Recent progress and open questions in ab initio simulations of nuclei


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## Ab initio vs. effective approach

## Ab initio approach

| A-body Hamiltonian |
| :---: |
| $H=T+V^{2 \mathrm{~N}}+V^{3 \mathrm{~N}}+\ldots+V^{A \mathrm{~N}}$ |$\quad H\left|\Psi_{k}^{A}\right\rangle=E_{k}^{A}\left|\Psi_{k}^{A}\right\rangle \quad \longrightarrow \quad A$-body wave-function

$\rightarrow$ Solve many-body Schrödinger equation in a controlled, systematically improvable way

## Effective approach

Ab initio valence space


Simplify $\left|\Psi_{k}^{A}\right\rangle \rightarrow \quad$ Energy density
functional
Accuracy

Reach across the mass table Predictive power/error estimate


## Evolution of ab initio nuclear chart

○ "Exact" approaches

- Since 1980's
- Monte Carlo, CI, ...
- Factorial/ exponential scaling

© Approximate approaches for closed-shell nuclei
- Since 2000's
- SCGF, CC, IMSRG
- Polynomial scaling


## © Ab initio shell model

- Since 2014
- Effective interaction via CC/IMSRG
- Mixed scaling


## Self-consistent Green's function approach

$\odot$ Solution of the $\boldsymbol{A}$-body Schrödinger equation $H\left|\Psi_{k}^{A}\right\rangle=E_{k}^{A}\left|\Psi_{k}^{A}\right\rangle$ achieved by

1) Rewriting it in terms of $\mathbf{1 -}, \mathbf{2 -}, \ldots . . A$-body objects $G_{1}=G, G_{2}, \ldots G_{A}$ (Green's functions)
2) Expanding these objects in perturbation (in practise $\mathbf{G} \rightarrow$ one-body observables, etc..)

- Self-consistent schemes resum (infinite) subsets of perturbation-theory contributions
$\odot$ Self-energy expansion

$\odot$ Access a variety of quantities
$\circ$ One-body GF $\rightarrow$ Ground-state properties of even-even $A+$ spectra of odd-even neighbours
- Two-body GF $\rightarrow$ Excited spectrum of even-even $A$
- Self-energy $\rightarrow$ Optical potential for nucleon-nucleus scattering


## Chiral effective field theory \& nuclear interactions

$\odot$ Chiral EFT aims to provide a systematic framework to construct $A \mathrm{~N}$ interactions ( $A=2,3, \ldots$ )
$\bigcirc$ Main features:

- High-energy physics unresolved $\rightarrow$ soft potentials $\rightarrow$ improved many-body convergence
- Many-body forces and currents consistently derived
- A theoretical error can be, in principle, assigned to each order in the expansion

$\Rightarrow$ Ideally: apply to the many-nucleon system (and propagate the theoretical error)


## Chiral effective field theory \& nuclear interactions

$\odot$ Renormalisability $\leftrightarrow$ independence of UV physics
© Most commonly used power counting scheme (Weinberg PC) not renormalisable
© Two alternatives:

## Fix-cutoff approach

- Phenomenological success
- A posteriori error estimate [e.g. Epelbaum et al. 2015]


## Renormalisable approach

- Work in progress [van Kolck, Pavon Valderrama, Long, ...]
- Non-trivial impact on/from many-body approximations used [Drissi et al. in preparation]


## First "standard" interaction [(EM) N3LO]


$\checkmark$ Successful benchmarks



$X$ Radii underestimated

## Testing interactions, pt. 1



## Testing interactions, pt. 2



## Regulator artefacts

$\bigcirc$ Regularisation scheme is a major source of variation among currently employed Hamiltonians

- Effects on 2 N phase space
- Effects on 3N phase space



[Dyhdalo, Furnstahl, Hebeler, Tews 2016]


## Emergence of magic numbers "ab initio"




## Doubly open-shell nuclei

© Currently, description of doubly open-shell nuclei quantitatively worsens with deformation

[Somà et al. in preparation]

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$\Rightarrow$ Correlation with deformation parameter $\beta \quad$ [Hilaire \& Girod 2007]


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## Charge radii



## The case of ${ }^{34} \mathrm{Si}$

© Unconventional depletion ("bubble") in the centre of $\rho_{\text {ch }}$ conjectured for certain nuclei
© Purely quantum mechanical effect

- $\ell=0$ orbitals display radial distribution peaked at $r=0$
$\circ \ell \neq 0$ orbitals are instead suppressed at small $r$
- Vacancy of $s$ states $(\ell=0)$ embedded in larger- $\ell$ orbitals might cause central depletion
© Conjectured associated effect on spin-orbit splitting
- Non-zero derivative at the interior
- Spin-orbit potential of "non-natural" sign
- Reduction of (energy) splitting of low- $\ell$ spin-orbit partners


๑ Bubbles predicted for hyper-heavy nuclei
[Dechargé et al. 2003]
$\odot$ In light/medium-mass nuclei the most promising candidate is ${ }^{34} \mathbf{S i}$

[Todd-Rutel et al. 2004, Khan et al. 2008, ...]

## The case of ${ }^{34} \mathrm{Si}$

$\odot$ Good reproduction of g.s. properties

| $E[\mathrm{MeV}]$ | ADC(1) | ADC(2) | ADC(3) | Experiment |
| :---: | :---: | :---: | :---: | :---: |
| ${ }^{34} \mathrm{Si}$ | -84.481 | -274.626 | -282.938 | -283.427 |
| ${ }^{36} \mathrm{~S}$ | -90.007 | -296.060 | -305.767 | -308.714 |


| $\left\langle r_{\mathrm{ch}}^{2}\right\rangle^{1 / 2}$ | $\mathrm{ADC}(1)$ | $\mathrm{ADC}(2)$ | $\mathrm{ADC}(3)$ | Experiment |
| :---: | :---: | :---: | :---: | :---: |
| ${ }^{34} \mathrm{Si}$ | 3.270 | 3.189 | 3.187 | - |
| ${ }^{36} \mathrm{~S}$ | 3.395 | 3.291 | 3.285 | $3.2985 \pm 0.0024$ |

## $\odot$ Mild central depletion predicted


h Charge density computed via folding with the finite charge of the proton
$\Rightarrow$ Folding smears out central depletion
b) Excellent agreement with experimental charge distribution of ${ }^{36} \mathrm{~S}$

## The case of ${ }^{34} \mathrm{Si}$

$\odot$ Addition and removal spectra compared to transfer and knock-out reactions

## One-neutron addition

[Thorn et al. 1984]
Exp. data: [Eckle et al. 1989]
[Burgunder et al. 2014]

## One-proton knock-out

[Khan et al. 1985]
Exp. data: [Mutschler et al. 2016 (PRC)]
[Mutschler et al. 2016 (Nature Phys.)]


## K spectra

$\Rightarrow K$ spectra show interesting g.s. spin inversion and re-inversion


Laser spectroscopy COLLAPS @ ISOLDE



## Doubly open-shell nuclei

๑ Approximate/truncated methods capture correlations via an expansion in ph excitations $\odot$ Open-shell nuclei are (near-)degenerate with respect to ph excitations

$\odot$ Solution: multi-determinantal or symmetry-breaking reference state

- Symmetry-breaking solution allows to lift the degeneracy


Developed and implemented

Quadrupole correlations
Deformation
$\uparrow$
Breaking of $\mathrm{SU}(2)$

Doubly open-shells

## Tensor decomposition of many-body formalism

$\Rightarrow$ Many-body methods require the handling (computation \& storage) of large tensors
c) Matrix elements of 3-body interaction represent current memory bottleneck

Use tensor decomposition techniques
$\bigcirc$ Two-body forces can be factorised as $v_{i j k l}=\sum_{a} \lambda_{a} g_{i k}^{a} g_{j l}^{a} \quad(\rightarrow$ Singular Value Decomposition)


Gain \#1: size $(\rightarrow$ storage and memory needs $)$

$$
\begin{aligned}
& \sum_{k l} v_{i j k l}=\sum_{a} \lambda_{a} \sum_{k} g_{i k}^{a} \sum_{l} g_{j l}^{a} \\
& \mathrm{~N}^{2} \quad \mathrm{~m} \quad(\mathrm{~N}+\mathrm{N})=\mathrm{mN} \\
& \text { Gain \#2: CPU speed-up }
\end{aligned}
$$

HF test: $\rightarrow 0.003 \%$ error and factor 10 speed-up
$\odot$ Higher-order tensors: exploit techniques from applied maths (e.g. tensor hypercontraction)

## Conclusions

© Not so good news

- Renormalisable approach $\rightarrow$ still a long way to go?
- Fix-cutoff approach $\rightarrow$ few issues hinder full phenomenological success
$\bigcirc$ Good news
- Many-body methods mature for applications in medium-mass nuclei
- Promising ideas for extension to heavy nuclei
© Extension of ab initio simulations to heavy nuclei
- Extension to doubly open shell requires new formal developments
- Computational challenges ahead: work in progress and more smart ideas needed

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