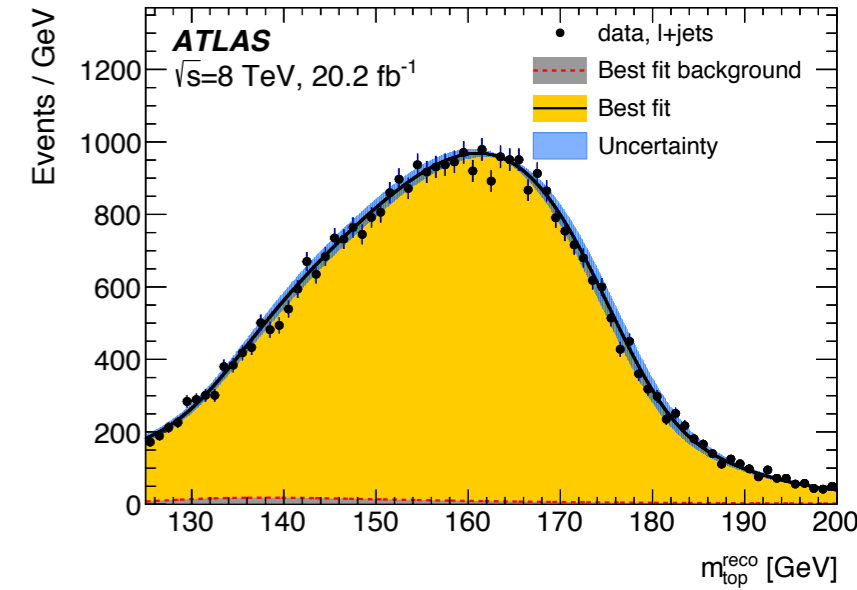
# Special reasons: Measure the top quark mass

*Eur. Phys. J. C79 (2019) 290*

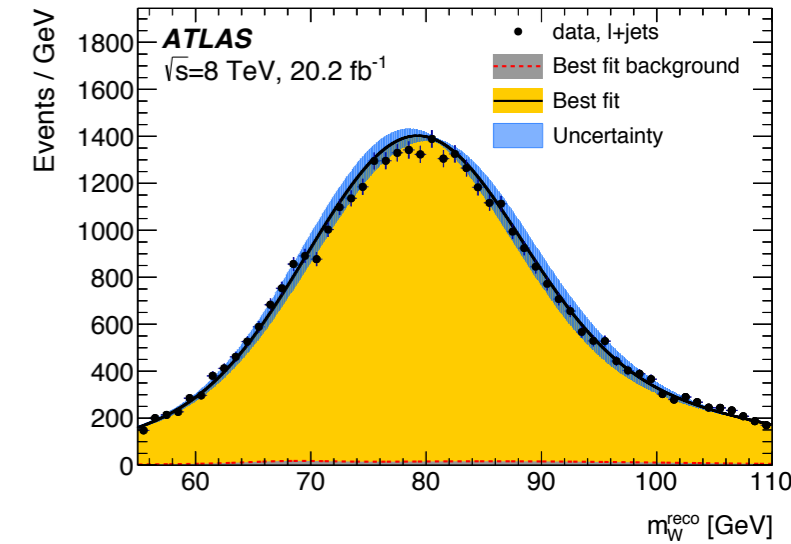
- **Optimize fit w.r.t. BDT: 19% improvement**



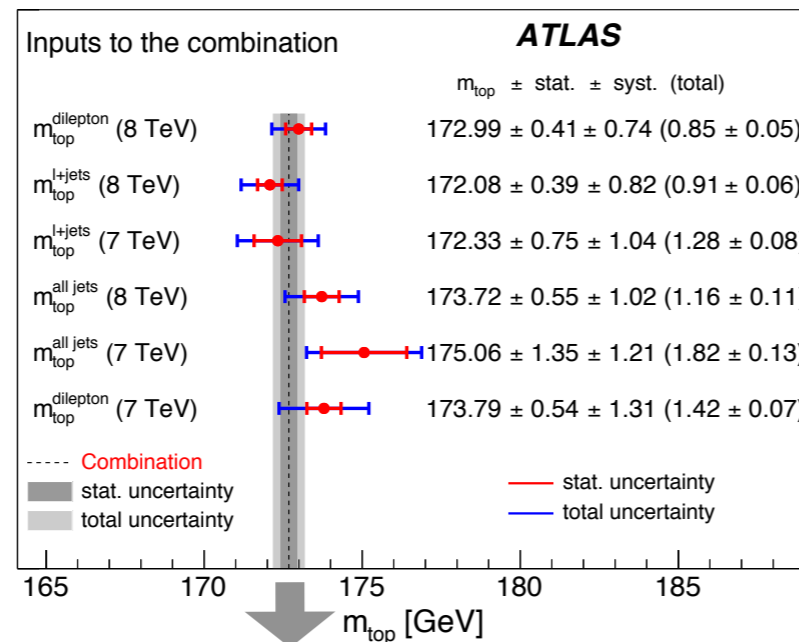
•  $m_{top}^{reco}$  sensitive to  $m_{top}$ , JSF, bJSF



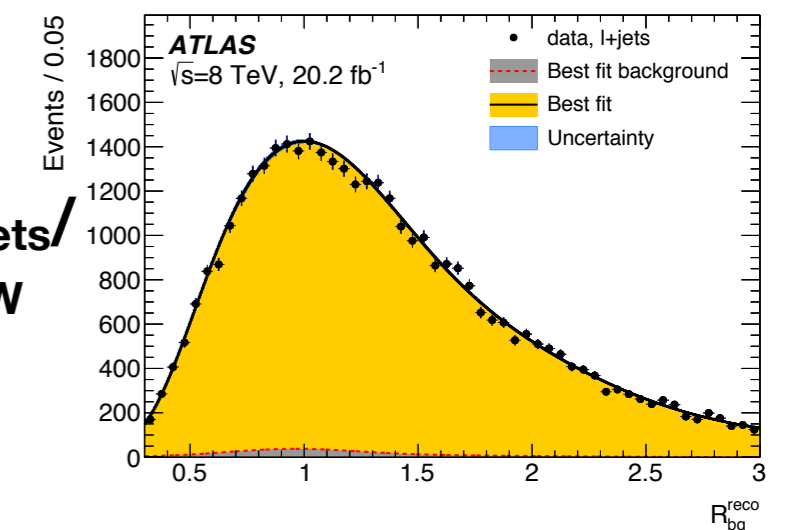
•  $m_W^{reco}$  sensitive to JSF



- **Combine with dilepton and all jets result**



•  $R = \frac{\sum p_{T,b-jets}}{\sum p_{T,jets\_in\_W}}$



$\delta m_{top} / m_{top} \sim 0.28\%$

**$172.69 \pm 0.25$  (stat)  $\pm 0.41$  (syst) GeV**

**Table 3** Systematic uncertainties in  $m_{\text{top}}$ . The measured values of  $m_{\text{top}}$  are given together with the statistical and systematic uncertainties in GeV for the standard and the BDT event selections. For comparison, the result in the  $t\bar{t} \rightarrow \text{lepton} + \text{jets}$  channel at  $\sqrt{s} = 7$  TeV from Ref. [9] is also listed. For each systematic uncertainty listed, the first value corresponds to the uncertainty in  $m_{\text{top}}$ , and the second to the statistical precision in this uncertainty. An integer value of zero means that the corresponding uncertainty is negligible and therefore not evaluated. Statistical uncertainties quoted as 0.00 are smaller than 0.005. The statistical uncertainty in the total systematic uncertainty is calculated from uncertainty propagation. The last line refers to the sum in quadrature of the statistical and systematic uncertainties

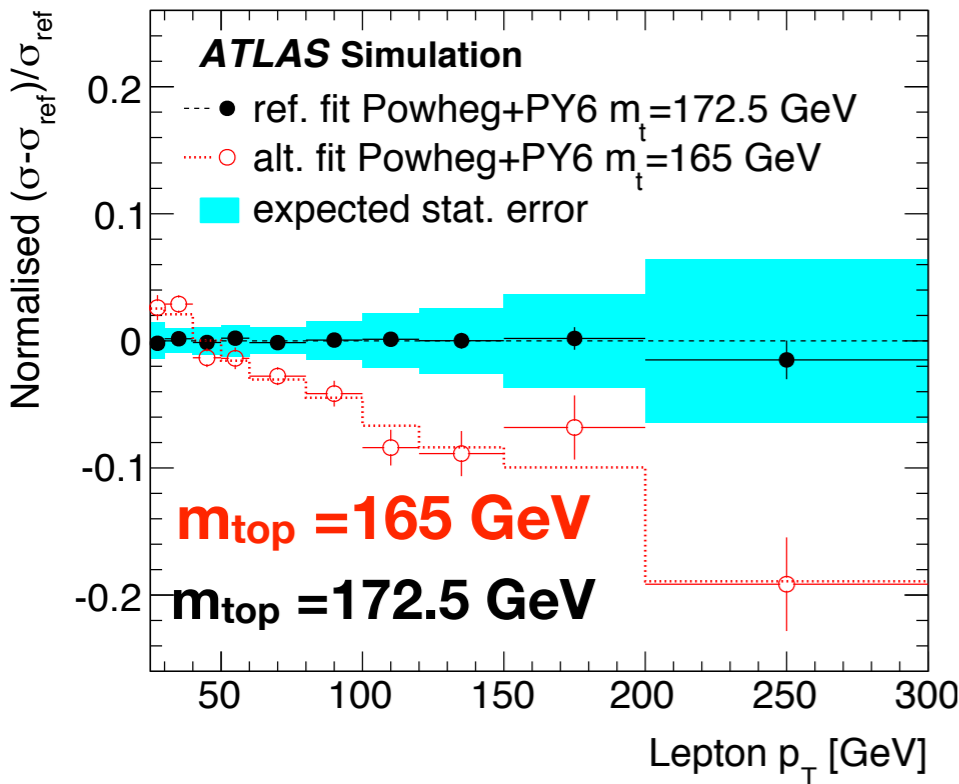
Event selection	$\sqrt{s} = 7$ TeV	$\sqrt{s} = 8$ TeV	
	Standard	Standard	BDT
$m_{\text{top}}$ result [GeV]	172.33	171.90	172.08
Statistics	0.75	0.38	0.39
– Stat. comp. ( $m_{\text{top}}$ )	0.23	0.12	0.11
– Stat. comp. (JSF)	0.25	0.11	0.11
– Stat. comp. (bJSF)	0.67	0.34	0.35
Method	$0.11 \pm 0.10$	$0.04 \pm 0.11$	$0.13 \pm 0.11$
Signal Monte Carlo generator	$0.22 \pm 0.21$	$0.50 \pm 0.17$	$0.16 \pm 0.17$
Hadronization	$0.18 \pm 0.12$	$0.05 \pm 0.10$	$0.15 \pm 0.10$
Initial- and final-state QCD radiation	$0.32 \pm 0.06$	$0.28 \pm 0.11$	$0.08 \pm 0.11$
Underlying event	$0.15 \pm 0.07$	$0.08 \pm 0.15$	$0.08 \pm 0.15$
Colour reconnection	$0.11 \pm 0.07$	$0.37 \pm 0.15$	$0.19 \pm 0.15$
Parton distribution function	$0.25 \pm 0.00$	$0.08 \pm 0.00$	$0.09 \pm 0.00$
Background normalization	$0.10 \pm 0.00$	$0.04 \pm 0.00$	$0.08 \pm 0.00$
$W$ +jets shape	$0.29 \pm 0.00$	$0.05 \pm 0.00$	$0.11 \pm 0.00$
Fake leptons shape	$0.05 \pm 0.00$	0	0
Jet energy scale	$0.58 \pm 0.11$	$0.63 \pm 0.02$	$0.54 \pm 0.02$
Relative $b$ -to-light-jet energy scale	$0.06 \pm 0.03$	$0.05 \pm 0.01$	$0.03 \pm 0.01$
Jet energy resolution	$0.22 \pm 0.11$	$0.23 \pm 0.03$	$0.20 \pm 0.04$
Jet reconstruction efficiency	$0.12 \pm 0.00$	$0.04 \pm 0.01$	$0.02 \pm 0.01$
Jet vertex fraction	$0.01 \pm 0.00$	$0.13 \pm 0.01$	$0.09 \pm 0.01$
$b$ -tagging	$0.50 \pm 0.00$	$0.37 \pm 0.00$	$0.38 \pm 0.00$
Leptons	$0.04 \pm 0.00$	$0.16 \pm 0.01$	$0.16 \pm 0.01$
Missing transverse momentum	$0.15 \pm 0.04$	$0.08 \pm 0.01$	$0.05 \pm 0.01$
Pile-up	$0.02 \pm 0.01$	$0.14 \pm 0.01$	$0.15 \pm 0.01$
Total systematic uncertainty	$1.04 \pm 0.08$	$1.07 \pm 0.10$	$0.82 \pm 0.06$
Total	$1.28 \pm 0.08$	$1.13 \pm 0.10$	$0.91 \pm 0.06$

# Measuring SM variables: top mass

*Eur. Phys. J. C 77 (2017) 804*

$d\sigma_{tt}/dX$  are sensitive to  $m_{top}$

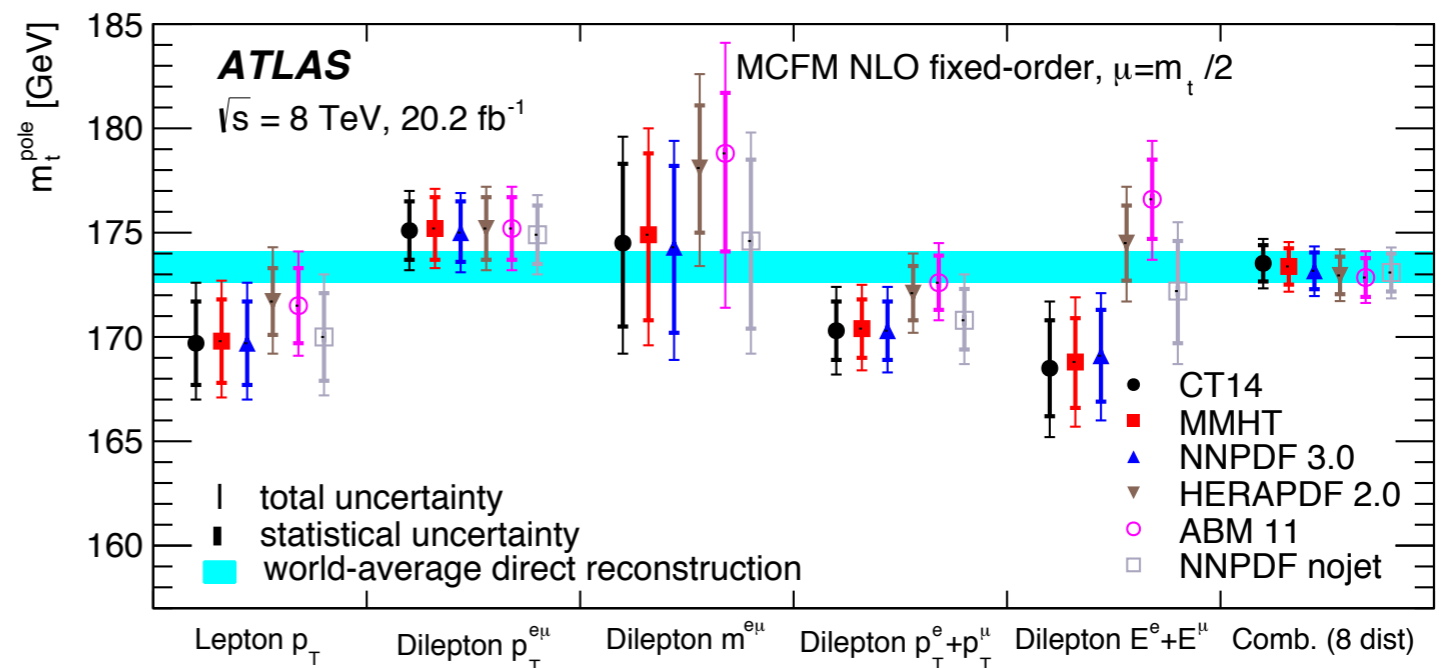
Extract  $m_{top}$  by minimising



$$\chi^2(m_{top}) = (\mathbf{X} - \mathbf{P}(m_{top}))^T \mathbf{C}^{-1} (\mathbf{X} - \mathbf{P}(m_{top}))$$

- $\mathbf{X}$ =measured  $d\sigma_{tt}/dX$ ,  $\mathbf{P}$ = predictions:NLO +PS or fixed order NLO (pole scheme)
- $\mathbf{C}$ = covariance matrix

@ particle level



- Final  $m_t^{pole}$ : combined (8 dist.) NNPDF fit
- Add  $\delta m_t^{pole}$ : PDF (0.3 GeV), dynamic vs static scale (1.1 GeV),  $\delta\alpha_s \sim$ negligible

$$m_t^{pole} = 173.2 \pm 0.9 \pm 0.8 \pm 1.2 \text{ GeV} \quad \delta m_t/m_t \sim 0.5\%$$

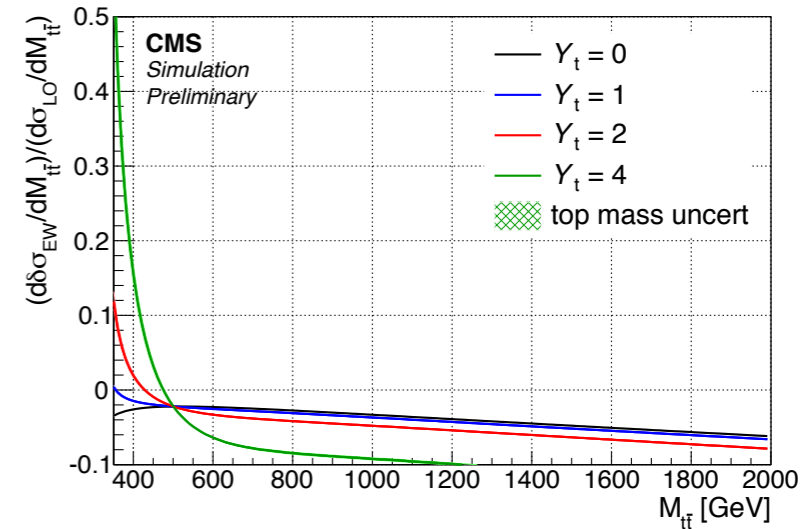
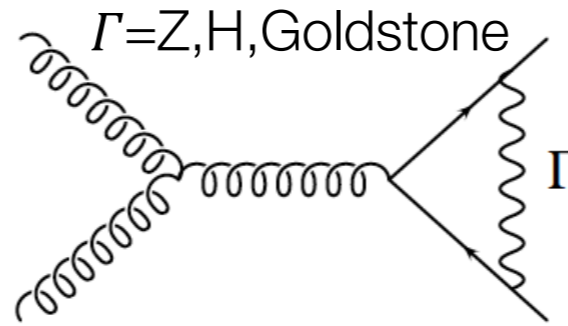
stat   exp   theo

Comparable to  $\delta m_t/m_t \sim 0.28\%$  from standard method (NLO+PS template @ reco level)



# Top Yukawa coupling

- @ threshold  $tt$ -production, Higgs boson mediator in  $Htt$  loop at  $\alpha_s^2 \alpha$   
 $\rightarrow d\sigma_{tt}/dX$  depends on  $Y_t^2$



- Require 1  $\ell$  (e, mu),  $\geq 3$  jets,  $\geq 2$  b-tag(s),  $m_T^W < 140$  GeV, if  $N_{\text{jets}} = 3$   $p_{T,\text{lead b-jet}} > 50$  GeV
- Bkg: data driven multi jets (from control region), simulated single top, W/Z+jets,
- Reconstruct  $tt$  system by likelihood discriminant: extend to 3-jets!**

- b-jets with largest b-tag weight  $\rightarrow$  b-quarks (2 possibilities)

$\geq 4$  jets

$= 3$  jets,

*Prob(missing jet is from W) : 93%*

- Derive  $\nu$  momentum: point on ellipse in 3-mom  $\nu$  space from  $(m_W, m_{\text{top}})$  intersection with minimum  $D_{\nu, \text{min}}$  distance of  $(x, y)$  projection from  $p_T^{\text{miss}}$

- find  $D_{\nu, \text{min}}$  for b-jet assignment
- Discard assignment with no  $D_{\nu, \text{min}}$
- Define discriminant

- $m(b, \ell) > m_{\text{top}} \rightarrow$  discard assignment

- Define discriminant

$$-\ln(\lambda_4) = -\ln(P_m(m_2, m_3)) - \ln(P_\nu(D_{\nu, \text{min}}))$$

prob to reconstruct  $m_W$  and  $m_{\text{top}}$  correct  $D_{\nu, \text{min}}$

$$-\ln(\lambda_3) = -\ln(P_{m_{t_h}}) - \ln(P_\nu(D_{\nu, \text{min}}))$$

prob to reconstruct  $m(b_h + W\text{-jet})$  correct  $D_{\nu, \text{min}}$

84% (69%) correct in 4-(5-)jets

80% correct in 3-jets

**choose jet assignment with maximum discriminant**

➔ Inclusion of 3-jets events: higher yield at sensitive low  $M_{tt}$ , reduce migration in  $N_{jet}$  → smaller JES/Had uncertainties

- Build binned likelihood for  $dN/dM_{tt} d(y_t - y_{anti-t})$  as function of EW correction strength,  $R = N_{tt}(Y_t) / N_{tt}(POWHEG)$ , a bin-dependent quadratic function of  $Y_t$  for 3,4,5 jets and all events

$$\mathcal{L} = \prod_{\text{bin} \in (M_{t\bar{t}}, |\Delta y_{t\bar{t}}|)} \mathcal{L}_{\text{bin}} = \prod_{\text{bin}} \text{Pois}(n_{\text{obs}}^{\text{bin}} | s^{\text{bin}}(\theta) \times R^{\text{bin}}(Y_t) + b^{\text{bin}}(\theta)) \times \rho(\theta | \tilde{\theta})$$

Gaussian for nuisance par → syst uncertainties

- Scan likelihood to find minimum and upper limit:  $Y_t < 1.67$  @ 95% CL

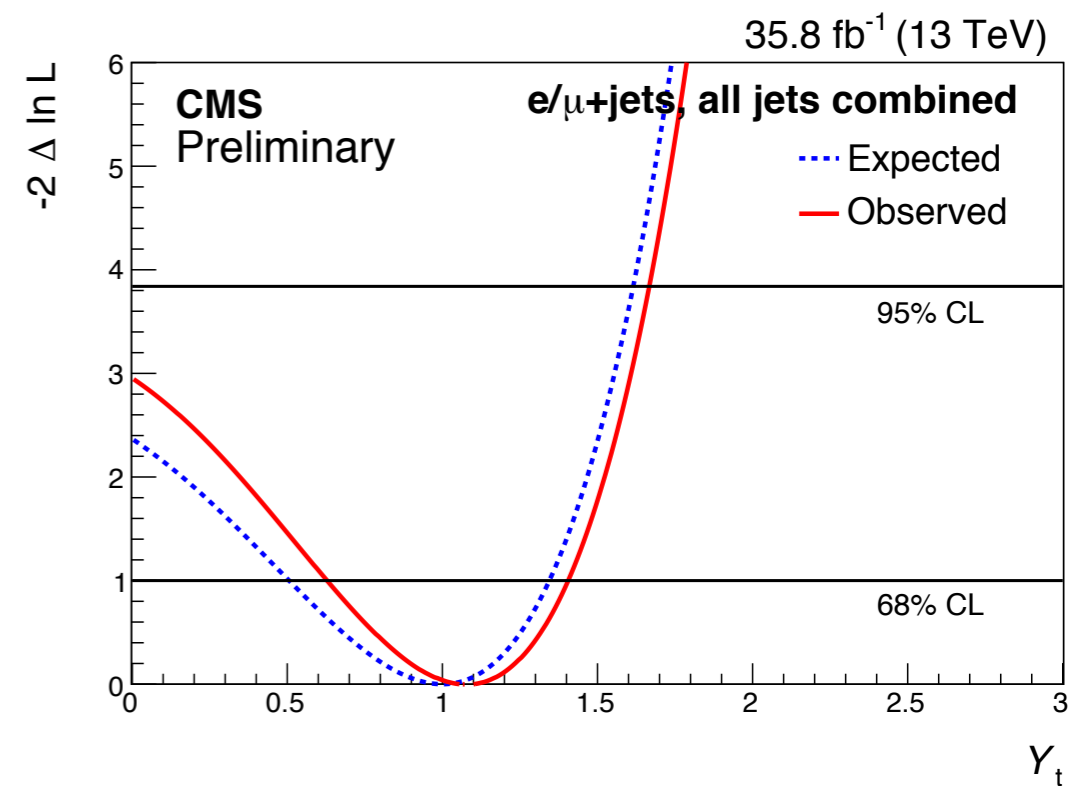
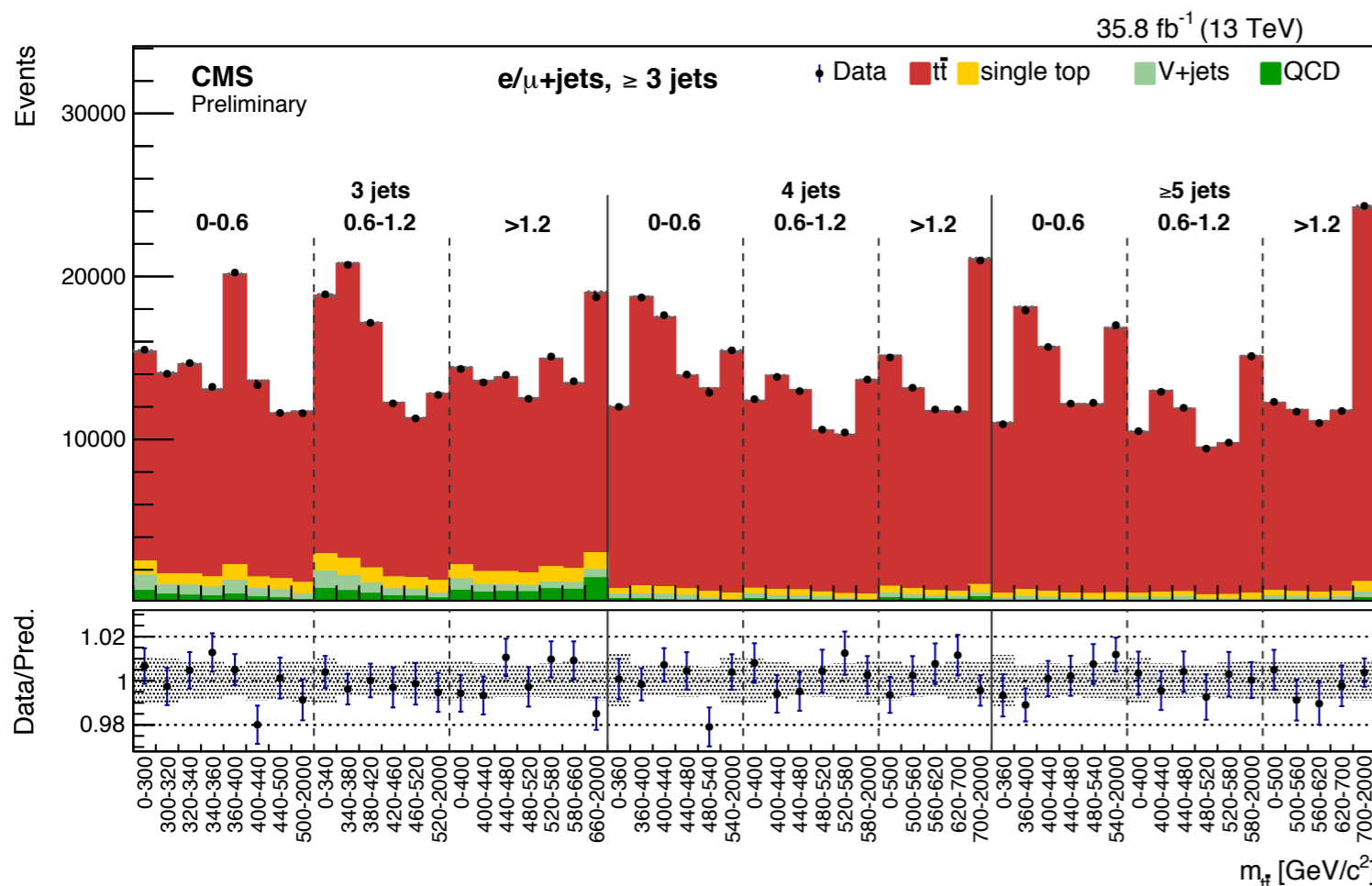


Table 3: The expected and observed 95% CL limits on  $Y_t$

Channel	Expected 95% CL	Observed 95% CL
3 jets	$Y_t < 2.17$	$Y_t < 2.59$
4 jets	$Y_t < 1.88$	$Y_t < 1.77$
5 jets	$Y_t < 2.03$	$Y_t < 2.23$
Combined	$Y_t < 1.62$	$Y_t < 1.67$

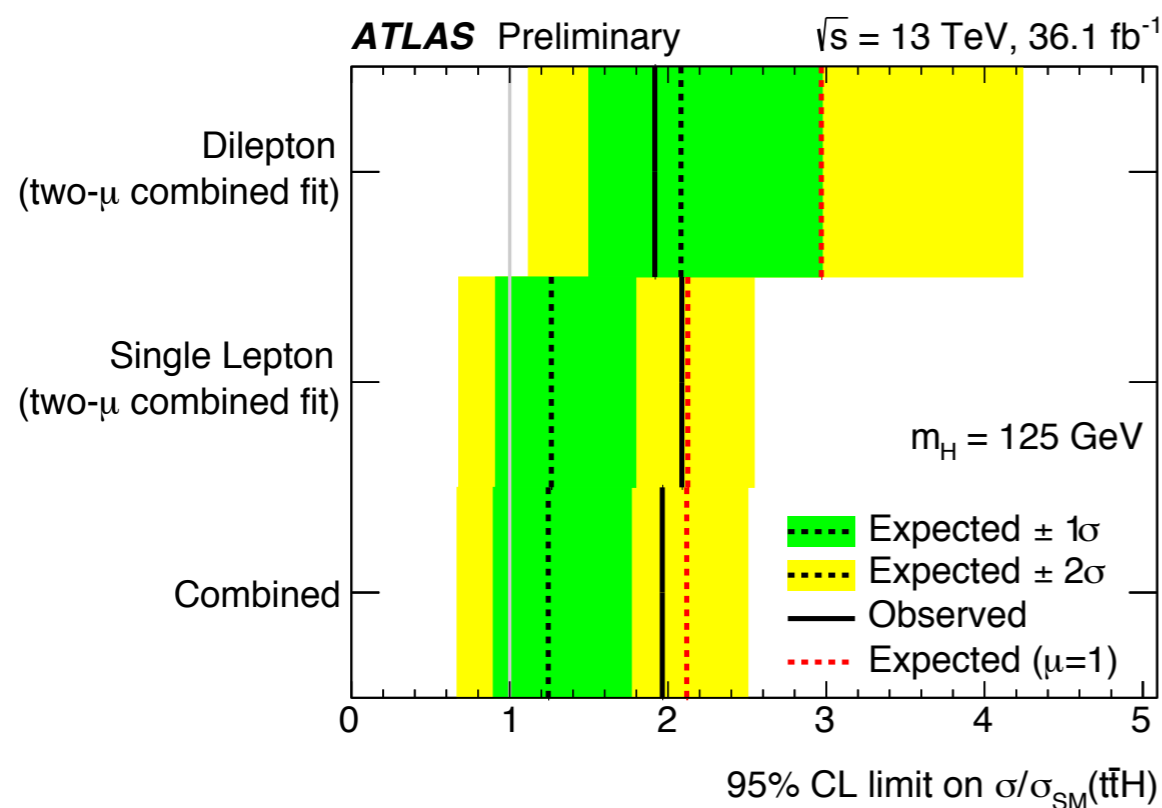
# Searches for BSM with top

# Example: Search for $t\bar{t}H$ , $H \rightarrow b\bar{b}$ @ $\sqrt{s} = 13\text{TeV}$

Given model for probability distributions for signal, bkg, systematic uncertainties ( $det \otimes theory$ ) build **likelihood-based variable  $f(\mu)$**  to

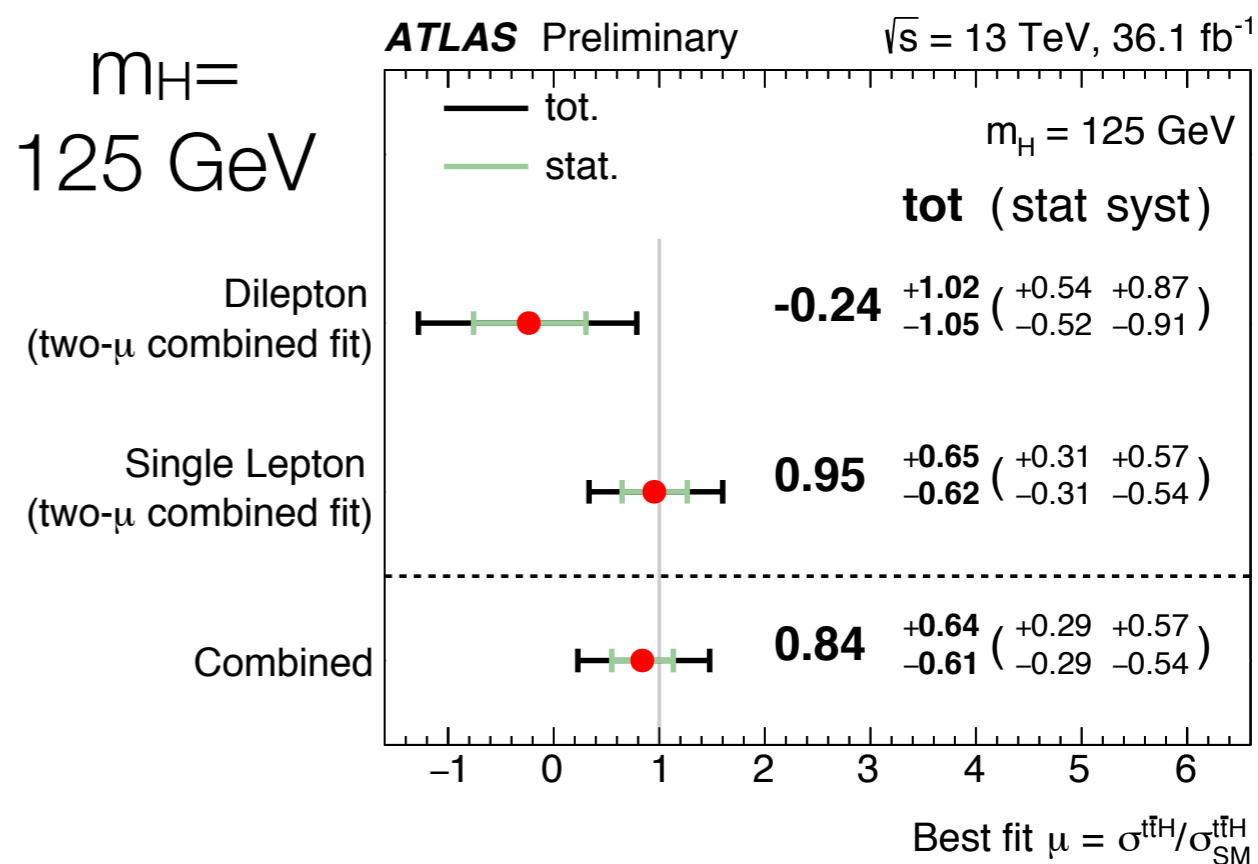
- in the bkg only hypothesis ( $\mu=0$ ), **derive** the probability to observe a more discrepant f-value than the observed one  $\rightarrow$  **test bkg hypothesis**
- in the sig+bkg hypothesis, **derive [0,upper limit] interval** that covers the “true”  $\mu$  value 95% of the times  $\rightarrow$  **test signal hypothesis**
- Estimate **signal strength** from maximum likelihood fit

**upper limit**  
*sig+bkg assumption*  
*bkg only assumption*



**signal strength**

assume  
 $m_H = 125\text{ GeV}$



**1.4 sigma significance w.r.t bkg only, excludes  $\mu_{t\bar{t}H} > 2$  at 95%CL**

## What is a “Confidence Interval?”

- you see them all the time:

Want to say there is a 68% chance that the true value of  $(m_W, m_t)$  is in this interval

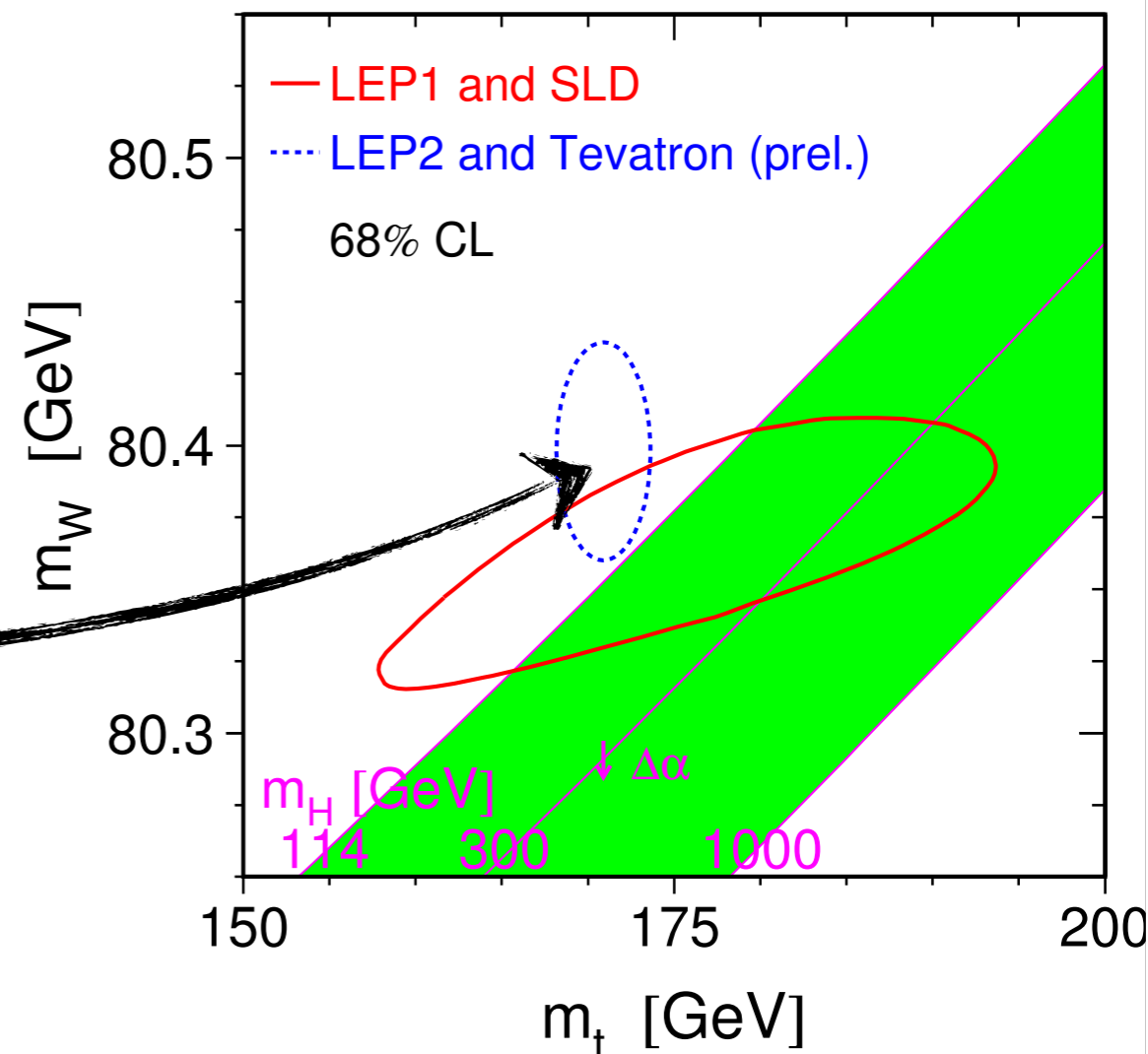
- but that’s  $P(\text{theory}|\text{data})!$

Correct frequentist statement is that the interval **covers** the true value 68% of the time

- remember, the contour is a function of the data, which is random. So it moves around from experiment to experiment

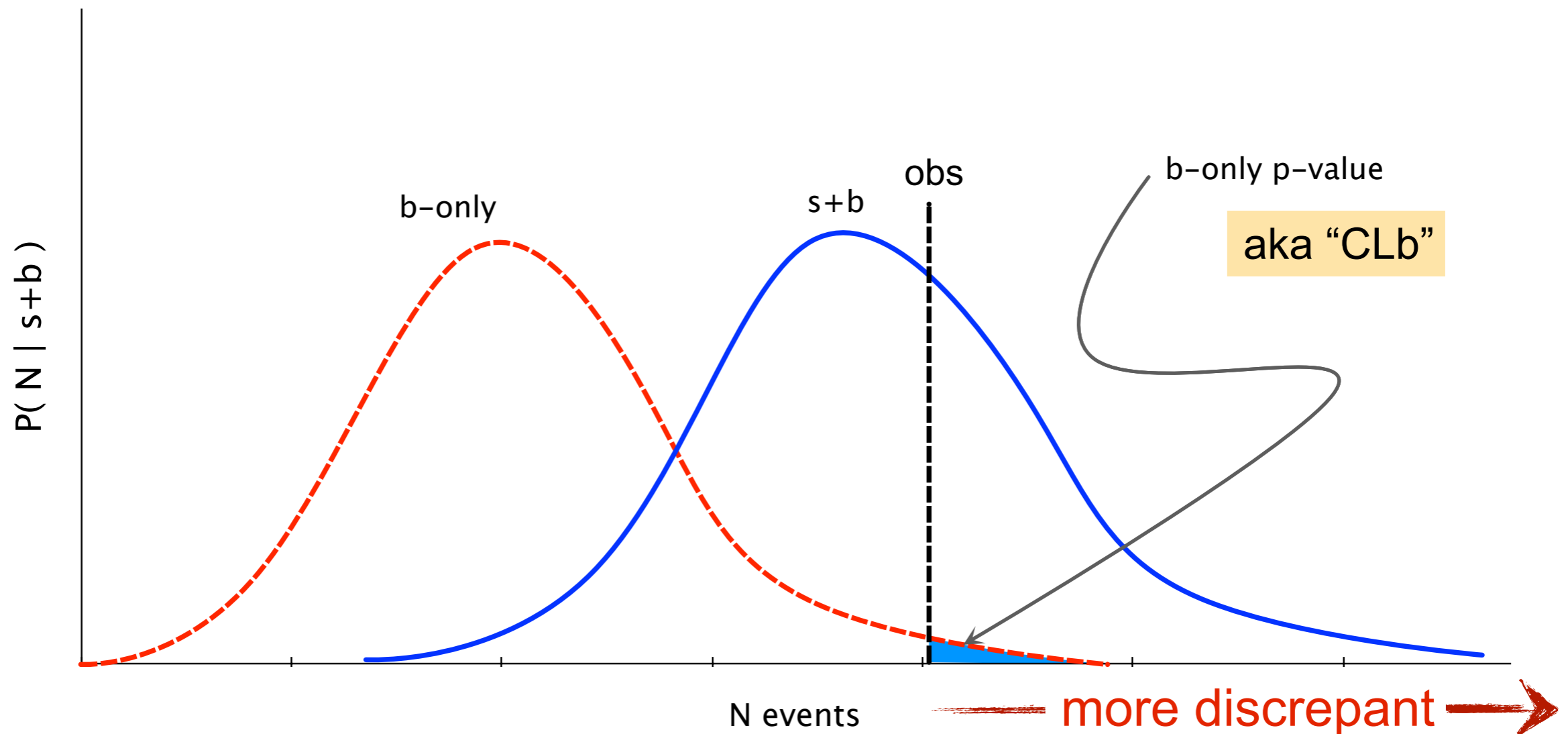
- Bayesian “credible interval” does mean probability parameter is in interval. The procedure is very intuitive:

$$P(\theta \in V) = \int_V \pi(\theta|x) = \int_V d\theta \frac{f(x|\theta)\pi(\theta)}{\int d\theta f(x|\theta)\pi(\theta)}$$



Discovery: test b-only (null:  $s=0$  vs. alt:  $s>0$ )

- note, **one-sided** alternative. larger  $N$  is “more discrepant”

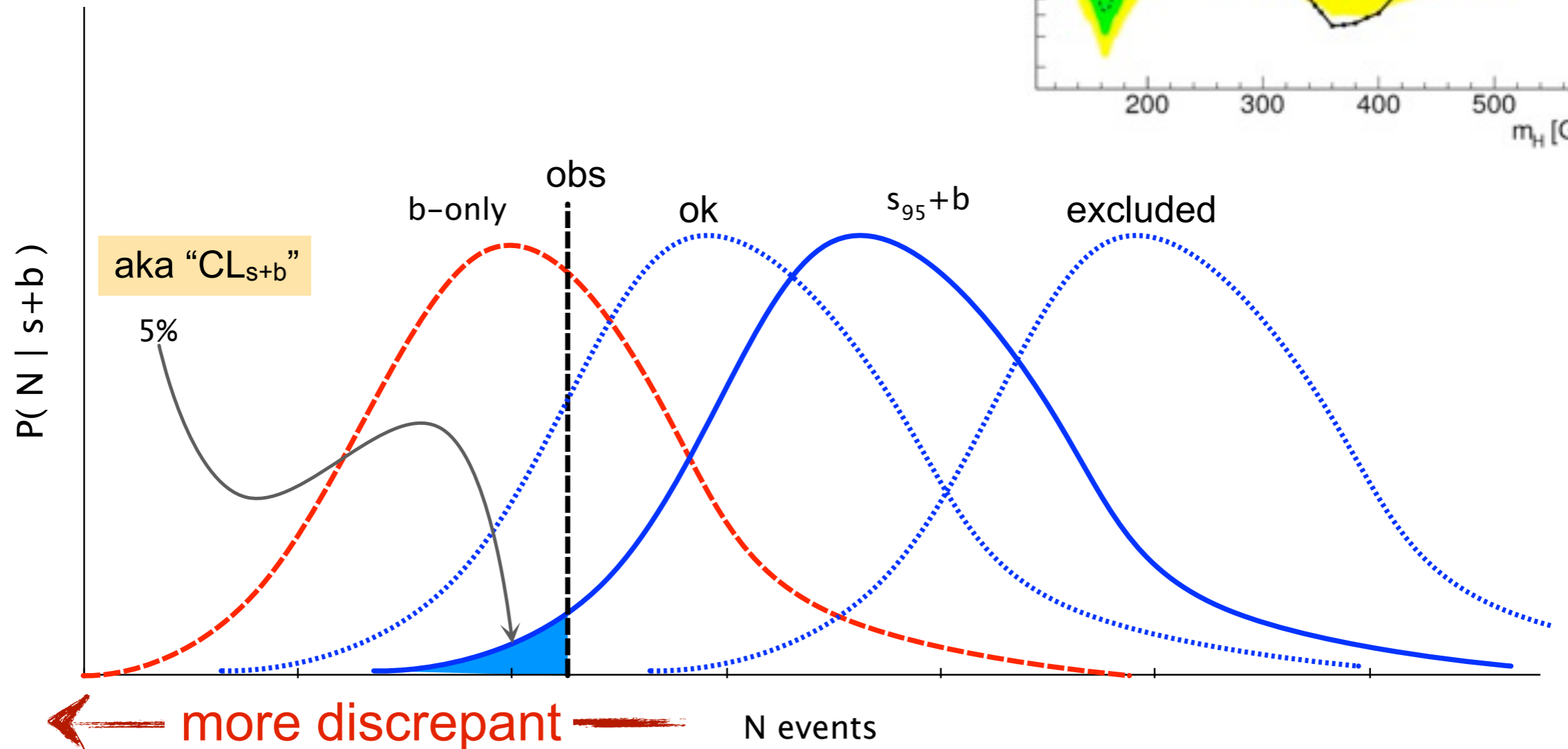
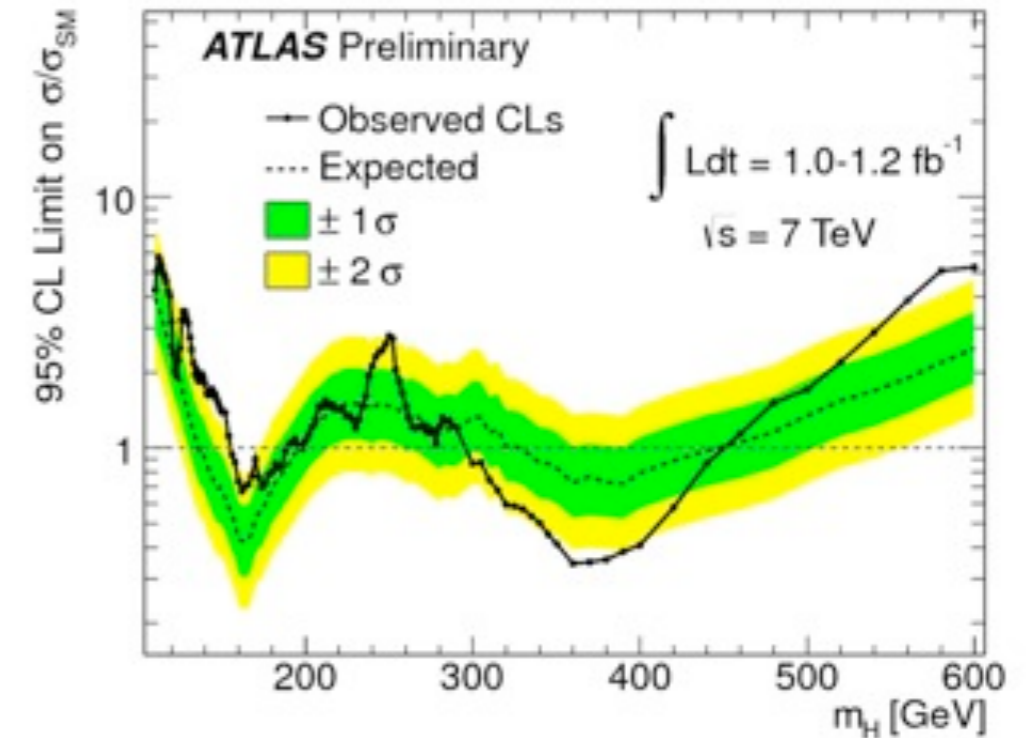


# Upper limits in pictures

What is meant by “95% upper limit” ?

See the picture below?

- ie. increase  $s$ , until the probability to have data “more discrepant” is  $< 5\%$

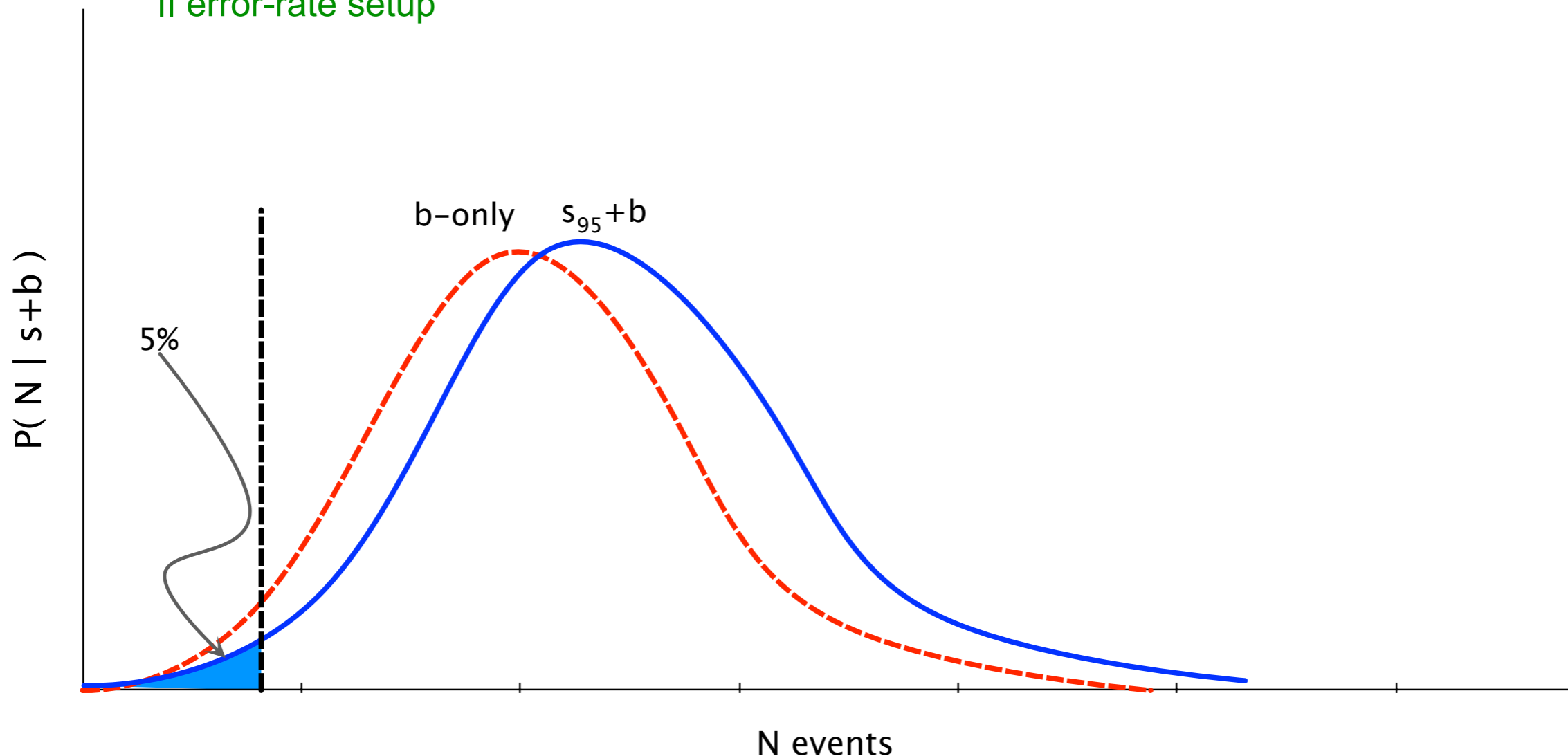




# The sensitivity problem

The physicist's worry about limits in general is that if there is a strong downward fluctuation, one might exclude arbitrarily small values of  $s$

- ▶ with a procedure that produces proper frequentist 95% confidence intervals, one should expect to exclude the true value of  $s$  5% of the time, no matter how small  $s$  is!
- ▶ This is not a problem with the procedure, but an undesirable consequence of the Type I / Type II error-rate setup

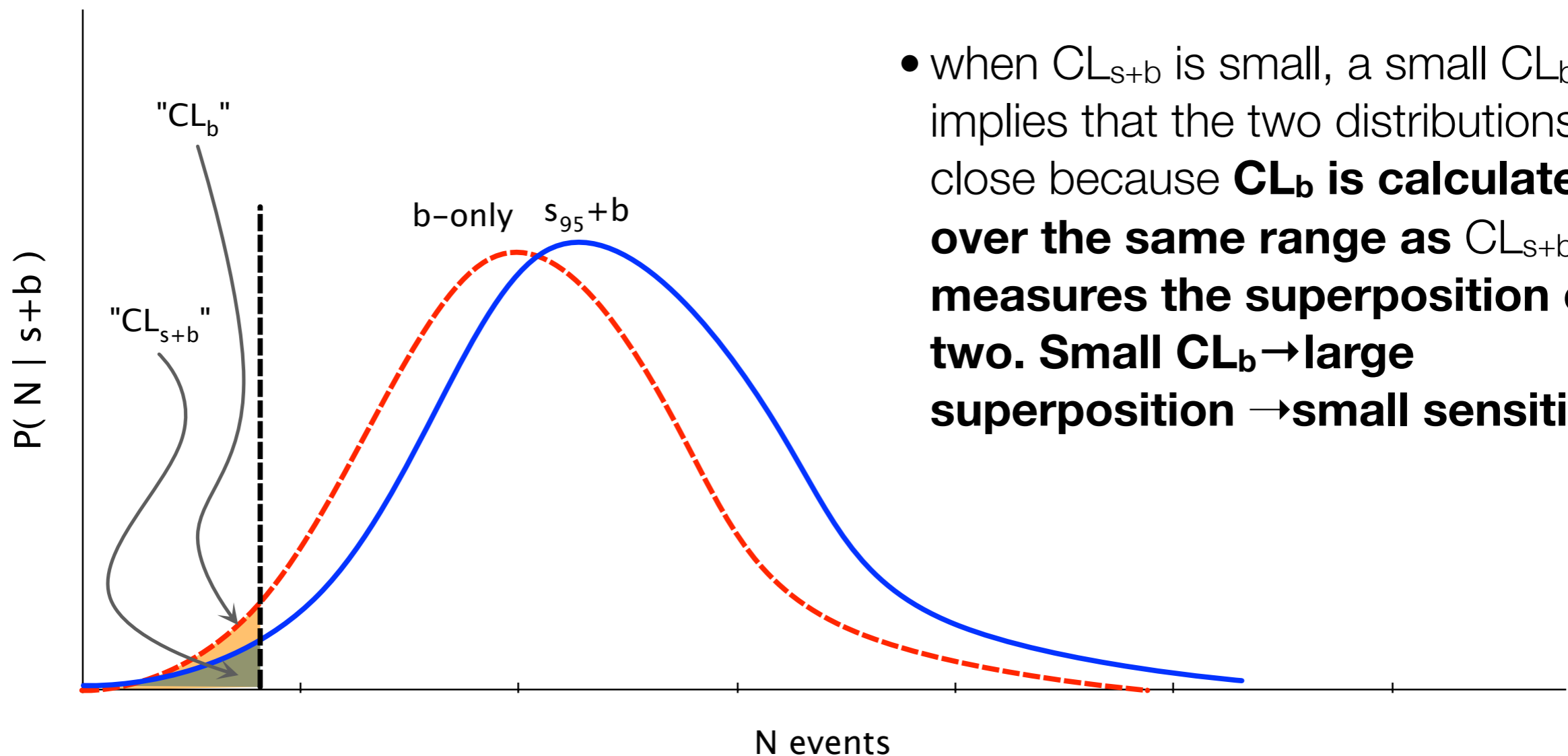


To address the sensitivity problem, CL<sub>s</sub> was introduced

- ▶ common (misused) nomenclature:  $CL_s = CL_{s+b}/CL_b$
- ▶ idea: only exclude if  $CL_s < 5\%$  (if  $CL_b$  is small,  $CL_s$  gets bigger)

CL<sub>s</sub> is known to be “conservative” (over-cover): expected limit covers with 97.5%

- Note: CL<sub>s</sub> is NOT a probability



- when  $CL_{s+b}$  is small, a small  $CL_b$  implies that the two distributions are close because **CL<sub>b</sub> is calculated over the same range as  $CL_{s+b}$  : it measures the superposition of the two. Small  $CL_b \rightarrow$  large superposition  $\rightarrow$  small sensitivity**

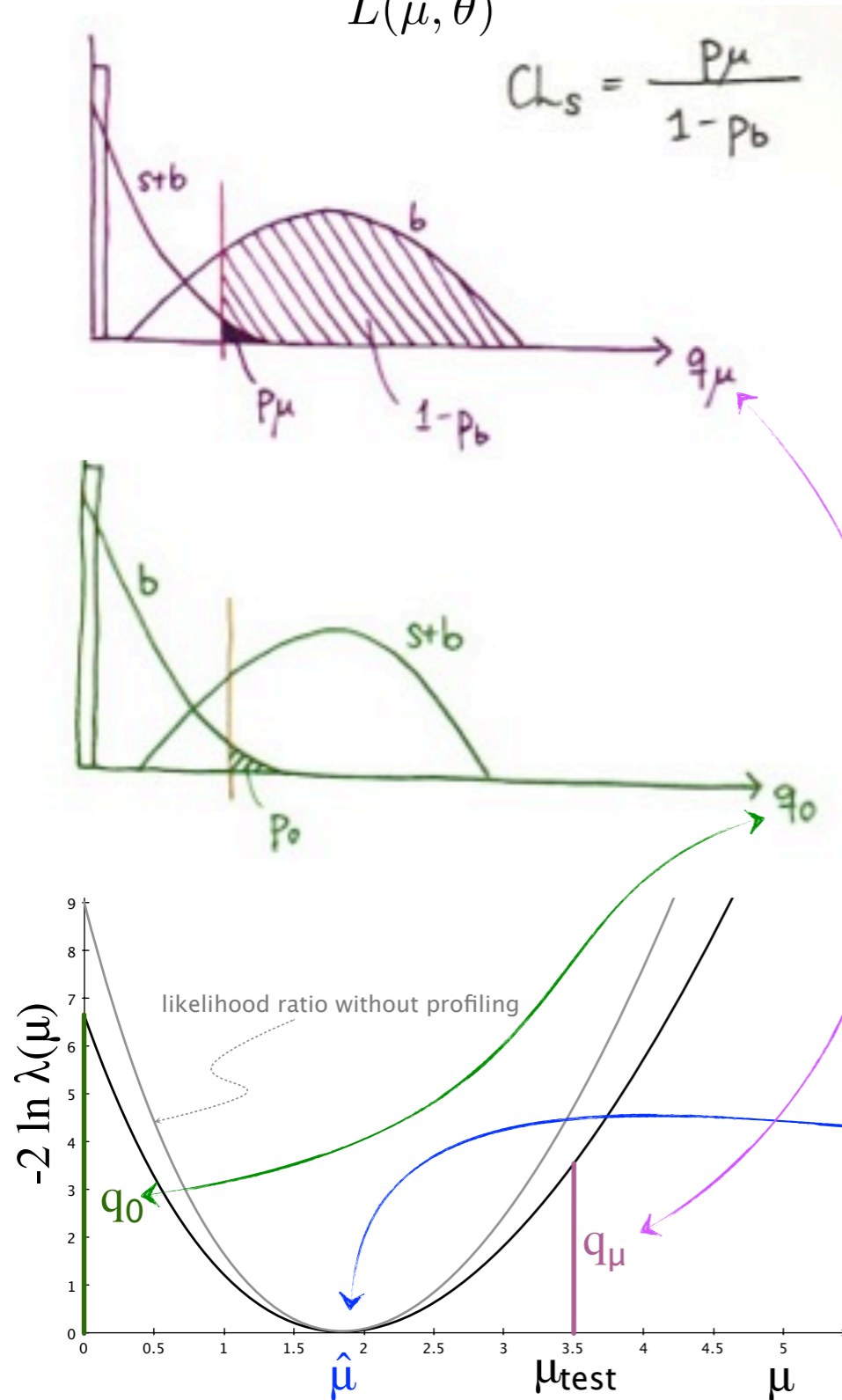
$$CL_s(\mu) = \frac{p_\mu}{1 - p_b}$$

the 95% CL for  $\mu$  is determined when prob to have higher value of test stat for the sign+bkg hypo is 5% of the prob to have higher test stat for the bkg option: so it means that the bkg only hypothesis can be rare i.e. its tail is small, but the sig+bkg is much rarer as its tail is only 5% of the bkg only tail.

$$\lambda(\mu) = \frac{L(\mu, \hat{\theta}(\mu))}{L(\hat{\mu}, \hat{\theta})}$$

$$CL_s = \frac{P_\mu}{1 - P_b}$$

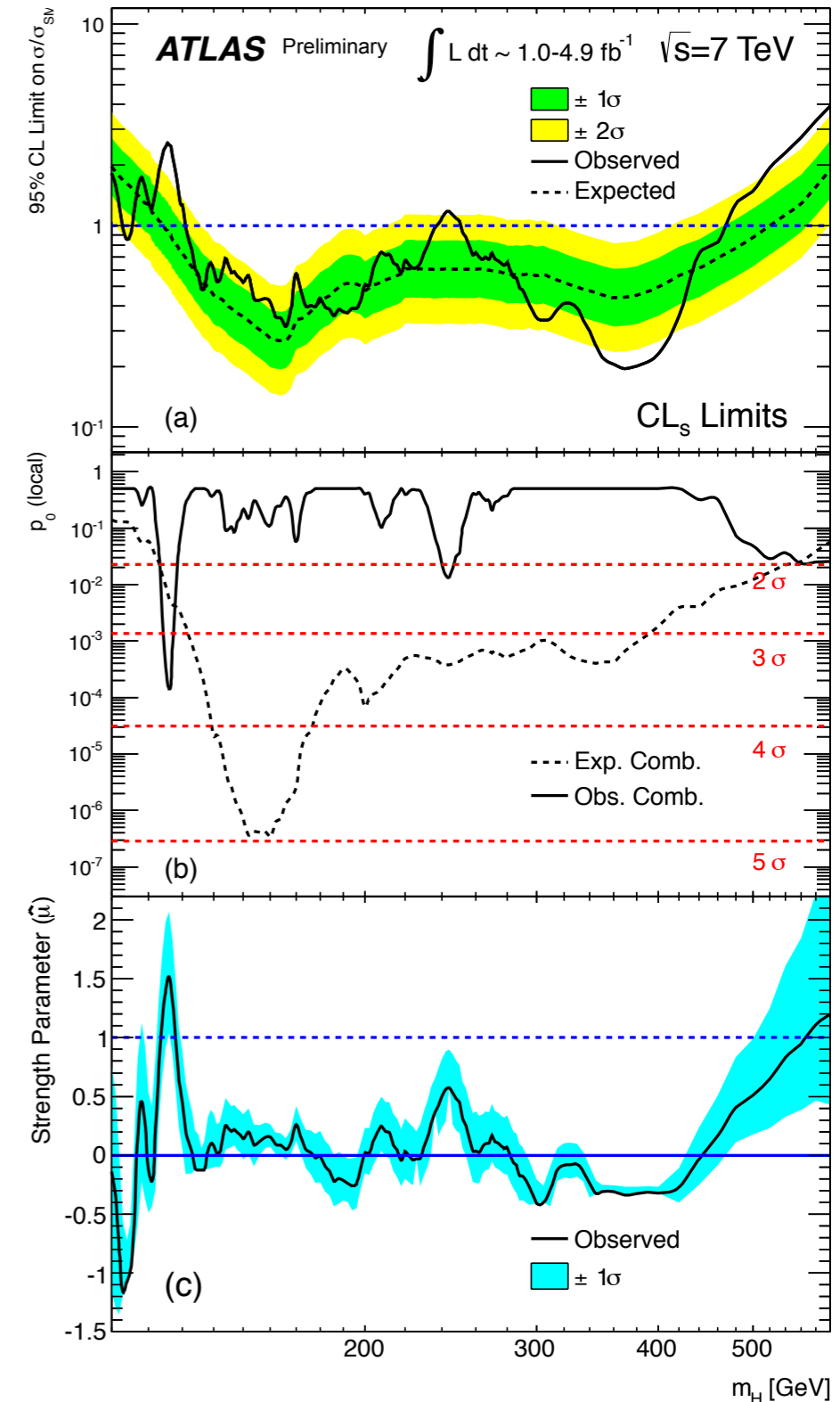
## Follow LHC-HCG Combination Procedures



$CL_s$  to test signal hypothesis

$p_0$  to test background hypothesis

$\hat{\mu}$  to estimate signal strength



# CMS top tagging

- **top tagged jet: anti- $k_T$  ( $R=0.8$ ) jet with**
  - ▶  **$\tau_{32} < 0.65$**
  - ▶  **$105 < m_{\text{jet,SoftDrop}} < 210 \text{ GeV}$** 
    - ❖ soft drop with  $\beta=0$  ,  $z_{\text{cut}}=0.1$
    - ❖  $R=0.8$

# Recognising “highly boosted “ top quarks

# Math Appendix : Mass, $P_T$ and DR

As we know that for any 4-momentum

$$E = m_T \cosh y, \quad p_x, p_y, p_z = m_T \sinh y$$

where  $m_T^2 = m^2 + p_x^2 + p_y^2$  and

The invariant mass  $M$  of the two-particle system

$$y = \frac{1}{2} \ln \left( \frac{E + p_z}{E - p_z} \right) = \ln \left( \frac{E + p_z}{m_T} \right) = \tanh^{-1} \left( \frac{p_z}{E} \right).$$

$$M^2 = m_1^2 + m_2^2 + 2[E_T(1)E_T(2) \cosh \Delta y - \mathbf{p}_T(1) \cdot \mathbf{p}_T(2)],$$

where

$$E_T(i) = \sqrt{|\mathbf{p}_T(i)|^2 + m_i^2},$$

This can be re-written as

$$M^2 = m_1^2 + m_2^2 + 2[E_T(1)E_T(2) \cosh(Dy) - p_T(1)p_T(2) \cos(DPhi)]$$

where  
 $DPhi = Phi(2) - Phi(1)$  is the angle between the two momenta in the transverse plane

Now if 1) the masses of the particles are small w.r.t. their momenta and 2) the splitting is quasi collinear i.e.  $\cos DPhi \sim 1 - (DPhi)^2/2$  and  $\cosh(Dy) \sim 1 + Dy^2/2$ , so  $E_T(i) \sim p_T(i)$

[http://en.wikipedia.org/wiki/Hyperbolic\\_function](http://en.wikipedia.org/wiki/Hyperbolic_function)

$$M^2 \sim 2[p_T(1)p_T(2) (1 + Dy^2/2 - 1 + (DPhi)^2/2)] = p_T(1)p_T(2) (Dy^2/2 + (DPhi)^2) = p_T(1)p_T(2)(DR(1,2))^2$$

So

Labelling  $i$  and  $j$  such that  $p_{tj} < p_{ti}$  and defining  $z = p_{tj}/p_t$

$$(p_t = p_{ti} + p_{tj}),$$

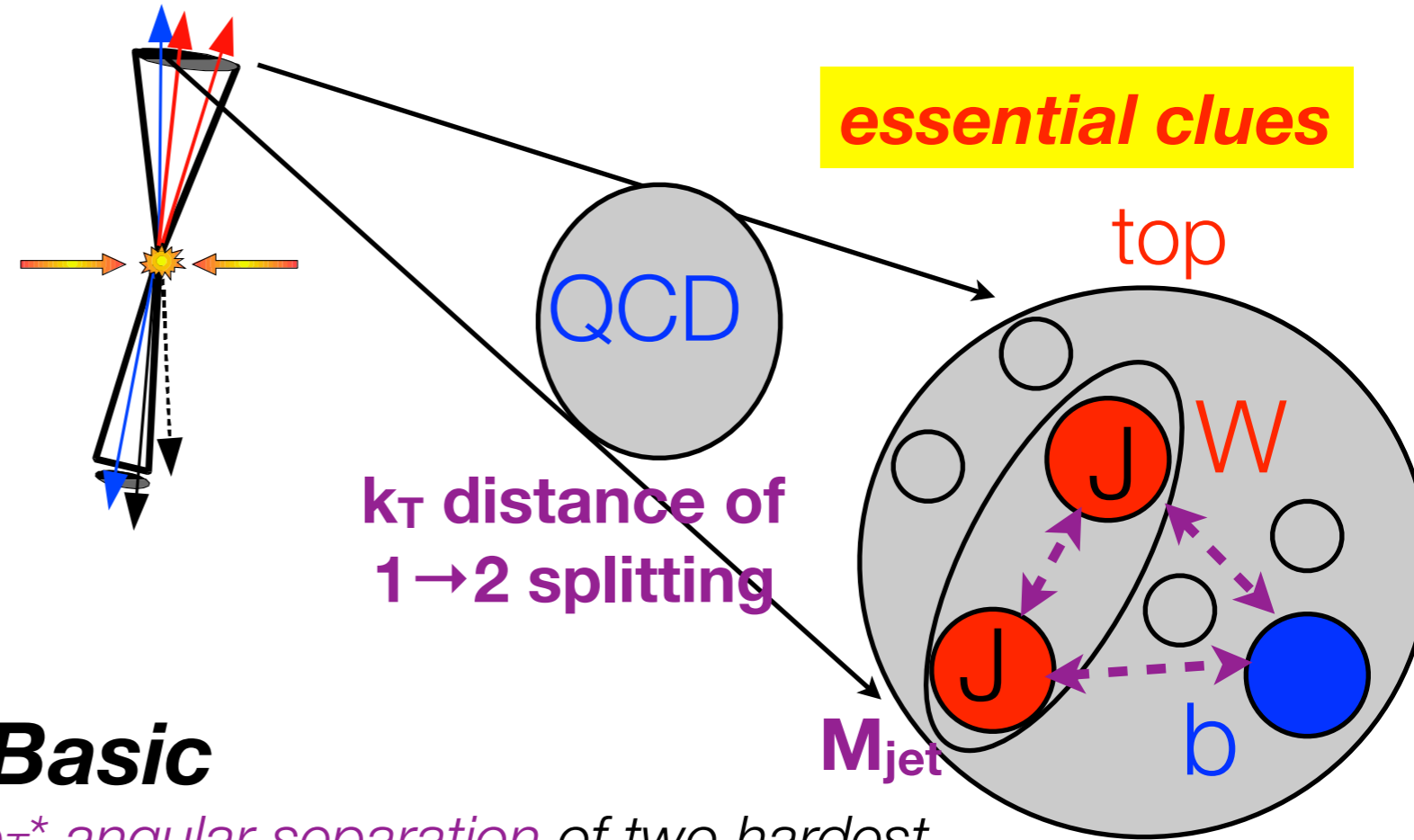
$$m^2 \simeq z(1-z)p_t^2 \Delta R_{ij}^2,$$

$$d_{ij} = z^2 p_t^2 \Delta R_{ij}^2 \simeq \frac{z}{(1-z)} m^2.$$

# How **to tag** a boosted hadronic top quark?

## Look into the jet substructure

(see Jose Juknevich, TOP2013)



## Basic

Use **jet mass** and *product of  $p_T^*$  angular separation* of two hardest jet constituents from jet algorithm

## Declustering

Discard soft coherent radiation (“grooming”) to reveal **boosted objects**:redefine jets

Soft-Drop mass, YSplitter, ATLAS TopTagger, Mass-Drop, CMS Top Tagger, HEPTopTagger, Trimming, Pruning...

## Pattern/Matrix El./Jet shapes

Recognize **energy pattern** in unchanged jet

TopTemplate Tagger, Shower deconstruction,  $N_{subjettiness}$  ratio...



# How **to tag** a boosted hadronic top quark? (II): Examples

*Declustering: redefine jet*

## Soft-drop Mass JHEP05(2014)146

- Revert jet making steps → at each iter. break jet J in 2 subjets  $j_1$  and  $j_2$

- if 
$$\frac{\min(p_{T1}, p_{T2})}{p_{T1} + p_{T2}} > z_{\text{cut}} \left( \frac{\Delta R_{12}}{R_0} \right)^\beta$$

tune  $\beta$   $z_{\text{cut}}$

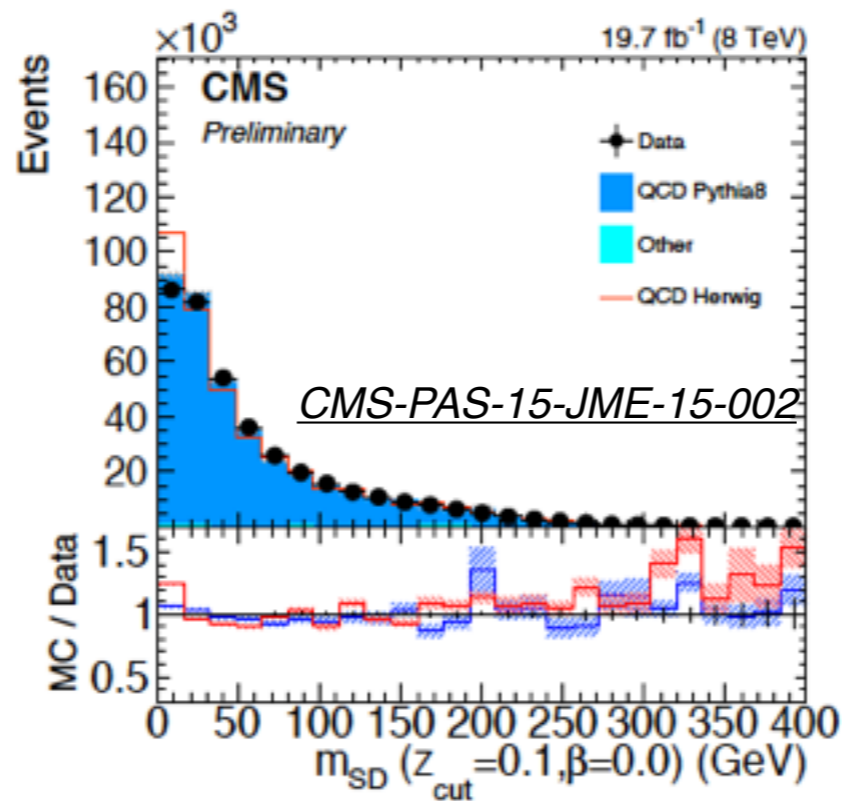
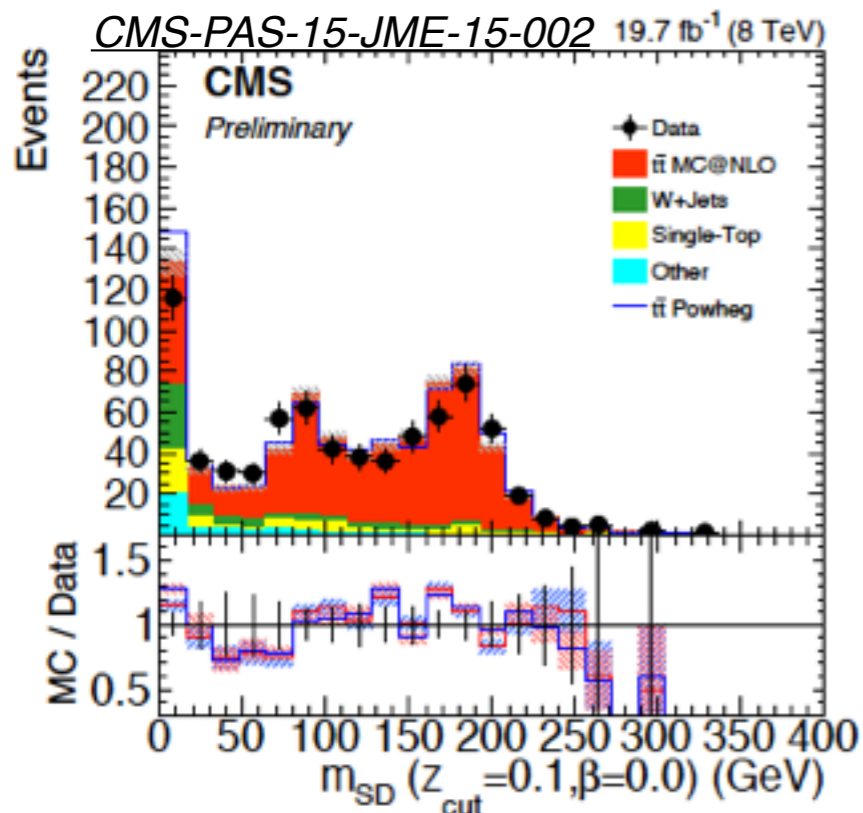
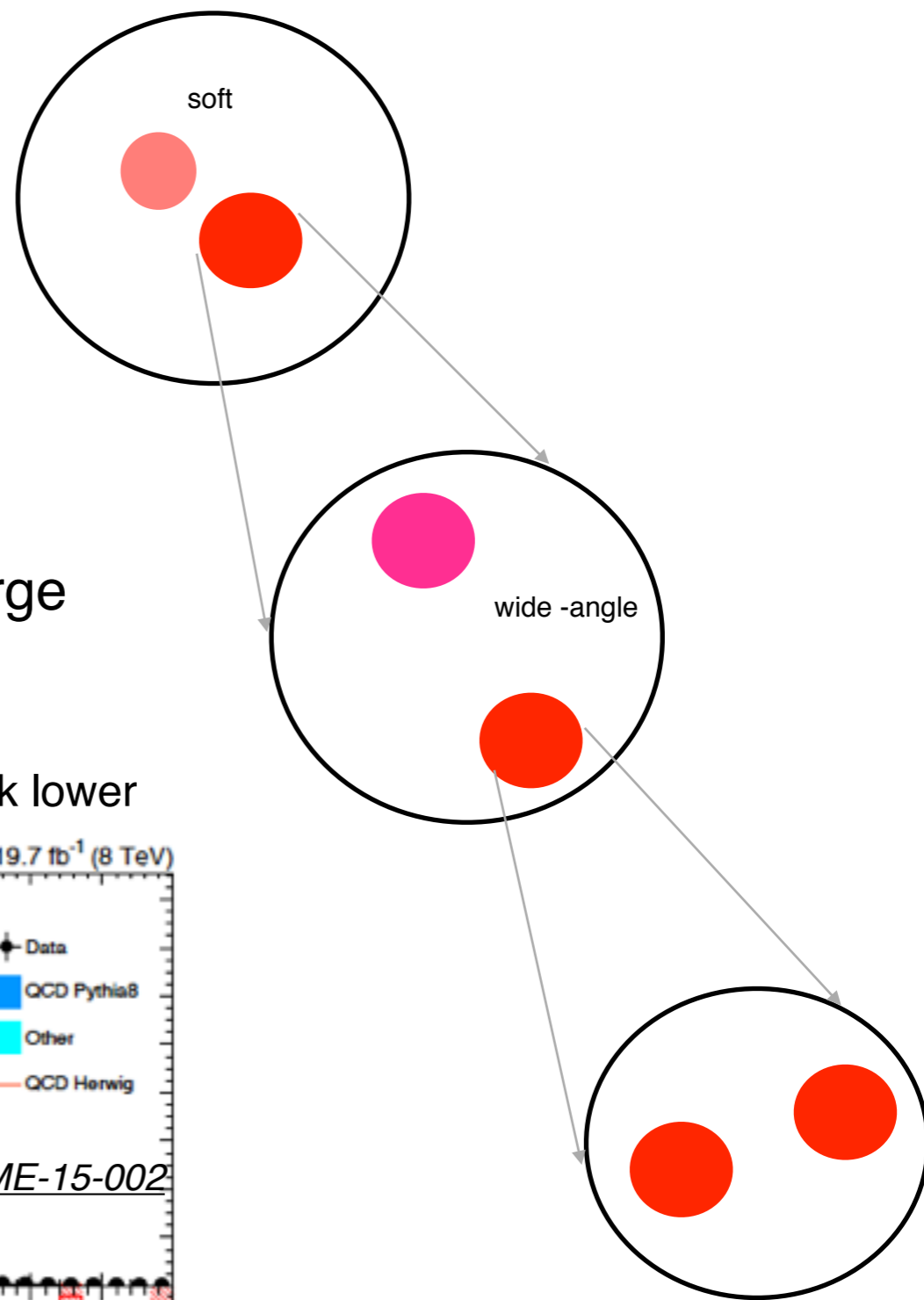
stop: J is the final jet.

Otherwise keep & decluster higher  $p_T$  subjet

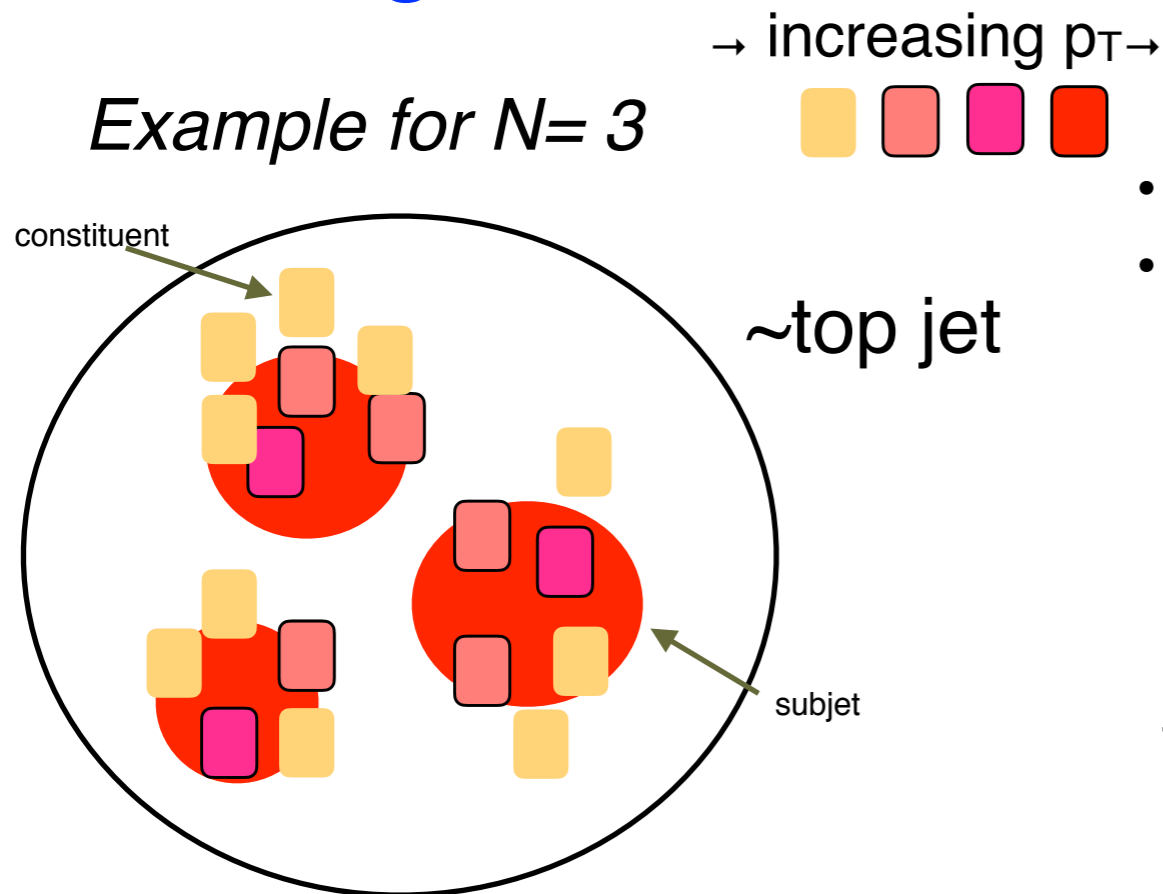
Recursively removes soft (small- $p_T$ ), wide angle (large  $\Delta R$ ) radiation from initial state, pile up, rest of event



Mass of final jet closer to top mass, light quark/gluon jet peak lower



# How to tag a boosted hadronic top quark? (II): Examples



- Make jet with  $N$  subjects ( $k_T$  algo)
- $\Delta R(\text{const}, \text{subject})$ -weighted sum of constituents  $p_T$  → alignment of constituents to subjects hypo

$$\tau_N = \frac{1}{d_0} \sum_k p_{T,k} \min \{ \Delta R_{1,k}, \Delta R_{2,k}, \dots, \Delta R_{N,k} \} . \quad d_0 = \sum_{k_{\text{jet radius}}} p_{T,k} R_0$$

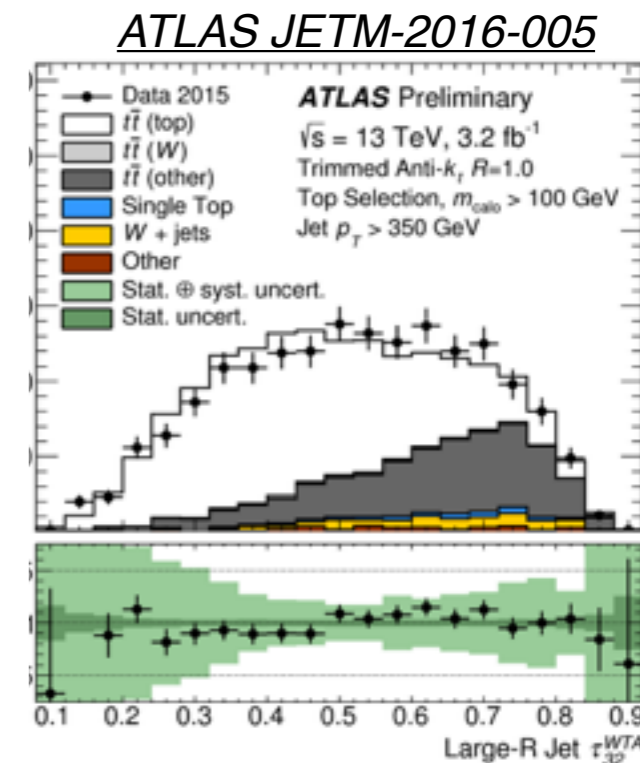
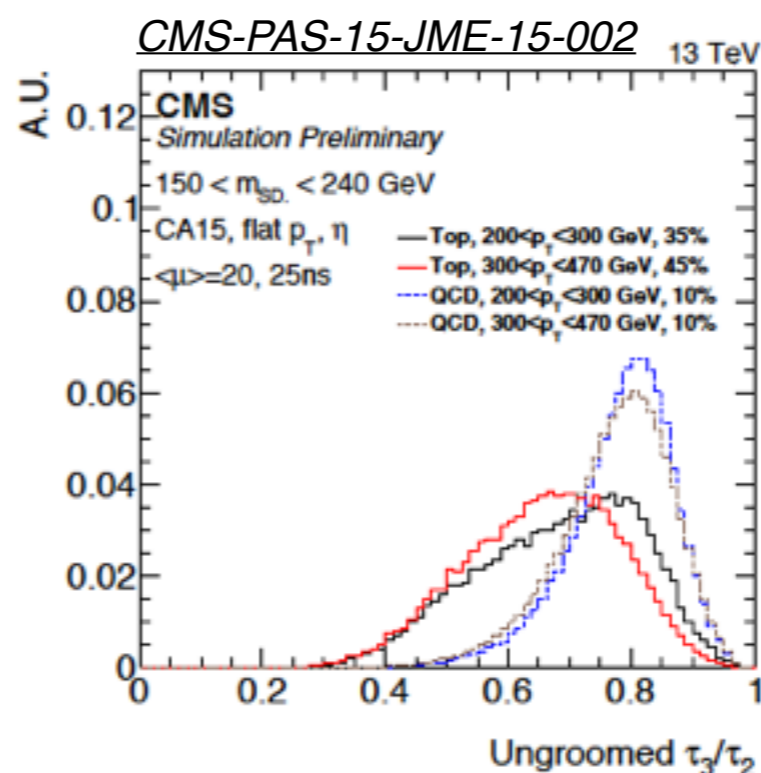
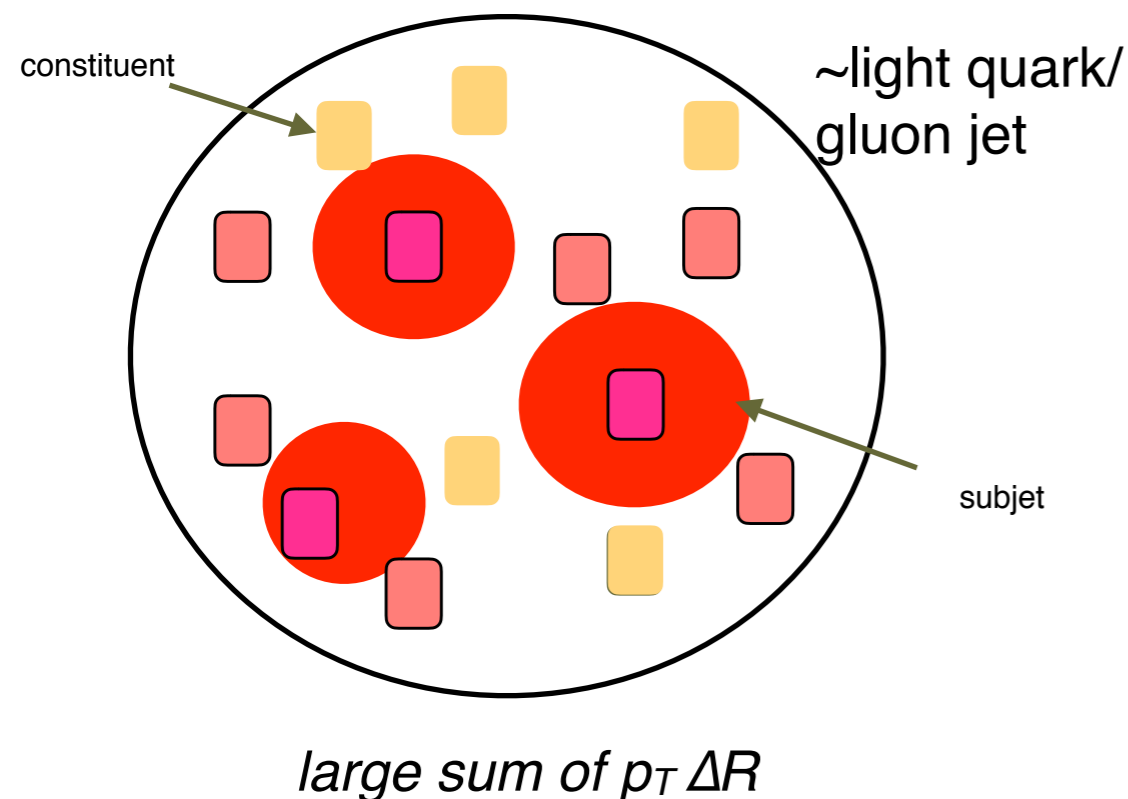
describes how well the jet contains  $N$  ( $\tau_N \rightarrow 0$ ) or fewer ( $\tau_N \rightarrow 1$ ) subjects



- Top-jet : 3 subjects ;
- light quark/gluon- jet  $< 3$

Use  $\tau_{32} = \tau_3 / \tau_2$

The lower, the more top-like (3-prong)



# How **to tag** a boosted hadronic top quark? (II): Performance

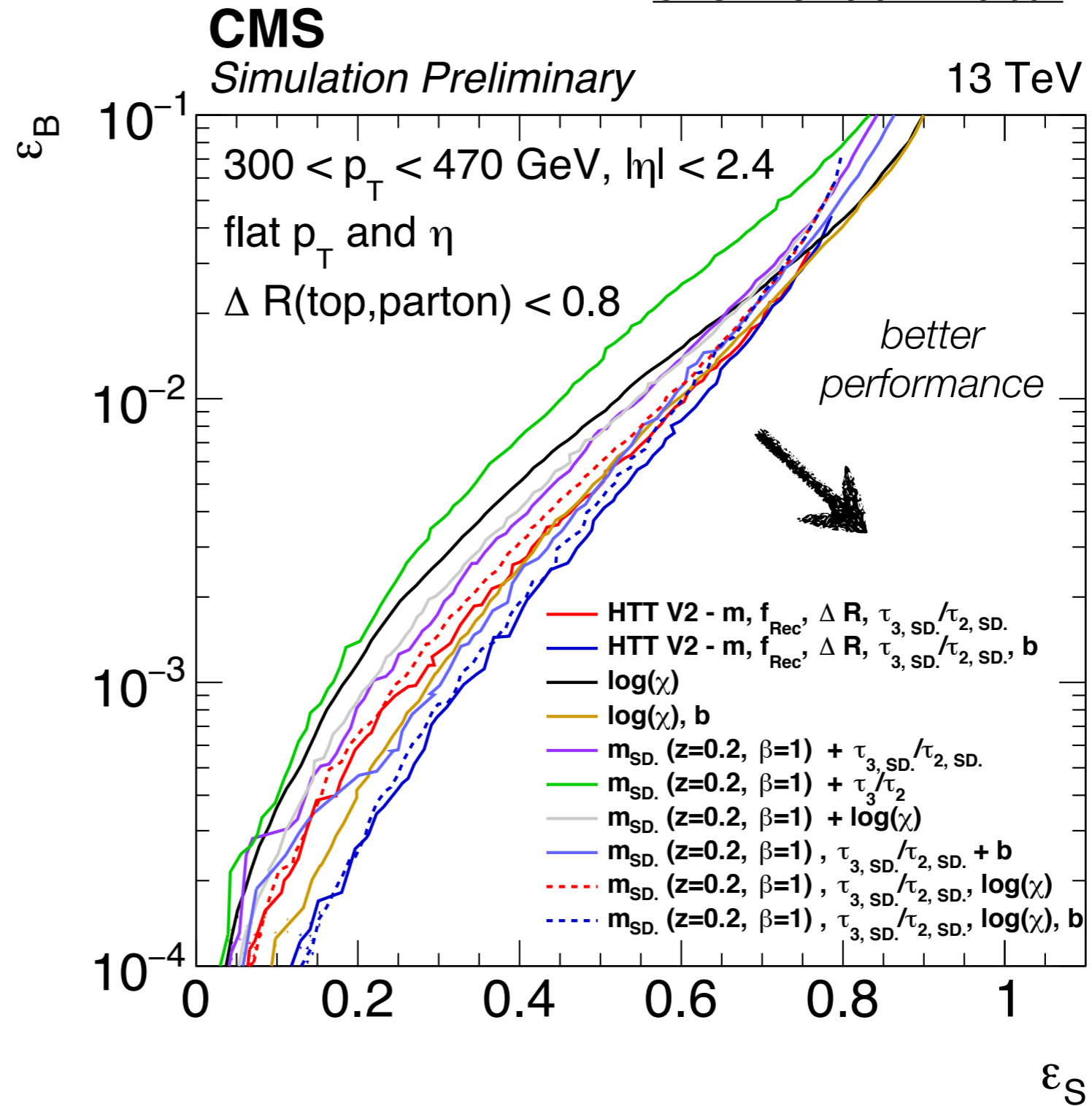
CMS-PAS-15-JME-15-002

**efficiency to tag top jet  $\epsilon_S$  VS efficiency to mistake a light quark/gluon jet  $\epsilon_B$**   
 (or rejection against light-gluon jet)

(Receiver Operating Characteristic curve)

**sensitivity to energy from superposed collisions (pile-up)**

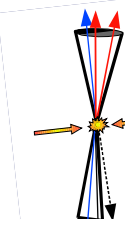
efficiency to **select the top final state** VS bkg



Combination of different tagging schemes improves performance

# Search for excess in “boosted” $t\bar{t}$ production vs $M_{t\bar{t}}$

*JHEP 07 (2017) 001*



▶ *single lepton trigger*

▶ **exactly 1 high  $p_T$  ( $>50$  GeV) central lepton ( $e, \mu$ ) with 2d isolation vs fake leptons:  $\Delta R(\text{lep}, j) > 0.4$  or  $p_{T, \text{rel}}(\text{lep}, j) > 20$  GeV**

▶ *~90% to 94% efficient*

$j = \Delta R$ -closest-to-lep anti- $k_T$  ( $R=0.4$ ) jet

▶ **very high  $E_T^{\text{miss}}$  ( $p_T^{\text{lep}} + E_T^{\text{miss}} > 120$  (150) GeV in  $e(\mu)$  channel**

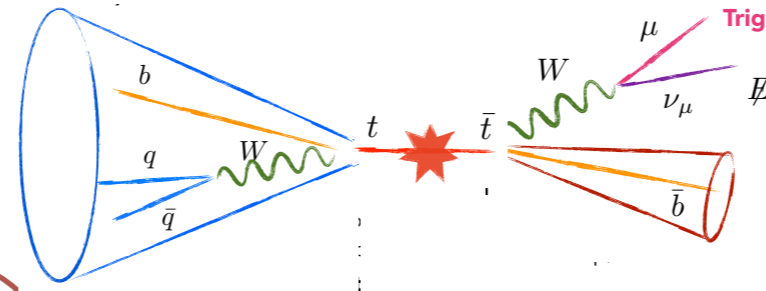
▶  **$\geq 2$  small- $R$  anti- $k_T$  ( $R=0.4$ ) jet with leading  $p_T > 250$  (150 GeV) in  $e(\mu)$  chan,  $p_T > 30$  GeV for all jets**

▶ **only zero or 1 large- $R$  top tagged jet: anti- $k_T$  ( $R=0.8$ ) with  $\tau_{32} < 0.69$  and  $110 < m_{\text{jet}, \text{SoftDrop}} < 210$  GeV (50%  $\epsilon_S$ , 3%  $\epsilon_B$ ) (orth. to fully had)**

▶ **b-tag on anti- $k_T$  ( $R=0.4$ ) (65% eff)**

**I+jets  $qq\ell vbb$**

*ingredients for leptonic W*

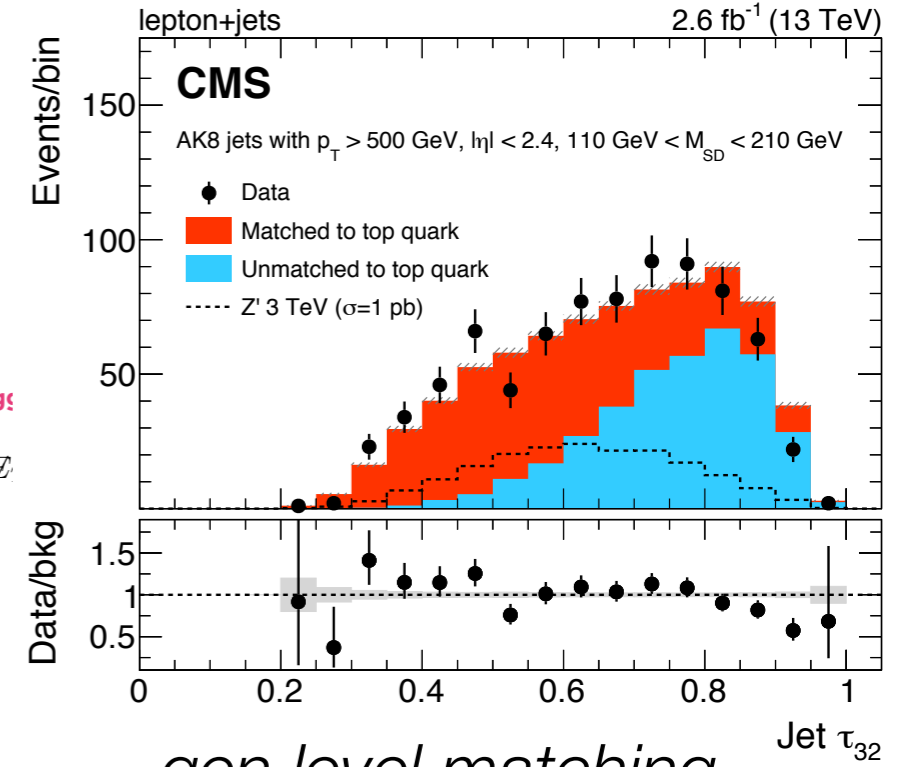


(figure by A Ovcharova (UCSB))

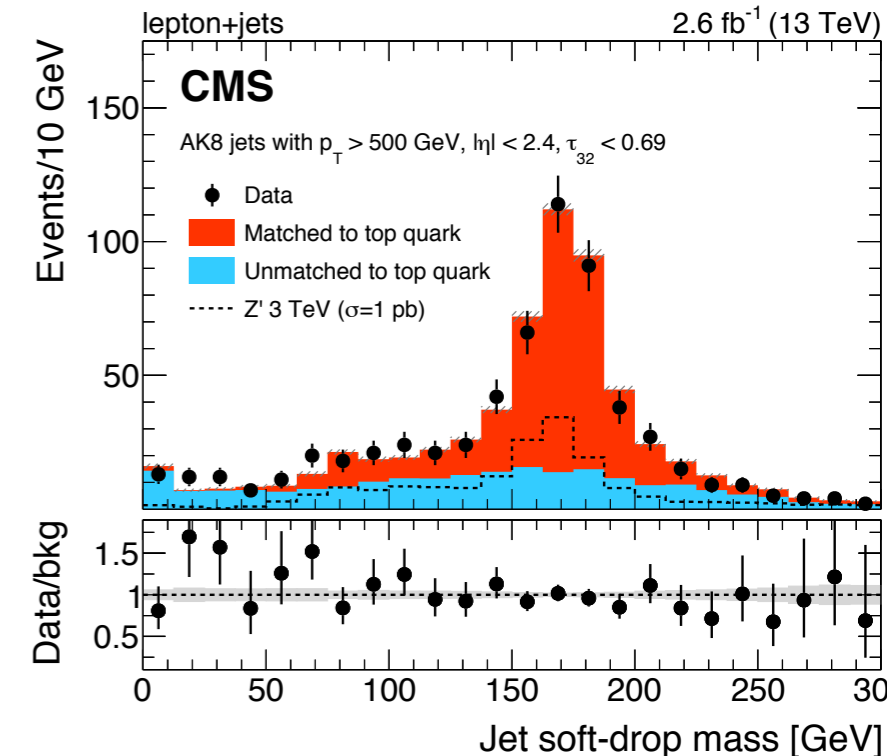
*ingredients for hadronic & leptonic top*

**6 selection regions**

	1 top tag	0 t-tag & 1 b-tag	0 t-tag & 0 b-tag
2 flavours x events	394	10447	8971
tt purity(%)	85	82	24



*gen level matching*



EFT

# Effective field theory in a nutshell

*We are all Wilsonians now ! (JPreskill, Caltech)*

- Current **absence of “light “ new states** in SM → possible **new physics at higher scales/masses** than observed
- Effective Field Theory: ultraviolet divergences = manifestation of new phys. Renormalisation= Lagrangian is “new phys scale”-dependent  
← absorb effects of fluctuations/momenta between observed and high new phys. scale in few parameters

*(F Maltoni, LHCTopWG open meeting Nov 2016)*

- SM measurements “searches for deviations predictions of SM in dim=4 .

$$\mathcal{L}_{SM}^{(6)} = \mathcal{L}_{SM}^{(4)} + \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i + \dots$$

*dimensionless Wilson coefficients*

- *Parametric new degrees of freedom in terms of old*

the BSM ambitions of the LHC Higgs/Top/SM physics programmes can be recast in a simple and powerful way in terms of one statement:

**DETERMINATION OF THE COUPLINGS OF THE SM  $\mathcal{L}$  UP TO DIM=6**

# Initial attempt: Top Fitter

Fit **~40 measurements** from LHC & Tevatron ( five 1-dim differential xsec for tt and single top too ) to **predictions** to derive 12 couplings

*Buckley et al,*  
*TopFitterColl,*  
*arxiv:1612.02294*

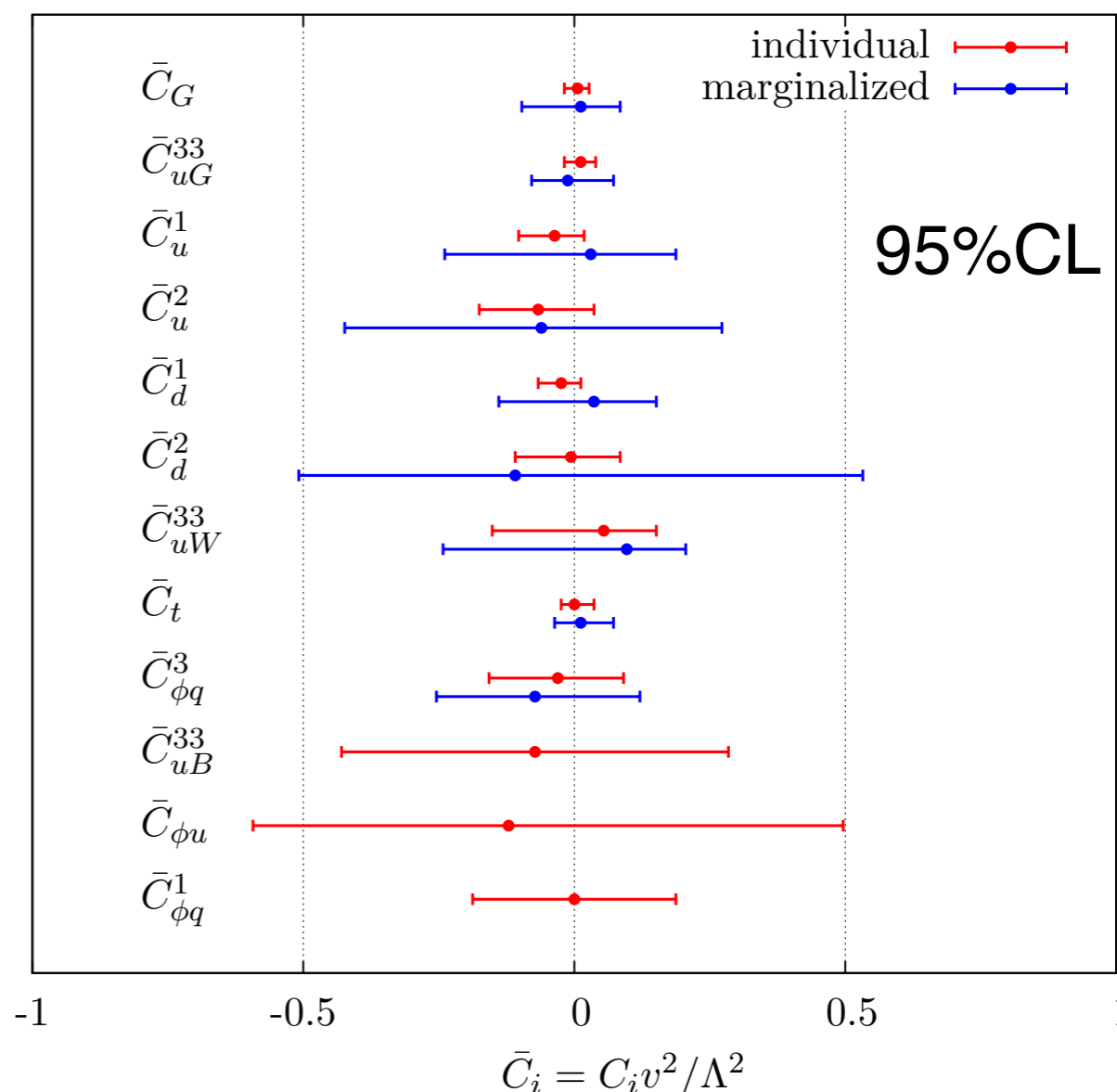
## Predictions as polynomials

$$f_b(\{C_i\}) = \alpha_0^b + \sum_i \beta_i^b C_i + \sum_{i \leq j} \gamma_{i,j}^b C_i C_j + \dots$$

$$\chi^2(\mathbf{C}) = \sum_{\mathcal{O}} \sum_{i,j} \frac{(f_i(\mathbf{C}) - E_i) \rho_{i,j} (f_j(\mathbf{C}) - E_j)}{\sigma_i \sigma_j}$$

Annotations:  
 -  $E_i$ : experimental value  
 -  $\rho_{i,j}$ : correlation  
 -  $\sigma_i \sigma_j$ : standard dev

Phys.Lett.B 763 (2016) 9



**Including covariances where provided by experiments otherwise  $\rho_{i,j} = \delta_{i,j}$**

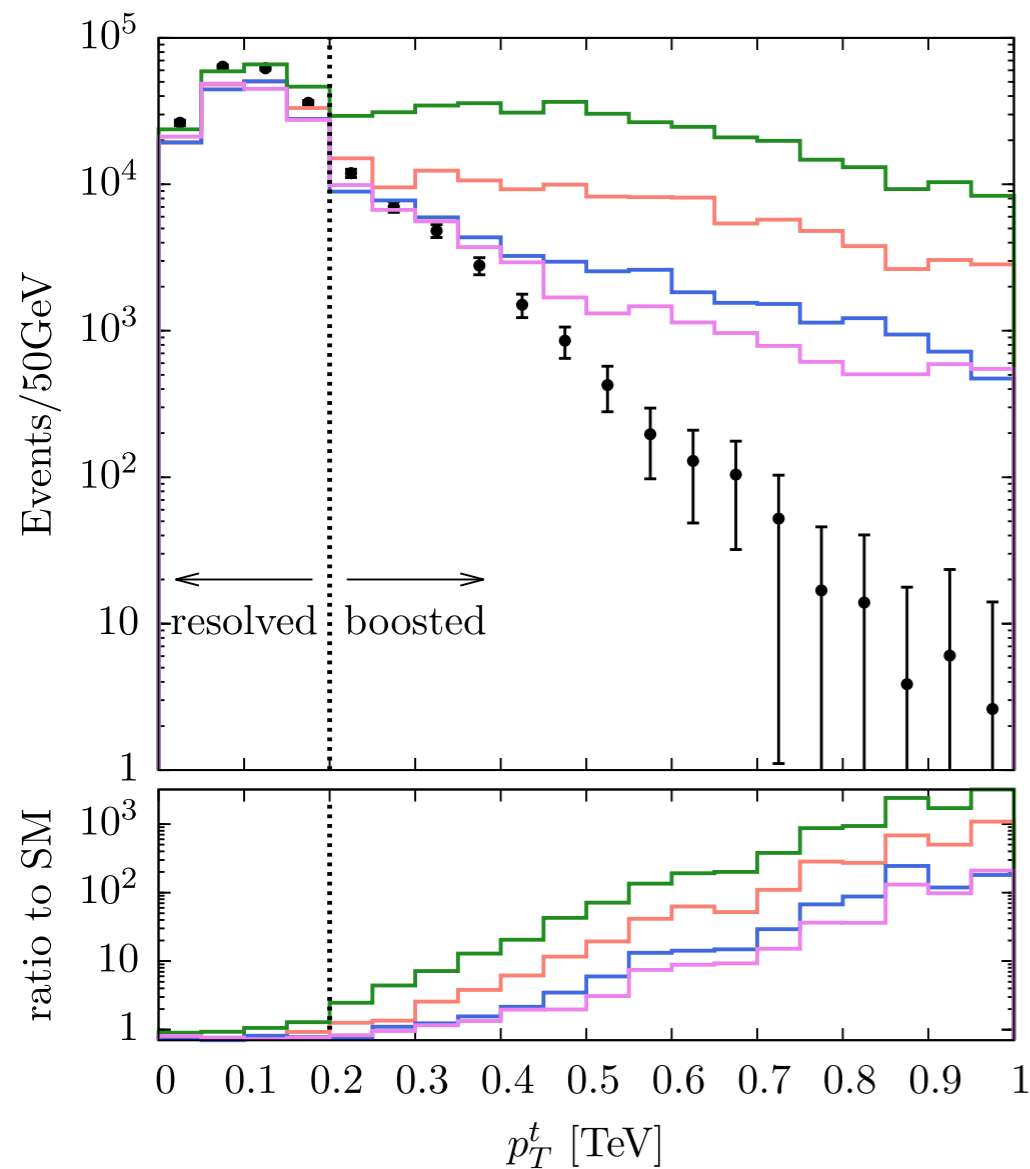
**fit each operator one at the time**  
**fit globally and marginalise over all other parameters**

**All operators are consistent with zero at 95% CL**

# Initial attempt: Top Fitter

*Buckley et al, TopFitterColl, arxiv:1612.02294*

Present bound on coefficients are weak : resulting Scale close to high energy range of  $d\sigma_{tt}/dX$  where EFT breaks down



SM @ 30 fb<sup>-1</sup> ●  
 $C_u^1$  —  
 $C_u^2$  —  
 $C_d^1$  —  
 $C_d^2$  —

20% sys improvement & 30/fb

20% sys improvement & 300/fb

20% sys improvement & 3000/fb

10% sys improvement & 30/fb

10% sys improvement & 300/fb

10% sys improvement & 3000/fb

Isolate region that are most sensitive to tails : fit resolved and boosted

**More data give modest gain in boosted.**

## Fractional improvement on 95% CL confidence interval

