

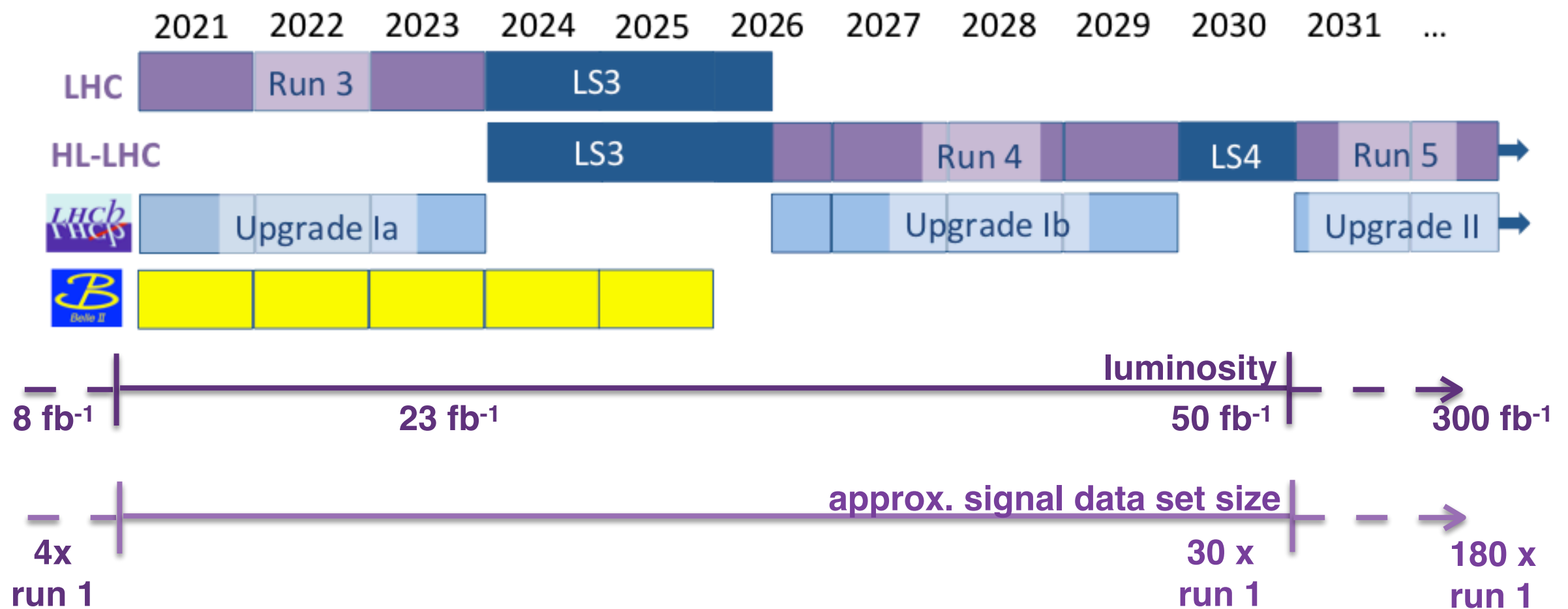
# Experimental prospects for baryon decays

Lyon, 6th September 2019

T. Blake



# LHCb Upgrades



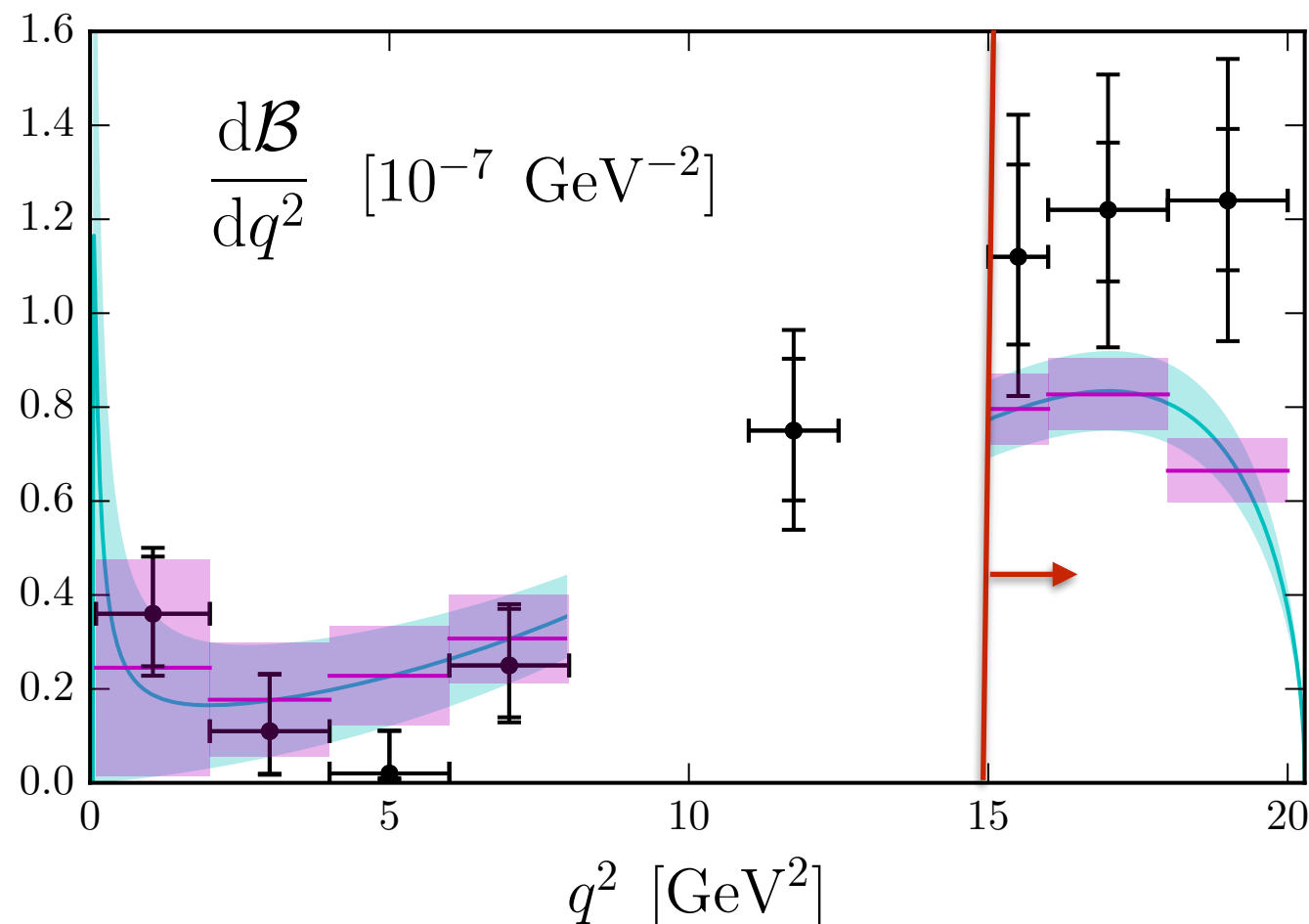
Expect a further factor of ~2 gain for channels with electrons/photons due to removal of the L0 hardware trigger.

- Focus on LHCb, but ATLAS and CMS can contribute to a number of the measurements mentioned in these slides (see talk by Greg).

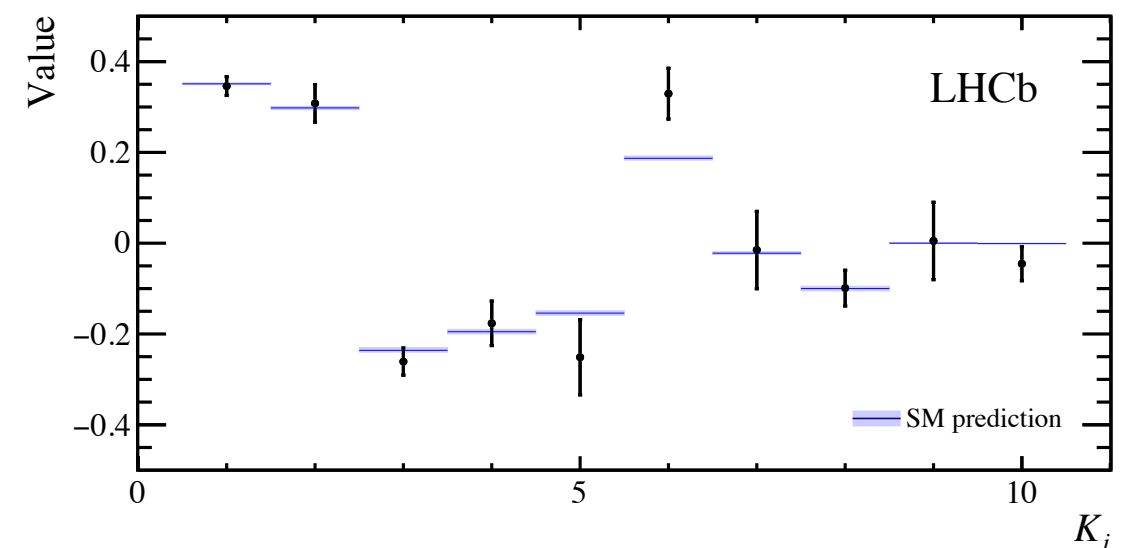
$$\Lambda_b \rightarrow \Lambda \mu^+ \mu^-$$

- Large number (34) of angular observables, 24 require  $\Lambda_b$  to be polarised at production and are consistent with zero in current dataset.

Detmold et al. arXiv:1602.01399



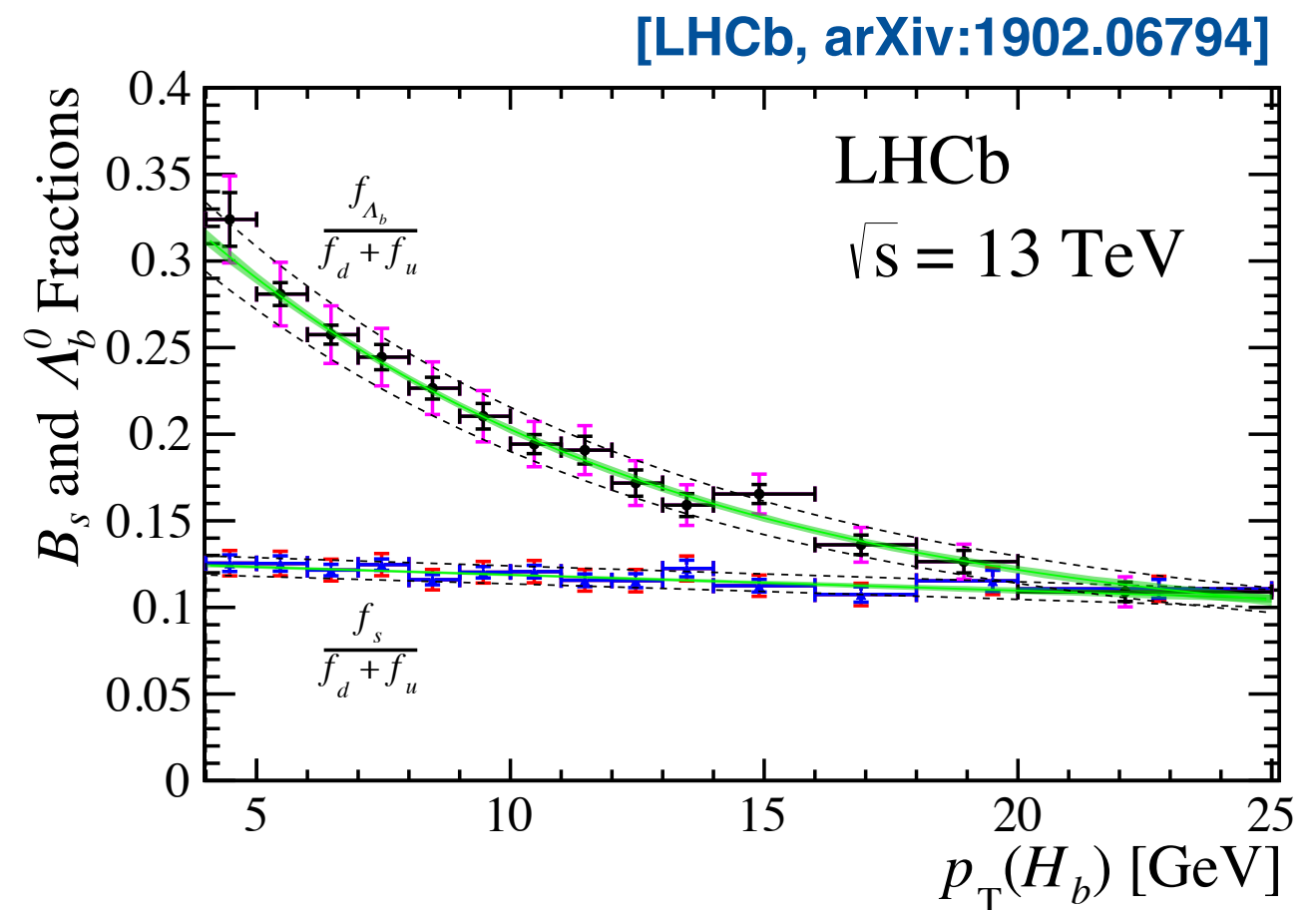
LHCb, arXiv:1808.00264



- SM predictions use external on  $\alpha_\Lambda$  and production polarisation (small at the LHC).
- Branching fraction uncertainty currently dominated by knowledge of  $\mathcal{B}(\Lambda_b \rightarrow J/\psi \Lambda)$ .

# $\Lambda_b \rightarrow \Lambda \mu^+ \mu^-$ branching fraction

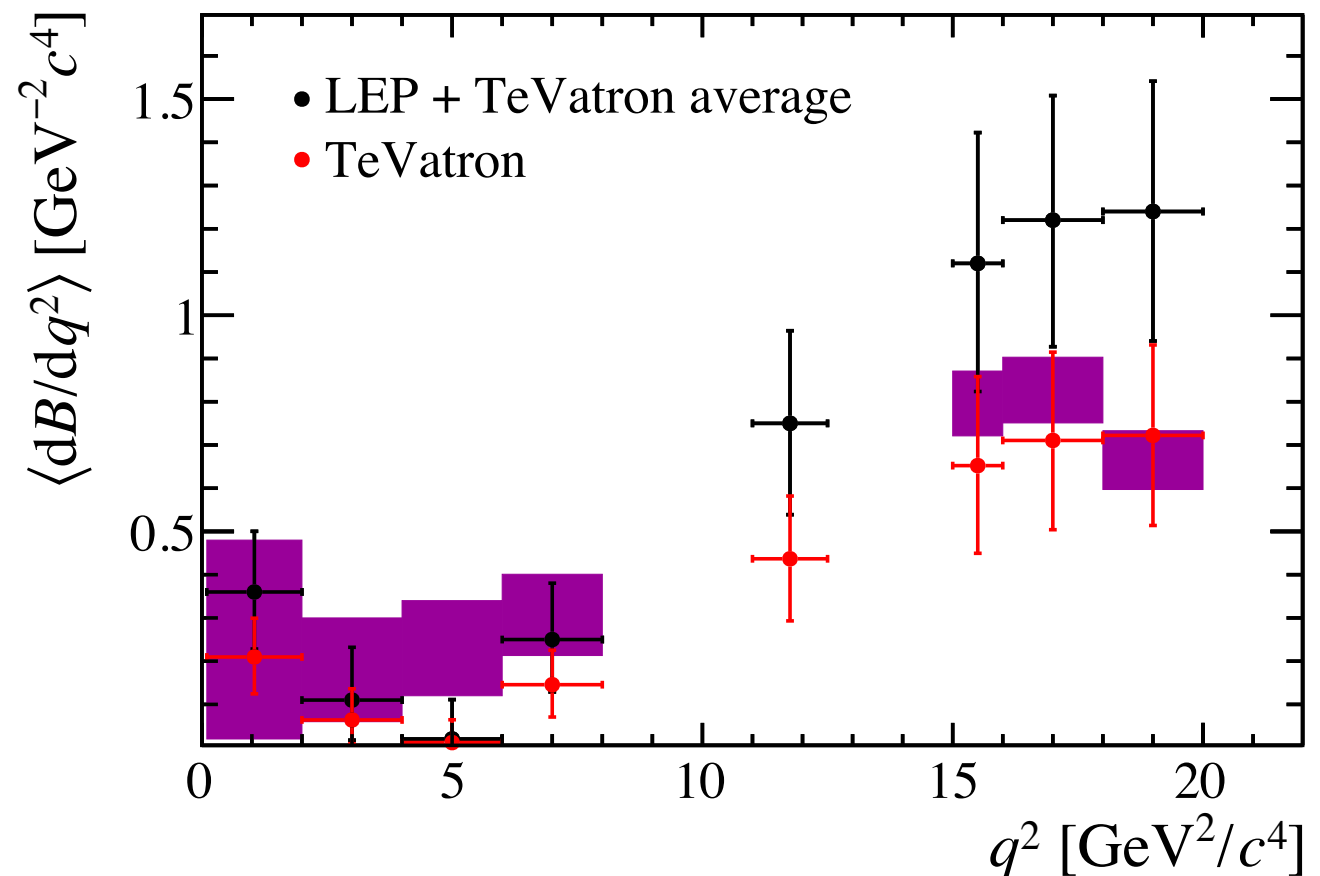
- Existing branching fraction measurement uses CDF/D0 average of measurement of  $f_{\Lambda_b} \times \mathcal{B}(\Lambda_b \rightarrow J/\psi \Lambda) = (5.8 \pm 0.8) \times 10^{-5}$ , with  $f_{\Lambda_b}$  taken as a LEP + TeVatron average.
- Now know that the baryon production fractions exhibit strong  $p_T$  dependences in  $pp$  collisions.
- $\Lambda_b$  baryons produced with lower average  $p_T$  at the TeVatron than LEP.





# $\Lambda_b \rightarrow \Lambda \mu^+ \mu^-$ branching fraction

- Re-evaluating the branching fraction using only TeVatron inputs significantly changes the picture.
- Data consistent with, but now below SM predictions.
- ➔ Consistent with pattern seen in other branching fraction measurements.



**Data [LHCb, JHEP 06 (2015) 115]**  
**SM [Detmold et al. arXiv:1602.01399]**

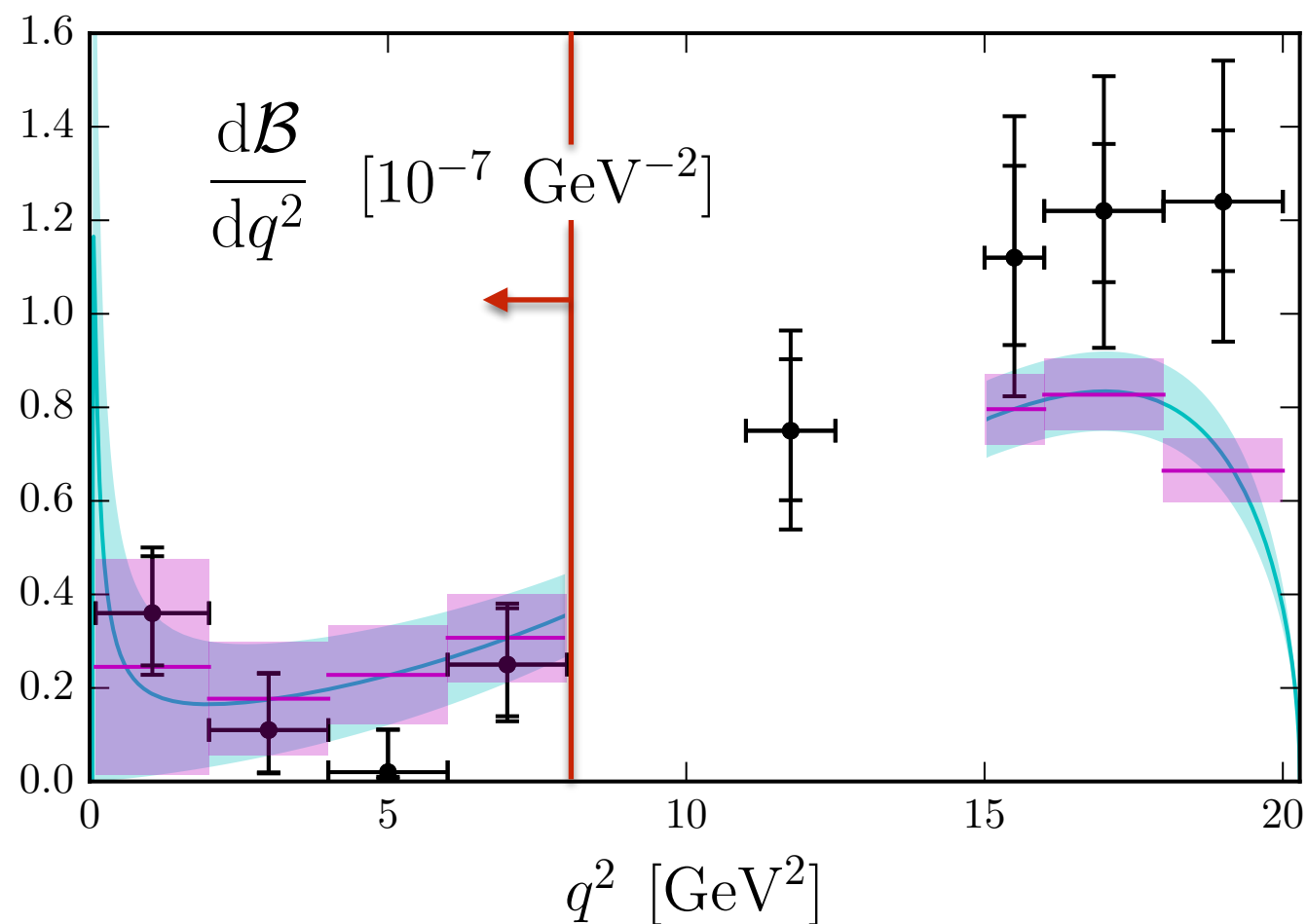
# $\Lambda_b \rightarrow \Lambda \mu^+ \mu^-$ branching fraction

- Plan to measure  $\mathcal{B}(\Lambda_b \rightarrow J/\psi \Lambda)$  at LHCb to improve the normalisation uncertainty.
- Ultimate precision on branching fraction measurements is limited by knowledge of  $\mathcal{B}(B \rightarrow J/\psi K)$  and  $f_{\Lambda_b}/f_d$ .
  - ➔ Bin-by-bin measurements will be systematically limited with an 8% ( $4\% \oplus 7\%$ ) systematic uncertainty.

$$\Lambda_b \rightarrow \Lambda \mu^+ \mu^-$$

- We should observe significant signal at low  $q^2$  in the run 2 data set.

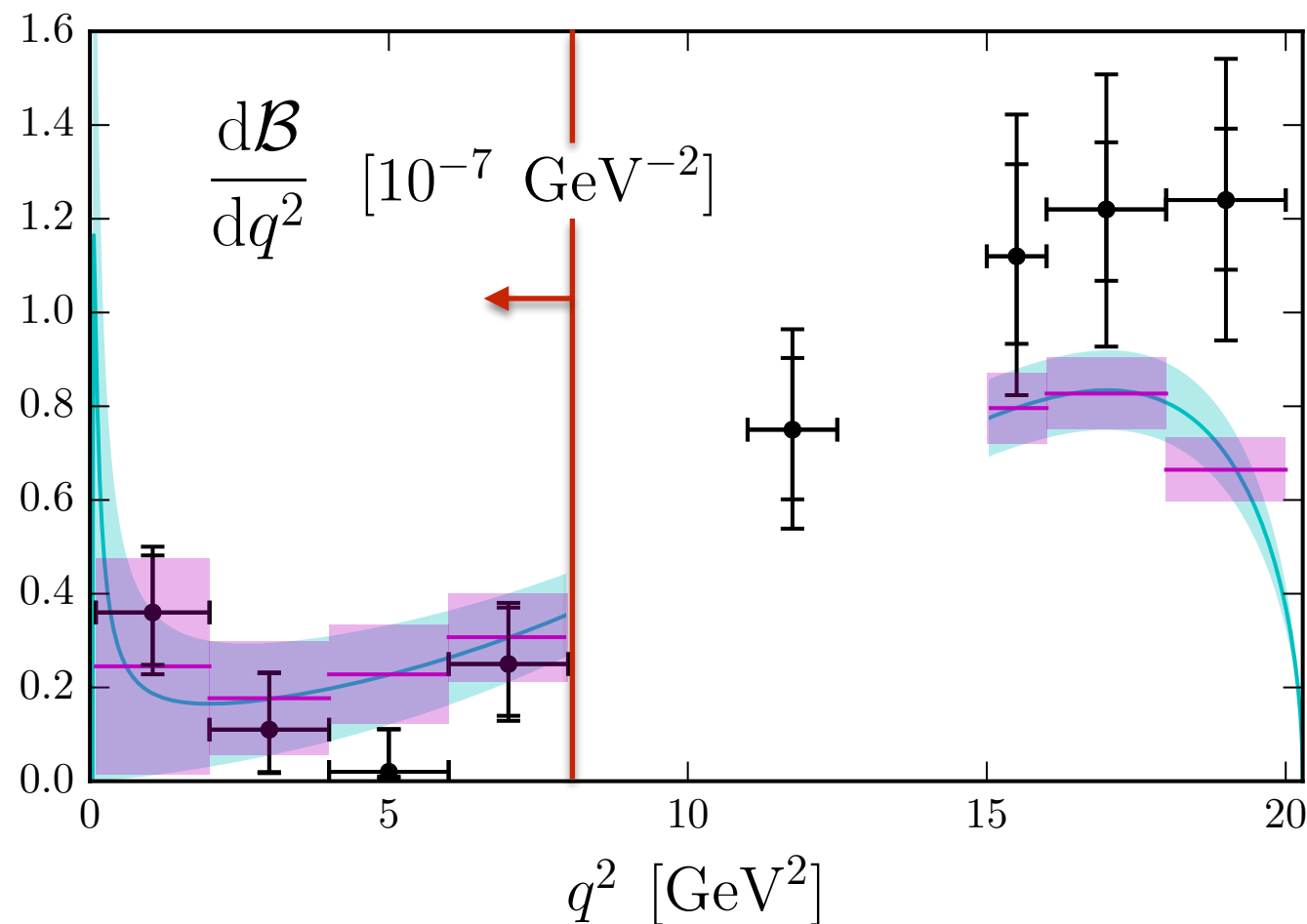
Detmold et al. arXiv:1602.01399



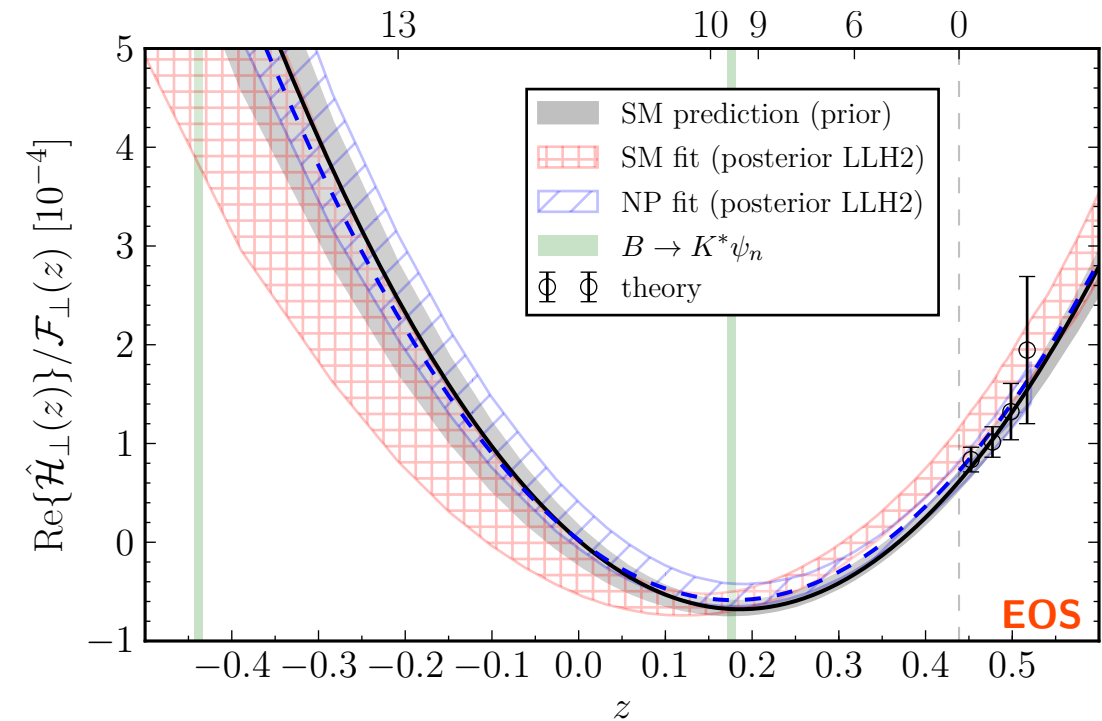
$$\Lambda_b \rightarrow \Lambda \mu^+ \mu^-$$

- We should observe significant signal at low  $q^2$  in the run 2 data set.

Detmold et al. arXiv:1602.01399



[Bobeth et al. EPJC 78 (2018) 451]



- Could also apply techniques from [EPJC 78 (2018) 451] to low  $q^2$  region.
- External input wanted on decay amplitudes for  $J/\psi \Lambda$  and  $\psi(2S) \Lambda$ .

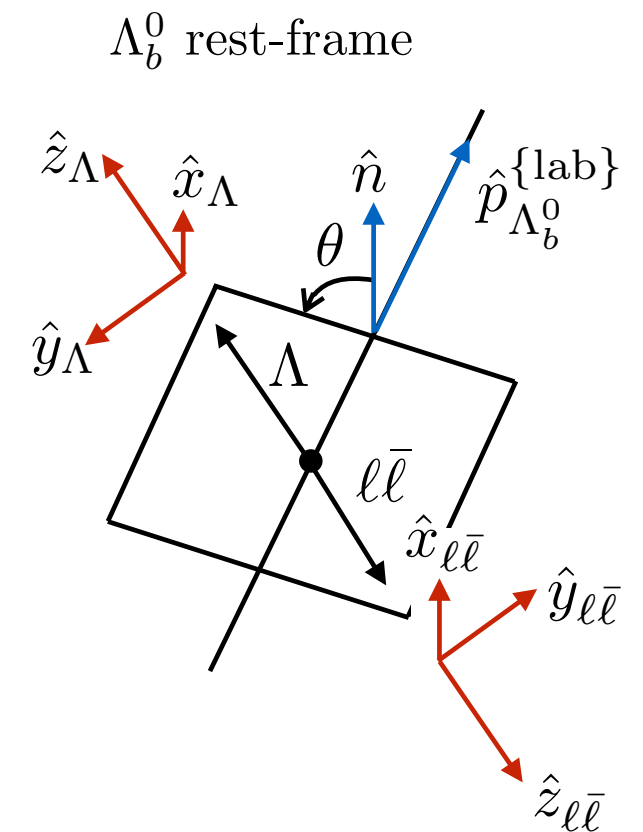
# $\Lambda_b \rightarrow J/\psi \Lambda$ angular distribution

- Parameterised by five decay angles:
  - ➔ Polar and azimuthal helicity angles for the  $\Lambda$  and  $J/\psi$  system and  $\theta$  defined by  $\hat{n}$ , where  $\hat{n} = \hat{p}_{\text{beam}} \times \hat{p}_{\Lambda_b}$ .
- Described by 4 helicity amplitudes.  $H_{\lambda_\Lambda, \lambda_\psi}$ , the  $\Lambda_b$  production polarisation and the  $\Lambda$  asymmetry parameter  $\alpha$ .
  - ➔ Amplitudes are
 
$$a_- = H_{-1/2, 0}$$

$$a_+ = H_{+1/2, 0}$$

$$b_- = H_{+1/2, +1}$$

$$b_+ = H_{-1/2, -1}$$



See, e.g.  
 T. Blake & M. Kreps arXiv:1710.00746,  
 J. Hrivnac et al. arXiv:hep-ph/9405231

# Observables

- Large number of observables vanish if the  $\Lambda_b$  is unpolarised:
- ✗ No longer have enough constraints to determine phases of the amplitudes.
- Even if the polarisation is large expect two amplitudes to be small:  $a_+ \approx b_- \approx 0$ .

Moments	Amplitude dependence
$M_1$	$\frac{1}{4}(2 a_+ ^2 + 2 a_- ^2 +  b_+ ^2 +  b_- ^2)$
$M_2$	$\frac{1}{2}( b_+ ^2 +  b_- ^2)$
$M_4$	$\frac{\alpha}{4}( b_- ^2 -  b_+ ^2 + 2 a_+ ^2 - 2 a_- ^2)$
$M_5$	$\frac{\alpha}{2}( b_- ^2 -  b_+ ^2)$
$M_7$	$\frac{\alpha}{\sqrt{2}}\text{Re}(-b_+^*a_+ + b_-a_-^*)$
$M_9$	$\frac{\alpha}{\sqrt{2}}\text{Im}(b_+^*a_+ - b_-a_-^*)$
$M_{11}$	$P_b\frac{1}{4}( b_+ ^2 -  b_- ^2 + 2 a_+ ^2 - 2 a_- ^2)$
$M_{12}$	$P_b\frac{1}{2}( b_+ ^2 -  b_- ^2)$
$M_{14}$	$P_b\frac{\alpha}{4}(- b_- ^2 -  b_+ ^2 + 2 a_+ ^2 + 2 a_- ^2)$
$M_{15}$	$-P_b\frac{\alpha}{2}( b_+ ^2 +  b_- ^2)$
$M_{17}$	$-P_b\frac{\alpha}{\sqrt{2}}\text{Re}(b_+^*a_+ + b_-a_-^*)$
$M_{19}$	$P_b\frac{\alpha}{\sqrt{2}}\text{Im}(b_+^*a_+ + b_-a_-^*)$
$M_{21}$	$-P_b\frac{1}{\sqrt{2}}\text{Im}(b_+^*a_- - b_-a_+^*)$
$M_{23}$	$P_b\frac{1}{\sqrt{2}}\text{Re}(b_+^*a_- - b_-a_+^*)$
$M_{25}$	$P_b\frac{\alpha}{\sqrt{2}}\text{Im}(b_+^*a_- + b_-a_+^*)$
$M_{27}$	$-P_b\frac{\alpha}{\sqrt{2}}\text{Re}(b_+^*a_- + b_-a_+^*)$
$M_{30}$	$P_b\alpha\text{Im}(a_+a_-^*)$
$M_{32}$	$-P_b\alpha\text{Re}(a_+a_-^*)$
$M_{33}$	$-P_b\frac{\alpha}{2}\text{Re}(b_+^*b_-)$
$M_{34}$	$P_b\frac{\alpha}{2}\text{Im}(b_+^*b_-)$

# $\Lambda$ asymmetry parameter

- Recent measurement by BESIII [[Nature Physics 15 \(2019\) 631–634](#)] is 17% larger than current world average value:

$$\alpha_{\Lambda} = 0.642 \pm 0.013 \quad \text{PDG}$$

$$\alpha_{\Lambda} = 0.750 \pm 0.010 \quad \text{BESIII}$$

- The larger BESIII value likely solves the problems with the existing LHCb, ATLAS and CMS analyses of  $\Lambda_b \rightarrow J/\psi \Lambda$ , which favour an unphysical solution [[LHCb, PLB 724 \(2013\) 27](#)][[ATLAS, PRD 89 \(2014\) 092009](#)][[CMS, PRD 97 \(2018\) 072010](#)].
- Old measurements of  $\alpha$  had to determine the proton polarisation from secondary scattering.
- Impacts interpretation of the  $\Lambda_b \rightarrow \Lambda \mu^+ \mu^-$  angular observables.

# Updated global fit

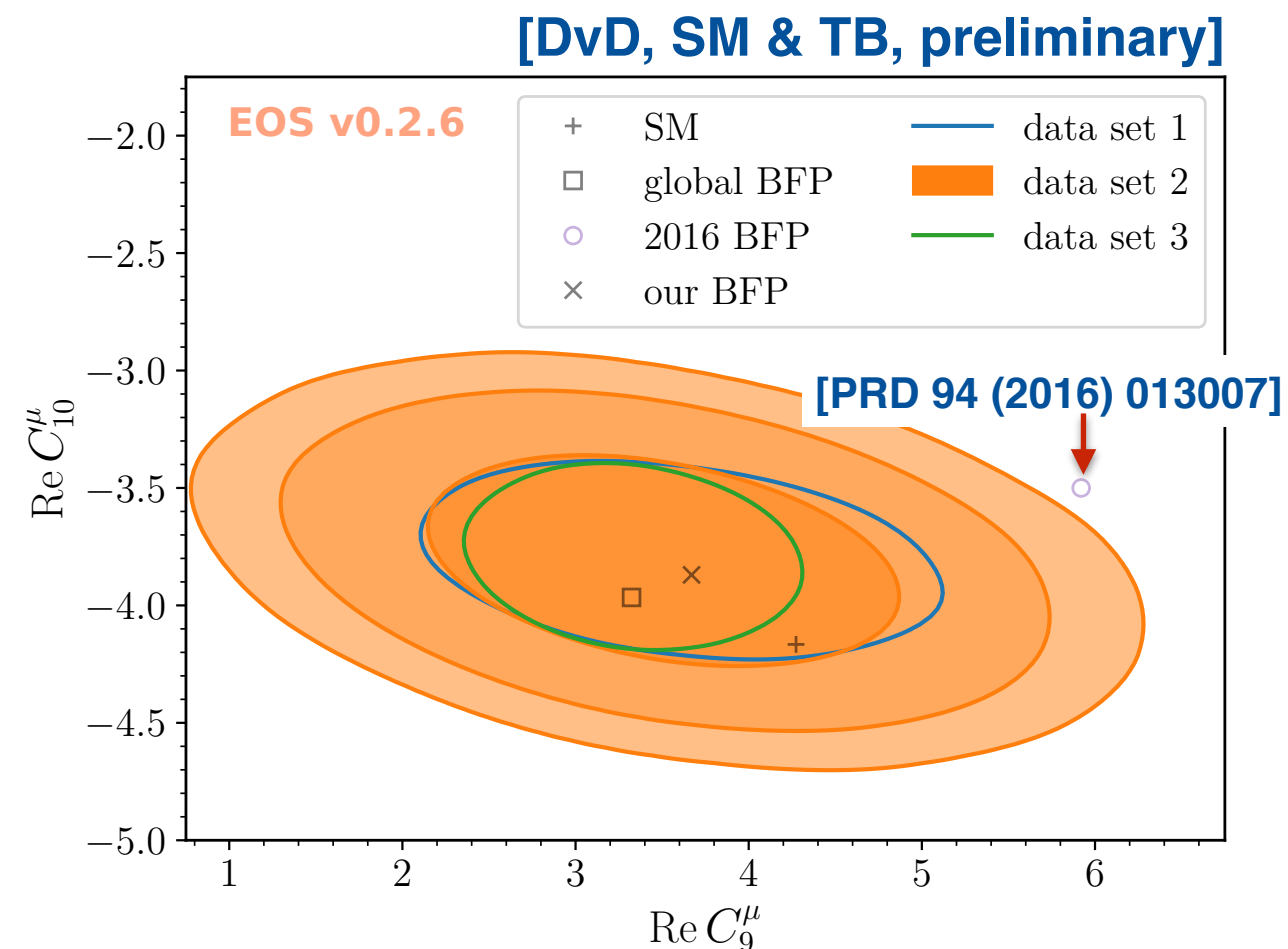
- Try three scenarios:

ATLAS, CMS & LHCb  $\mathcal{B}(B_s \rightarrow \mu^+ \mu^-)$  + unpolarised  $\Lambda_{\mu^+ \mu^-}$  angular observables.

+ Polarised angular observables.

+ Updated branching fraction.

- SM point has good p-value.
- Data are consistent with the anomalies (best-fit point is close to the one obtained using meson data).



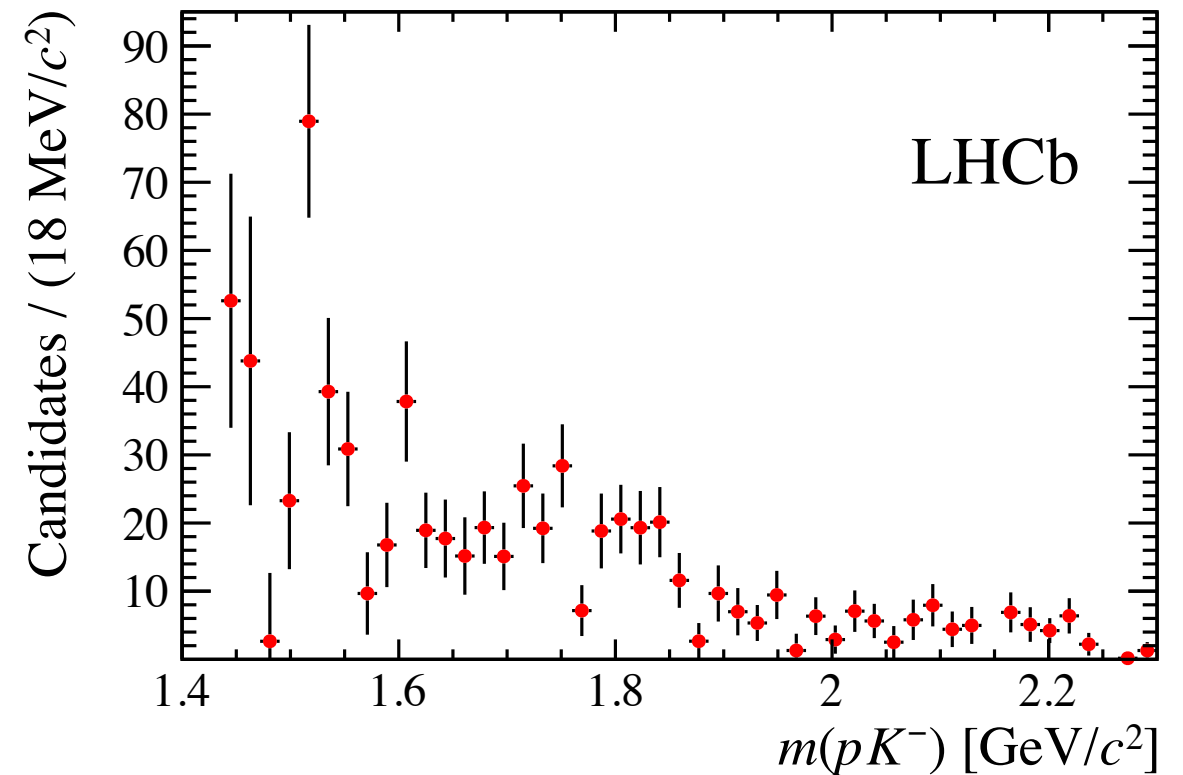
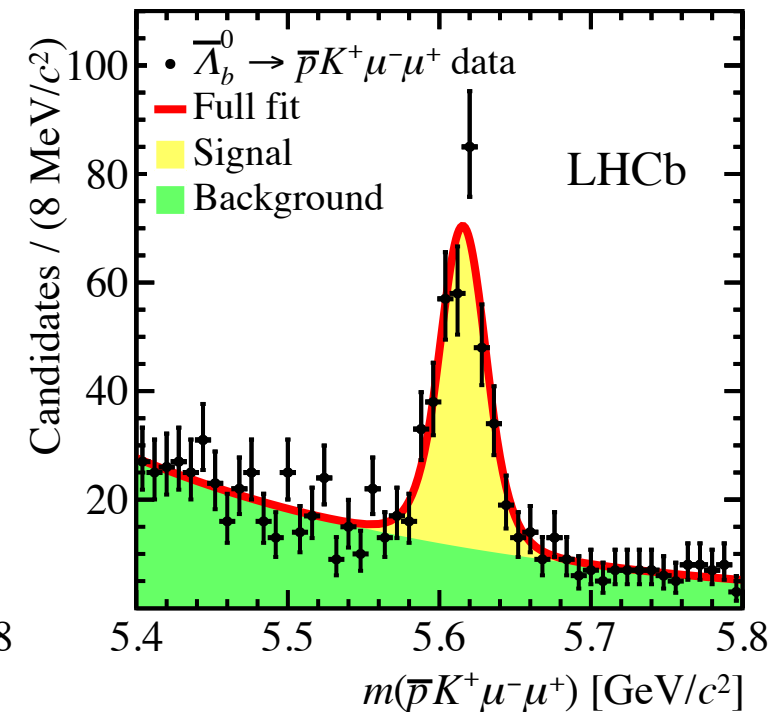
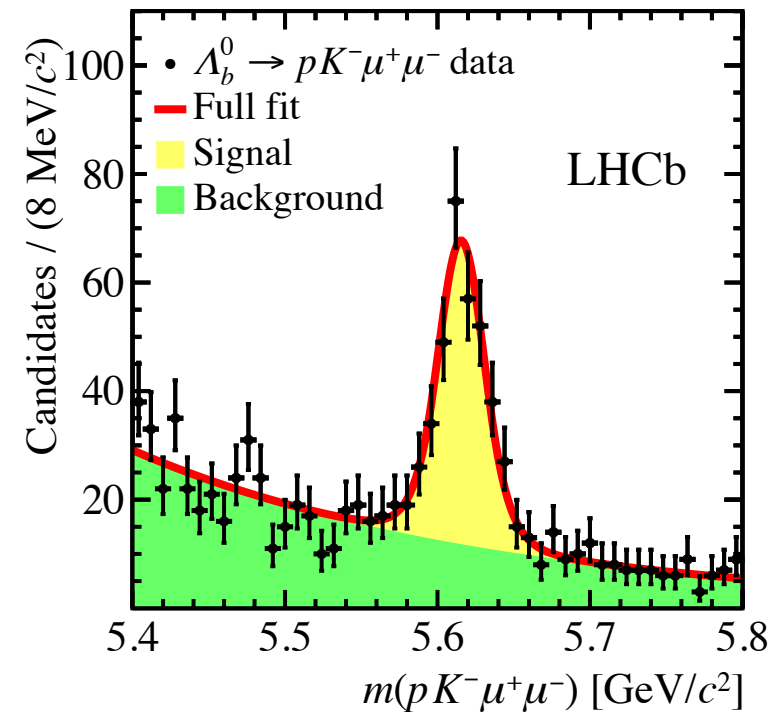


# Polarised $\Lambda_b$ baryons?

- Large number of new observables would be available if the  $\Lambda_b$  baryons were produced polarised, see [\[T. Blake & M. Kreps JHEP 11 \(2017\) 138\]](#).
- Polarisation is small at LHC, but is large in  $Z^0$  decays [\[ALEPH, PLB 365 \(1996\) 437-447\]](#), [\[OPAL, PLB 444 \(1998\) 539-554\]](#), [\[DELPHI, PLB 474 \(2000\) 205-222\]](#)
- Long-term, could perform these measurements at a future  $e^+e^-$  collider, e.g. expect  $5 \times 10^{12}$   $Z^0$ 's at FCC-ee [\[A. Abada et al, EPJC 79 \(2019\) 474\]](#).
  - ➔  $O(1000)$  reconstructible  $\Lambda_b \rightarrow \Lambda \mu^+ \mu^-$  decays at low-recoil.
- Could we also exploit the large  $t\bar{t}$  cross-section at ATLAS and CMS as a source of polarised  $b$ -baryons?

$$\Lambda_b \rightarrow p K \mu^+ \mu^-$$

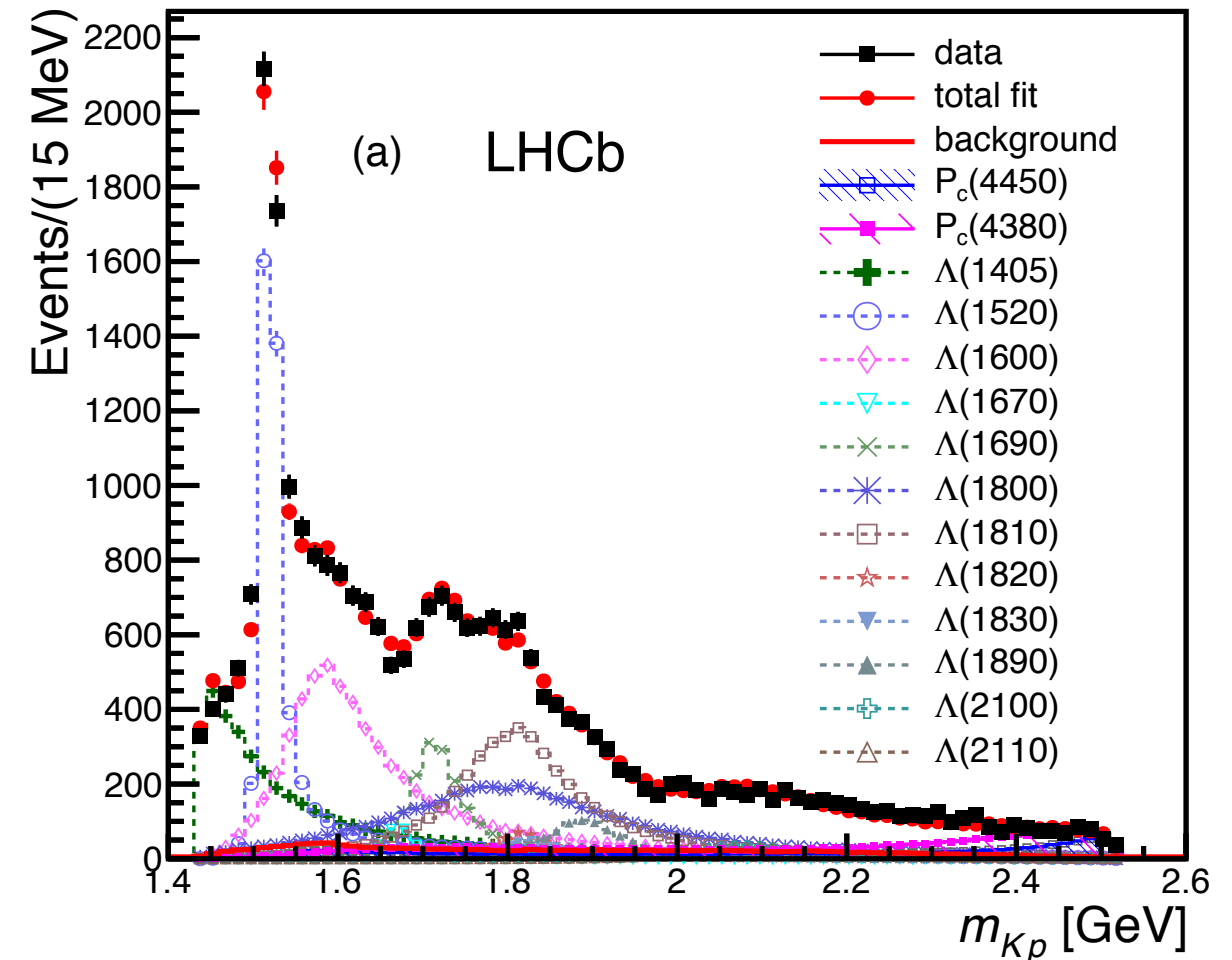
- First observed in the LHCb Run 1 data set.
- Measured CP and triple product asymmetries (null result).
- Branching fraction measurement and angular analysis complicated by the presence of a large number of overlapping  $\Lambda^*$  resonances with different  $J^{PC}$ .



$$\Lambda_b \rightarrow J/\psi p K$$

[LHCb, PRL 115 (2015) 072001]

- An amplitude analysis of  $J/\psi p K$  was carried out as part of the original pentaquark observation paper.  
[PRL 115 (2015) 072001]
- Data dominated by  $\Lambda(1520)$  (with  $J^{PC} = 3/2^-$ ) at low  $pK$  mass. However, there are still large contributions from other resonances that are difficult to disentangle.

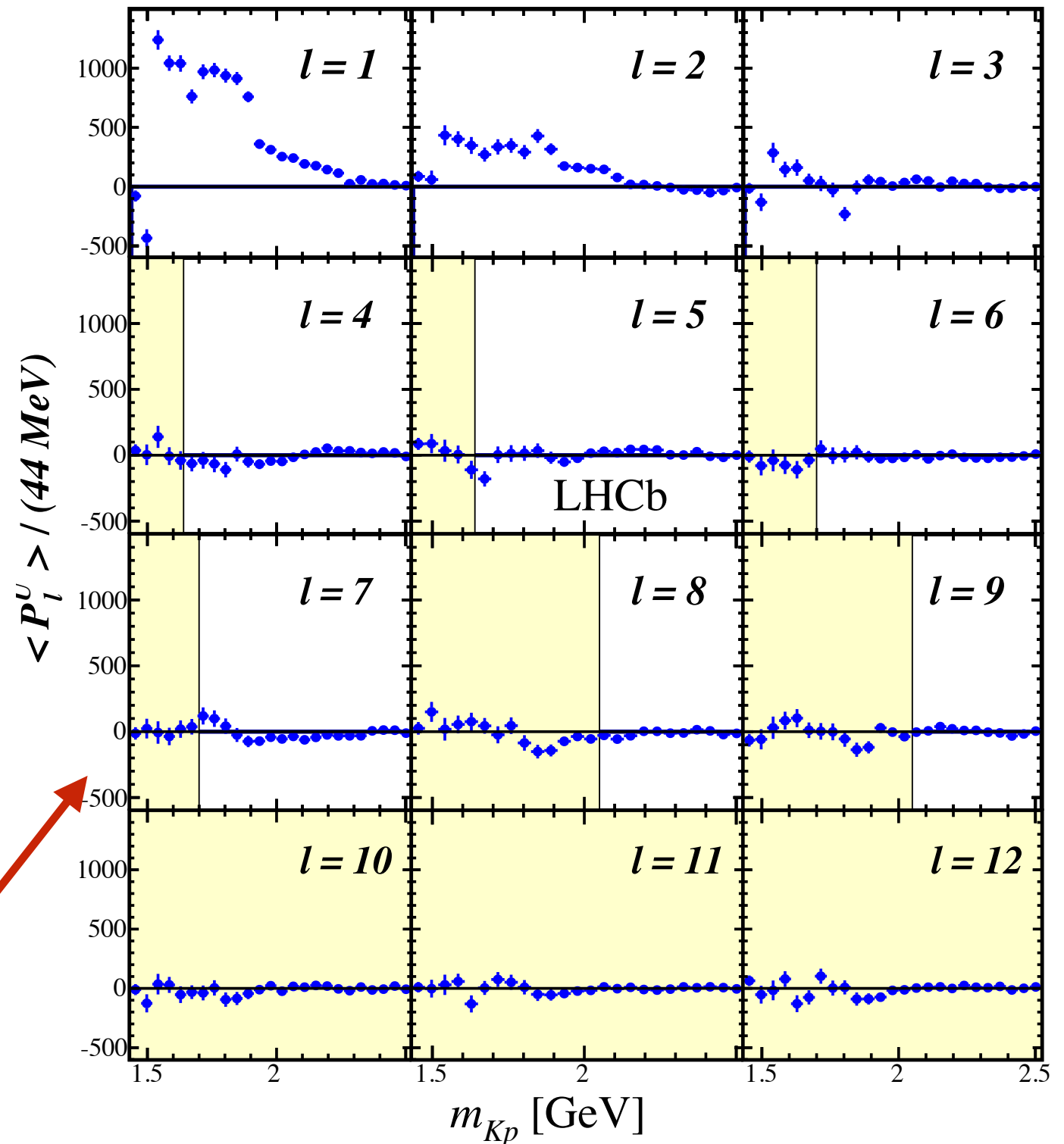


$$\Lambda_b \rightarrow p K \mu^+ \mu^-$$

[LHCb, PRL 117 (2016) 082002]

- Large number of overlapping  $\Lambda^*$  resonances with different  $J^{PC}$ .
- Can separate using an amplitude analysis of the  $pK\mu^+\mu^-$  system.
- Simple first step, perform a moment analysis in bins on  $pK$  mass (and  $q^2$ ).

c.f.  $\Lambda_b \rightarrow J/\psi p K$  model independent analysis in [LHCb, PRL 117 (2016) 082002]



$$\Lambda_b \rightarrow p K \mu^+ \mu^-$$

- Even if we cannot separate the different  $\Lambda^*$  resonances, can still perform clean tests of the SM through LFU ratios.
  - ✓ Decay has different sources of experimental background to  $B \rightarrow K^{*0} \mu^+ \mu^-$ .

# Other $b$ -baryon decays

- What other  $b \rightarrow s \mu^+ \mu^-$  baryonic decays could we look for in the Run II/upgrade data sets?

- ➔  $\Xi_b^- \rightarrow \Xi^- \mu^+ \mu^-$  (with  $\Xi^- \rightarrow \Lambda \pi^-$ )
- ➔  $\Omega_b^- \rightarrow \Omega^- \mu^+ \mu^-$  (with  $\Omega^- \rightarrow \Lambda K^-$ )

Decay chains involve  
two long-lived particles  
⇒ low efficiency


- From TeVatron:

$$\frac{f_{\Xi_b^-}}{f_{\Lambda_b}} = 0.167 \pm 0.037 \pm 0.012 ,$$

$$\frac{f_{\Omega_b^-}}{f_{\Lambda_b}} = 0.045 \pm 0.017 \pm 0.004 ,$$

[CDF , PRD 80 (2009) 072003]  
[D0, PRL 99 (2007) 052001]

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**Decay chains involve two long-lived particles  
⇒ low efficiency**
  - Based on LHCb measurements of  $\Xi_b^- \rightarrow \Xi^- J/\psi$  and  $\Omega_b^- \rightarrow \Omega^- J/\psi$  [\[PLB 736 \(2014\) 154\]](#) expect:
    - $O(10) \Xi_b^- \rightarrow \Xi^- \mu^+ \mu^-$
    - $O(2) \Omega_b^- \rightarrow \Omega^- \mu^+ \mu^-$
- decays to be reconstructible in the Run 2 data set.

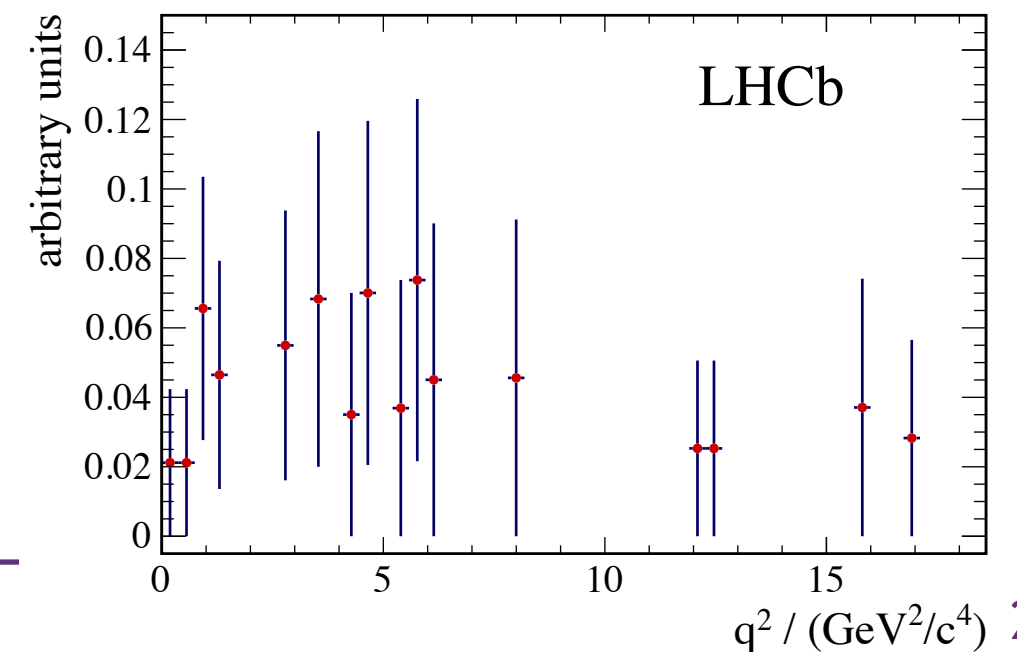
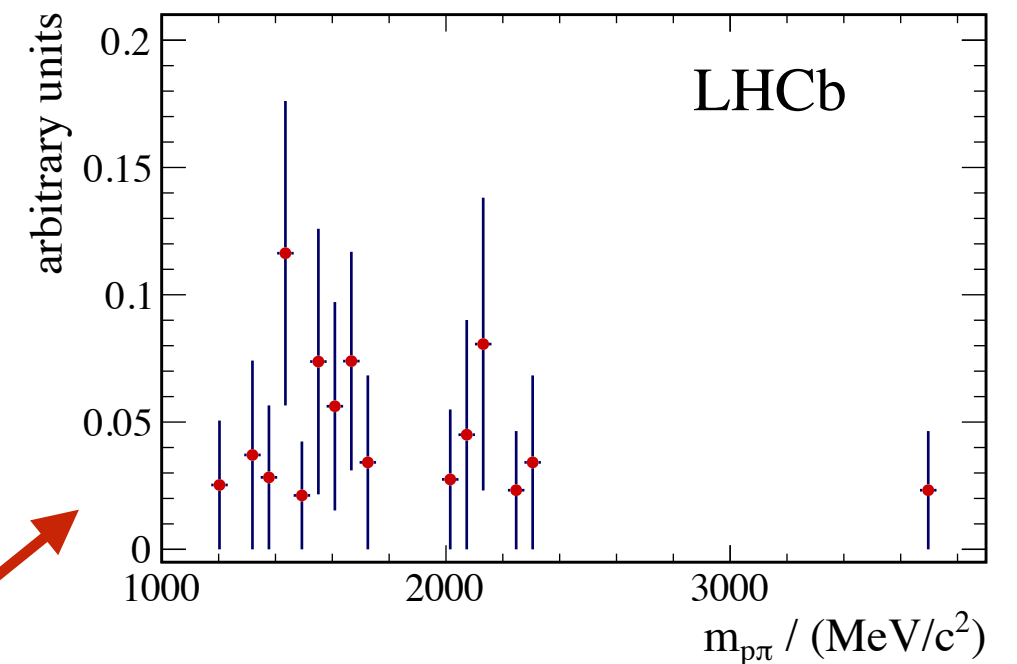
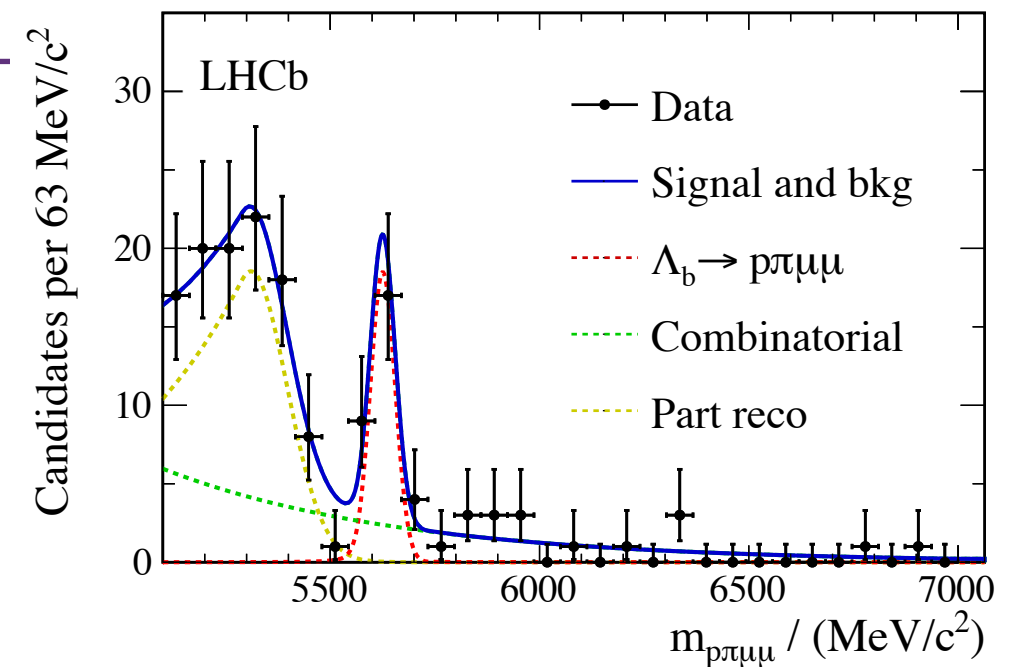
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    - ➔  $\Omega_b^- \rightarrow \Omega^- \mu^+ \mu^-$  (with  $\Omega^- \rightarrow \Lambda K^-$ )
- Decay chains involve two long-lived particles  $\Rightarrow$  low efficiency**
- Should be theoretically clean due to weak decays of  $\Xi$  and  $\Omega$ .
  - Described by a 5 or 7 (when considering the subsequent  $\Lambda$  decay) dimensional angular distribution:
    - ➔ Do we get additional NP sensitivity by probing the  $\Lambda$  helicity?



$$\Lambda_b \rightarrow \rho \pi \mu^+ \mu^-$$

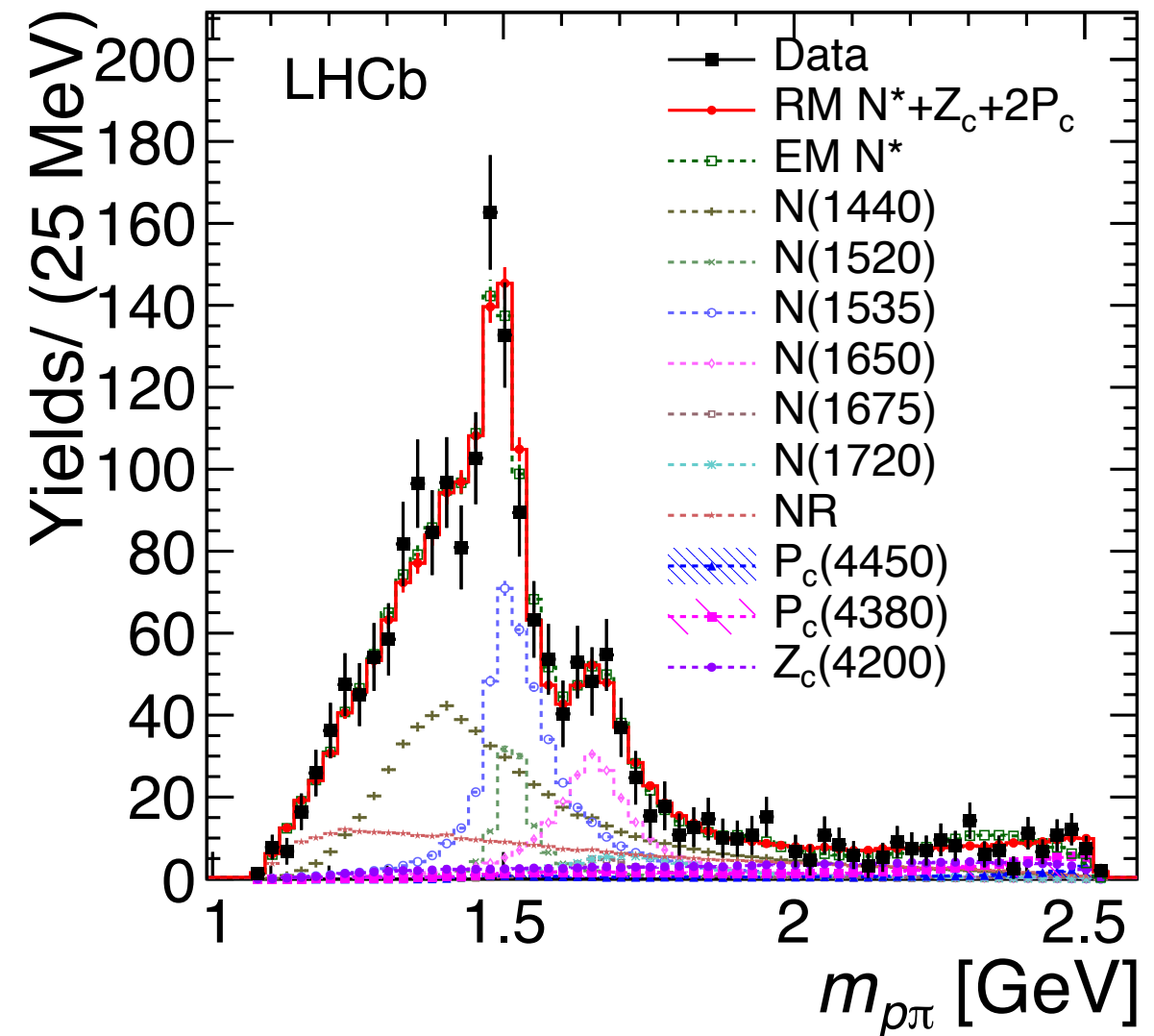
- Mediated by a  $b \rightarrow d \mu^+ \mu^-$  FCNC.
- First observation with  $5.5\sigma$  significance ( $22 \pm 6$  decays) in the LHCb Run 1 dataset.
- Same conceptual issue as  $p K \mu^+ \mu^-$  when comparing to theory predictions:
  - ➔ Large number of overlapping  $N^*$  resonances with different  $J^{PC}$ .



$$\Lambda_b \rightarrow J/\psi \rho \pi$$

[LHCb, PRL 117 (2016) 082003]

- An amplitude analysis of  $\Lambda_b \rightarrow J/\psi \rho \pi$  has been carried out to search for exotic states in [LHCb, PRL 117 (2016) 082003].
- To describe the data need to consider a number of overlapping  $N^*$  resonances along with a broad non-resonant contribution.



# Other $b$ -baryon decays

- What about  $b \rightarrow d \mu^+ \mu^-$  decays to ground-state baryons?

→  $\Lambda_b \rightarrow n \mu^+ \mu^-$     **x neutron in final-state**

→  $\Xi_b^- \rightarrow \Lambda \mu^+ \mu^-$

→ ...

Decays are isospin  
violating, expect  
negligible rate.

# Radiative $\Lambda_b$ decays

- Photon polarisation can be determined from the angular distribution of baryonic decays,

$$\alpha_\gamma = \frac{P(\gamma_L) - P(\gamma_R)}{P(\gamma_L) + P(\gamma_R)}$$

- $\Lambda_b \rightarrow \Lambda \gamma$  **[G. Hiller & A. Kagan, PRD 65 (2002) 074038]**

$$\frac{d\Gamma}{d \cos \theta_\gamma} = \frac{1}{2}(1 - \alpha_\gamma P_b \cos \theta_\gamma) \qquad \frac{d\Gamma}{d \cos \theta_p} = \frac{1}{2}(1 - \alpha_\gamma \alpha_\Lambda \cos \theta_p)$$

- $\Lambda_b \rightarrow \Lambda^* \gamma$  **[F. Legger & T. Schietinger, PLB 645 (2007) 204-212]**

$$\frac{d\Gamma}{d \cos \theta_\gamma} = \frac{1}{2}(1 - \alpha_\gamma \alpha_{3/2} P_b \cos \theta_\gamma) \qquad \frac{d\Gamma}{d \cos \theta_p} \propto \frac{1}{2}(1 - \alpha_{p,3/2} \cos^2 \theta_p)$$

- Only have sensitivity from  $\theta_p$  in  $\Lambda_b \rightarrow \Lambda \gamma$  given due to small size of  $P_b$ .

# Radiative $\Lambda_b$ decays

- Using  $\Lambda_b \rightarrow \Lambda \gamma$  expect sensitivity to  $\alpha_\gamma$  of:

- ➔ 25% with Run 2 dataset

- ➔ 15% with 23fb<sup>-1</sup>

- ➔ 4% with 300fb<sup>-1</sup>

**[LHCb, Upgrade II physics case, arXiv:1808.08865].**

Note, this uses the old PDG value of  $\alpha_\Lambda$ .

Expect a 17% improvement with the larger BESIII value.


- We will also measure the  $CP$  asymmetry of  $\Lambda_b \rightarrow p K \gamma$  and  $\Lambda_b \rightarrow p \pi \gamma$ .

# Other radiative decays

- Could also look at  $\Xi_b^- \rightarrow \Xi^- \gamma$  (with  $\Xi^- \rightarrow \Lambda \pi^-$ ):
  - ➔ Involves two weak decays and will be challenging to reconstruct experimentally.
  - ➔ Expect  $O(10)$  candidates in the LHCb Run 2 data set.

- More complex angular distribution given by:

$$\frac{d^2\Gamma}{d\cos\theta_\Lambda d\cos\theta_p} = \frac{1}{4} (1 - \alpha_\gamma \alpha_\Xi \cos\theta_\Lambda + \alpha_\Lambda \cos\theta_p (\alpha_\Xi - \alpha_\gamma \cos\theta_\Lambda))$$



$-0.39 \pm 0.01$                        $+0.75 \pm 0.01$

- Expect sensitivity of 40% (10%) on  $\alpha_\gamma$  with  $23\text{fb}^{-1}$  ( $300\text{fb}^{-1}$ )  
[\[LHCb, Upgrade II physics case, arXiv:1808.08865\]](#).

# Summary

- Large number of different processes could be studied with the data collected during Runs 2+ of the LHC.
- Often have a trade-off between ease of the experimental measurement and theoretical complexity,  
e.g. in dealing with overlapping states in  $ph-\mu^+\mu^-$ .





# $\Lambda_b \rightarrow \Lambda \mu^+ \mu^-$ branching fraction

- New numbers use:

$$f_{\Lambda_b} \times \mathcal{B}(\Lambda_b \rightarrow J/\psi \Lambda) = (5.8 \pm 0.8) \times 10^{-5} , \quad \text{[PDG 2018]}$$

with

$$f_{\text{baryon}} = 0.218 \pm 0.047 \quad \text{[HFLAV 2017]}$$

assuming

$$f_{\text{baryon}} = f_{\Lambda_b} \left( 1 + 2 \frac{f_{\Xi_b^-}}{f_{\Lambda_b}} + \frac{f_{\Omega_b^-}}{f_{\Lambda_b}} \right) ,$$

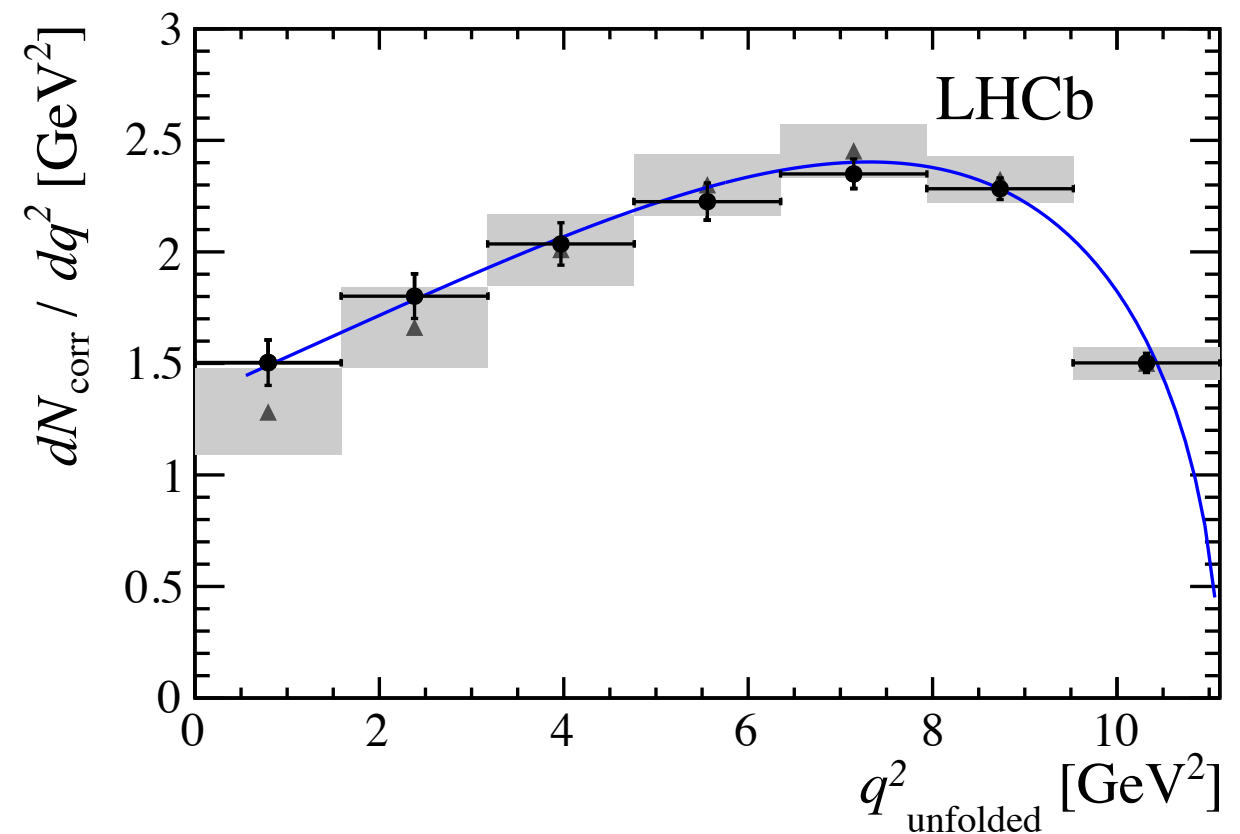
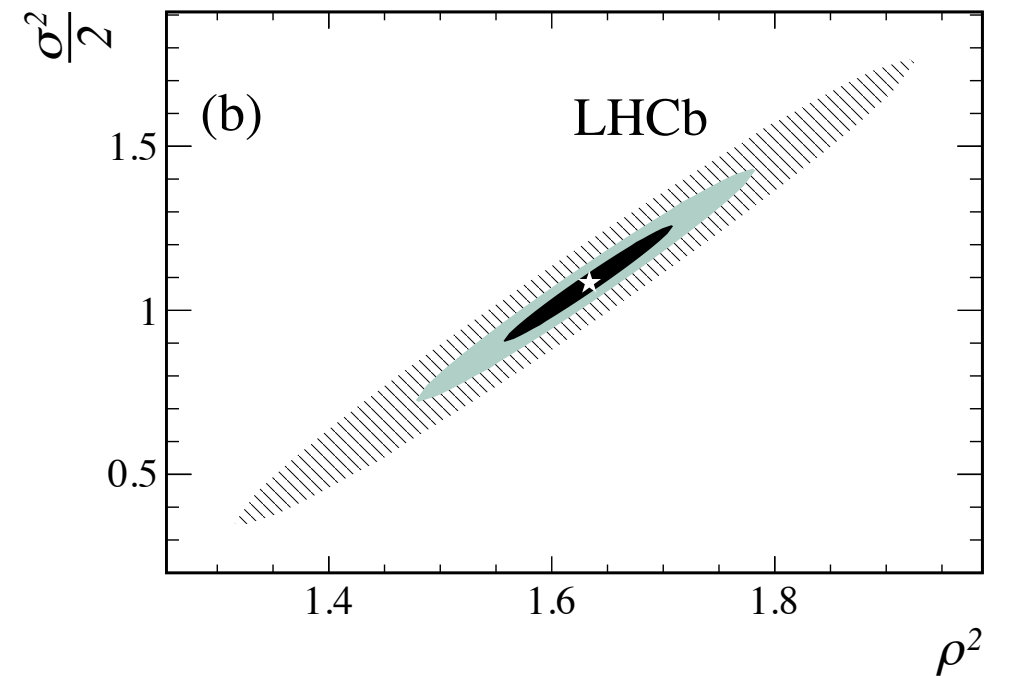
where

$$\frac{f_{\Xi_b^-}}{f_{\Lambda_b}} = 0.167 \pm 0.037 \pm 0.012 , \quad \text{[PDG 2018]}$$

$$\frac{f_{\Omega_b^-}}{f_{\Lambda_b}} = 0.045 \pm 0.017 \pm 0.004 ,$$

$$\Lambda_b \rightarrow \Lambda_c^+ l^- \nu$$

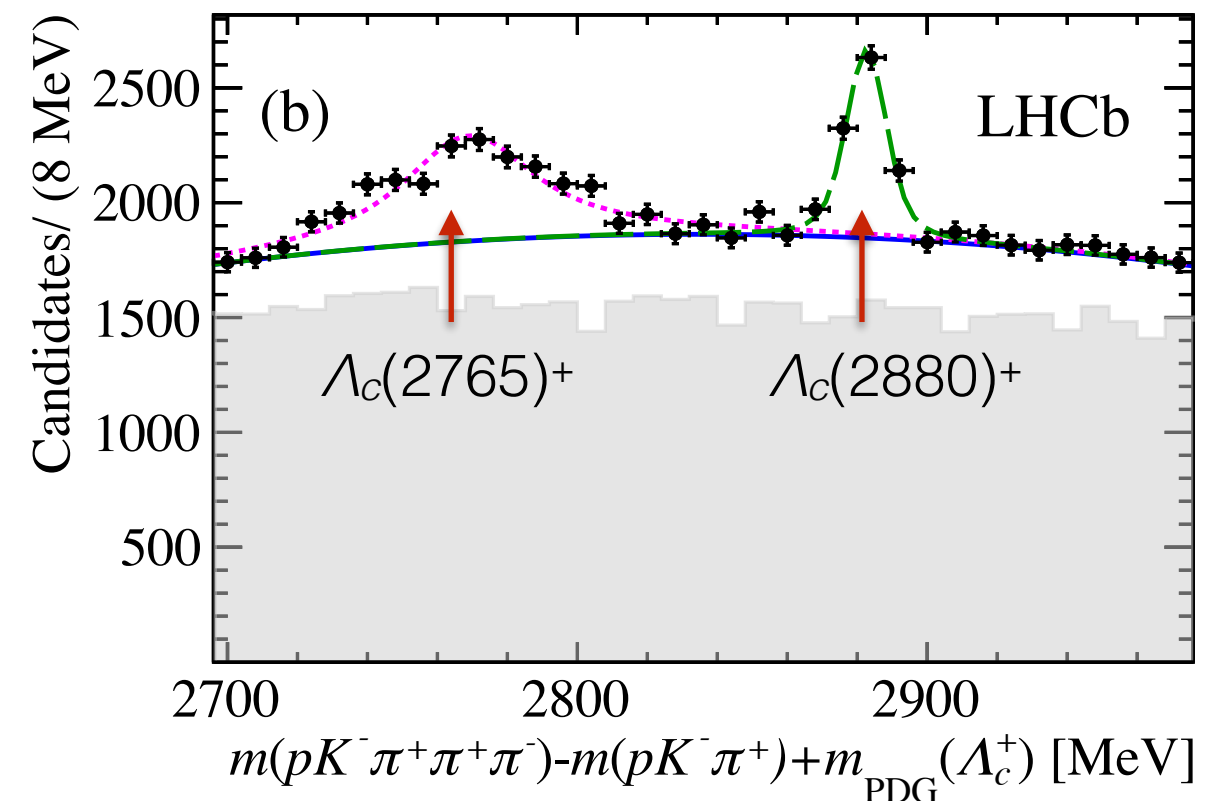
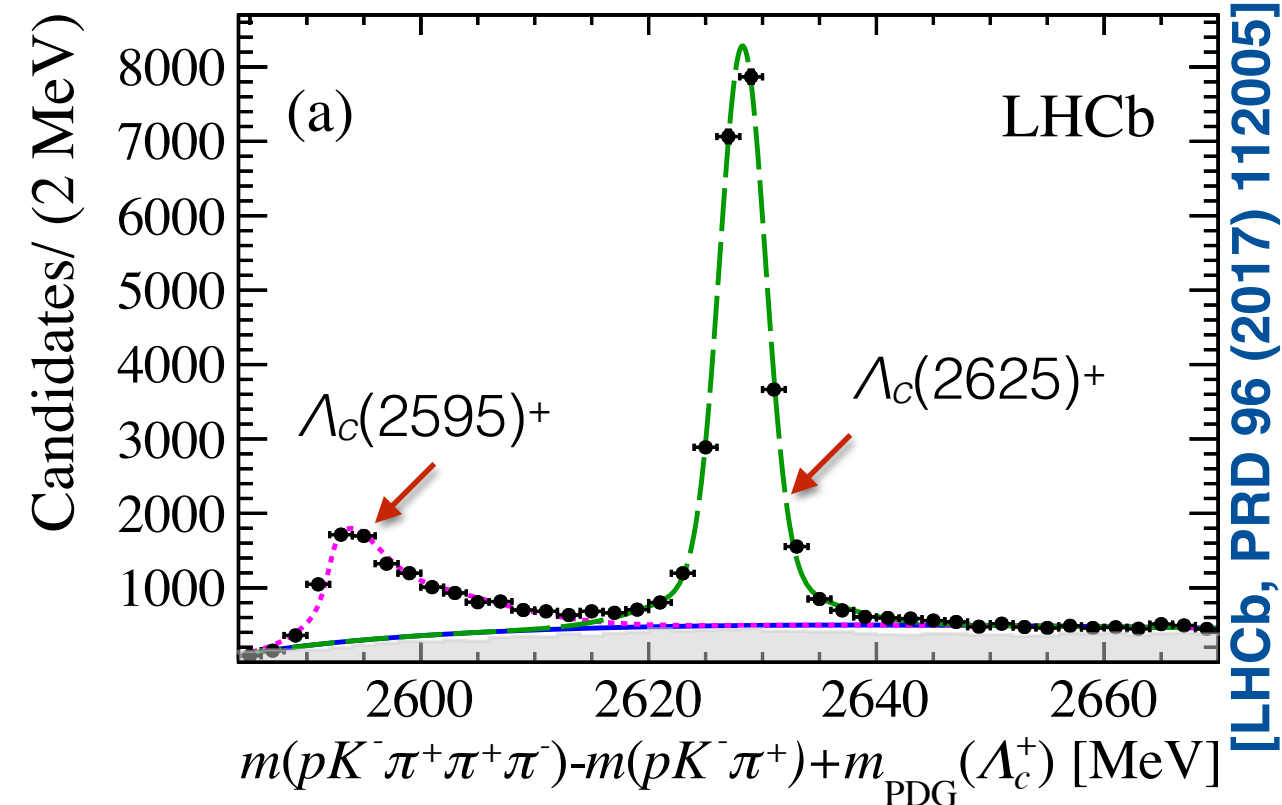
- Only existing measurement is of  $\Lambda_b \rightarrow \Lambda_c^+ l^- \nu$  form-factors.
- Reconstruct neutrino momentum up-to two-fold ambiguity by exploiting large  $\Lambda_b$  flight distance.
- Unfold  $q^2$  and Isgur-Wise parameters.
- Good agreement between data/HQET/lattice.



■ Lattice predictions from  
**[Detmold et al. PRD 92 (2015) 034503]**

$$\Lambda_b \rightarrow \Lambda_c^{(*)} l \nu$$

- Background studies show large number of different  $\Lambda_c^*$  states in the data.
- Next steps are to:
  - ➔ perform an angular analysis of  $\Lambda_b \rightarrow \Lambda_c^{(*)} l \nu$ .
  - ➔ measure  $R(\Lambda_c)$ .

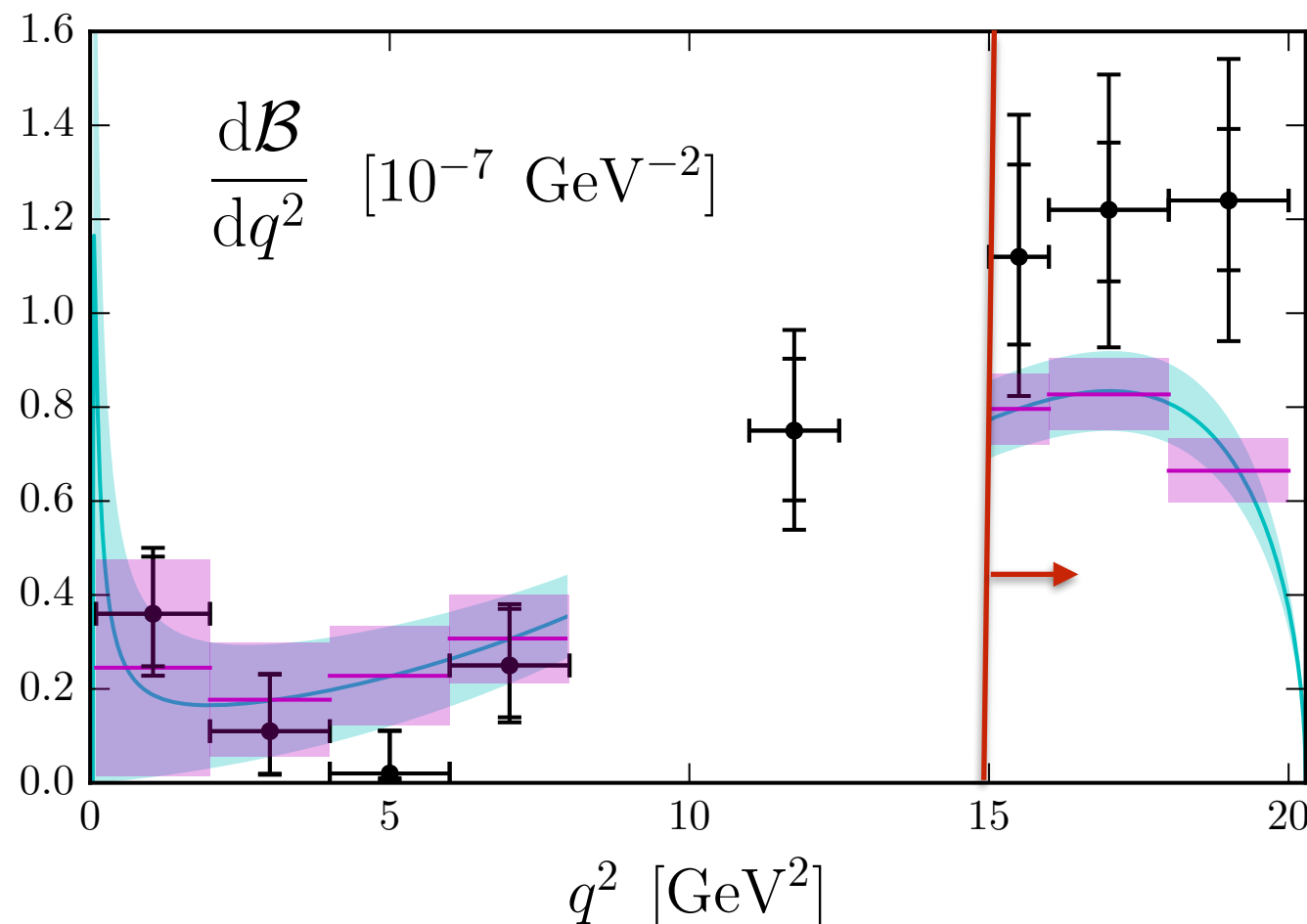


[LHCb, PRD 96 (2017) 112005]

$$\Lambda_b \rightarrow \Lambda \mu^+ \mu^-$$

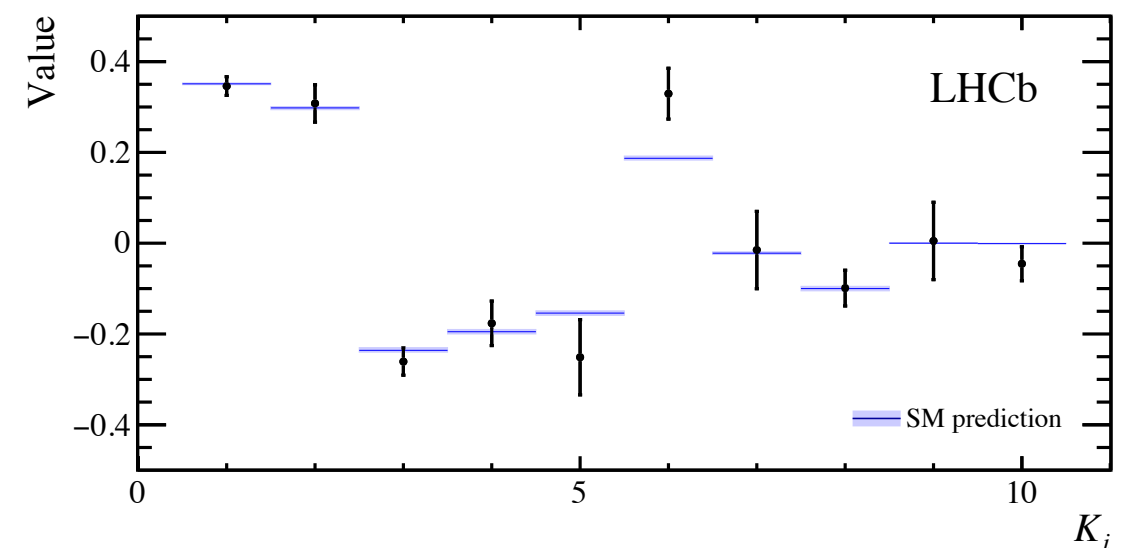
- Large number (34) of angular observables, 24 require  $\Lambda_b$  to be polarised at production and are consistent with zero in current dataset.

Detmold et al. arXiv:1602.01399



- $\Lambda_b \rightarrow D\bar{D}\Lambda$  can provide useful input on long-distance contributions.

LHCb, arXiv:1808.00264



- SM predictions use external on  $\alpha_\Lambda$  and production polarisation (small at the LHC).
- Branching fraction uncertainty currently dominated by knowledge of  $\mathcal{B}(\Lambda_b \rightarrow J/\psi \Lambda)$ .