Stochastic and resolvable gravitational waves from ultralight bosons

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> > (mainly) based on

arxiv: 1706.05097; 1706.06311; 1805.08229; see also arXiv:1404.7149; 1811.07786



## Outline

- A review of boson condensate formation around spinning BHs and their GW emission
- Astrophysical models for spinning BHs
- Constraints on boson masses in the LISA and LIGO bands by
  - Direct detections
  - Stochastic backgrounds
  - "Holes" in Regge plane
- Constraints on BH mimickers

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## Why light bosons?

- Scalars ubiquitous in string theory, inflation, dark matter models (e.g. fuzzy/axionic dark matter)
- Useful as toy models for unknown phenomena/ interactions (e.g. modifications of GR)
- "Light" means <~ 1.e-10 eV</li>
- Effect of mass term expected to be qualitatively similar for all boson degrees of freedom

# Self-gravitating scalar configurations

- Scalars can form self-gravitating configurations, especially if complex, massive (to avoid dispersion to infinity) and time dependent (to provide pressure): boson stars, oscillatons
- Around BHs, massive real (complex) scalars can form quasi-stationary (stationary) configurations: boson clouds or condensates, hairy BHs

## BH-boson condensates

- Formation linked to superradiant instabilities/Penrose process (amplification of scattered waves with  $\omega < m\Omega_H$ )
- BH with high enough spin and "mirror" are superradiance unstable (BH bomb; Zeldovich 71, Press & Teukolsky 72, Cardoso et al 04)
- In ergoregion, negative energy modes can be produced but are confined (only positive energy modes can travel to infinity)
- By energy conservation, more and more negative energy modes can be produced, which may cause instability according to boundary conditions (at horizon and spatial infinity)





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#### Superradiance from near horizon physics



- Deviations away from Kerr geometry near horizon (e.g. firewalls, gravastars, wormholes, Lorentz violations, etc) can produce significant changes in QNM spectrum
- Delays  $\Delta t \sim \log[r_0/(2M) 1]$



Cardoso, Franzin & Pani 2016

EB, Cardoso & Pani 2014

## BH-boson condensates

 Same instability of spinning BH + massive boson (mass acts as "mirror" and allows for bound states), but NOT for fermions. Cf Damour, Deruelle & Ruffini 76



## Instability end point

 BH sheds excess spin (and to a lesser degree mass) into a mostly dipolar rotating boson cloud ...

$$m_s \equiv \mu \hbar$$
,  $\omega_R \sim \mu - rac{M^2 \mu^3}{8}$ 

$$\Phi = A_0 g(r) \cos(m_\phi \phi - \omega_R t) \sin \theta$$
,



• ... till instability saturates

 $\mu \sim m \Omega_{
m H}$ 

$$\left| \tau_{\rm inst} \sim 0.07 \, \chi^{-1} \left( \frac{M}{10 \, M_\odot} \right) \left( \frac{0.1}{M \mu} \right)^9 \, {\rm yr} \, , \right.$$

(for Mµ<<1 and  $\chi$ <<1; max instability for Mµ=0.42)

## GW emission

 Long-lived rotating scalar dipole produces almost monochromatic GWs via quadrupole formula on timescale

$$au_{\rm GW} \sim 6 \times 10^4 \, \chi^{-1} \left( \frac{M}{10 \, M_\odot} \right) \left( \frac{0.1}{M \mu} \right)^{15} \, {\rm yr}$$

$$h = \sqrt{rac{2}{5\pi}} rac{GM}{c^2 r} \left(rac{M_S}{M}
ight) A(\chi, f_s M),$$
 rms strain amplitude



# Background from isolated spinning BHs

energy emission efficiency

 $f_{\rm ax} \sim \mathcal{O}(1\%)$ 

$$\Delta \ln f \sim 1$$

LISA band massive BHs ~  $10^4\text{--}10^7\ M_{sun},\ m_s\text{--}10^{-16}\text{--}10^{-18}\ eV$ 

$$\begin{split} \rho_{\rm BH} &\sim \mathcal{O}(10^4) M_{\odot}/{\rm Mpc}^3 \\ \Omega_{\rm GW,\,ax} &= (1/\rho_{\rm c}) (d\rho_{\rm GW}/d\ln f) \sim f_{\rm ax} \rho_{\rm BH}/\rho_{\rm c} \\ \Omega_{\rm GW,\,ax}^{\rm LISA} &\sim 10^{-9} \end{split}$$

# Background from isolated spinning BHs

energy emission efficiency

 $f_{\rm ax} \sim \mathcal{O}(1\%)$ 

monochromatic GW in source frame

$$\Delta \ln f \sim 1$$

LIGO/Virgo band stellar-mass BHs ~ 10-50 M<sub>sun</sub>, m<sub>s</sub>~10<sup>-13</sup> - 10<sup>-12</sup> eV  $\Omega_{\rm GW, \, bin} \sim f_{\rm GW} f_{\rm m} \rho_{\rm BH} / \rho_c$  $f_{\rm GW} \sim \mathcal{O}(1\%) \quad f_{\rm m} \sim \mathcal{O}(1\%)$  $\Omega_{\rm GW, \, ax} / \Omega_{\rm GW, \, bin} \sim f_{\rm ax} / (f_{\rm GW} f_{\rm m}) \sim 10^2$ 

$$\Omega_{\rm GW, \, bin} \sim 10^{-9} - 10^{-8} \ \Omega_{\rm GW, \, ax}^{\rm LIGO} \sim 10^{-7} - 10^{-6}$$

# Background from isolated spinning BHs



### Background and resolved sources



Need to account for effect of stochastic background on sensitivity (cf e.g. WD binaries for LISA)



## Regge plane "holes"



Look for "accumulation" near instability threshold to avoid having to make assumptions on astrophysical model

### Bounds on BH mimickers

BH mimickers with no horizon are unstable to superradiance



EB, Brito, Cardoso, Dvorkin, Pani 2018

- Lorentz violations introduce "asymmetry" between space and time (e.g. Horava, Einstein-aether)
- Dispersion relations imply diverging group velocity in UV:

- Event horizon definition still possible because of preferred time foliation: universal horizon (EB, Jacobson and Sotiriou 2011; Blas, Sibiryakov 2011
- Universal horizon may be unstable and form finite area curvature singularity away from spherical symmetry: echoes/near horizon exotica? (Blas, Sibiryakov 2011; Bhattacharyya, Colombo, & Sotiriou 2016; Ramos & EB 2018)



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### Conclusions

- Ultralight bosons can induce superradiant instabilities in spinning black holes, tapping their rotational energy to trigger the growth of a bosonic condensate
- Boson condensates emit almost monochromatic GWs
- GWs are LISA/LIGO band if boson's Compton wavelength is Gm/km scale
- Main observable is stochastic background, but resolved sources and Regge plane "holes" also possible
- LIGO rules out already masses ~ a few x 10<sup>-12</sup> eV, LISA will extend to ~ 10<sup>-18</sup> eV
- Similar bounds on BH mimickers (and possibly Lorentz symmetry) from stochastic background