

Signals from the structure@horizon

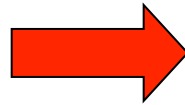
Iosif Bena

IPhT, CEA Saclay

Nick has hopefully convinced you that there must exist **STRUCTURE @ horizon** scale

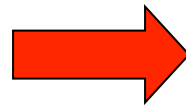
Analogy with ideal gas:

Thermodynamics
(Air = ideal fluid)
 $P V = n R T$
 $dE = T dS + P dV$



Statistical Physics
(Air -- molecules)
 e^S microstates
typical
atypical

Thermodynamics
Black Hole Solution



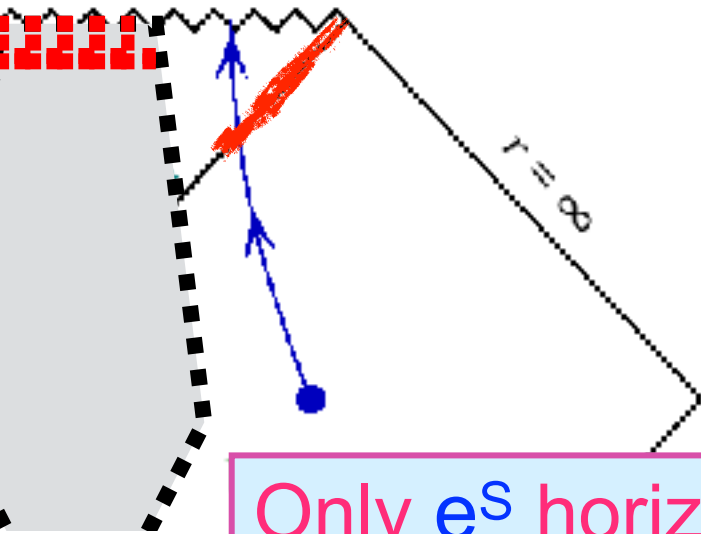
Statistical Physics
Microstate geometries

Long distance physics
Gravitational lensing

Physics at horizon
Information loss
Gravity waves ?

Avoid forming a horizon

- Collapsing shell forms horizon Oppenheimer and Snyder (1939)
- If curvature is low, no reason not to trust classical GR
- By the time shell becomes curved-enough for quantum effects to become important, horizon in causal past (60 hours for NGC 4889 BH)



Backwards in time - illegal !

BH has e^S microstates with no horizon
Small tunneling probability = e^{-S}
Will tunnel with probability ONE !!!

Kraus, Mathur; Bena, Mayerson, Puhm, Vercnocke

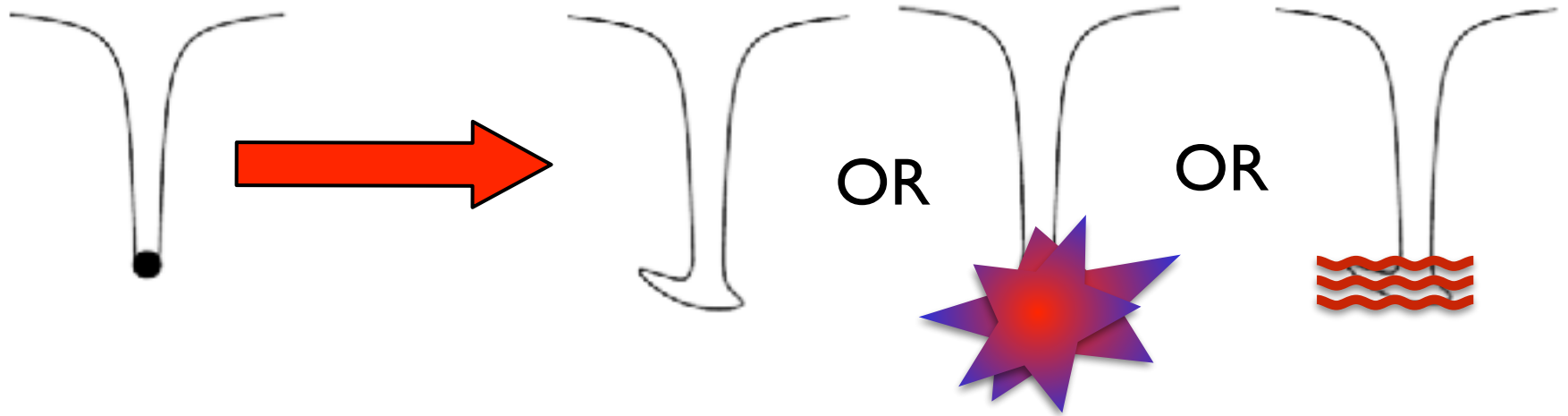
Only e^S horizon-sized microstates can do it !

Black hole entropy the structure must have



MICROSTATE GEOMETRIES

whose construction we can trust



**HOW DO WE GET FROM
String Theory BH to Real World BH ?**

**HOW DOES STRUCTURE AFFECT
HORIZON PHYSICS ?**

Important feature

- Have the same size as the BH for **any** mass
↔ size grows with G_{Newton}
- Boson stars need **different scalar fields of different masses** to replace various BH's:
One field for M_{\odot} , another for $30 M_{\odot}$, etc
- String theory automatically produces fields whose **mass decreases for larger BH**
 - Non-perturbative physics
- They are the reason why microstate geometries have BH size for ANY BH mass

Simplest Black Hole:

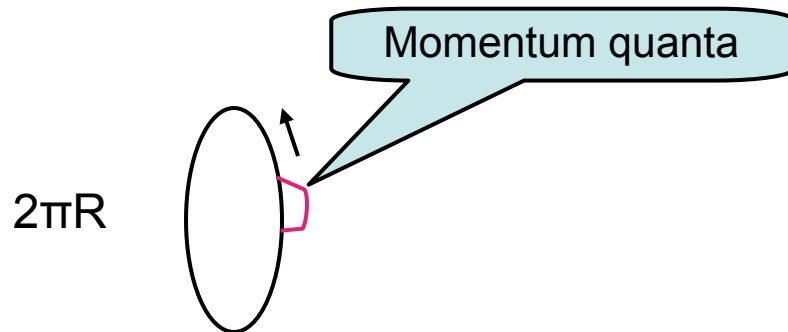
Strominger, Vafa

D1 branes (strings), D5 branes, momentum P

one D1 brane, $2\pi R$



Momentum quanta
 $\delta P = 1/R$



Simplest Black Hole:

Strominger, Vafa

D1 branes (strings), D5 branes, momentum P

one D1 brane, $2\pi R$

N_1 D1 branes, $2\pi R$

one D1 brane, $2\pi N_1 R$

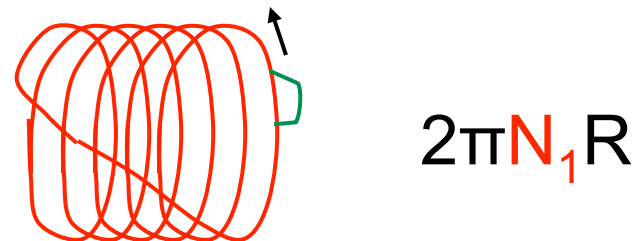
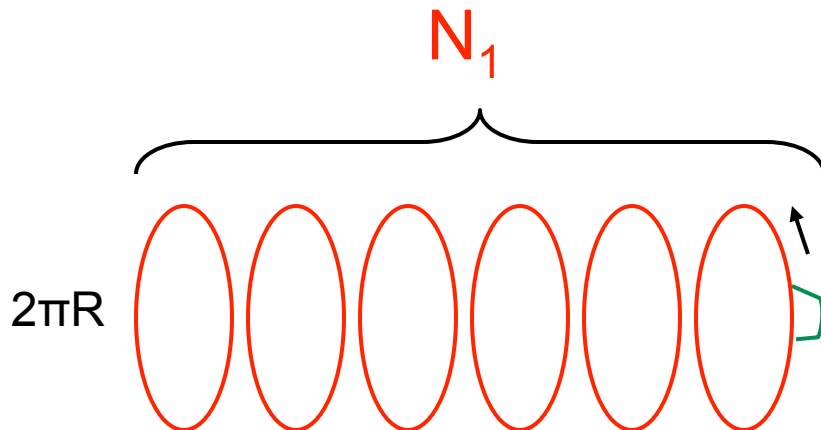


$$\delta P = 1/R$$



$$\delta P = 1/N_1 R$$

Momentum quanta



Simplest Black Hole:

Strominger, Vafa

D1 branes (strings), **D5** branes, momentum **P**

one D1 brane, $2\pi R$

N_1 D1 branes, $2\pi R$

1 D1 brane, $2\pi N_1 R$

N_1 D1 + N_5 D5 branes

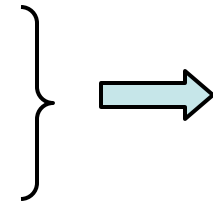
effective string, $2\pi N_1 N_5 R$



$$\delta P = 1/R$$

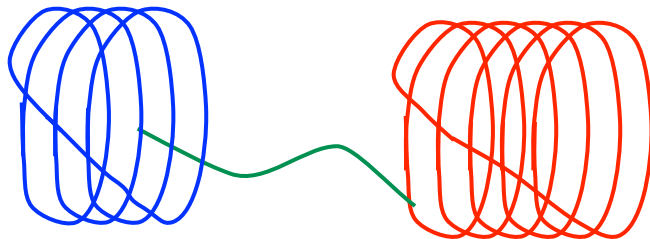


$$\delta P = 1/N_1 R$$



$$\delta P = 1/N_1 N_5 R$$

Momentum quanta



Microstate Counting

Strominger, Vafa

- Total momentum N_P / R carried by quanta of

$$1/ N_1 N_5 R$$

- Total = $N_1 N_5 N_P$ quanta

- Number of states \Leftrightarrow partitions of $N_1 N_5 N_P$

- How many states (partitions) ?

$$N_1 N_5 N_P = 2 : (1,1) (2)$$

$$N_1 N_5 N_P = 3 : (1,1,1) (2,1) (3)$$

$$N_1 N_5 N_P = 5 : (1,1,1,1,1) (1,1,1,2) (1,1,3) (1,4) (5) (1,2,2) (2,3)$$

$N_1 N_5 N_P$ quanta: e^S states,

$$S_{\text{MICRO}} = 2\pi(N_1 N_5 N_P)^{1/2} = S_{\text{BLACK HOLE}} !!!$$

Low-mass degrees of freedom

Given total energy budget:
most entropy obtained by making
brane-antibrane pairs:

$$S = 2\pi (N_1^{1/2} + \underline{N}_1^{1/2})(N_2^{1/2} + \underline{N}_2^{1/2})(N_3^{1/2} + \underline{N}_3^{1/2})$$

Horowitz, Maldacena, Strominger

$$\text{Mass gap} = 1/N_1 N_2 N_3$$

Larger BH \rightarrow lower mass-gap \rightarrow larger size

Essentially same result as boson star whose scalar mass decreases with increasing BH size !

Easy (relatively...)

Hard

Supersymmetric

Non-supersymmetric

Extremal

Non-extremal

Classical

Quantum

Stationary

Time-dependent

Universal features:

YugeTM amount of new degrees of freedom @ horizon

Two new scales: size of bubbles

distance from the horizon - E_{GAP}

Horizon viscosity

Microstate mountains

Distortion of Love numbers / Kerr multipoles

Horizon viscosity

Lots of low-mass-gap degrees of freedom: $e^{10^{90}}$
(**entropy** of BH @ center of galaxy = 10^{90})

- can cause dissipative effects (like for electron moving parallel to a conducting plate)
- any distortion of structure@horizon immediately dissipated by these degrees of freedom
- hard to get any reflection - unless one can **convince** these degrees of freedom **not to get excited**

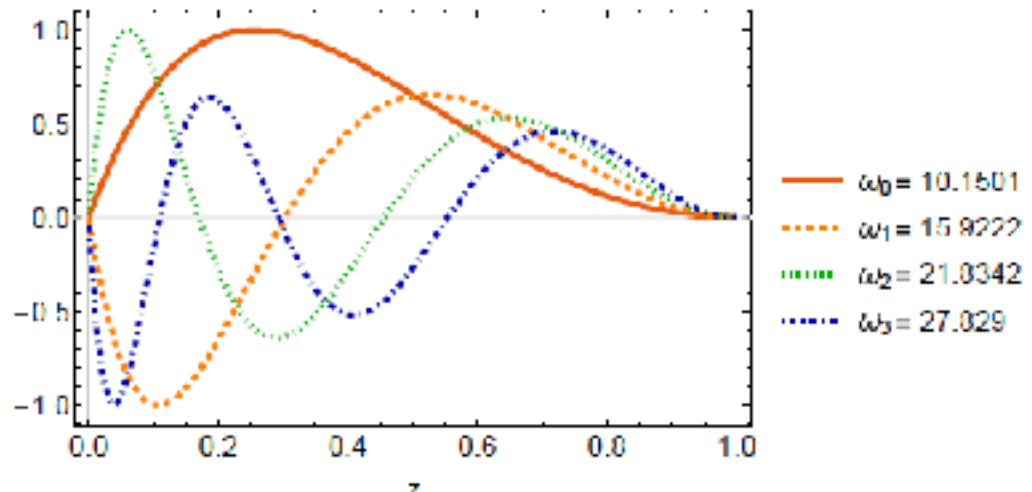
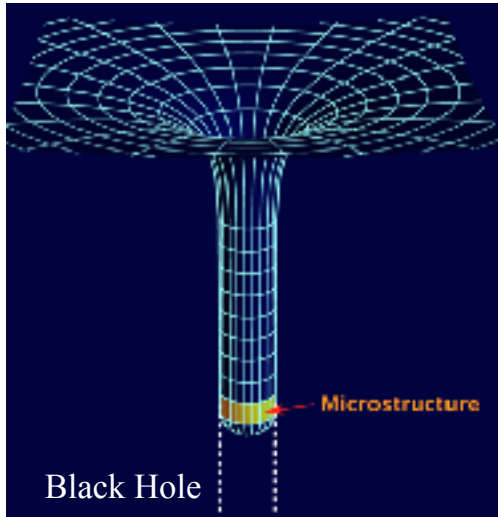


Already taken into account by **perfectly-absorbing boundary conditions** of GR BH ?

Why do people get reflection / echoes ?

Ashfordi & friends

Solve wave equation in one particular microstate geometry



Beautiful gapped spectrum.

Infalling wave reflected \rightarrow echoes

Ignoring **huge number** of low-energy excitations that move this geometry to the other geometries

Microstate mountains

- microstate geometries **break** spherical symmetry
 - similar to **molecules breaking translational symmetry** of fluid approx.
 - visible at the scale of the **mean free path** (shot noise)
- analogous to mountains on neutron stars
 - **create jitters on infalling stuff**

Questions:

- What is the scale where spherical symmetry is broken ? (**width of microstate mountains ?**)
- By how much ? (**height of mountains ?**)
- Can people who do signals from neutron stars implement them in their calculations ?

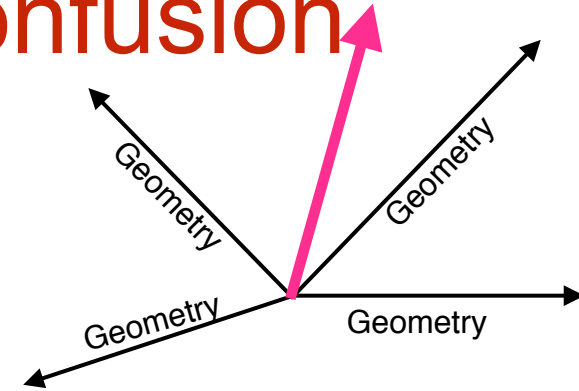


Typical microstates may be **very spherical ?**

Planck-size bumps ?

Quantum - Classical confusion

- Quantum Tunneling important
- Hilbert space dimension = $e^{10^{90}}$
- Microstate geometries = **coherent states**
 - **basis** for Hilbert space (overcomplete) Glauber '63
- Typical state in Hilbert space will be **quantum superposition of geometries** →
expectation values $\sim e^{-S/2}$???
- Decoherence → one will get result from **typical microstate** Don Page, private communication



Multipole moments

- spinning ball of dust, liquid, solid shell, boson star
- Kerr BH with vacuum at horizon
- spinning **STRUCTURE@horizon** scale

Gravity waves from **EMRI with non-aligned spins:**

Need 3 moments to rule out Kerr BH, 4 to rule out spinning boson star
(Ryan '95)

What we **need**: Kerr microstate geometries

What we **have**: NHEK microstates (Heidmann '18)

Next steps: Glue NHEK to Kerr
go non-extremal

Multipole moments

Love numbers : response to tidal stress

zero for BH Damour, Nagar '09

finite for other objects

log dependence on scale above horizon

Cardoso, Franzin, Maselli, Pani, Raposo '17

Computable once we have **Kerr microstates**



Fuzzball complementarity: unusual matter forming **STRUCTURE @ horizon** gives “classical horizon experience” to infalling observer. **Same as vacuum.**

Mathur, Turton

What if String Theory does not describe real world ?

- *Tant pis pour le monde réel* 😞
- More seriously: structure@horizon **must be there** in any QG theory
- Must have **unusual** properties:
 - lighter-mass effective fields for increasing BH size
 - **non-perturbative** degrees of freedom
 - nontrivial 4D matter to **stay@horizon** (no **Pixie dust**)
 - huge entropy: **quantum tunneling**
- “top-down predictions” ignoring this physics = b.s.
- Universal features - like low QGP viscosity predicted by AdS-CFT, measured at RHIC

What would help

- Solve second-order PDE's with >10 functions of 2 variables
 - (rod structure \rightarrow pentagon and more)
 - other spectral methods
- Essential for constructing systematically **non-extremal** BH microstates
- Signals for microstate mountains (same code as for neutron stars)
- Extra d.o.f. modify viscosity - can one take that into account ?