## Supersymmetry: Status 2019

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## Supersymmetry can address the following shortcomings of the SM:

- The Hierarchy Problem
- Dark Matter relic density (incl. constraints from spin dependent/independent direct detection experiments)
- Unification of gauge couplings (quarks and leptons fill complete $\operatorname{SU}(5)$ representations, but gauge couplings do not quite unify in the Standard Model without Supersymmetry)
- Possibly:
- The $\sim 3 \sigma$ deviation of the measured anomalous magnetic moment of the muon $a_{\mu}$ w.r.t. the Standard Model,
- galactic center gamma ray excess (dark matter interpretation still viable!),
- AMS-02 anti-proton excess,
- ATLAS 2/3 lepton + ISR excess $(\sim 2 / 3 \sigma)$


## To do

1) Better not rely on interpretations of limits within simplified models (simplified decay cascades, typically 1 step), but recast limits within realistic versions of the MSSM (pMSSM) or NMSSM, preferably including the dark matter relic density
2) Try to fit at least some of the excesses without violating existing constraints
3) Provide experimentalists with new promising search channels

## Recasting limits in the pMSSM:

## GAMBIT collaboration (1705.07917), pMSSM7:

- Gaugino mass ratios motivated by GUT: $M_{3} / \alpha_{s}=M_{2} / \alpha_{2}=M_{1} / \alpha_{1}$
- Degenerate soft squark/slepton masses, but free $A_{t} \neq A_{b}$ (the muon anomalous magnetic moment $a_{\mu}$ cannot be fitted)
- Free soft Higgs masses $M_{H_{u}}, M_{H_{d}}, \tan \beta\left(\rightarrow \mu, B_{\mu}\right.$ fixed by $\left.M_{Z}, \tan \beta\right)$
- $\Omega h^{2} \lesssim 0.1189$ (smaller $\Omega h^{2}$ alleviates constraints from direct DM detection, but requires additional sources of dark matter)

MasterCode (1710.11091), pMSSM11:

- Free gaugino masses
- Different soft squark/slepton masses for the first two/third generations, free $A_{t}=A_{b}$
- Free $\mu, \tan \beta, M_{A}$
- $\Omega h^{2}=0.1186 \pm 0.004$
- With or without fits to the muon anomalous magnetic moment $a_{\mu}$


## Limits on sparticle masses

(Within $2 \sigma$ of the "best fit point" to numerous search signal regions mainly from sparticle searches by ATLAS/CMS)

MasterCode: Some of the parameters of the "best fit points" with or w/o $a_{\mu}$ are completely different!

|  | GAMBIT | MasterCode with $a_{\mu}$ | MasterCode w/o $a_{\mu}$ |
| :---: | :---: | :---: | :---: |
| $M_{\chi_{1}^{\circ}}$ | $\gtrsim 60 \mathrm{GeV}$ (H'ino) | $90-500 \mathrm{GeV}$ (bino) | $>90 \mathrm{GeV}$ (H'ino) |
| $M_{\chi_{1}^{ \pm}}$ | $\gtrsim 90 \mathrm{GeV}$ | $\gtrsim 90 \mathrm{GeV}$ | $\gtrsim 90 \mathrm{GeV}$ |
| $M_{\tilde{\varepsilon}}$ | $\gtrsim 1,0 \mathrm{TeV}$ | $\gtrsim 1,8 \mathrm{TeV}$ | $\gtrsim 1,0 \mathrm{TeV}$ |
| $M_{\tilde{q}}$ | $\gtrsim 1,2 \mathrm{TeV}$ | $\gtrsim 1,9 \mathrm{TeV}$ | $\gtrsim 800 \mathrm{GeV}$ |
| $M_{\tilde{L}}$ | $\gtrsim 0,5 \mathrm{TeV}$ | $\sim 500 \mathrm{GeV} / \gtrsim 1,0 \mathrm{TeV}$ | $\gtrsim 500 \mathrm{GeV}$ |
| $M_{\tilde{\tau}}$ | $\gtrsim 1,3 \mathrm{TeV}$ | $\gtrsim 110 \mathrm{GeV}(\mathrm{LEP})$ | $\gtrsim 110 \mathrm{GeV}(\mathrm{LEP})$ |
| $M_{\tilde{\mu}}$ | $\gtrsim 1,3 \mathrm{TeV}$ | $110-770 \mathrm{GeV}$ | $\gtrsim 110 \mathrm{GeV}$ |
| $M_{A}$ | $\gtrsim 500 \mathrm{GeV}$ | $\gtrsim 800 \mathrm{GeV}$ | $\gtrsim 800 \mathrm{GeV}$ |

No sign for "dark spots" in the combined signal regions (light sparticles escaping detection)
$\rightarrow$ Limits on squarks/gluinos depend strongly on assumptions (possible decay cascades)

## Recast limits in the NMSSM:

- Extra neutral CP-even and CP-odd scalars $H_{S}, A_{S}$ (not degenerate!) on top of the MSSM-like heavy $\sim$ degenerate $\operatorname{SU}(2)$ doublets $H / A$
- Extra singlino $\tilde{S}$ on top of the MSSM-like charged/neutral bino/wino/higgsinos
- $H_{S}, A_{S}, \tilde{S}$ have small couplings to SM particles/MSSM sparticles, except to the Higgs sector from a coupling $\lambda \widetilde{H}_{u} \widetilde{H}_{d} \widetilde{H}_{S}$ in the superpotential (in terms of superfields)
$\rightarrow$ Small direct production cross sections proportional to mixing angles ${ }^{2} \sim \lambda^{2}$, but singlets can be possible decay products of Higgs bosons or sparticles
$\rightarrow$ Still: $H_{S}, A_{S}$ decay into $S M$ particles like $H_{125}$ due to mixing

Searches for $g g F \rightarrow H_{S} \rightarrow \gamma \gamma$ with $M_{H_{s}}<125 \mathrm{GeV}$


From CMS-HIG-17-013 (13 TeV)


Possible Xsect $\times$ BR in the NMSSM for 13 TeV using limits from 8 TeV , from 1512.04281
$\rightarrow$ Sensitivity to viable cross sections $\times \mathrm{BR}$ in the NMSSM!

## Searches for $H_{125} \rightarrow A_{S} A_{S} / H_{S} H_{S}$

Many possible final states, many recent and ongoing searches by ATLAS/CMS

Compilation by R. Aggleton et al., JHEP 1702 (2017) 035:

Light green/blue points: viable in the NMSSM after 2017 LEP/LHC constraints

CMS-PAS-HIG-18-011:


Significant improvement in the $\mu \mu b b$ channel!
$\rightarrow$ Sensitivity to BSM branching fractions of $H_{125}$ allowed by indirect constraints from measured $H_{125}$ couplings!

## If the singlino $\tilde{S}$ is the LSP (I)

A good DM candidate: a relic density $\Omega h^{2} \sim 0.119$ is possible even if $\tilde{S}$ is very light (a few GeV ) through annihilation via $A_{S}$ funnel ( $\neq \mathrm{MSSM}$ )
Coloured region: NMSSM points allowed by constraints from LUX/PandaX-II/Xenon (from 1806.09478 with C. Hugonie)

$\rightarrow$ Xsection possibly below the neutrino floor (black curve)!

## If the singlino $\tilde{S}$ is the LSP (II)

Every NLSP (neutralino, chargino, slepton, stop...) will decay into

$$
N L S P \rightarrow \tilde{S}+H_{125} / H_{S} / A_{S} / Z, W, \text { lepton, top } \ldots,
$$

the only available decay channels due to R -parity conservation


Notably if $\tilde{S}$ is light (a few GeV) AND $M_{N L S P} \approx M_{\tilde{S}}+M_{X}, X=H_{125} / H_{S} / A_{S} / Z$ :
Little energy is given to $\tilde{S}$ in any decay $N L S P_{\text {heavy }} \rightarrow X_{\text {heavy }}+\tilde{S}_{\text {light }}$
$\rightarrow$ Little $E_{T}^{m i s s}$ from $\tilde{S}$ in all Susy searches
$\rightarrow$ Reduced lower limits on sparticle masses (A.Teixeira, U.E., 1406.7221, 1412.6394)

Recast limits from CMS squark search via jets and $E_{T}^{m i s s}$ (1802.02110) (A. Titterton et al., 1807.10672)

MSSM with bino LSP
Assume $\tilde{q} \rightarrow q+$ bino

NMSSM with $\tilde{q} \rightarrow q+$ bino $\rightarrow q+\tilde{S}+H_{125}$

$$
M_{\text {bino }}=M_{\tilde{s}}+M_{H_{125}}+2 \mathrm{GeV}
$$


(red/black curves: expected/observed limits)
$\rightarrow$ strong reduction of the lower limit on $M_{\text {Squark }}$ for small

$$
M_{L S P}=M_{\tilde{S}}!
$$

## Recast searches by ATLAS/CMS for trileptons:

At the LHC, neutralinos/charginos can be produced via $W^{ \pm *} \rightarrow \chi_{i}^{0}+\chi_{j}^{ \pm}$ (or $Z^{*} \rightarrow \chi_{i}^{ \pm}+\chi_{j}^{\mp}, \chi_{i}^{0}+\chi_{j}^{0}$ ):


Results are typically interpreted for wino-like $\chi_{2}^{0}+\chi_{1}^{ \pm}$: Largest cross sections $\rightarrow$ strongest constraints But: Higgsinos have only half the cross section (even adding $\chi_{2}^{0}, \chi_{3}^{0}$ ) $\rightarrow$ weaker constraints

For limits on the NMSSM singlino-higgsino sector (with C. Hugonie, 1806.09478): Scan the parameter space with singlino LSP, require a viable relic density consistent with constraints from direct DM detection, apply bounds from the CMS trilepton search in 1801.03957 (the strongest ones)

## Comparison of limits the in the $M_{\chi_{1}^{0}} / M_{\chi_{1}^{ \pm}} \sim M_{\chi_{2}^{0}}$ plane:

CMS, assuming wino-like $\chi_{2}^{0}$ and $\chi_{1}^{ \pm}$:
NMSSM, singlino LSP and higgsino-like $\chi_{2}^{0}, \chi_{3}^{0}$ and $\chi_{1}^{ \pm}$, bino-like $\chi_{4}^{0}$ :



Red: Excluded by constraints on DM and by CMS

Blue: Excluded iff the bino mass satisfies $M_{1}>300 \mathrm{GeV}$ as motivated by the GUT relation $M_{1} \approx M_{\text {Gluino }} / 6$ and $M_{\text {Gluino }} \gtrsim 1.8 \mathrm{TeV} \rightarrow$ no bino/higgsino mixing
$\rightarrow$ Substantial reduction of limits!

Allowed regions in the plane $M_{\tilde{\chi}_{1}^{ \pm}}-M_{\tilde{\chi}_{1}^{0}}$ in the constrained NMSSM: universal soft susy breaking terms at the GUT scale, but non-universal soft Higgs mass terms (allows to estimate the necessary amount of finetuning):

$\rightarrow$ Relatively low finetuning for $M_{\tilde{\chi}_{1}^{0}} \sim M_{z} / 2, M_{\tilde{\chi}_{1}^{0}} \sim M_{H 125} / 2$ or $M_{\tilde{\chi}_{1}^{0}} \sim M_{\tilde{\chi}_{1}^{ \pm}}$ where s-channel annihilation or co-annihilation is possible Otherwise: s-channel annihilation via $A_{S}$ with $M_{\tilde{\chi}_{1}^{0}} \sim M_{A_{s}} / 2$ $\rightarrow$ Many regions with relatively low fine-tuning $\approx 100$ remain to be tested

## Dark Spots for neutralino/chargino searches:

- Mixed bino - higgsino NLSP $\chi_{2}^{0}$ : reduces production cross section further
- $\chi_{2,3}^{0}$ cascade decays via light $H_{S}$ or $A_{S}$ (escape searches for $H_{125}$ via $b \bar{b}$ )
- Light staus $\tilde{\tau}$ as NLSP: Hardly constrained by the LHC (limits from LEP), $\rightarrow$ less "Trileptons" in the final state


## Attempts to fit excesses

Searches for neutralinos/charginos by ATLAS using recursive jigsaw reconstruction (1806.02293), ISR jet, sensitive to small $\chi_{2}^{0}-\chi_{1}^{0}$ mass differences:


Local $2-3 \sigma$ excesses in the signal regions $\mathrm{SR} 2 \ell_{\text {low }}+\mathrm{ISR}$ and $\mathrm{SR} 3 \ell_{\text {low }}+\mathrm{ISR}$
If interpreted in terms of simplified models:
No significant deviations from observed w.r.t. expected limits

Combine 4 ATLAS and 4 CMS electroweakino searches after $39 \mathrm{fb}^{-1}(\approx 10$ signal regions each, up to $\sim 40$ bins), simulations within a pMSSM electroweakino sector (bino, wino, higgsinos), allowing for cascade decays
$\rightarrow$ local $3,2 \sigma$ excess for $M_{\chi_{1}^{0}} \sim 50 \mathrm{GeV}, M_{\chi_{1}^{ \pm}} \sim 150 \mathrm{GeV}$ via contributions from $\chi_{2}^{0}, \chi_{3}^{0}$ and $\chi_{2}^{ \pm}$multi- $W / Z$ cascade decays ( $M_{\chi_{1}^{0}} \sim 8-155 \mathrm{GeV}, M_{\chi_{1}^{ \pm}} \sim 104-259 \mathrm{GeV}$ within $95 \% \mathrm{CL}$ )
(Missing covariance matrices for stat. analysis including more search results)

> M. Carena et al., 1809.11082:

Require a viable dark matter relic density and a fit of $a_{\mu}$ :
Bino-like $\chi_{1}^{0}$, resonant pair annihilation via $H_{S M}$ funnel (requires some higgsino component for $H_{S M}-\chi_{1}^{0}-\chi_{1}^{0}$ coupling) Large $\tan \beta$ to suppress $H_{S M} \rightarrow \chi_{1}^{0} \chi_{1}^{0}$ decay

Suppress spin independent direct detection cross section via $\mu \cdot M_{1}<0$ But: Assume $\mu \cdot M_{2}>0$ (and large $\tan \beta, M_{\tilde{\nu}_{\mu}}<400 \mathrm{GeV}$ ) to fit $a_{\mu}$
M. Carena et al., 1905.03768:

Include fit to galactic center gamma ray excess: Need $\chi_{1}^{0} \chi_{1}^{0}$ annihilation (Higgs funnel) via s-wave; then: simultaneous explanation of AMS-02 anti-proton excess from $\chi_{1}^{0} \chi_{1}^{0} \rightarrow b \bar{b}$.
MSSM: $M_{1}$ complex $\rightarrow$ CP-violating $H_{S M}-\chi_{1}^{0}-\chi_{1}^{0}$ coupling allows for annihilation via s-wave; but: constraints from electric dipole moments require heavy sleptons $\rightarrow$ fit of $a_{\mu}$ impossible
NMSSM: Relic density from $\chi_{1}^{0}$ pair annihilation with a pseudoscalar in the s-channel
$\rightarrow$ NMSSM benchmark point with common fit of ATLAS $3 \ell_{\text {low }}+$ ISR excess, $a_{\mu}$, galactic center and AMS-02 anti-proton excesses!

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- Due to $M_{\text {Higgs }} \sim 125 \mathrm{GeV}$ the MSSM has a "little" finetuning problem of at least $\mathcal{O}(1 \%)$, of $\mathcal{O}(1 \%$ ) with (grand) unified soft Supersymmetry breaking terms, somewhat less in the NMSSM
- Of course: even with $M_{\text {Squark }}, M_{G / u i n o}>1-2$ TeV Supersymmetry still solves the "BIG" hierarchy problem
- To derive definite constraints on the high dimensional parameter space is a challenging task, notably in the NMSSM ( $\rightarrow$ dark spots), but a MUST for the future
- The dark matter relic density and some (mild) excesses in particle/astroparticle physics can be explained with still viable parameters in the MSSM, notably the NMSSM

