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# Towards LHC Phenomenology beyond Leading Order

Gudrun Heinrich



Institute for Particle Physics Phenomenology



GDR Terascale, Heidelberg, 16.10.09

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on behalf of the

# GOLEM collaboration

( General One-Loop Evaluator of Matrix elements )

T. Binoth, G. Cullen, J.-Ph. Guillet, GH, S. Karg,  
N. Kauer, E. Pilon, T. Reiter,  
M. Rodgers, I. Wigmore

# The LHC ...

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... has been planned long time ago ...



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Linear Colliders also seem to have been supported . . .



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... so why do we say we are entering a  
New Era in Particle Physics?

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New Era in Particle Physics?



... because instead of hunting buffaloes,  
we are now hunting Higgs bosons ...

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process	events/sec	
QCD jets $E_T > 150$ GeV	100	background
$W \rightarrow e\nu$	15	background
$t\bar{t}$	1	background
Higgs, $m_H \sim 130$ GeV	0.02	signal
gluinos, $m \sim 1$ TeV	0.001	signal

# LHC phases

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- early phase:  
have to "rediscover" the Standard Model, control underlying event, jet energy scale, . . .  
do we control the SM predictions well enough to claim a discovery soon?

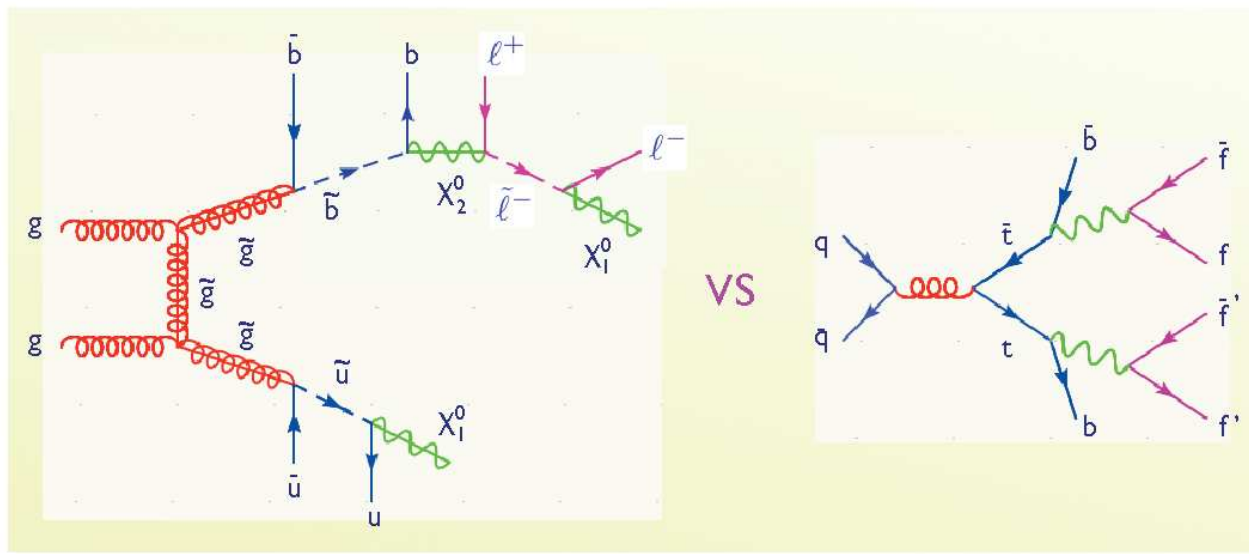
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- **"model discrimination phase":**  
have to measure particle masses, widths, couplings, CP properties, ...  
do we control the SM/MSSM predictions well enough?

# (heavy) SUSY particles:

- decay through cascades emitting quarks and leptons
- signatures: energetic jets and leptons, missing  $E_T$
- QCD radiation generates additional hard jets  
⇒ multi-particle final states



F. Maltoni

# identifying New Physics at hadron colliders

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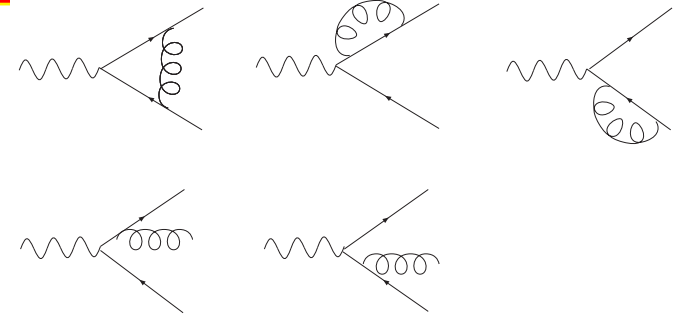
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 $\Rightarrow$  calculations of higher orders increasingly difficult  
until recently: LO tools highly automated, whereas  
NLO calculations tedious case-by case exercises
- **now paradigm change:**  
we are moving towards **automated NLO tools**

# ingredients for $m$ -particle observable at NLO

virtual part (one-loop integrals):

$$\mathcal{A}_{NLO}^V = A_2/\epsilon^2 + A_1/\epsilon + A_0^{(v)}$$

$$d\sigma^V \sim \text{Re} \left( A_{LO}^\dagger \mathcal{A}_{NLO}^V \right)$$



real radiation part: soft/collinear emission of massless particles

⇒ need subtraction terms

$$\Rightarrow \int_{\text{sing}} d\sigma^S = -A_2/\epsilon^2 - A_1/\epsilon + A_0^{(r)}$$

$$\sigma^{NLO} = \underbrace{\int_{m+1} \left[ d\sigma^R - d\sigma^S \right]_{\epsilon=0}}_{\text{numerically}} + \underbrace{\int_m \left[ \underbrace{d\sigma^V}_{\text{cancel poles}} + \underbrace{\int_s d\sigma^S}_{\text{analytically}} \right]_{\epsilon=0}}_{\text{numerically}}$$

# Modular structure

Tree Modules

One-Loop Module

IR Modules

$$|\mathcal{A}^{LO}|^2 \oplus$$

$$2 \operatorname{Re}(\mathcal{A}^{LO\dagger} \mathcal{A}^{NLO,V}) \oplus$$

$$\sum_j \int_j \mathcal{S}_j$$

$$|\mathcal{A}^{NLO,R}|^2$$

$$\ominus \sum_j \mathcal{S}_j$$



has been bottleneck in past years

# progress

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- more **efficient** techniques to calculate **loop amplitudes**
  - unitarity-based methods  
e.g. BlackHat, Rocket, CutTools, analytic, ...
  - improved methods based on Feynman diagrams  
e.g. GOLEM, Denner et. al, ...

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- **automatisation** of IR modules
- **use existing technology** from leading order tools  
LO tools can provide:
  - event generation
  - phase space integration
  - histogramming tools
  - subtraction terms for soft/collinear radiation

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  - phase space integration
  - histogramming tools
  - subtraction terms for soft/collinear radiation
- **matching NLO** amplitudes with parton showers



# 2009 status of NLO wishlist for LHC

$pp \rightarrow W W \text{ jet}$	Denner/Dittmaier/Kallweit/Uwer, Ellis/Campbell/Zanderighi
$pp \rightarrow Z Z \text{ jet}$	Binoth/Guillet/Karg/Kauer/Sanguinetti
$pp \rightarrow t\bar{t} b\bar{b}$	Bredenstein/Denner/Dittmaier/Pozzorini
$pp \rightarrow t\bar{t} + 2 \text{ jets}$	
$pp \rightarrow Z Z Z$	Lazopoulos/Melnikov/Petriello, Hankele/Zeppenfeld
$pp \rightarrow V V V$	Binoth/Ossola/Papadopoulos/Pittau, Zeppenfeld et al.
$pp \rightarrow V V b\bar{b}$	
$pp \rightarrow V V + 2 \text{ jets}$	VBF: Bozzi/Jäger/Oleari/Zeppenfeld, VBFNLO coll.
$pp \rightarrow W + 3 \text{ jets}$	BlackHat coll.; Ellis/Giele/Kunszt/Melnikov/Zanderighi*
$pp \rightarrow b\bar{b}b\bar{b}$	Binoth/Guffanti/Guillet/Reiter/Reuter
$pp \rightarrow t\bar{t} \text{ jet}$	Dittmaier/Uwer/Weinzierl
$pp \rightarrow t\bar{t} Z$	Lazopoulos/McElmurry/Melnikov/Petriello
$pp \rightarrow b\bar{b} Z, b\bar{b} W$	Febres Cordero/Reina/Wackerroth

● done ● partial results \* leading colour only

# Interface

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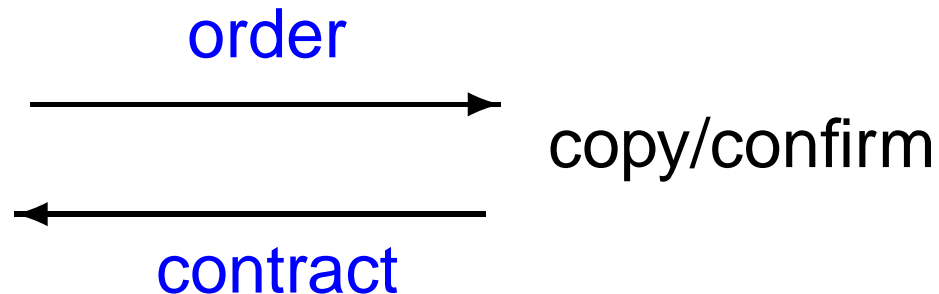
details worked out at Les Houches 2009 workshop on TeV colliders

Monte Carlo tool (MC)

One-Loop-Provider (OLP)

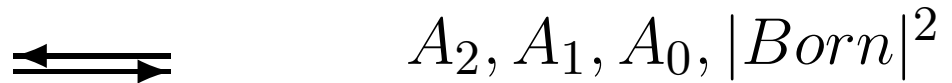
initialisation:

process info  
CH summed  
model parameters  
fix scheme  
...



runtime:

events



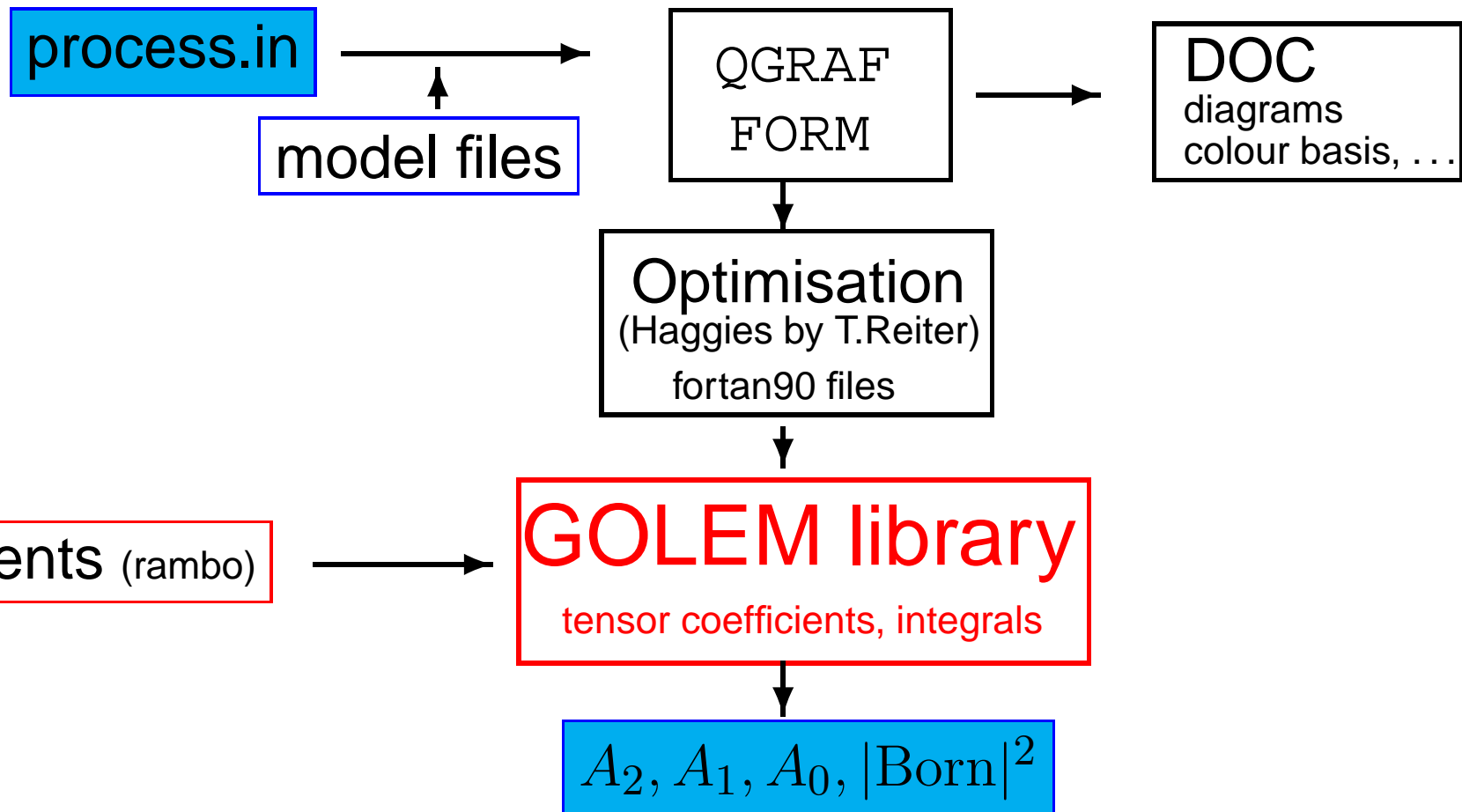
standard interface

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# GOLEM

## General One-Loop Evaluator of Matrix elements

[ Binoth, Cullen, Guillet, GH, Karg, Kauer, Pilon, Reiter, Rodgers, Wigmore ]



# Golem strong points

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- can deal with an **arbitrary** number of **mass scales**  
link LoopTools for finite massive boxes
- **colour** does not add additional complexity
- rational parts are **"for free"**
- efficient use of recursive structure  
caching system
- projection onto **helicity** states  
exploit spinor helicity techniques, gauge cancellations, smaller building blocks
- collaboration has several independent programs  
⇒ strong checks  
alternatively based on FeynArts diagram generator ⇒ MSSM model file  
different levels of analytic simplifications
- can avoid spurious singularities from  
inverse **Gram determinants**

# Gram determinants

- $\det(G)$  ,  $G_{ij} = 2 r_i \cdot r_j$  ,  $r_i$  external momenta  
can become very small in certain regions of phase space  $\Rightarrow 1/\det(G)$  can lead to numerical instabilities
- reduction  $N \geq 5 \rightarrow N = 4$ : inverse Gram determinants **completely absent**  
[Binoth, Guillet, GH, Pilon, Schubert '05, Denner/Dittmaier '05]
- reduction of  $N \leq 4$  tensor integrals:  
introduces spurious  $1/\det(G) \sim 1/B$

$$I_4^{n+2}(j) = \frac{1}{B} \left\{ b_j I_4^{n+2} + \frac{1}{2} \sum_k \mathcal{S}_{jk}^{-1} I_3^{n,\{k\}} - \frac{1}{2} \sum_{k \in S \setminus \{j\}} b_k I_3^{n,\{k\}}(j) \right\}$$

$$I_4^{n+2}(j_1, j_2) \sim \frac{1}{B^2} , I_4^{n+2}(j_1, j_2, j_3) \sim \frac{1}{B^3} \dots$$

$$B = \det(G)/\det(\mathcal{S}) (-1)^{N+1} ; \quad \mathcal{S}_{ij} = (r_i - r_j)^2 - m_i^2 - m_j^2$$

# Gram determinants

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solution:

do not reduce if  $B$  is small

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- calculate integral numerically
- special feature:  
1-dimensional integral representation for all massless basis integrals  
⇒ numerical integration fast

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do not reduce if  $B$  is small

- calculate integral **numerically**
- **special feature:**  
1-dimensional integral representation for all massless basis integrals  
⇒ numerical integration **fast**
- **under construction:**  
same feature for integrals with internal masses



# recent developments

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- under construction:
  - allow for complex masses  $\Rightarrow$   
deal with unstable particles (SUSY!)
  - renormalisation

# recent developments

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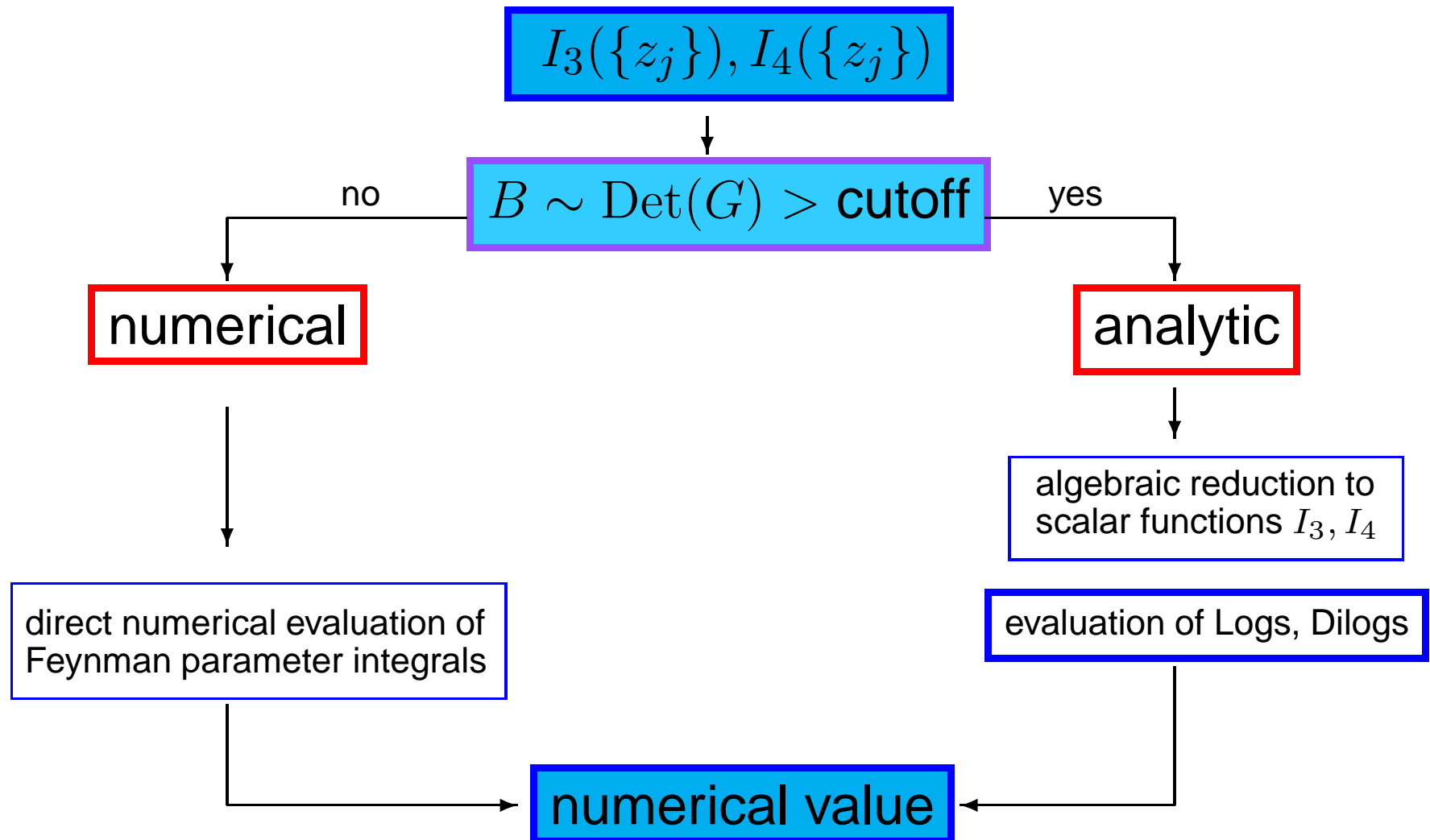
- **under construction:**
  - allow for **complex masses**  $\Rightarrow$  deal with **unstable** particles (**SUSY!**)
  - renormalisation
- **done:** implementation of **all** IR divergent integrals, scalar and tensor  
(massless propagators: publicly available; massive: will become public very soon)

## note:

LoopTools: no dim.reg. for IR divergent integrals

QCDLoop: only scalar integrals

# treatment of 3-and 4-point integrals



# example six-photon amplitude

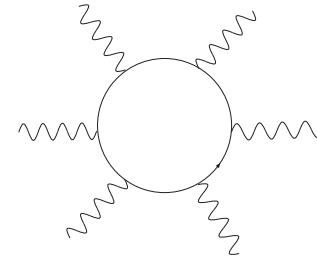
[Mahlon 94] (special helicity configurations only)

[Nagy, Soper 06; Gong, Nagy, Soper 08] (purely numerically)

[Binoth, Gehrmann, GH, Mastrolia 07]

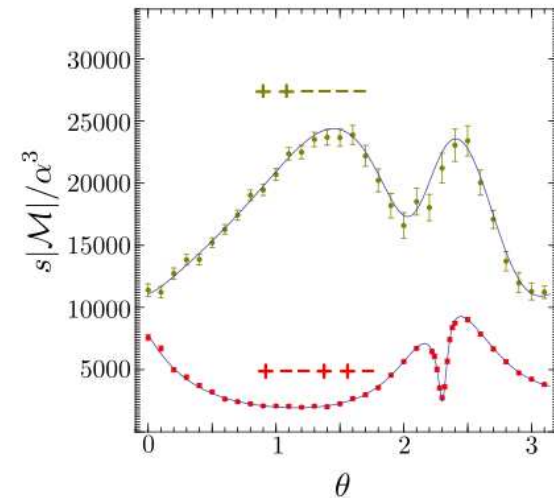
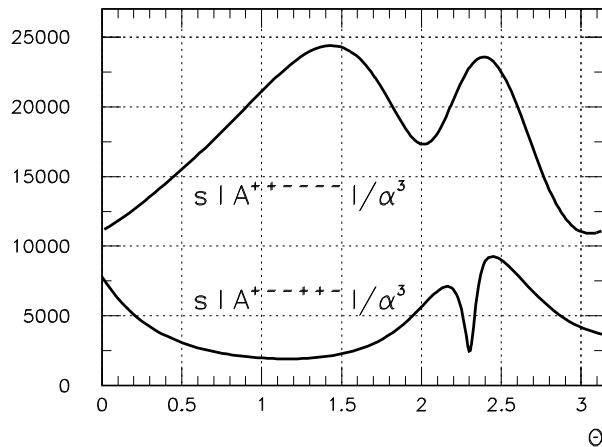
[Ossola, Pittau, Papadopoulos 07]

[Bernicot, Guillet 08]



● rational parts shown to be zero [Binoth, Guillet, GH 06]

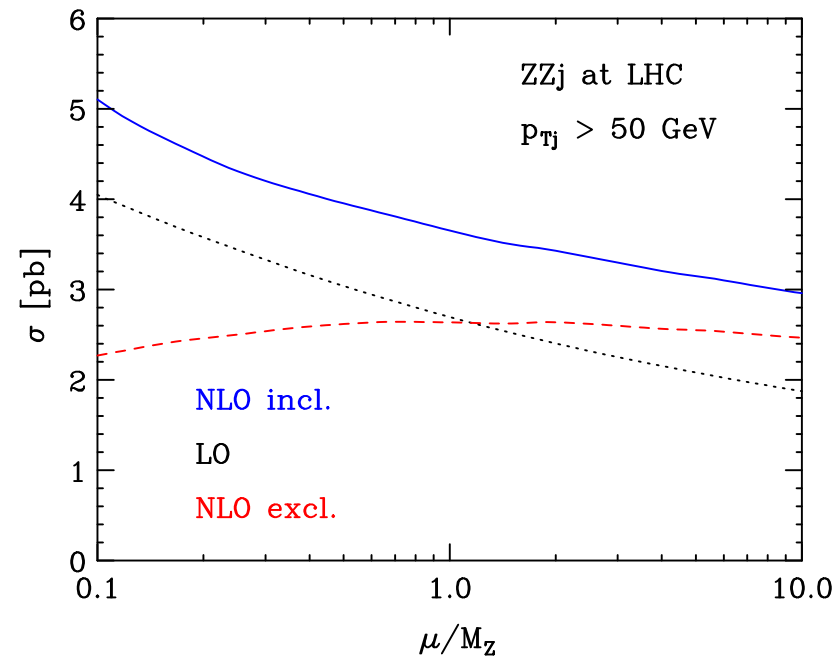
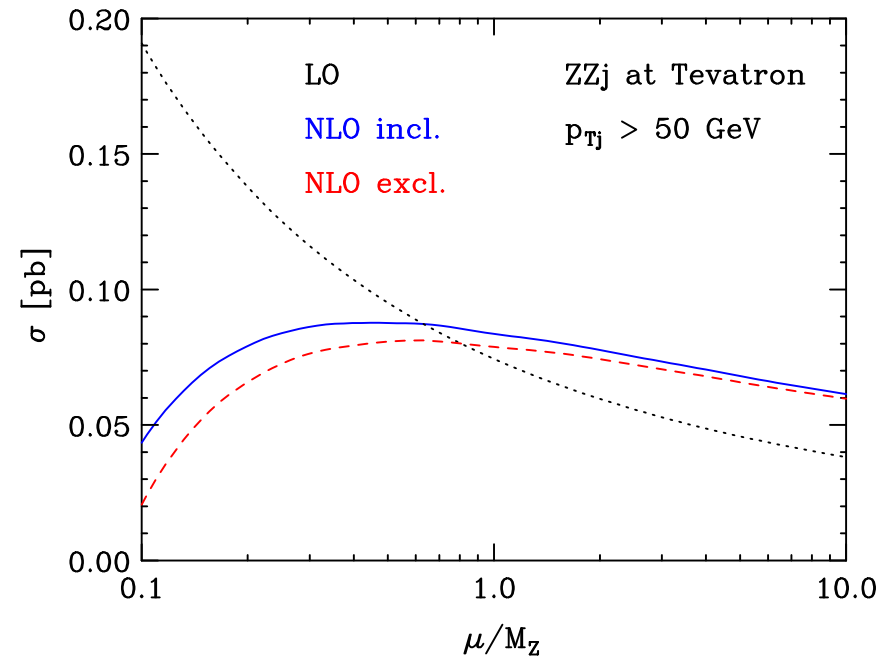
● used both unitarity cuts and Golem



Gong, Nagy, Soper

# ZZ + jet production: scale dependence

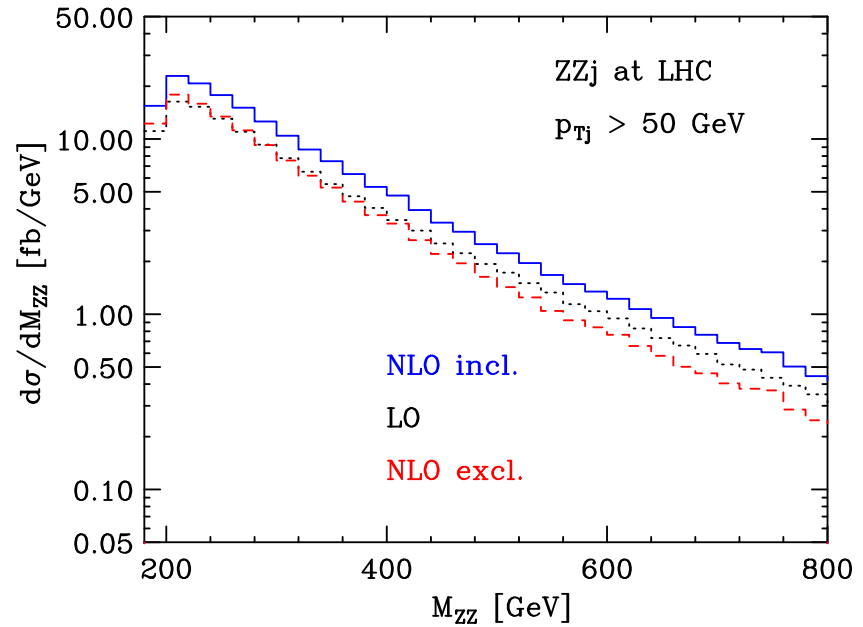
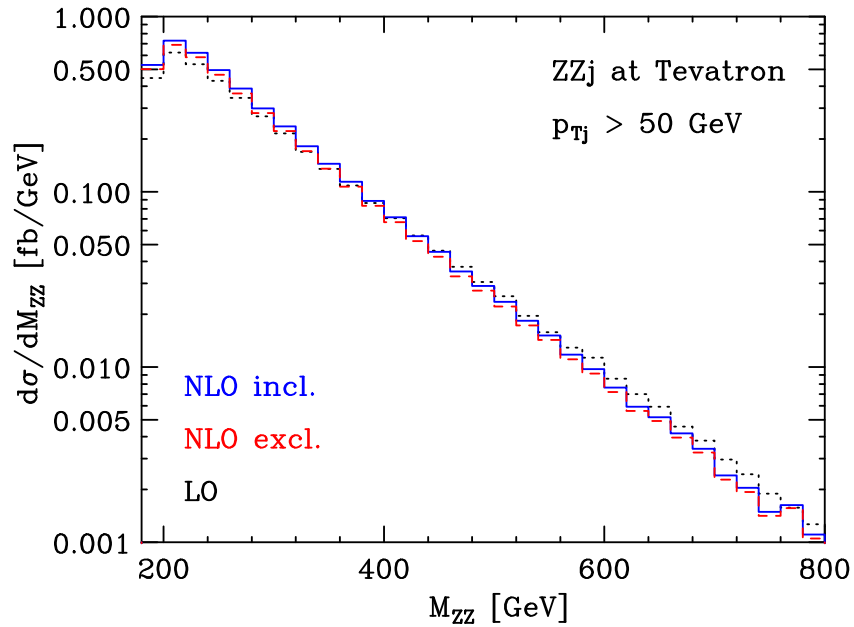
GOLEM collaboration+T. Gleisberg, plots by N. Kauer



**NLO excl.:** jet veto: no additional jets with  $p_T > 50$  GeV

# ZZ + jet production

GOLEM collaboration+T. Gleisberg, plots by N. Kauer



# $pp \rightarrow b\bar{b}b\bar{b}$ one-loop amplitude

$q\bar{q} \rightarrow b\bar{b}b\bar{b}$  completed

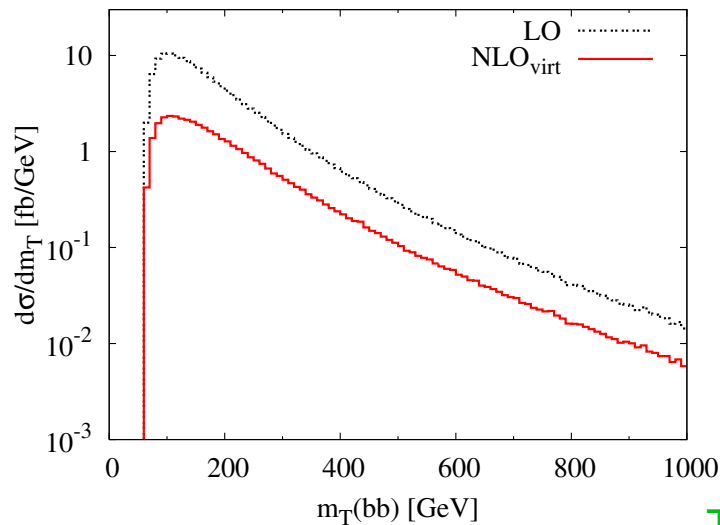
[Ph.D. thesis of Thomas Reiter, Dec. '08]

$gg \rightarrow b\bar{b}b\bar{b}$  virtual part completed

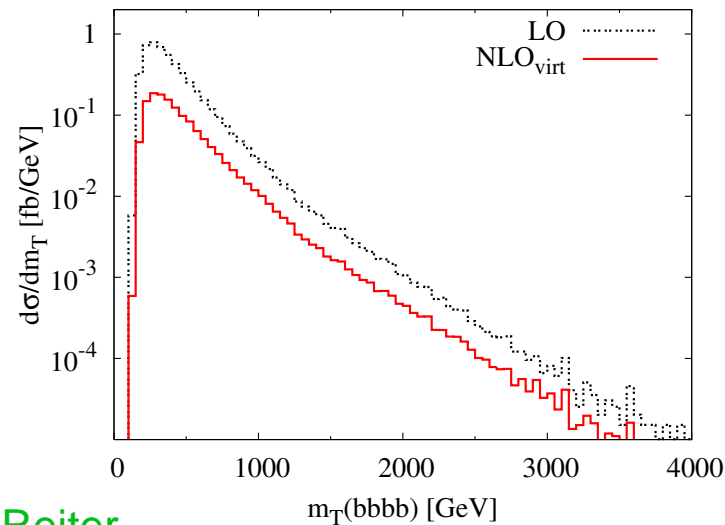
[Binoth, Guillet, Reiter '09]

results for finite combination

$|\mathcal{A}_{\text{LO+NLO}_{\text{virt}}}|^2 - \text{UV counterterms} - \text{IR subtraction terms}$



T.Reiter



# **golem-2.0 demo** (*T. Reiter*)

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# Summary

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- in order to understand "New Physics" at TeV colliders:  $\Rightarrow$  need accuracy beyond LO

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- in order to understand "New Physics" at TeV colliders:  $\Rightarrow$  need accuracy beyond LO
- we are moving towards automated NLO tools

## GOLEM approach:

- setup valid for massive and massless particles
- keeps spin information
- efficient extraction of IR singularities
- numerically robust as inverse Gram determinants can be avoided
- tensor integral library publicly available at <http://lappweb.in2p3.fr/lapth/Golem/golem95.html>
- full package **golem-2.0** available soon



# backup slides

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# asymptotic complexity

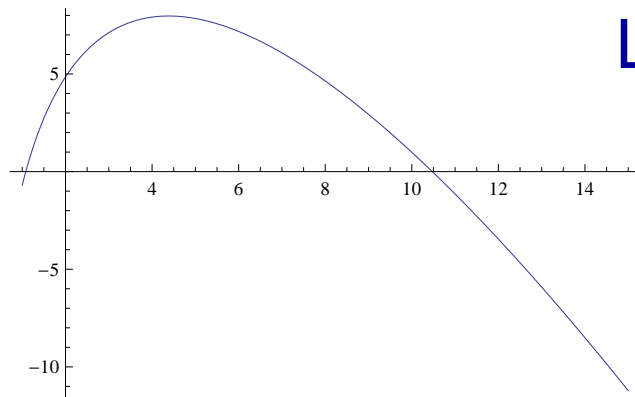
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- unitarity based methods:  
complexity of **colour ordered** amplitudes:

$$\tau_{\text{tree}} \times \tau_{\text{cuts}} \sim N^4 \times \binom{N}{5} \xrightarrow{N \text{ large}} N^9$$

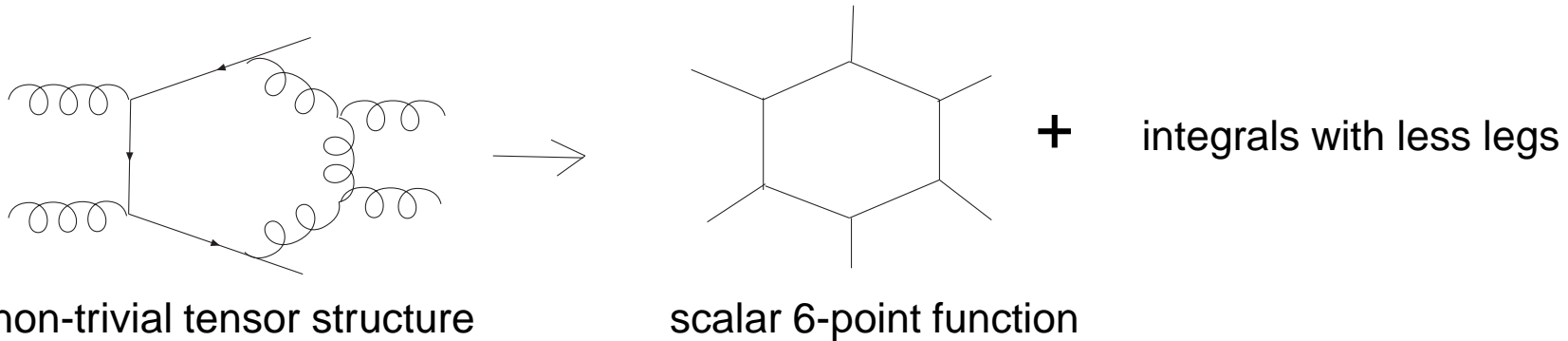
- Feynman diagram reduction:

$$\tau_{\text{diagrams}} \times \tau_{\text{form factors}} \sim 2^N \times \Gamma(N)$$



$$\text{Log}(N^9 / (\Gamma(N) 2^N))$$

# reduction to basis integrals



$$\text{scalar 6-point function} = \sum_{i=1}^6 b_i \text{ (5-point function with leg } i \text{)} \dots \text{ factorial growth in complexity!}$$

reduction to set of **basis integrals** (4-, 3- and 2-point funcs.)

$$\mathcal{A} = C_4 \text{ (4-point function)} + C_3 \text{ (3-point function)} + C_2 \text{ (2-point function)} + \mathcal{R}$$

# form factor representation

---

$$\begin{aligned} I_N^{n, \mu_1 \dots \mu_r}(S) = & \\ & \sum_{l_1 \dots l_r \in S} p_{l_1}^{\mu_1} \dots p_{l_r}^{\mu_r} A_{l_1, \dots, l_r}^{N, r}(S) \\ & + \sum_{l_1 \dots l_{r-2} \in S} \left[ g^{\cdot\cdot} p_{l_1}^{\cdot} \dots p_{l_{r-2}}^{\cdot} \right]^{\{\mu_1 \dots \mu_r\}} B_{l_1, \dots, l_{r-2}}^{N, r}(S) \\ & + \sum_{l_1 \dots l_{r-4} \in S} \left[ g^{\cdot\cdot} g^{\cdot\cdot} p_{l_1}^{\cdot} \dots p_{l_{r-4}}^{\cdot} \right]^{\{\mu_1 \dots \mu_r\}} C_{l_1, \dots, l_{r-4}}^{N, r}(S) \end{aligned}$$

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 \end{aligned}$$

**important:** more than two metric tensors  $g^{\mu\nu}$  **never** occur!

for  $N \geq 6$  : simultaneous reduction of rank  $r$  and number of legs  $N$

$$I_N^{n, \mu_1 \dots \mu_r}(S) = - \sum_{j \in S} C_{j6}^{\mu_1} I_{N-1}^{n, \mu_2 \dots \mu_r}(S \setminus \{j\})$$

$$\mathcal{S}_{ij} = (r_i - r_j)^2 - m_i^2 - m_j^2$$



# Gram determinants

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to avoid spurious  $1/\det(G)$  terms: **do not reduce**

**golem95:**

define dimensionless quantity  $\hat{B} = B \times$  (largest entry of  $S$ )

**if  $\hat{B} < \hat{B}^{\text{cut}}$  : switch to direct numerical evaluation**

**(default:  $\hat{B}^{\text{cut}} = 0.005$ )**

file **demo\_detg.f90** contains example where  $\hat{B} \rightarrow 0$

in rank 3 box integral  $I_4^{n+2}(1, 2, 2; S)$  with two massive legs

# Real part for $B \rightarrow 0$

