
Supersymmetric QCD Effects on Neutralino Dark Matter Annihilation

Björn Herrmann

Institut für Theoretische Physik und Astrophysik
Universität Würzburg

Work realized in collaboration with Karol Kovařík and Michael Klasen
[accepted by *Phys. Rev. D* (2009); *Phys. Rev. D*79:061701 (2009)]

GDR Terascale
Heidelberg, October 15, 2009

Outline

1. Motivation
2. Neutralino annihilation into heavy quarks
3. Supersymmetric QCD corrections
4. Numerical results
5. Conclusion

Supersymmetric Dark Matter

Supersymmetry provides interesting candidates for “WIMPs”

- Minimal Supersymmetric Standard Model (MSSM) with R-parity conservation
- Lightest supersymmetric particle (LSP) stable
- Good candidates: Neutralino or Gravitino (depending on SUSY-breaking mechanism)

Supersymmetric Dark Matter

Supersymmetry provides interesting candidates for “WIMPs”

- Minimal Supersymmetric Standard Model (MSSM) with R-parity conservation
- Lightest supersymmetric particle (LSP) stable
- Good candidates: Neutralino or Gravitino (depending on SUSY-breaking mechanism)

In “mSUGRA-like” models based on gravity-mediation the LSP is a neutralino

$$\tilde{\chi}_1^0 = \mathcal{N}_{11}\tilde{B}^0 + \mathcal{N}_{12}\tilde{W}^0 + \mathcal{N}_{13}\tilde{H}_1^0 + \mathcal{N}_{14}\tilde{H}_2^0$$

Universal parameters at the grand unification scale: $m_0, m_{1/2}, A_0, \tan\beta, \text{sgn}(\mu)$

Supersymmetric Dark Matter

Supersymmetry provides interesting candidates for “WIMPs”

- Minimal Supersymmetric Standard Model (MSSM) with R-parity conservation
- Lightest supersymmetric particle (LSP) stable
- Good candidates: Neutralino or Gravitino (depending on SUSY-breaking mechanism)

In “mSUGRA-like” models based on gravity-mediation the LSP is a neutralino

$$\tilde{\chi}_1^0 = \mathcal{N}_{11}\tilde{B}^0 + \mathcal{N}_{12}\tilde{W}^0 + \mathcal{N}_{13}\tilde{H}_1^0 + \mathcal{N}_{14}\tilde{H}_2^0$$

Universal parameters at the grand unification scale: $m_0, m_{1/2}, A_0, \tan\beta, \text{sgn}(\mu)$

Cosmology allows to constrain the (large) MSSM parameter space

- Additional information w.r.t. collider and precision data
- Identify (dis)favoured regions of parameter space
- Scenario should include viable candidate for cold dark matter
- Relic density of cold dark matter required to agree with WMAP+SN+BAO data

$$0.1097 < \Omega_{\text{CDM}} h^2 < 0.1165 \quad (\text{at } 2\sigma)$$

[Hinshaw et al. (WMAP) 2008]

Relic Density Calculation

Number density of relic particle governed by the Boltzmann equation

$$\frac{dn}{dt} = -3Hn - \langle \sigma_{\text{ann}} v \rangle (n^2 - n_{\text{eq}}^2)$$

$$\Omega_{\text{CDM}} h^2 \propto n_0 \propto \frac{1}{\langle \sigma_{\text{ann}} v \rangle}$$

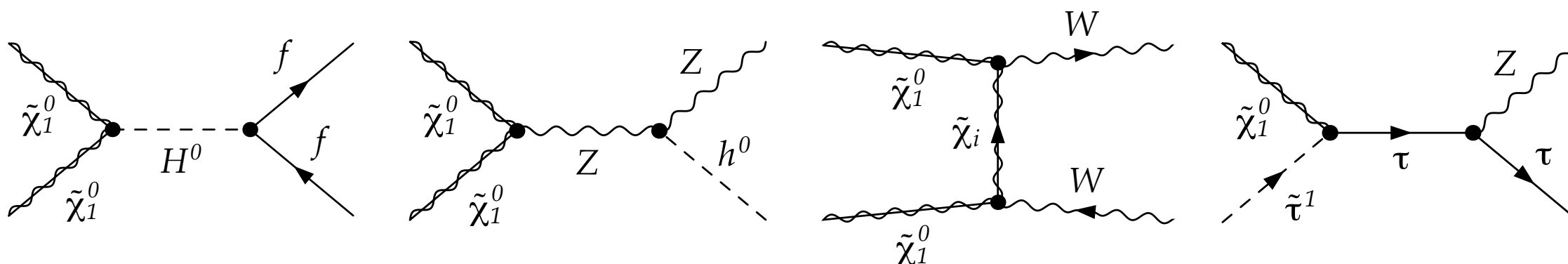
Relic Density Calculation

Number density of relic particle governed by the Boltzmann equation

$$\frac{dn}{dt} = -3Hn - \langle \sigma_{\text{ann}} v \rangle (n^2 - n_{\text{eq}}^2)$$

$$\Omega_{\text{CDM}} h^2 \propto n_0 \propto \frac{1}{\langle \sigma_{\text{ann}} v \rangle}$$

Cross section σ_{ann} includes all annihilation and coannihilation processes



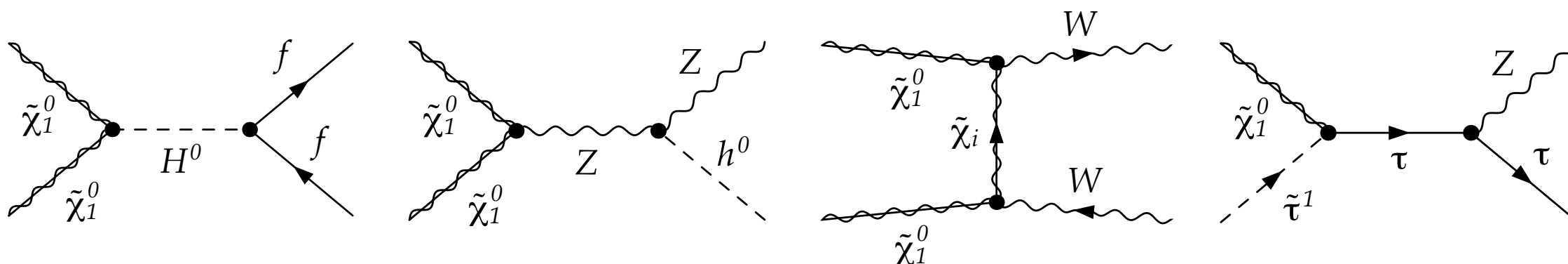
Relic Density Calculation

Number density of relic particle governed by the Boltzmann equation

$$\frac{dn}{dt} = -3Hn - \langle \sigma_{\text{ann}} v \rangle (n^2 - n_{\text{eq}}^2)$$

$$\Omega_{\text{CDM}} h^2 \propto n_0 \propto \frac{1}{\langle \sigma_{\text{ann}} v \rangle}$$

Cross section σ_{ann} includes all annihilation and coannihilation processes



Thermal average involves velocity distribution of the non-relativistic relic particle

$$\langle \sigma_{\text{ann}} v \rangle = \int dv \tilde{f}(v) \sigma_{\text{ann}} v = \int ds f(s) \sigma_{\text{ann}}(s)$$

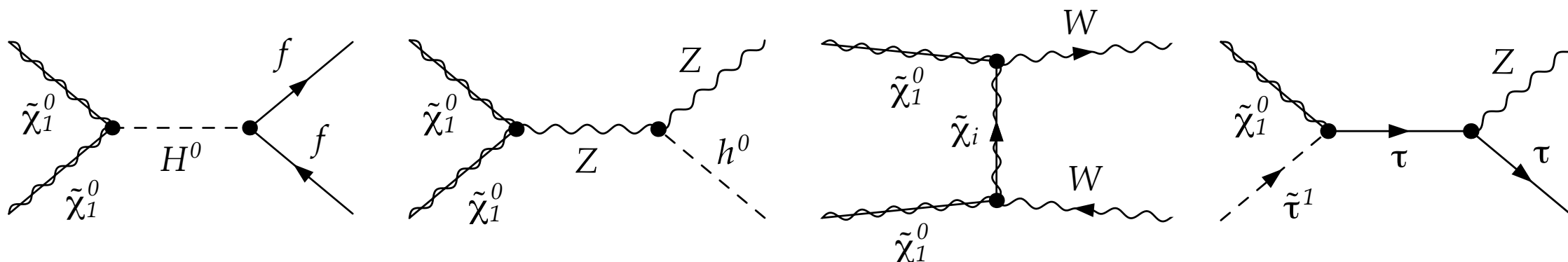
Relic Density Calculation

Number density of relic particle governed by the Boltzmann equation

$$\frac{dn}{dt} = -3Hn - \langle \sigma_{\text{ann}} v \rangle (n^2 - n_{\text{eq}}^2)$$

$$\Omega_{\text{CDM}} h^2 \propto n_0 \propto \frac{1}{\langle \sigma_{\text{ann}} v \rangle}$$

Cross section σ_{ann} includes all annihilation and coannihilation processes



Thermal average involves velocity distribution of the non-relativistic relic particle

$$\langle \sigma_{\text{ann}} v \rangle = \int dv \tilde{f}(v) \sigma_{\text{ann}} v = \int ds f(s) \sigma_{\text{ann}}(s)$$

Public computer codes perform a calculation of the relic density

- DarkSUSY (neutralino in the MSSM)
- micrOMEGAs (all kinds of LSP in any model)
- SuperIso Relic (possibility to modify cosmological model)

[Gondolo *et al.* 2004]

[Bélanger *et al.* 2003]

[Arbey 2009]

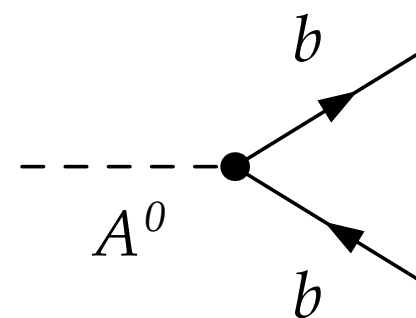
Motivation for Radiative Corrections

(Co)annihilation processes only implemented in public codes at tree-level

→ Corrections included only for some very sensitive quantities

→ e.g. bottom Yukawa-coupling

$$h_{Abb} \propto \frac{\bar{m}_b(Q)}{1 + \Delta_b} \tan \beta$$

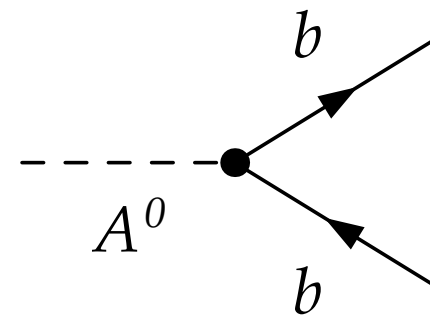


Motivation for Radiative Corrections

(Co)annihilation processes only implemented in public codes at tree-level

- Corrections included only for some very sensitive quantities
- e.g. bottom Yukawa-coupling

$$h_{Abb} \propto \frac{\bar{m}_b(Q)}{1 + \Delta_b} \tan \beta$$



Higher order corrections can have important impact on cross-sections

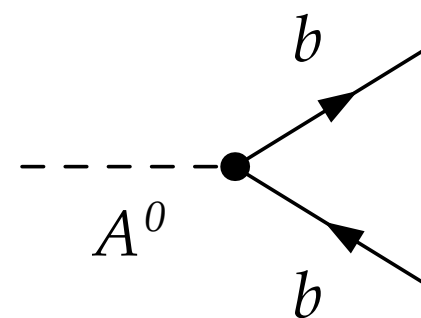
- QCD corrections significant due to strong coupling constant
- Quark final states give sizeable contributions to neutralino annihilation
- Radiative corrections affect the prediction of the relic density

Motivation for Radiative Corrections

(Co)annihilation processes only implemented in public codes at tree-level

- Corrections included only for some very sensitive quantities
- e.g. bottom Yukawa-coupling

$$h_{Abb} \propto \frac{\bar{m}_b(Q)}{1 + \Delta_b} \tan \beta$$

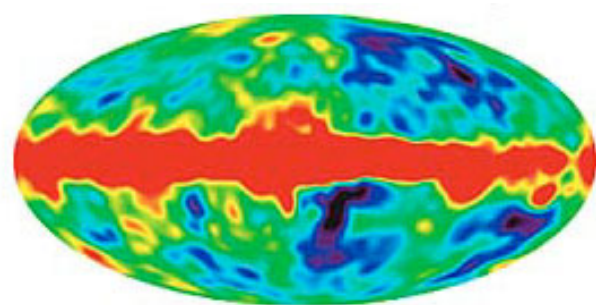


Higher order corrections can have important impact on cross-sections

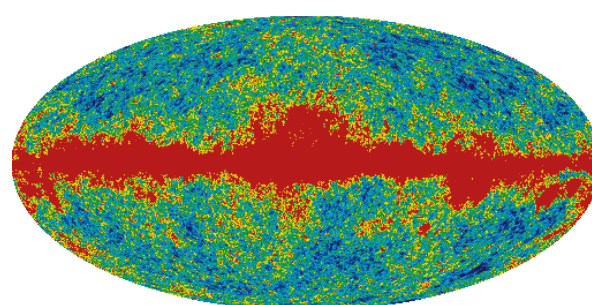
- QCD corrections significant due to strong coupling constant
- Quark final states give sizeable contributions to neutralino annihilation
- Radiative corrections affect the prediction of the relic density

Planck satellite will deliver new cosmological data in near future

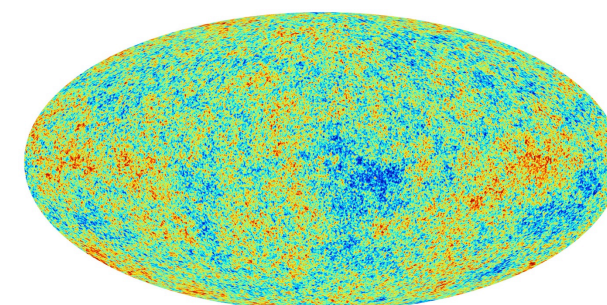
- Higher precision in theoretical predictions needed to match experimental improvements



COBE 1989



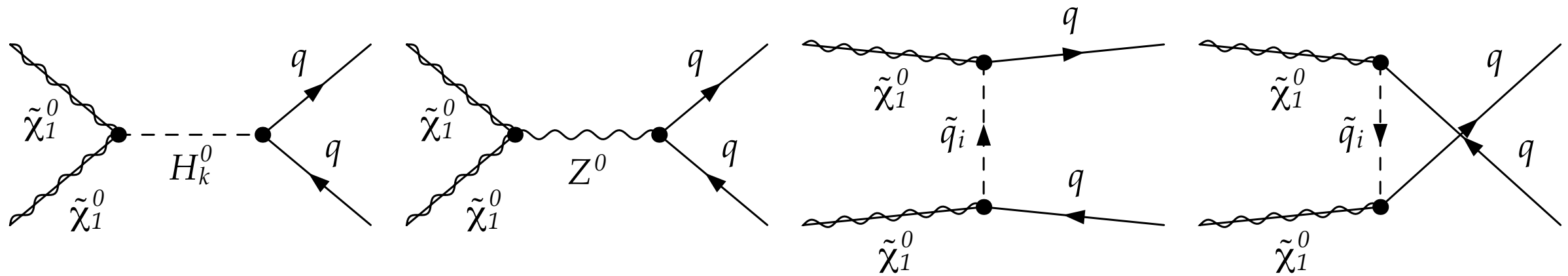
WMAP 2002



Planck 2009

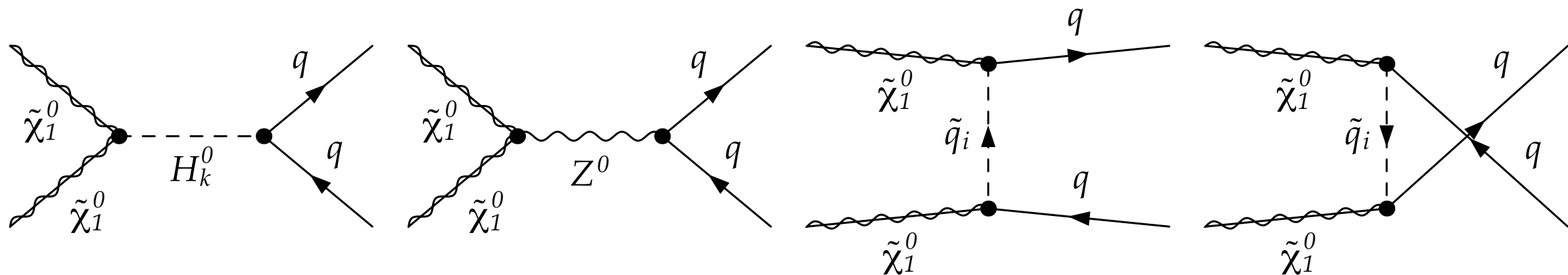
Neutralino Annihilation into (Heavy) Quarks

Annihilation cross-section receives sizeable contributions from quark final states



Neutralino Annihilation into (Heavy) Quarks

Annihilation cross-section receives sizeable contributions from quark final states



Higgs-exchange dominant in mSUGRA

- ▶ Low $m_{1/2}$ (if not excluded by LEP)

$$\tilde{\chi}\tilde{\chi} \rightarrow h^0 \rightarrow b\bar{b}$$

- ▶ A-Funnel at high $\tan\beta$

$$\tilde{\chi}\tilde{\chi} \rightarrow A^0 \rightarrow b\bar{b}$$

- ▶ Focus point region

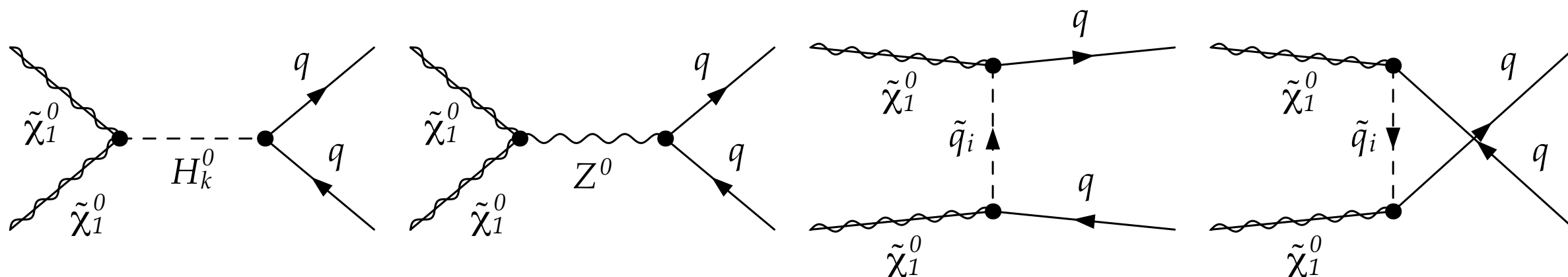
$$\tilde{\chi}\tilde{\chi} \rightarrow H^0 \rightarrow t\bar{t}$$

[Herrmann & Klasen 2007]

[Herrmann, Klasen, Kovařik 2009]

Neutralino Annihilation into (Heavy) Quarks

Annihilation cross-section receives sizeable contributions from quark final states



Higgs-exchange dominant in mSUGRA

- ▶ Low $m_{1/2}$ (if not excluded by LEP)

$$\tilde{\chi}\tilde{\chi} \rightarrow h^0 \rightarrow b\bar{b}$$

- ▶ A-Funnel at high $\tan\beta$

$$\tilde{\chi}\tilde{\chi} \rightarrow A^0 \rightarrow b\bar{b}$$

- ▶ Focus point region

$$\tilde{\chi}\tilde{\chi} \rightarrow H^0 \rightarrow t\bar{t}$$

[Herrmann & Klasen 2007]

[Herrmann, Klasen, Kovařik 2009]

Relax scalar or gaugino mass unification

- ▶ Non-universal Higgs masses (NUHM) or “compressed SUSY” (non-univ. gaugino masses)

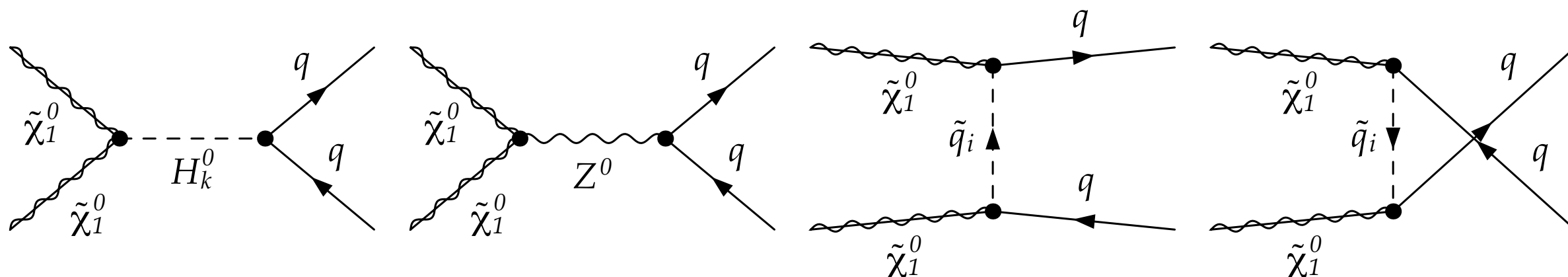
- ▶ A-Funnel already at low $\tan\beta$

- ▶ Larger higgsino-comp. favours Z^0 -exchange

$$\tilde{\chi}\tilde{\chi} \rightarrow Z^0 \rightarrow t\bar{t}$$

Neutralino Annihilation into (Heavy) Quarks

Annihilation cross-section receives sizeable contributions from quark final states



Higgs-exchange dominant in mSUGRA

- ▶ Low $m_{1/2}$ (if not excluded by LEP)

$$\tilde{\chi}\tilde{\chi} \rightarrow h^0 \rightarrow b\bar{b}$$

- ▶ A-Funnel at high $\tan\beta$

$$\tilde{\chi}\tilde{\chi} \rightarrow A^0 \rightarrow b\bar{b}$$

- ▶ Focus point region

$$\tilde{\chi}\tilde{\chi} \rightarrow H^0 \rightarrow t\bar{t}$$

[Herrmann & Klasen 2007]

[Herrmann, Klasen, Kovařík 2009]

Relax scalar or gaugino mass unification

- ▶ Non-universal Higgs masses (NUHM) or “compressed SUSY” (non-univ. gaugino masses)
- ▶ A-Funnel already at low $\tan\beta$
- ▶ Larger higgsino-comp. favours Z^0 -exchange

$$\tilde{\chi}\tilde{\chi} \rightarrow Z^0 \rightarrow t\bar{t}$$

Large mass splitting favours squark exchange

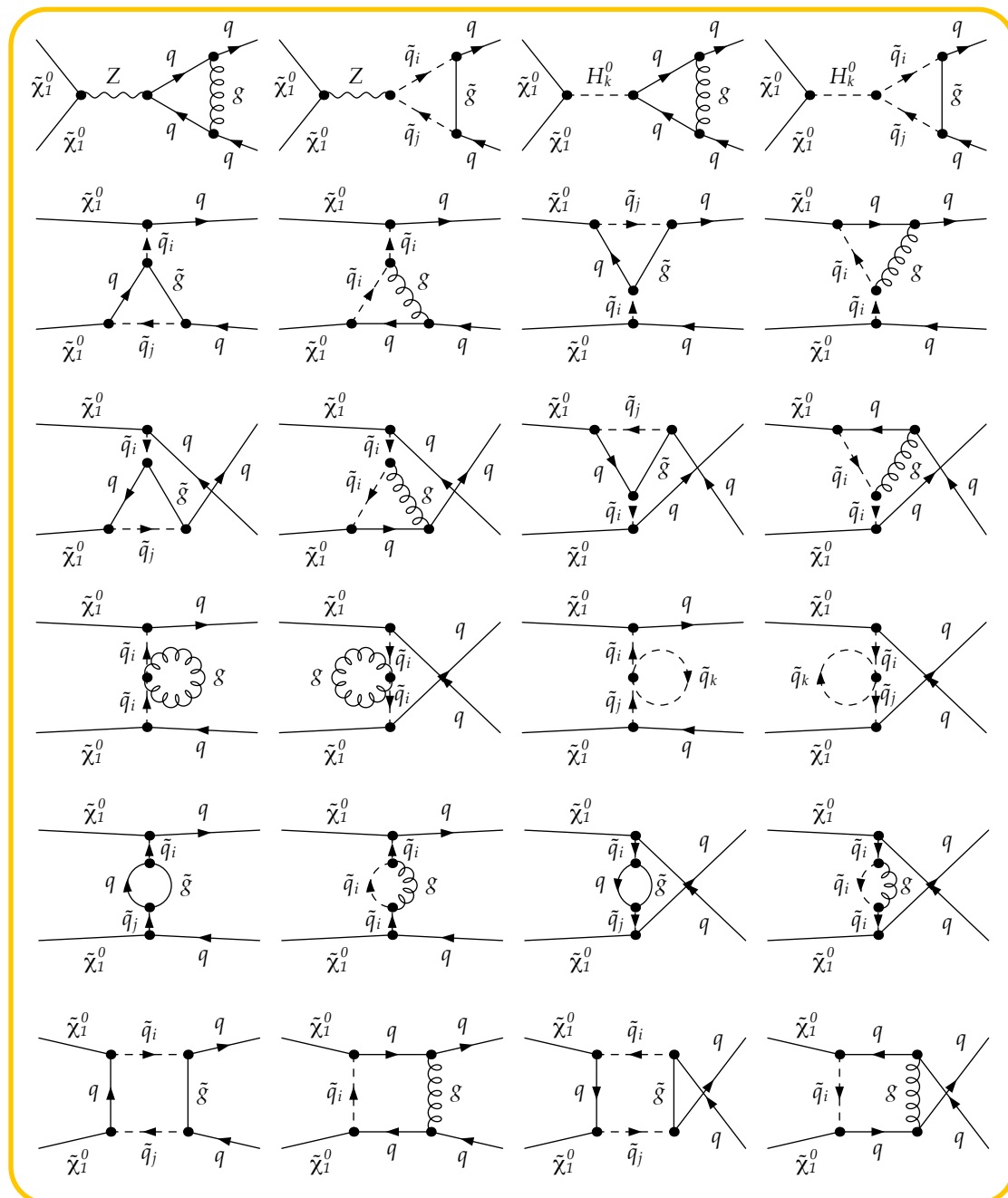
- ▶ Large trilinear coupling A_0

$$\tilde{\chi}\tilde{\chi} \rightarrow t\bar{t}$$

[Herrmann, Klasen, Kovařík 2009]

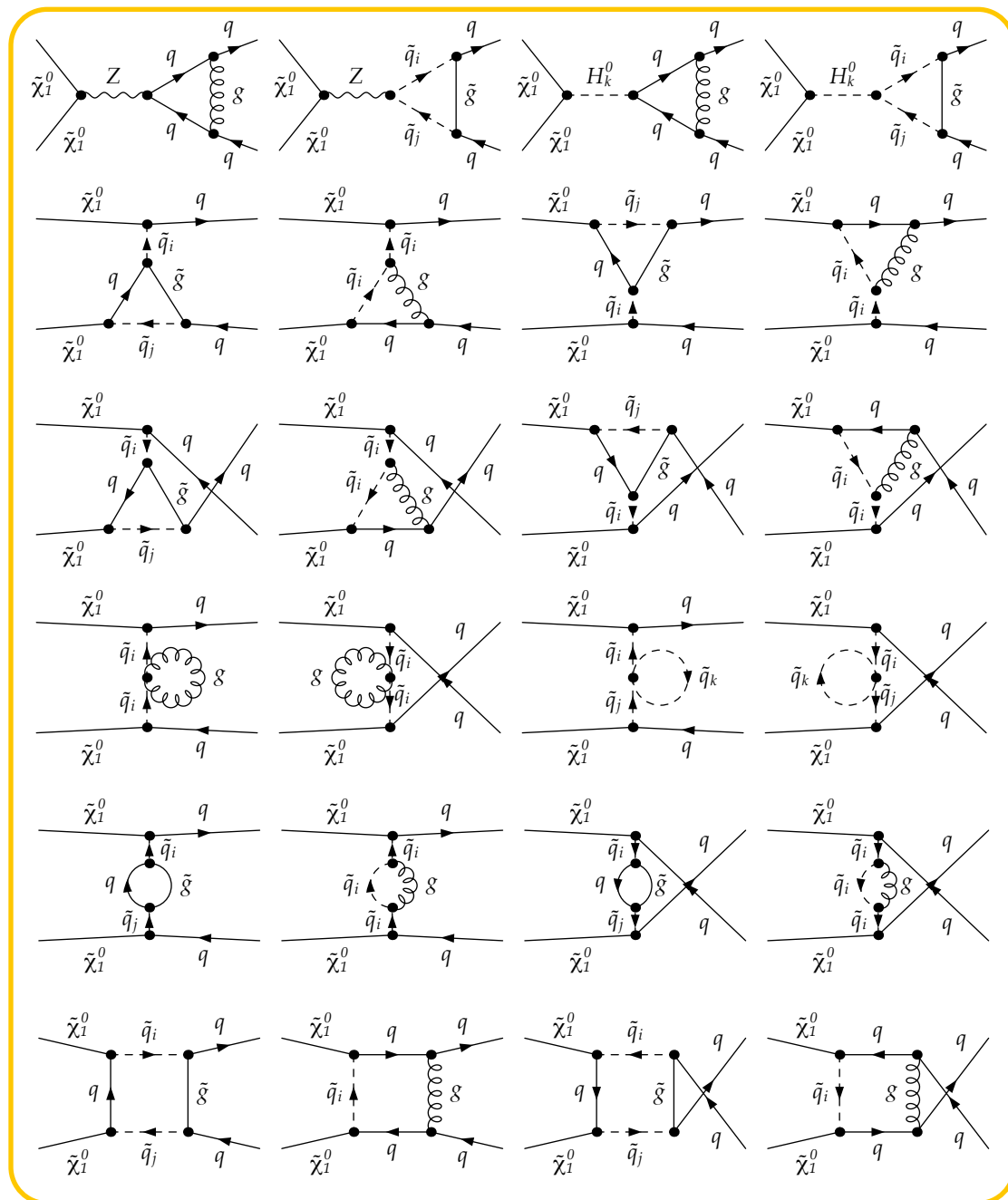
Supersymmetric QCD Corrections

One-loop contributions

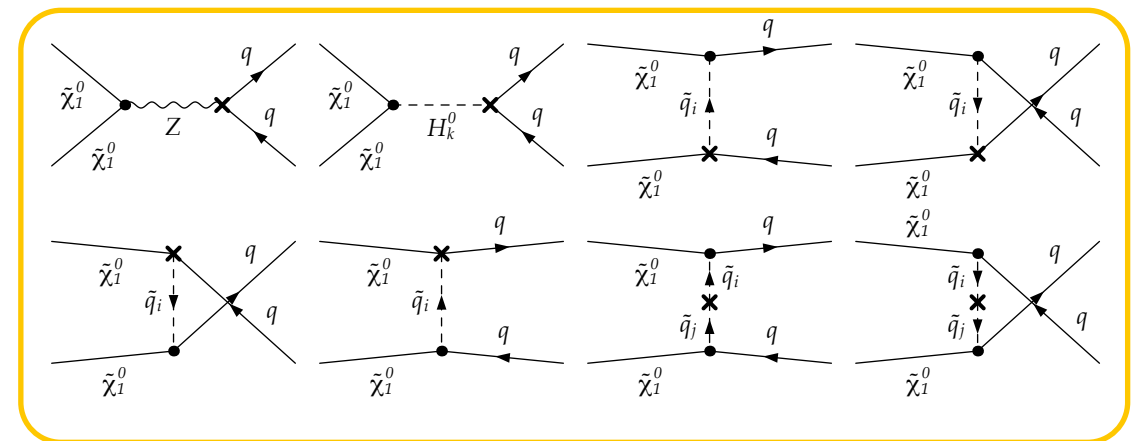


Supersymmetric QCD Corrections

One-loop contributions

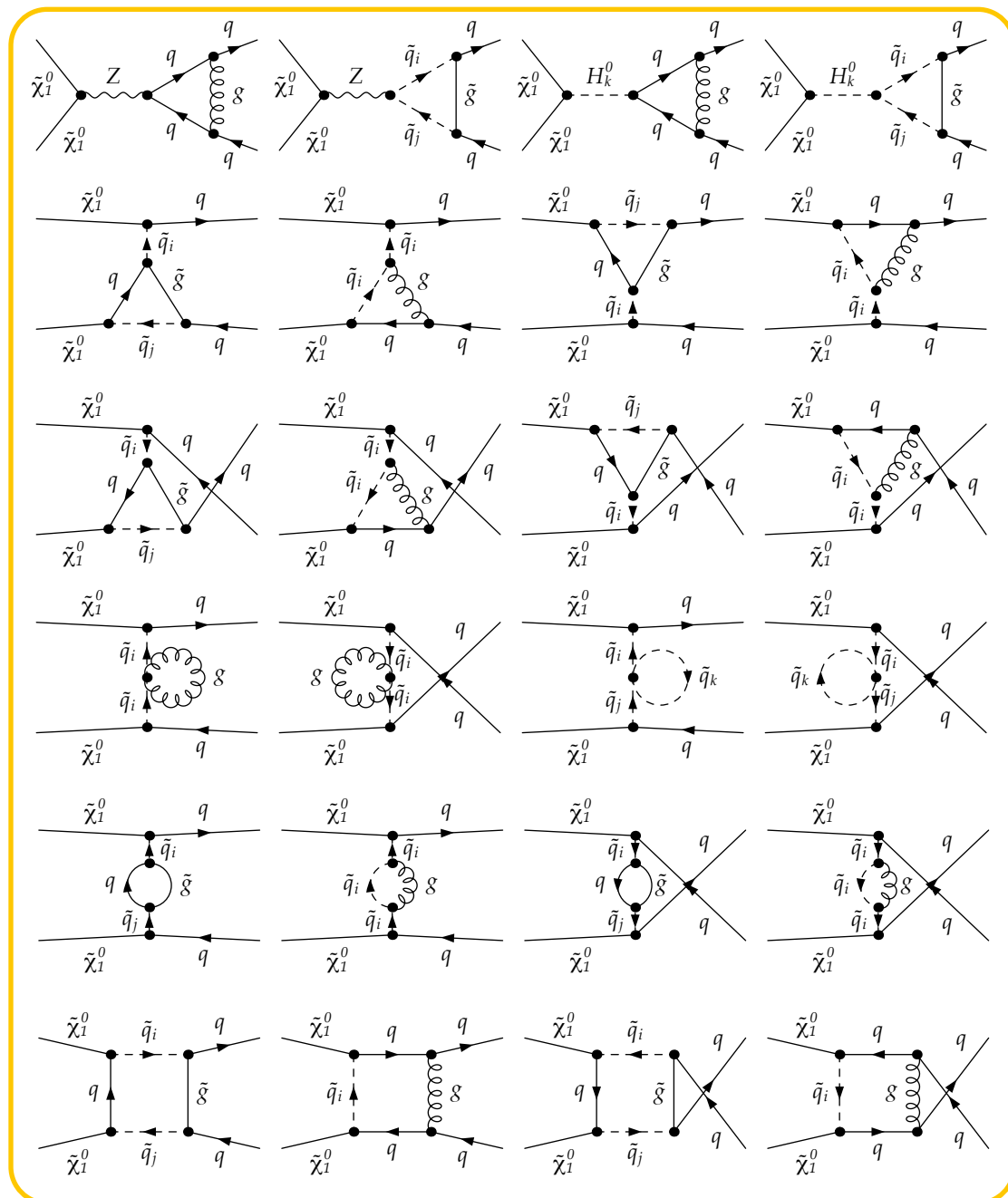


Counterterms [Kovařík et al. 2005]

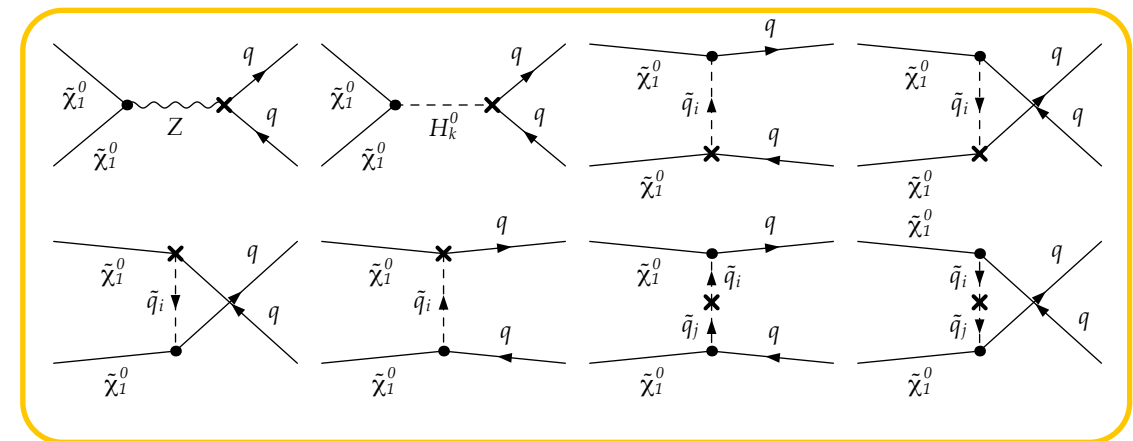


Supersymmetric QCD Corrections

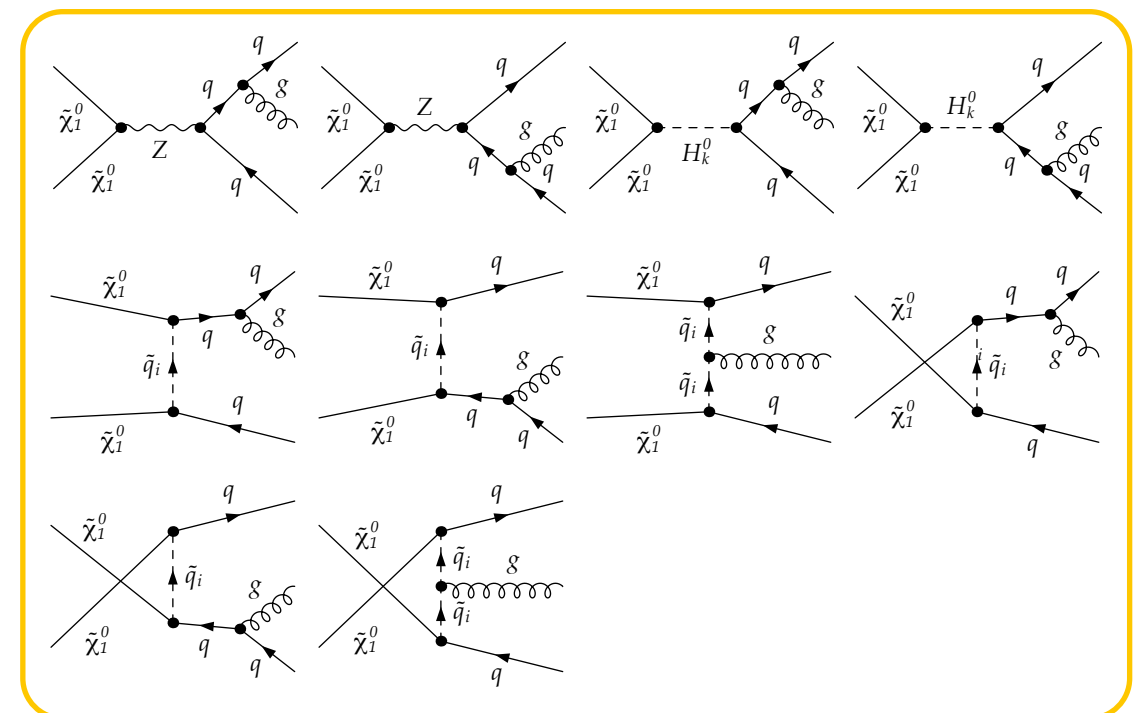
One-loop contributions



Counterterms [Kovařík et al. 2005]

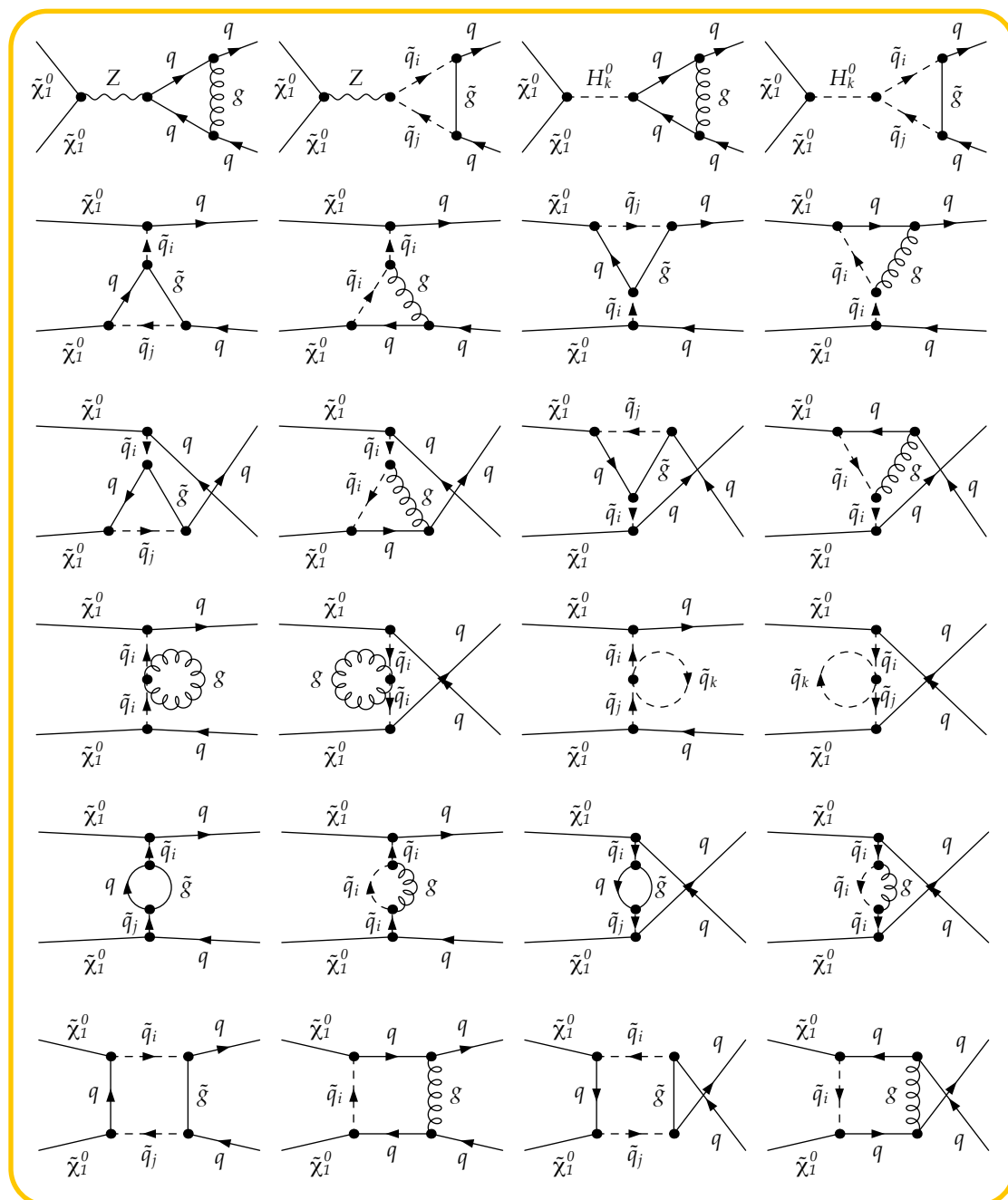


Real gluon emission

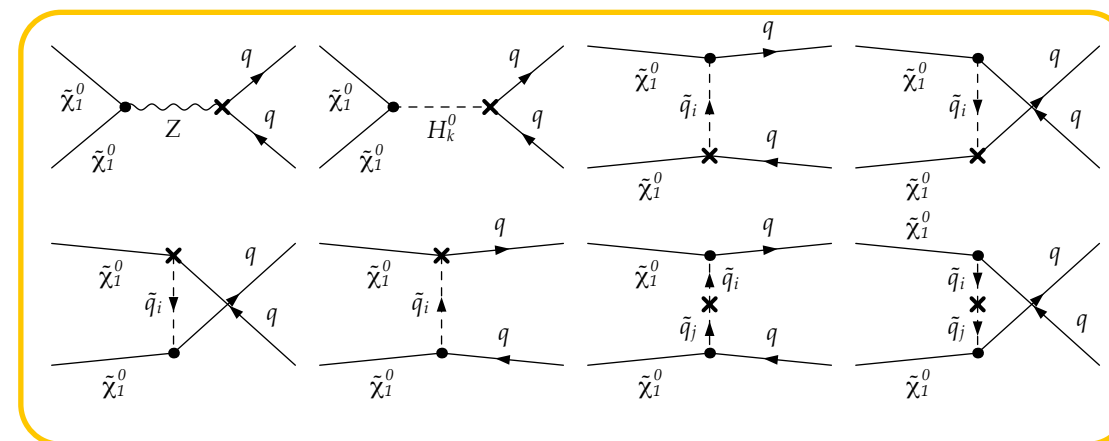


Supersymmetric QCD Corrections

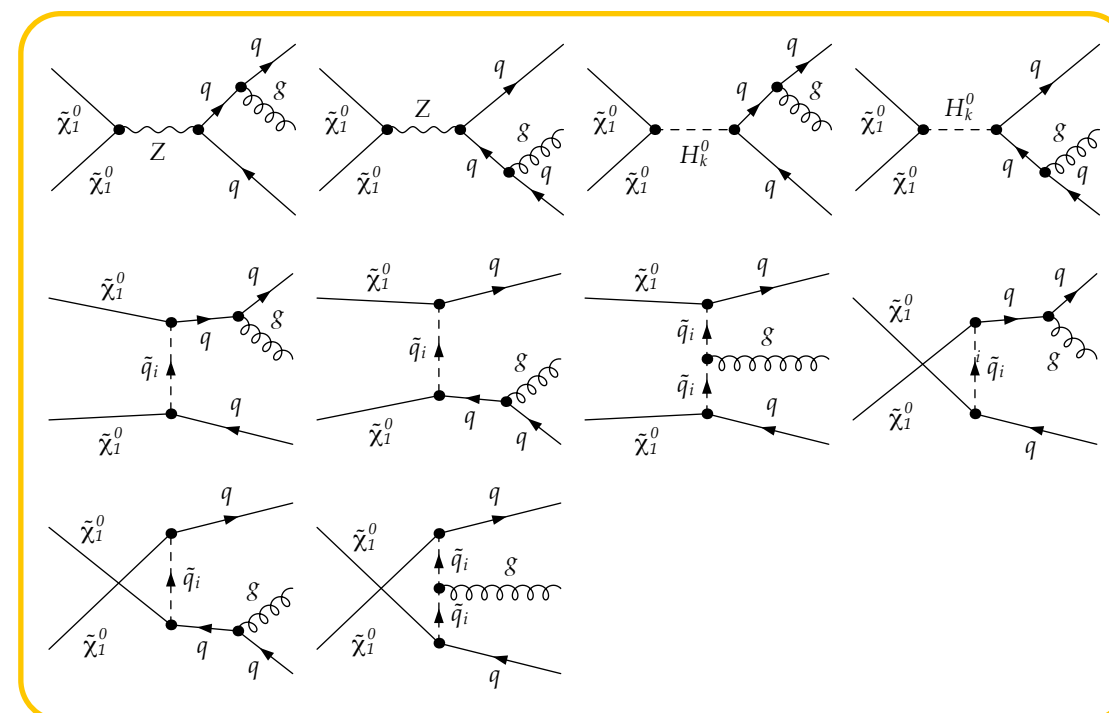
One-loop contributions



Counterterms [Kovařík et al. 2005]



Real gluon emission



Dipole Subtraction Method

[Catani et al. 2000-2002]

$$\sigma_{\text{NLO}} = \int_2 [d\sigma^{\text{V}} + \int_1 d\sigma^{\text{A}}] + \int_3 [d\sigma^{\text{R}} - d\sigma^{\text{A}}]$$

Scenario with dominant Z^0 -exchange

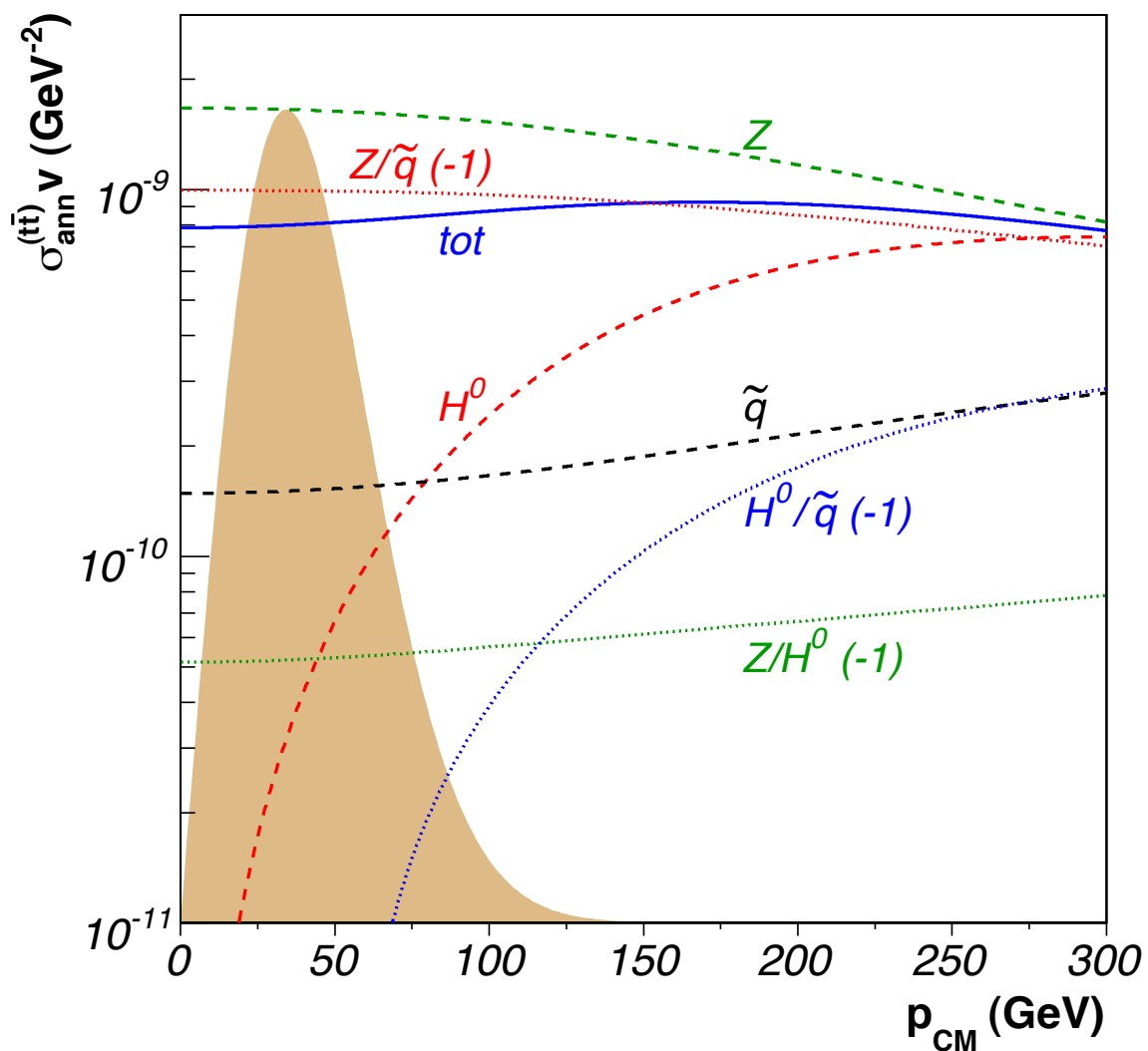
[Herrmann, Klasen, Kovařík 2009]

A_0	0
$\tan \beta$	10
$\text{sgn}(\mu)$	+

m_0	1500 GeV
$M_1 = M_2$	600 GeV
M_3	266 GeV
$\Omega_{\tilde{\chi}} h^2$	0.104
$t\bar{t}$	50.4%
$m_{\tilde{\chi}}$	235.6 GeV
$m_{\tilde{t}}$	939.0 GeV

Scenario with dominant Z^0 -exchange

[Herrmann, Klasen, Kovařík 2009]



A_0	0
$\tan \beta$	10
$\text{sgn}(\mu)$	+

m_0	1500 GeV
$M_1 = M_2$	600 GeV
M_3	266 GeV
$\Omega_{\tilde{\chi}} h^2$	0.104
$t\bar{t}$	50.4%
$m_{\tilde{\chi}}$	235.6 GeV
$m_{\tilde{t}}$	939.0 GeV

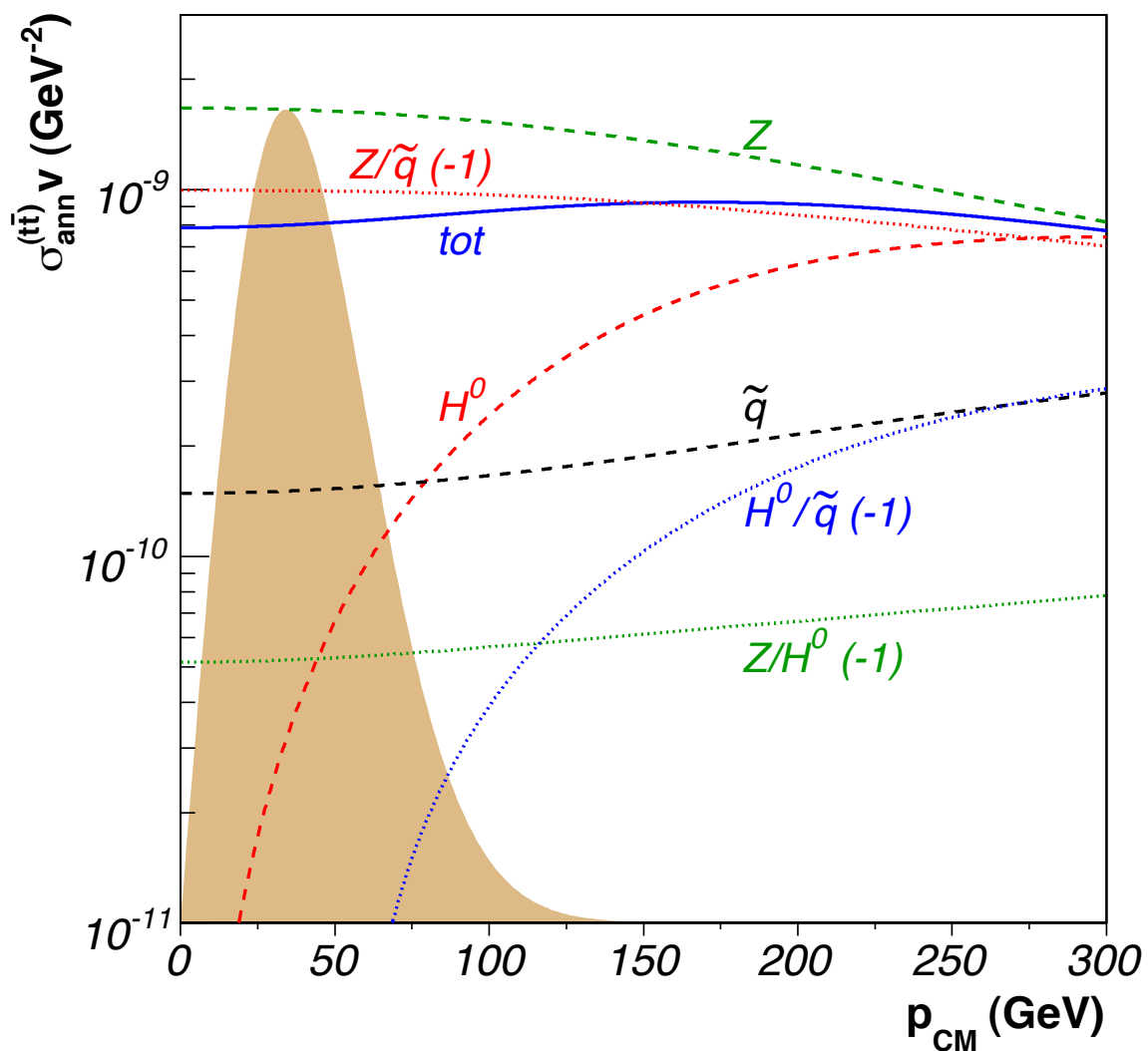
Z^0 - and squark-exchanges related by interference

→ Corrections to both diagrams important

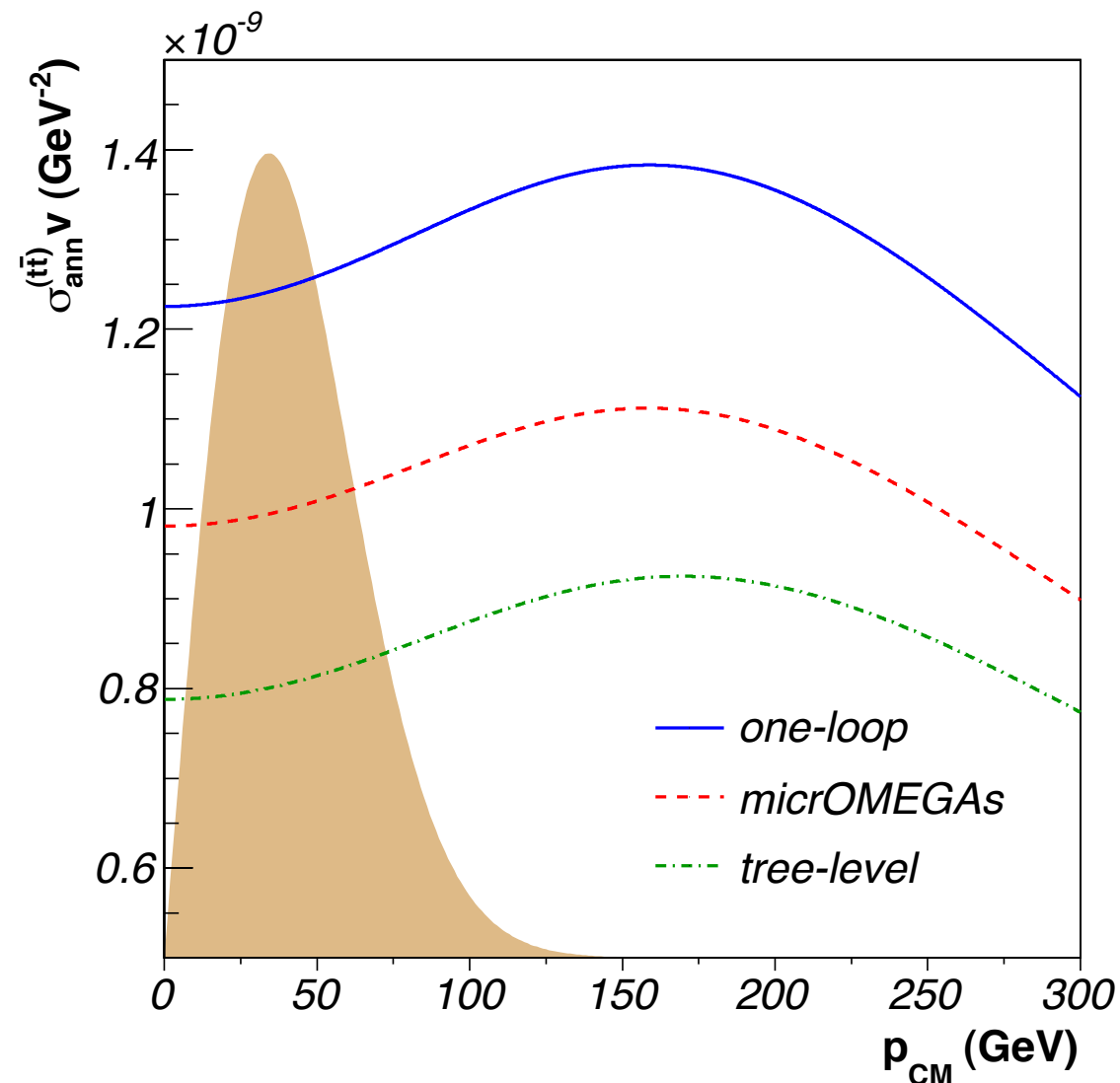
Scenario with dominant Z^0 -exchange

[Herrmann, Klasen, Kovařík 2009]

m_0	1500 GeV
$M_1 = M_2$	600 GeV
M_3	266 GeV
$\Omega_{\tilde{\chi}} h^2$	0.104
$t\bar{t}$	50.4%
$m_{\tilde{\chi}}$	235.6 GeV
$m_{\tilde{t}}$	939.0 GeV



A_0	0
$\tan \beta$	10
$\text{sgn}(\mu)$	+



Z^0 - and squark-exchanges related by interference

→ Corrections to both diagrams important

Full one-loop corrections increase cross-section

→ about 50% w.r.t. tree-level approximation

→ about 25% w.r.t. micrOMEGAs calculation

Scenario with dominant squark-exchange

[Herrmann, Klasen, Kovařík 2009]

A_0	-1200 GeV
$\tan \beta$	10
$\text{sgn}(\mu)$	+

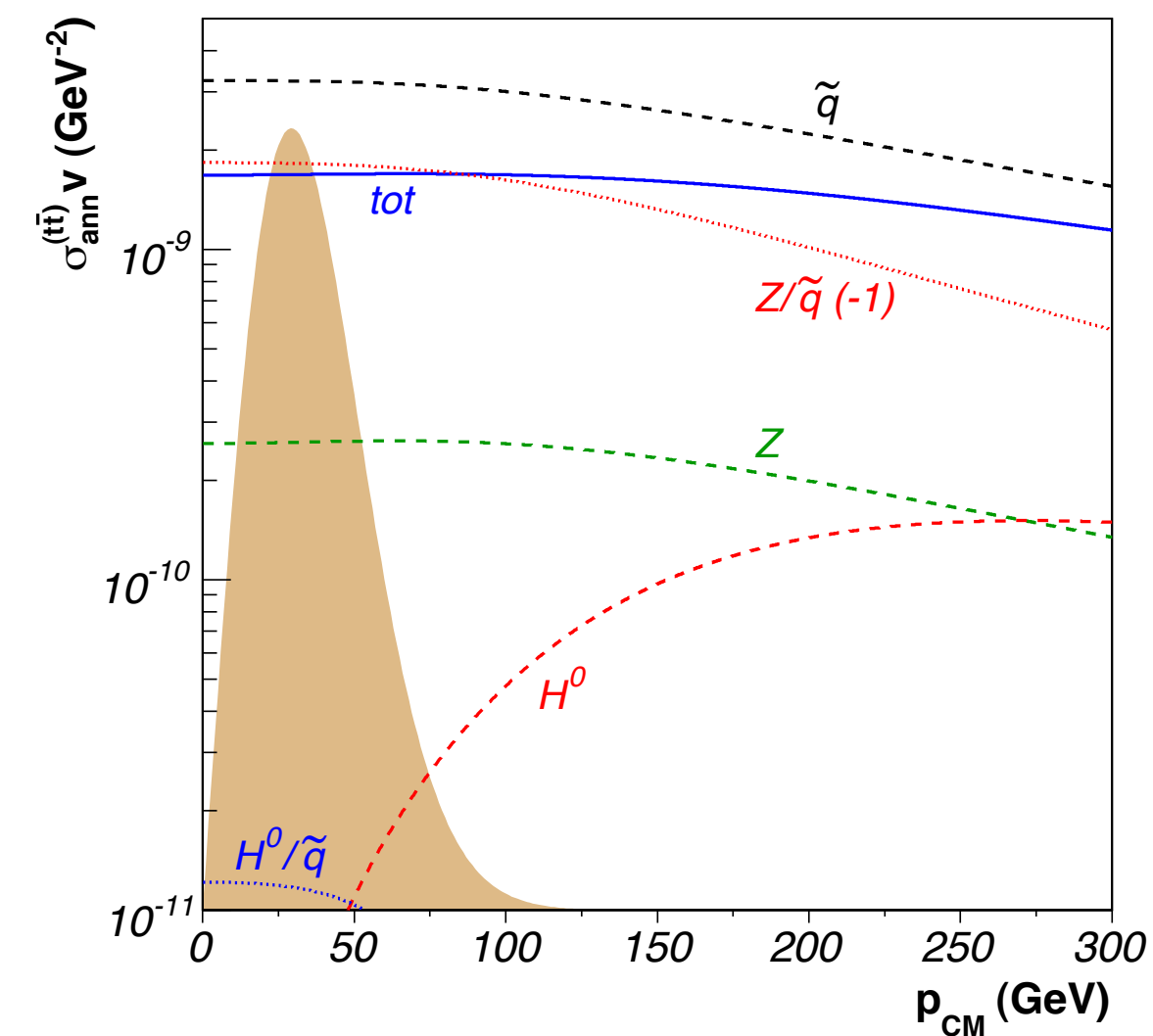
m_0	500 GeV
m_{H_u}	1250 GeV
m_{H_d}	2290 GeV
$m_{1/2}$	500 GeV
$\Omega_{\tilde{\chi}} h^2$	0.113
$t\bar{t}$	93.4%
$m_{\tilde{\chi}}$	200.7 GeV
$m_{\tilde{t}}$	259.3 GeV

Scenario with dominant squark-exchange

[Herrmann, Klasen, Kovařík 2009]

m_0	500 GeV
m_{H_u}	1250 GeV
m_{H_d}	2290 GeV
$m_{1/2}$	500 GeV
$\Omega_{\tilde{\chi}} h^2$	0.113
$t\bar{t}$	93.4%
$m_{\tilde{\chi}}$	200.7 GeV
$m_{\tilde{t}}$	259.3 GeV

A_0	-1200 GeV
$\tan \beta$	10
$\text{sgn}(\mu)$	+



Z^0 - and squark-exchanges related by interference

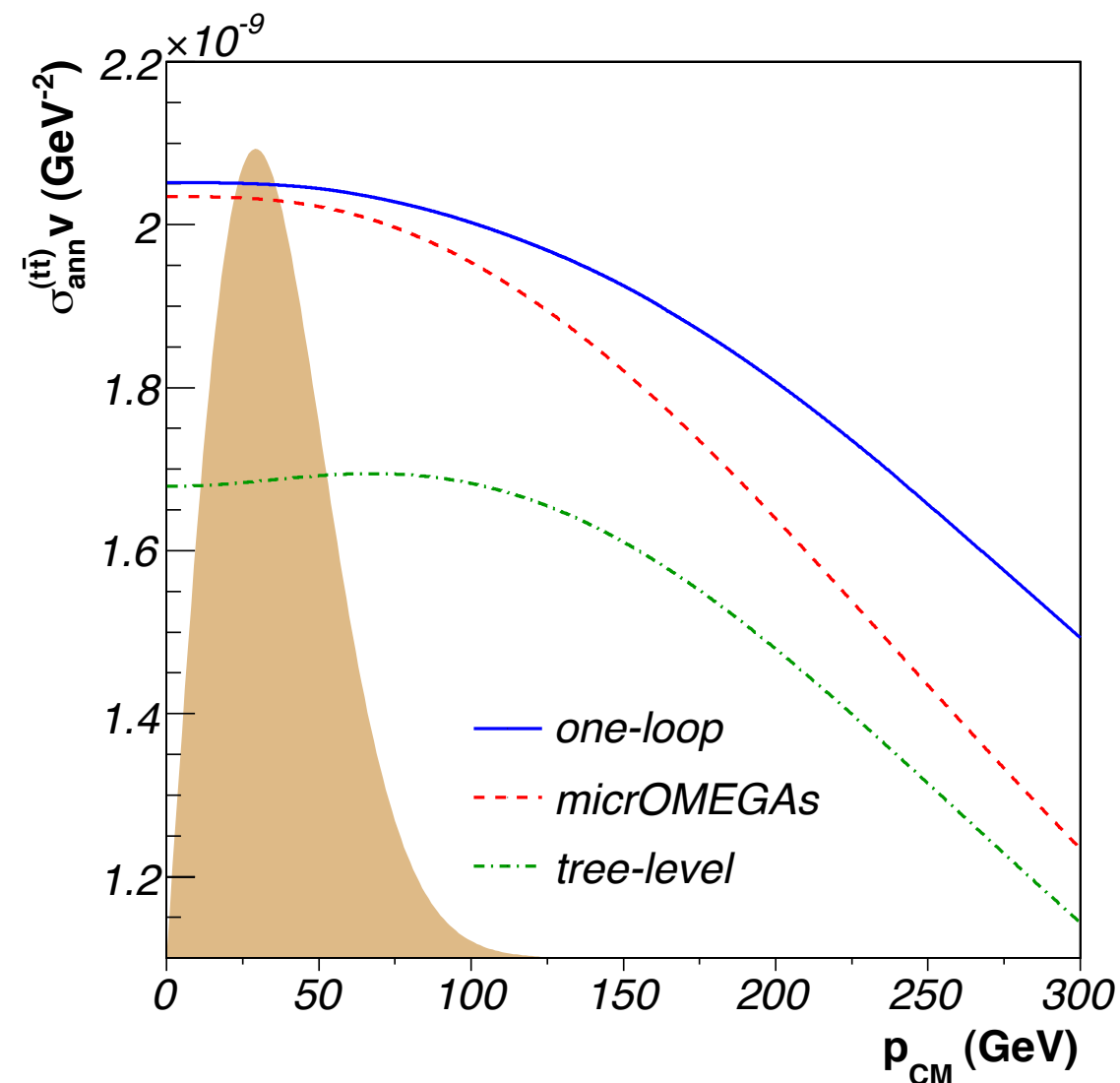
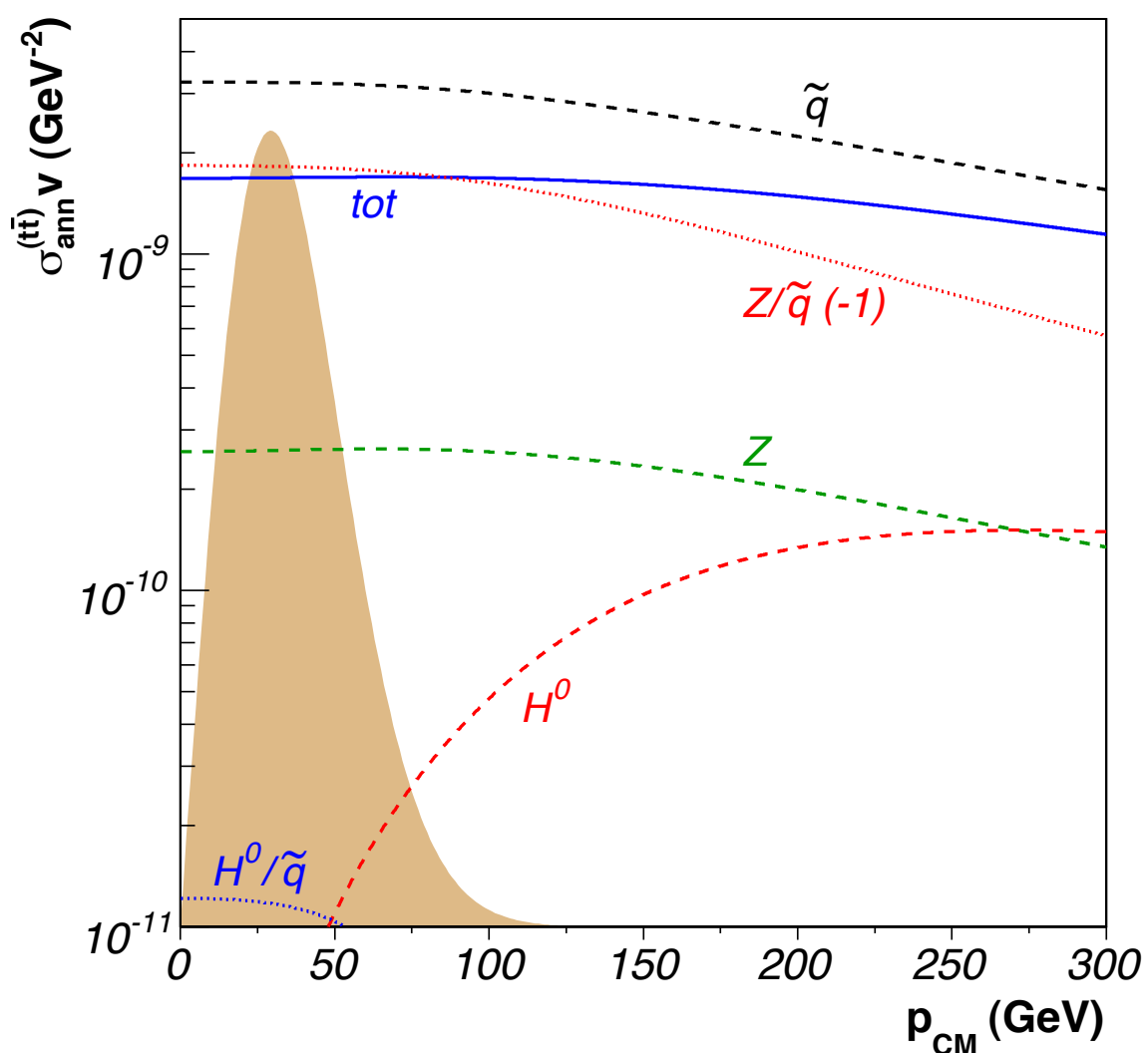
→ Corrections to both diagrams important

Scenario with dominant squark-exchange

[Herrmann, Klasen, Kovařík 2009]

m_0	500 GeV
m_{H_u}	1250 GeV
m_{H_d}	2290 GeV
$m_{1/2}$	500 GeV
$\Omega_{\tilde{\chi}} h^2$	0.113
$t\bar{t}$	93.4%
$m_{\tilde{\chi}}$	200.7 GeV
$m_{\tilde{t}}$	259.3 GeV

A_0	-1200 GeV
$\tan \beta$	10
$\text{sgn}(\mu)$	+



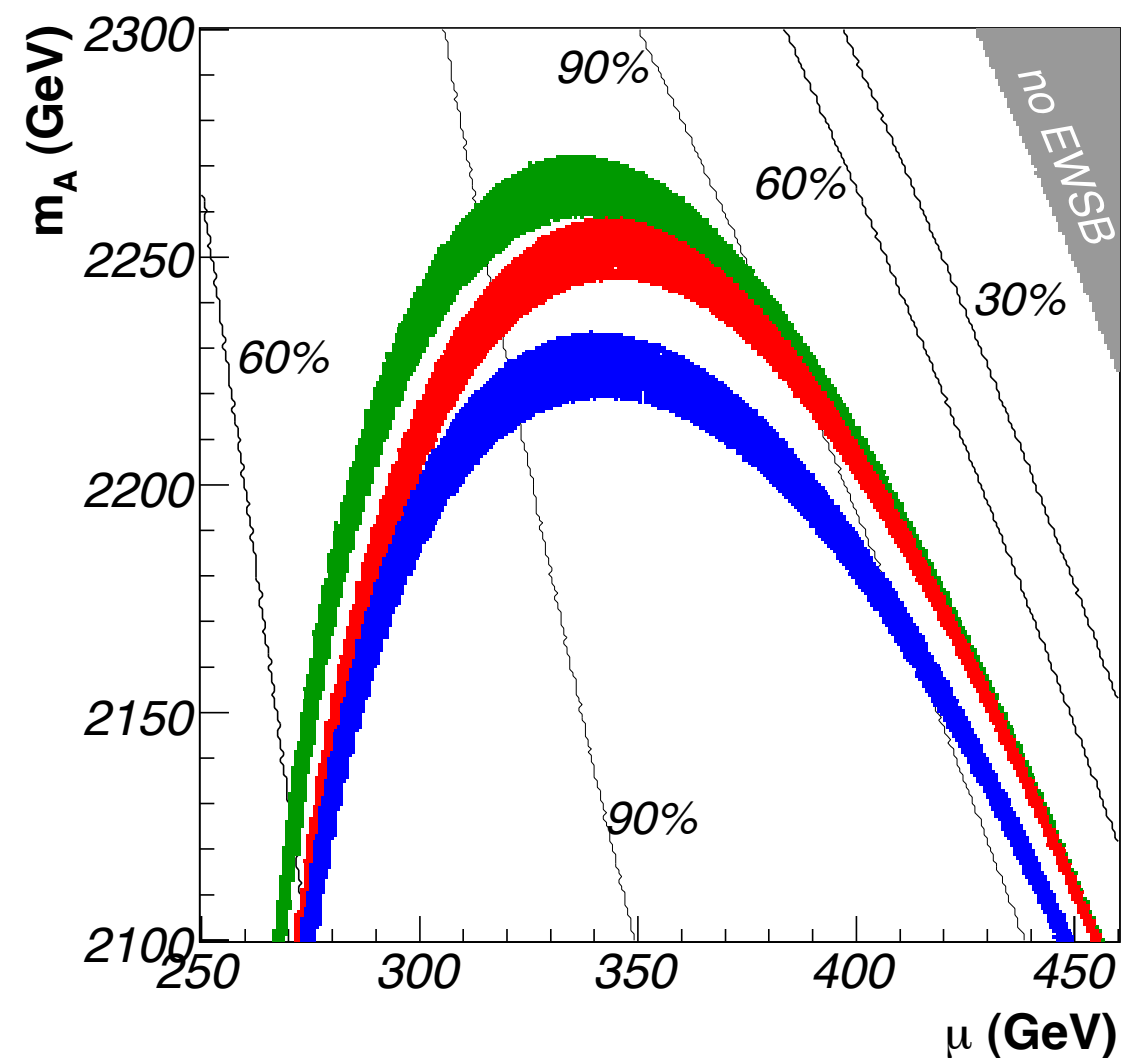
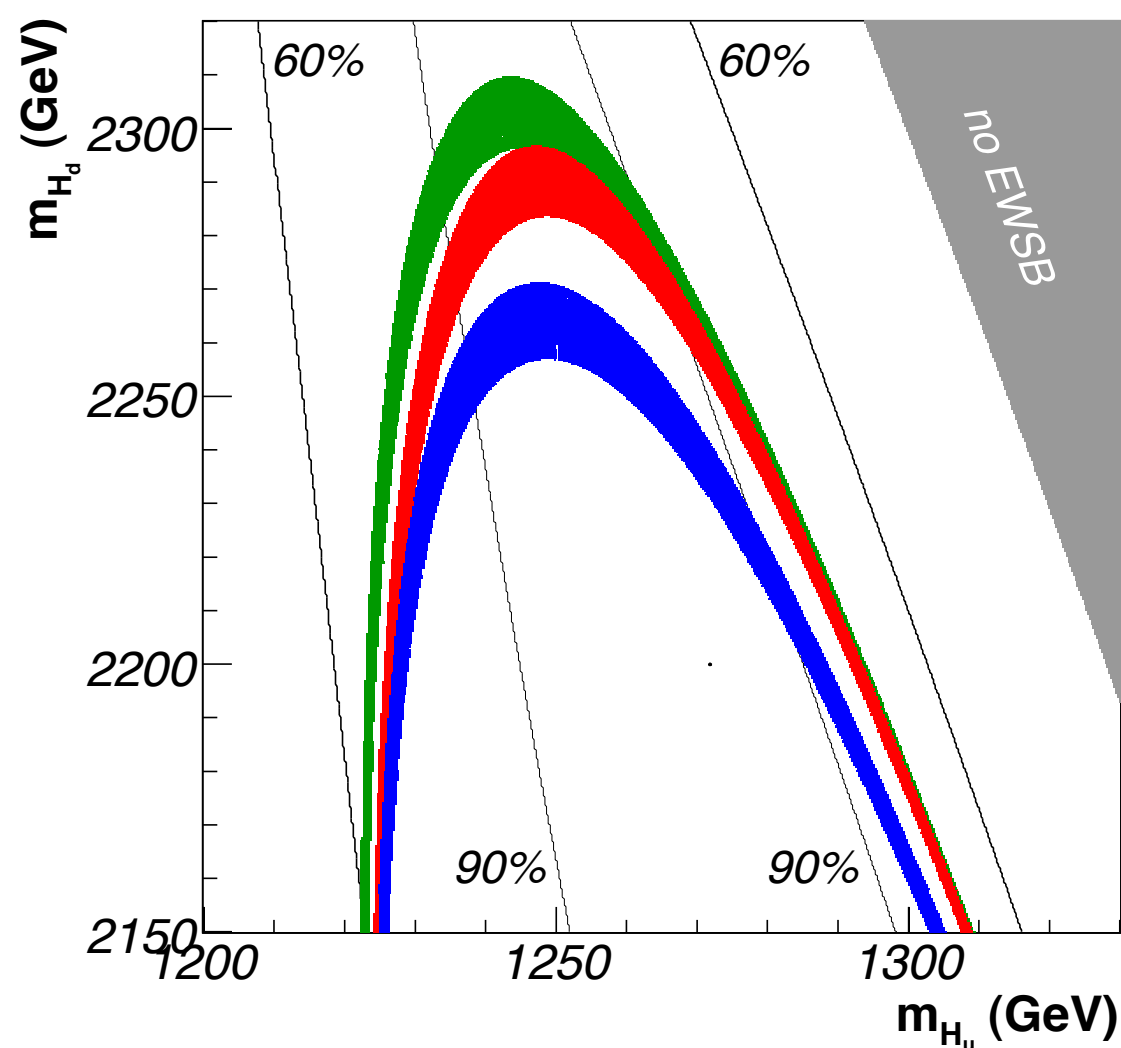
Z^0 - and squark-exchanges related by interference

→ Corrections to both diagrams important

Full one-loop corrections increase cross-section

→ about 25% w.r.t. tree-level approximation

Favoured regions of parameter space



Scenario with dominant squark-exchange

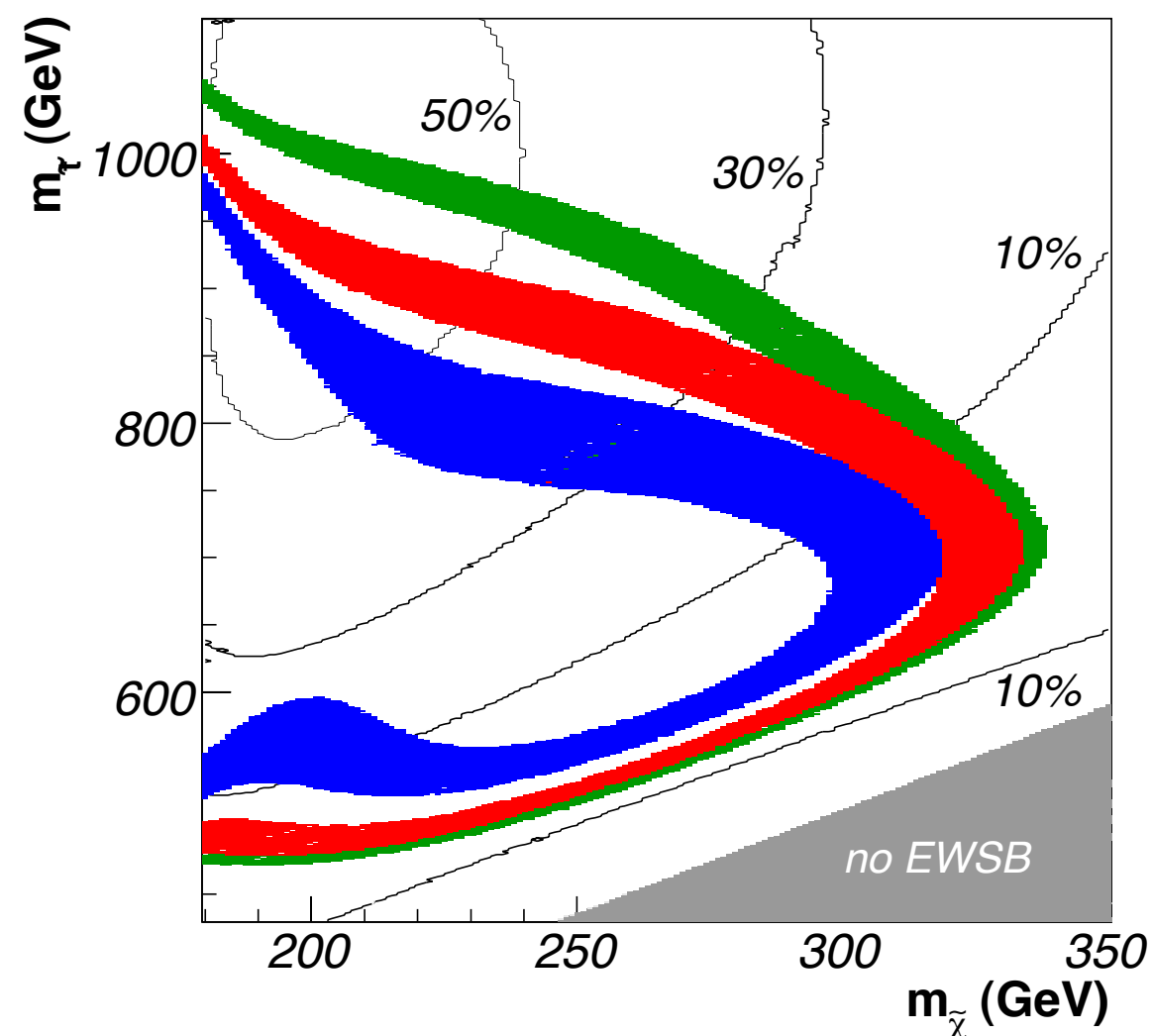
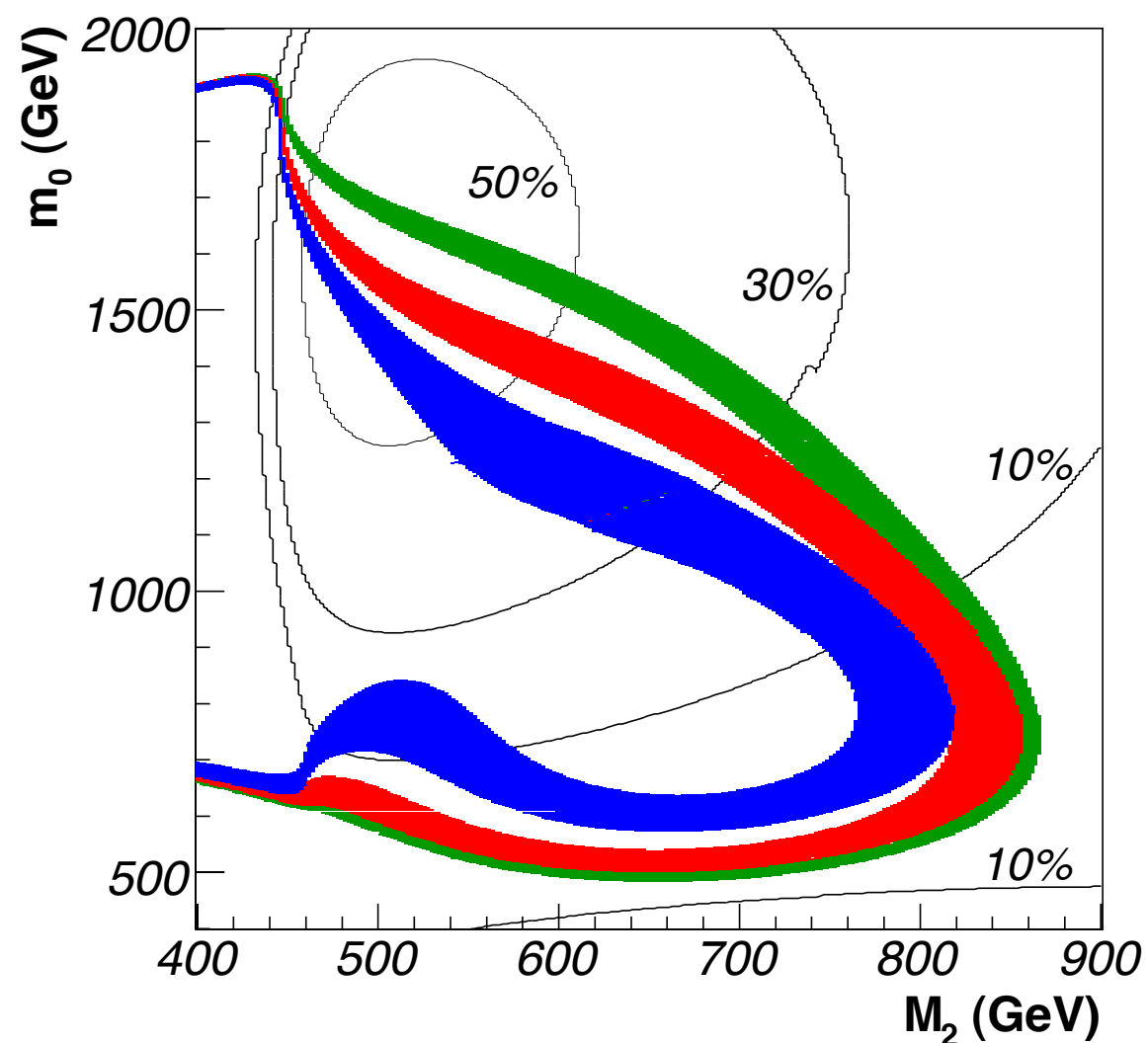
→ cross-section increased by about 25% w.r.t. tree-level calculation

Effect on annihilation cross-section directly reflected in the prediction of the relic density

→ about 20% decrease w.r.t. tree-level approximation

→ important shifts (~ 50 GeV) of the favoured regions of parameter space

Favoured regions of parameter space



Scenario with dominant Z^0 -exchange

- cross-section increased by about 50% w.r.t. tree-level calculation
- important shifts (up to 200 GeV) of the favoured regions of parameter space

Effect of SUSY-QCD corrections is larger than experimental uncertainty

- distinct bands corresponding to different levels of correction
- corrections even more important when Planck will deliver new cosmological data

Conclusions

Relic density calculation interesting tool to constrain SUSY models

- Additional information w.r.t. collider and precision data
- Identify (dis)favoured regions of MSSM parameter space

Conclusions

Relic density calculation interesting tool to constrain SUSY models

- Additional information w.r.t. collider and precision data
- Identify (dis)favoured regions of MSSM parameter space

SUSY-QCD corrections strongly influence neutralino annihilation cross-section

- Quark final states give sizeable contributions
- Effect on prediction of dark matter relic density
- Shift of favoured regions of the MSSM parameter space

Conclusions

Relic density calculation interesting tool to constrain SUSY models

- Additional information w.r.t. collider and precision data
- Identify (dis)favoured regions of MSSM parameter space

SUSY-QCD corrections strongly influence neutralino annihilation cross-section

- Quark final states give sizeable contributions
- Effect on prediction of dark matter relic density
- Shift of favoured regions of the MSSM parameter space

We have analyzed annihilation into “heavy” quark final states

- Higgs-resonances in mSUGRA scenarios [Herrmann & Klasen 2007; Herrmann, Klasen, Kovařík 2009]
- Z^0 -boson and squark exchanges in NUHM and “compressed SUSY” scenarios
- SUSY-QCD effects larger than experimental uncertainty [Herrmann, Klasen, Kovařík 2009]

Conclusions

Relic density calculation interesting tool to constrain SUSY models

- Additional information w.r.t. collider and precision data
- Identify (dis)favoured regions of MSSM parameter space

SUSY-QCD corrections strongly influence neutralino annihilation cross-section

- Quark final states give sizeable contributions
- Effect on prediction of dark matter relic density
- Shift of favoured regions of the MSSM parameter space

We have analyzed annihilation into “heavy” quark final states

- Higgs-resonances in mSUGRA scenarios [Herrmann & Klasen 2007; Herrmann, Klasen, Kovařík 2009]
- Z^0 -boson and squark exchanges in NUHM and “compressed SUSY” scenarios
- SUSY-QCD effects larger than experimental uncertainty [Herrmann, Klasen, Kovařík 2009]

Outlook

- Study “light” quark final states in “EW-scale” MSSM [Herrmann, Klasen, Kovařík *in preparation*]
- Study effects on indirect dark matter detection
- Consider other uncertainties in relic density calculation [Bélanger *et al.* 2005, Arbey *et al.* 2009]
- Include also EW corrections [Baro, Boudjema, Chalons, Hao, Semenov 2008-2009; Talk by Guillaume Chalons]